Gravitational Waves as Strong Field Probes of the Universe

Nicolas Yunes Montana State University

APC Invited Talk, June 12th, 2014 Yunes & Siemens, Living Reviews in Relativity 2014, http://arxiv.org/abs/1304.3473

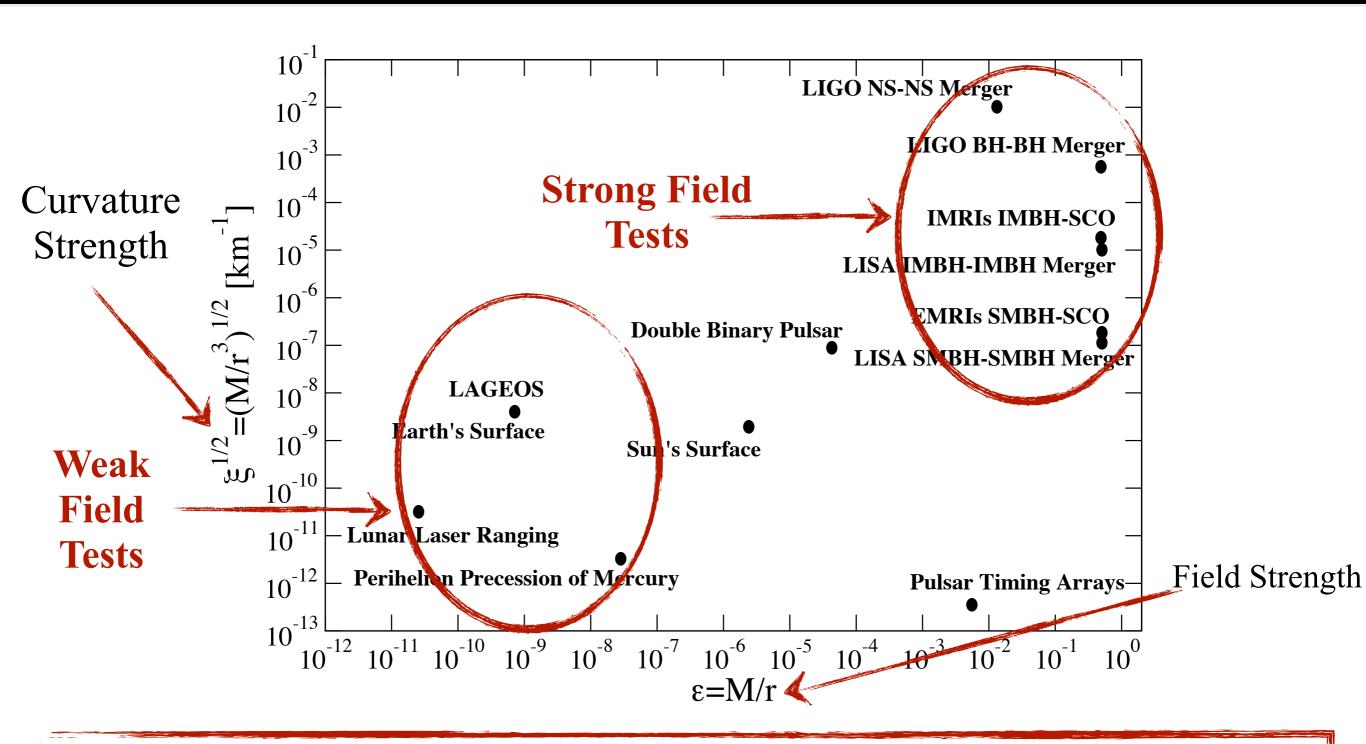
Standing on the Shoulders of...

Clifford Will, Jim Gates, Stephon Alexander, Abhay Ashtekar, Sam Finn, Ben Owen, Pablo Laguna, Emanuele Berti, Uli Sperhake, Dimitrios Psaltis, Avi Loeb, Vitor Cardoso, Leonardo Gualtieri, Daniel Grumiller, David Spergel, Frans Pretorius, Neil Cornish, Scott Hughes, Carlos Sopuerta, Takahiro Tanaka, Jon Gair, Paolo Pani, Antoine Klein, Kent Yagi, Laura Sampson, Luis Lehner, Masaru Shibata, Curt Cutler, Haris Apostolatos,

An <u>incomplete</u> summary of what GWs will tell us about gravity in the cosmology context

Leo Stein, Sarah Vigeland, Katerina Chatziioannou, Philippe Jetzer, Leor Barack, Kostas Glampedakis, Stanislav Babak, Ilya Mandel, Chao Li, Eliu Huerta, Chris Berry, Alberto Sesana, Carl Rodriguez, Georgios Lukes-Gerakopoulos, George Contopoulus, Chris van den Broeck, Walter del Pozzo, Jon Veitch, Nathan Collins, Deirdre Shoemaker, Sathyaprakash, Devin Hansen, Enrico Barausse, Carlos Palenzuela, Marcelo Ponce, etc.

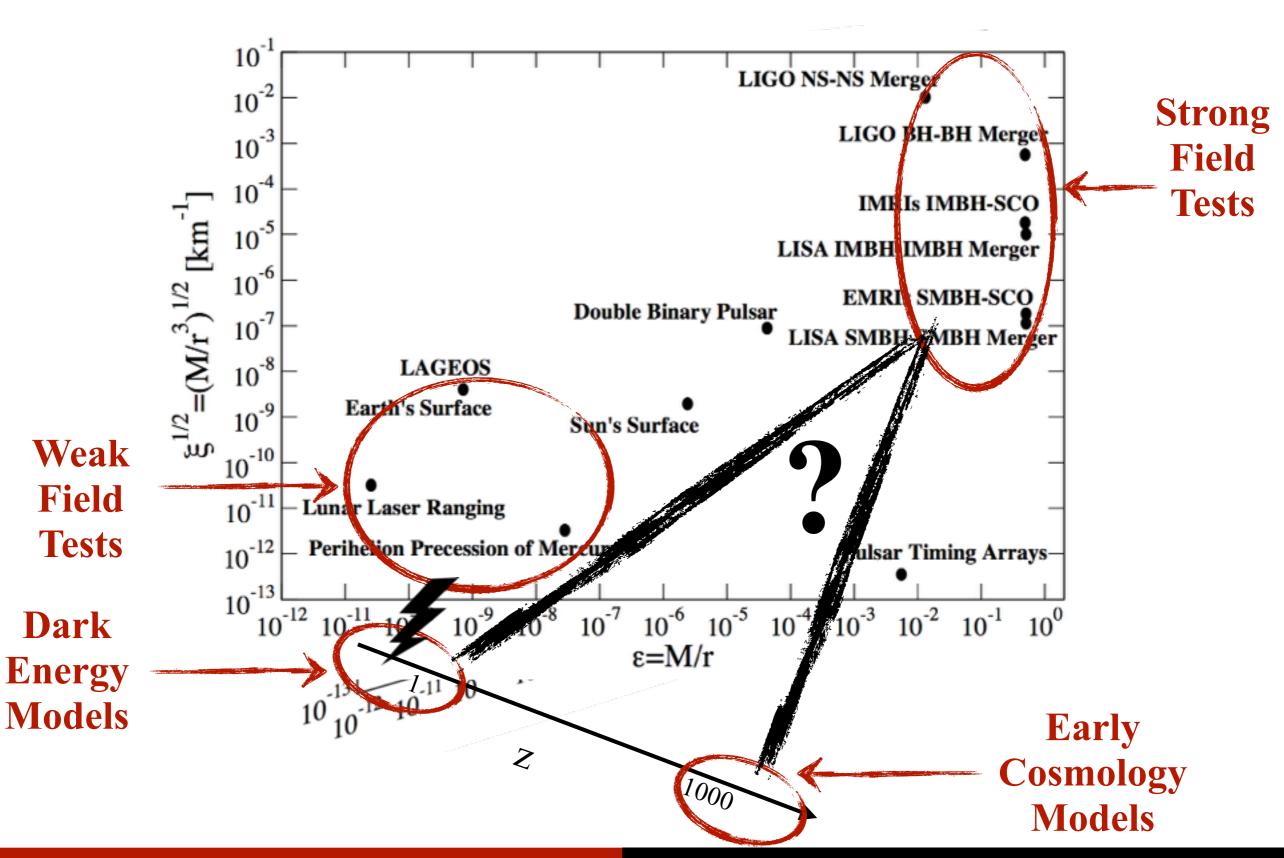
Testing Phase Space



GWs can probe the non-linear, dynamical, strong-field regime

Will, Liv. Rev., 2005, Psaltis, Liv. Rev., 2008, Siemens & Yunes, Liv. Rev. 2013.

Enlarging Phase Space



Divide and Conquer

Nico's (GW-Biased) GW Modified Theory Classification:



Strong-Field

Well-constrained by binary pulsars, so need screening Eg, Scalar-Tensor theories

Constrainable with GW observations, natural suppression without screening Eg, Chern-Simons, Gauss-Bonnet, etc.

Nico's (GW-Biased) Cosmological Modified Theory Classification:



Unscreened

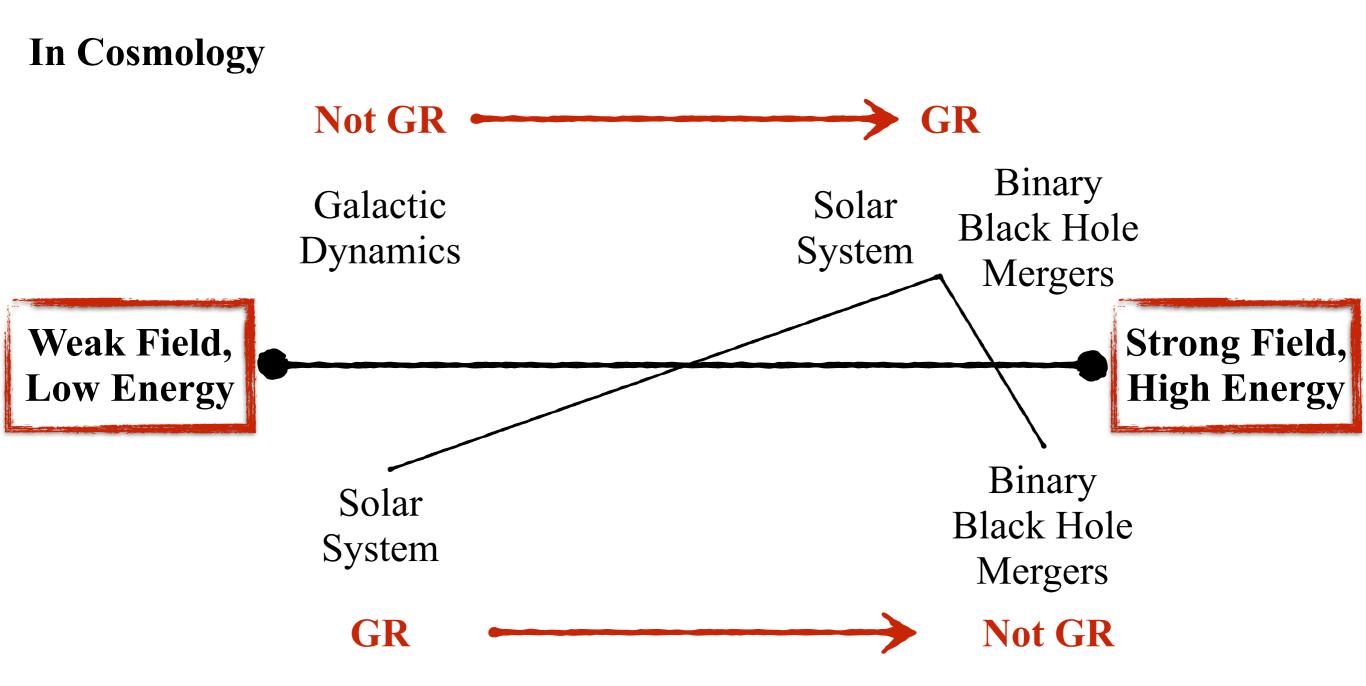
Late-time expansion, DE

Eg, chameleon, Vainshtein, etc.

Early-time cosmology, inflation

Eg, Chern-Simons, Gauss-Bonnet, etc.

Screening in Cosmology ≠ Screening in GWs



In Gravitational Wave Physics

Roadmap



Weak Field Theories

Example: Scalar Tensor Theories

$$S_{\mathrm{Jordan}} \sim \int d^4x \sqrt{-g} \left[\phi R - \frac{\omega(\phi)}{\phi} \left(\partial^{\mu} \phi \right) \left(\partial_{\mu} \phi \right) + \mathcal{L}_{\mathrm{matter}} \right]$$

$$\phi \to g_{\mu\nu} \to T_{\mu\nu}$$

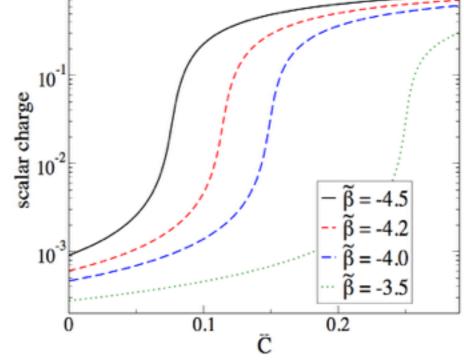
Effective Coupling to Matter:

$$\alpha \sim \frac{1}{\sqrt{3+2\omega_{BD}}} \phi \beta (\phi - \phi_0)$$

Main Effect:

Stars acquire scalar charge

Spontaneous Scalarization



Dominant Observables:

Grav. and Inertial center of mass — do not coincide

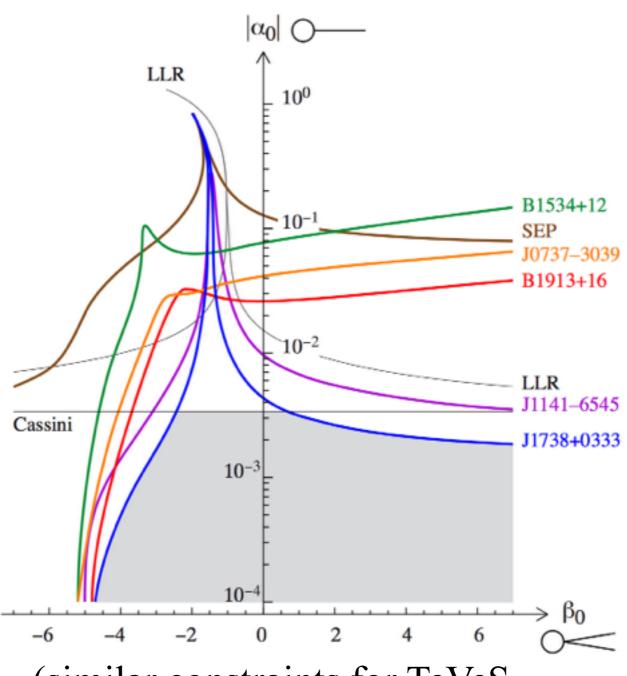
Screened Dipole Gravitational Wave Emission

Faster Orbital
Decay

Damour+Esposito-Farese, PRD 54 ('96) Palenzuela, et al, PRD 97 ('13), 89 ('14).

Constraints on Weak Field Theories

Scalarizable Scalar-Tensor:



(similar constraints for TeVeS and for massive Brans-Dicke)

Freire, et al, MRAS 18 ('12). Alsing, et al, PRD 85, ('12).

Strong Field Theories

Example: Quadratic Gravity

Definition:

$$S_{\mathrm{Quad}} \sim \int d^4x \sqrt{-g} \left[R - \frac{1}{2} \left(\partial_{\mu} \vartheta \right) \left(\partial^{\mu} \vartheta \right) + \alpha_1 \vartheta R^2 + \alpha_2 \vartheta R_{\mu\nu} R^{\mu\nu} + \alpha_3 \vartheta R_{\mu\nu\delta\sigma} R^{\mu\nu\delta\sigma} + \alpha_4 \vartheta R_{\mu\nu\delta\sigma} \right. \\ \left. * R^{\mu\nu\delta\sigma} \right] \left[R - \frac{1}{2} \left(\partial_{\mu} \vartheta \right) \left(\partial^{\mu} \vartheta \right) + \alpha_1 \vartheta R^2 + \alpha_2 \vartheta R_{\mu\nu} R^{\mu\nu} + \alpha_3 \vartheta R_{\mu\nu\delta\sigma} R^{\mu\nu\delta\sigma} + \alpha_4 \vartheta R_{\mu\nu\delta\sigma} \right] \right] \left[R - \frac{1}{2} \left(\partial_{\mu} \vartheta \right) \left(\partial^{\mu} \vartheta \right) + \alpha_1 \vartheta R^2 + \alpha_2 \vartheta R_{\mu\nu} R^{\mu\nu} + \alpha_3 \vartheta R_{\mu\nu\delta\sigma} R^{\mu\nu\delta\sigma} + \alpha_4 \vartheta R_{\mu\nu\delta\sigma} \right] \right]$$

certain choices of couplings lead to Einstein-Dilaton-Gauss-Bonnet theory or dynamical Chern-Simons gravity.

Main Effects: dCS. Gravitational Parity Violation, inverse no-hair theorem.

Dominant Observables:

Chirping of Gravitational
Wave Phase

Requires observation of late inspiral & merger

Alexander & Yunes, Phys. Rept 480 ('09) Yunes & Stein, PRD 83 ('11)

Constraints on Strong Field Theories

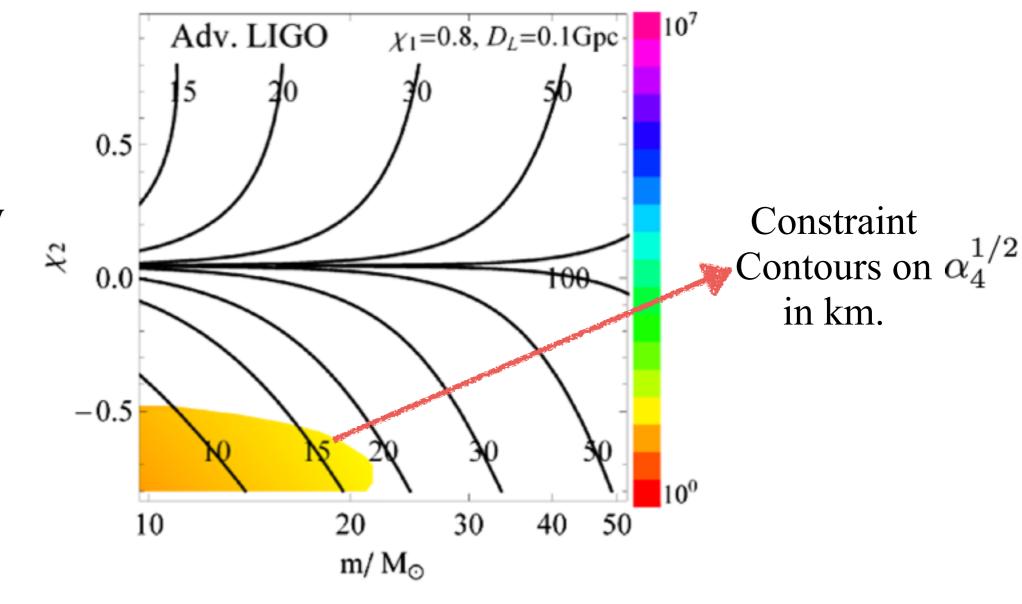
dCS

Current constraints

Extremely weak from Solar System (GPB)

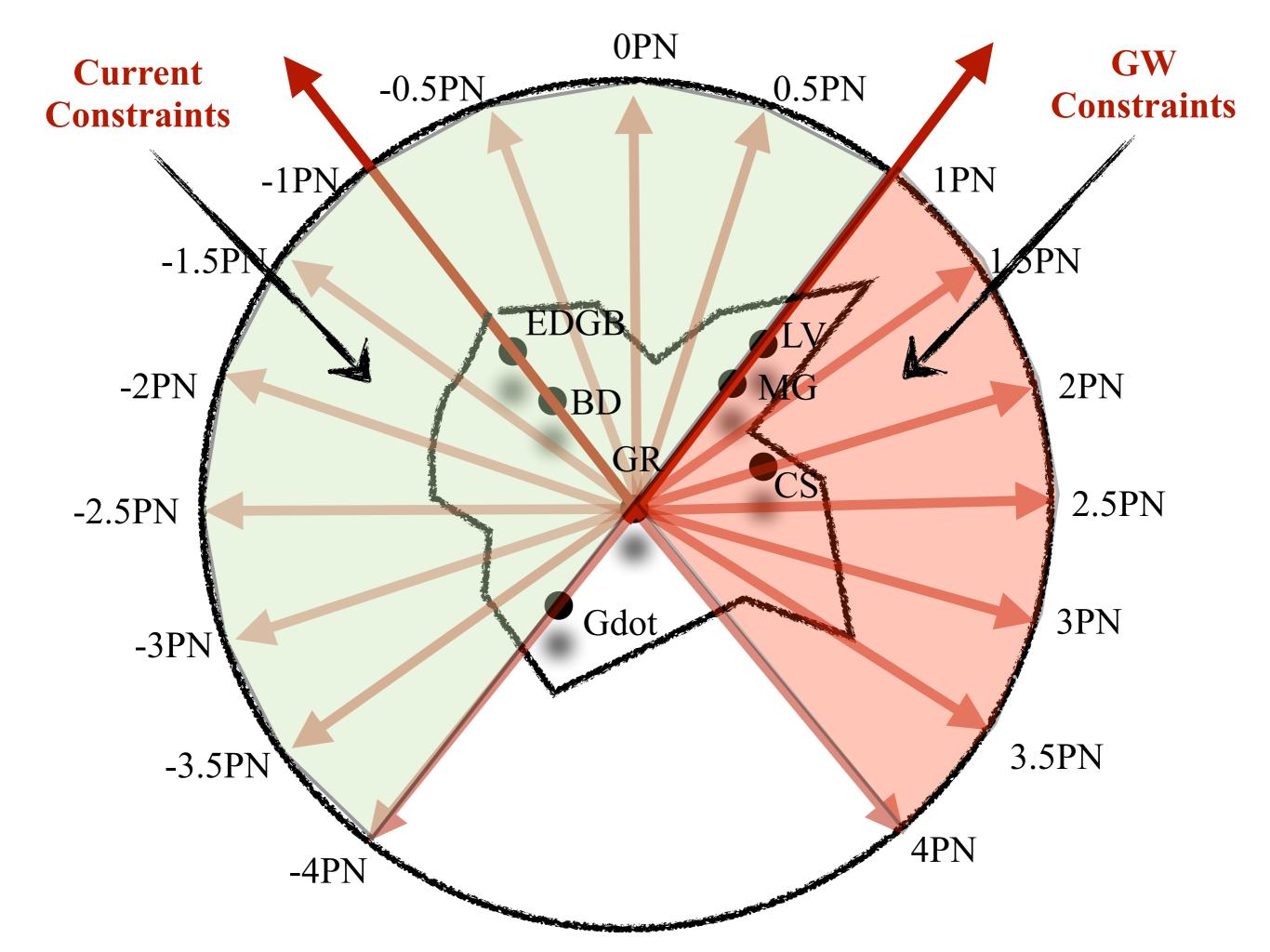
$$\alpha_4^{1/2} < \mathcal{O}(10^8~\mathrm{km})$$

Projected GW constraints



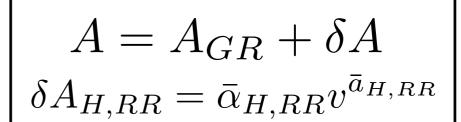
Yagi, Yunes & Tanaka, PRL 109 ('12)

Parametrized Post-Einsteinian



Parameterized post-Einsteinian Framework

- I. Parametrically deform the Hamiltonian.
- II. Parametrically deform the RR force.



III. Deform waveform generation.

$$h = F_+ h_+ + F_\times h_\times + F_s h_s + \dots$$

IV. Parametrically deform g propagation.

$$\left| E_g^2 = p_g^2 c^4 + \tilde{\alpha} p_g^{\tilde{\alpha}} \right|$$

Result: To leading PN order and leading GR deformation

$$\tilde{h}(f) = \tilde{h}_{GR}(f) \left(1 + \alpha f^a\right) e^{i\beta f^b}$$

Yunes & Pretorius, PRD 2009 Mirshekari, Yunes & Will, PRD 2012 Chatziioannou, Yunes & Cornish, PRD 2012

Questions to Ask

Templates/ Theories	GR	ppE
GR	Business as usual	Quantify the statistical significance that the detected event is within GR. Anomalies?
Not GR	Quantify fundamental bias introduced by filtering non-GR events with GR templates	Can we measure deviations from GR characterized by non-GR signals? Model Evidence.

$$\tilde{h}(f) = \tilde{h}_{GR}(f) \left(1 + \alpha f^a\right) e^{i\beta f^b}$$

[Yunes & Pretorius, PRD 2009, Mirshekari, Yunes & Will, PRD 2012, Chatziioannou, Yunes & Cornish, PRD 2012]

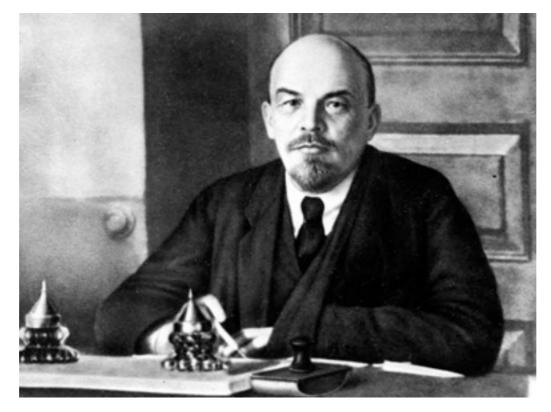
$a_{ m ppE}$	$b_{ m ppE}$	Interpretation
1		Parity Violation
-8	-13	Anomalous Acceleration, Extra
		Dimensions, Violation of Posi-
		tion Invariance
	-7	Dipole Gravitational Radiation,
		Electric Dipole Scalar Radiation
	-3	Massive Graviton Propagation
$\propto \text{spin}$	-1	Magnetic Dipole Scalar Radia-
		tion, Quadrupole Moment Cor-
		rection, Scalar Dipole Force

What does it all mean?

Clear difference between "Cosmological Modified Theories" and "GW Modified Theories". More cross-pollination needed.

Weak Field Modified Theories best constrained with binary pulsars (they will be hard to constrain with gravitational waves)

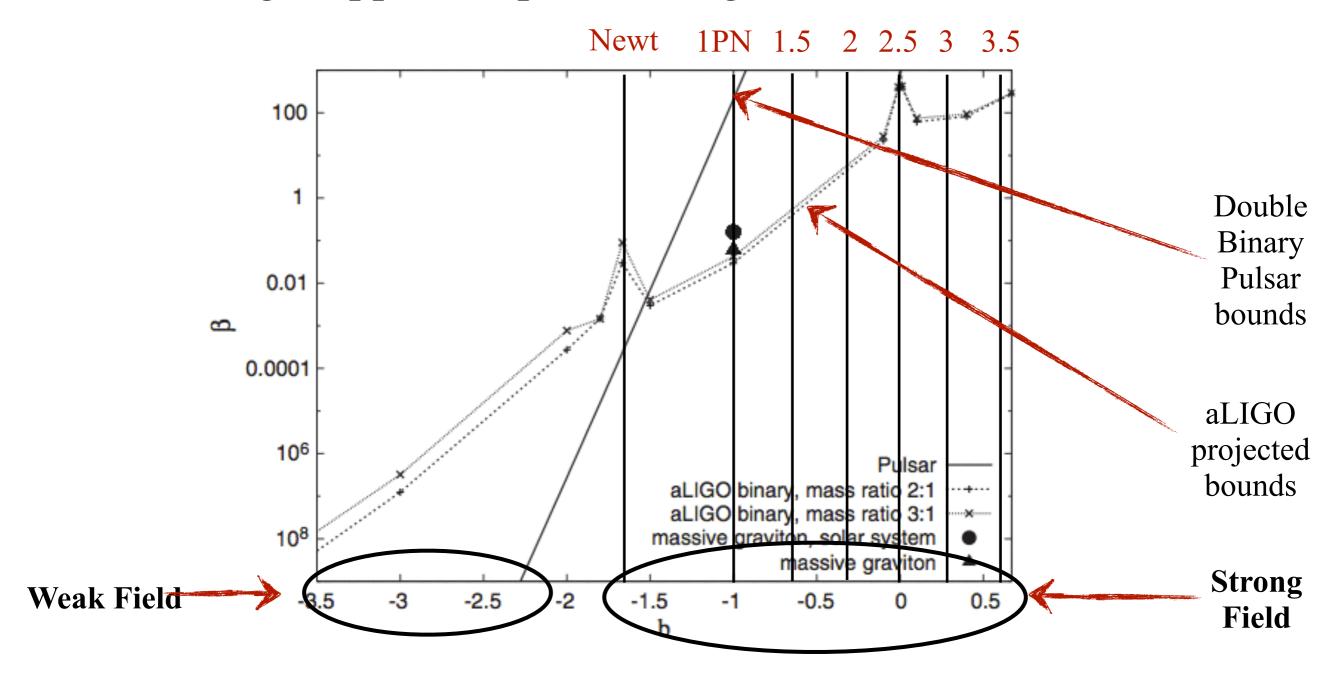
Strong Field Theories are only strongly constrained by gravitational waves from binary mergers.



Doveryai, no proveryai

Projected Gravitational Wave Constraints

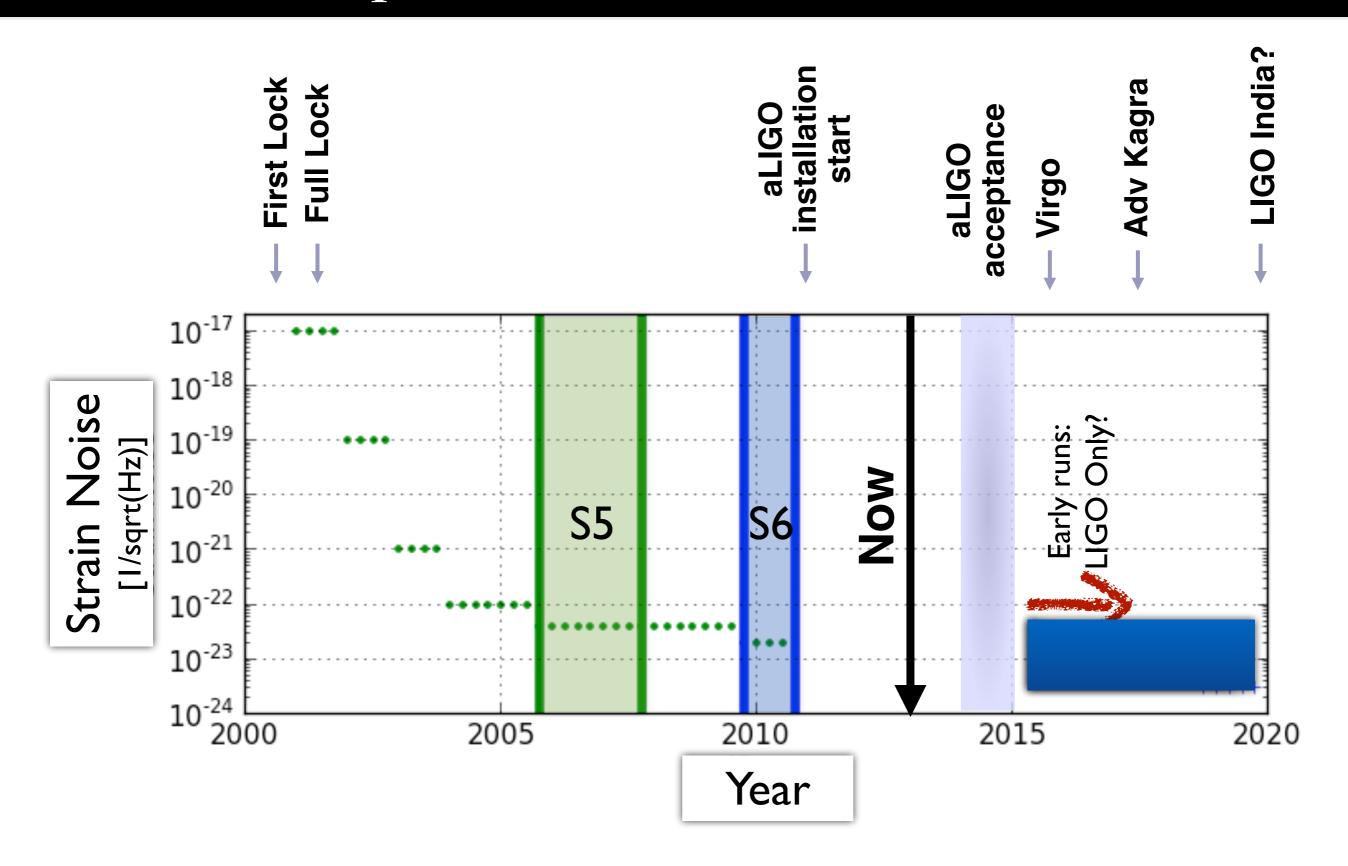
GR Signal/ppE Templates, 3-sigma constraints, SNR = 20



$$\tilde{h}(f) = \tilde{h}_{GR}(f) \left(1 + \alpha f^a\right) e^{i\beta f^b}$$

Yunes & Hughes, 2010, Cornish, Sampson, Yunes & Pretorius, 2011 Sampson, Cornish, Yunes 2013.

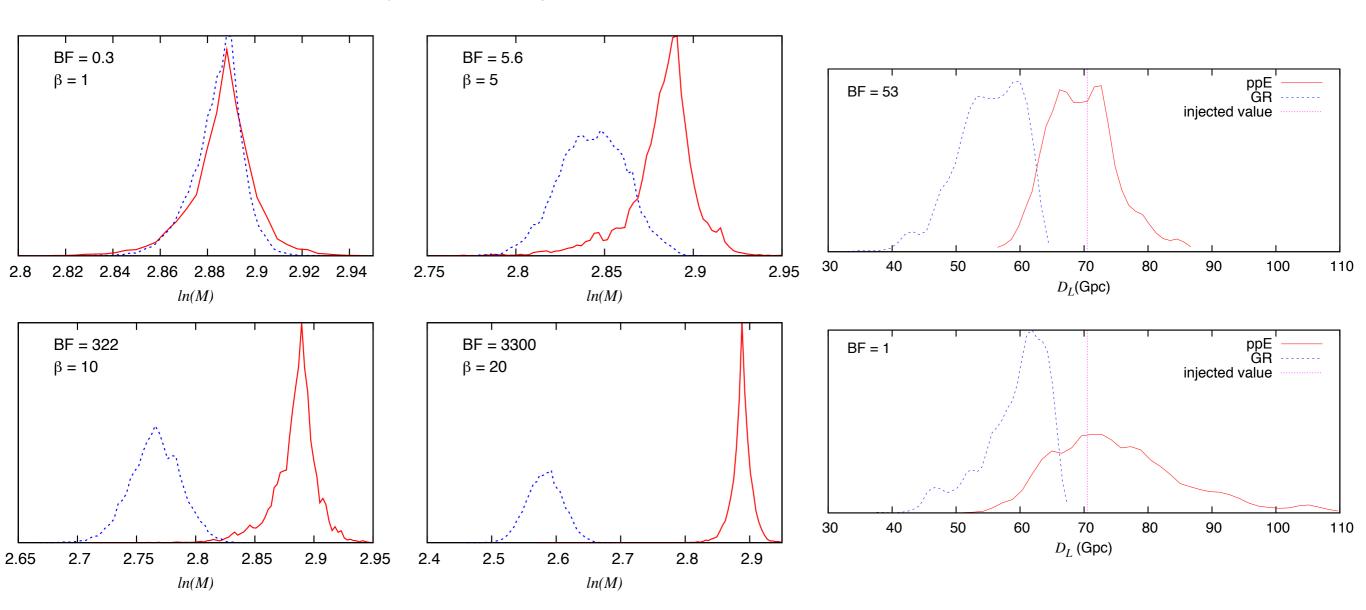
At our doorstep...



Fundamental Bias

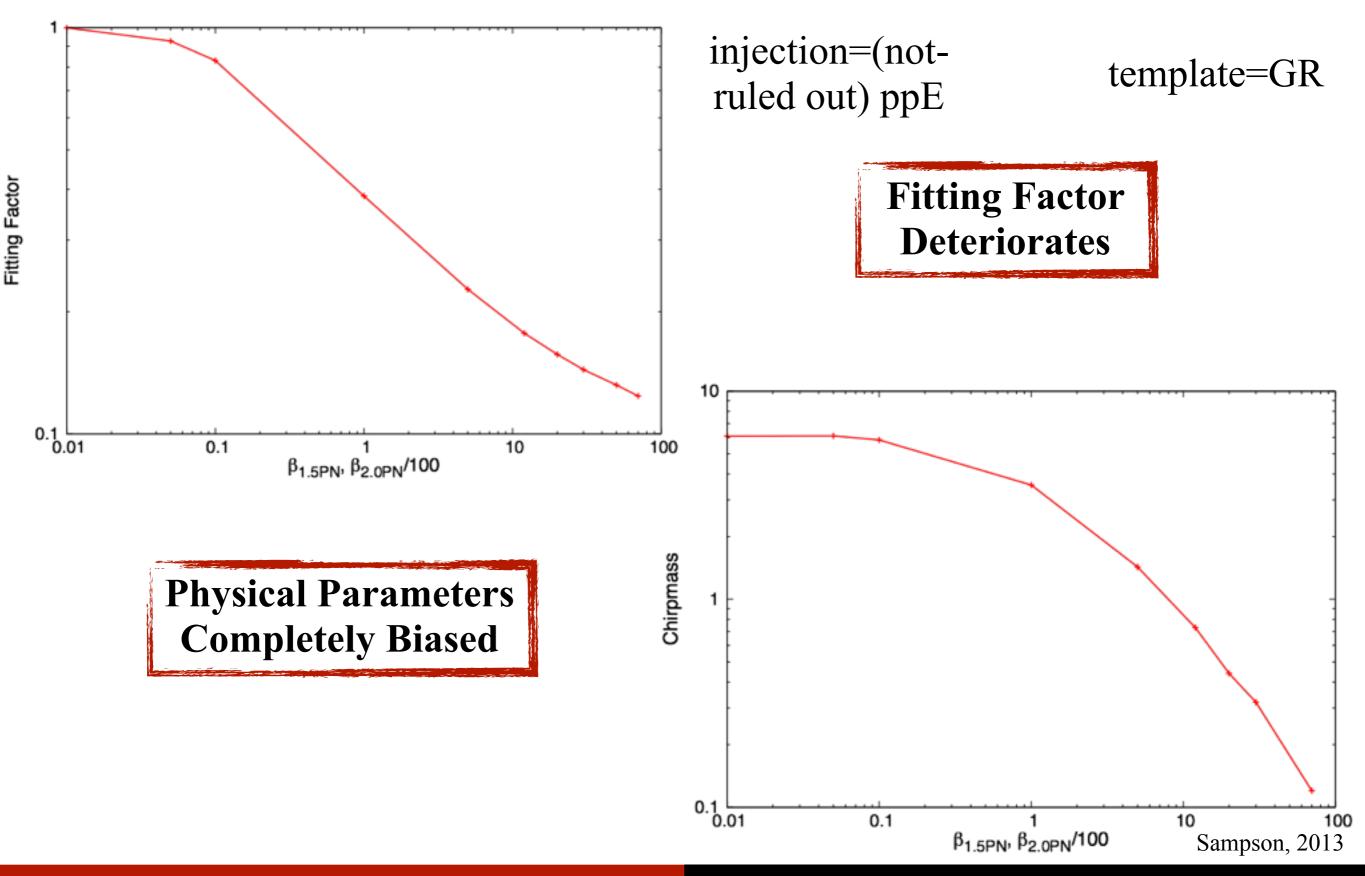
Non-GR Signal/GR Templates, SNR = 20

Non GR injection, extracted with GR templates (blue) and ppE templates (red). GR template extraction is "wrong" by much more than the systematic (statistical) error. "Fundamental Bias"



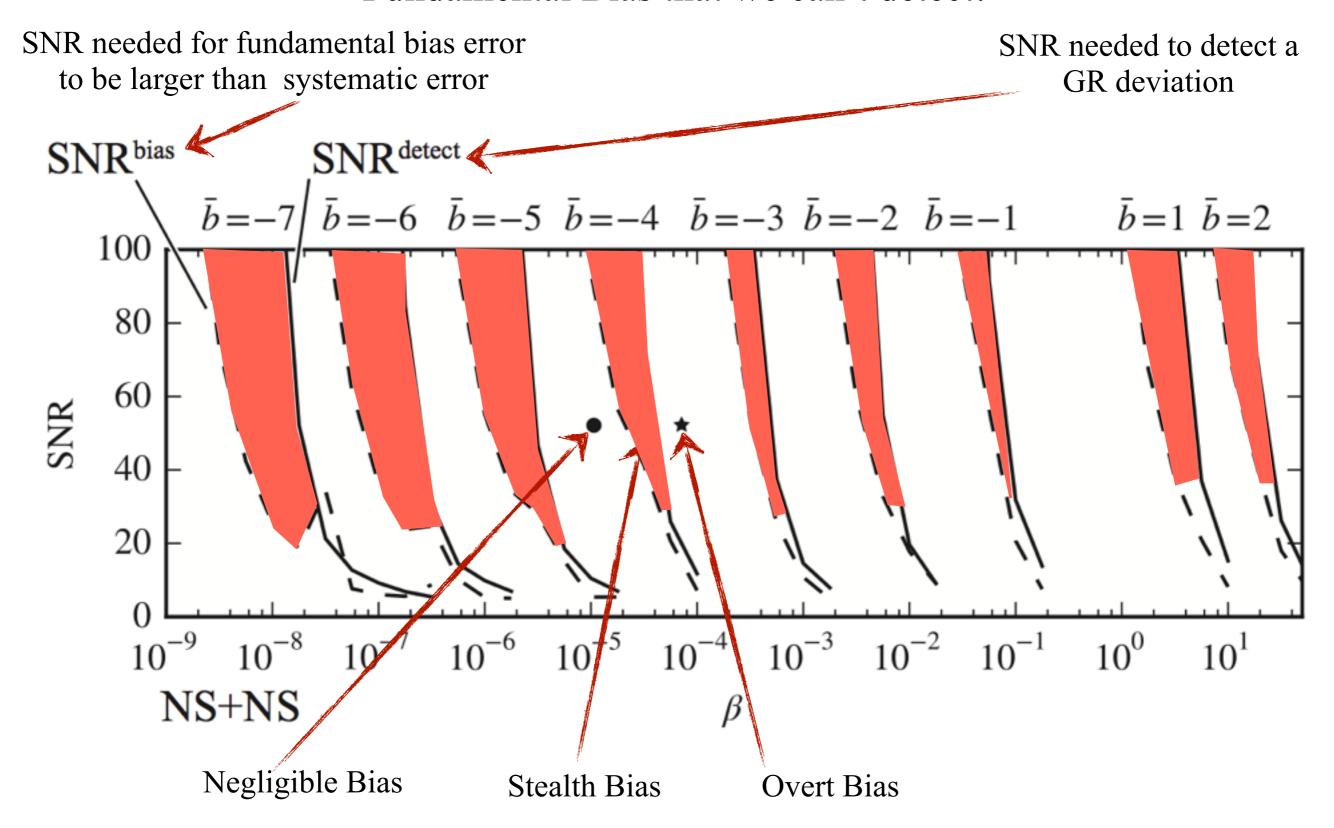
Cornish, Sampson, Yunes & Pretorius, 2011

Ignoring Fundamental Bias...

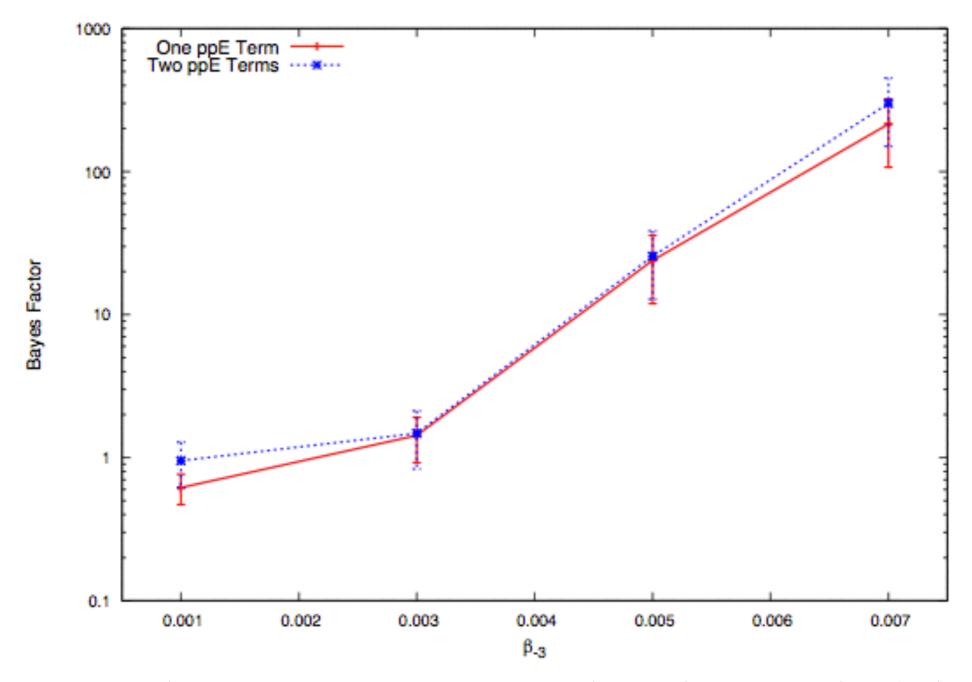


Stealth Bias

Fundamental Bias that we can't detect!



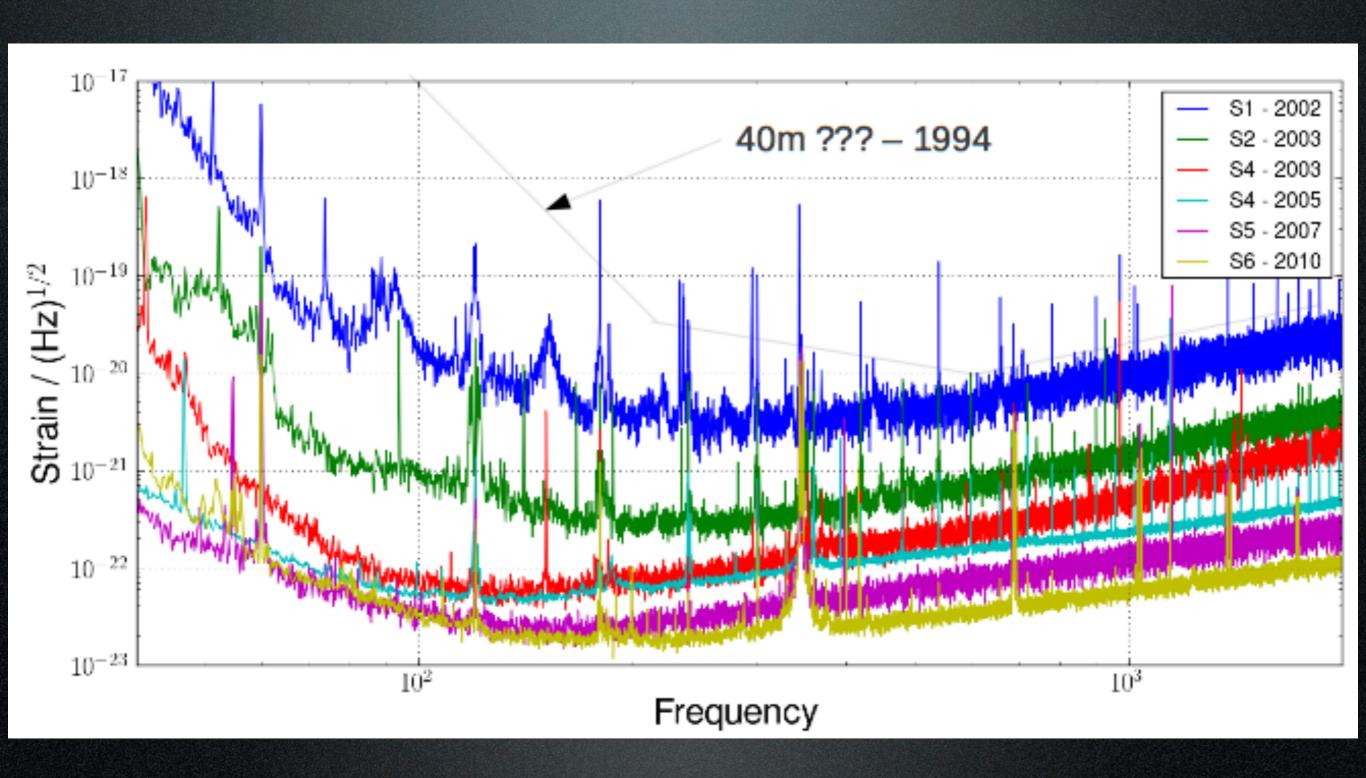
Simple ppE Performance



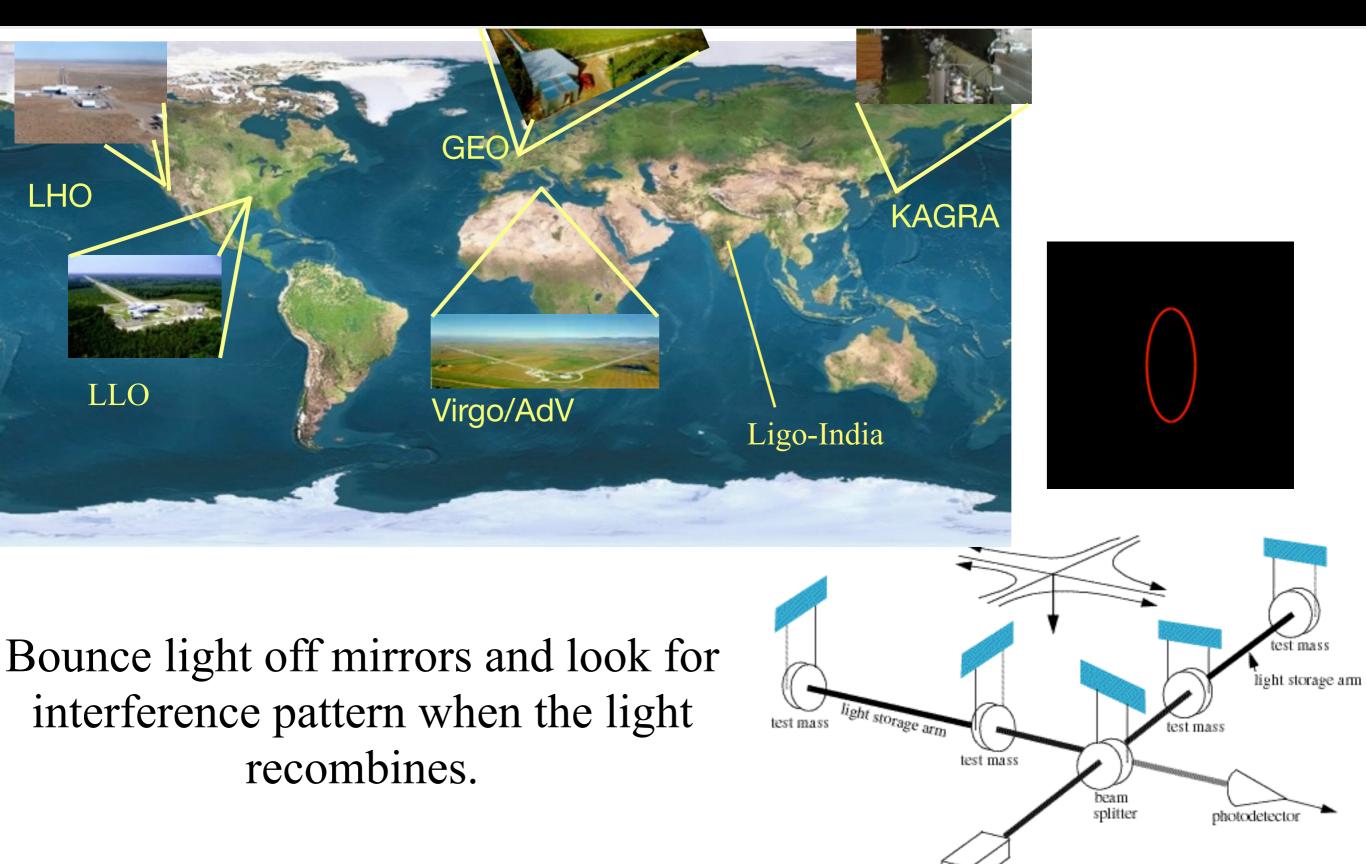
Bayes Factor between a 1-parameter ppE template and a GR template (red) and between a 2-parameter ppE template and a GR template (blue), given a non-GR injection with 3 phase deformations, as a function of the magnitude of the leading-order phase deformation.

Sampson, Cornish & Yunes, 2013

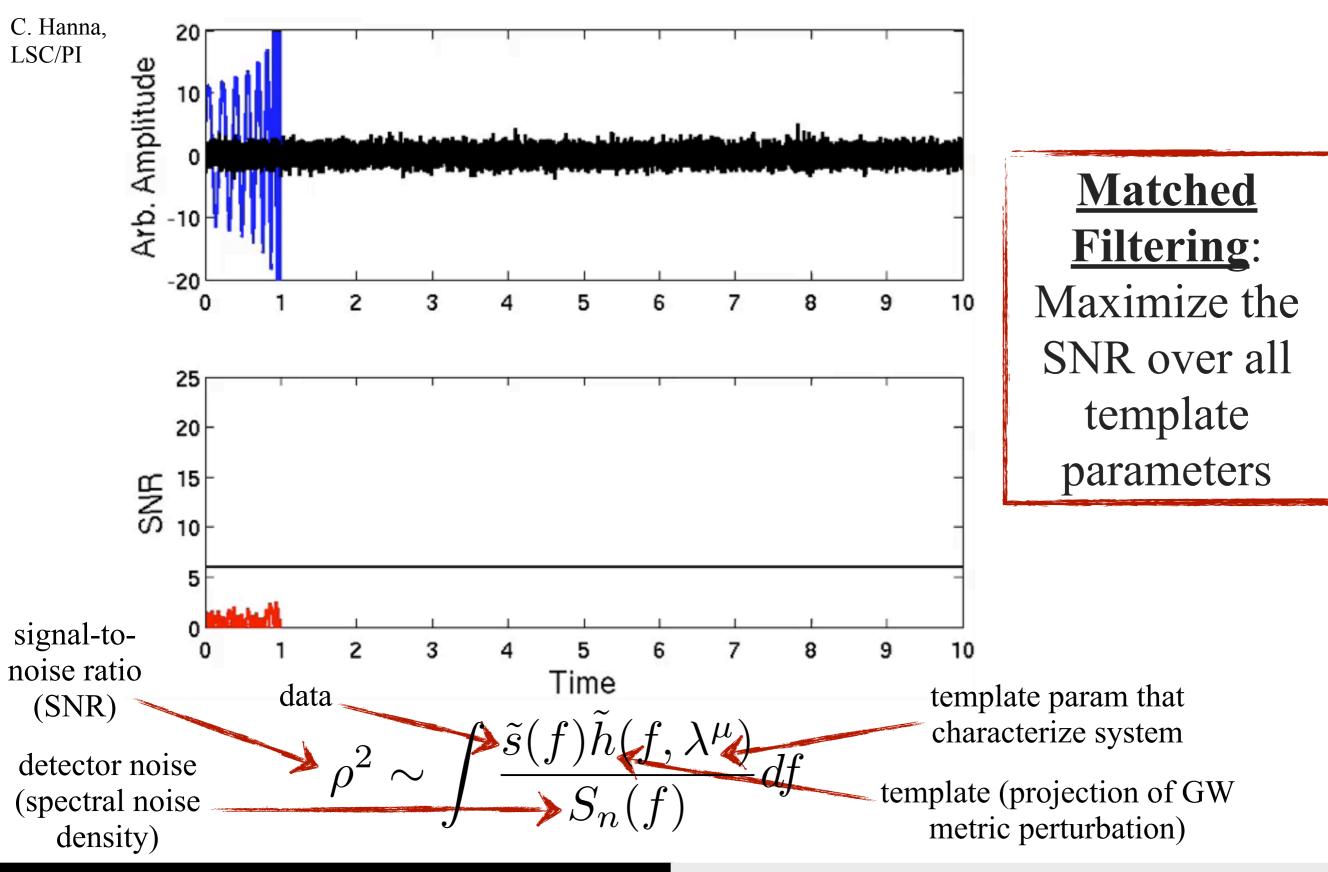
The Need for Accuracy



Detectors



Data Analysis at work



GW Strong-Field Tests

Yunes