

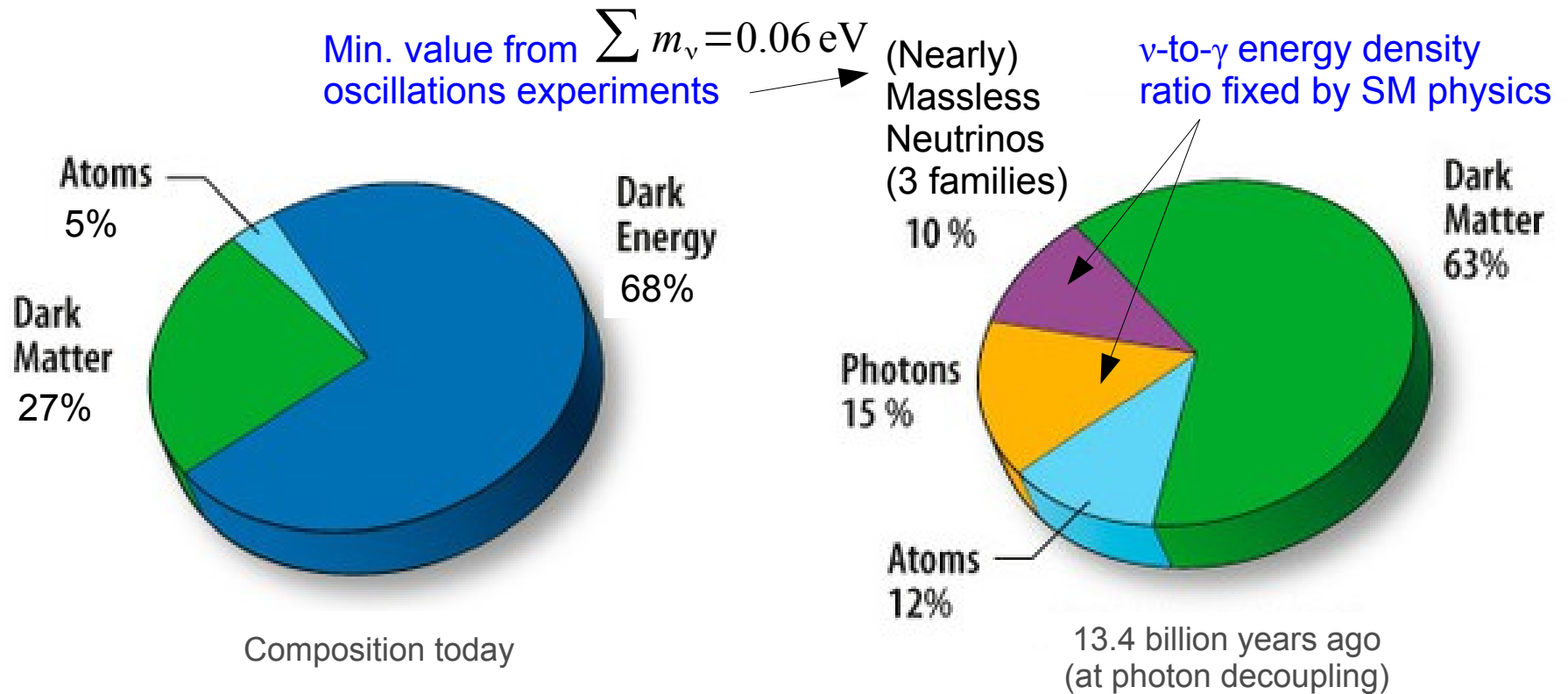
Neutrino properties from cosmology

Yvonne Y. Y. Wong
The University of New South Wales
Sydney, Australia

Rencontres de Moriond EW 2014, La Thuile, March 15 – 22, 2014

The concordance flat Λ CDM model...

The **simplest** model consistent with **present observations**.



Plus flat spatial geometry+initial conditions from single-field inflation

The neutrino sector beyond Λ CDM...

There are many ways in which the neutrino sector can be **extended beyond the standard picture**.

Neutrino dark matter

$$\Omega_{\nu,0} h^2 = \sum \frac{m_\nu}{94 \text{ eV}} = ??$$

- **Masses** larger than 0.06 eV.

- No reason to fix at the minimum mass.
- Laboratory upper limit $\Sigma m_\nu < 7 \text{ eV}$ from β -decay endpoint.

- **More than three flavours.** $N_{\text{eff}} \neq 3 ??$

- **Sterile neutrinos** and discrepancies potentially solved by them?

- **Hidden interactions**

- Neutrino-neutrino, neutrino-dark matter, neutrino-dark energy.

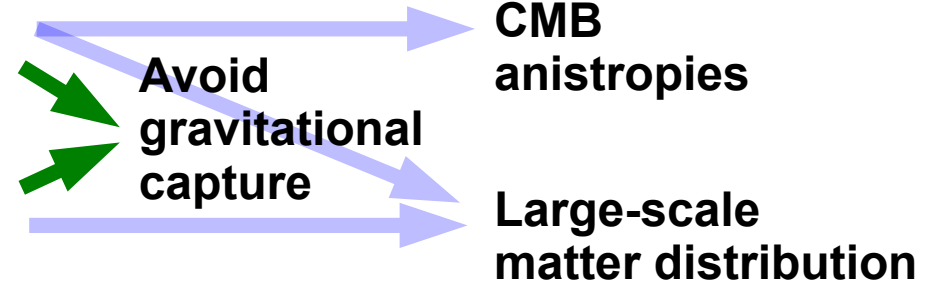
This talk

Measuring neutrino masses with
cosmology...

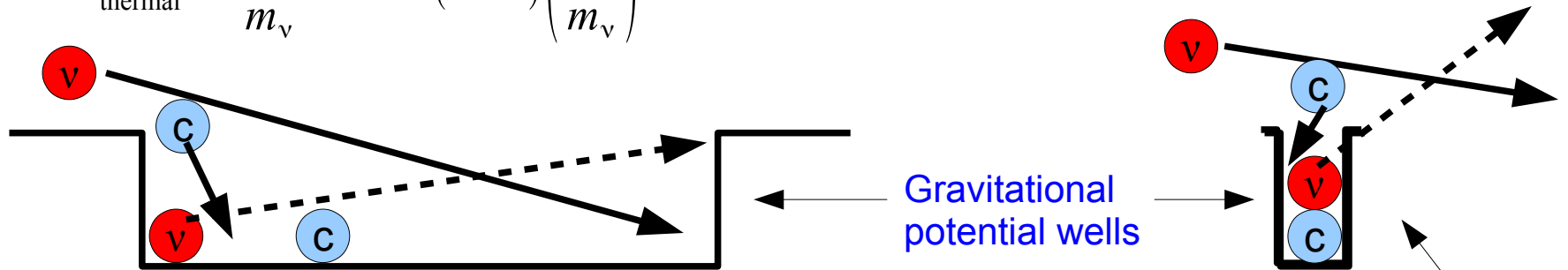
Free-streaming neutrinos...

For most of the observable history of the universe **neutrinos have significant speeds.**

- eV-mass neutrinos **become nonrelativistic** near γ decoupling.
- Even when nonrelativistic, neutrinos have large **thermal motion.**



$$v_{\text{thermal}} = \frac{T_\nu}{m_\nu} \simeq 50.4(1+z) \left(\frac{\text{eV}}{m_\nu} \right) \text{ km s}^{-1}$$



Free-streaming scale:

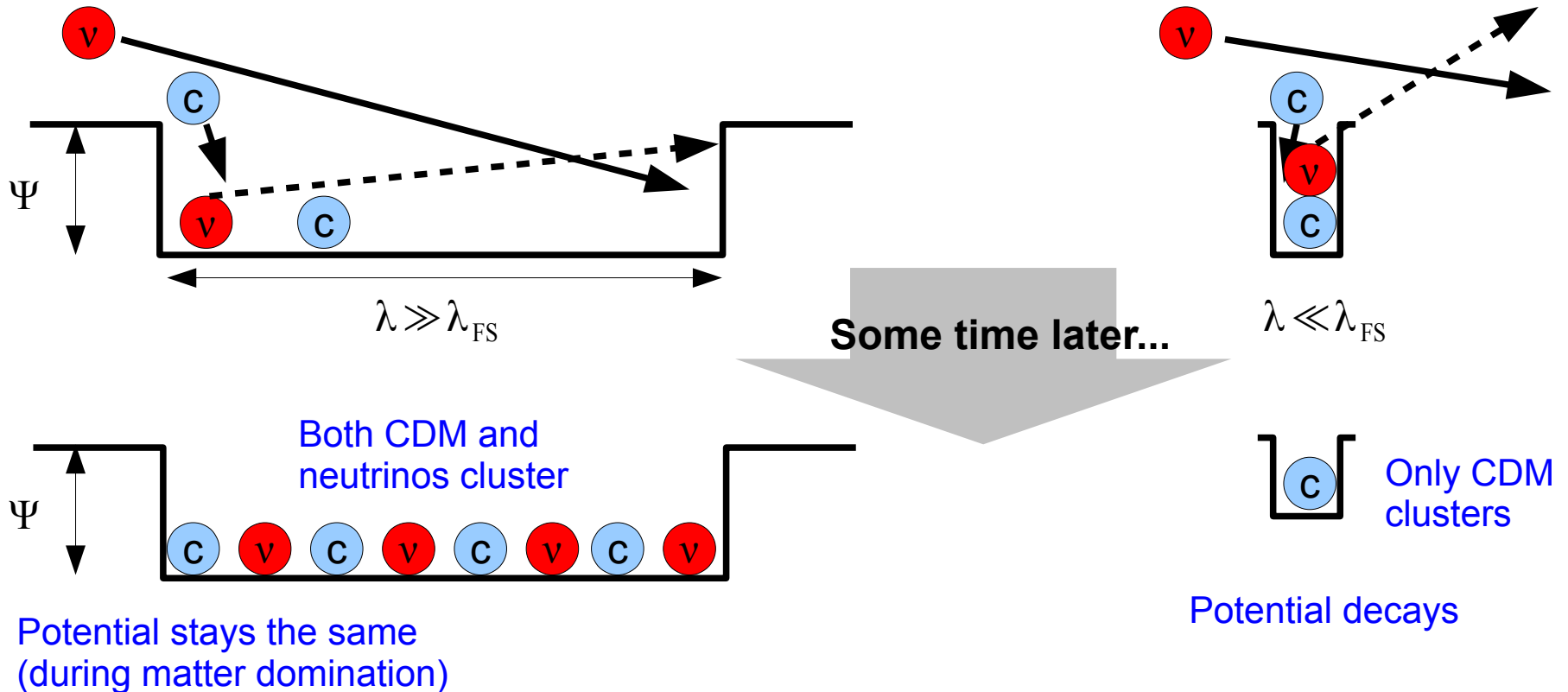
$$\lambda_{\text{FS}} \equiv \sqrt{\frac{8 \pi^2 v_{\text{thermal}}^2}{3 \Omega_m H^2}} \simeq 4.2 \sqrt{\frac{1+z}{\Omega_{m,0}}} \left(\frac{\text{eV}}{m_\nu} \right) h^{-1} \text{ Mpc}; \quad k_{\text{FS}} \equiv \frac{2 \pi}{\lambda_{\text{FS}}}$$

Non-clustering

$$\lambda \ll \lambda_{\text{FS}}$$

$$k \gg k_{\text{FS}}$$

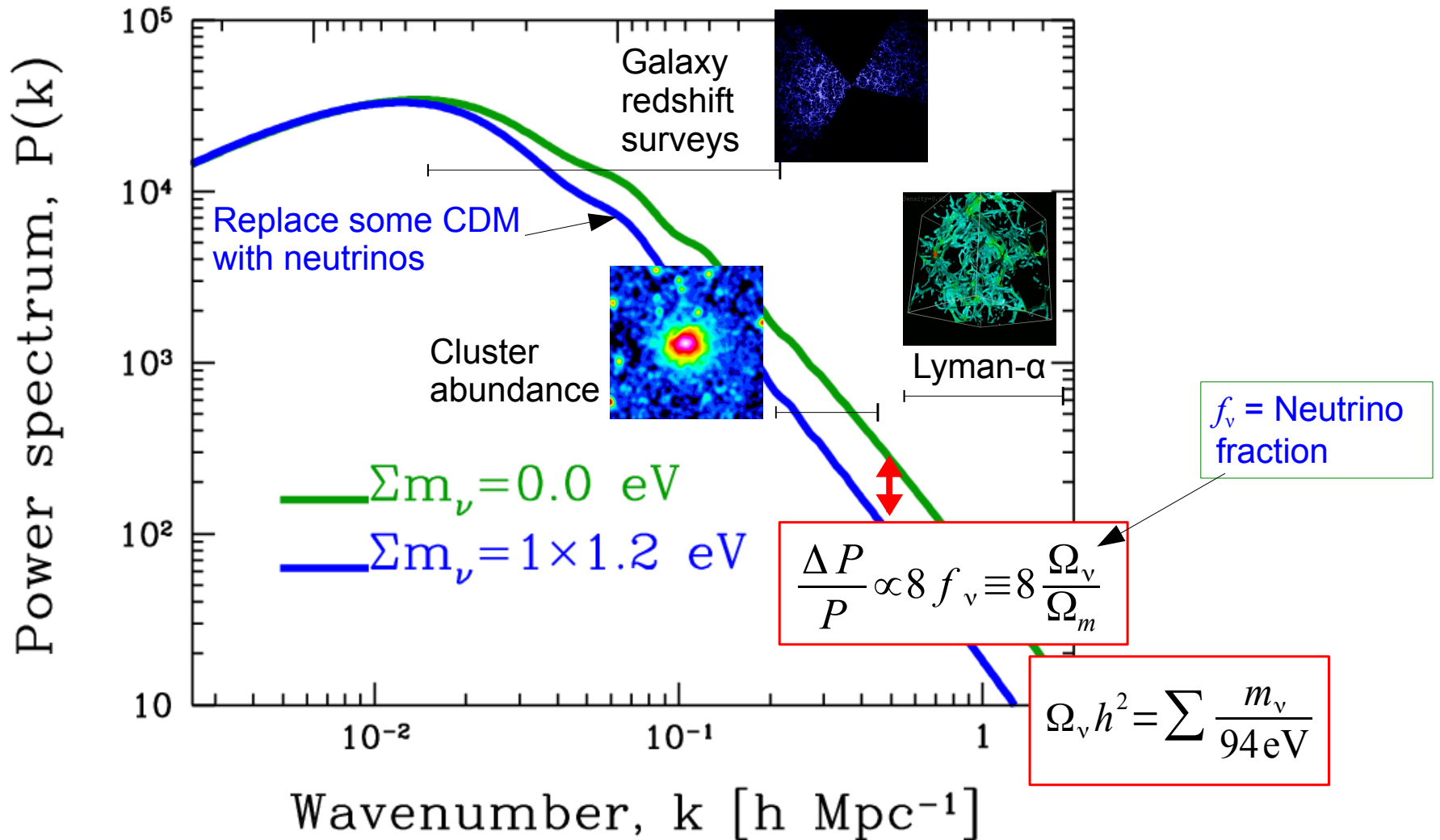
Consider a **neutrino** and a **cold dark matter particle** encountering two gravitational potential wells of different sizes in an expanding universe:



→ **Cosmological neutrino mass measurement** is based on observing this **free-streaming induced potential decay** at $\lambda \ll \lambda_{\text{FS}}$.

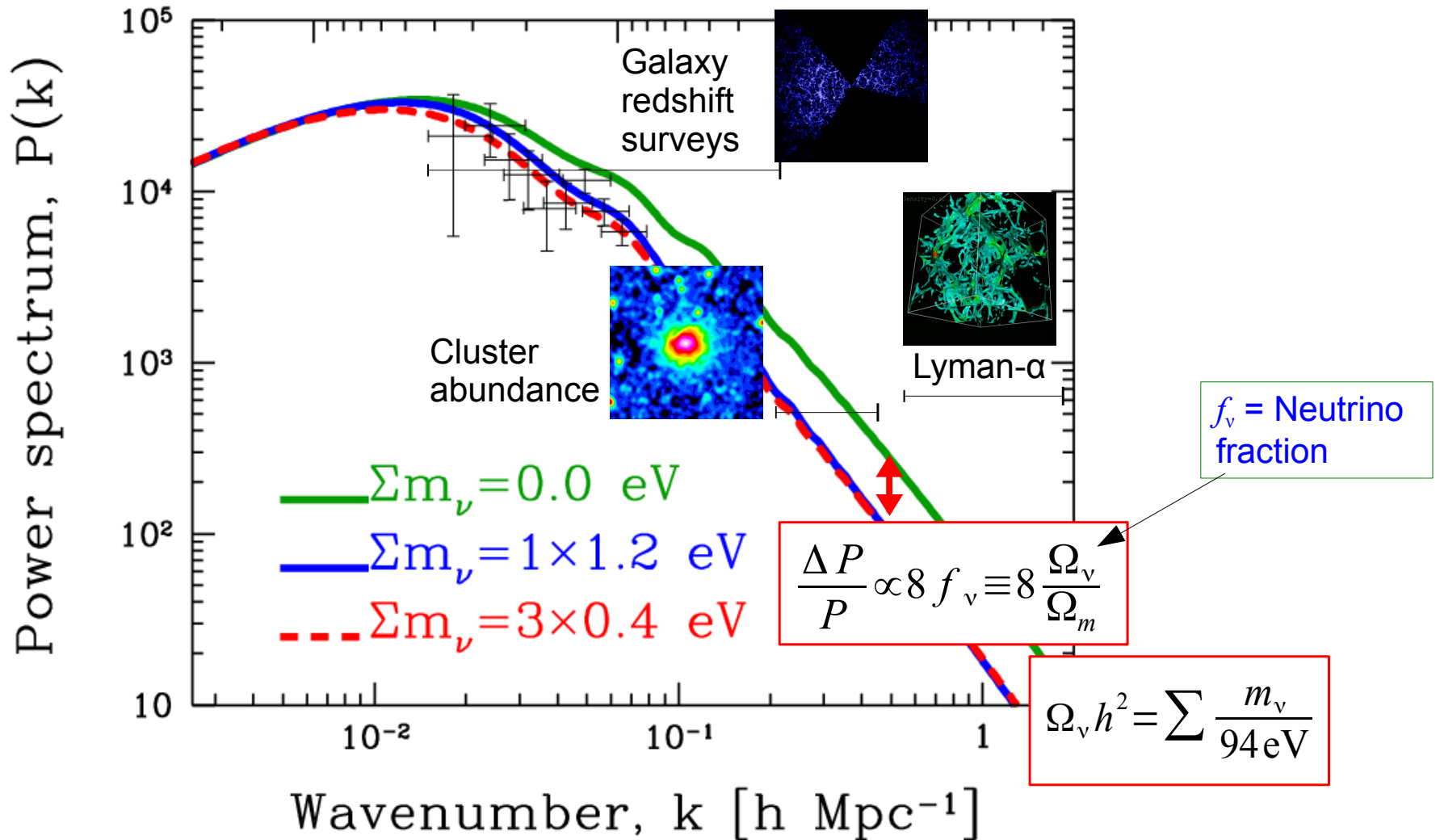
Large-scale matter distribution...

$$P(k) = \langle |\delta(k)|^2 \rangle$$

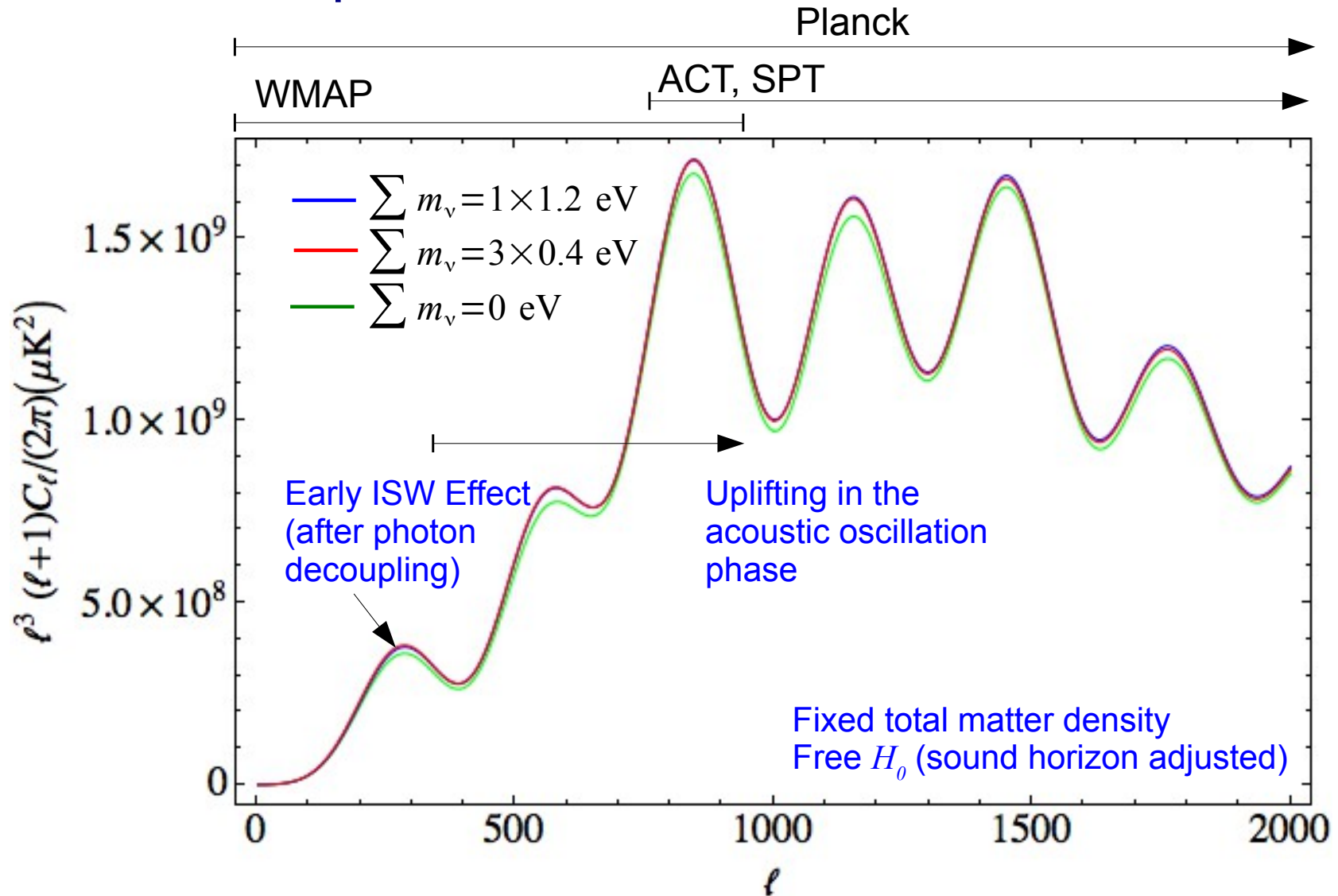


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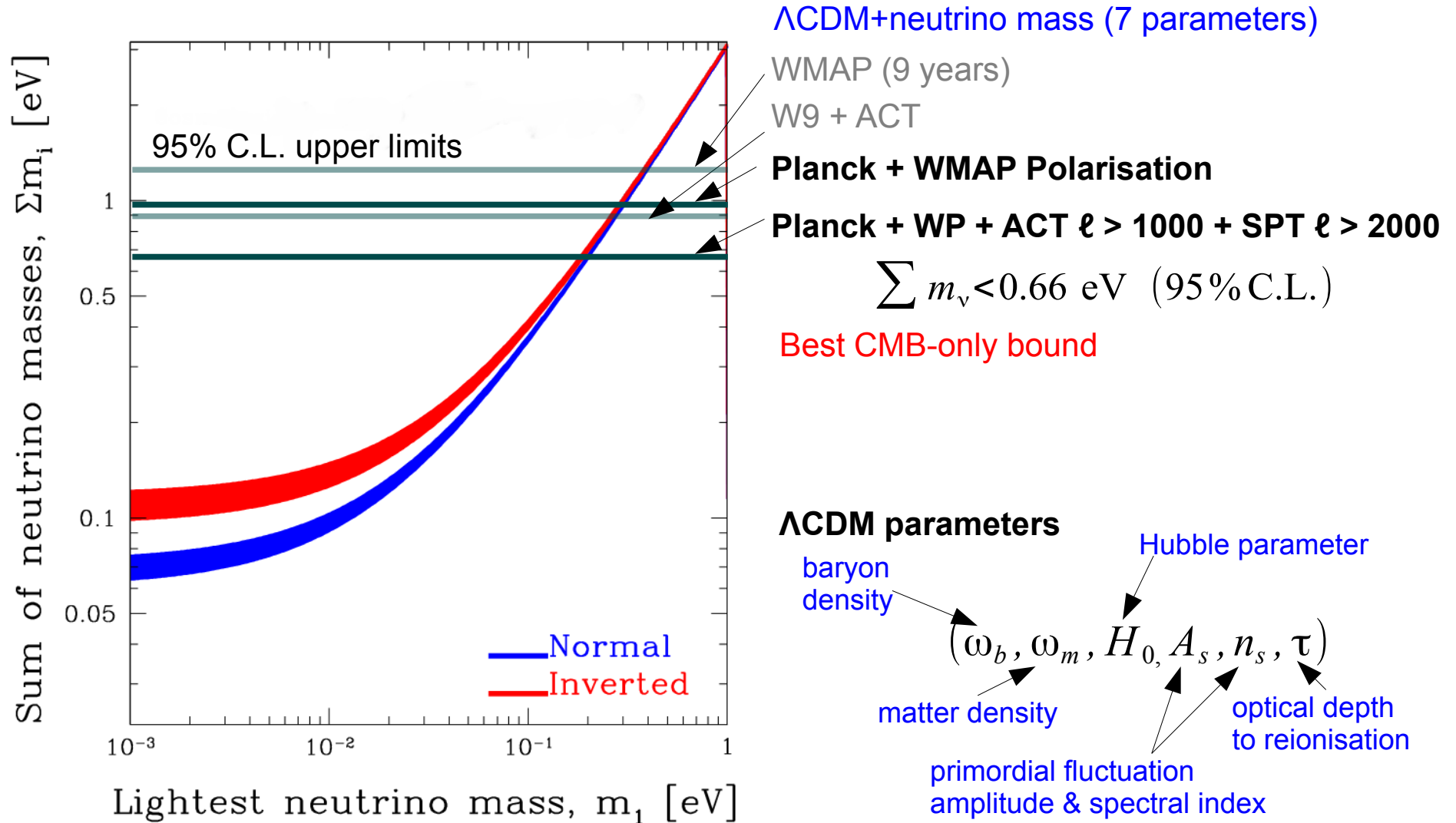


CMB anisotropies...



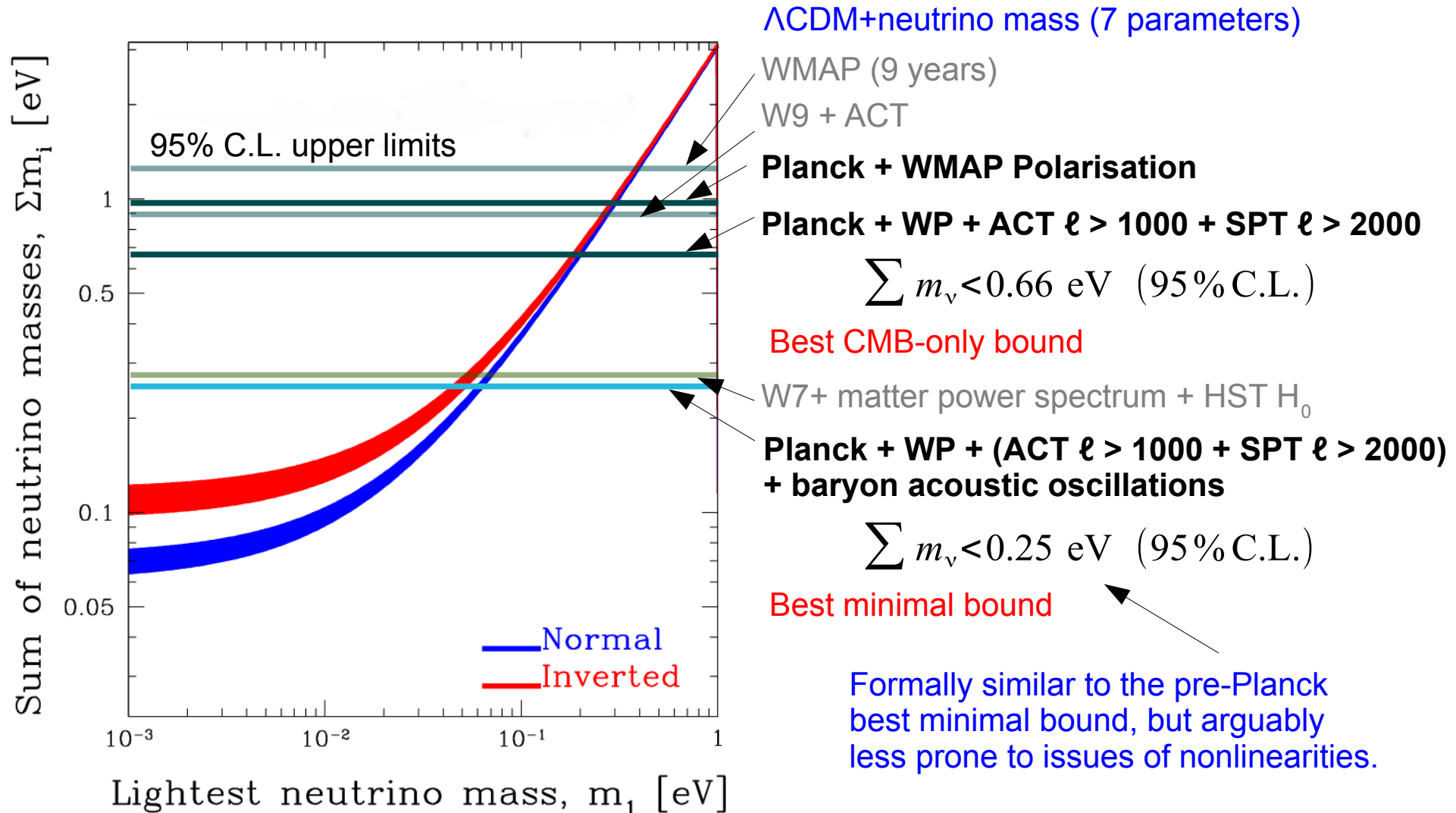
Post-Planck constraints...

Ade et al.[Planck] 2013

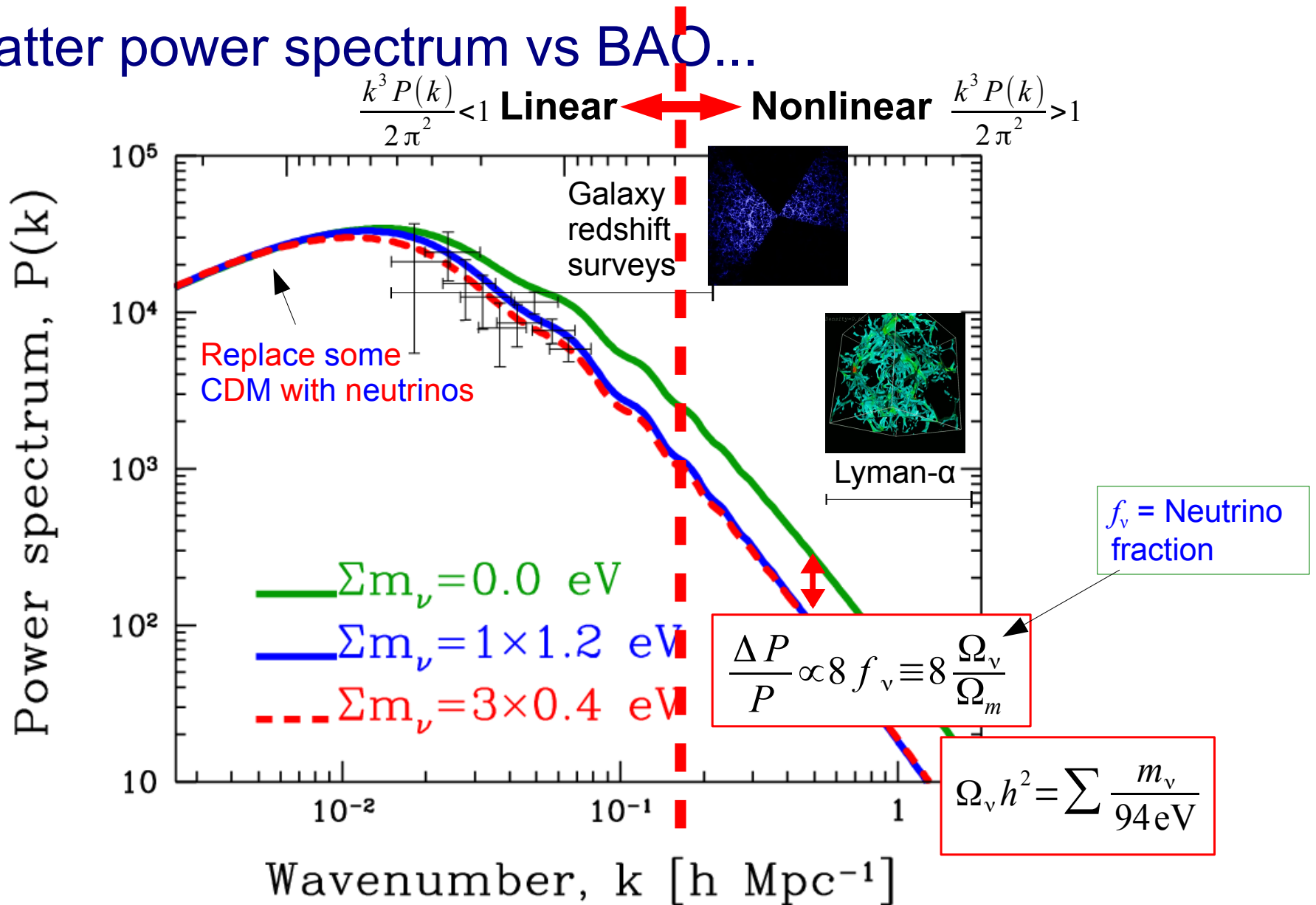


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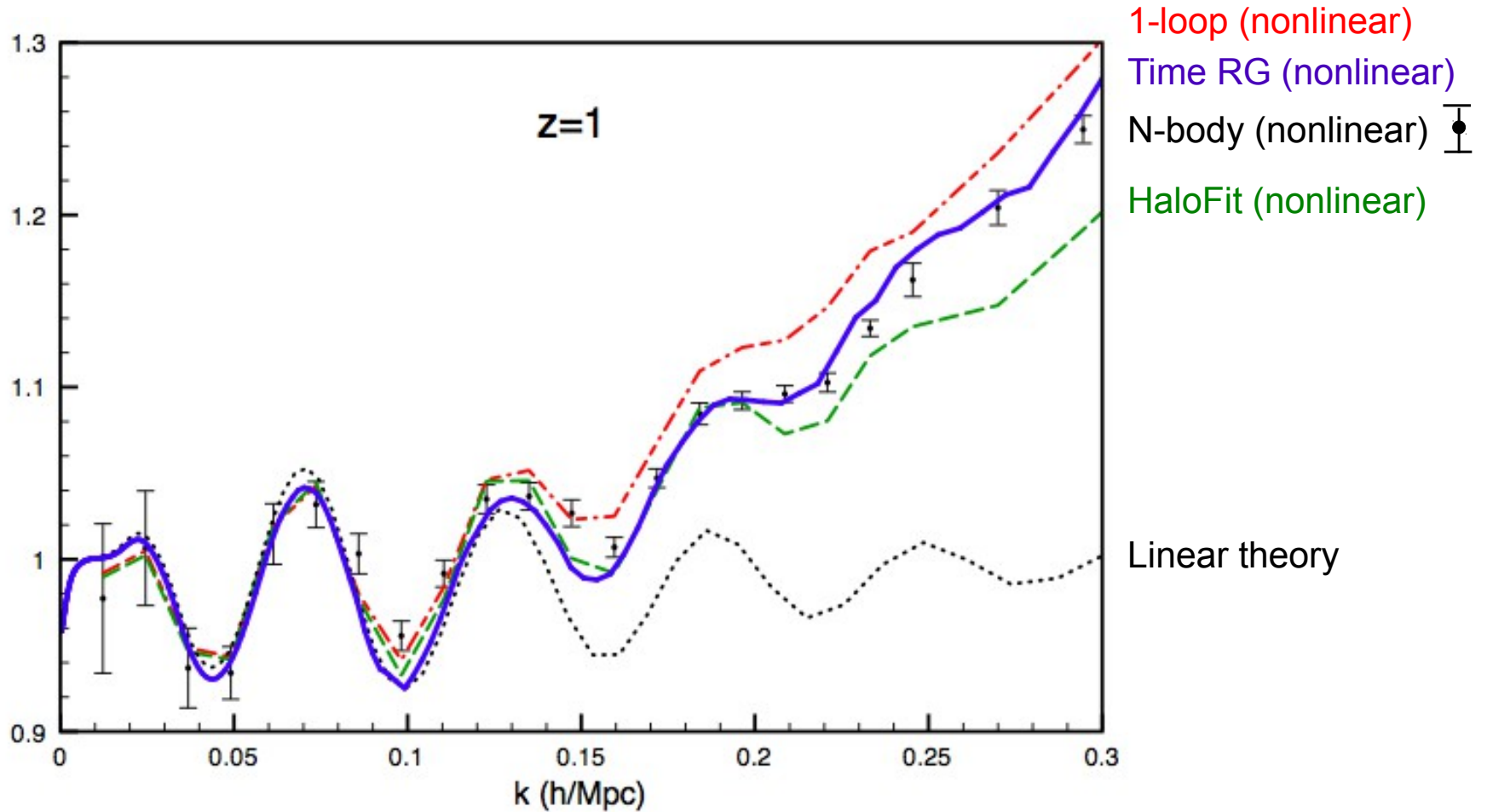


Matter power spectrum vs BAO...



Matter power spectrum = Shape
Baryon acoustic oscillations = Location of oscillatory features

Matter power spectrum (normalised to smooth spectrum)



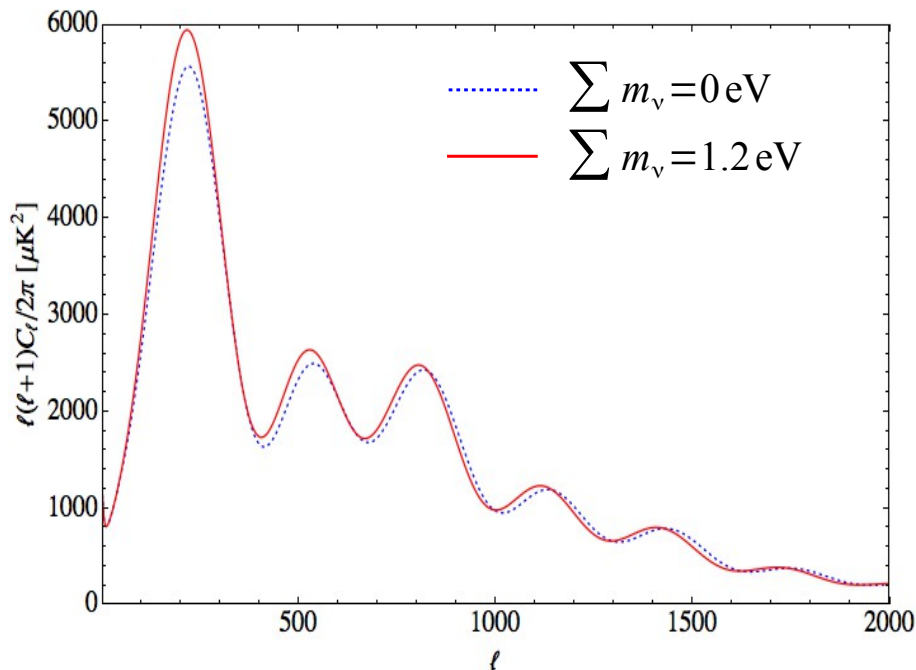
Pietroni 2008

The take-home message...

- Formally, the best minimal (7-parameter) upper bound on Σm_ν is **still hovering around 0.3 eV** post-Planck.
- The bound has however become **more robust against uncertainties**:
 - Less nonlinearities in BAO than in the matter power spectrum.
 - Does not rely on local measurement of the Hubble parameter...
 - ... or on the choice of lightcurve fitters for the Supernova Ia data.
- **Dependence on cosmological model** used for inference?

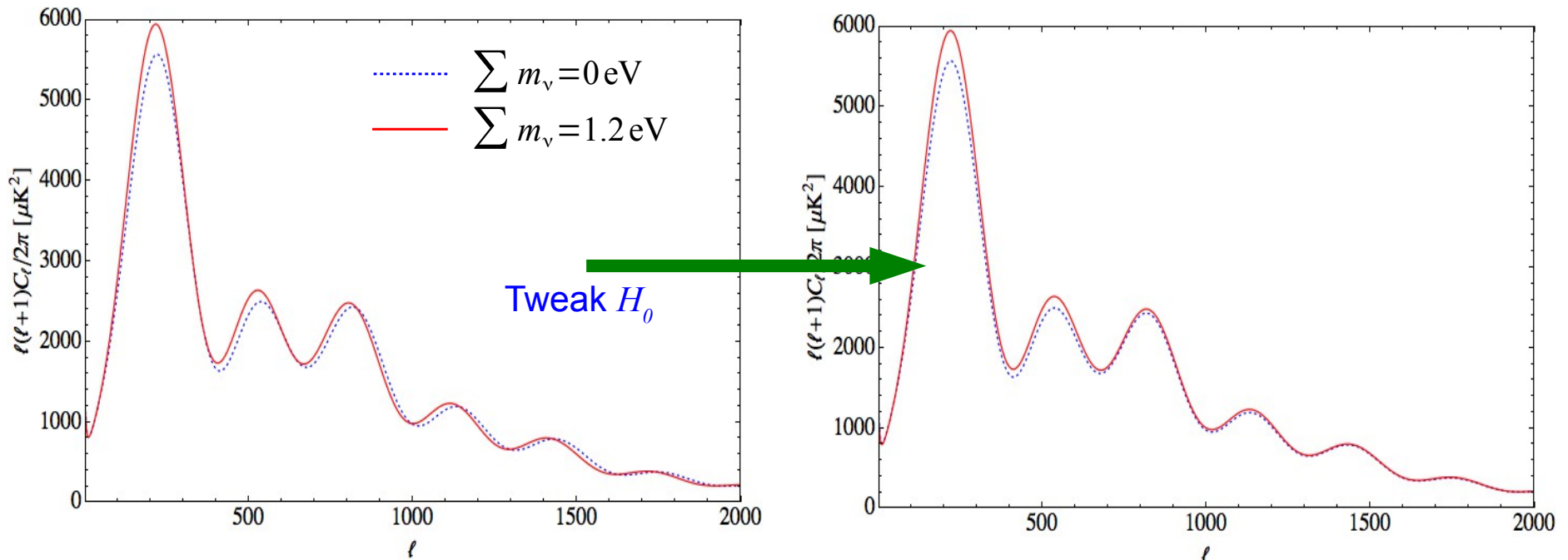
Model dependence: parameter degeneracies...

- We **do not** measure the neutrino mass *per se*, but rather its **indirect effect** on the clustering statistics of the CMB/large-scale structure.
 - It is **not impossible** that **other cosmological parameters** could give rise to **similar effects** (within measurement errors/cosmic variance).



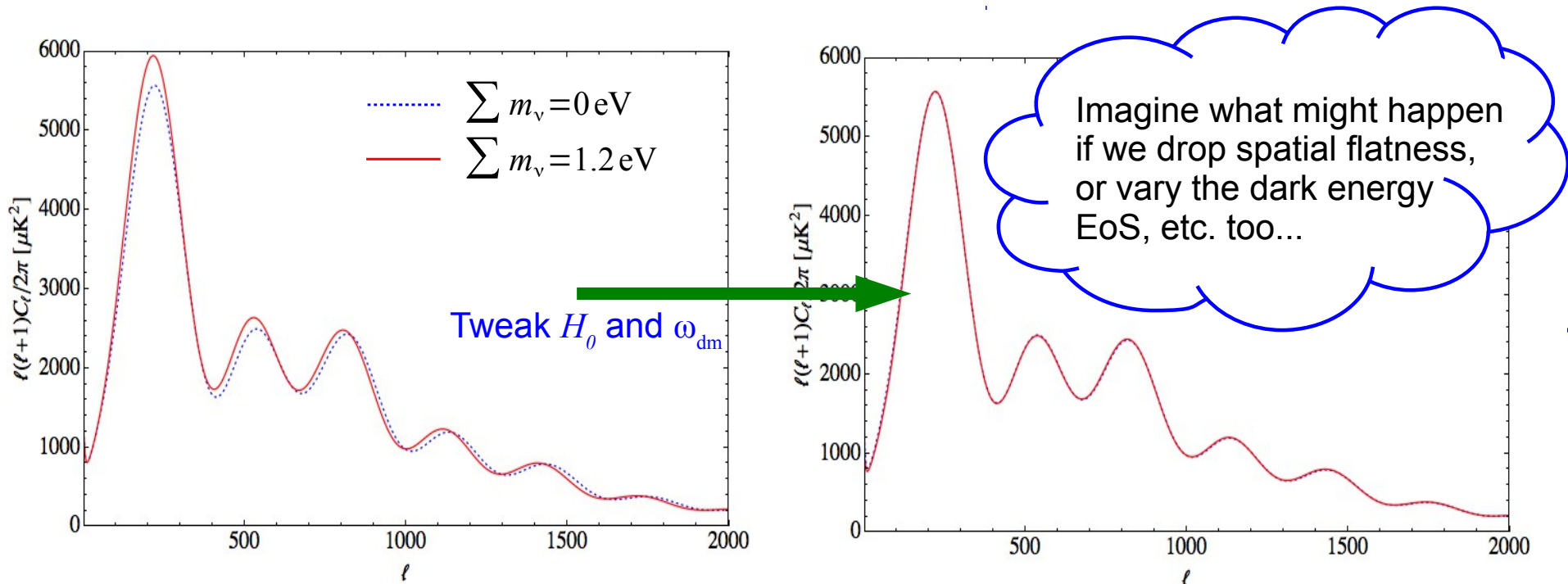
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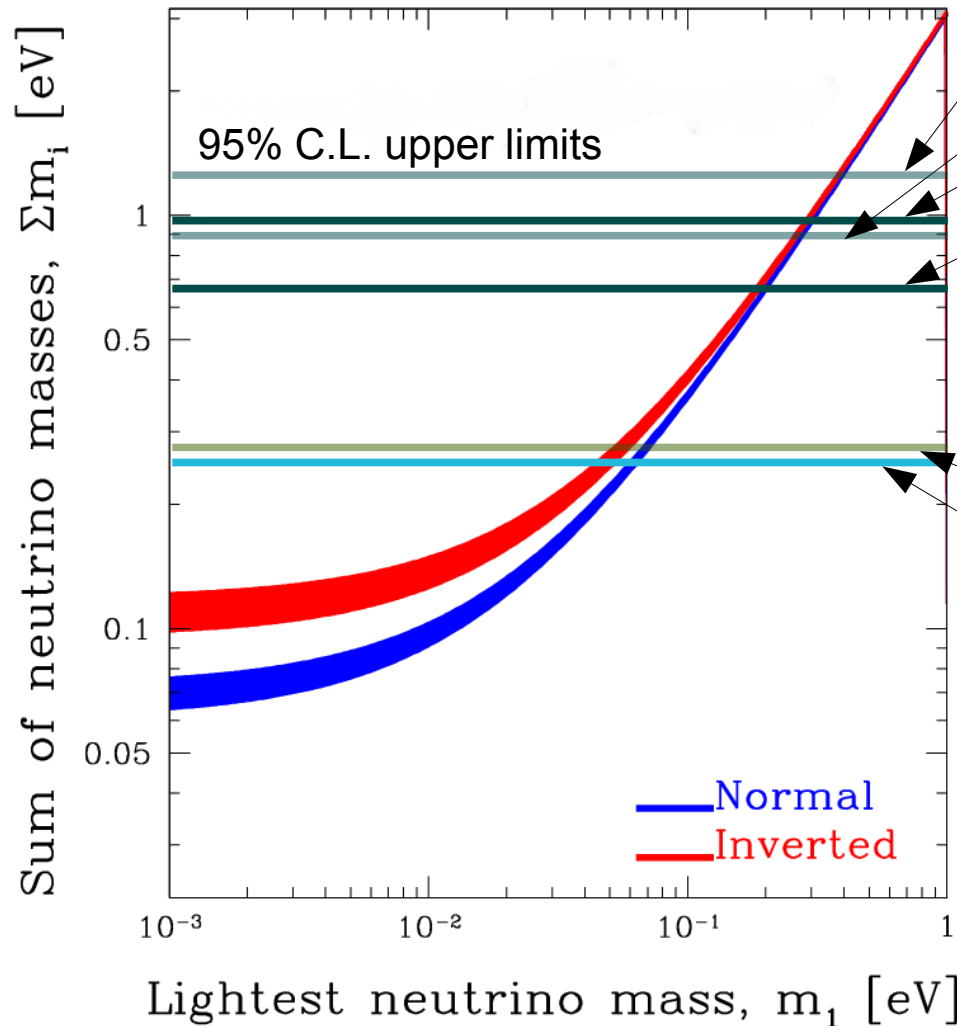
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Post-Planck...

Ade et al.[Planck] 2013



Λ CDM+neutrino mass (7 parameters)

WMAP (9 years)

W9 + ACT

Planck + WMAP Polarisation

Planck + WP + ACT $\ell > 1000$ + SPT $\ell > 2000$

$$\sum m_\nu < 0.66 \text{ eV (95\% C.L.)}$$

Best CMB-only bound

W7+ matter power spectrum + HST H_0

Planck + WP + (ACT $\ell > 1000$ + SPT $\ell > 2000$) + baryon acoustic oscillations

$$\sum m_\nu < 0.25 \text{ eV (95\% C.L.)}$$

Best minimal bound

Dropping assumption of spatial flatness:

$$\sum m_\nu < 0.32 \text{ eV (95\% C.L.)}$$

Other extensions?? Still to be checked!

A fourth neutrino??

It doesn't even have to be a real neutrino...

Any particle species that

- decouples **while ultra-relativistic** and **before $z \sim 10^6$**
- does **not** interact with itself after decoupling

will behave (more or less) like a neutrino as far as the CMB and LSS are concerned.

Smallest relevant
scale enters the horizon

$$\begin{aligned} \sum_i \rho_{\nu,i} + \rho_X &= N_{\text{eff}} \left(\frac{7}{8} \frac{\pi^2}{15} T_{\nu}^4 \right) \\ &= (3.046 + \Delta N_{\text{eff}}) \rho_{\nu}^{(0)} \end{aligned}$$

Three SM neutrinos

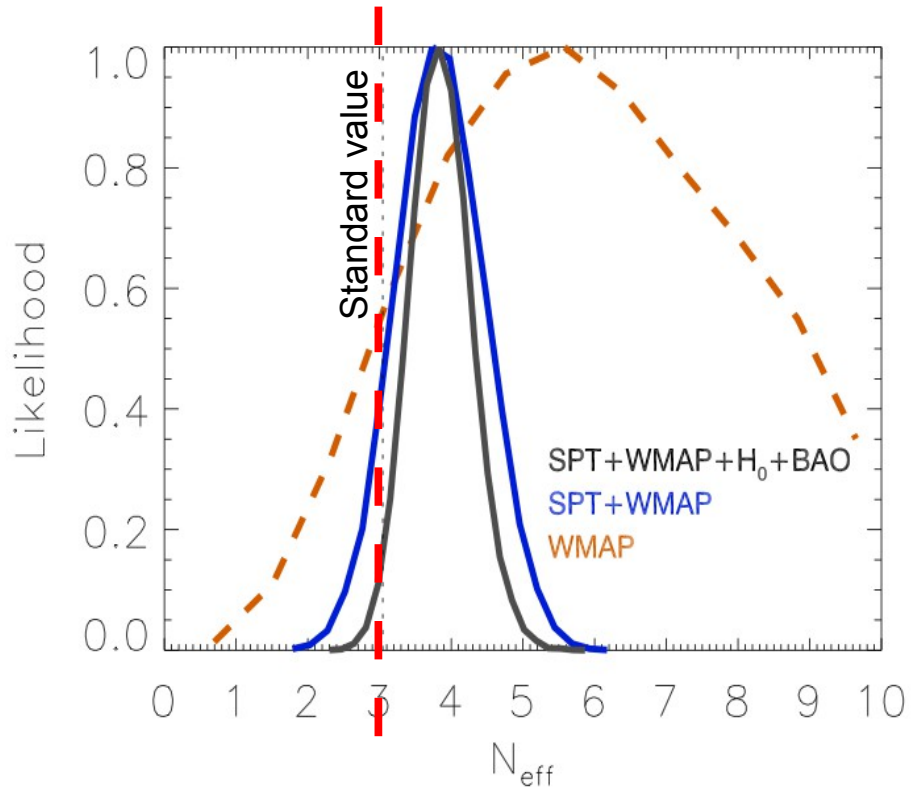
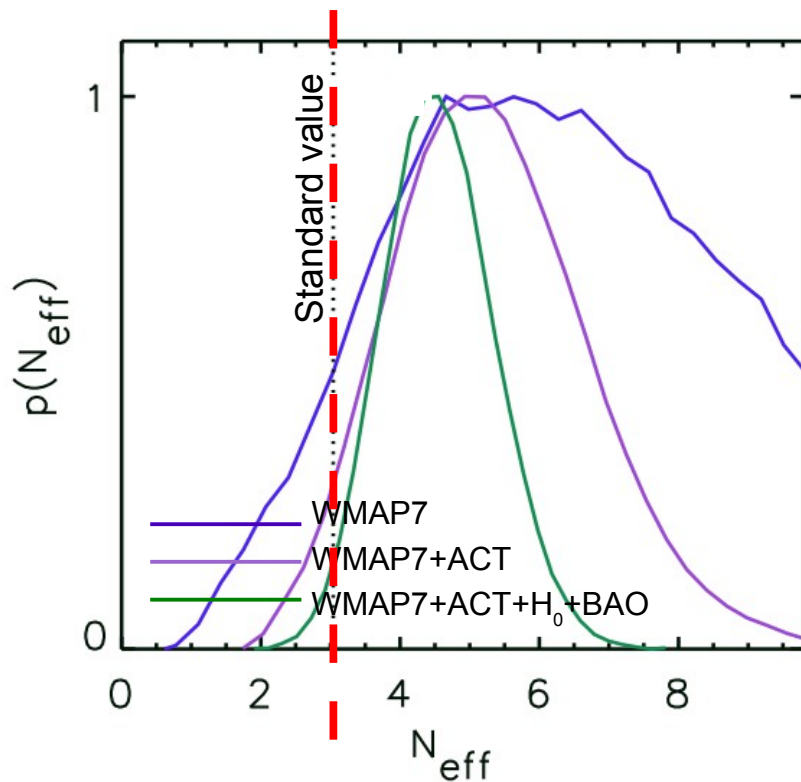
Other non-interacting relativistic energy densities, e.g., sterile neutrinos, axions, hidden photons, etc.

Neutrino temperature per definition

Corrections due to non-instantaneous decoupling, finite temperature effects, and flavour oscillations

Evidence for $N_{\text{eff}} > 3$ circa 2011...

Some pre-Planck observations preferred an **excess of non-interacting relativistic energy density** \rightarrow “**extra neutrinos**”.



Dunkley et al. [Atacama Cosmology Telescope] 2010

Keisler et al. [South Pole Telescope] 2011

Then the evidence disappeared again... largely...

New data from WMAP, ACT and SPT in late 2012 – early 2013 favour an N_{eff} value compatible with the standard value of 3.046.

WMAP 9 years, 1212.5226;
 ACT 3 seasons, 1301.0824
 SPT (2540 deg²), 1212.6267

1 σ error bars

	W9+ACT	W9+ACT	W9+ACT	W9+ACT	W9+ACT	W9+ACT
		+ HST	+BAO	+SNLS3	+BAO+HST	+BAO+SNLS3
N_{eff}	2.74 ± 0.47	3.12 ± 0.38	2.77 ± 0.49	2.79 ± 0.47	3.43 ± 0.36	2.83 ± 0.47
	W9+SPT	W9+SPT	W9+SPT	W9+SPT	W9+SPT	W9+SPT
		+ HST	+BAO	+SNLS3	+BAO+HST	+BAO+SNLS3
N_{eff}	3.93 ± 0.68	3.59 ± 0.39	3.50 ± 0.59	3.96 ± 0.69	3.83 ± 0.41	3.55 ± 0.63

Archidiacono, Giusarma, Melchiorri & Mena, 1303.0143

$N_{\text{eff}} > 3$ at 2 σ +

Post-Planck N_{eff} ...

Planck-inferred N_{eff} **compatible with 3.046** at better than 2σ .

2 σ error bars	<i>Planck</i> +WP		<i>Planck</i> +WP+BAO		<i>Planck</i> +WP+highL		<i>Planck</i> +WP+highL+BAO	
	Best fit	95% limits	Best fit	95% limits	Best fit	95% limits	Best fit	95% limits
Ω_K	-0.0105	-0.037 ^{+0.043} _{-0.049}	0.0000	0.0000 ^{+0.0066} _{-0.0067}	-0.0111	-0.042 ^{+0.043} _{-0.048}	0.0009	-0.0005 ^{+0.0065} _{-0.0066}
Σm_ν [eV]	0.022	< 0.933	0.002	< 0.247	0.023	< 0.663	0.000	< 0.230
N_{eff}	3.08	3.51 ^{+0.80} _{-0.74}	3.08	3.40 ^{+0.59} _{-0.57}	3.23	3.36 ^{+0.68} _{-0.64}	3.22	3.30 ^{+0.54} _{-0.51}
Y_P	0.2583	0.283 ^{+0.045} _{-0.048}	0.2736	0.283 ^{+0.043} _{-0.045}	0.2612	0.266 ^{+0.040} _{-0.042}	0.2615	0.267 ^{+0.038} _{-0.040}
$dn_s/d \ln k$	-0.0090	-0.013 ^{+0.018} _{-0.018}	-0.0102	-0.013 ^{+0.018} _{-0.018}	-0.0106	-0.015 ^{+0.017} _{-0.017}	-0.0103	-0.014 ^{+0.016} _{-0.017}
$r_{0.002}$	0.000	< 0.120	0.000	< 0.122	0.000	< 0.108	0.000	< 0.111
w	-1.20	-1.49 ^{+0.65} _{-0.57}	-1.076	-1.13 ^{+0.24} _{-0.25}	-1.20	-1.51 ^{+0.62} _{-0.53}	-1.109	-1.13 ^{+0.23} _{-0.25}

Very possibly the end of the N_{eff} story...

Or maybe not...
More later...

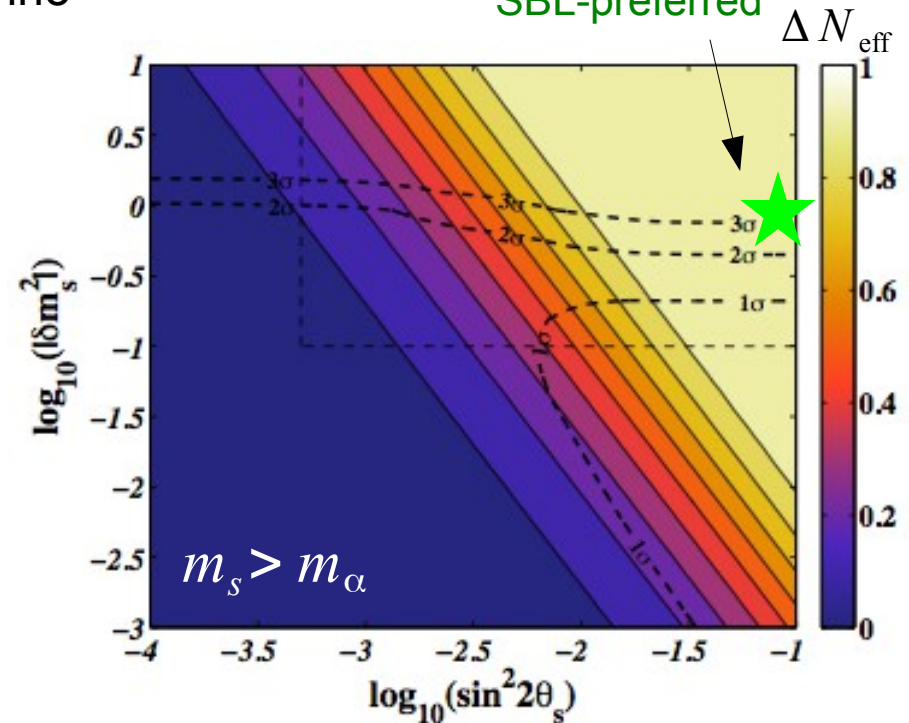
Implications for the short baseline sterile neutrino...

The **LSND/MiniBooNE/Reactor anomalies** can be explained by oscillations into a sterile neutrino with oscillation parameters:

$$\begin{aligned}
 & \Delta m_{\text{SBL}}^2 \sim 1 \text{ eV}^2 \\
 (\nu_e \leftrightarrow \nu_\mu) \quad & \sin^2 2\theta_{\text{SBL}} \sim 3 \times 10^{-3}
 \end{aligned}
 \left. \vphantom{\begin{aligned} \Delta m_{\text{SBL}}^2 \sim 1 \text{ eV}^2 \\ \sin^2 2\theta_{\text{SBL}} \sim 3 \times 10^{-3} \end{aligned}} \right\} \Delta N_{\text{eff}} = 1$$

Fully thermalised sterile neutrino population!

SBL-preferred



Hannestad, Tamborra & Tram 2012
 also older works of Abazajian, Di Bari,
 Foot, Kainulainen, etc. from 1990s-early 2000s

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$(\nu_e \leftrightarrow \nu_\mu)$

$m_{\text{sterile}} > \sqrt{\Delta m_{\text{SBL}}^2} \sim 1 \text{ eV}$

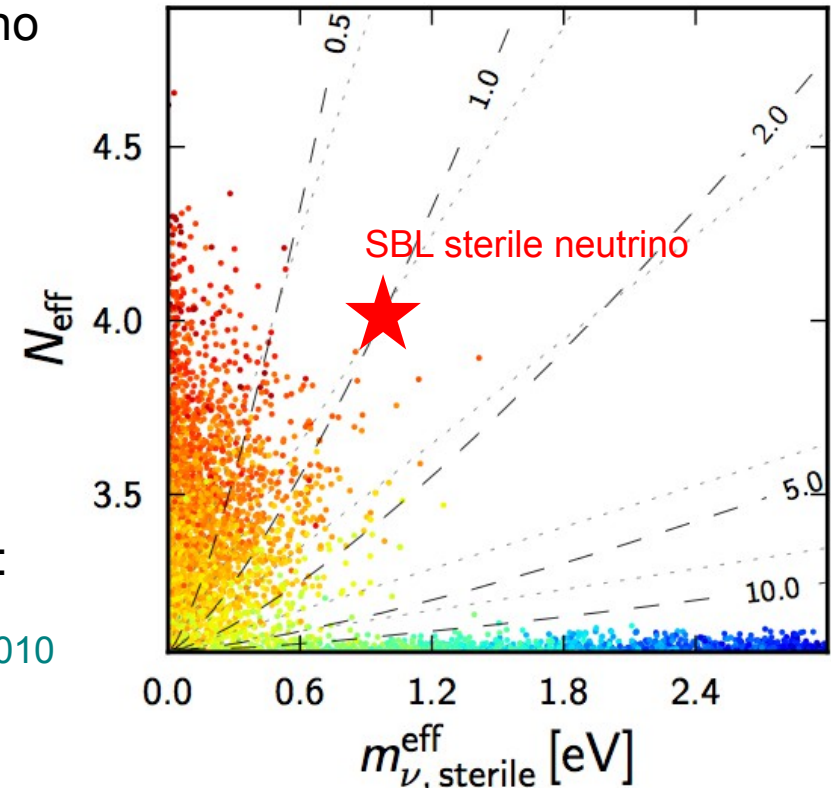
- **Already a problem** for WMAP (+LSS+HST):

$$m_{\text{sterile}} < 0.45 \text{ eV} \quad (95\% \text{ C.L.}) \quad \text{Hamann et al. 2010}$$

- **Post-Planck:**

$m_{\text{sterile}} < 0.42 \text{ eV}; \quad N_{\text{eff}} < 3.8 \quad (95\% \text{ C.L.})$

Ade et al. [Planck collaboration] 2013



Planck+WP+highL+BAO
(Λ CDM+ N_{eff} + m_{sterile})

Reconciling the SBL sterile neutrinos with cosmology??

The SBL sterile neutrino is problematic for cosmology **only because it is produced** in abundance in the early universe.

→ If production can be **suppressed**, then there is no conflict.

- **Some possible mechanisms:**

- A **large lepton asymmetry** ($L \gg B \sim 10^{-10}$) generates an effective mass for the active neutrino to suppress effective active-sterile mixing; $L \sim 10^{-2}$ will do.

Foot & Volkas 1995

- **Hidden sterile neutrino self-interaction** generates an effective mass for the sterile neutrino.

Dasgupta & Kopp 2014
Hannestad, Hansen & Tram 2014

- A **low reheating temperature** ($T_R < 10$ MeV) → incomplete thermalisation of even the SM neutrinos.

Discrepancies potentially resolved by
a fourth neutrino??

Planck discrepancies with other observations... 1.

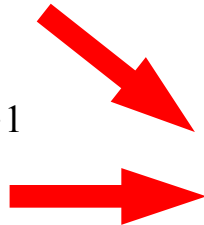
- **Hubble parameter H_0** : Planck-inferred value lower than **local HST measurement**.
- **Small-scale RMS fluctuation σ_8** : Planck CMB prefers a higher value than galaxy cluster count and galaxy shear from CFHTLenS.

Parameter	Planck		Planck+lensing		Planck+WP	
	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits
σ_8	0.8344	0.834 ± 0.027	0.8285	0.823 ± 0.018	0.8347	0.829 ± 0.012
z_{re}	11.35	$11.4^{+4.0}_{-2.8}$	11.45	$10.8^{+3.1}_{-2.5}$	11.37	11.1 ± 1.1
H_0	67.11	67.4 ± 1.4	68.14	67.9 ± 1.5	67.04	67.3 ± 1.2

Hubble space telescope

$$H_0 = 73.8 \pm 2.4 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Riess et al. 2011



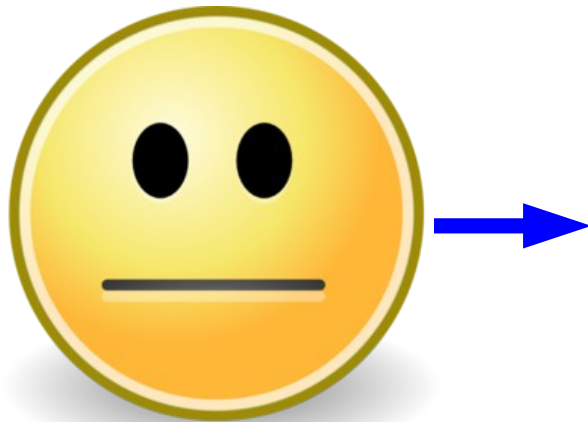
Exploit the $N_{\text{eff}} - H_0$ degeneracy and introduce to a large N_{eff} to bring HST and Planck in line with one another.

Planck + HST

$$N_{\text{eff}} = 3.62 \pm 0.25 (1 \sigma)$$

(Λ CDM + ΔN_{eff}
7-parameter model)

... Not quite the SBL sterile neutrino because this fit assumes massless neutrinos...



The impact of additional astrophysical data is particularly complex in our investigation of neutrino physics (Sect. 6.3). We will use the effective number of relativistic degrees of freedom, N_{eff} as an illustration. From the CMB data alone, we find $N_{\text{eff}} = 3.36 \pm 0.34$. Adding BAO data gives $N_{\text{eff}} = 3.30 \pm 0.27$. Both of these values are consistent with the standard value of 3.046. Adding the H_0 measurement to the CMB data gives $N_{\text{eff}} = 3.62 \pm 0.25$ and *relieves the tension between the CMB data and H_0 at the expense of new neutrino-like physics (at around the 2.3σ level)*. It is possible to alleviate the tensions between the CMB, BAO, H_0 and SNLS data by invoking new physics such as an increase in N_{eff} . However, *none of these cases are favoured significantly over the base Λ CDM model by the Planck data (and they are often disfavoured)*. Any preference for new physics comes almost entirely from the astrophysical data sets. It is up to the reader to decide how to interpret such results, but it is simplistic to assume that all astrophysical data sets have accurately quantified estimates of systematic errors. We have therefore tended to place greater weight on the CMB and BAO measurements in this paper rather than on more complex astrophysical data.

Planck discrepancies with other observations... 2.

- **Hubble parameter H_0** : Planck-inferred value lower than local HST measurement.
- **Small-scale RMS fluctuation σ_8** : Planck CMB prefers a higher value than **galaxy cluster counts** and **galaxy shear from CFHTLens**.

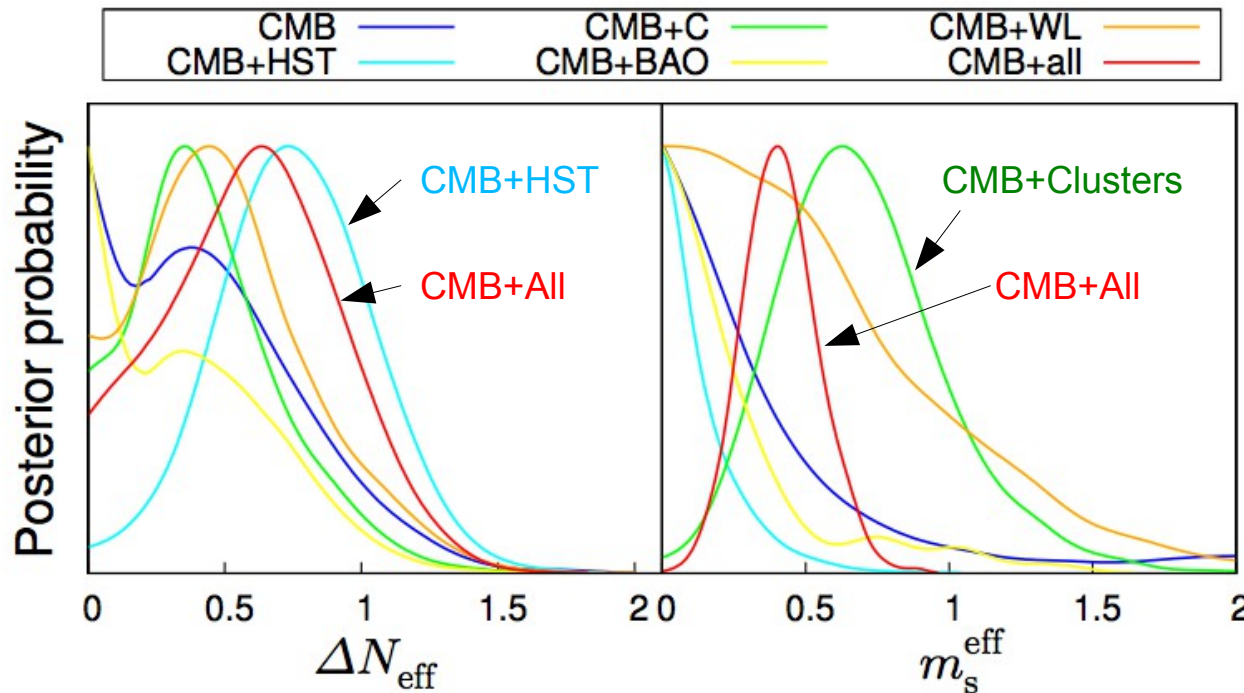
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Planck SZ clusters $\sigma_8 (\Omega_m / 0.27)^{0.3} = 0.782 \pm 0.01$ Ade et al. [Planck collaboration] 2013

CFHTLens galaxy shear $\sigma_8 (\Omega_m / 0.27)^{0.46} = 0.774 \pm 0.04$ Heymans et al. 2013

Solved by a fourth, **massive** neutrino??

At face value a fourth, massive neutrino is a possible solution.



CMB+all
(Λ CDM+ $\Delta N_{\text{eff}}+m_s$
8-parameter model)

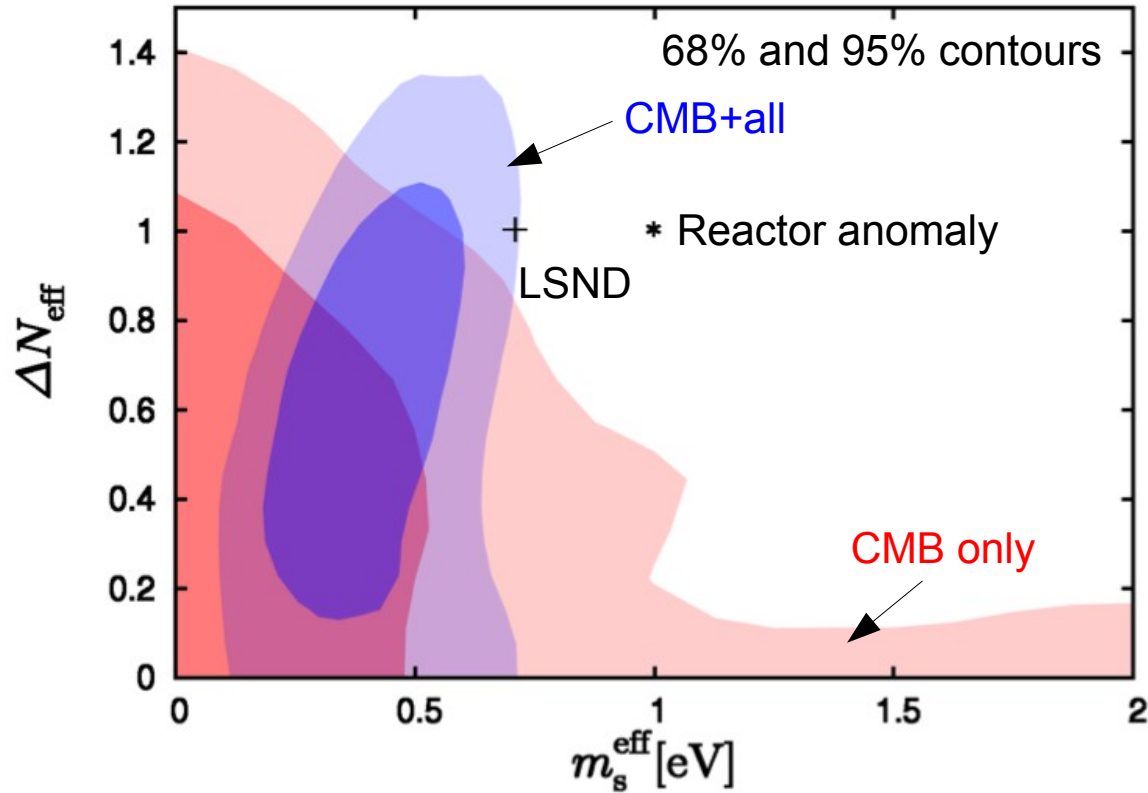
$$\Delta N_{\text{eff}} = 0.61 \pm 0.30$$

$$m_s = (0.41 \pm 0.13) \text{ eV}$$

Hamann & Hasenkamp 2013
also Wyman et al. 2013,
Battye & Moss 2013

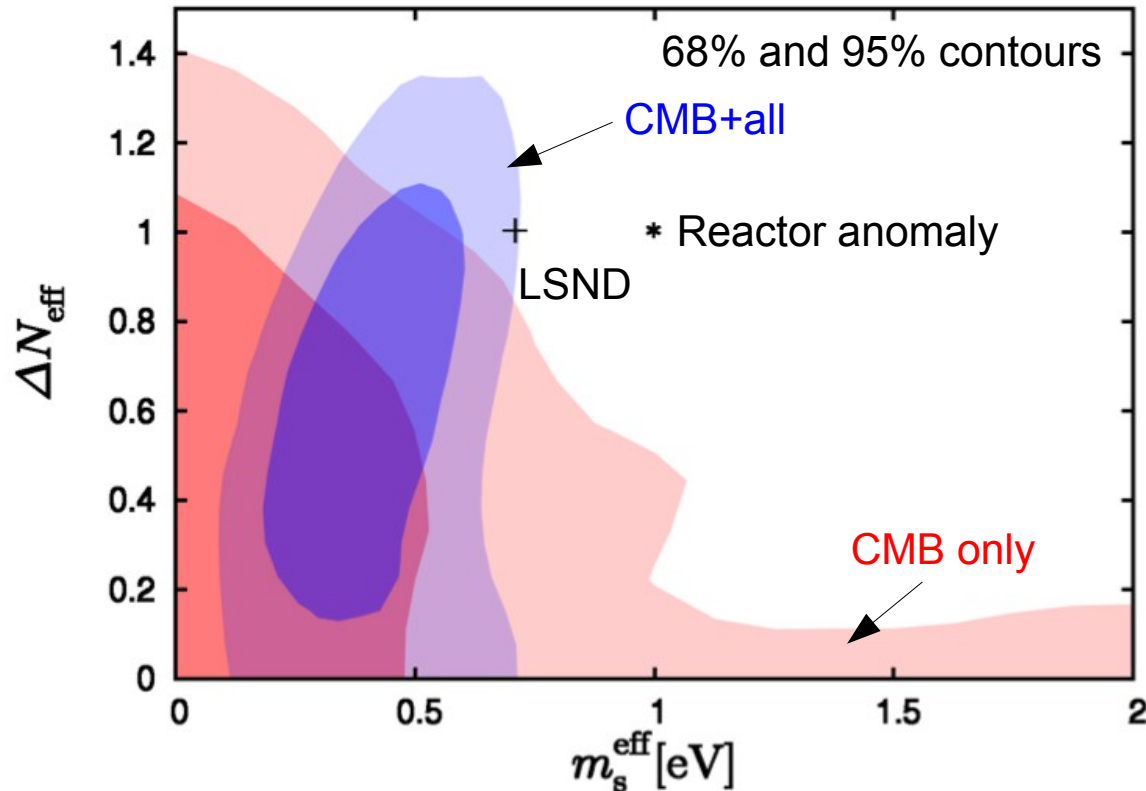
- Large N_{eff} driven mainly by HST
- Large m_s driven mainly by cluster counts.

Solved by a fourth, **massive** neutrino??



Hamann & Hasenkamp 2013
also Wyman et al. 2013
Battye & Moss 2013

Solved by a fourth, **massive** neutrino??



Hamann & Hasenkamp 2013
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My take: discrepancies are most likely due to **poorly understood nonlinearities** (cluster counts are particularly difficult to model).

- Take the fourth neutrino solution *cum magno grano salis!*

Summary...

- Precision cosmological observables can be used to “measure” the absolute neutrino mass scale based on the effect of neutrino free-streaming.
- Existing precision cosmological data already provide strong constraints on the neutrino mass sum.
 - No significant formal improvement between the best pre-Planck and post-Planck upper bounds (at least not for the minimal 7-parameter model).
 - But the post-Planck bound is arguably more robust.
- The fourth neutrino?? There are outstanding discrepancies between Planck and measurements from HST, clusters, and cosmic shear.
 - Taken at face value these discrepancies can be resolved by a fourth neutrino (although not necessarily the same one in all cases...).
 - But personally I'd take it *cum magno grano salis*.