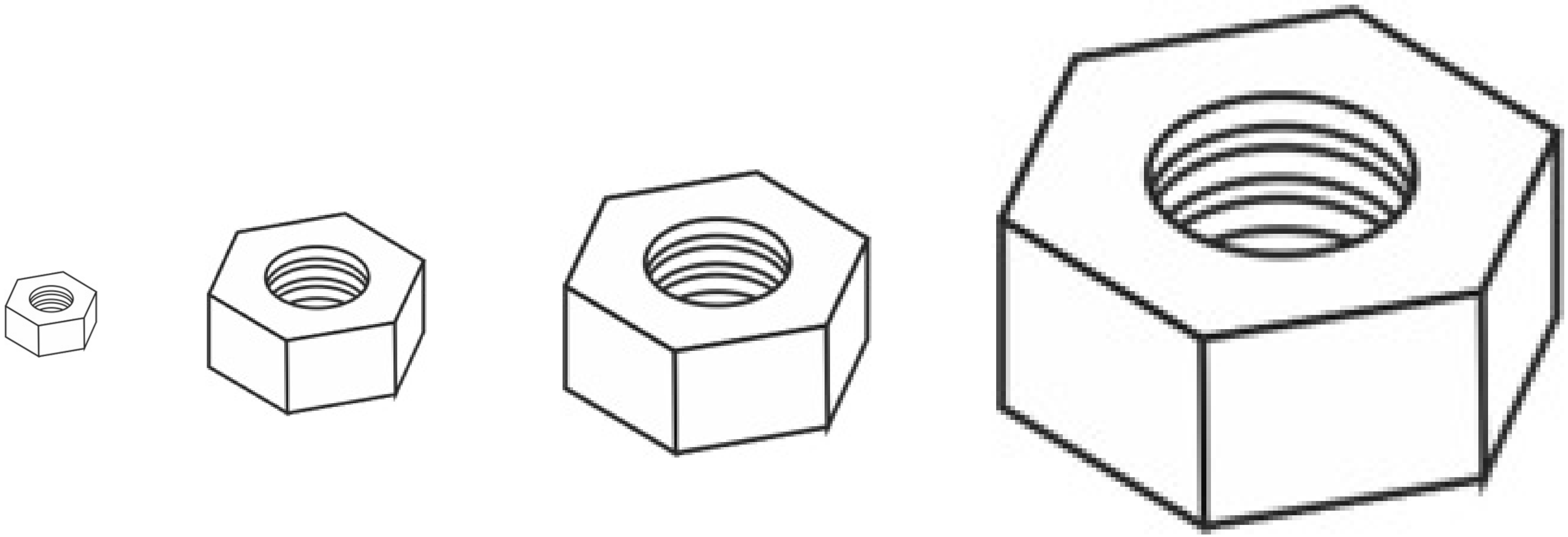
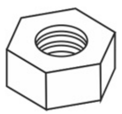
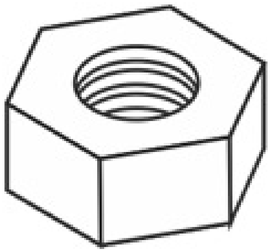


1 eV



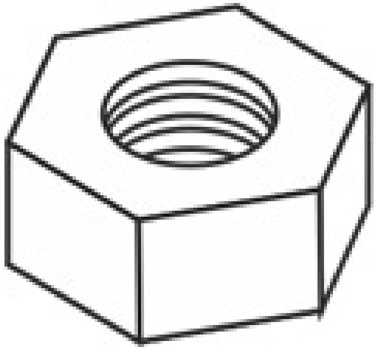
Talk of
Michele
Maltoni

1 KeV
1 GeV



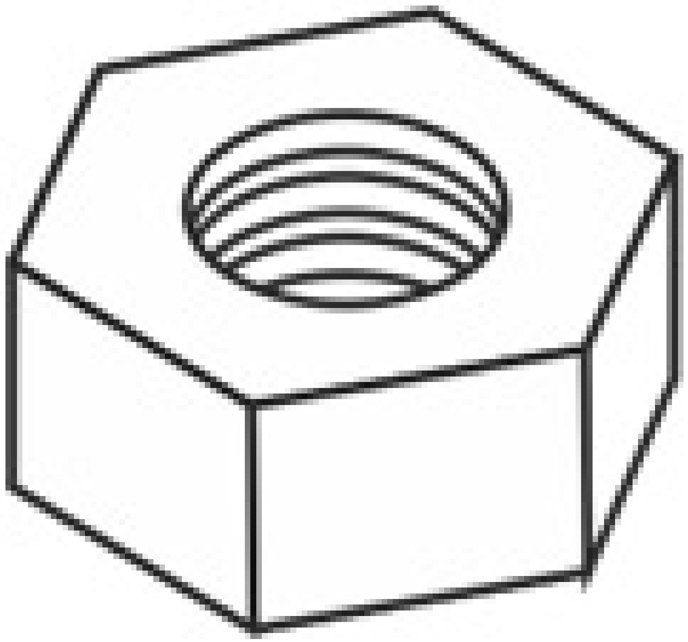
Talk of
Marco
Drewes

1 TeV



This talk

1 EeV



Talk of
Guido
Altarelli

Impact of TeV-scale sterile neutrinos

on precision low-energy observables

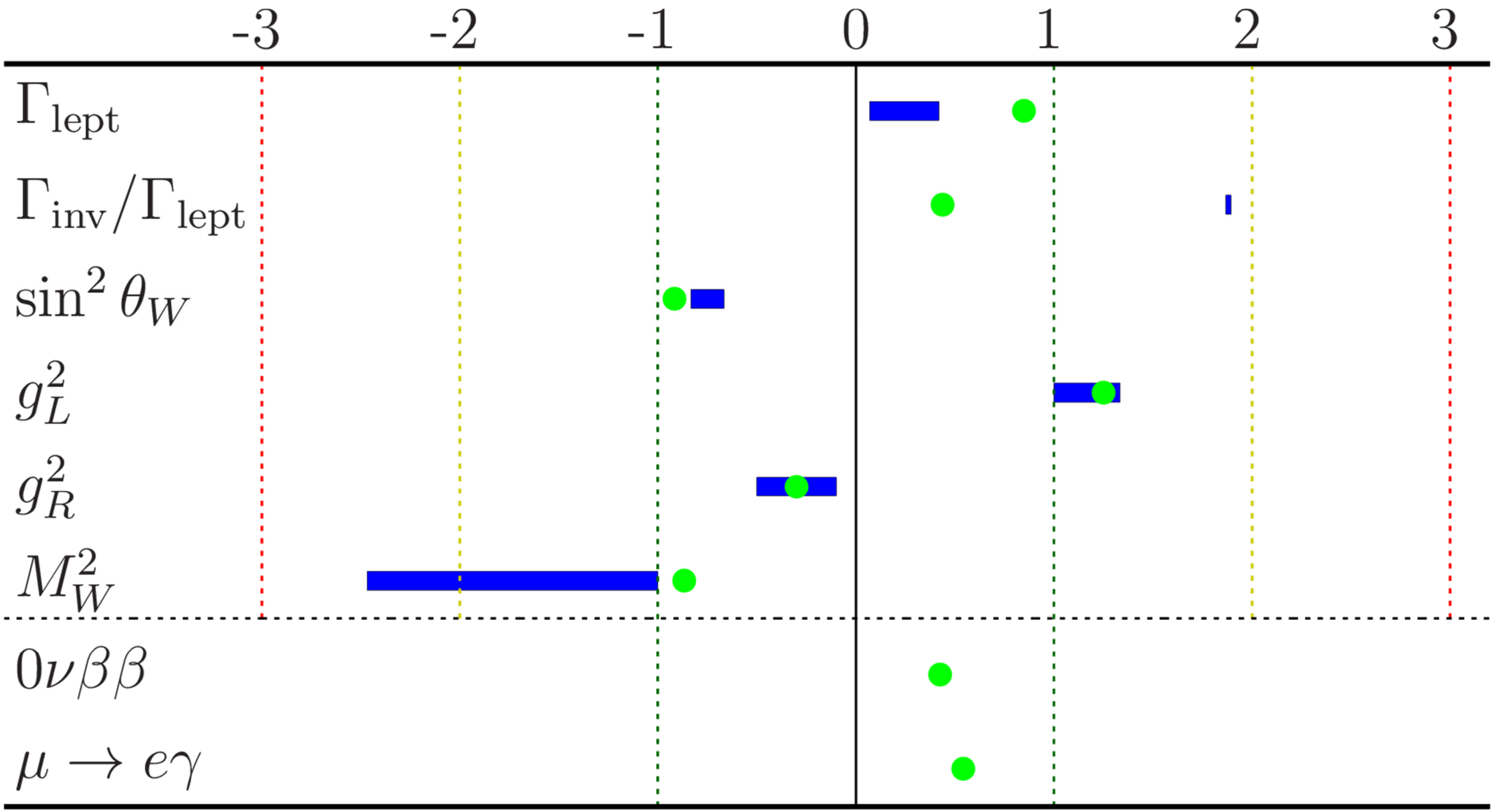
E. Akhmedov, A. Kartavtsev*, M. Lindner, L. Michaels, J. Smirnov

EWPO	Theory (Standard Model)	Experiment
Γ_{lept} (MeV)	84.005 ± 0.015	83.984 ± 0.086
$\Gamma_{\text{inv}}/\Gamma_{\text{lept}}$	5.9721 ± 0.0002	5.942 ± 0.016
$\sin^2\theta_{\text{W}}$	0.23150 ± 0.0001	0.2324 ± 0.0012
g_{L}^2	0.3040 ± 0.0002	0.3026 ± 0.0012
g_{R}^2	0.0300 ± 0.0002	0.0303 ± 0.0010
M_{W} (GeV)	80.359 ± 0.011	80.385 ± 0.015

*Max-Planck Institute for Physics, Munich, Germany
alexander.kartavtsev@mpp.mpg.de

Based on arXiv:1302.1872
DOI: 10.1007/JHEP05(2013)081

Improved electroweak fit



Outline

1. Why are low energy observables affected ?
2. Which observables are affected ?
3. Parameter scan
4. What comes out ?
5. Summary

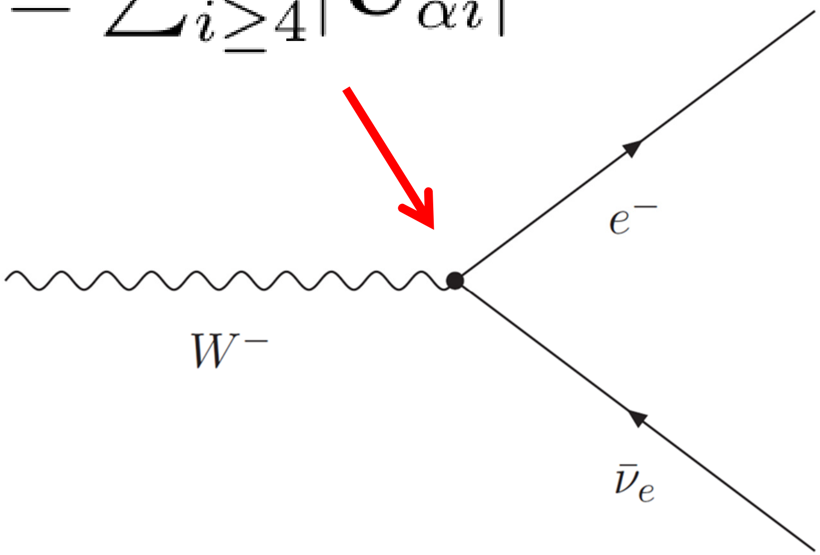
Why are SM observables affected?

PMNS matrix → $\begin{pmatrix} \mathscr{U} & \mathscr{R} \\ \mathscr{W} & \mathscr{V} \end{pmatrix}$

$$\begin{aligned}\mathscr{L}_{\text{int}} = &- \frac{e}{2c_w s_w} Z_\mu \sum_{i,j=1}^{3+n} \sum_{\alpha=e,\mu,\tau} \bar{\nu}_i \mathbf{U}^\dagger_{i\alpha} \gamma^\mu P_L \mathbf{U}_{\alpha j} \nu_j \\ &- \frac{e}{\sqrt{2} s_w} W_\mu \sum_{i=1}^{3+n} \sum_{\alpha=e,\mu,\tau} \bar{\nu}_i \mathbf{U}^\dagger_{i\alpha} \gamma^\mu P_L e_\alpha + \text{h.c.}\end{aligned}$$

Lepton universality constraints

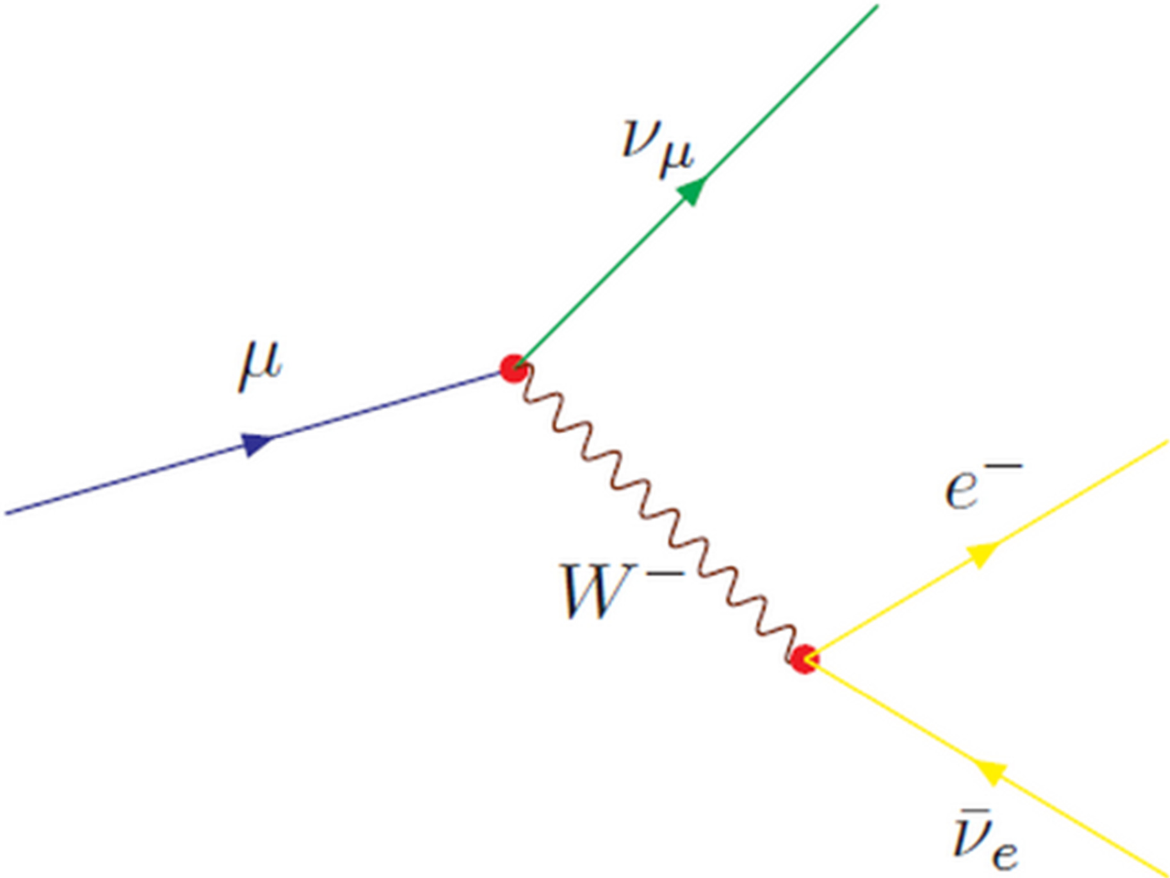
$$\epsilon_\alpha \equiv \sum_{i\geq 4} |\mathbf{U}_{\alpha i}|^2$$



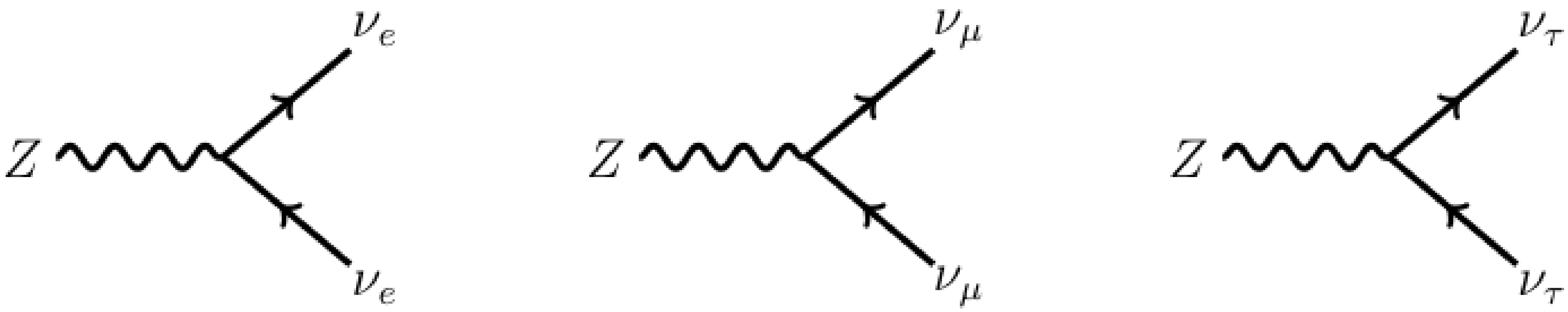
$$\epsilon_e - \epsilon_\mu = 0.0022 \pm 0.0025$$
$$\epsilon_\mu - \epsilon_\tau = 0.0017 \pm 0.0038$$
$$\epsilon_e - \epsilon_\tau = 0.0039 \pm 0.0040$$

Loinaz et. al., hep-ph/0403306

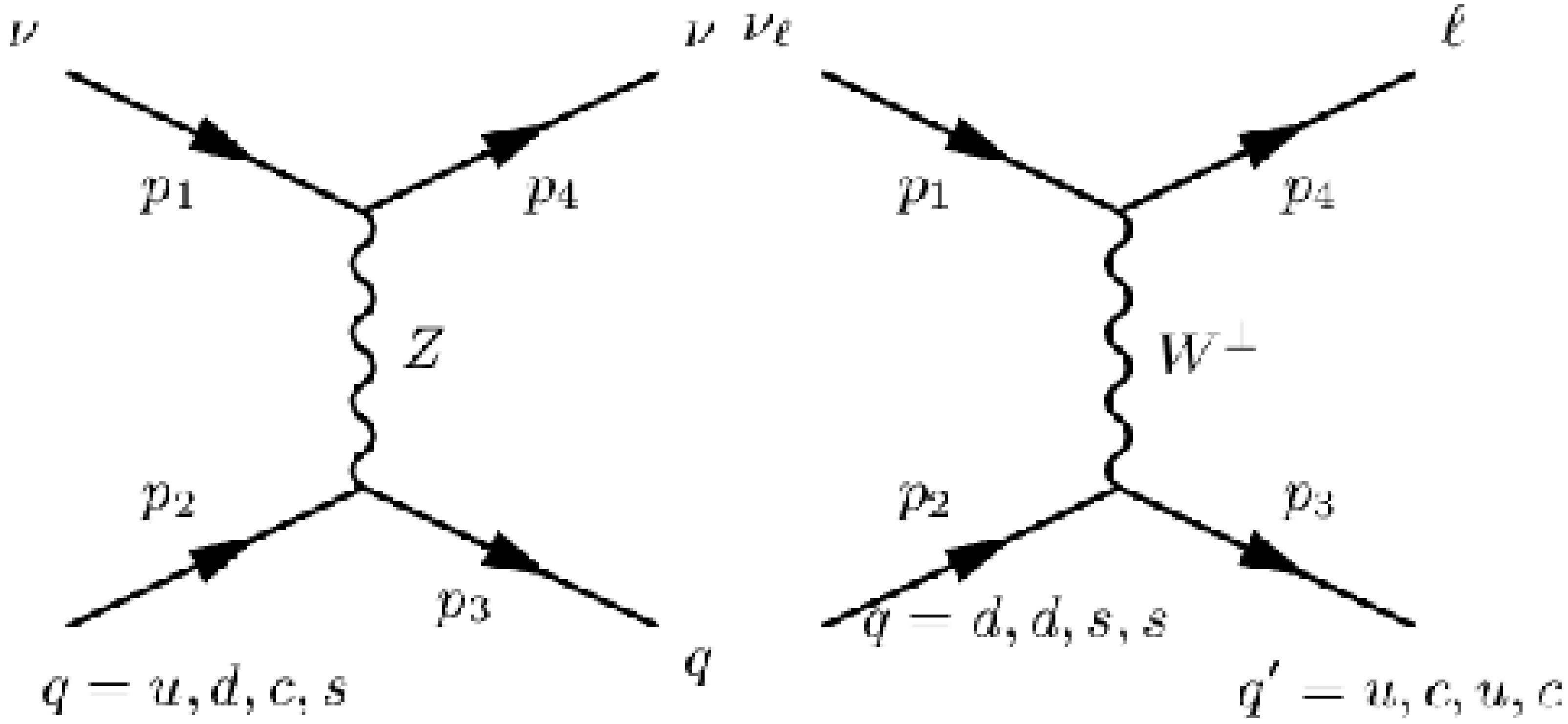
W-boson mass



Invisible Z-decay width

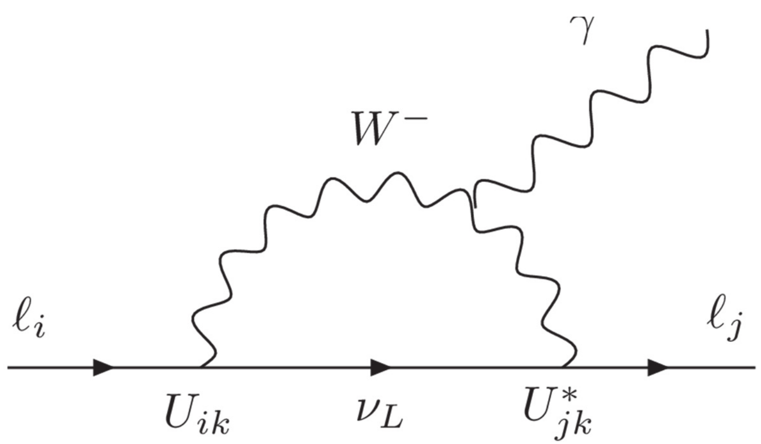


Charged to neutral current ratio



Park, Kwangwoo et al. arXiv:0910.5013

Lepton-flavor violating decays



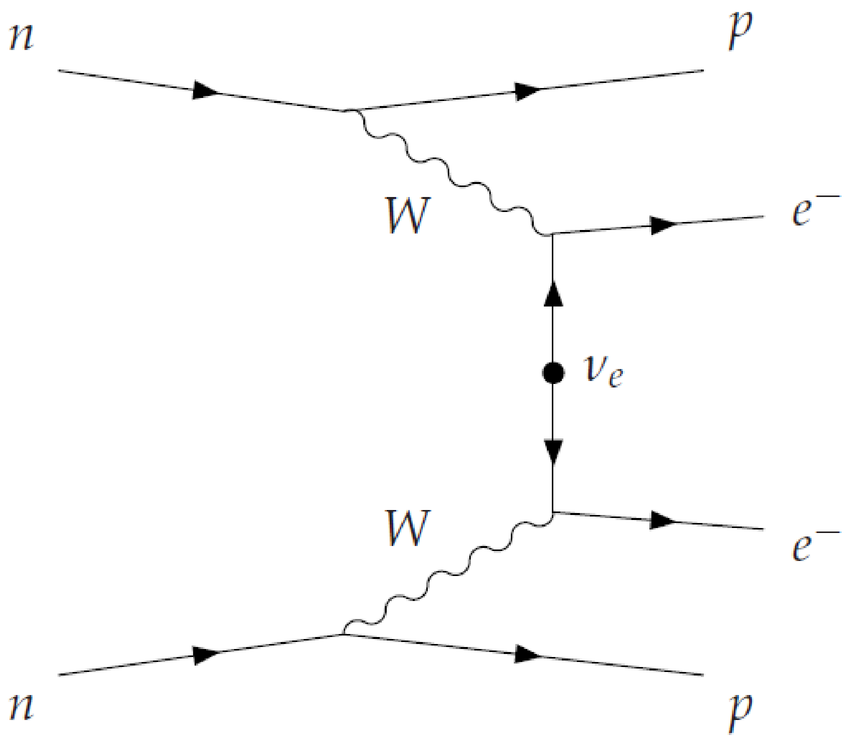
Abada, arXiv:1110.6507

$$\mathrm{BR}(\mu \to e\gamma) = \frac{\Gamma(\mu \to e\gamma)}{\Gamma(\mu \to e\nu\bar{\nu})} = \frac{3\alpha}{32\pi}|\delta_\nu|^2$$

$$\delta_\nu = 2\sum_{i=4}^{3+n} \mathbf{U}^*_{ei}\mathbf{U}_{\mu i}\left[g\left(m_i^2/M_W^2\right) - 5/3\right]$$

$$\mathrm{BR}(\mu^+ \to e^+\gamma) \leq 5.7 \cdot 10^{-13}$$

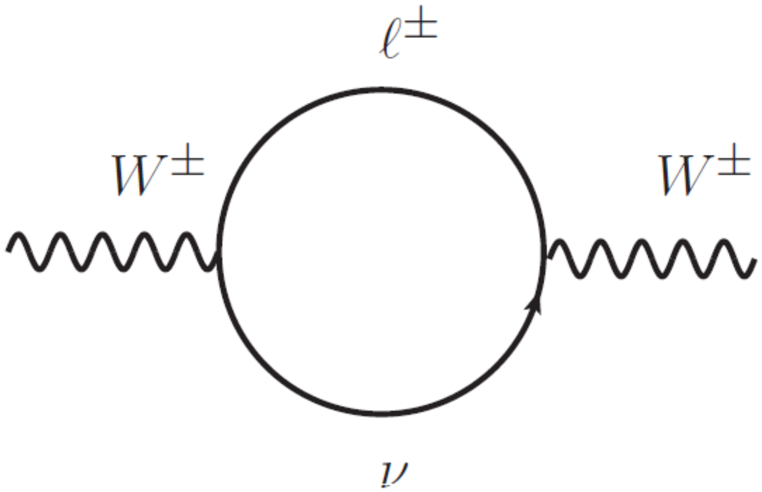
Neutrinoless double-beta decay

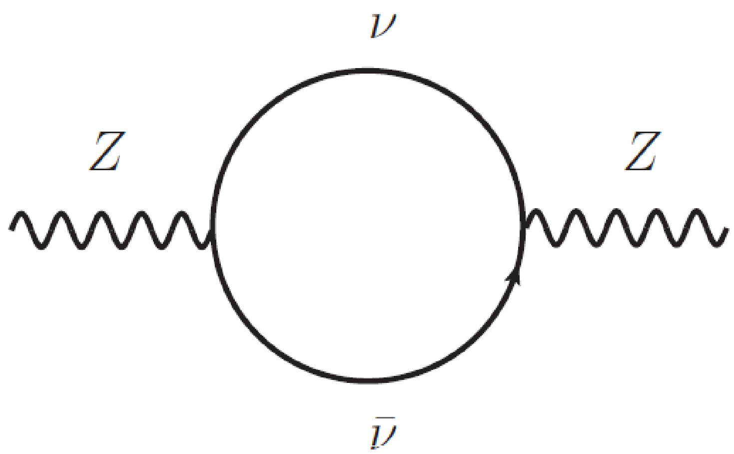


$$|\langle m_{ee}\rangle| \approx \Big|\sum_{i=1}^{3}\mathbf{U}_{ei}^2 m_i - \sum_{i=4}^{3+n} F(A,M_i)\mathbf{U}_{ei}^2 m_i\Big|$$

$$|\langle m_{ee}\rangle| < 0.2 - 0.4\ \mathrm{eV}$$

STU parameters





Tree-level and loop corrections

$$\frac{\Gamma_{\rm lept}}{[\Gamma_{\rm lept}]_{\rm SM}} = 1 + 0.6\,(\epsilon_e + \epsilon_\mu + 0.0145\,T) - 0.0021\,S\,,$$

$$\frac{\Gamma_{\rm inv}/\Gamma_{\rm lept}}{[\Gamma_{\rm inv}/\Gamma_{\rm lept}]_{\rm SM}} = 1 - 0.67\,(\epsilon_e + \epsilon_\mu + \epsilon_\tau) + 0.0021\,S - 0.0015\,T\,,$$

$$\frac{\sin^2\theta_{\rm w}^{\rm lept}}{\left[\sin^2\theta_{\rm w}^{\rm lept}\right]_{\rm SM}} = 1 - 0.72\,(\epsilon_e + \epsilon_\mu + 0.0145\,T) + 0.0016\,S\,,$$

$$\frac{g_L^2}{[g_L^2]_{\rm SM}} = 1 + 0.41\,\epsilon_e - 0.59\,\epsilon_\mu - 0.0090\,S + 0.0022\,T\,,$$

$$\frac{g_R^2}{[g_R^2]_{\rm SM}} = 1 - 1.4\,\epsilon_e - 2.4\,\epsilon_\mu + 0.031\,S - 0.0067\,T\,,$$

$$\frac{M_{\rm W}}{[M_{\rm W}]_{\rm SM}} = 1 + 0.11\,\epsilon_e + 0.11\,\epsilon_\mu - 0.0036\,S + 0.0056\,T + 0.0042\,U\,.$$

Loinaz et. al., hep-ph/0403306

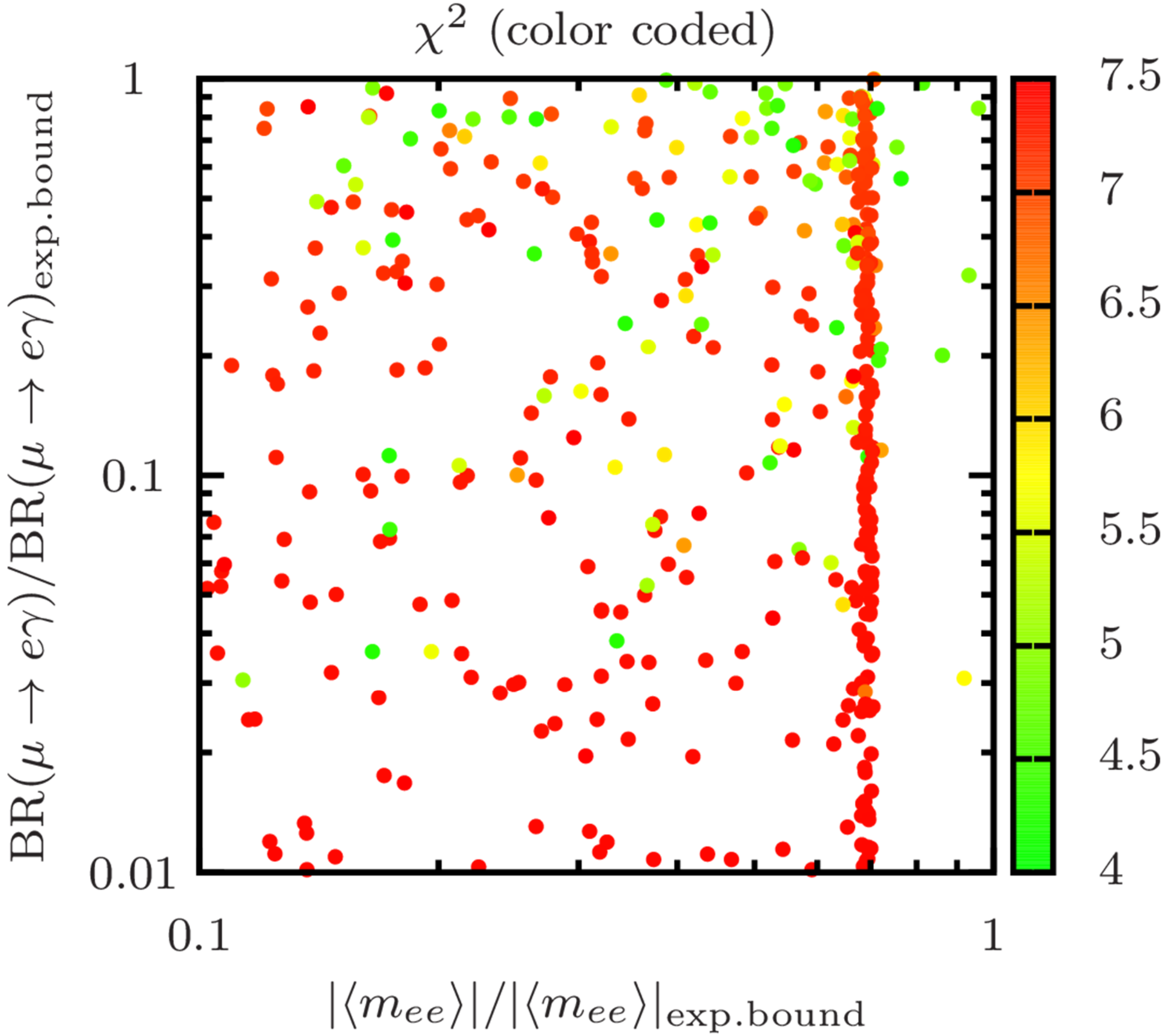
Parameter scan

$$\mathscr{R} = -i\,\mathcal{U}\,\hat{m}_{\text{light}}^{\frac{1}{2}}\,O^{*}\,\hat{m}_{\text{heavy}}^{-\frac{1}{2}}\,,$$

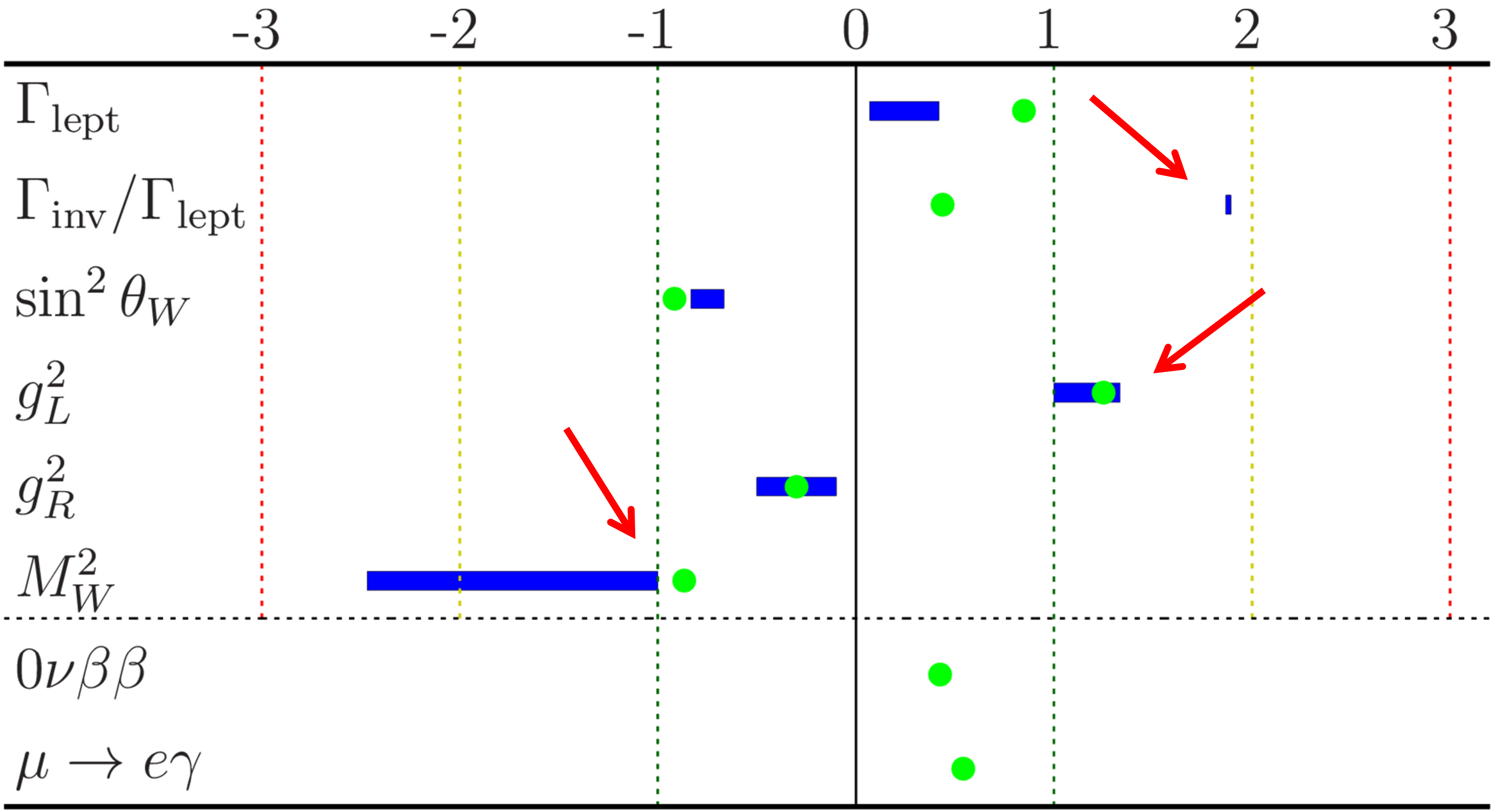
$$\mathscr{U} = \left(1 - \mathscr{R}\,\mathscr{R}^{\dagger}\right)^{\frac{1}{2}}\,\mathcal{U}\,,$$

M1, M2, M3	3	S, T, U	6	~~S~~, T, ~~U~~	4
α1, α2, α3	6	ε1, ε2, ε3		ε1, ε2, ε3	

Normal hierarchy: rare processes



Normal hierarchy: EW fit



Conclusions

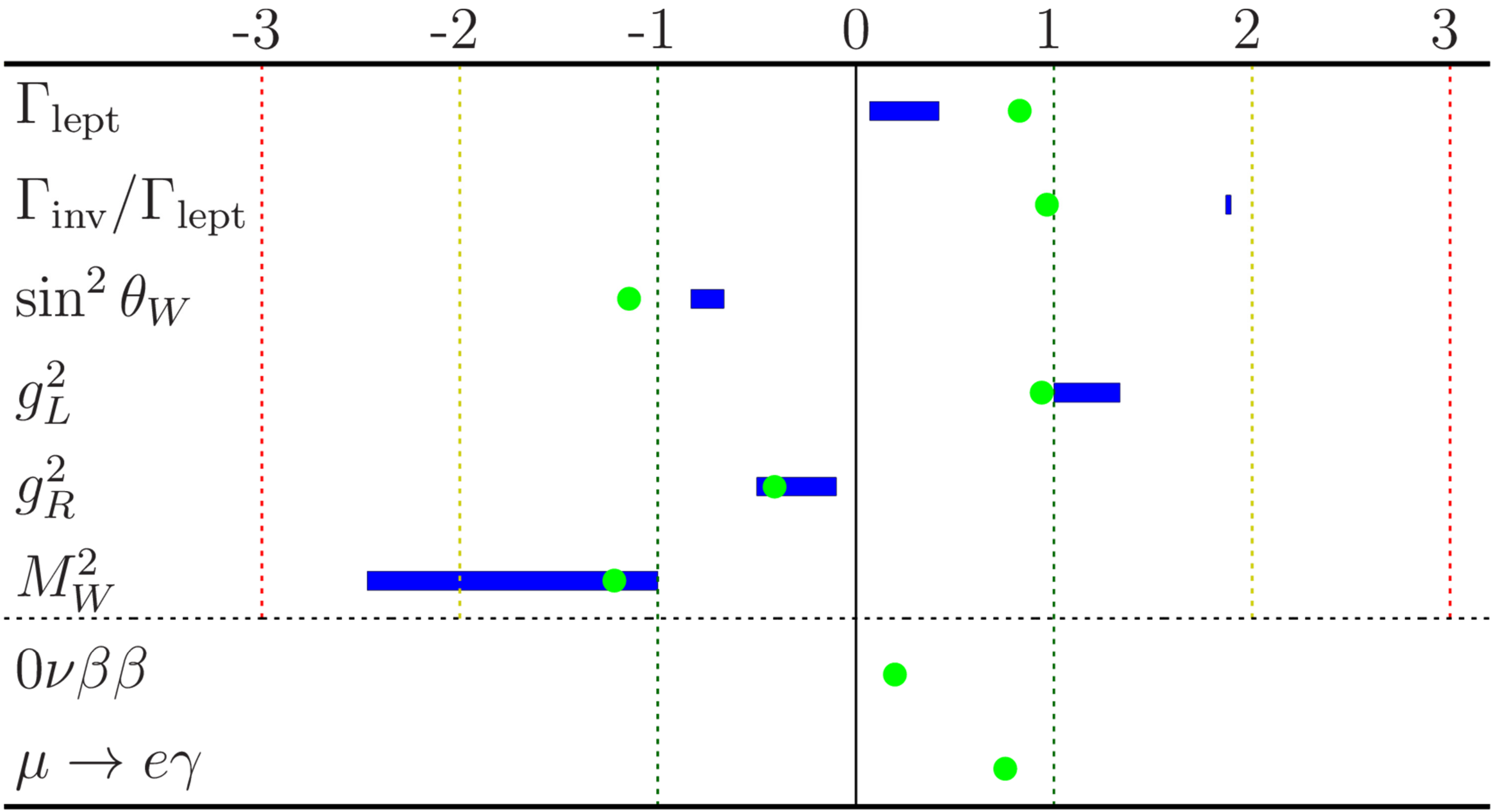
RH neutrinos affect low energy observables

- via induced non-unitarity of PMNS matrix
- via loop corrections

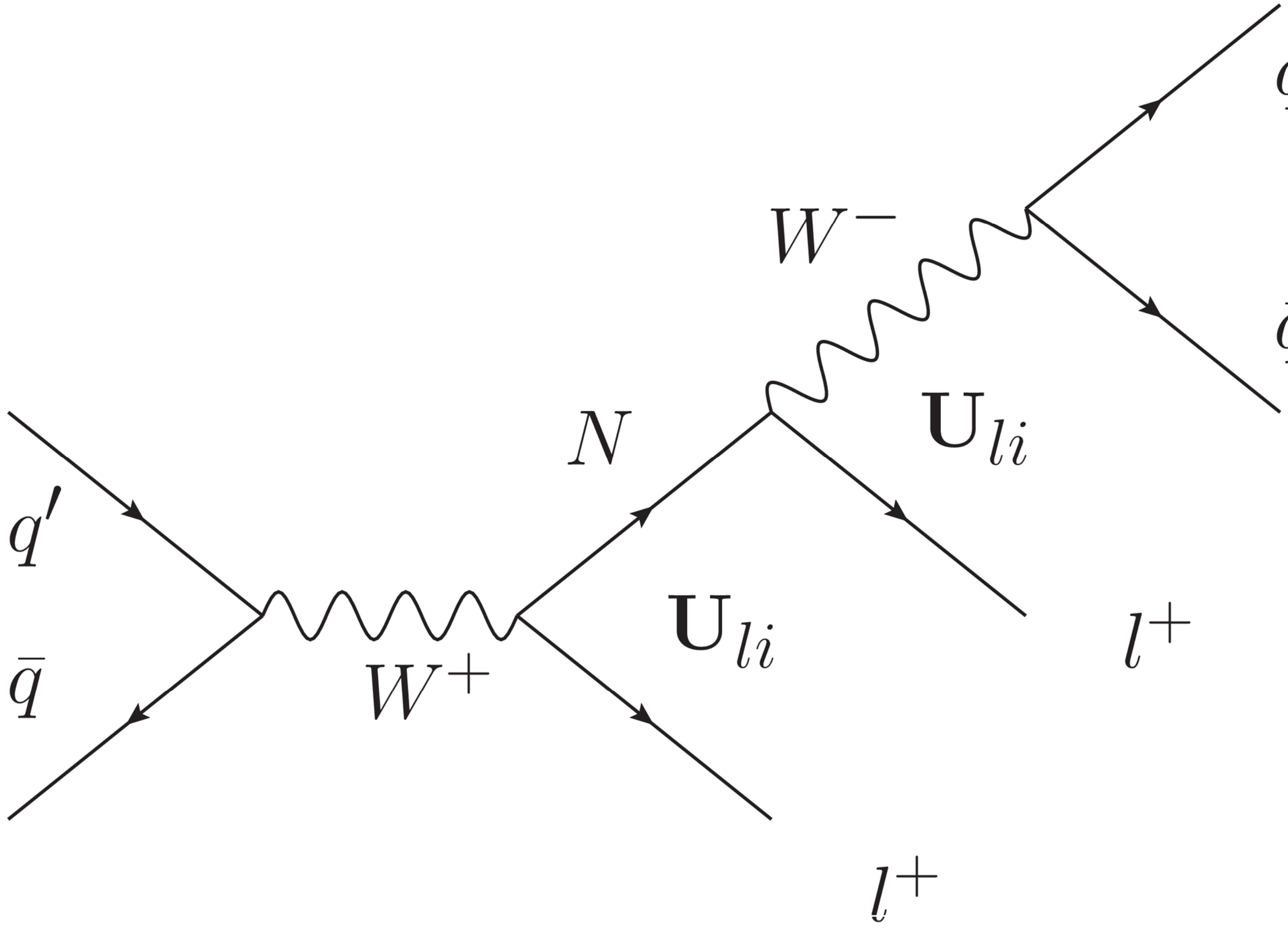
With some fine tuning the fit can be improved

- neutral-to-charged current ratio
- W-boson mass
- to a lesser extent invisible Z-decay width

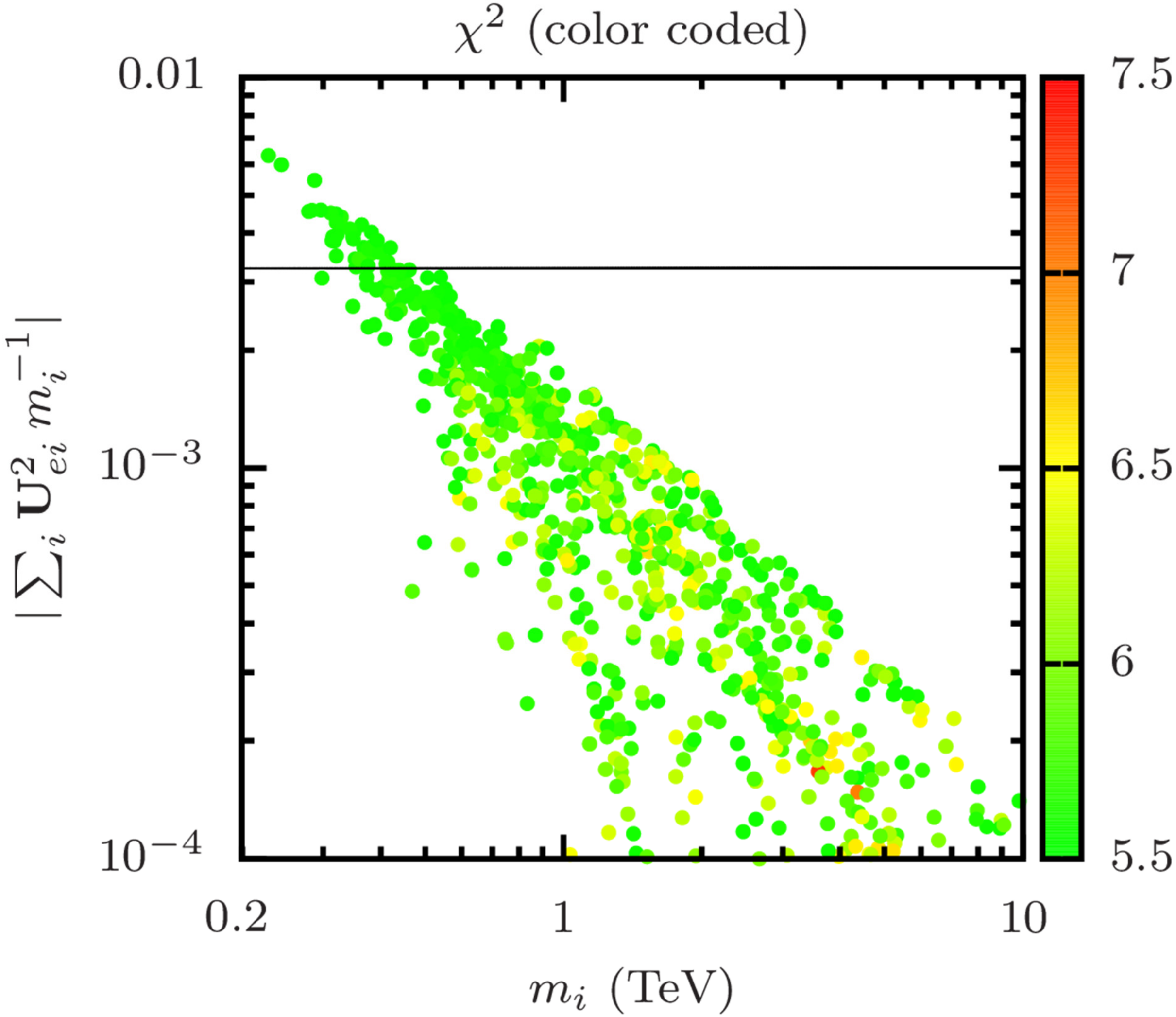
Inverse hierarchy: EW fit



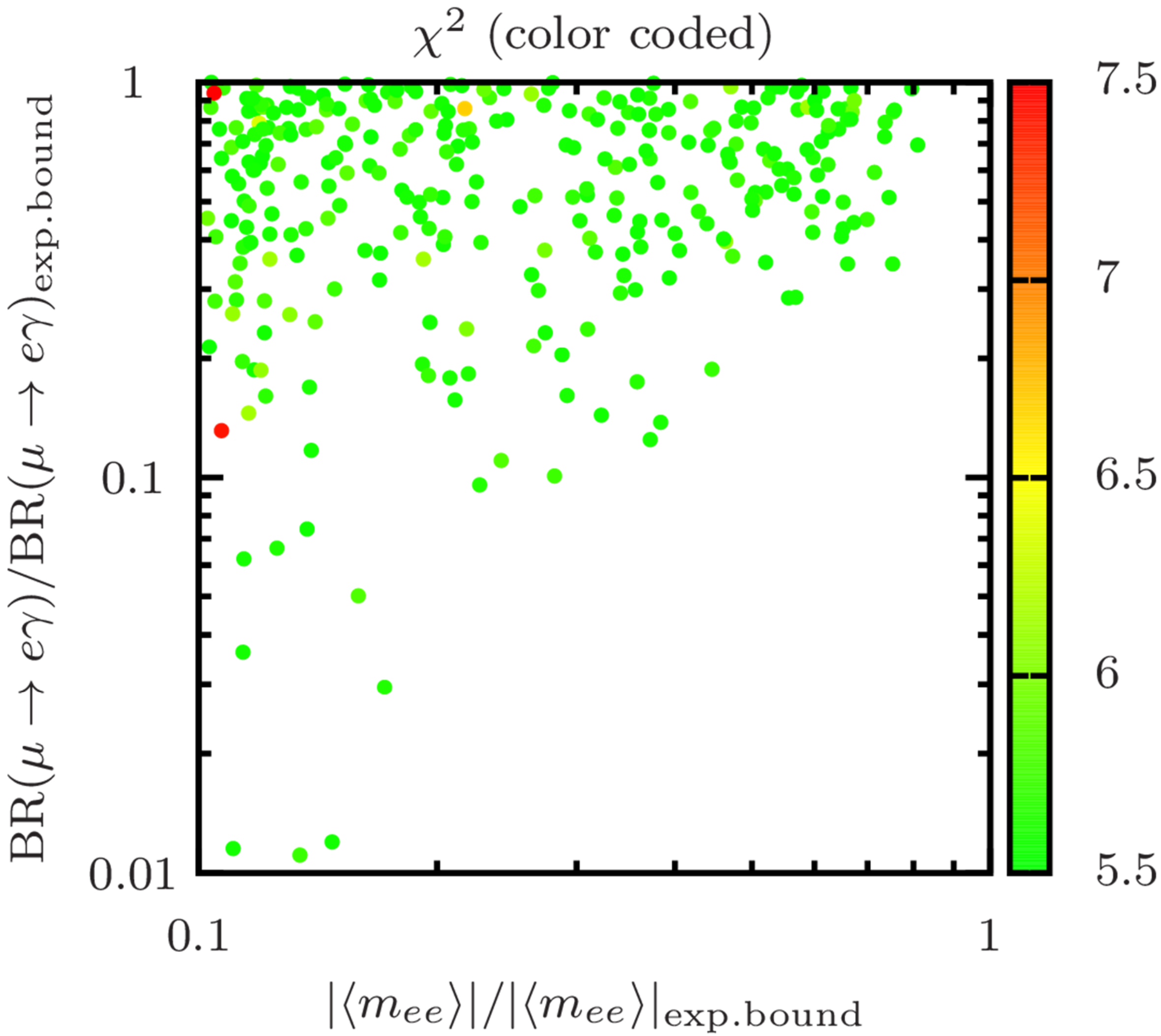
Heavy neutrinos in colliders



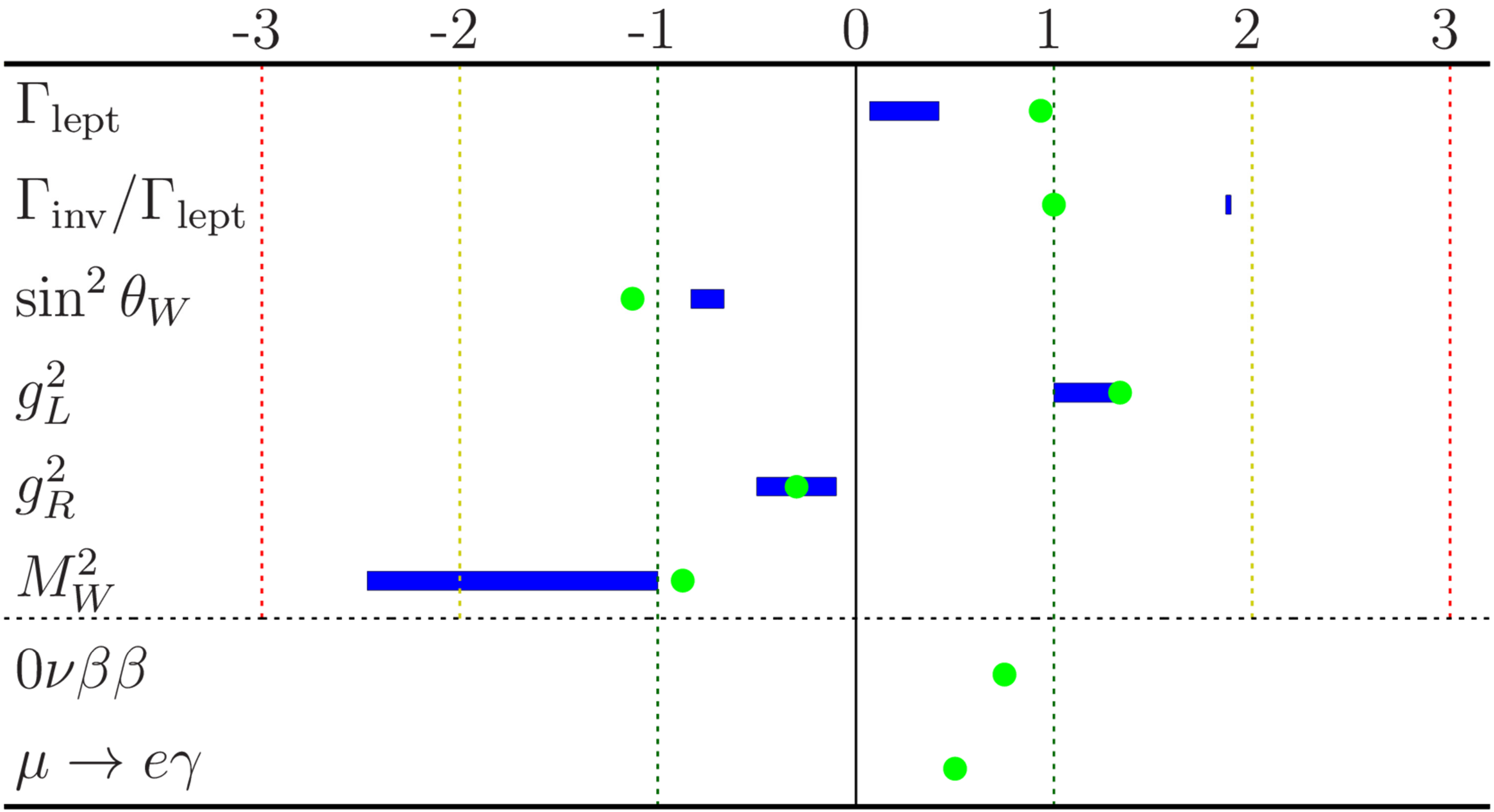
Inverse hierarchy: LHC



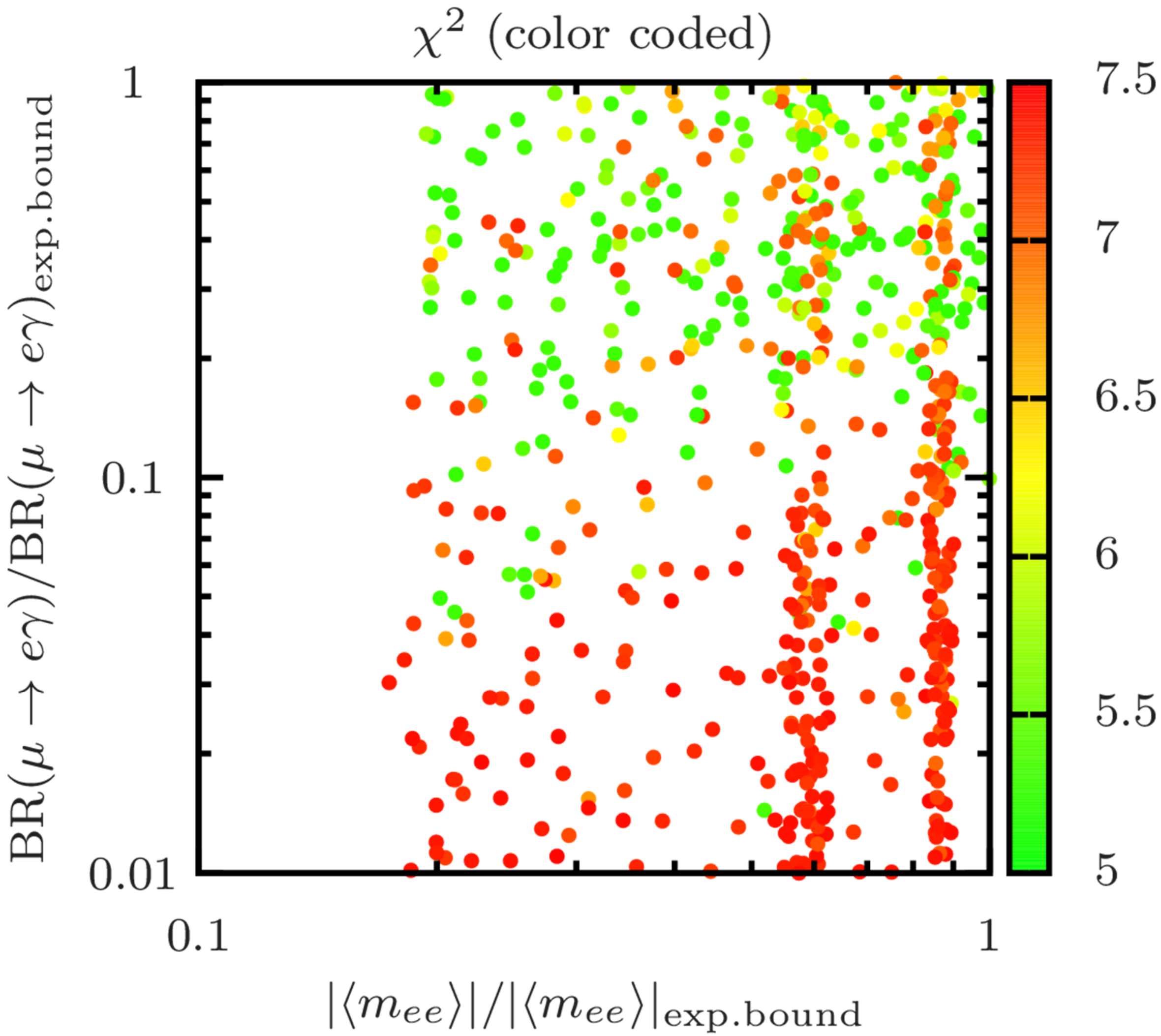
Inverse hierarchy: rare processes



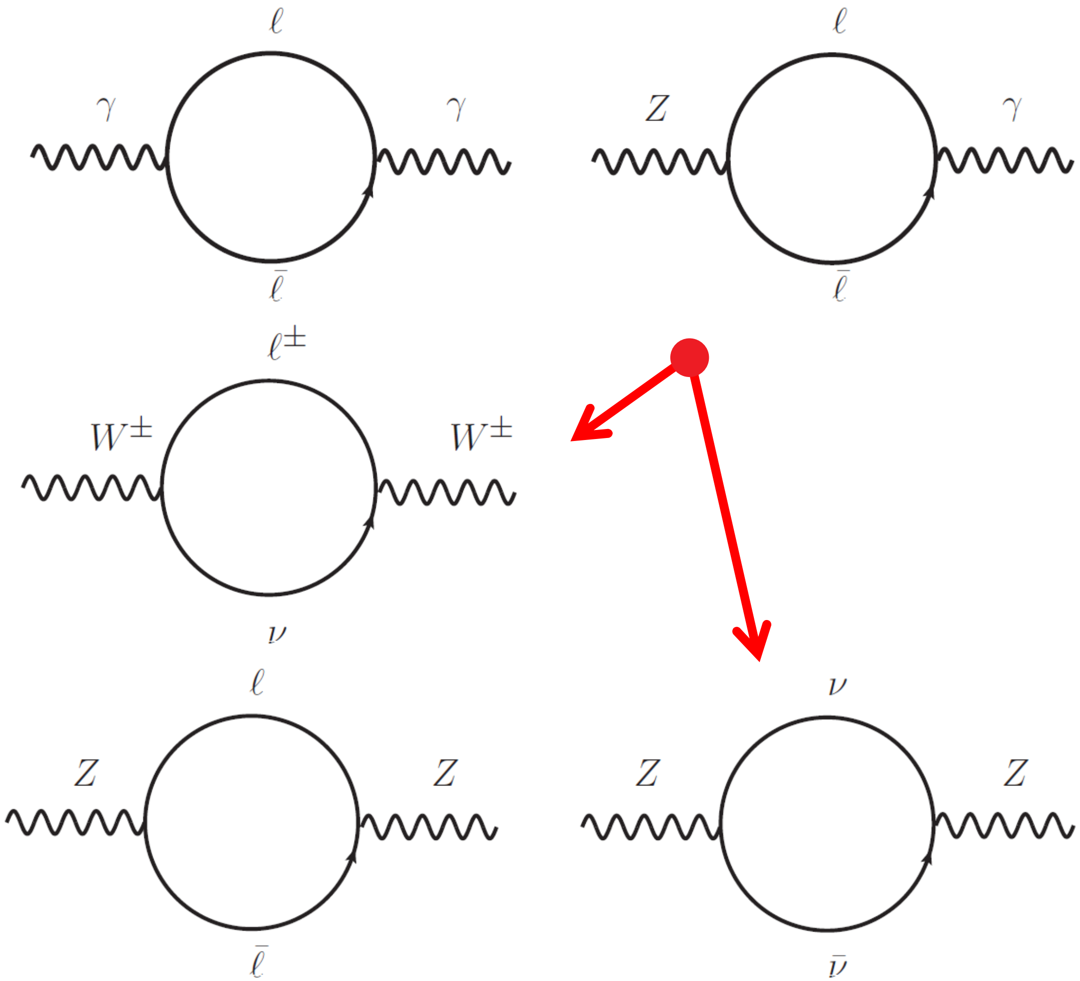
Quasidegenerate: EW fit



Quasidegenerate: rare processes



Self-energy contributions



STU parameters

$$S = \frac{4s_w^2c_w^2}{M_Z^2}\left[\Pi_{ZZ}(0) - \Pi_{ZZ}(M_Z^2) + \Pi_{\gamma\gamma}(M_Z^2) - \frac{c_w^2 - s_w^2}{c_w s_w}\Pi_{Z\gamma}(M_Z^2)\right]$$

$$T = \frac{\Pi_{ZZ}(0)}{M_Z^2} - \frac{\Pi_{WW}(0)}{M_W^2}$$

$$U = 4s_w^2c_w^2\left[\frac{\Pi_{WW}(0) - \Pi_{WW}(M_W^2)}{c_w^2 M_W^2} - \frac{\Pi_{ZZ}(0) - \Pi_{ZZ}(M_Z^2)}{M_Z^2} + \frac{s_w^2}{c_w^2}\frac{\Pi_{\gamma\gamma}(M_Z^2)}{M_Z^2} - 2\frac{s_w}{c_w}\frac{\Pi_{Z\gamma}(M_Z^2)}{M_Z^2}\right]$$

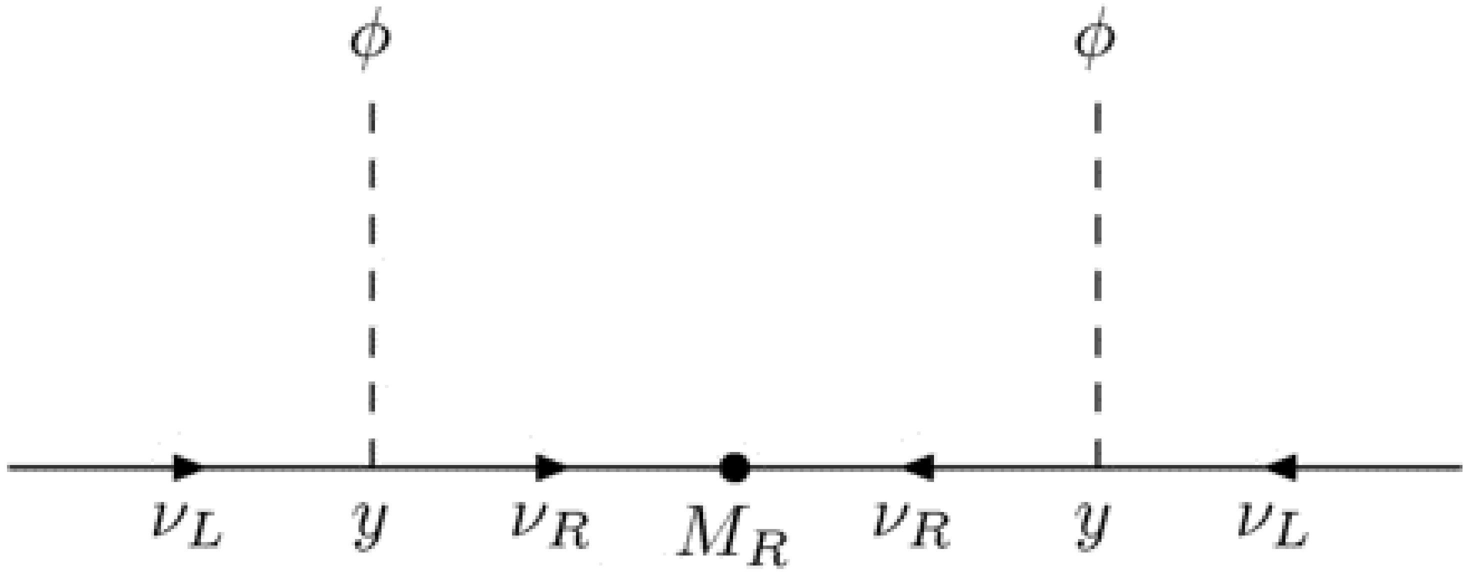
Explicit expressions for T

$$T_{\text{tot}} = T_N + T_{\text{SM}} = -\frac{1}{8\pi s_w^2 M_W^2}$$
$$\times \left[\sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}^\dagger_{i\alpha} \mathbf{U}_{\alpha j} \mathbf{U}^\dagger_{j\beta} \mathbf{U}_{\beta i}\, Q(0, m_i^2, m_j^2)\right.$$
$$+ \sum_{i,j=1}^{3+n} \sum_{\alpha\beta} \mathbf{U}^\dagger_{i\alpha} \mathbf{U}_{\alpha j} \mathbf{U}^\dagger_{i\beta} \mathbf{U}_{\beta j} m_i m_j B_0(0, m_i^2, m_j^2)$$
$$- 2 \sum_{i=1}^{3+n} \sum_{\alpha} \mathbf{U}^\dagger_{i\alpha} \mathbf{U}_{\alpha i}\, Q(0, m_i^2, m_\alpha^2)$$
$$\left. + \sum_{\alpha} m_\alpha^2 B_0(0, m_\alpha^2, m_\alpha^2)\right],$$

Non-unitarity in oscillations

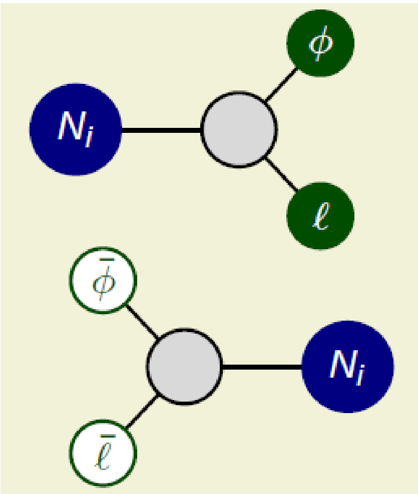
$$P_{\alpha\beta}(L=0) = \frac{\delta_{\alpha\beta}\,(1-2\epsilon_\alpha) + \epsilon_\alpha\epsilon_\beta}{(1-\epsilon_\alpha)(1-\epsilon_\beta)}$$

See-saw mechanism

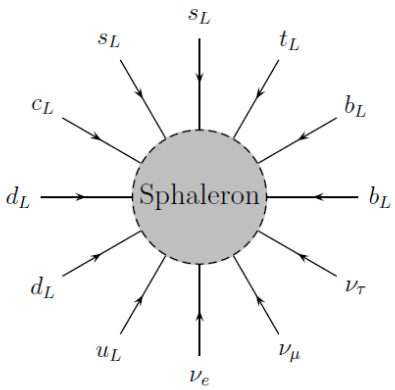


$$\mathcal{M}_\nu = \mathcal{M}_D^T \mathcal{M}_M^{-1} \mathcal{M}_D$$

Baryogenesis via leptogenesis



→ ΔL →



→ ΔB →