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**Rare decays and
MSSM phenomenology**

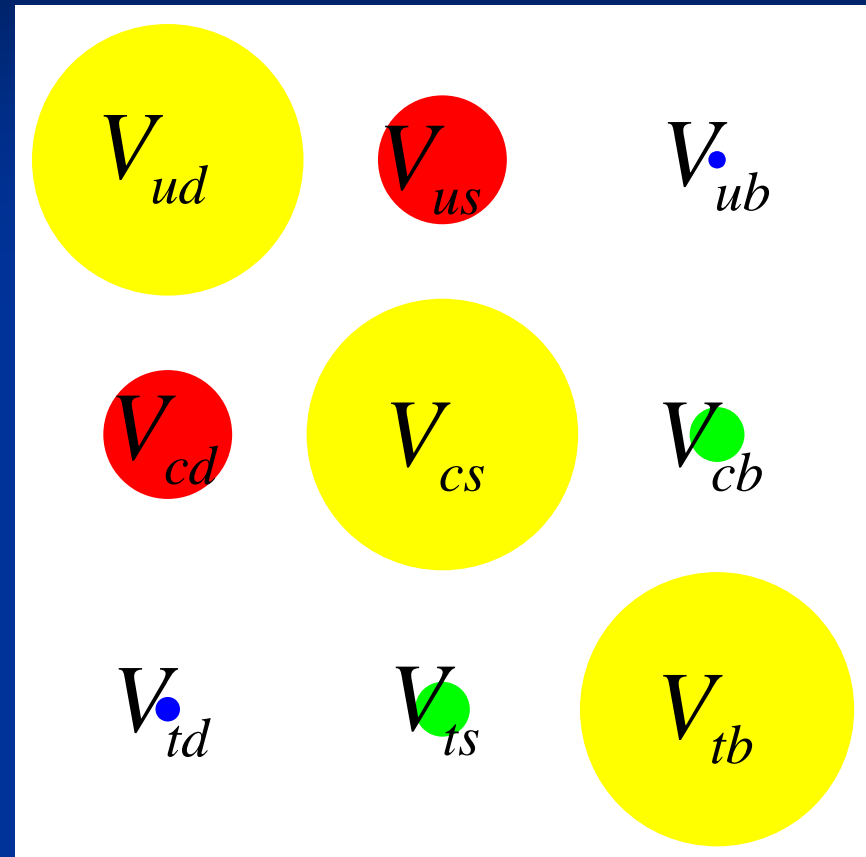
Outline:

- The SUSY flavor and CP problem
- Constraints from flavour observables
- Where are large effects still possible?
- SUSY_FLAVOUR v2.0

CKM matrix

- CKM matrix is the only source of flavor and CP violation in the SM.
- No tree-level FCNCs.
- Off-diagonal CKM elements are small

$$V_{\text{CKM}} =$$



➡ Flavor-violation is very suppressed in the SM.

SUSY flavor (CP) problem

- The sfermion mass matrices are not necessarily diagonal (and real) in the same basis as the fermion mass matrices.
- Especially the trilinear A-terms can induce dangerously large flavor-mixing (and complex phases) since they don't necessarily respect the hierarchy of the SM fermions (CKM matrix).
- The MSSM possesses two Higgs-doublets: Flavor-changing charged Higgs and (loop-induced) neutral Higgs interactions.
- **Possible solutions:**
 - MFV [D'Ambrosio, Giudice, Isidori, Strumia hep-ph/0207036](#)
 - Flavor-symmetries
 - effective SUSY [Barbieri et al hep-ph/10110730](#)
 - **Radiative flavour violation (RFV)** [A.C., Hofer, Nierste, Scherer, hep-ph/11052818](#)

Squark mass matrix

$$M_{\tilde{q}}^2 = \begin{pmatrix} M_{LL}^{\tilde{q}2} & \Delta^{\tilde{q}LR} \\ \Delta^{\tilde{q}LR\dagger} & M_{RR}^{\tilde{q}2} \end{pmatrix}$$

■ Bilinear Terms

$$M_{LL}^{\tilde{q}2} = \begin{pmatrix} (M_{LL}^{\tilde{q}2})_{11} & \Delta_{12}^{\tilde{q}LL} & \Delta_{13}^{\tilde{q}LL} \\ \Delta_{12}^{\tilde{q}LL*} & (M_{LL}^{\tilde{q}2})_{22} & \Delta_{23}^{\tilde{q}LL} \\ \Delta_{13}^{\tilde{q}LL*} & \Delta_{23}^{\tilde{q}LL*} & (M_{LL}^{\tilde{q}2})_{33} \end{pmatrix} \quad M_{RR}^{\tilde{q}2} = \begin{pmatrix} (M_{RR}^{\tilde{q}2})_{11} & \Delta_{12}^{\tilde{q}RR} & \Delta_{13}^{\tilde{q}RR} \\ \Delta_{12}^{\tilde{q}RR*} & (M_{RR}^{\tilde{q}2})_{22} & \Delta_{23}^{\tilde{q}RR} \\ \Delta_{13}^{\tilde{q}RR*} & \Delta_{23}^{\tilde{q}RR*} & (M_{RR}^{\tilde{q}2})_{33} \end{pmatrix}$$

■ Left-Right Terms

$$\Delta_{ij}^{dLR} = -v_d \left(\mu \tan(\beta) Y_i^{d(0)} \delta_{ij} + A_{ij}^d \right)$$

$$\Delta_{ij}^{uLR} = -v_u \left(\mu \cot(\beta) Y_i^{u(0)} \delta_{ij} + A_{ij}^u \right)$$

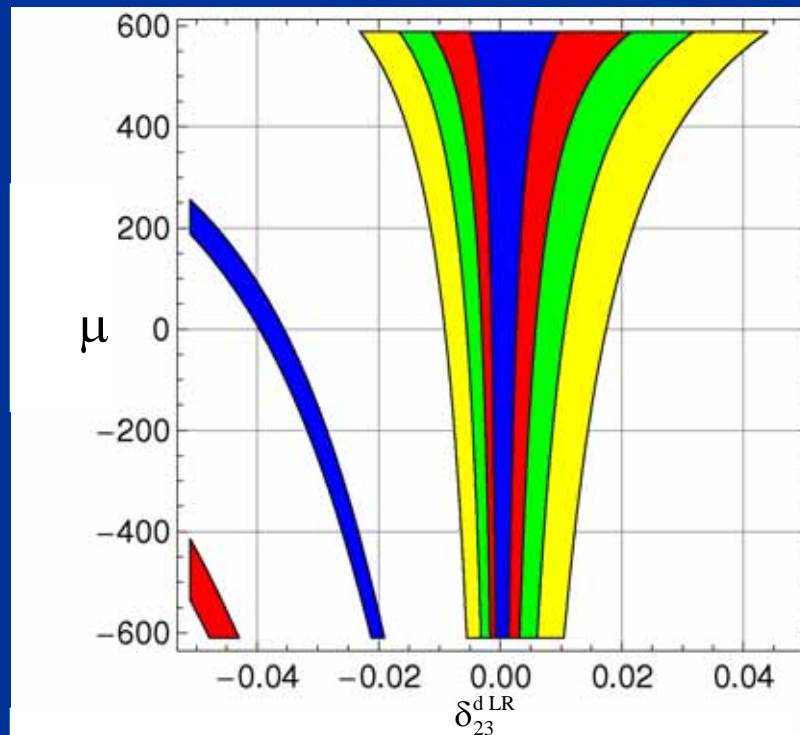
$$\tan(\beta) = \frac{v_u}{v_d}$$

$$\delta_{ij}^q = \frac{\Delta_{ij}^q}{m_{\tilde{q}}^2}$$

Constraints on the sfermion mass matrices from flavour-observables

$B \rightarrow X_{s(d)} \gamma$

- Very strong constraints on $\delta_{23,13}^{dLR}, \delta_{32,31}^{dLR}$
- Strong constraints on $\delta_{23,13}^{dLL,RR}$ if $\tan(\beta)$ is large
- $M_{H^+} > 300 \text{ GeV}$ if no cancellations occur. Misiak et al. arXiv:0609232



$$m_{g\%} = 1000 \text{ GeV}$$

$$\text{Yellow box} \quad m_{g\%} = 2000 \text{ GeV}$$

$$\text{Green box} \quad m_{g\%} = 1500 \text{ GeV}$$

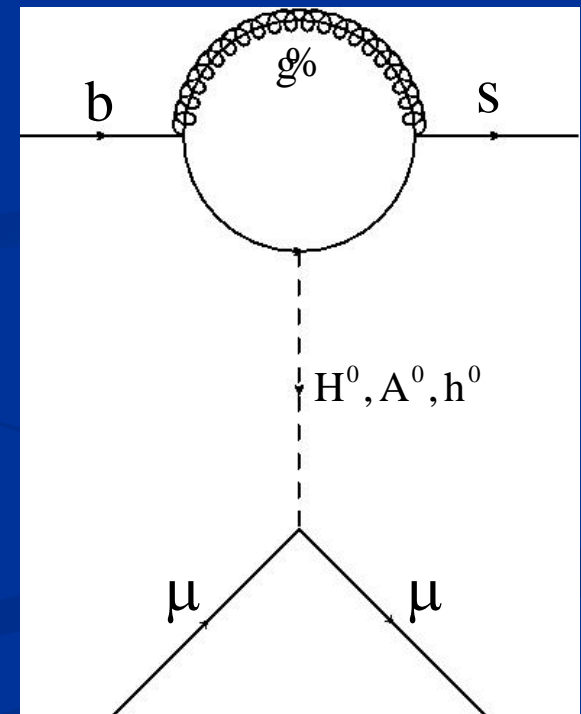
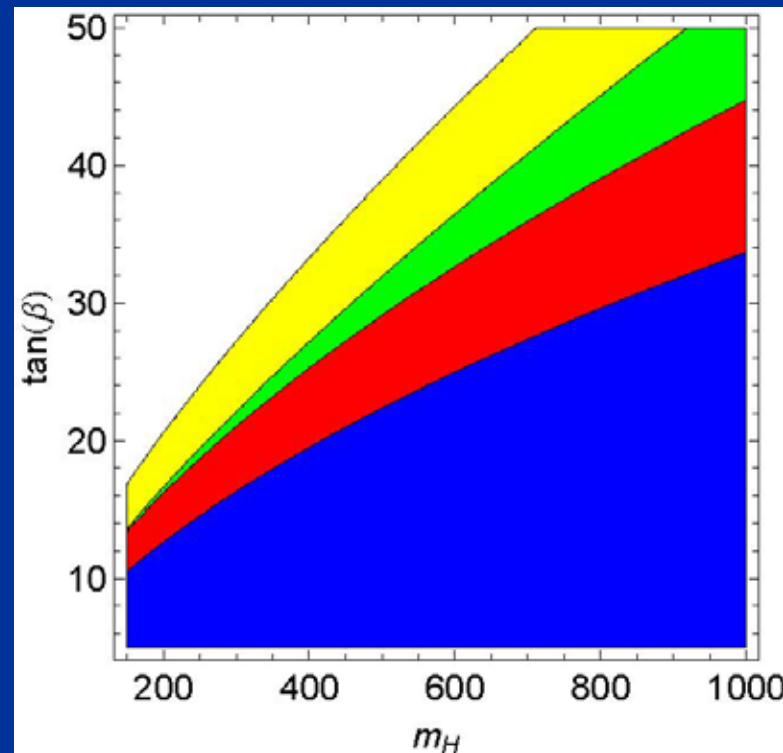
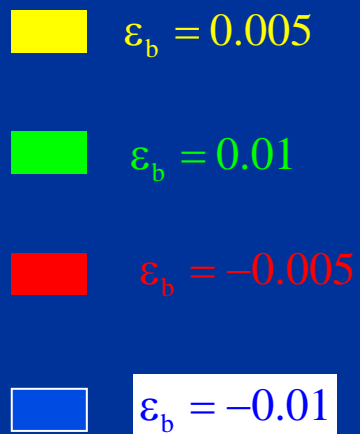
$$\text{Red box} \quad m_{g\%} = 1000 \text{ GeV}$$

$$\text{Blue box} \quad m_{g\%} = 500 \text{ GeV}$$

$B_{s,d} \rightarrow \mu^+ \mu^-$

- Neutral Higgs contribution dominant for large $\tan(\beta)$ (non-decoupling)
- Even for MFV sizable contributions

Allowed regions for RFV



$B_{s,d}$ mixing

- At low $\tan(\beta)$ constraints on $\delta_{13,23}^{dRR}, \delta_{13,23}^{dLL}$
- Double-Higgs penguin contributions at large $\tan(\beta)$
(if the Peccei-Quinn symmetry is broken)

Kaon and D mixing

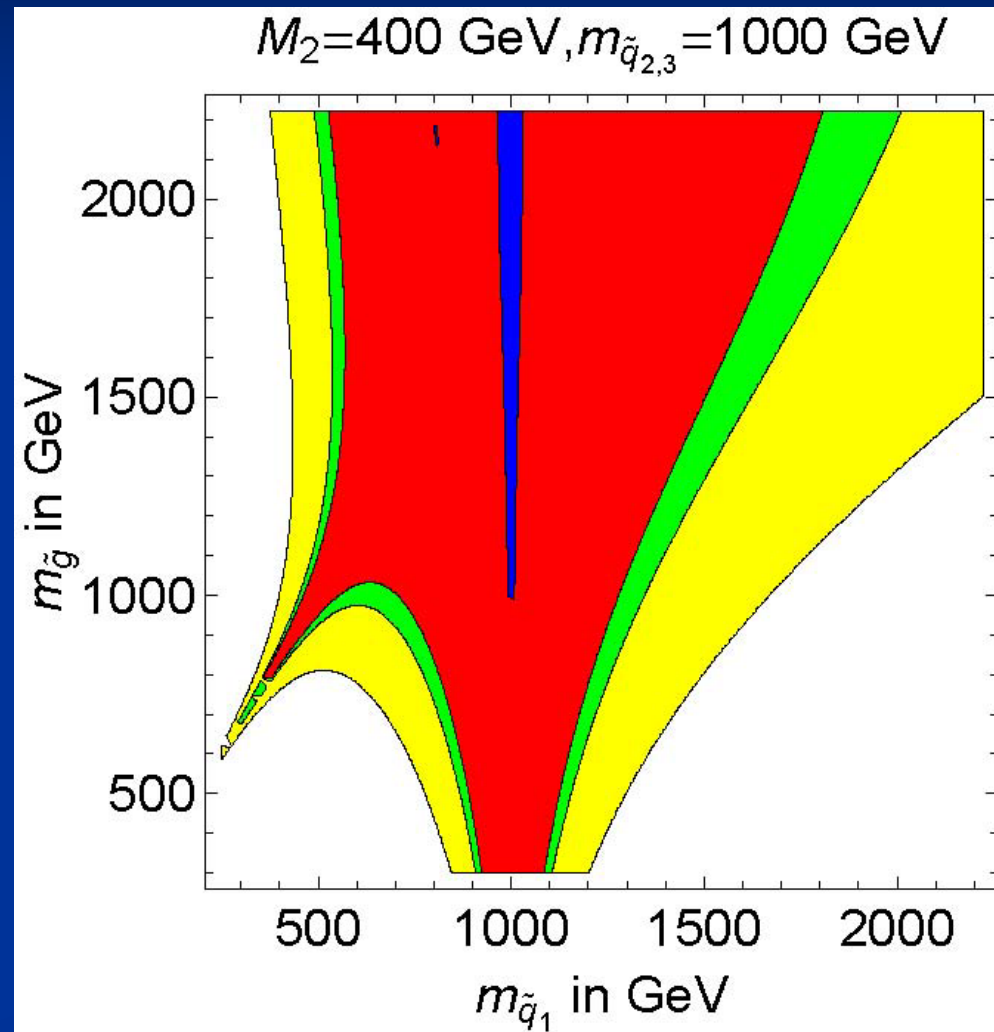
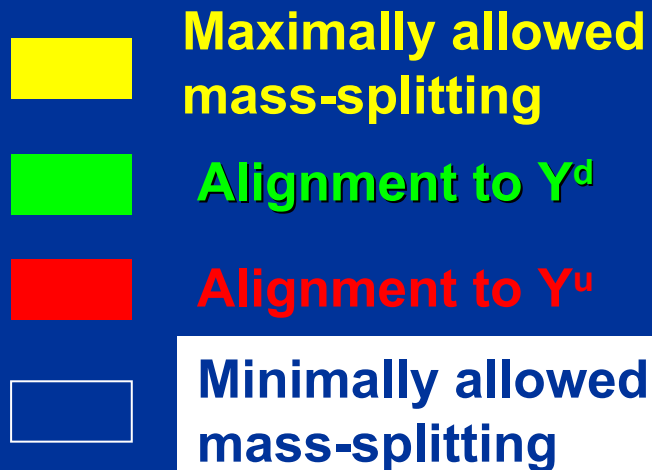
- Stringent constraints on $\delta_{12,21}^{dLR}, \delta_{12,21}^{uLR}, \delta_{12}^{dRR}, \delta_{12}^{uRR}, \delta_{12}^{dLL}, \delta_{12}^{uLL}$
- Constraints on the mass splitting of left-handed squarks.
Because of the SU(2) relation

$$M_{LL}^{\frac{q}{2}} = V^\dagger M_{LL}^{\frac{q}{2}} V$$

effects in Kaon and D mixing cannot be simultaneously avoided for non-degenerate squarks.

Allowed mass splitting

- Non-degenerate squark masses are allowed.
- More space for models with abelian flavor symmetries.
- Interesting for LHC benchmark scenarios.



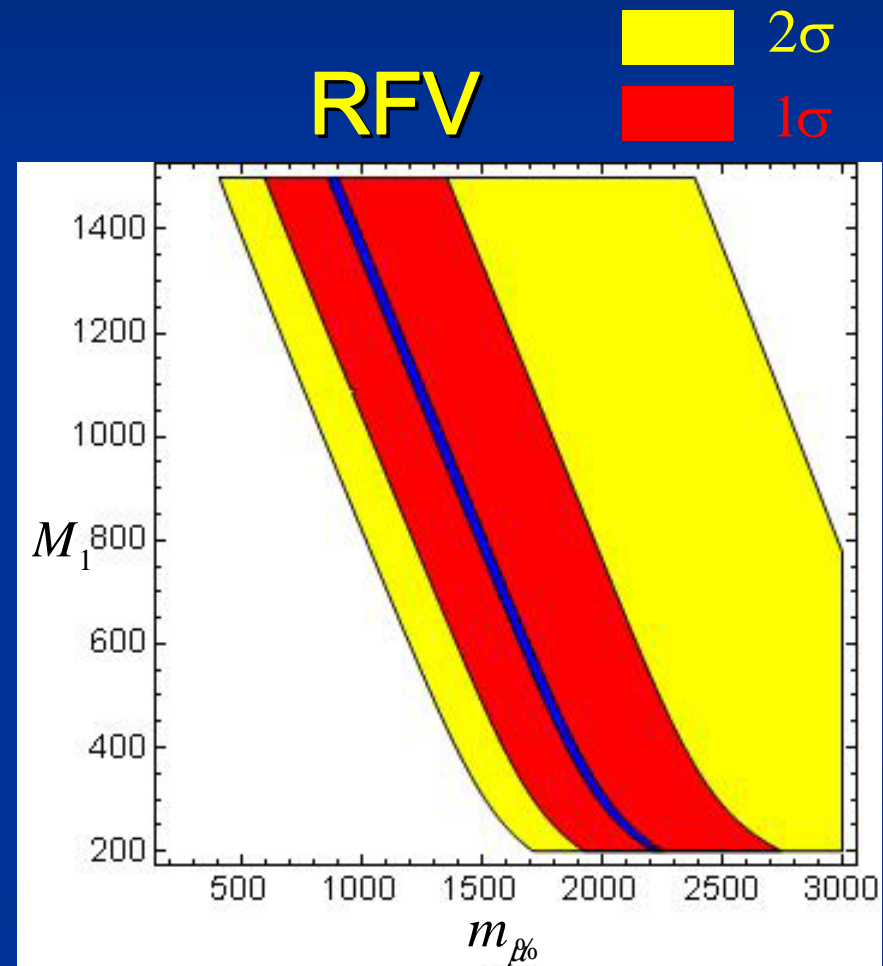
Anomalous magnetic moments and EDMs

- Constraints on

$$\text{Re}\left[\delta_{11,22}^{l\text{ LR}}\right], \text{Re}\left[\delta_{11}^{d,u\text{ LR}}\right]$$

$$\text{Im}\left[\delta_{11,22}^{l\text{ LR}}\right], \text{Im}\left[\delta_{11}^{d,u\text{ LR}}\right]$$

- The deviation from the SM expectation of the anomalous magnetic moment of the muon can be explained with large $\tan(\beta)$ or large A-terms.

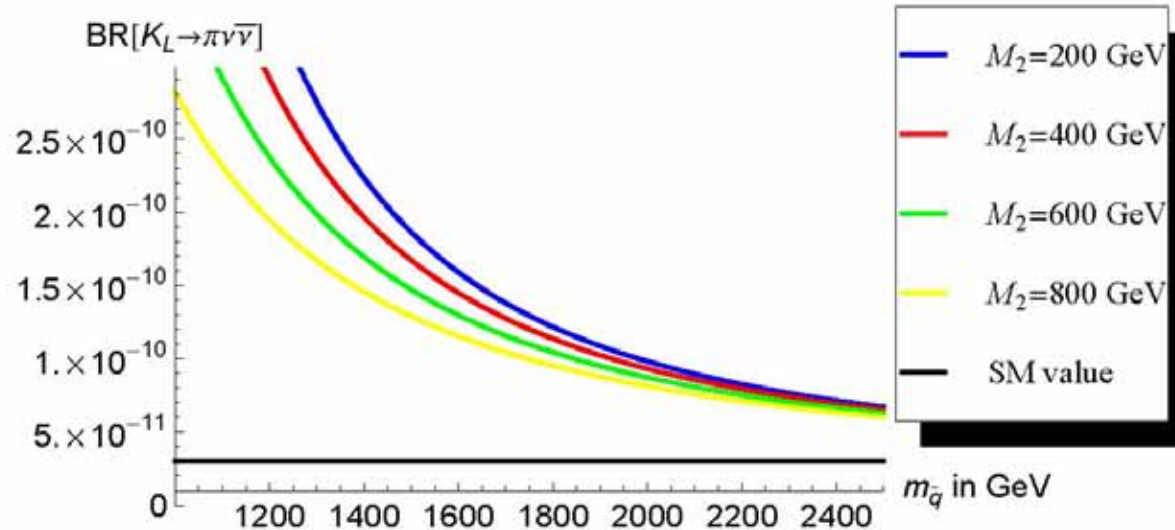
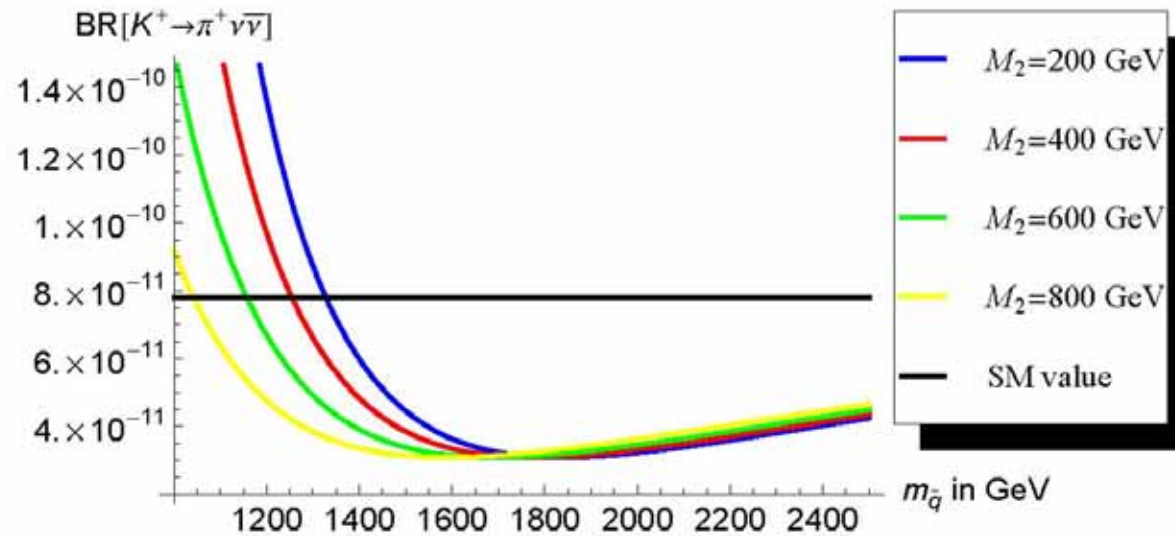


Where are large effects still possible?

$$\delta_{13}^{u LR} \times \delta_{23}^{u LR}$$

- RFV-effects in $K \rightarrow \pi \nu \bar{\nu}$

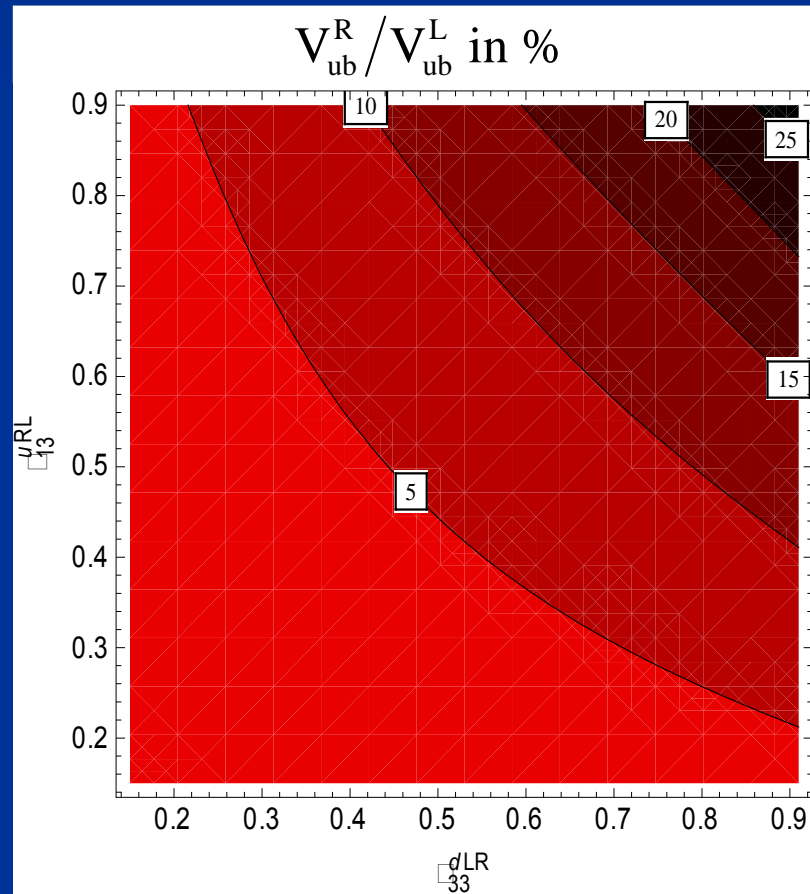
- Verifiable predictions for NA62



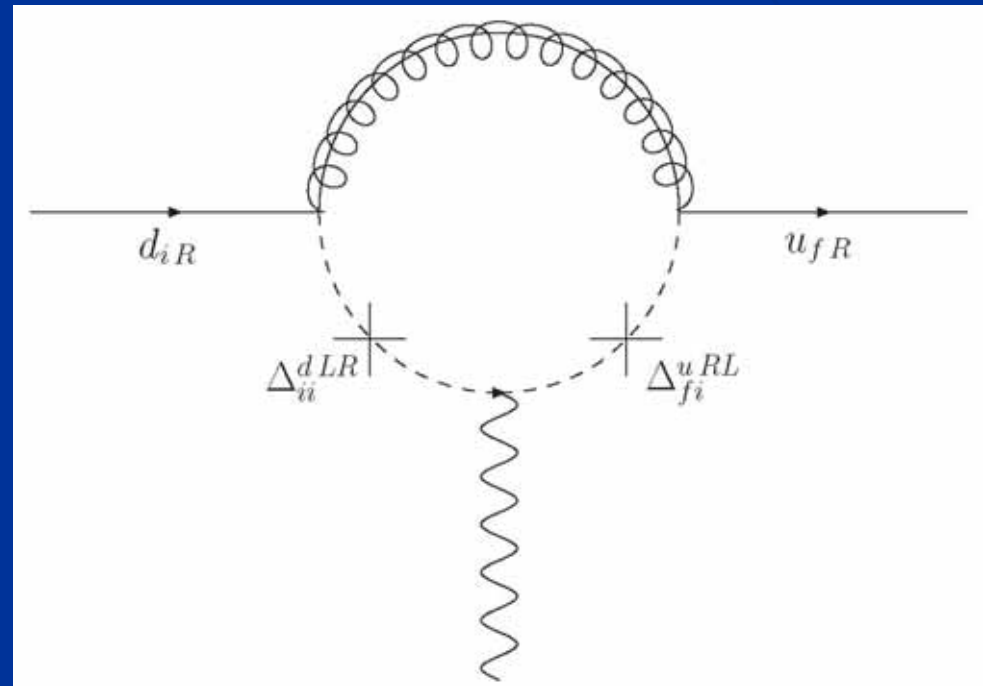
Right-handed W-coupling

AC, arXiv:0907.2461

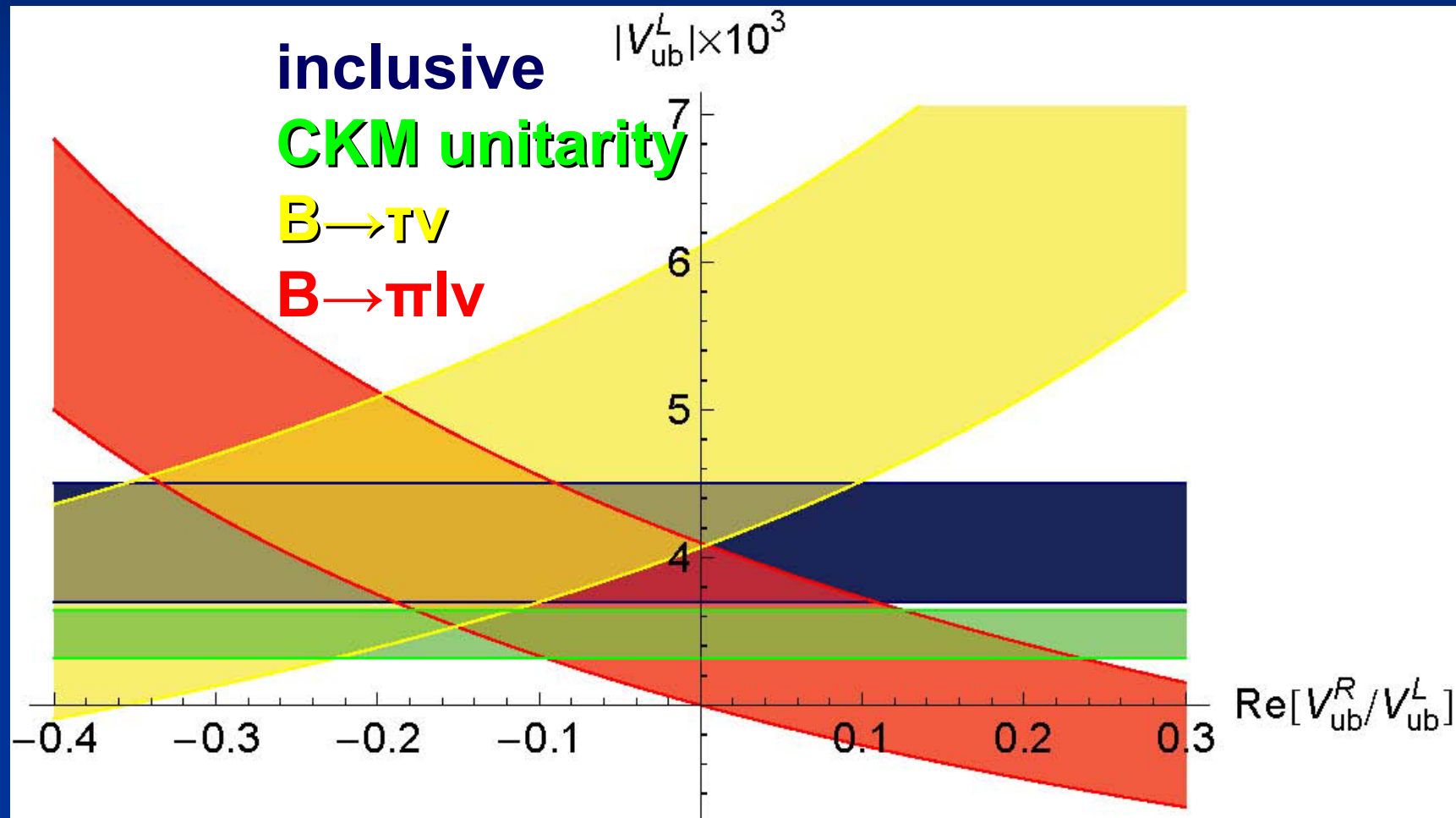
$\delta_{31}^{uLR} \times \delta_{33}^{dLR}$ generates V_{ub}^R



$$-i \frac{g_W}{\sqrt{2}} \gamma^\mu \left(P_L V_{fi}^L + P_R V_{fi}^R \right)$$



Determination of V_{ub}



SUSY_FLAVOUR v2.0

A.C., Janusz Rosiek, arXiv:1203.XXXX

Calculates in the generic MSSM:

- EDMs and anomalous magnetic moments
- $\bar{D} - D, \bar{B}_{s,d} - B_{s,d}, \bar{K} - K$
- $b \rightarrow s\gamma, \mu \rightarrow e\gamma, \tau \rightarrow \mu\gamma, \tau \rightarrow e\gamma$
- $K_L \rightarrow \pi\nu\bar{\nu}, K^+ \rightarrow \pi^+\nu\bar{\nu}$
- $B_{s,d} \rightarrow 1^+1^-, K_L \rightarrow 1^+1^-$
- $B \rightarrow \tau\nu, B \rightarrow D\tau\nu$

Including the important resummation of all chirally enhanced effects A.C., L. Hofer, J. Rosiek arXiv:1103.4272

Conclusions

- The flavour structure of the MSSM is constrained from many flavour-observables.
- Except $\delta_{13,23}^{u\text{RL,RR}}$ all flavour-changing elements of the squark mass matrices are stringently constrained
- Radiative-flavour-violation is an interesting alternative to MFV.
- $\delta_{13,23}^{u\text{LR}}$ might generate large effects in $K \rightarrow \pi \nu \bar{\nu}$ or an **anomalous W-coupling** which affects the determination of V_{ub} and enhanced single top production.
- Chirally enhanced corrections are numerically important.

 SUSY_FLAVOUR v2.0 is a valuable tool for calculating FCNC processes in the MSSM