

Search for new resonances in ttbar final states with the ATLAS detector

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ttbar heavy resonances

- Benchmark models (2HDM)
- Search in lepton + jets channel
 - Strategy
 - Selection
 - mtt reconstruction
- Impact of detector resolution
 - Jet rescaling method

Introduction

- Several Beyond-the-SM (BSM) theories seeking to solve a variety of open questions in SM:
 - Neutrino mass
 - Mass spectrum of the particles
 - Inclusion of gravitation
 - Hierarchy problem, ...
- BSM theories predict new heavy particles which could couple strongly with the top quark

Benchmark models

- Searches for new particles in ttbar final states
 - Warped extra-dimensions (Randall-Sundrum):
 - Kaluza-Klein gluon (g_{kk})
 - KK Bulk-RS graviton (G_{kk})
 - Top-color model (TC):
 - Leptophobic TC heavy Z' boson
 - 2 Higgs Doublet Model (2HDM):
 - 4 scalars and 1 pseudo-scalar:

Benchmark models

- Searches for new particles in ttbar final states
 - Warped extra-dimensions (Randall-Sundrum):
 - Kaluza-Klein gluon (g_{kk}) spin-1, width: 10-40% (broad resonance)
 - KK Bulk-RS graviton (G_{kk}) spin-2, width: 3-6%
 - Top-color model (TC):
 - Leptophobic TC heavy Z' boson spin-1, width: 1.2% (narrow resonance)
 - 2 Higgs Doublet Model (2HDM):
 - 4 scalars and 1 pseudo-scalar: spin-0, width: from ~0% until 100%

Width have an impact in the potential discovery!

Top quark pairs

- Top quark decays before hadronizing
 - BR(t \rightarrow bW⁺) ~ 100%
- At the LHC, top pair production (LO) at 8 TeV.
 - 80% Gluon fusion: (a) and (b)
 - 20% quark-antiquark annihilation: (c)



- Top pair branching ratio:
 "alljets" 46%
 t+jets 15%
 "dileptons"
- Search for new physics in the Lepton + jets channel: Best balance between the BR and background rejection



Experimental setup

The ATLAS detector





Two Higgs doubles model (2HDM)

2HDM motivation

 Discovery of a SM Higgs-like particle at LHC: m_h ~ 125 GeV



- Further investigations into the scalar sector
 - It may have richer structure
- Two-Higgs doublet are the simplest extension in the scalar sector of SM

2HDM: Scalar sector

- Two identical SU(2) doublets scalar fields Φ_1 and Φ_2
- Most general scalar potential invariant under electroweak symmetry (tree-level)

$$\begin{split} V &= m_{11}^2 \, \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \, \Phi_2^{\dagger} \Phi_2 - m_{12}^2 \, \left(\Phi_1^{\dagger} \Phi_2 + \Phi_2^{\dagger} \Phi_1 \right) + \frac{\lambda_1}{2} \left(\Phi_1^{\dagger} \Phi_1 \right)^2 + \frac{\lambda_2}{2} \left(\Phi_2^{\dagger} \Phi_2 \right)^2 \\ &+ \lambda_3 \, \Phi_1^{\dagger} \Phi_1 \, \Phi_2^{\dagger} \Phi_2 + \lambda_4 \, \Phi_1^{\dagger} \Phi_2 \, \Phi_2^{\dagger} \Phi_1 + \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + \left(\Phi_2^{\dagger} \Phi_1 \right)^2 \right], \end{split}$$

- 8 degrees of freedom $\,\rightarrow\,$ 3 correspond to the W, Z bosons masses
- The other 5 are manifested as physical particles
- Free parameters of the model:
 - tanβ: ratio between the two vev's
 - α : mixing angle of CP-even fields
 - Masses of the new particles



• 4 types of 2HDM: different coupling to fermions

Туре	u_R	d_R	e_R	
Type I Type II Lepton-specific Flipped	$egin{array}{c} \Phi_2 \ \Phi_2 \ \Phi_2 \ \Phi_2 \ \Phi_2 \end{array}$	$egin{array}{c} \Phi_2 \ \Phi_1 \ \Phi_2 \ \Phi_1 \ \Phi_2 \ \Phi_1 \end{array}$	Φ_2 Φ_1 Φ_1 Φ_2	→ MSSM

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• Couplings with SM particles modified by factors $(g_{\Phi,\bar{q}/V,qV})$ depending of α and β

Φ	$g_{\Phi ar{u} u}$		9	$g_{\Phi VV}$	
	Type I	Type II	Type I	Type II	Type I/II
h	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$-\sin lpha / \cos eta$	$\sin(eta-lpha)$
H	$\sin lpha / \sin eta$	$\sin lpha / \sin eta$	$\sin lpha / \sin eta$	$\cos lpha / \cos eta$	$\cos(eta-lpha)$
A	\coteta	\coteta	\coteta	aneta	0

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• $|\sin(\beta - \alpha)| = 1 \rightarrow h$ scalar looks like SM Higgs boson

 $H^{SM} = h * sin(\beta - \alpha) + H * cos(\beta - \alpha)$





- $tan\beta = 1.0$ and $sin(\beta \alpha) = 0.1$
- WW and ZZ channels are dominant wrt ttbar

Couplings: top quark, b quark and vector boson

Φ	$g_{\Phi ar{u} u}$		9	$g_{\Phi VV}$	
	Type I	Type II	Type I	Type II	Type I/II
h	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$-\sinlpha/\coseta$	$\sin(eta-lpha)$
H	$\sin lpha / \sin eta$	\sinlpha/\sineta	$\sin lpha / \sin eta$	$\cos lpha / \cos eta$	$\cos(eta-lpha)$
A	\coteta	\coteta	\coteta	aneta	0





- $tan\beta = 1.0$ and $sin(\beta \alpha) = 0.5$
- WW and ZZ channels are dominant
 - \rightarrow Coupling with top quark is weaker

Couplings: top quark, b quark and vector boson

Φ	$g_{\Phi ar{u} u}$		g	$g_{\Phi VV}$	
	Type I	Type II	Type I	Type II	Type I/II
h	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$-\sinlpha/\coseta$	$\sin(eta-lpha)$
H	$\sin lpha / \sin eta$	\sinlpha/\sineta	$\sin lpha / \sin eta$	$\cos lpha / \cos eta$	$\cos(eta-lpha)$
A	\coteta	\coteta	\coteta	aneta	0

2HDMs: Type II



- $tan\beta = 1.0$ and $sin(\beta \alpha) = 1.0$
- SM-like limit
- ttbar channel is dominant
 - \rightarrow Coupling W, Z vanish
- Couplings: top quark, b quark and vector boson

Φ	$g_{\Phi ar{u} u}$		g	$g_{\Phi VV}$	
	Type I	Type II	Type I	Type II	Type I/II
h	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$-\sinlpha/\coseta$	$\sin(eta-lpha)$
H	$\sin lpha / \sin eta$	\sinlpha/\sineta	$\sin lpha / \sin eta$	$\cos lpha / \cos eta$	$\cos(eta-lpha)$
A	\coteta	\coteta	\coteta	aneta	0

2HDMs: Type I



- $tan\beta = 1.0$ and $sin(\beta \alpha) = 1.0$
- SM-like limit
- Same couplings with u-type quraks than Type I \rightarrow ttbar channel is dominant
- Couplings: top quark, b quark and vector boson

Φ	$g_{\Phi ar{u} u}$		g	$g_{\Phi VV}$	
	Type I	Type II	Type I	Type II	Type I/II
h	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$\cos lpha / \sin eta$	$-\sinlpha/\coseta$	$\sin(eta-lpha)$
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ttbar resonances in the lepton+jets channel



- Top-antitop events selected in e/mu + jets channels
- Two types of jets topology: resolved and boosted
- Invariant mass of ttbar candidates are reconstructed
- Background estimation
- Invariant mass spectrum is searched for local excesses/deficits
- Limits are set in the cross-section*BR if no significant excesses/deficits
- Cross-section*BR limits translated into bounds on the allowed mass for new particles

Event selection

- Common selection:
 - Exactly 1 lepton (e or μ) in the acceptance of the detector
 - MET > 20 GeV
 - MTW (transverse mass computed using MET and lepton four-momentum)
 - MET + MTW > 60 GeV
- Boosted selection:
 - ≥ 1 small-R jet
 - ≥ 1 large-R jet
 - ≥ 1 b-tagged small-R
- Resolved selection:
 - Fail boosted selection
 - ≥ 4 small-R jets
 - ≥ 1 b-tagged small-R jet
 - $Log_{10}\chi^2 < 0.9$



Neutrino's reconstruction

- Neutrinos do not interact with the detector material
 - \rightarrow inferred from the total transverse momentum balance
- Neutrino's momentum: transverse component from MET
 - Longitudinal component of the momentum is unknown (pZ)
 - pZ computed imposing an on-shell W mass constrain: lepton + MET system
 - Quadratic equation for pZ:
 - Real solutions:
 - Smallest pZ

$$p_{z,
u}^{\pm} = rac{\mu p_{z,l}}{p_{\mathrm{T},l}} \pm \sqrt{rac{\mu^2 p_{z,l}^2}{p_{\mathrm{T},l}^4} - rac{E_l^2 p_{\mathrm{T},
u}^2 - \mu^2}{p_{\mathrm{T},l}^2}}$$

- Non-real solutions:
 - MET is rescaled and rotated applying the minimum variation necessary to find one real solution

 \rightarrow Imperfect resolution on MET is the most likely explanation for non-real solutions to pZ equation

Chi2 minimization: algorithm used to select the right jet combination to reconstruct mtt



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$$\chi^{2} = \left[\frac{m_{jj} - m_{W}}{\sigma_{W}}\right]^{2} + \left[\frac{m_{jjb} - m_{jj} - m_{th-W}}{\sigma_{th-W}}\right]^{2} + \left[\frac{m_{jl\nu} - m_{tl}}{\sigma_{tl}}\right]^{2} + \left[\frac{(p_{T,jjb} - p_{T,jl\nu}) - (p_{T,th} - p_{T,tl})}{\sigma_{P_{T,th} - P_{T,tl}}}\right]^{2}$$

- All possible jet permutation are tried
 - Only the permutation with lowest chi2 are used
- Constrains determined from the reconstructed objets matched to the partons



• **Chi2 minimization**: algorithm used to select the right jet combination to reconstruct mtt

$$\chi^{2} = \left[\frac{m_{jj} - m_{W}}{\sigma_{W}}\right]^{2} + \left[\frac{m_{jjb} - m_{jj} - m_{th-W}}{\sigma_{th-W}}\right]^{2} + \left[\frac{m_{jl\nu} - m_{tl}}{\sigma_{tl}}\right]^{2} + \left[\frac{(p_{T,jjb} - p_{T,jl\nu}) - (p_{T,th} - p_{T,tl})}{\sigma_{P_{T,th} - P_{T,tl}}}\right]^{2}$$

Chi2 performance: The algorithm selects the correct combination ~70% of the reconstructible events



mtt reconstruction

- Invariant mass of the ttbar pair is the discriminant variable
- ttbar pair is build by summing up the four momentum of jets, lepton and neutrino
- Reconstruction of mtt for the signal of benchmark models considered at different generated mass





Impact of the detector resolution in the search for 2HDM signal into ttbar

2 HDM signal into tt

- Decay channel: $H/A \rightarrow t \bar{t}$
- Signal have same initial/final states than the SM bkg
 - Interference (peak/dip structure)

H signal + interference



Smearing

- Detector smearing as a tool for phenomenology
 - Useful to understand the limitations induced by detectors on an theoretical well motivated analysis
 - Smearing techniques can help to produce rough estimates
- Detector resolution: assuming a gaussian impact
 - ~8% of the reconstructed mass for the ATLAS detector
 - Convolution of the mtt spectra with a gaussian of 8% width
 - σ/μ is interpreted as the detector resolution

Impact of detector resolution

 The smearing will help to check the impact of the resolution of ATLAS on the scalar signal w/ interference

Loss of sensitivity \rightarrow Signal more difficult to detect

• Improving the mtt resolution to balance the detector resolution:

 \rightarrow Let use the information of the ttbar system





Jet rescaling

Rescaling the jets from W_h

- Considering reconstructible events
- Rescale the jets 4-momentum $(J_1 + J_2)$ to get the Mw value at the PDG
 - M_w = 80.4 GeV
 - $k^{2}(\mathbf{J}_{1} + \mathbf{J}_{2})^{2} = (M_{W})^{2}$
- Top mass and mtt resolutions are improved





Rescaling the b-jets

- Rescale the b-jet 4-momentum:
 - Constraining the top mass ($M_{top} = 172.5 \text{ GeV}$)
- The jet rescaling absorbs the uncertainties on the mass reconstruction
- Improvement on the mtt resolution





Improvement of mtt resolution

• Relative mass difference:

$$rac{\Delta_{m_{tt}}}{m_{tt}^{true}} = rac{m_{tar{t}}^{ extsf{reco}} - m_{tar{t}}^{ extsf{true}}}{m_{tar{t}}^{ extsf{true}}}$$

• Gaussian fit to extract the mass resolution



mtt resolution vs scalar mass (assuming perfect reconstruction)

- σ (mtt diff rel) for all the mass points:
 - Low mass: improvement after each rescaling
 - High mass: degradation after the rescaling of the leptonic top
 - Uncertainty in the MET direction has large impact at high mass





Improvement wrt the no-rescaling case

$$1 - rac{\sigma_{Cal} - \sigma_{noCal}}{\sigma_{noCal}}$$

 ~25% improvement after W_h and b_h rescaling

mtt resolution vs scalar mass (χ² algorithm)

- Same tendency using the standard reconstruction algorithm
 - High mass: degradation after the rescaling of the leptonic top





- ~20% improvement after W_h and b_h rescaling
- Reasonable considering the improvement in the ideal case

Conclusion

- ttbar decay mode for the scalar signal is dominant in a region of the parameters space
- Improving the mass resolution will help to balance the impact of the detector resolution

Rescaling of Wh and bh:

- → ~25% in the assuming perfect reconstruction
- → ~20% using the chi2 minimization
- The calibration of the bl degrades the mtt resolution (investigations ongoing)
 Bad energy/direction resolutions of MET
- Next prospect: Rescaling the large-R jet for the boosted topology This study will be used in the search for tt resonance at 13 TeV

Back up



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Various BSM models predict resonances in the $m_{t\bar{t}}$ spectrum

Warped extra-dimensions (Randall-Sundrum)

- Kaluza-Klein (KK) gluon (g_{KK}) Spin-1, width: 10-40%
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Top-colour (TC) model:

Leptophobic TC Z' boson (Z') Spin-1, width: 1.2%

Searches in all decay channels:

- All-hadronic ۰
- 1 lepton+jets 0
- (Dileptonic)



Bulk-RS Graviton Branching Ratios

2 HDM signal into Tt

- Decay channel: $H/A \rightarrow t \bar{t}$
- Signal have same initial/final states than the SM bkg
 - Interference (peak/dip structure)



- $m_{t\bar{t}}$ at truth level:
 - SM bkg
 - Signal + bkg
 - Signal + bkg + interference

2 HDM signal into tt

• Impact of the detector resolution (Smearing by 8%)

Loss of sensitivity \rightarrow Signal more difficult to detect

- Improving the mtt resolution to balance the detector resolution:
 - → Jet Rescaling



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• $|sin(\beta - \alpha)| = 1 \rightarrow h$ scalar looks like SM Higgs boson

Н SM =	h*	$sin(\beta)$	$-\alpha$)	+H	* COS	s (β –	α)
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- Reconstructible events
 - Partons matched to the physics objects
 - Perfect reconstruction assumption
- Efficiencies of the reconstructible events
 - Between 50% and 60% after the matching of all the objects





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H signal

H signal + interference



Lepton+jet channel: Physics objects

- Only considered electron/muon in the final state
- Mini-isolation: [Σ pT of tracks within a variable cone size]/pT
 - Stable performance with efficiency close to 100%





 a) Fixed-cone isolation: Cut on ∑p_T (∑E_T) in cone with fixed radius R



b) Mini-isolation: Cut on ∑p_T (∑E_T) in cone with radius R=k/p_T¹

Lepton+jet channel: Physics objects

• Jets:



• B-tagging is applied to small-R jets

2HDM production mode



• Gluon fusion $(gg \rightarrow A/H)$ is the dominant channel