

# **Discovering Walking Technicolor at LHC**

**-- 125 GeV Techni-dilaton at LHC --**

**Koichi Yamawaki  
(KMI, Nagoya)**

**March 8, 2013 @Moriond**



# Kobayashi-Maskawa Institute (KMI)

for the Origin of Particles and the Universe

Nagoya University

Started April 2010





M. Kobayashi

# Disciples of Sakata at Nagoya



Shoichi Sakata (1911-1970)  
Nagoya Univ. Professor

Composite Model Approach



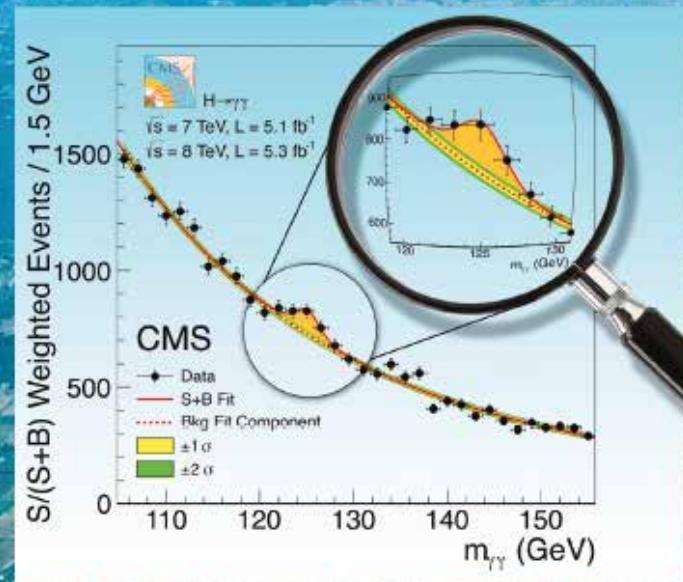
T. Maskawa

Sakata Model (1965)  
Maki-Nakagawa-Sakata  
(1962)

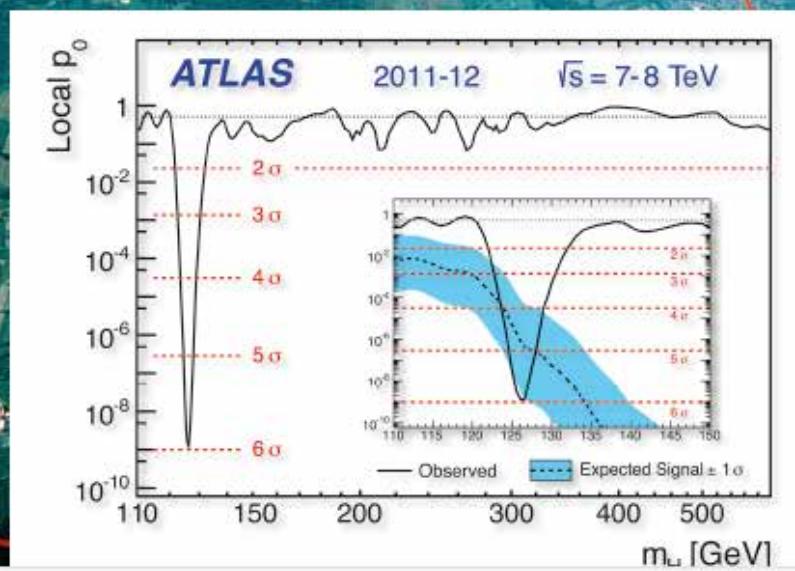
# PHYSICS LETTERS B

Available online at [www.sciencedirect.com](http://www.sciencedirect.com)

SciVerse ScienceDirect



## Discovery of 125 GeV Boson

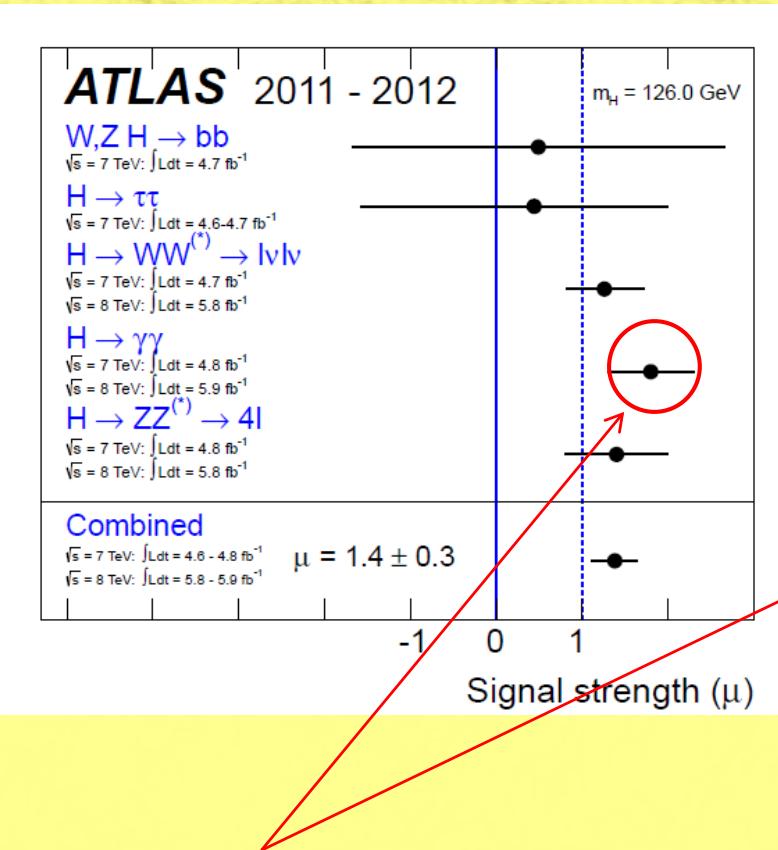


Is this the SM Scalar  
or something else?

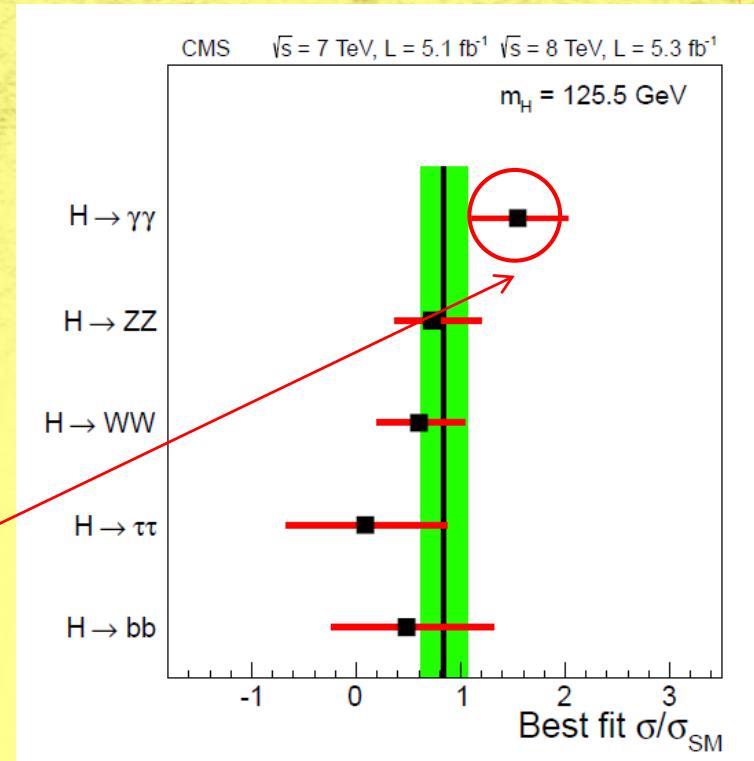
Roughly consistent with the SM Scalar,  
but .....

# The signal strengths ( $\mu = \sigma/\sigma_{SM}$ )

ATLAS (PLB 716 (2012) 1)



CMS (PLB 716 (2012) 30)



**Somewhat large diphoton event rate:**  
 $\mu$  (diphoton)  $\sim 2$  implies a “new scalar boson” (impostor)  
 beyond the SM !

# Standard Model is incomplete

- | No Dark matter candidates
- | Baryogenesis: KM CP violation not enough,  
No 1<sup>st</sup> order phase transition
- | Strong CP Problem: neutron EDM
- | ...
- | Naturalness Problem  $\longleftrightarrow$  BSM on TeV  
Technicolor (QCD – like Theory)  
Hierarchy & tachyon  
$$|\delta M_H^2| \sim \Lambda^2 \text{ composite} \pi \rightarrow m_{W,Z}$$
$$M_H^2 + \delta M_H^2 = -\mathcal{O}((10^2 \text{ GeV})^2)$$

**Folklore:**

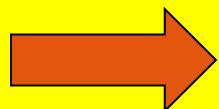
~~Technicolor = Higgsless~~

~~(No light scalar)~~

**Walking Technicolor**

KY-Bando-Matsumoto (1986)

Approx. Scale Symmetry



**Techni-dilaton**



**125 GeV Boson**

## Scale-Invariant Hypercolor Model and a Dilaton

Koichi Yamawaki, Masako Bando,<sup>(a)</sup> and Ken-iti Matumoto<sup>(b)</sup>

*Department of Physics, Nagoya University, Nagoya 464, Japan*

(Received 24 December 1985)

We propose a scale-invariant hypercolor model with a nontrivial ultraviolet fixed point having large anomalous dimension, which resolves the notorious flavor-changing neutral-current problem in hypercolor models, and at the same time predicts a  $J^{PC} = 0^{++}$  Nambu-Goldstone boson (dilaton) associated with the spontaneous breakdown of the scale invariance.

%\cite{Yamawaki:1985zg}

\bibitem{Yamawaki:1985zg}

K.~Yamawaki, M.~Bando and K.~Matumoto,

``Scale Invariant Technicolor Model and a Technidilaton,''

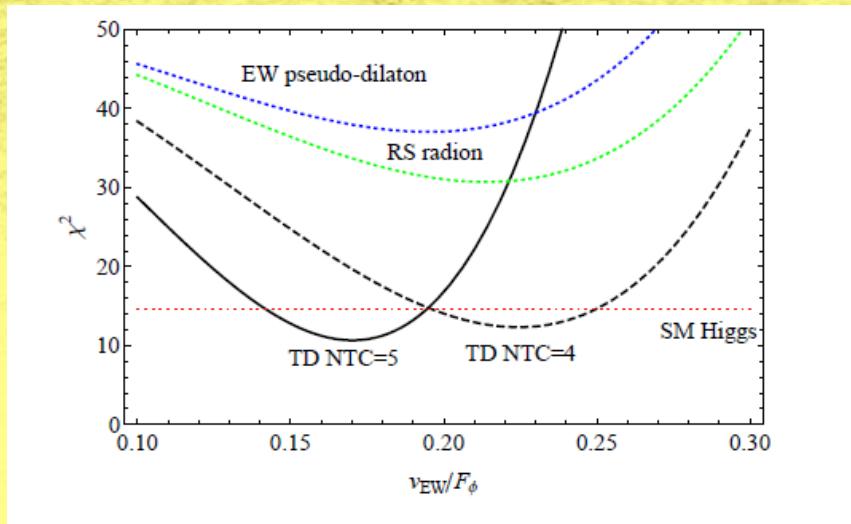
Phys.\ Rev.\ Lett.\ {\bf 56}, 1335 (1986).

%%CITATION = PRLTA,56,1335;%%

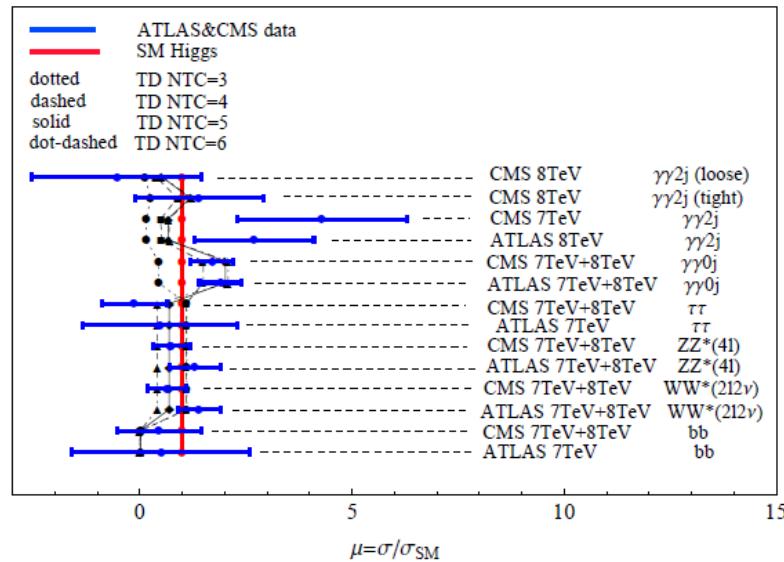
%609 citations counted in INSPIRE as of 07 Mar 2013

# 125 GeV Tchni-dilaton(TD) at LHC

S.Matsuzaki and K. Y. ,  
 PLB719 (2013) 378  
 PRD86 (2012) 115004



As of July 2012



$$\chi^2 = \sum_{i \in \text{events}} \left( \frac{\mu_i - \mu_i^{\text{exp}}}{\sigma_i} \right)^2$$

***TD (in 1FM) is favored by the current data !!***

\* ***diphoton rate enhaced by techni-fermions (> W loop contribution)***

\* ***goodness-of-fit performed for each search category***

***TD can be better than the SM Scalar***

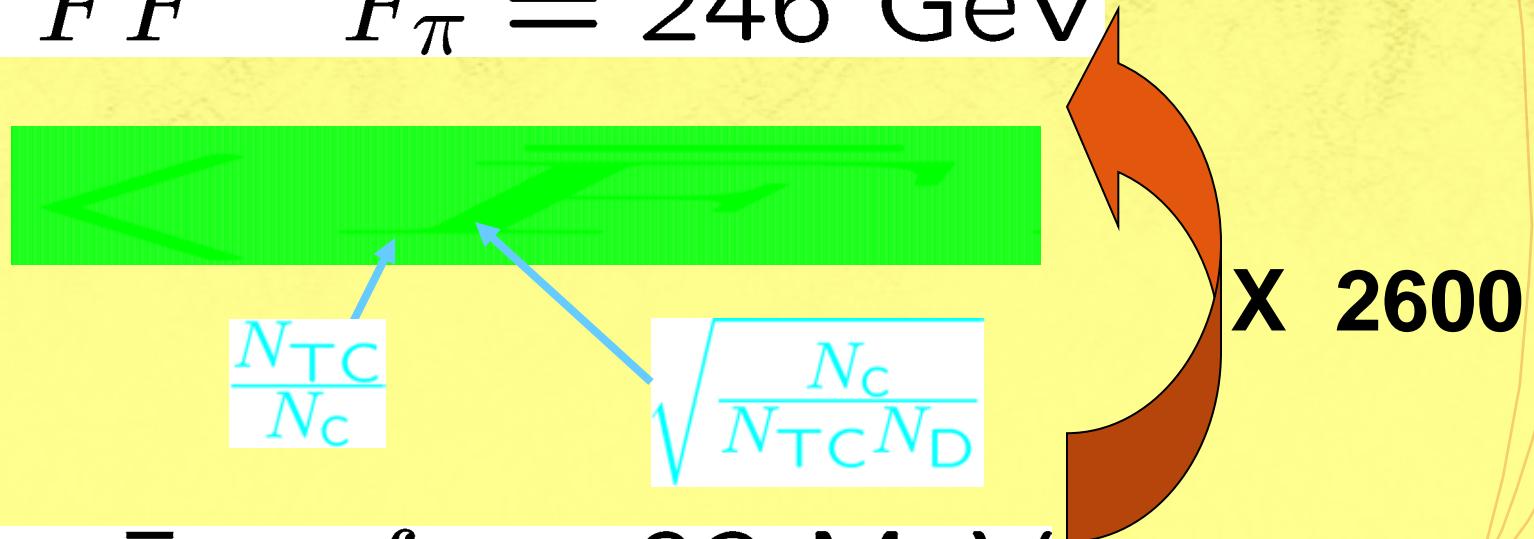
# CONTENTS

- | Technicolor: QCD-Scale-up
- | Walking Technicolor and Techni-dilaton
- | Discovering Walking Technicolor at LHC
  - Techni-dilaton at 125 GeV
- | Discovering Walking Technicolor at Lattice
  - KMI Lattice Project

# Technicolor: a Scale-Up of QCD

S. Weinberg (1976)  
L. Susskind (1979)

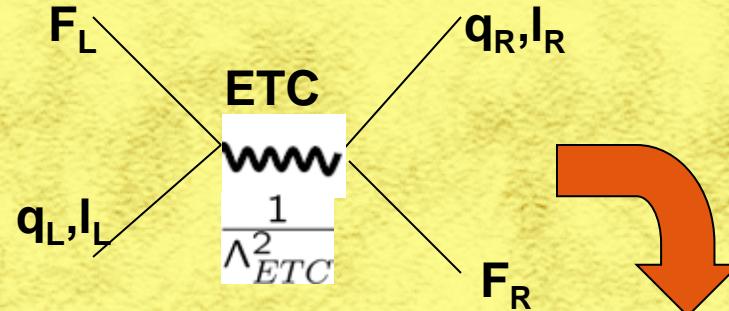
Composite  $\pi \Rightarrow$  Composite  $\pi_{\text{TC}}$   
 $H \sim \bar{F}F \quad F_\pi = 246 \text{ GeV}$   
 $\rightarrow m_{W,Z}$



$\sigma \sim \bar{q}q \quad f_\pi = 93 \text{ MeV}$

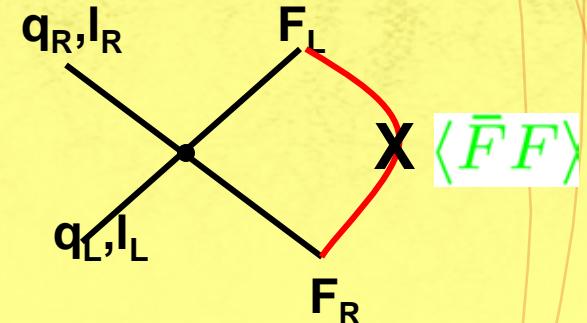


# FCNC Problems:



## Mass of Quarks/Leptons

$$m_{q/l} \sim \frac{1}{\Lambda_{ETC}^2} \langle \bar{F} F \rangle$$



FCNC

$$\frac{1}{\Lambda_{ETC}^2} \bar{s}d\bar{s}d < (10^3 \text{ TeV})^{-2}$$

$$m_s < (10^3 \text{ TeV})^{-2} \times (0.7 \text{ TeV})^3 \sim 10^{-1} \text{ MeV}$$

Needs  $10^3$  enhancement

# By Large Anomalous Dimension $\gamma_m$

**Holdom (1981)**

Pure Assumption of  
Existence of Large  $\gamma_m$   
No Concrete Dynamics  
No Concrete Value  $\gamma_m$

$$m_{q/l} = \frac{1}{\Lambda_{\text{ETC}}^2} \langle (\bar{T}T)_{\Lambda_{\text{ETC}}} \rangle$$

$$\langle \bar{F}F \rangle|_{\Lambda_{ETC}} = Z_m^{-1} \cdot \langle \bar{F}F \rangle|_{\Lambda_{EW}}$$

$$Z_m^{-1} = (\Lambda_{\text{ETC}}/\Lambda_{\text{EW}})^{\gamma_m} \simeq (10^3)^{\gamma_m}$$

$$\gamma_m > 1 \quad \longrightarrow \quad > 10^3$$

# Walking Technicolor

K.Y., Bando, Matumoto (Dec. 24, 1985)

## Ladder Schwinger-Dyson Equation

Scale Invariance  $\Leftarrow (\alpha(p) = \text{constant})$

$\gamma_m = 1$             FCNC Sol.

Techni-dilaton

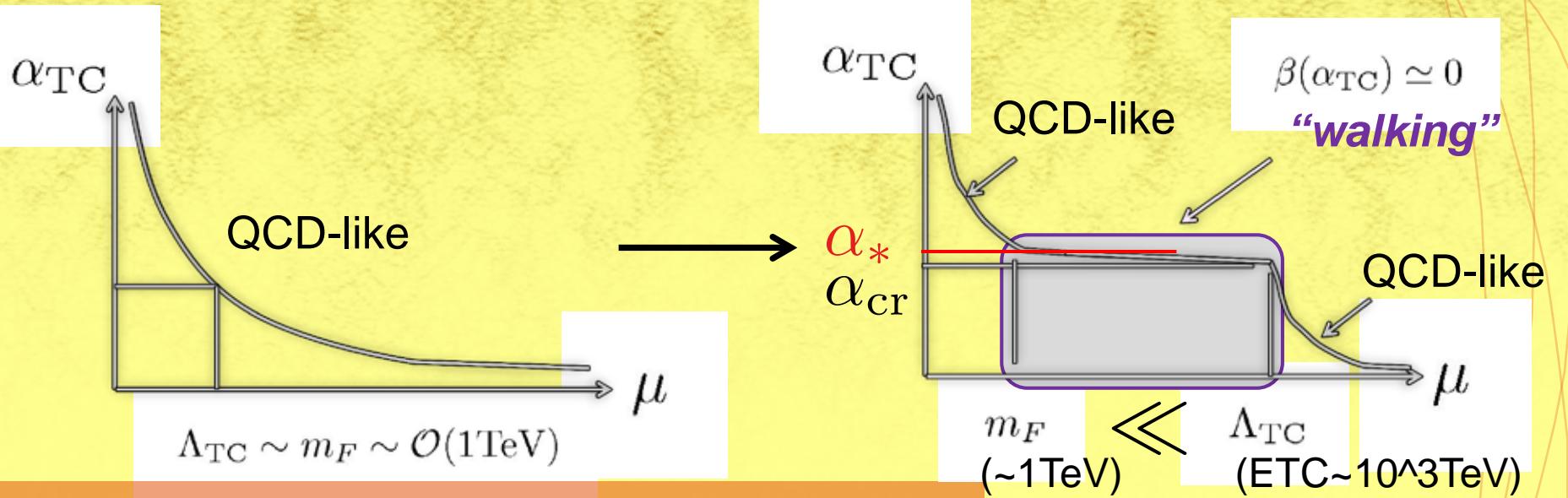
Similar FCNC Sol. without notion of  $\gamma_m$ , Scale Invariance, Techni-dilaton :

Akiba, Yanagida (Jan. 3, 1986)

Appelquist, Karabali, Wijewardhana (June 2, 1986)

( Holdom (Oct. 12, 1984), pure numerical )

# ★ A schematic view of Walking TC



\*Dynamical TF mass generation by WTC

$$m_F \simeq \Lambda_{\text{TC}} \cdot \exp \left( -\frac{\pi}{\sqrt{\alpha/\alpha_{\text{cr}} - 1}} \right) \ll \Lambda_{\text{TC}}$$

**nonperturbative  
scale anomaly  
due to  $m_F$**

$$\langle \partial_\mu D^\mu \rangle = \frac{\beta(\alpha)}{4\pi^2} \langle \alpha G_{\mu,\nu}^2 \rangle \ll m_F^4 \quad (\ll \Lambda_{\text{TC}}^4)$$

**Pseudo NG Boson: Techni-dilaton**

# ★ 125 GeV Techni-dilaton at LHC

S. Matsuzaki and K.Y.

PRD85 (2012) 095020 (Ladder)

PRD86 (2012) 035025 (Basic Formulation, Ladder)

PLB719 (2013) 378 (Ladder, Latest fit)

PRD86 (2012) 115004 (Holography, Latest fit)

*Model-Independent Result:*

*Just a simple scaling from the SM Scalar:*

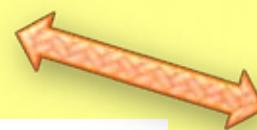
$$\frac{g_{\phi WW/ZZ}}{g_{h_{\text{SM}}WW/ZZ}} = \frac{v_{\text{EW}}}{F_\phi},$$
$$\frac{g_{\phi ff}}{g_{h_{\text{SM}}ff}} = \frac{v_{\text{EW}}}{F_\phi}, \quad \text{for } f = t, b, \tau.$$

*Model-Dependence:*

$\phi\gamma\gamma, \phi gg$  and  $F_\phi$  depending  
on particle contents of WTC models.

$$\frac{v_{\text{EW}}}{F_\phi} = \mathcal{O}(10^{-1})$$

Suppression



Enhancement !

$$\mathcal{L}_{\phi\gamma\gamma, gg} = \frac{\phi}{F_\phi} \left[ \frac{\beta_F(e)}{2e^3} F_{\mu\nu}^2 + \frac{\beta_F(g_s)}{2g_s^3} G_{\mu\nu}^2 \right]$$

$\beta_F$ : TF-loop contribution  
to beta function

We consider the one-family WTC model

# One-doublet model (1DM)

Weinberg(1976), Susskind(1979)

$TF_{EW}$	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$
$\begin{pmatrix} U \\ D \end{pmatrix}_L$	1	2	0
$U_R$	1	1	$1/2$
$D_R$	1	1	$-1/2$

# One-family model (1FM)

Farhi-Susskind (1981)

$TF_{EW}$	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$
$Q_L = \begin{pmatrix} U \\ D \end{pmatrix}_L$	3	2	$1/6$
$L_L = \begin{pmatrix} N \\ E \end{pmatrix}_L$	1	2	$-1/2$
$U_R$	3	1	$2/3$
$D_R$	3	1	$-1/3$
$N_R$	1	1	0
$E_R$	1	1	-1

Total # of techni-fermions

$$N_{\text{TF}} = (N_{\text{TF}})_{\text{EW-singlet}} + 2N_D$$

w/ critical # for mass generation  
in WTC

$$N_{\text{TF}} \simeq 4N_{\text{TC}}$$

Appelquist et al (1996)

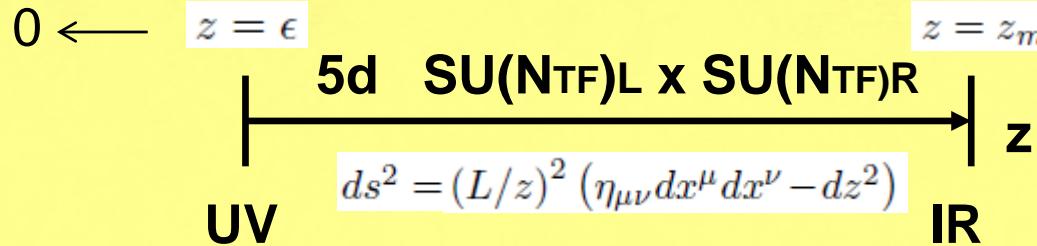
4 weak-doublets

# Holographic estimate w/ techni-gluonic effects

Haba-Matsuzaki-KY, PRD82 (2010) 055007  
 Matsuzaki- K.Y., PRD86 (2012) 115004

$z_m, \xi, G$

- \* Ladder approximation : gluonic dynamics is neglected
- \* Deformation of successful AdS/QCD model (Bottom-up approach)  
Da Rold and Pomarol (2005); Erlich, Katz, Son and Stephanov (2005)
- incorporates nonperturbative gluonic effects



$$S_5 = \int d^4x \int_{\epsilon}^{z_m} dz \sqrt{-g} \frac{1}{g_5^2} e^{cg_5^2 \Phi_X(z)} \left( -\frac{1}{4} \text{Tr} [L_{MN} L^{MN} + R_{MN} R^{MN}] \right. \\ \left. + \text{Tr} [D_M \Phi^\dagger D^M \Phi - m_\Phi^2 \Phi^\dagger \Phi] + \frac{1}{2} \partial_M \Phi_X \partial^M \Phi_X \right)$$

$$m_\Phi^2 = -(3 - \gamma_m)(1 + \gamma_m)/\tilde{L}^2$$

QCD      WTC

$$\gamma_m = 0$$

$$\gamma_m = 1$$



# \* QCD-fit w/ $\gamma_m \simeq 0$

input

$$\begin{aligned} f_\pi &= 92.4 \text{ MeV} \\ M_\rho &= 775 \text{ MeV} \\ \langle \alpha G \mu u^2 \rangle / \pi &= 0.012 \text{ GeV}^4 \end{aligned}$$

fix  $\longrightarrow$

model parameters

$$\begin{aligned} \xi &= 3.1 \\ G &= 0.25 \\ zm^{-1} &= 347 \text{ MeV} \end{aligned}$$

Model predictions

		measured
Ma1	[a1 meson]	: 1.3 GeV
Mf <sub>0</sub> (1370)	[qqbar bound state]	: 1.2 GeV
M <sub>G</sub>	[glueball ]	: 1.3 GeV
S = - 16 π L <sub>10</sub>	[S parameter]	: 0.31
[- <qbar q>] <sup>(1/3)</sup> [chiral condensate]	:	277 MeV

*Monitoring QCD works well!*

# \*WTC-case with $\gamma_m = 1$

$$G \sim \frac{\langle \alpha G_{\mu\nu}^2 \rangle}{F_\pi^4}$$

--- TD mass (lowest pole of dilatation current correlator)

$$\frac{M_\phi}{4\pi F_\pi} \simeq \sqrt{\frac{3}{N_{\text{TC}}}} \frac{\sqrt{3}/2}{1+G} \rightarrow 0 \text{ as } G \rightarrow \infty$$

125 GeV TD is realized by a large gluonic effect :  $G \sim 10$   
for one-family model w/  $F\pi = 123$  GeV (c.f. QCD case,  $G \sim 0.25$ )

--- TD decay constant (pole residue)

$$\begin{aligned} \frac{F_\phi}{F_\pi} &\simeq \sqrt{2N_{\text{TF}}} \cdot \sqrt{J_0^2(x) + J_1^2(x)} \Big|_{x=(M_\phi z_m) \ll 1} \\ &\simeq \underline{\sqrt{2N_{\text{TF}}}} . \end{aligned}$$

free from holographic-parameters !!

Massless NGB limit (“conformal limit”) is realized:

$$\frac{M_\phi}{F_\pi} \rightarrow 0 \quad \text{and} \quad \frac{F_\phi}{F_\pi} \rightarrow \text{finite}, \quad \text{as} \quad G \rightarrow \infty. \quad (\langle \Phi(z_m) \rangle \sim \xi \rightarrow 0)$$

in contrast to ladder approximation

# ★ Estimate of $\frac{v_{\text{EW}}}{F_\phi}$ -- Holographic approach

\* TD decay constant for the light TD case w/  $G \sim 10$ :

$$\frac{F_\phi}{F_\pi} \simeq \sqrt{2N_{\text{TF}}} \cdot \sqrt{J_0^2(x) + J_1^2(x)} \Big|_{x=(M_\phi z_m) \ll 1}$$

$$\simeq \underline{\sqrt{2N_{\text{TF}}}}.$$

**holographic-parameter free !!**

**Theoretical Uncertainties:  $1/N_{\text{TC}}$  corr. (20% ~ 30%)**

$$\frac{v_{\text{EW}}}{F_\phi} \Big|_{\text{holo}}^{+1/N_{\text{TC}}} \sim 0.2 - 0.4$$

**This is consistent with ladder estimate:**

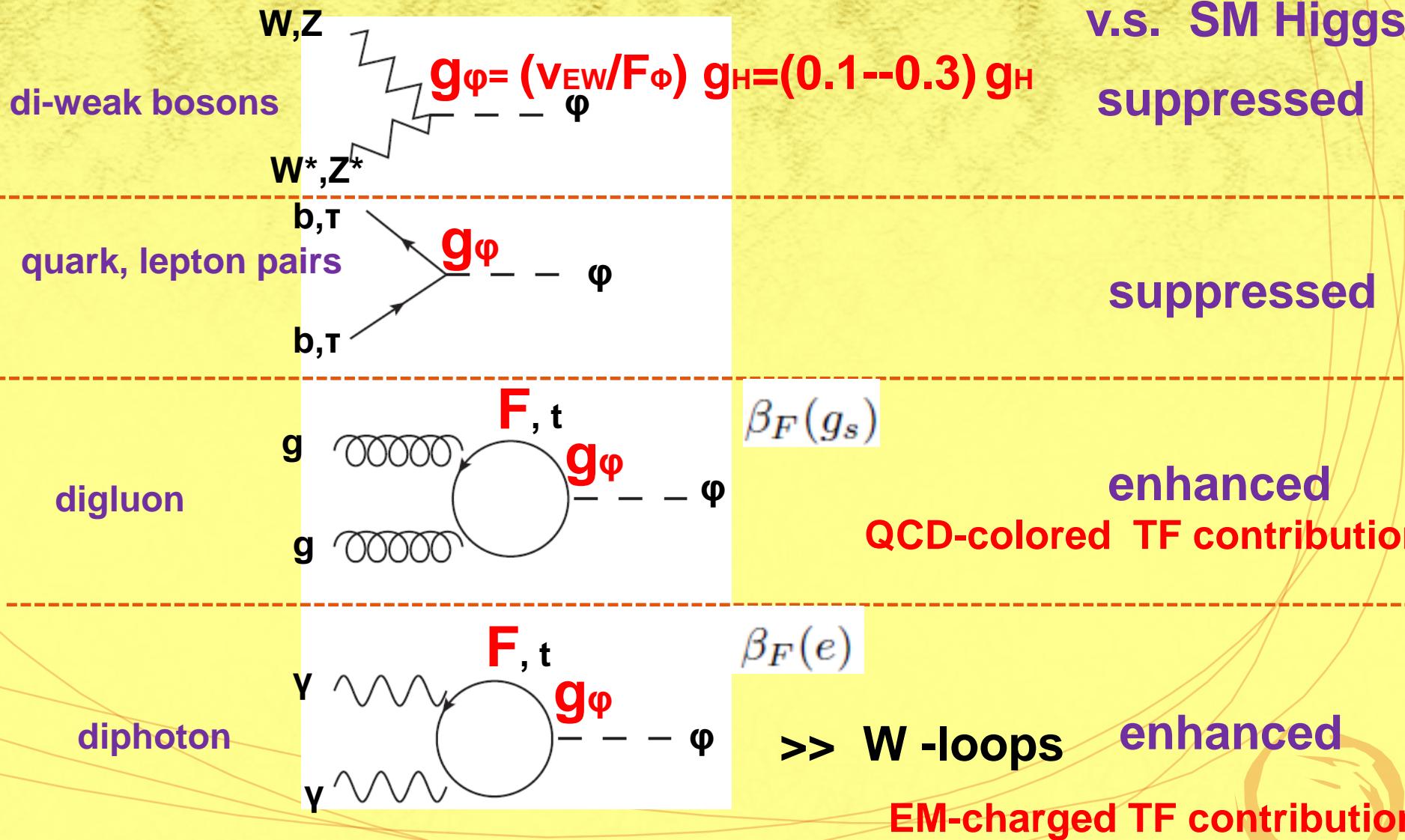
LHC best fit (to be shown)

$$\frac{v_{\text{EW}}}{F_\phi} \simeq 0.22(N_{\text{TC}} = 4)$$

$$\simeq 0.17(N_{\text{TC}} = 5)$$

$$\frac{v_{\text{EW}}}{F_\phi} \simeq \underline{(0.1 - 0.3)} \times \left( \frac{N_D}{4} \right) \left( \frac{M_\phi}{125 \text{ GeV}} \right)$$

# Characteristic features of ★ 125 GeV TD in 1FM (w/ $N_{TC}=4,5$ ) at LHC



# ★ The 125 GeV TD signal strengths

\* decays to bb, tau tau, WW\*, ZZ\*, diphoton

**bb:**

$$\begin{aligned}\mu_{bb} &= \frac{\sigma_{\text{VBA}}^\phi(s)}{\sigma_{\text{VBA}}^{h_{\text{SM}}}(s)} \frac{BR(\phi \rightarrow b\bar{b})}{BR(h_{\text{SM}} \rightarrow b\bar{b})} \\ &= \frac{\sigma_{W\phi}(s) + \sigma_{Z\phi}(s)}{\sigma_{Wh_{\text{SM}}}(s) + \sigma_{Zh_{\text{SM}}}(s)} \frac{BR(\phi \rightarrow b\bar{b})}{BR(h_{\text{SM}} \rightarrow b\bar{b})}\end{aligned}$$

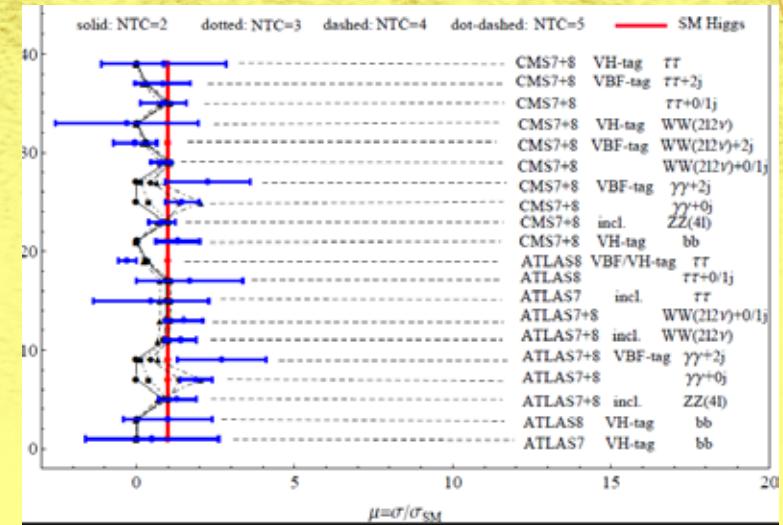
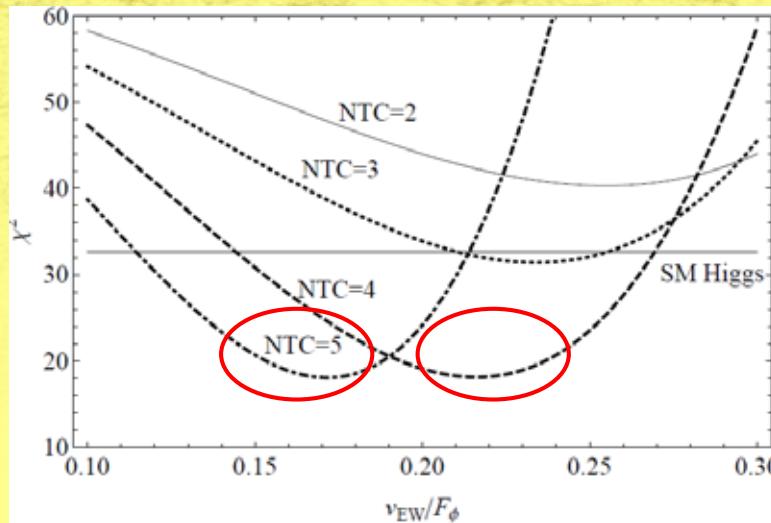
X= tautau, WW\*, ZZ\*:

$$\mu_X = \frac{\sigma_{\text{GF}}^\phi(s) + \sigma_{\text{VBF}}^\phi(s)}{\sigma_{\text{GF}}^{h_{\text{SM}}}(s) + \sigma_{\text{VBF}}^{h_{\text{SM}}}(s)} \frac{BR(\phi \rightarrow X)}{BR(h_{\text{SM}} \rightarrow X)}$$

Diphoton:

$$\begin{aligned}\mu_{\gamma\gamma 0j} &= \frac{\sigma_{\text{GF}}^\phi(s)}{\sigma_{\text{GF}}^{h_{\text{SM}}}(s)} \frac{BR(\phi \rightarrow X)}{BR(h_{\text{SM}} \rightarrow X)}, \\ \mu_{\gamma\gamma 2j} &= \frac{\xi_{\text{GF}} \cdot \sigma_{\text{GF}}^\phi(s) + \xi_{\text{VBF}} \cdot \sigma_{\text{VBF}}^\phi(s)}{\xi_{\text{GF}} \cdot \sigma_{\text{GF}}^{h_{\text{SM}}}(s) + \xi_{\text{VBF}} \cdot \sigma_{\text{VBF}}^{h_{\text{SM}}}(s)} \\ &\quad \times \frac{BR(\phi \rightarrow \gamma\gamma)}{BR(h_{\text{SM}} \rightarrow \gamma\gamma)},\end{aligned}$$

# The 125 GeV TD signal fitting \*updated after HCP2012 to the current Higgs search data



$$\chi^2 = \sum_{i \in \text{events}} \left( \frac{\mu_i - \mu_i^{\text{exp}}}{\sigma_i} \right)^2$$

$N_{TC}$	$[v_{EW}/F_\phi]_{best}$	$\chi^2 \text{ min } / \text{d.o.f.}$
4	0.22	18/19 = 0.95
5	0.17	18/19 = 0.95

\* TD can be better than the SM Scalar( $\chi^2/\text{d.o.f.} = 33/20 = 1.6$ ), due to the enhanced diphoton rate, by extra BSM (TF) contributions!

## \* The TD characteristic signal strengths for each category

$\mu_{zz} = 0.7 \text{ -- } 1.0$  (inclusive)

$\mu_{bb} = 0.006 \text{ -- } 0.01$  (VH-tag)

$\mu_{ww0j} = 0.8 \text{ -- } 1.1$  (ggF-tag)

$\mu_{ww2j} = 0.2 \text{ -- } 0.3$  (VBF-tag)

$\mu_{ww} = 0.006 \text{ -- } 0.01$  (VH-tag)

$\mu_{\tau\tau0j} = 0.8 \text{ -- } 1.1$  (ggF-tag)

$\mu_{\tau\tau2j} = 0.2 \text{ -- } 0.3$  (VBF-tag)

$\mu_{\tau\tau} = 0.006 \text{ -- } 0.01$  (VH-tag)

$\mu_{\gamma\gamma0j} = 1.4 \text{ -- } 2.0$  (ggF-tag)

$\mu_{\gamma\gamma2j} = 0.5 \text{ -- } 0.7$  (VBF-tag)

VH & VBF-tags :  
suppressed

$\gamma\gamma0j$  :  
enhanced

# Discovering the Walking Technicolor at LHC

## Road Map

- | 125 GeV Techni-dilaton vs SM Scalar:  
 $\gamma\gamma$  excess/others suppressed
- | Techni-pions around 300-500 GeV
- |  $63 - 3 = 60$  pseudo NG Bosons
- | Techni-baryon as a dark matter
- | Techni-rho/a\_1 on TeV

Jia-Matuszaki-K.Y.  
PRD 87 (2012) 016006

# Theoretical Issues

- | Walking Dynamics beyond Ladder/Holography ?
- | More Precise Quantitative Predictions?

$F_\pi, F_\phi, M_\phi, M_\rho, M_{a_1}, M_{\text{baryon}}, \text{etc.}$   
 $S, T, U$  – Parameters

Lattice !

# Discovering the Walking Technicolor on the Lattice

## KMI Lattice Project (LatKMI Collaboration)

- | Finding a candidate of WTC on the Lattice
- | Finding a light scalar composite on the Lattice
- | Calculating the composite spectra on the Lattice

# LatKMI collaboration members



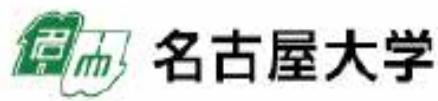
M. Kurachi T. Maskawa K. Nagai K. Yamawaki



Y. Aoki



T. Aoyama



E. Rinaldi

A. Shibata



# KMI Computer

φ

(March 02, 2011~)

## Only for Beyond SM Physics

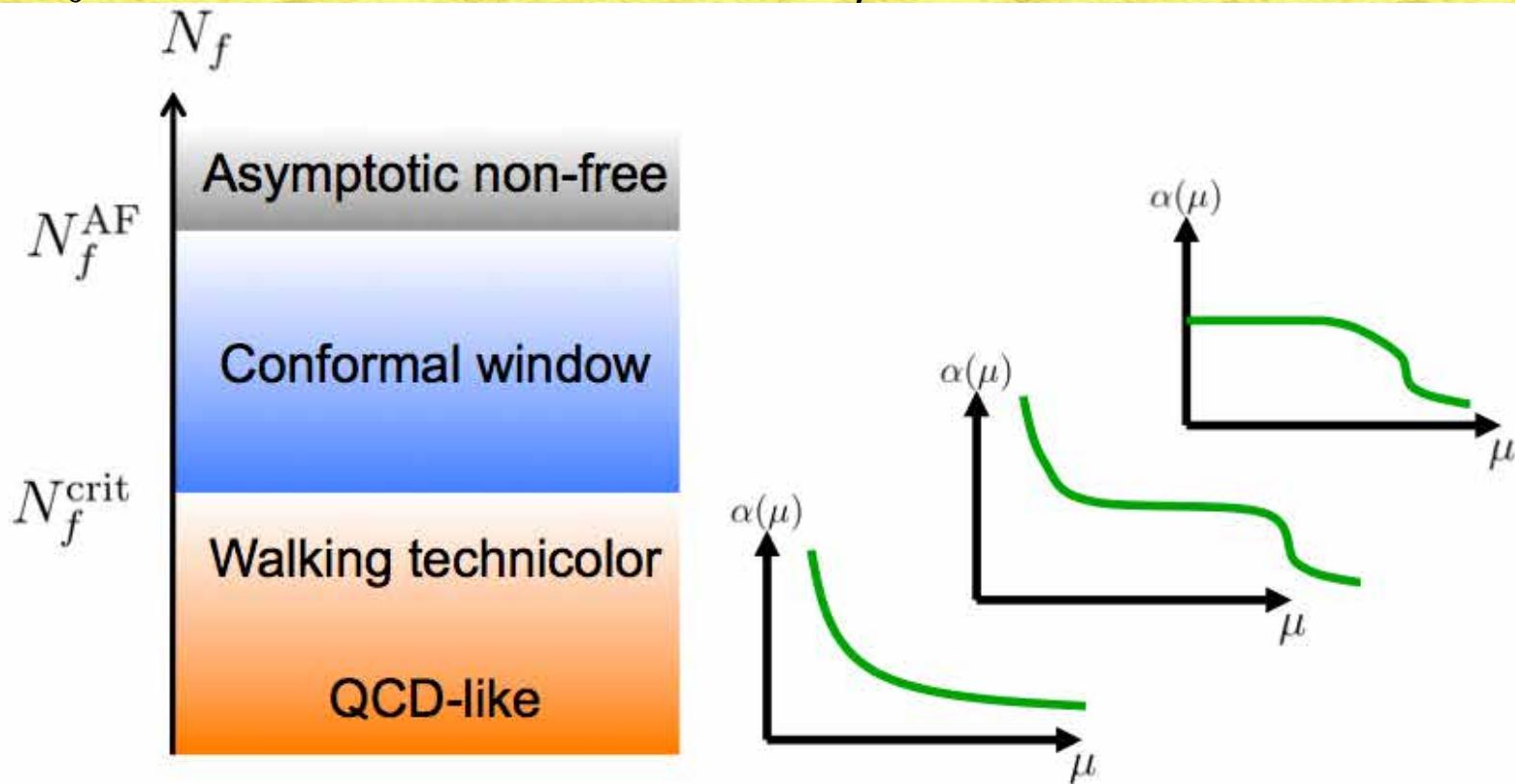
62.4 TFLOPS

26.88 TFLOPS (128 nodes)

35.53 TFLOPS (23 nodes /w GPGPU)

$SU(3); N_c = 3$

$N_f = 4, 8, 12, 16 (< N_f^{\text{AF}} = 11N_c/2 = 16.5)$



2 – loop :  $N_f^{\text{crit}} = 8.05$

2 – loop + ladder SD equation :  $N_f^{\text{crit}} = 11.9$

# The scalar spectrum of many-flavour QCD\*

Yasumichi Aoki<sup>a</sup>, Tatsumi Aoyama<sup>a</sup>, Masafumi Kurachi<sup>a</sup>, Toshihide Maskawa<sup>a</sup>, Kei-ichi Nagai<sup>a</sup>, Hiroshi Ohki<sup>a</sup>, Enrico Rinaldi<sup>a,b†</sup>, Akihiro Shibata<sup>c</sup>, Koichi Yamawaki<sup>a</sup>, Takeshi Yamazaki<sup>a</sup>

LatKMI collaboration

arXiv: 1302.4577 [hep-lat]

Preliminary

N<sub>f</sub>=8: in progress

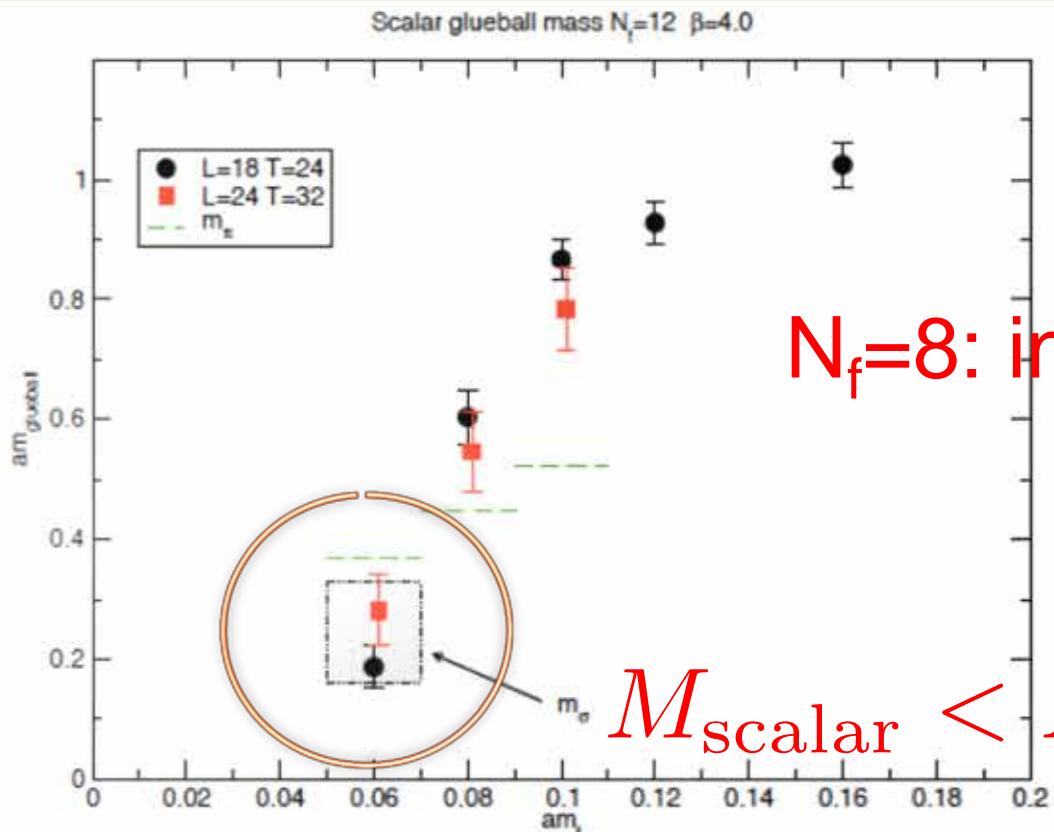


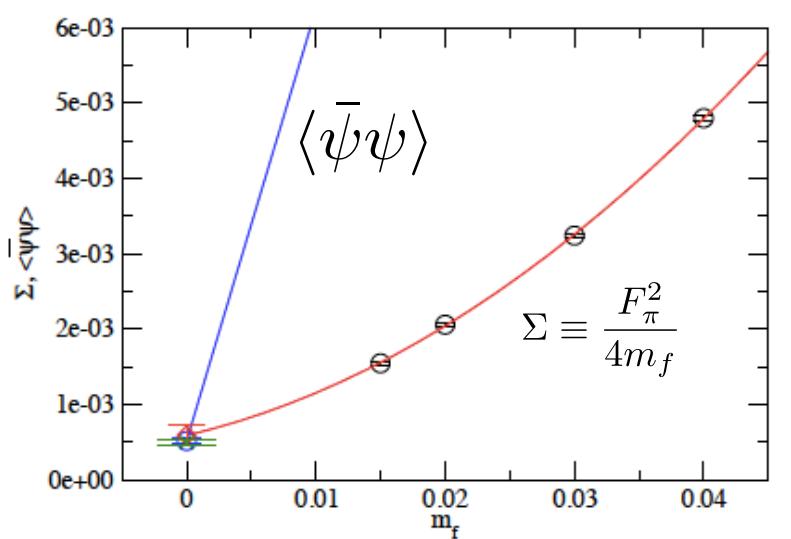
Fig. 5. Summary of the fitted scalar glueball masses for  $N_f = 12$  QCD at  $\beta = 4.0$  and for several bare fermion masses. Some values of the pion mass are shown for comparison. The point on different volumes have been slightly displaced for clarity. The grey box indicates the location

# Walking signals in $N_f = 8$ QCD on the lattice

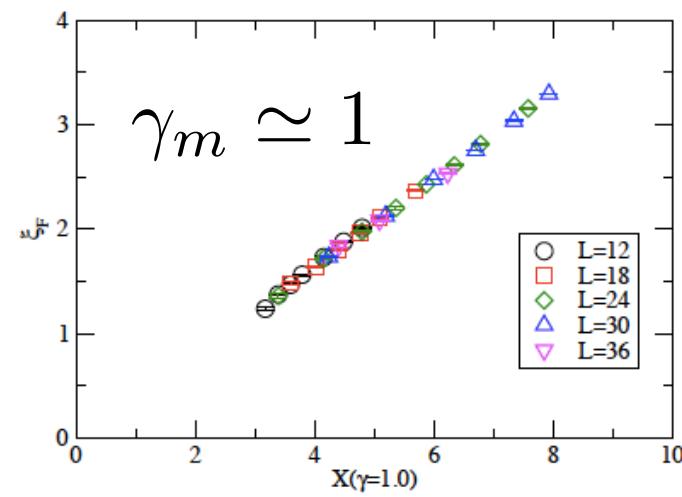
Yasumichi Aoki,<sup>1</sup> Tatsumi Aoyama,<sup>1</sup> Masafumi Kurachi,<sup>1</sup> Toshihide Maskawa,<sup>1</sup>  
 Kei-ichi Nagai,<sup>1</sup> Hiroshi Ohki,<sup>1</sup> Akihiro Shibata,<sup>2</sup> Koichi Yamawaki,<sup>1</sup> and Takeshi Yamazaki<sup>1</sup>  
 (LatKMI Collaboration)

[arXiv:1302.6859 \[hep-lat\]](https://arxiv.org/abs/1302.6859)

S<sub>XSB</sub>     $m_f = 0.015 - 0.04$



“Conformal”  $m_f = 0.05 - 0.16$



$F_\pi \rightarrow \neq 0, M_\pi \rightarrow 0, M_\rho \rightarrow \neq 0$   
 at  $m_f \rightarrow 0$

$$\xi_H \equiv LM_H = f_H(Lm_f^{\frac{1}{1+\gamma}})$$

# Thank you

# Backup Slides

# ★ Other characteristic spectrum in one-family WTC

Chiral symmetry breaking:  $SU(8)L \times SU(8)R \rightarrow SU(8)V$

3 “would-be” NGBs: eaten by W, Z  
60 (pseudo) NGB = techni-pions (TPs)

$$\theta_a^{\pm,3,0}$$

: color-octet scalars (# 32)

$$\sim \bar{Q} \gamma_5 \lambda_a \tau^{\pm,3,0} Q$$

$$T_c^{\pm,3,0}(\bar{T}_c^{\pm,3,0})$$

: color-triplet scalars (“leptoquark”) (#24)

$$\sim \bar{Q} \gamma_5 \tau^{\pm,3,0} L \text{ (h.c.)}$$

$$P^{\pm,3,0}$$

: color-singlet scalars (# 4)

$$\sim \bar{Q} \gamma_5 \tau^{\pm,3,0} Q - 3 \bar{L} \gamma_5 L$$

Expected masses :

300 ~ 500 GeV

J.Junji, S.Matsuzaki, K.Y.,  
PRD87.016006. (2012)

Techni-pions: **couplings to WW and ZZ highly suppressed** (NO NGB-NBG-NGB)

à narrow resonances (tot.width ~ 5--10GeV)

J.Junji, S.Matsuzaki, K.Y., PRD87.016006. (2012)

## Discovering isospin singlet TPs: $P^0$ and $\theta_a^0$

- \* produced ONLY from ggF :  
~~VBF, VH,~~
- \* Predominantly decaying to SM fermions

$$\sigma_{ggF}(pp \rightarrow P^0 \rightarrow \tau^+\tau^-) \sim 0.1 \text{ pb} \quad @ \sim 300 \text{ GeV}$$

$$\sigma_{ggF}(pp \rightarrow P^0 \rightarrow \bar{t}t) \sim 10 \text{ pb} \quad @ \sim 400 \text{ GeV}$$

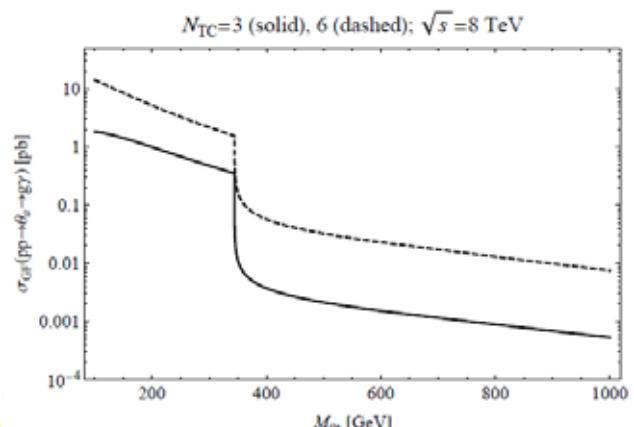
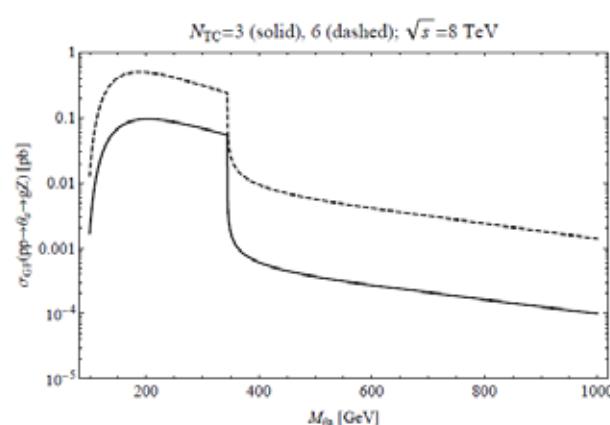
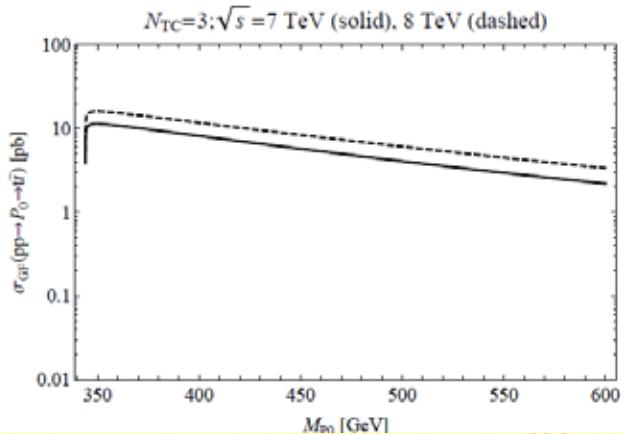
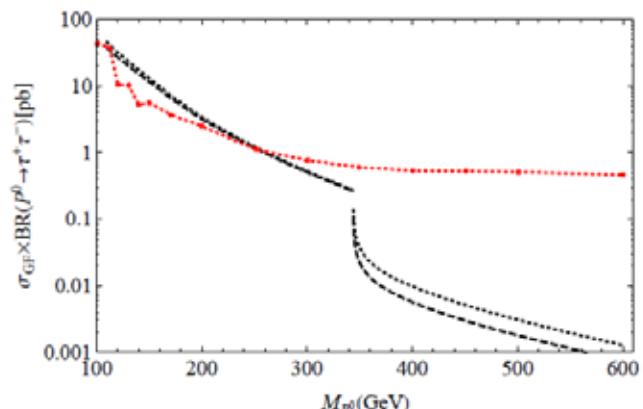
$$\sigma_{GF}(pp \rightarrow \theta_a) \times \text{BR}(\theta_a \rightarrow t\bar{t}) \simeq 60 \text{ pb} \quad @ \sim 500 \text{ GeV}$$

- \* decays to glue-gamma/Z

$$\sigma_{GF}(pp \rightarrow \theta_a \rightarrow gZ/\gamma) :$$

~ 1 fb

@ ~ 500 GeV



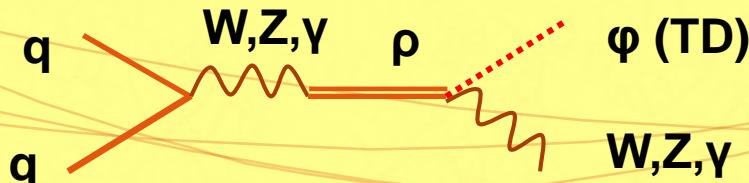
## \* One-family walking techni-rho mesons (#63)

$\rho_{\theta_a}^{\pm,3,0}$	: color-octet vectors (# 32)	Expected masses: 1 ~ 4 TeV  consistent w/ EW precision tests  S.Matsuzaki and K.Y., PRD86.115004. (2012)
$\rho_{T_c}^{\pm,3,0}$ ( $\bar{\rho}_{\bar{T}_c}^{\pm,3,0}$ )	: color-triplet vectors (# 24)	
$\rho_P^{\pm,3,0}$	: color-singlet vectors (# 4)	
$\rho_\pi^{\pm,3}$	: color-singlet vectors (# 3) [corresponding to vector states for eaten NGBs]	

- Typical discovery channels: decays to **WLWL (ZLZL) or TP and WL**



- Novel discovery channel: decays to  **$W(Z,\gamma)$  and TD!**



S.Matsuzaki and K.Y., in progress

# ★ TD phenomenological Lagrangian

S.M.atsuzaki and K. Y.,  
PRD86(2012)

## \* Nonlinear realization of scale symmetry

Nonlinear base  $\chi$ , TD field  $\Phi$

$$\chi = e^{\phi/F_\phi}, \quad \delta\chi = (1 + x^\nu \partial_\nu)\chi$$

$$\delta\phi = F_\phi + x^\nu \partial_\nu \phi$$

i) The scale anomaly-free part:

$$\mathcal{L}_{\text{inv}} = \frac{F_\pi^2}{4} \chi^2 \text{Tr}[\mathcal{D}_\mu U^\dagger \mathcal{D}^\mu U] + \frac{F_\phi^2}{2} \partial_\mu \chi \partial^\mu \chi$$

ii) The anomalous part (made invariant by including spurion field “S”):

$$\begin{aligned} \mathcal{L}_S = & -m_f \left( \left( \frac{\chi}{S} \right)^{2-\gamma_m} \cdot \chi \right) \bar{f} f \\ & + \log \left( \frac{\chi}{S} \right) \left\{ \frac{\beta_F(g_s)}{2g_s} G_{\mu\nu}^2 + \frac{\beta_F(e)}{2e} F_{\mu\nu}^2 \right\} + \dots \end{aligned}$$

iii) The scale anomaly part:

$$V_\chi = \frac{F_\phi^2 M_\phi^2}{4} \chi^4 \left( \log \chi - \frac{1}{4} \right)$$

which correctly reproduces the PCDC relation:

$$\langle \theta_\mu^\mu \rangle = -\delta_D V_\chi \Big|_{\text{vacuum}} = -\frac{F_\phi^2 M_\phi^2}{4} \langle \chi^4 \rangle \Big|_{\text{vacuum}} = -\frac{F_\phi^2 M_\phi^2}{4}$$

The TD couplings are exactly reproduced by these terms.

# ★ Ladder estimate of TD mass

$$\Lambda_{\text{TC}} (\sim \Lambda_{\text{ETC}}) \\ N_f (\rightarrow N_f^{\text{cr}})$$

- \* LSD + BS in large Nf QCD

*Harada-Kurachi-K.Y. (1989)*

- \* LSD via gauged NJL

*Shuto-Tanabashi-K.Y. (1990);  
Carena-Wagner (1992); Hashimoto (1998)*

- \* This is reflected in PCDC (partially conserved dilatation current)

$$F_\phi^2 M_\phi^2 = -4\langle\theta_\mu^\mu\rangle = \frac{\beta(\alpha)}{\alpha} \langle G_{\mu\nu}^2 \rangle \simeq 3\eta m_F^4$$

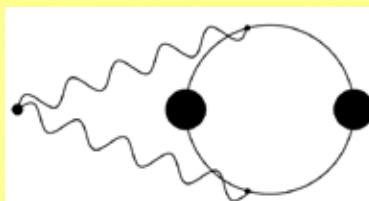
where  $\eta \simeq \frac{N_{\text{TC}} N_{\text{TF}}}{2\pi^2} = \mathcal{O}(1)$

→  $\frac{F_\phi^2}{m_F^2} \cdot \frac{M_\phi^2}{m_F^2} = \text{finite}$

A composite Higgs mass

$$M_\phi \sim 4F_\pi \ll M_\rho, M_{a_1}$$

~ 500 GeV  
for one-family model (1FM)  
still larger than ~ 125 GeV



*Miransky-Gusynin (1989);  
Hashimoto-K.Y. (2011):*

No exactly massless NGB limit:

$$M_\phi/m_F \rightarrow 0$$

only when  $F_\phi/m_F \rightarrow \infty$ , i.e., a decoupled limit.

# ★ Estimate of $\frac{v_{\text{EW}}}{F_\phi}$ : - Ladder approximation

\* PCDC (partially conserved dilatation current)

$$F_\phi^2 M_\phi^2 = -4 \langle \theta_\mu^\mu \rangle$$

$$\langle \theta_\mu^\mu \rangle = 4\mathcal{E}_{\text{vac}} = -\kappa_V \left( \frac{N_{\text{TC}} N_{\text{TF}}}{2\pi^2} \right) m_F^4$$

\* criticality condition

$$N_{\text{TF}} \simeq 4N_{\text{TC}}$$

Appelequist-Terning-Wijewardhana (1996)

$$\leftarrow N_f^{\text{cr}} \simeq 4N_c$$

\* Pagels-Stokar formula

$$F_\pi^2 = \kappa_F^2 \frac{N_{\text{TC}}}{4\pi^2} m_F^2$$

$$F_\pi = v_{\text{EW}} / \sqrt{N_D}$$

# of EW doublets

$$\frac{v_{\text{EW}}}{F_\phi} \simeq \frac{1}{8\sqrt{2}\pi} \sqrt{\frac{\kappa_F^4}{\kappa_V}} N_D \frac{M_\phi}{v_{\text{EW}}}$$

\* Recent ladder SD analysis  
(large Nf QCD)

$$\kappa_V \simeq 0.7, \quad \kappa_F \simeq 1.4$$

Hashimoto-K.Y. (2011)

## \* Theoretical uncertainties

Ladder approximation is subject to **about 30% uncertainty** for estimate of critical coupling and QCD hadron spectrum

critical coupling : T. Appelquist et al (1988);

Hadron spectrum : K.-I. Aoki et al (1991); M. Harada et al (2004).

$$\frac{N_{\text{TF}}}{4N_{\text{TC}}} \simeq 1 \pm 0.3$$

$$\langle \theta_\mu^\mu \rangle = 4\mathcal{E}_{\text{vac}} = -\frac{\kappa_V}{30\%} \left( \frac{N_{\text{TC}} N_{\text{TF}}}{2\pi^2} \right) m_F^4$$

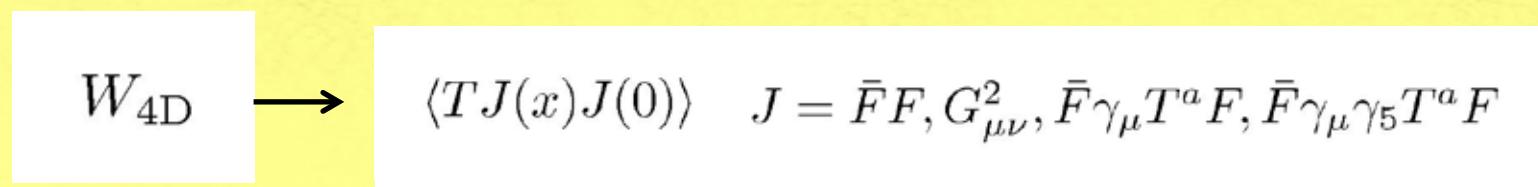
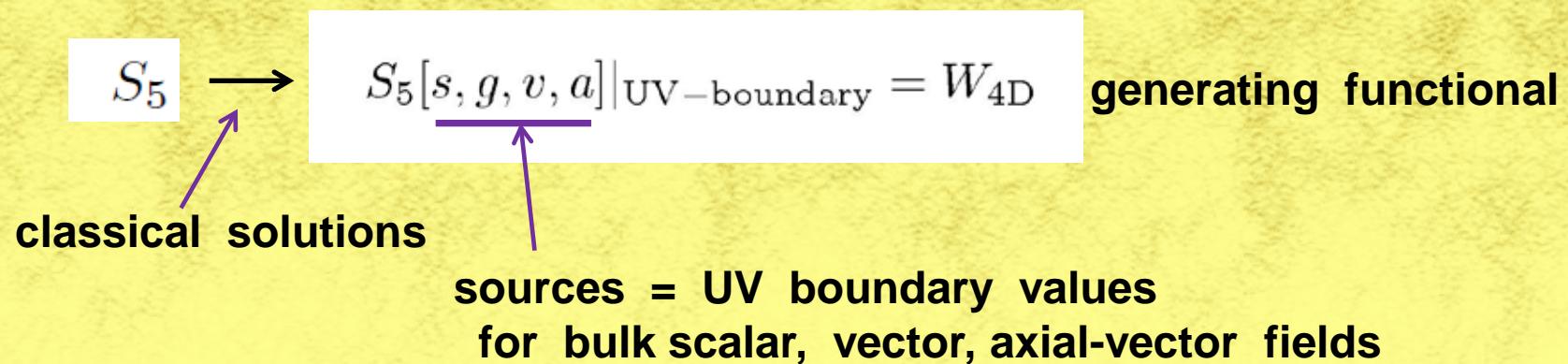
$$F_\pi^2 = \frac{\kappa_F^2}{4\pi^2} \frac{N_{\text{TC}}}{m_F^2}$$

Estimate  
w/ uncertainty included

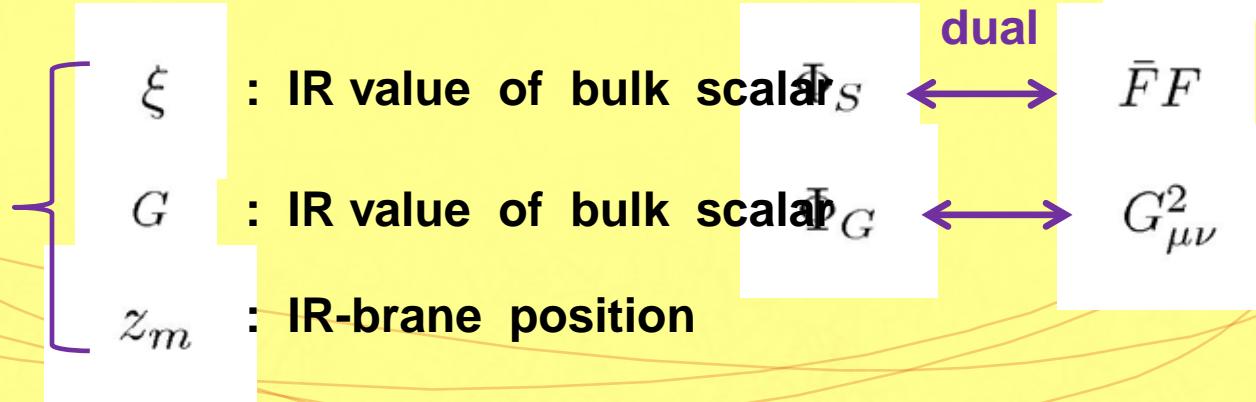
30%

$$\frac{v_{\text{EW}}}{F_\phi} \simeq (0.1 - 0.3) \times \left( \frac{N_D}{4} \right) \left( \frac{M_\phi}{125 \text{ GeV}} \right)$$

\* AdS/CFT recipe:



Current collerators  
are calculated as a function of three IR –boundary values and



$$S_5 = \int d^4x \int_{\epsilon}^{z_m} dz \sqrt{-g} \frac{1}{g_5^2} e^{cg_5^2 \Phi_X(z)} \left( -\frac{1}{4} \text{Tr} [L_{MN} L^{MN} + R_{MN} R^{MN}] \right. \\ \left. + \text{Tr} [D_M \Phi^\dagger D^M \Phi - m_\Phi^2 \Phi^\dagger \Phi] + \frac{1}{2} \partial_M \Phi_X \partial^M \Phi_X \right)$$

$$\Phi(x, z) = \frac{1}{\sqrt{2}}(v(z) + \sigma(x, z)) \exp[i\pi(x, z)/v(z)] \\ \Phi_X(z) = v_X(z),$$

**AdS/CFT dictionary:**

\* **UV boundary values = sources**

$$\alpha M = \lim_{\epsilon \rightarrow 0} Z_m \left( \frac{L}{z} v(z) \right) \Big|_{z=\epsilon}, \quad Z_m = Z_m(L/z) = \left( \frac{L}{z} \right)^{\gamma_m}$$

$$M' = \lim_{\epsilon \rightarrow 0} L v_X(z) \Big|_{z=\epsilon}$$

\* **IR boundary values:**

$$\xi = L v(z) \Big|_{z=z_m}$$



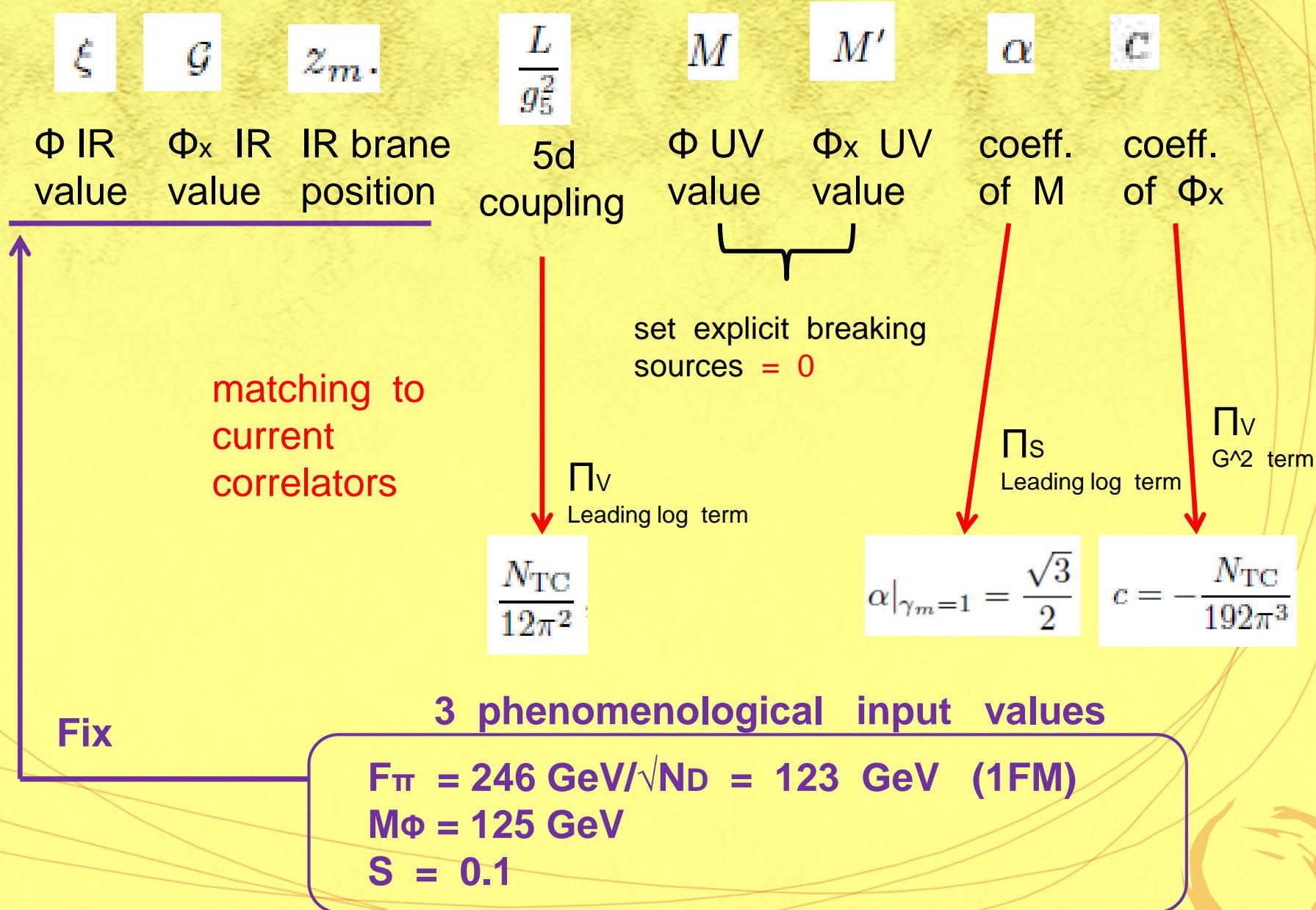
**chiral condensate**  $\langle \bar{T}T \rangle$

$$\mathcal{G} = L v_X(z) \Big|_{z=z_m}$$



**gluon condensate**  $\langle \alpha G_{\mu\nu}^2 \rangle$

## The model parameters:



## Other holographic predictions (1FM w/ S=0.1)

NTC = 3

Techni- $\rho$ , a1 masses	:	$M_\rho = M_{a1} = 3.5 \text{ TeV}$
Techni-glueball (TG) mass	:	$M_G = 19 \text{ TeV}$
TG decay constant	:	$F_G = 135 \text{ TeV}$
dynamical TF mass $m_F$	:	$m_F = 1.0 \text{ TeV}$

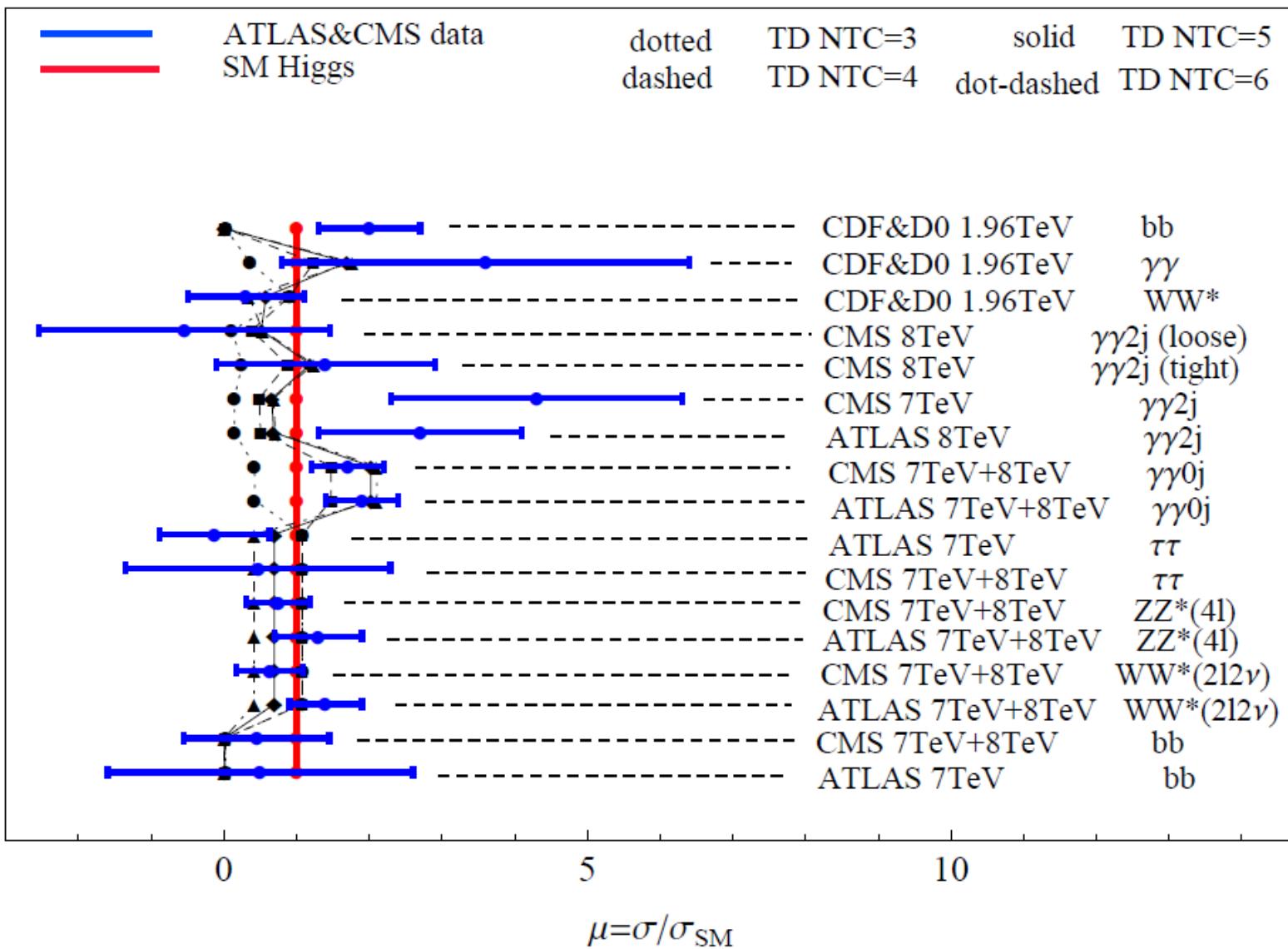
NTC = 4

Techni- $\rho$ , a1 masses	:	$M_\rho = M_{a1} = 3.6 \text{ TeV}$
Techni-glueball (TG) mass	:	$M_G = 18 \text{ TeV}$
TG decay constant	:	$F_G = 156 \text{ TeV}$
dynamical TF mass $m_F$	:	$m_F = 0.95 \text{ TeV}$

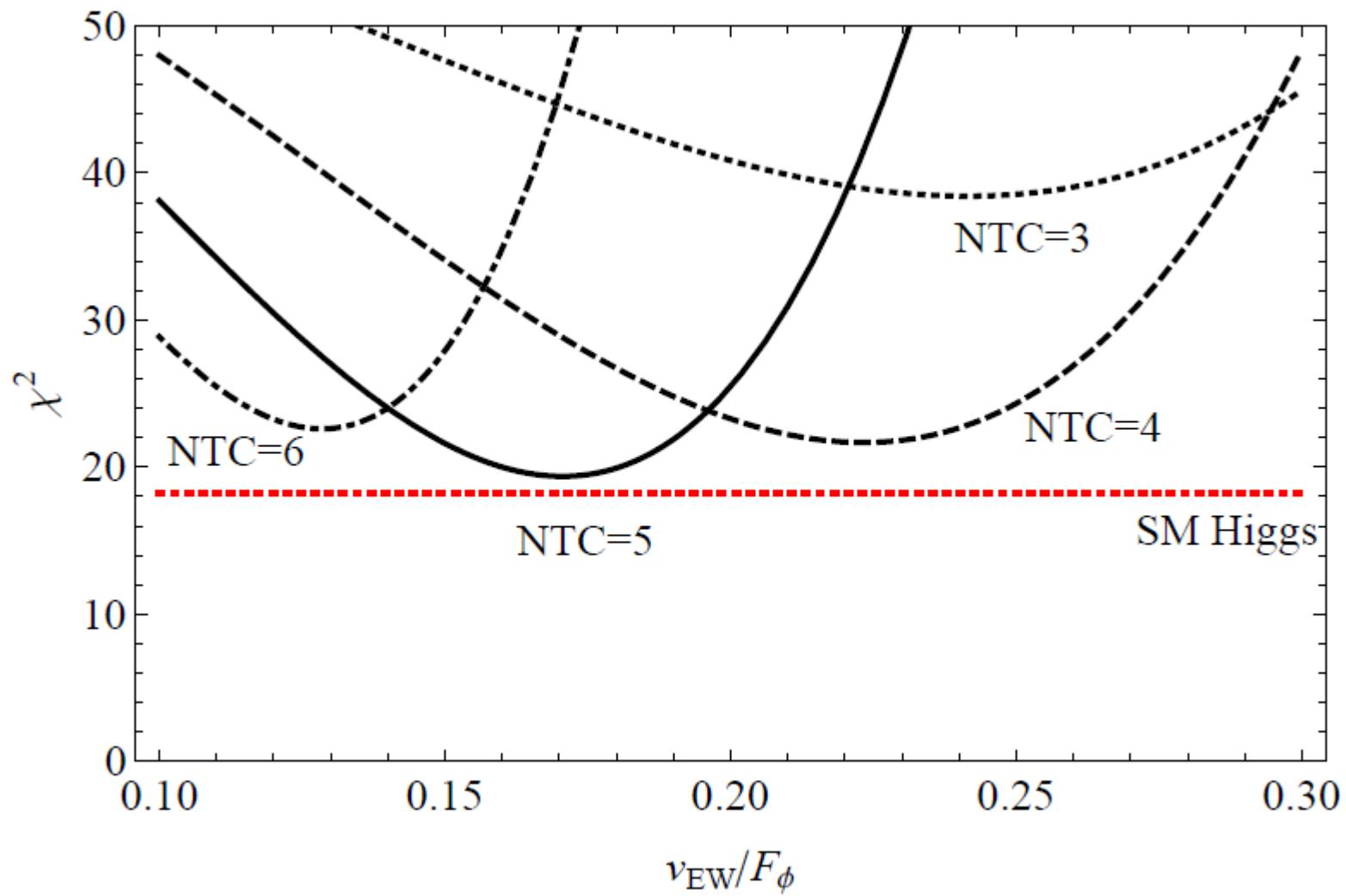
NTC = 5

Techni- $\rho$ , a1 masses	:	$M_\rho = M_{a1} = 3.9 \text{ TeV}$
Techni-glueball (TG) mass	:	$M_G = 18 \text{ TeV}$
TG decay constant	:	$F_G = 174 \text{ TeV}$
dynamical TF mass $m_F$	:	$m_F = 0.85 \text{ TeV}$

# *W/ Tevatron data included:*



*W/ Tevatron data included:*



# ★ Other pheno. issues in TC scenarios

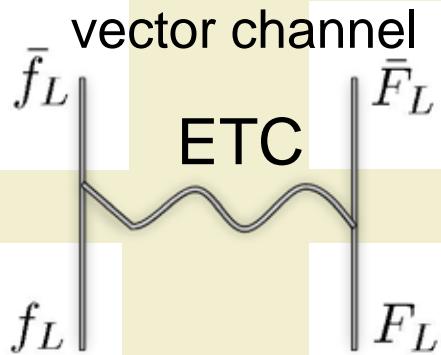
## S parameter

$$S \approx N_D \cdot \frac{8\pi F_\pi^2}{M_\rho^2} \simeq \underline{0.3 \cdot N_D} \quad (\text{for QCD-like})$$

$N_D$  : # EW doublets

too large! Cf:  $S(\text{exp}) < 0.1$  around  $T = 0$

One resolution: *ETC-induced “delocalization” operator*



$$-\frac{1}{\Lambda_{\text{ETC}}^2} J_{\mu \text{SM}_L}^a J_{\text{TC}_L}^{\mu a}$$

in low-energy

$$J_{\text{TC}_L}^{\mu a} \rightarrow \text{Tr}[U^\dagger \frac{\sigma^a}{2} i D^\mu U]$$

Chivukula-Simmons-He-Kurachi-Tanabashiet al (2005)

$$\text{w/ } U = e^{2i\pi_{\text{eaten}}/v_{\text{EW}}}$$

$\ni g_W W_\mu - g_Y B_\mu$  modifies SM f-couplings to W, Z

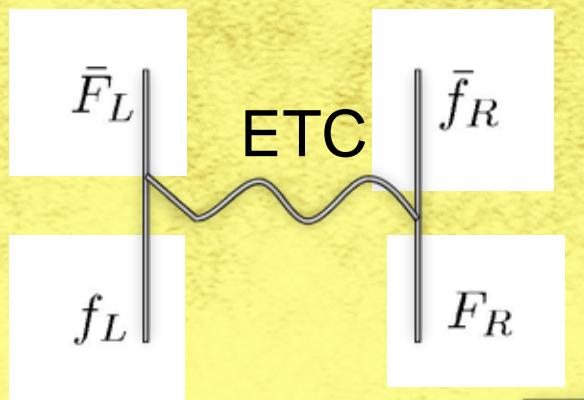
contributes to S “negatively”



$$\Delta S \sim \cancel{-\frac{8\pi}{g_W^2}} \left( \frac{v_{\text{EW}}}{\Lambda_{\text{ETC}}} \right)^2$$

$S_{\text{total}} \rightarrow 0$  (“ideal delocalization”)

## Top quark mass generation



$$m_t \approx \frac{\langle \bar{U}U \rangle_{\text{ETC}}}{\Lambda_{\text{ETC}}^2} \approx \left( \frac{\Lambda_{\text{TC}}}{\Lambda_{\text{ETC}}} \right)^2 \Lambda_{\text{TC}}$$

ETC scale associated w/ top mass

$$\Lambda_{\text{ETC}}^{\text{top}} \approx 1 \text{TeV} \left( \frac{\Lambda_{\text{TC}}}{1 \text{TeV}} \right)^{3/2} \left( \frac{172 \text{GeV}}{m_t} \right)^{1/2}$$

**too small!**

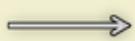
**One resolution:** **Strong ETC**

Miransky-K.Y. (1989), Matumoto(1989), Appelquist-Einhorn-Takeuchi-Wijewardhana (1989)

--- makes induced 4-fermi ( $t\bar{t} U\bar{U}$ ) coupling large enough to trigger chiral symm. breaking (almost by NJL dynamics)

$$\langle \bar{U}U \rangle_{\text{ETC}} \approx \left( \frac{\Lambda_{\text{ETC}}}{\Lambda_{\text{TC}}} \right)^{\gamma_m} \langle \bar{U}U \rangle_{\text{TC}} \quad 1 < \gamma_m \leq 2$$

**boost-up**



$$m_t \approx \left( \frac{\Lambda_{\text{TC}}}{\Lambda_{\text{ETC}}} \right)^{2-\gamma_m} \Lambda_{\text{TC}} \leq \Lambda_{\text{TC}} \sim 1 \text{TeV}$$

**T parameter** (Strong) ETC generates large isospin breaking  
à highly model-dependent issue