

# $\Upsilon$ production: an example of heavy-ion physics with the extracted 2.76 TeV lead LHC beam



**Jean-Philippe Lansberg**



Plenary meeting of the French GDR PH-QCD  
*SPhN- CEA/Saclay*  
*26 November, 2013*

Use LHC beams on fixed target :

- LHC 7 TeV proton beam

$$\sqrt{s} \sim 115 \text{ GeV: } pp, pd, pA$$

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spin physics

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Cold Nuclear Matter effects

W, Z prod. near threshold

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QGP studies, high precision heavy  
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between SPS and top AA RHIC energies

- benefit from the typical advantages of a fixed target experiment

- ▶ high luminosity, high boost ( $y_{cms}^{lab} = 4.84 @ 115 \text{ GeV}$ ), target versatility

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- benefit from the typical advantages of a fixed target experiment
  - high luminosity, high boost ( $y_{cms}^{lab} = 4.84 @ 115 \text{ GeV}$ ), target versatility
- multipurpose experiment, modern detection techniques

spin physics

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# More details

▶ on the website:  
[after.in2p3.fr](http://after.in2p3.fr)

▶ in Physics Reports :



Contents  
Aller à la page 242

1.	Introduction.....	
2.	Key numbers and features.....	
3.	Nucleon partonic structure.....	
3.1.	Drell-Yan.....	
3.2.	Gluons in the proton at large $x$ .....	
3.2.1.	Quarkonia.....	
3.2.2.	Jets.....	
3.2.3.	Direct/isolated photons.....	
3.3.	Gluons in the deuteron and in the neutron.....	
3.4.	Charm and bottom in the proton.....	
3.4.1.	Open-charm production.....	
3.4.2.	$J/\psi + D$ meson production.....	
3.4.3.	Heavy-quark plus photon production.....	
4.	Spin physics.....	
4.1.	Transverse SSA and DY.....	
4.2.	Quarkonium and heavy-quark transverse SSA.....	
4.3.	Transverse SSA and photon.....	
4.4.	Spin asymmetries with a final state polarization.....	
5.	Nuclear matter.....	
5.1.	Quark nPDF: Drell-Yan in $pA$ and $PbPb$ .....	
5.2.	Gluon nPDF.....	
5.2.1.	Isolated photons and photon-jet correlations.....	
5.2.2.	Precision quarkonium and heavy-flavour studies.....	
5.3.	Color filtering, energy loss, Sudakov suppression and hadron break-up in the nucleus.....	
6.	Deconfinement in heavy-ion collisions.....	
6.1.	Quarkonium studies.....	
6.2.	Jet quenching.....	
6.3.	Direct photon.....	
6.4.	Deconfinement and the target rest frame.....	
6.5.	Nuclear-matter baseline.....	
7.	$W$ and $Z$ boson production in $pp$ , $pd$ and $pA$ collisions.....	250
7.1.	First measurements in $pA$ .....	250
7.2.	$W/Z$ production in $pp$ and $pd$ .....	251
8.	Exclusive, semi-exclusive and backward reactions.....	251
8.1.	Ultra-peripheral collisions.....	251
8.2.	Hard diffractive reactions.....	251
8.3.	Heavy-hadron (diffractive) production at $x_F \rightarrow -1$ .....	251
8.4.	Very backward physics.....	252
8.5.	Direct hadron production.....	252
9.	Further potentialities of a high-energy fixed-target set-up.....	252
9.1.	$D$ and $B$ physics.....	252
9.2.	Secondary beams.....	252
9.3.	Forward studies in relation with cosmic shower.....	253
10.	Conclusions.....	253
	Acknowledgments.....	253
	References.....	253

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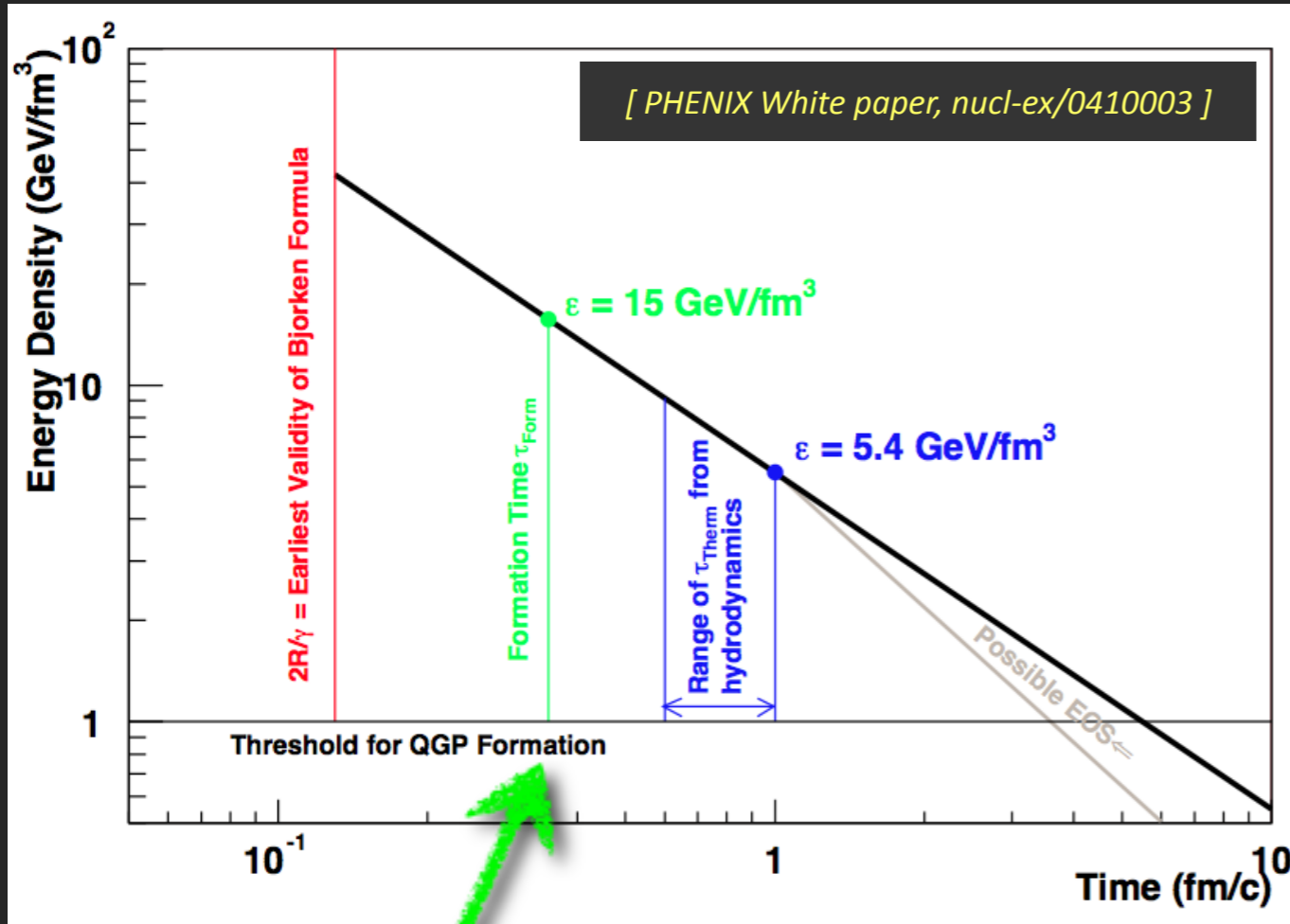
Physics opportunities of a fixed-target experiment using LHC beams

S.J. Brodsky<sup>a</sup>, F. Fleuret<sup>b</sup>, C. Hadjidakis<sup>c</sup>, J.P. Lansberg<sup>c,\*</sup>

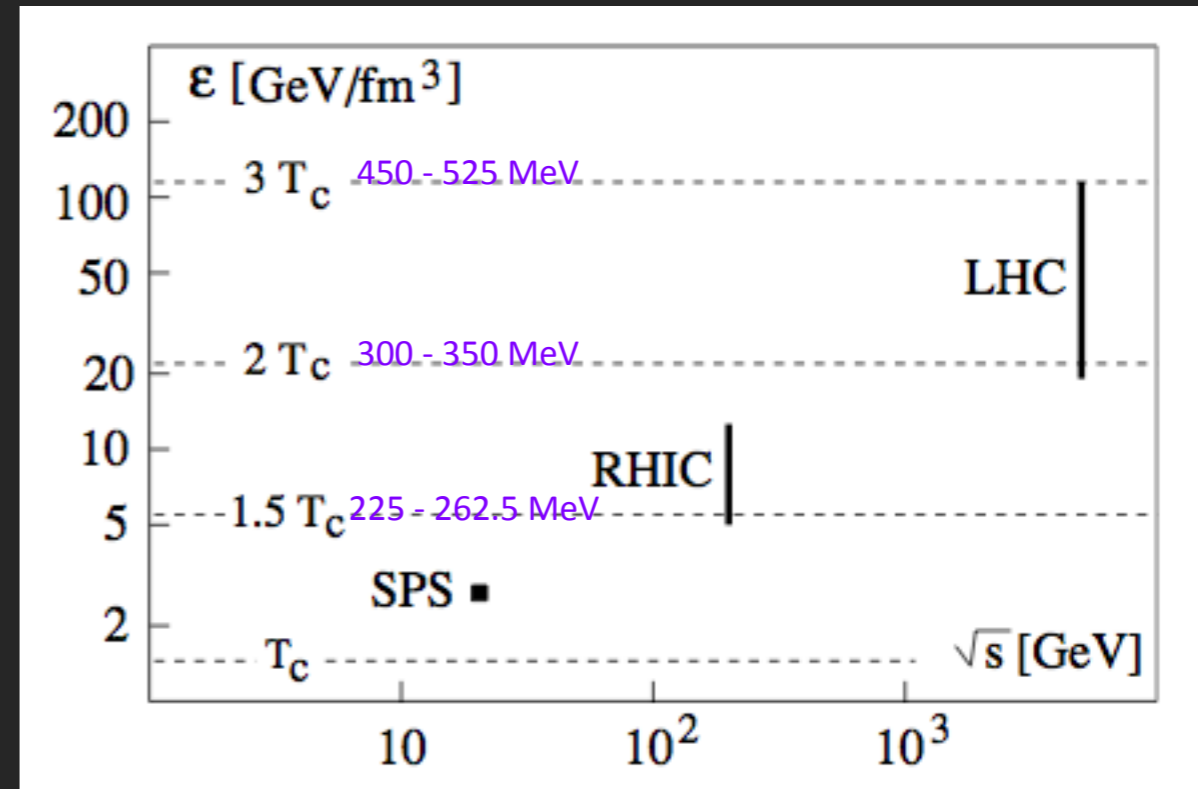
<sup>a</sup> SLAC National Accelerator Laboratory, Stanford University, Menlo Park, CA 94025, USA  
<sup>b</sup> Laboratoire Leprince Ringuet, Ecole polytechnique, CNRS/IN2P3, 91128 Palaiseau, France  
<sup>c</sup> IPNO, Université Paris-Sud, CNRS/IN2P3, 91406 Orsay, France

# Energy density and temperature

Energy density vs time @ RHIC



Energy density, max. collision energy, and temperature



[ Satz, J.Phys. G32 (2006) R25 ]

$T_c \sim 150 - 175 \text{ MeV}$

$T_{\text{initial}}$   
370 - 450 MeV

$T_{\text{avg}} = 221 \pm 19 \text{ (stat)} \pm 19 \text{ (syst)}$   
MeV (0-20% AuAu)

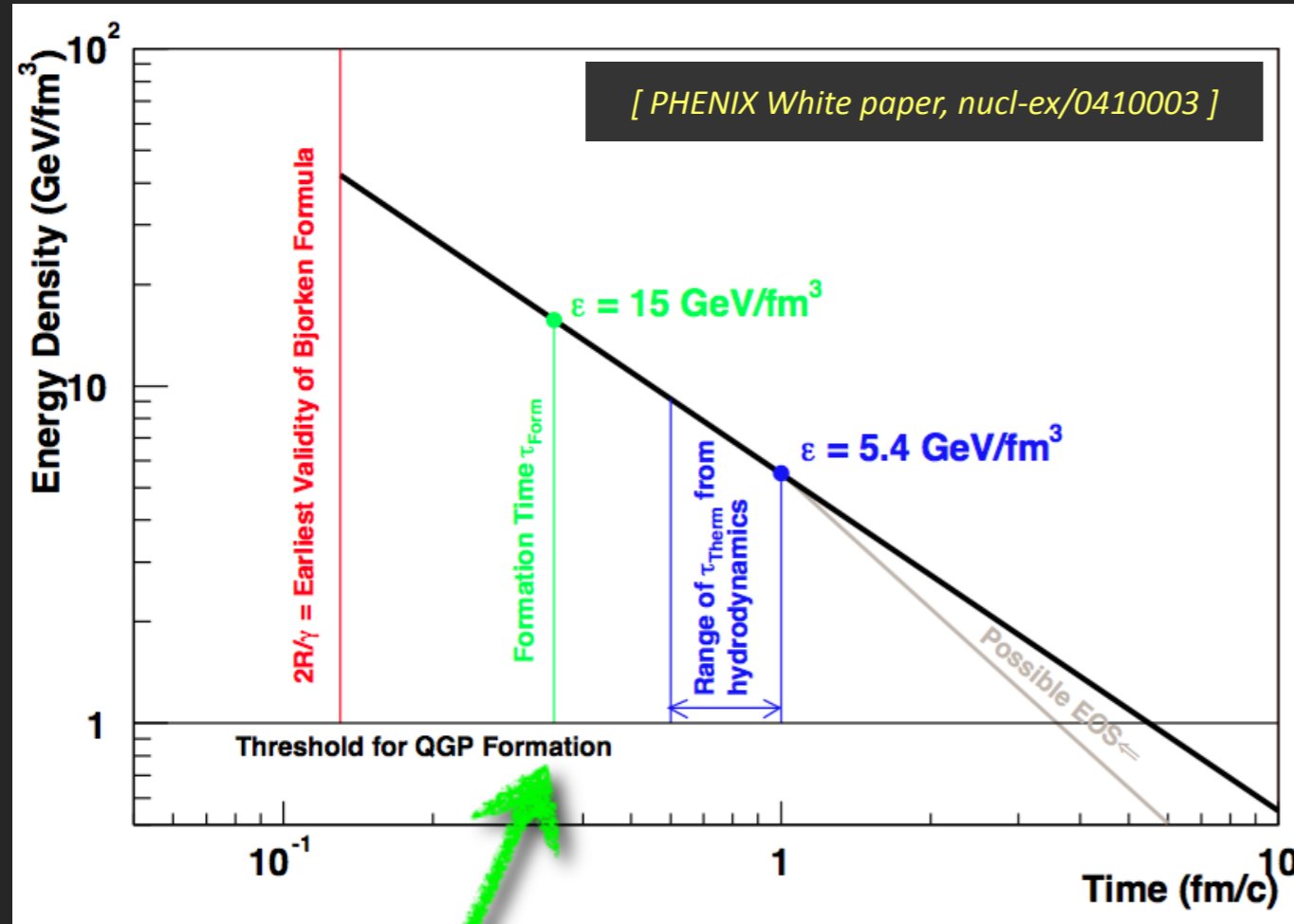
[ Strickland et al., NPA 879 (2012) 25-58 ]

[ Turbide et al., PRC 69 (2004) 014903 ]

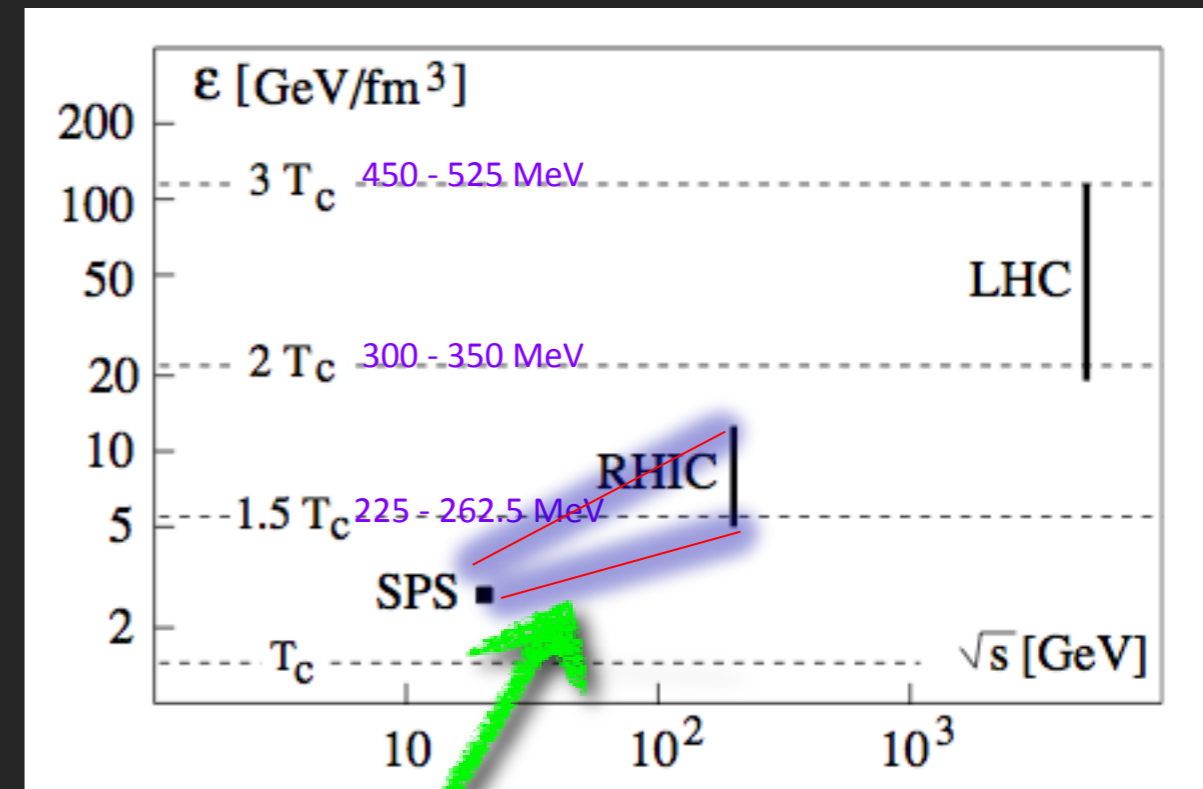
[ PHENIX, PRL. 104 (2010) 132301 ]

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AFTER in PbA

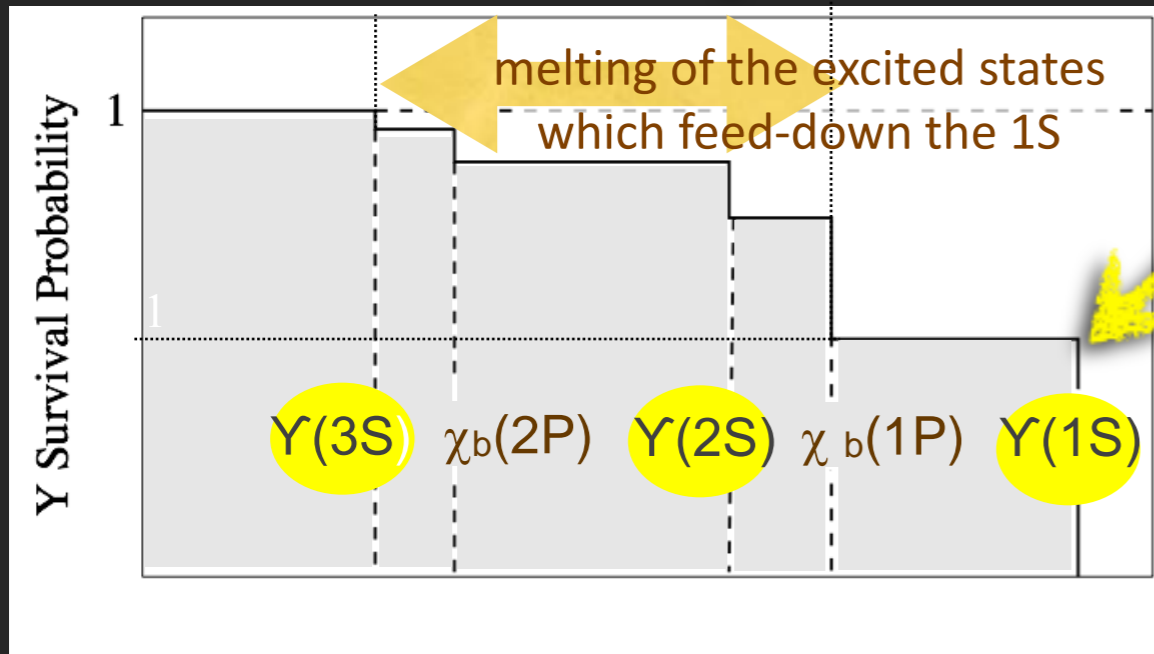
$\sqrt{s} = 72 \text{ GeV}$

[ Strickland et al., NPA 879 (2012) 25-58 ]

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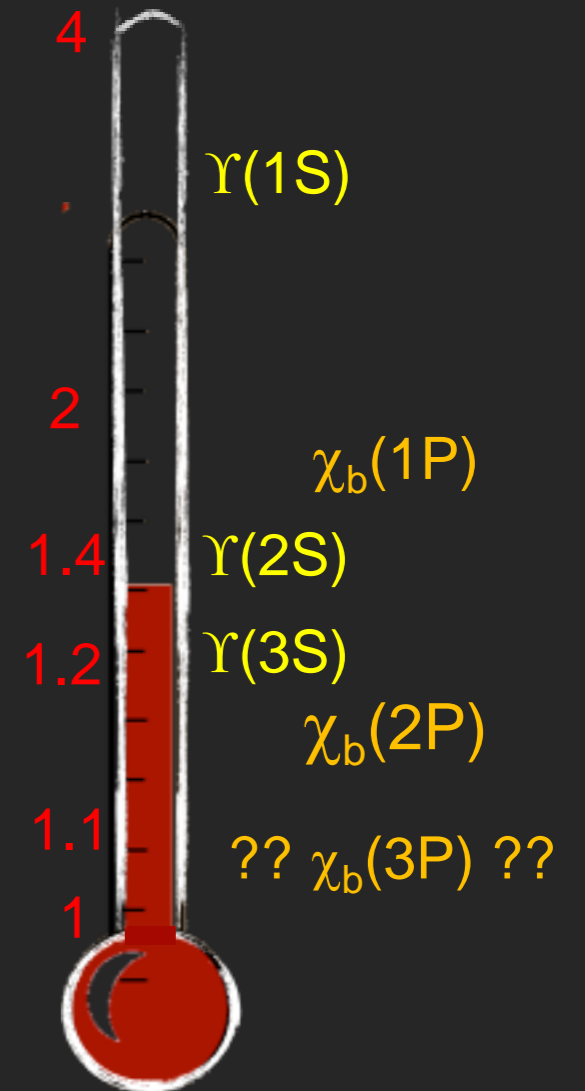
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# Sequential melting in QGP



Dissociation temperatures from lattice QCD (+hydro)

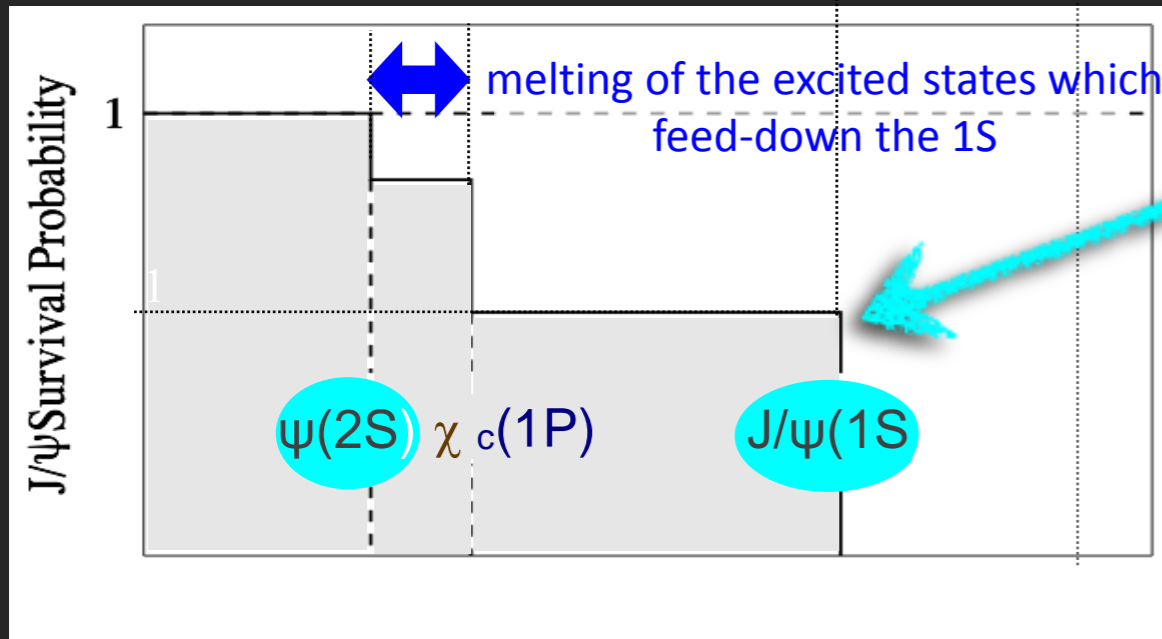
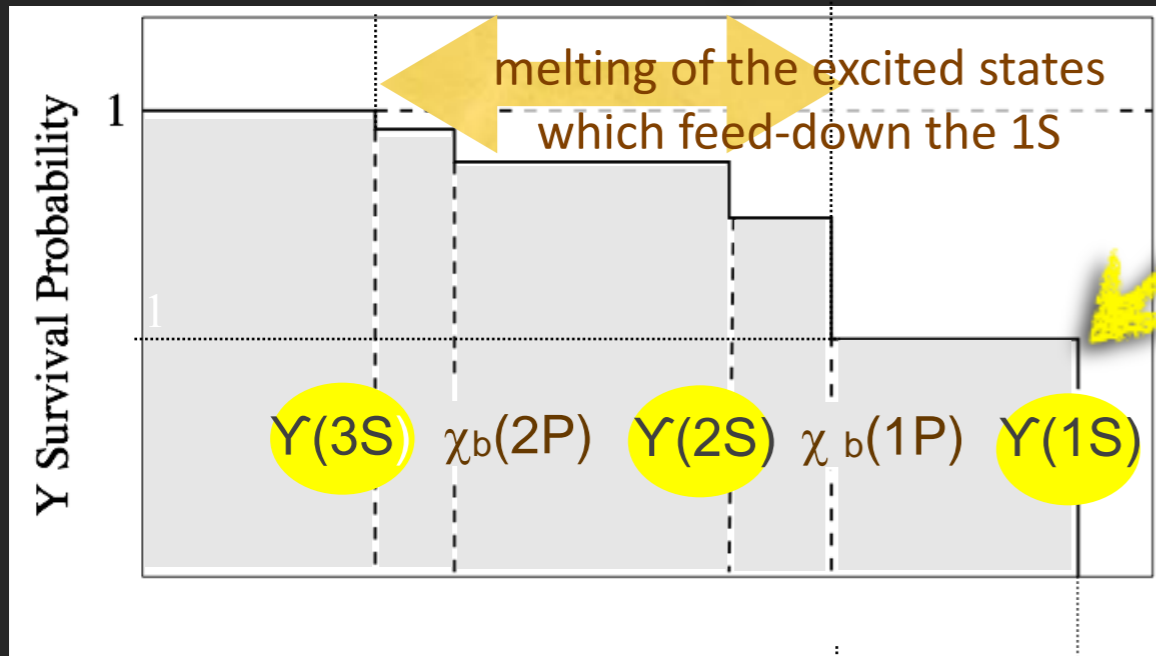
$T_d/T_c$



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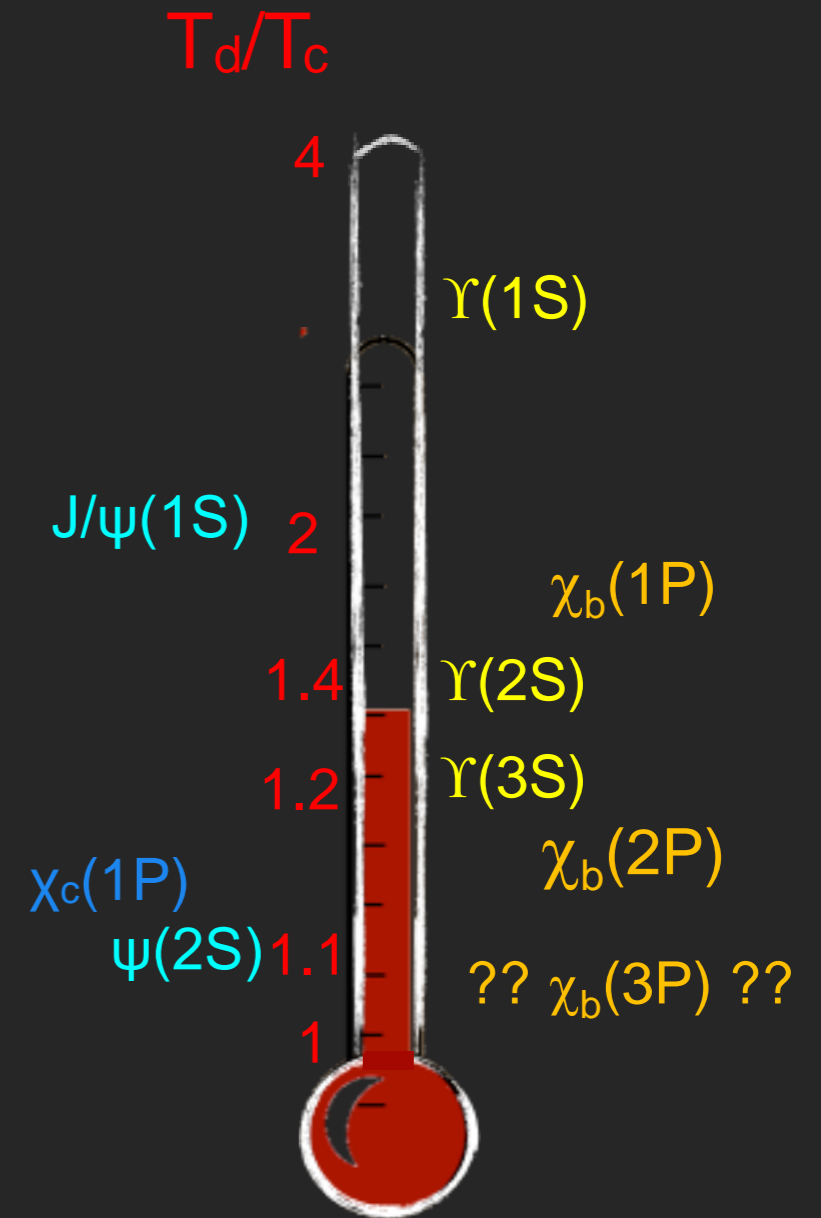
[ Mocsy et al., *Int.J.Mod.Phys. A28* (2013) 1340012 ]

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[ Satz, *Int.J.Mod.Phys. A28* (2013) 1330043 ]

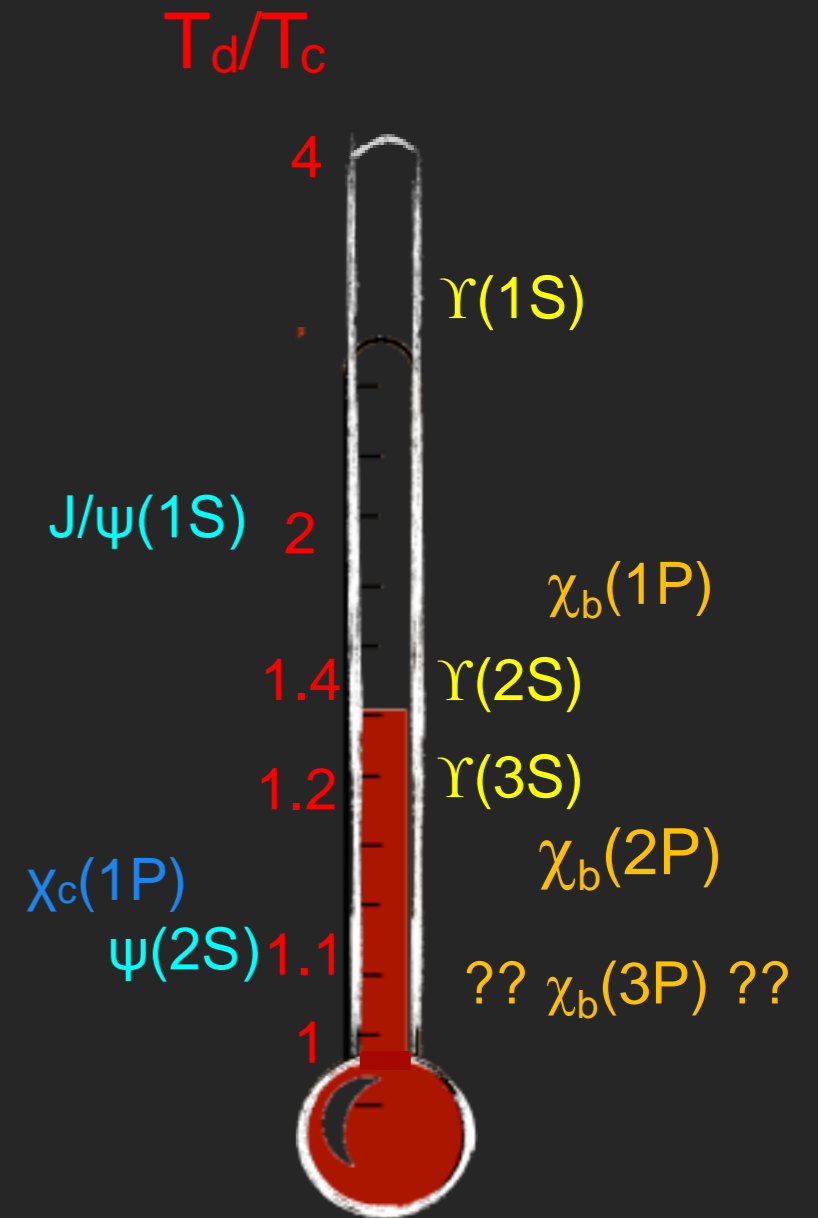
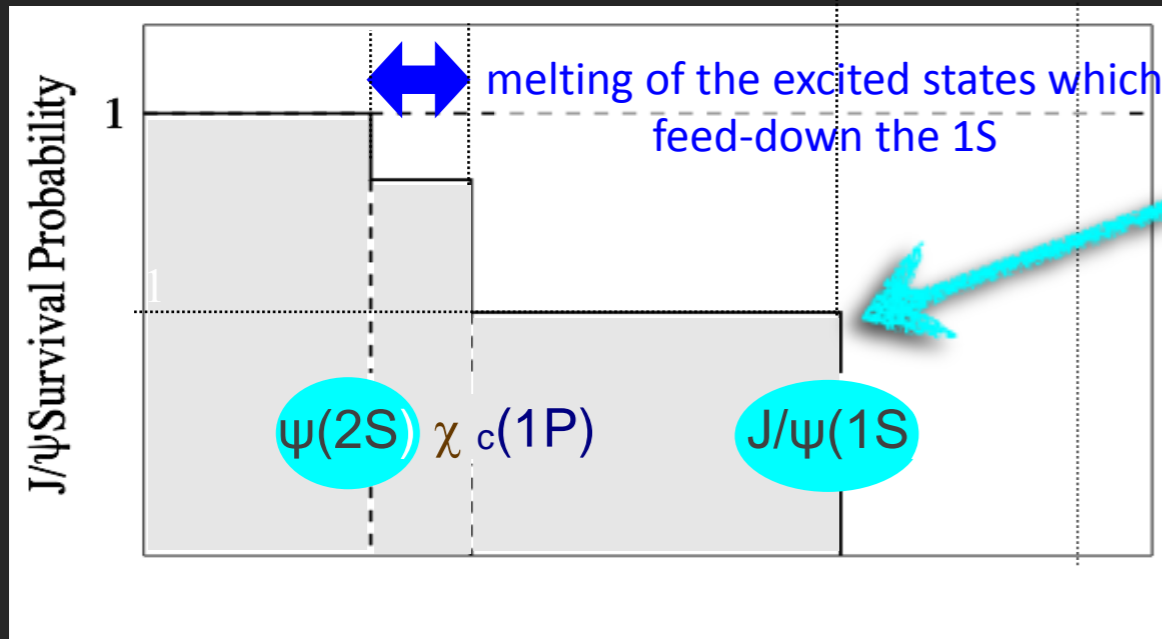
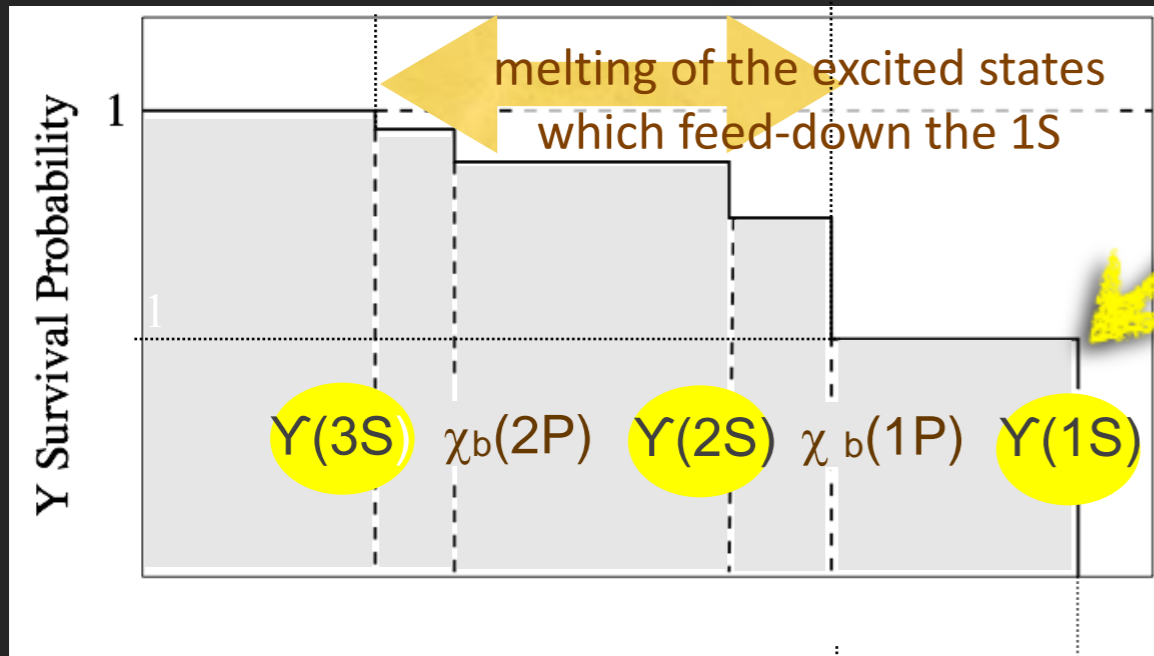
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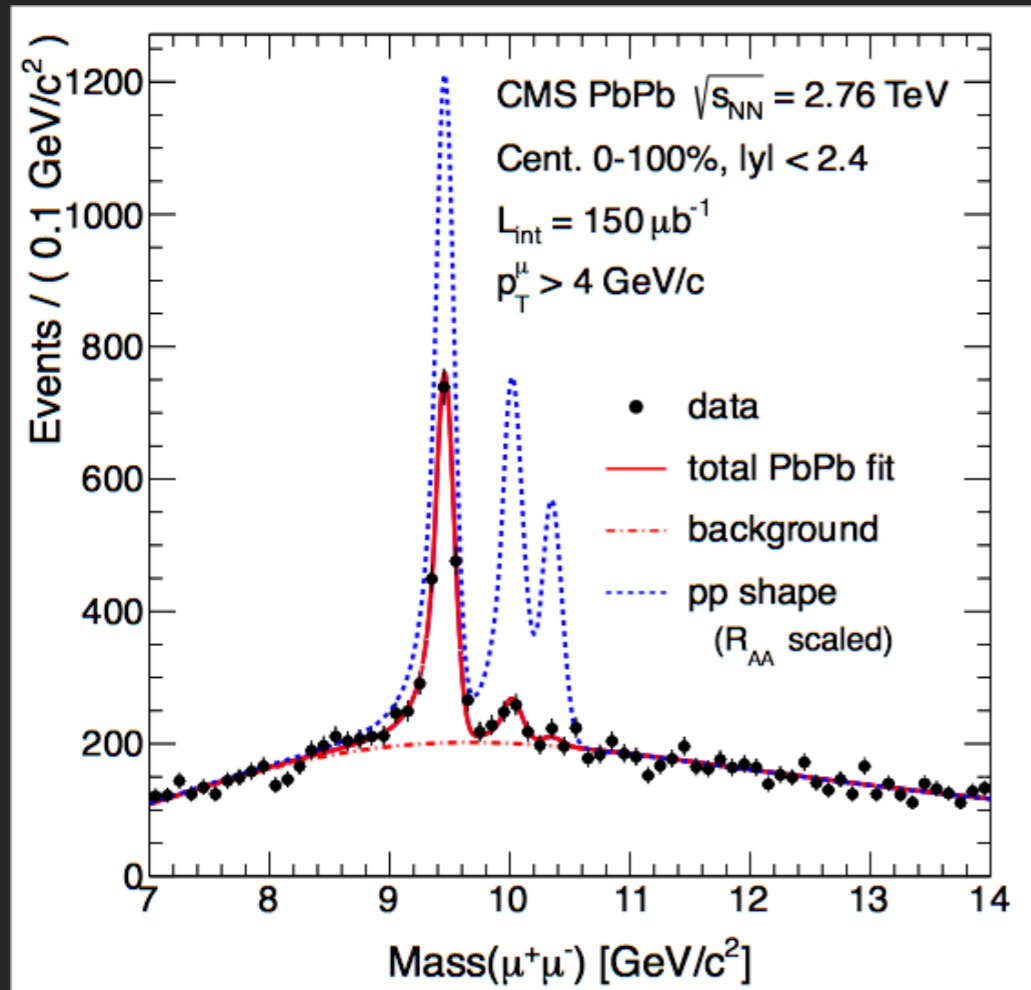
[ Satz, *Int.J.Mod.Phys. A28* (2013) 1330043 ]

- ▶ Bottomonium family : richer, broader range in  $T$  (compared to charmonium)
- ▶ Less necessary to measure the  $\chi_b(nP)$

[ Mocsy et al., *Int.J.Mod.Phys. A28* (2013) 1340012 ]

# Bottomonium sequential suppression @ LHC

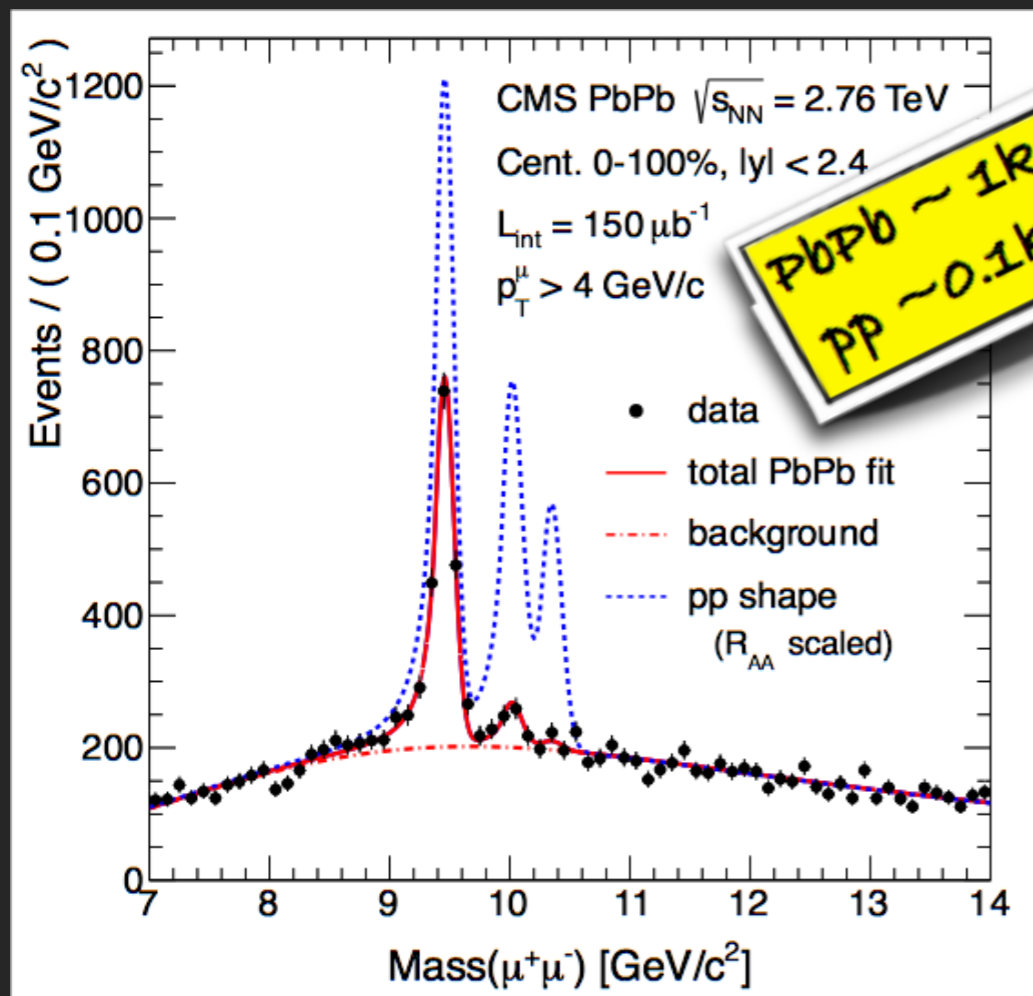
Serious candidate for a « textbook-like » plot at the recent Hard Probes 2013 conference



[ CMS, PRL 109 (2012) 222301 ]

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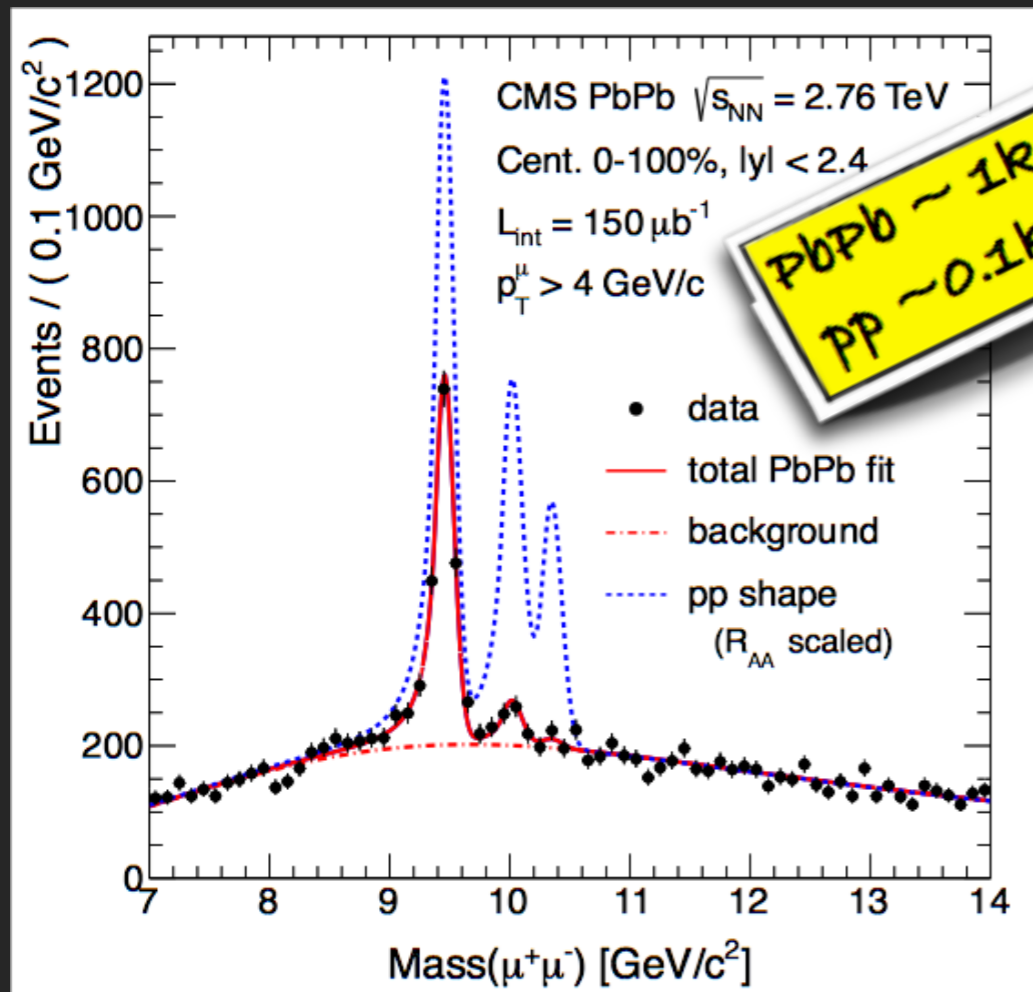
PbPb ~ 1k events  
PP ~ 0.1k events

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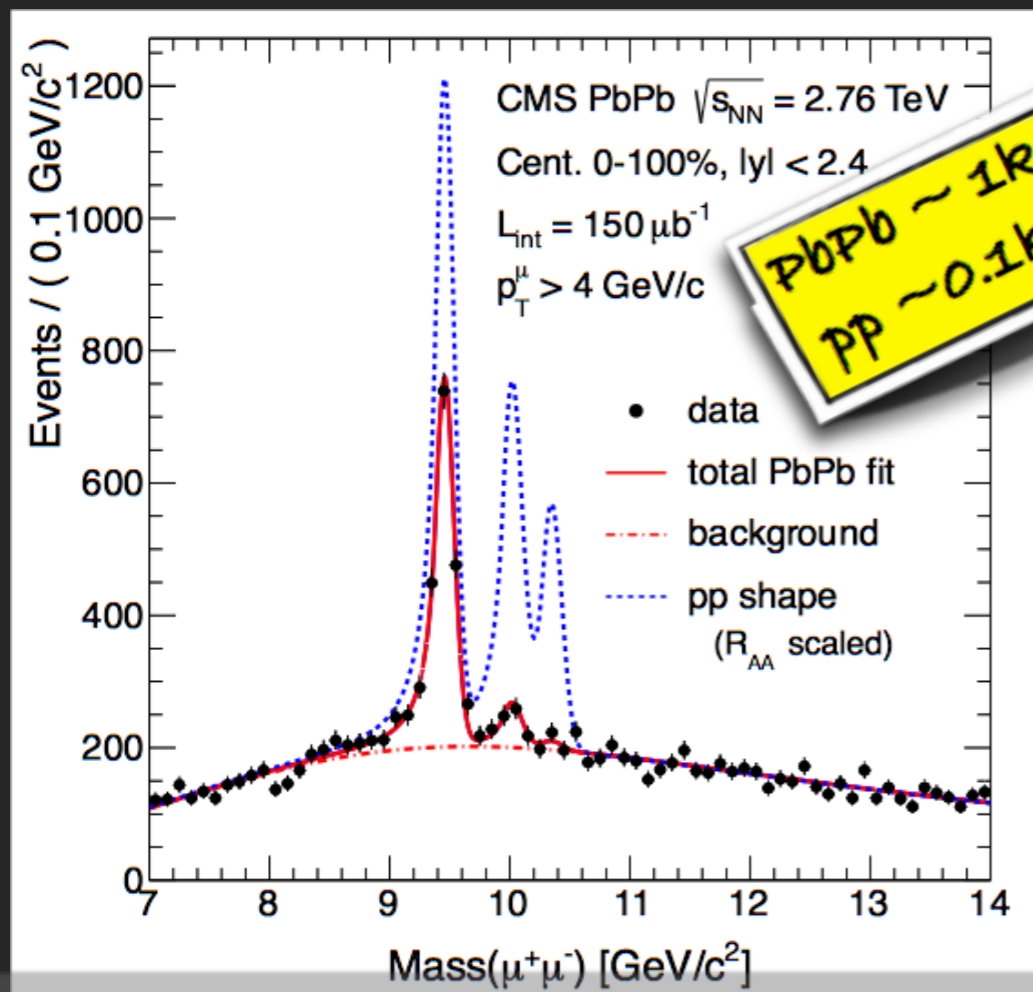


[ CMS, PRL 109 (2012) 222301 ]

- ▶ necessary ingredients :
- high inv. mass resolution in pp and PbPb + background under control
- ▶ Sequential suppression seen :
  - 3S completely melted ?
  - 2S very suppressed
  - direct 1S not affected ?

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PRL 109, 222301 (2012)

Selected for a **Viewpoint** in *Physics*  
PHYSICAL REVIEW LETTERS

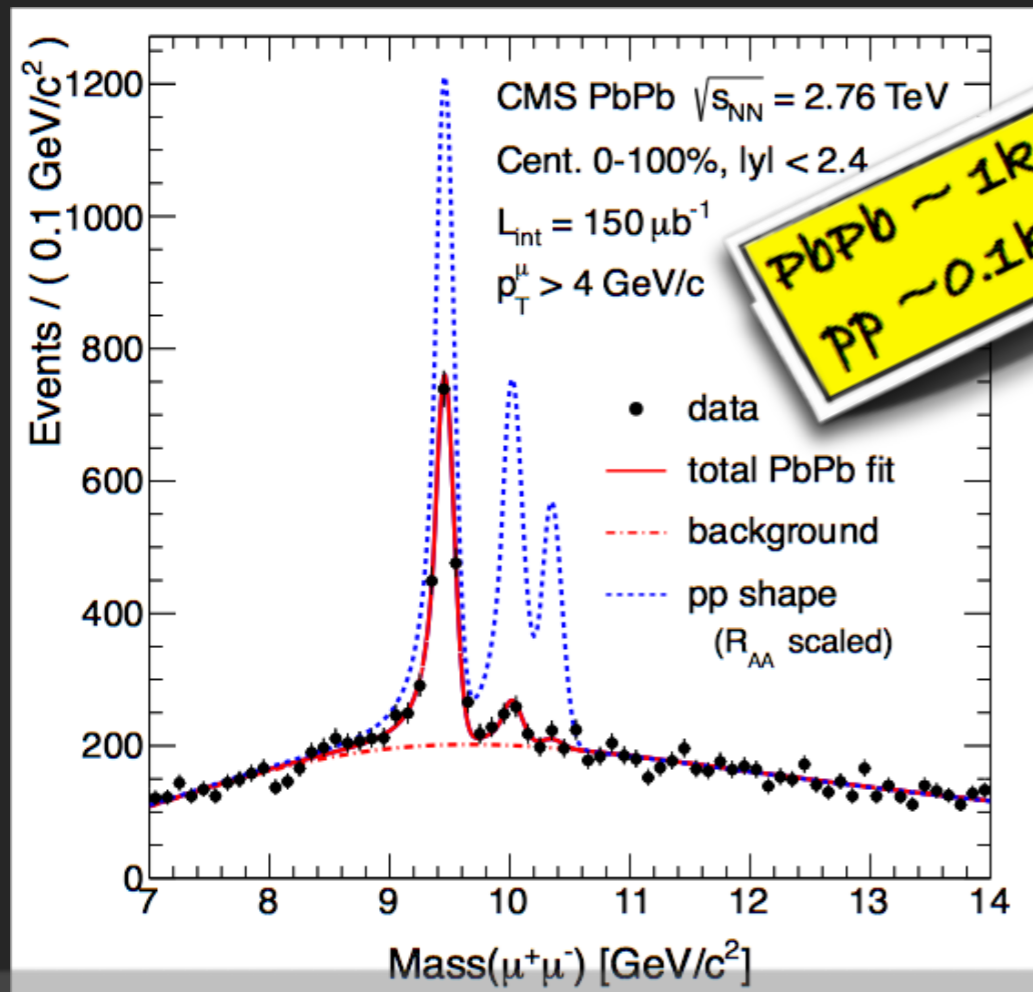
week ending  
30 NOVEMBER 2012

## Observation of Sequential $\Upsilon$ Suppression in PbPb Collisions

S. Chatrchyan *et al.*\*  
(CMS Collaboration)

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Physics

Viewpoint

**New Temperature Probe for Quark-Gluon Plasma**

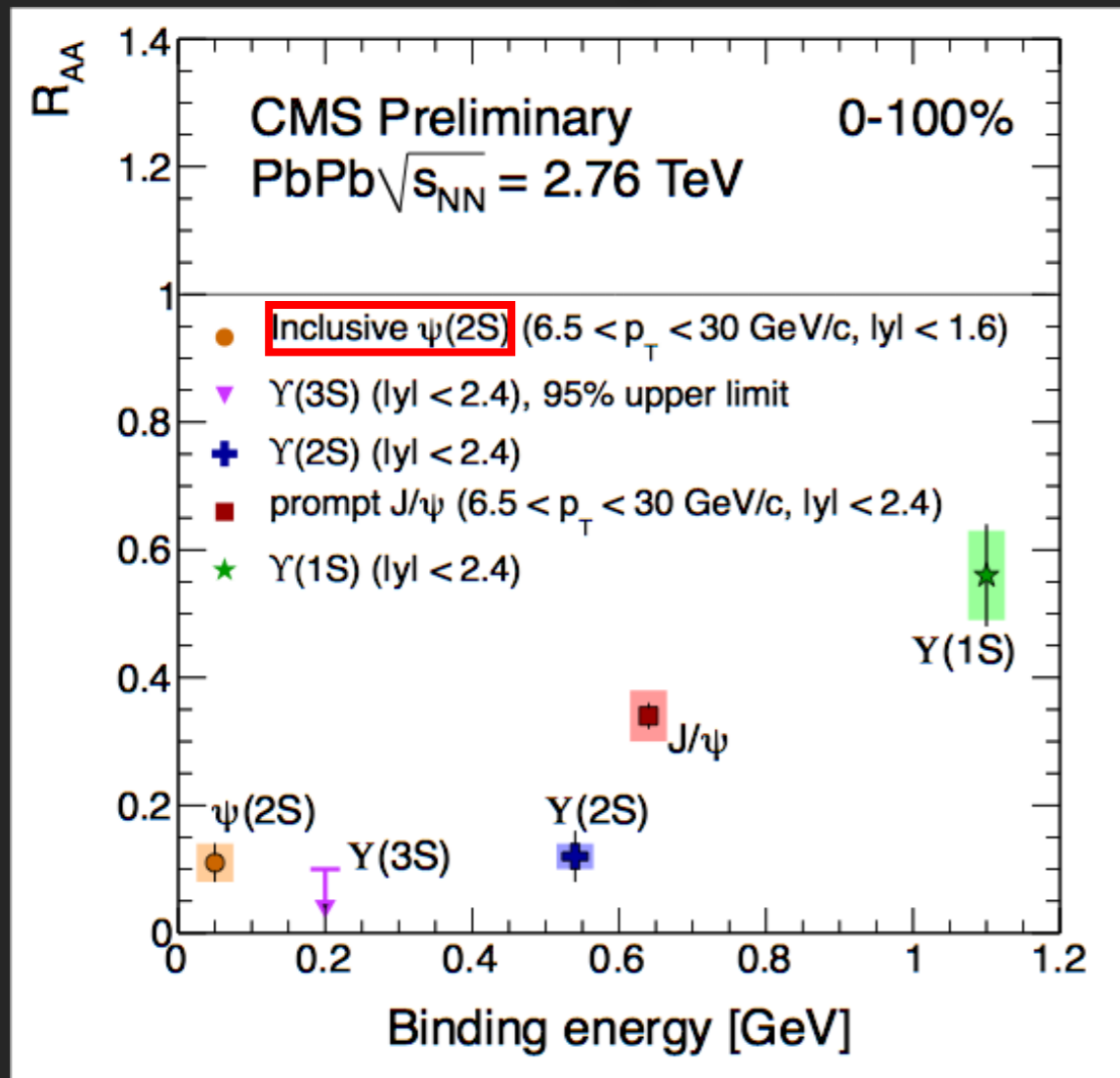
Ramona Vogt

Lawrence Livermore National Laboratory, Livermore, CA 94551, USA and  
Physics Department, University of California at Davis, Davis, CA 95616, USA

Published November 26, 2012  
The population of Upsilon mesons in quark-gluon plasma can be used to measure the plasma's temperature.

Physics 5, 132 (2012)

# Sequential *melting* @ LHC ?



- 3S completely melted ?
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[ Velkovska for CMS, HP2013 ]

$$R_{AA} \equiv \frac{\sigma_{AA}}{\langle N_{coll} \rangle \sigma_{pp}}$$

state  $R_{AA} \pm \text{stat} \pm \text{syst}$

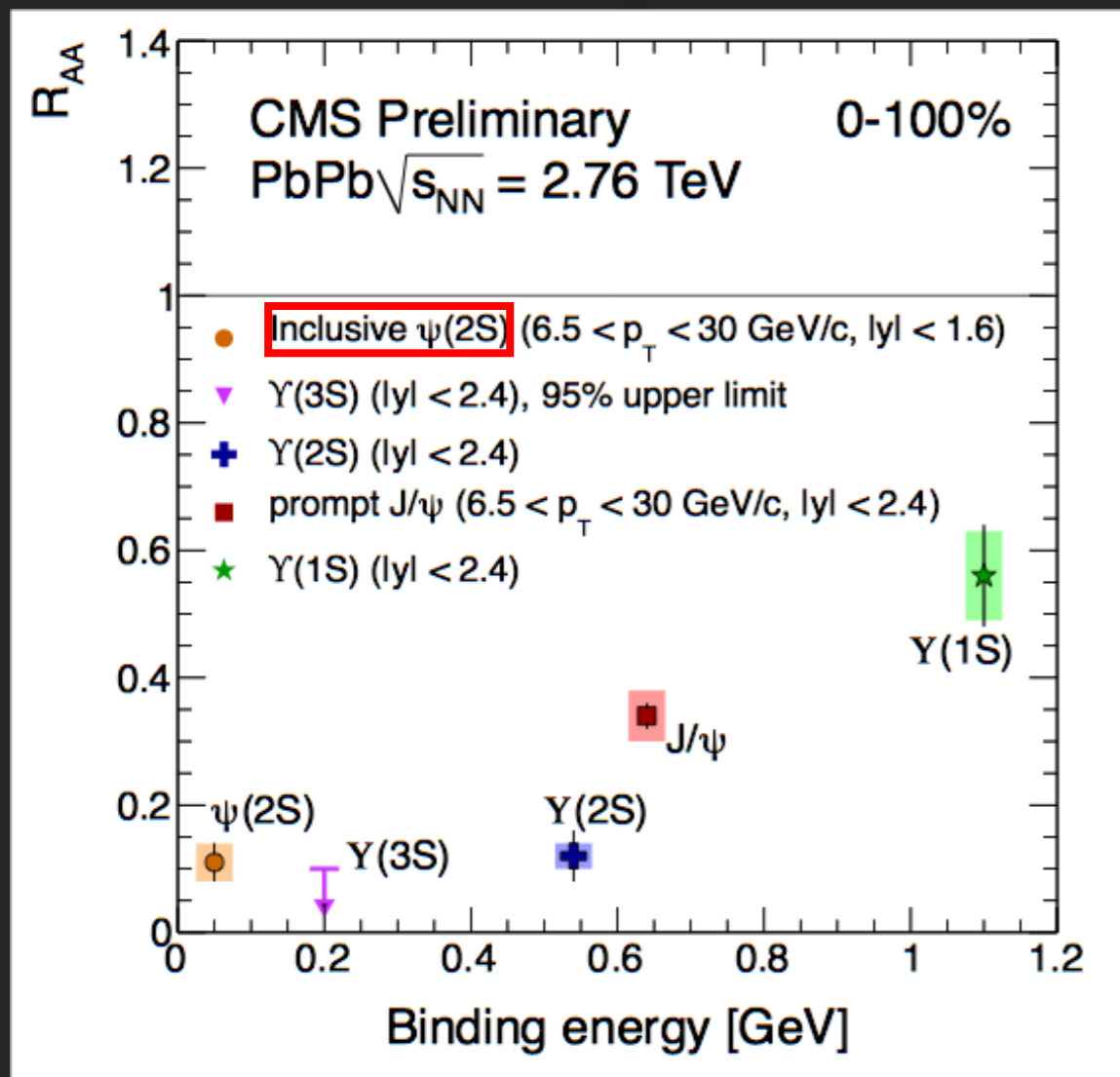
$Y(1S)$   $0.56 \pm 0.08 \pm 0.07$

$Y(2S)$   $0.12 \pm 0.04 \pm 0.02$

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[ CMS, PRL 109 (2012) 222301 ]

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If the sequential suppression is due to QGP effects *only*, what is the temperature reached @ LHC ?

▸ rough first estimate:

$1.4 T_c (\sim 230 \text{ MeV}) < T < 4 T_c (\sim 600 \text{ MeV})$

▸ lattice QCD + hydro evolution :

$T_{\text{initial}} \sim 550 \text{ MeV} > T$

[ Strickland et al., NPA 879 (2012) 25-58 ]

Measurement (thermal photons, dominant at low  $p_T$ ) :  $T_{\text{avg}} \sim 304 \pm 51 \text{ MeV}$  (0-40% PbPb)

[ Alice, NPA 904 (2013) 573c ]

# Lessons from SPS and RHIC

*Same  $J/\psi$  suppression observed at SPS & RHIC*

Two widely spread interpretations:

→ Melting of **excited** states of SPS & RHIC energies (1P & 2S)  
→ Induced suppression by feed-down  
→ **No additional melting** of the **direct** yield at RHIC  
→ Temperature between RHIC and SPS is somewhere between  $1.2$  &  $2 T_c$

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→ **Direct  $J/\psi$ 's partially melt**  
→ This additional suppression is compensated by **regeneration**  
→ Temperature between RHIC and SPS range from  $1.2 T_c$  up to  $2 T_c$

Nota: The fact that the  $\Upsilon(2S)$  would only be partially suppressed at LHC energies does not fit well with the hypothesis that the  $J/\psi$  already partially melts at RHIC [Theory predictions (lattice, ...) sometimes disagree on this, though]

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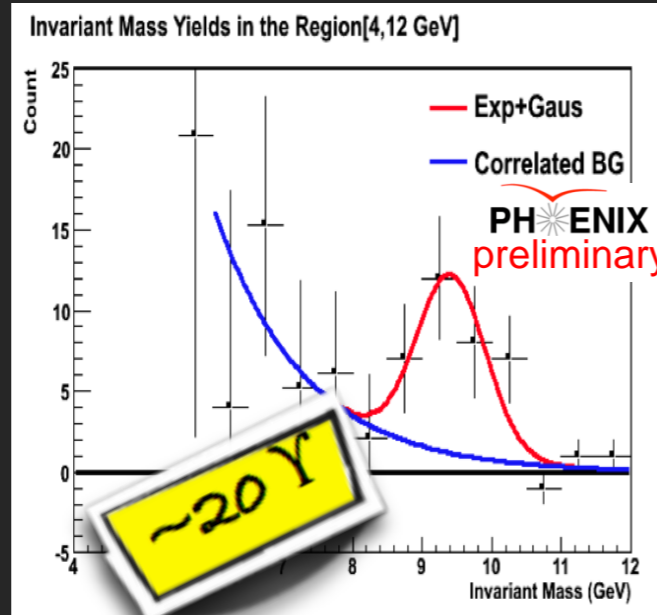
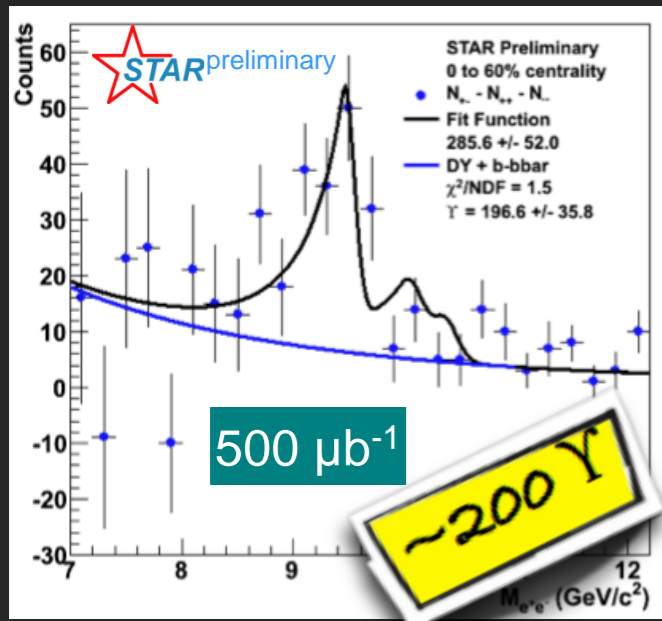
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In both case, the temperature expected for AFTER@LHC is likely around where the 2S and 3S bb states are expected to melt

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# Another hint: $\Upsilon(1S + 2S + 3S)$ suppression @ RHIC

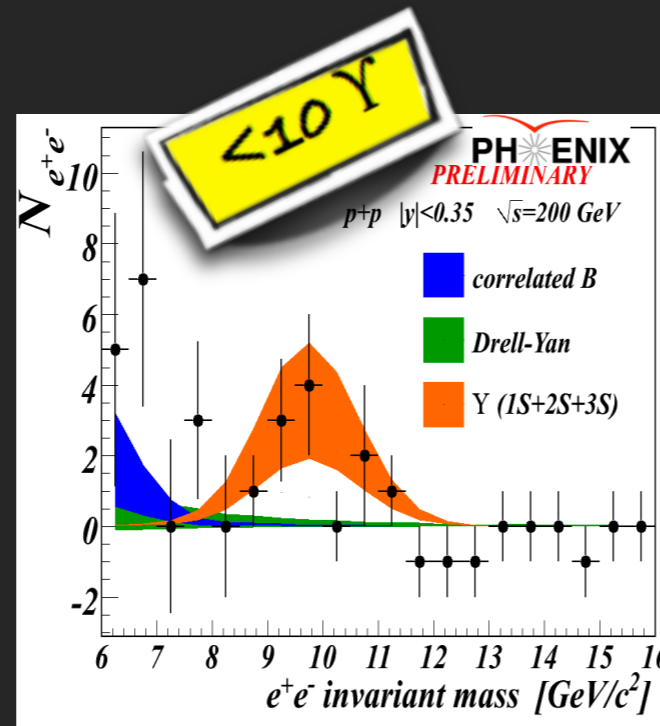
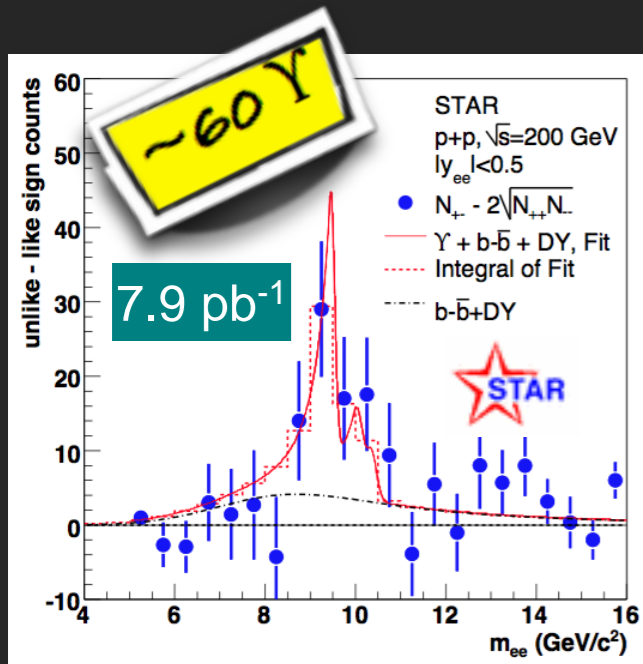
AuAu@200GeV (STAR run 2007, PHENIX run 2010)



[ Reed for STAR, JPG 38 (2011) 124185 ]

[ Whitaker for PHENIX, poster at QM2012 ]

pp@200GeV (run 2006)



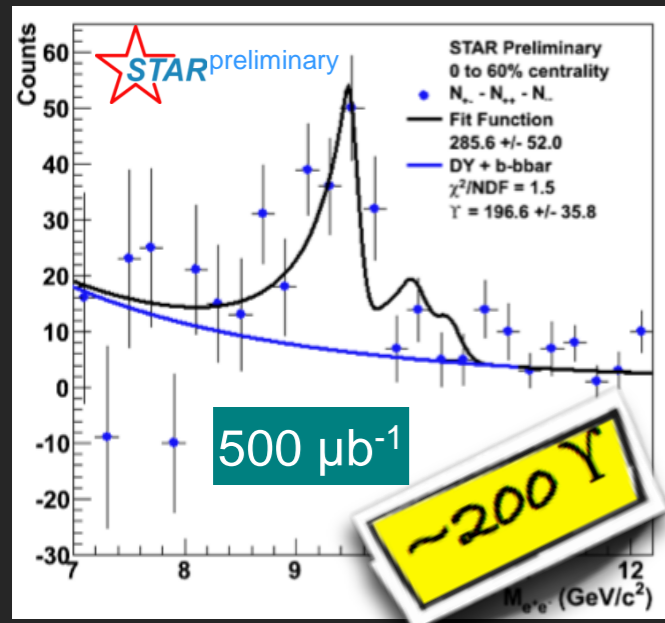
[ STAR, PRD 82 (2010) 012004 ]

[ Leitch for PHENIX, QM2009 ]

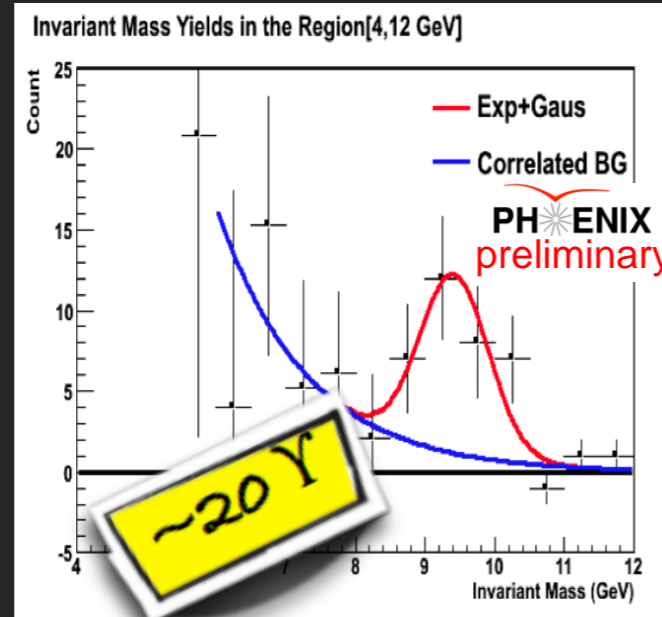


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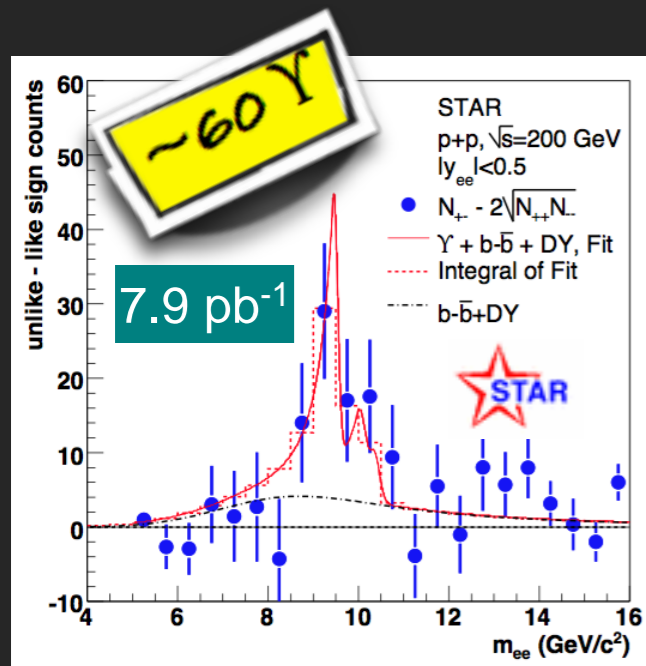
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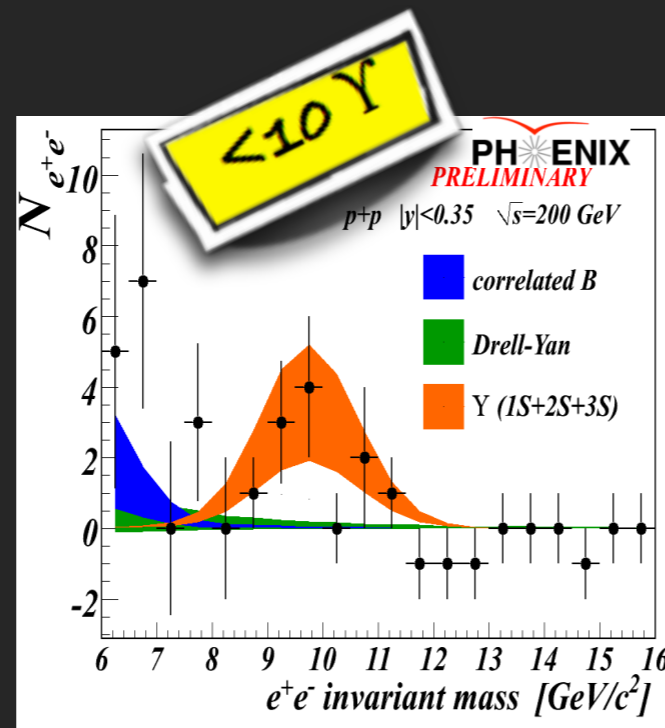
[ Whitaker for PHENIX, poster at QM2012 ]

► Not enough stat. (and resolution) to get separate results for the 3 states

pp@200GeV (run 2006)



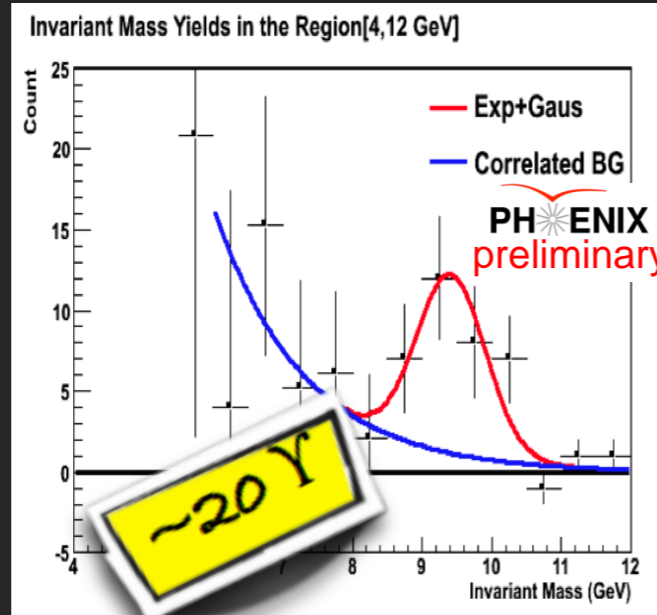
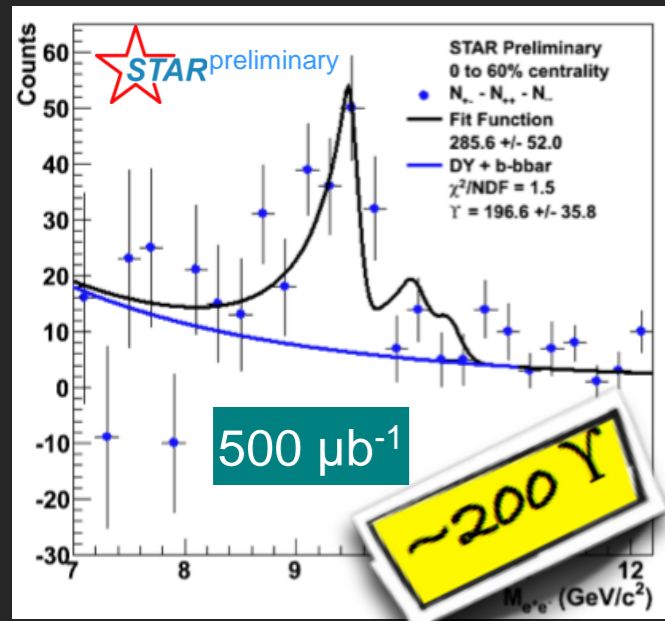
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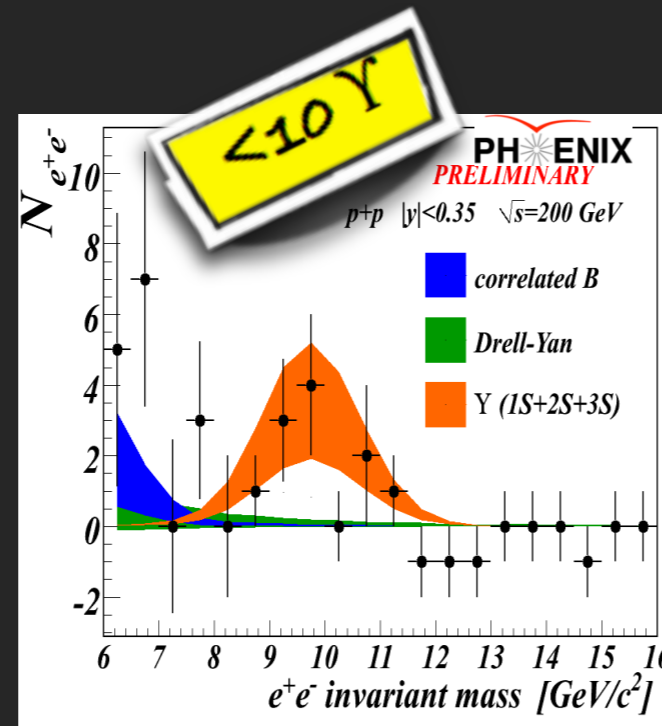
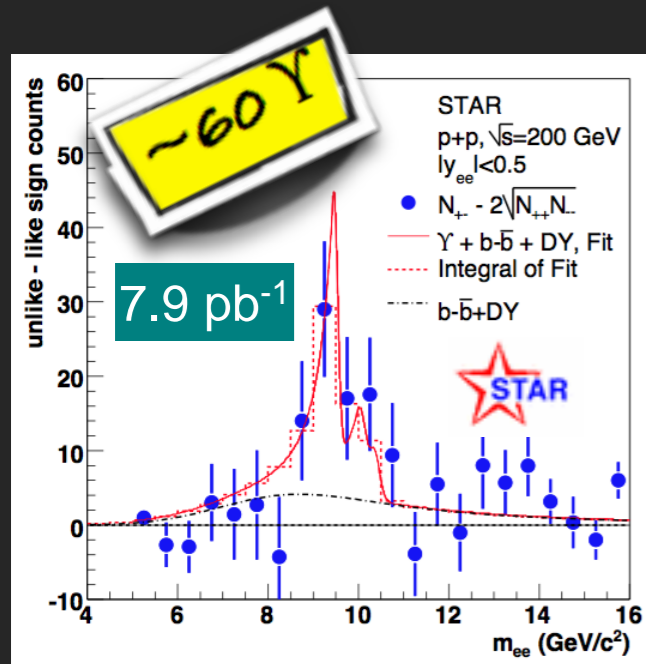
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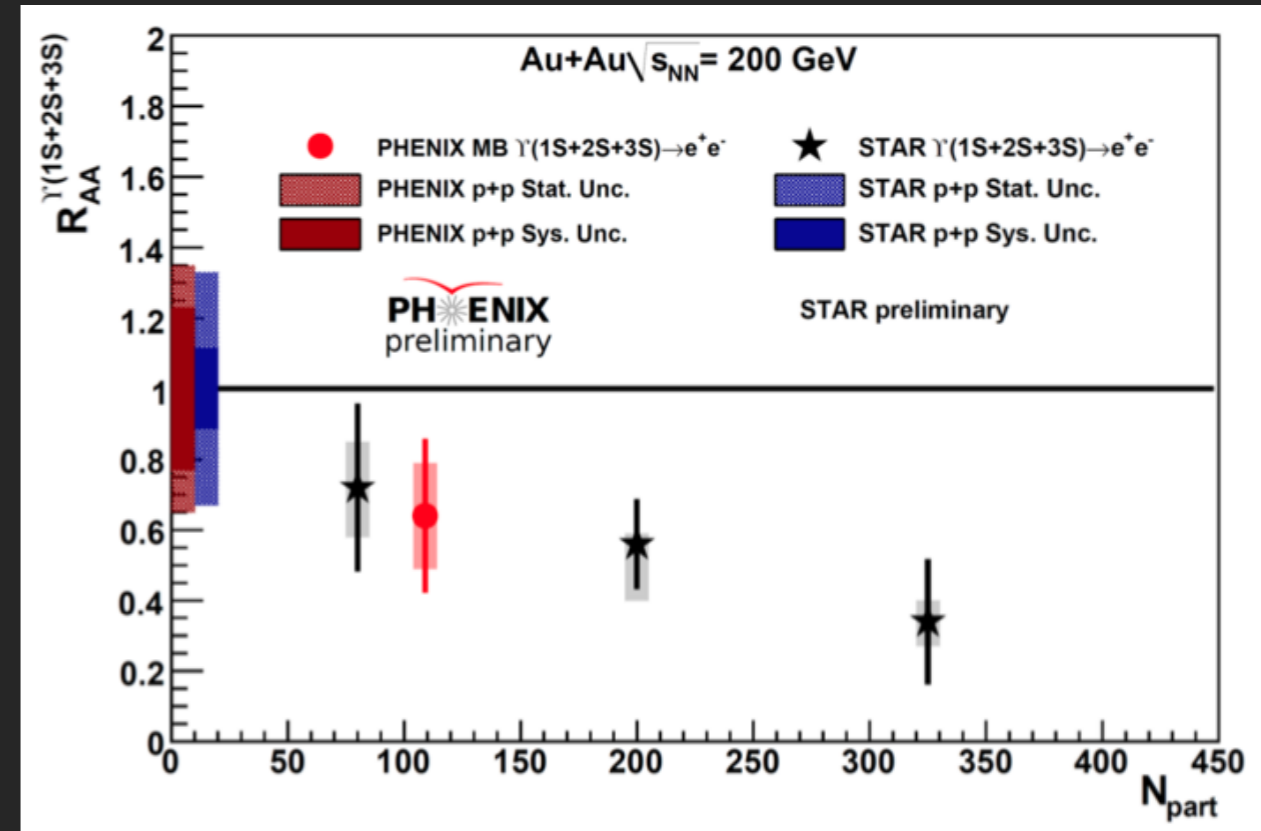
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specific  $R_{AA}$  computation for PHENIX:  
pp  $J/\psi$ ,  $\Upsilon$  run 2006  
AuAu  $\Upsilon$  run 2010

$$R_{AA}(\Upsilon) = \frac{[N(\Upsilon)/N(J/\psi)]_{AA}}{[N(\Upsilon)/N(J/\psi)]_{pp}} \times R_{AA}(J/\psi)$$



[ Whitaker for PHENIX, poster at QM2012 ]

# Luminosities

Instantaneous luminosity :

$$\mathcal{L} = N_{\text{beam}} \times N_{\text{target}} = N_{\text{beam}} \times (\rho \cdot e \cdot N_A) \text{ with } e = \text{target thickness}$$

Planned luminosity for PHENIX :

- @ 200 GeV run14pp 12 pb<sup>-1</sup>, run14dAu 0.15 pb<sup>-1</sup>
- @ 200 GeV run15AuAu 2.8 pb<sup>-1</sup> ( 0.13 nb<sup>-1</sup> @ 62 GeV)

Nominal LHC luminosity PbPb 0.5 nb<sup>-1</sup>

# Luminosities

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7 TeV proton beam

$pp, pd, pA \sqrt{s} = 115 \text{ GeV}$

[S. Brodsky, F. Fleuret, C. Hadjidakis, J.P. Lansberg, Phys. Rep. 522 (2013) 239 ]

Target (1 cm thick)	$\rho$ (g cm <sup>-3</sup> )	A	$\mathcal{L}$ ( $\mu\text{b}^{-1} \text{s}^{-1}$ )	$\int \mathcal{L}$ ( $\text{pb}^{-1} \text{yr}^{-1}$ )
solid H	0.088	1	26	260
liquid H	0.068	1	20	200
liquid D	0.16	2	24	240
Be	1.85	9	62	620
Cu	8.96	64	42	420
W	19.1	185	31	310
Pb	11.35	207	16	160

Table 1: Instantaneous and yearly luminosities obtained with an extracted beam of  $5 \times 10^8 \text{ p}^+/\text{s}$  with a momentum of 7 TeV for various 1cm thick targets

extracted beam  $N_{\text{beam}} = 5 \cdot 10^8 \text{ p}^+/\text{s}$   
9 months running / year  $\Leftrightarrow 10^7 \text{ s}$

Planned luminosity for PHENIX :

- @ 200 GeV run14pp 12  $\text{pb}^{-1}$ , run14dAu 0.15  $\text{pb}^{-1}$
  - @ 200 GeV run15AuAu 2.8  $\text{pb}^{-1}$  ( 0.13  $\text{nb}^{-1}$  @ 62 GeV)
- Nominal LHC luminosity PbPb 0.5  $\text{nb}^{-1}$

# Luminosities

Instantaneous luminosity :

$$\mathcal{L} = N_{\text{beam}} \times N_{\text{target}} = N_{\text{beam}} \times (\rho \cdot e \cdot N_A) \text{ with } e = \text{target thickness}$$

7 TeV proton beam

$pp, pd, pA \sqrt{s} = 115 \text{ GeV}$

2.76 TeV lead beam

$Pbp, Pbd, PbA \sqrt{s} = 72 \text{ GeV}$

[S. Brodsky, F. Fleuret, C. Hadjidakis, J.P. Lansberg, Phys. Rep. 522 (2013) 239 ]

Target (1 cm thick)	$\rho$ (g cm <sup>-3</sup> )	A	$\mathcal{L}$ ( $\mu\text{b}^{-1} \text{s}^{-1}$ )	$\int \mathcal{L}$ ( $\text{pb}^{-1} \text{yr}^{-1}$ )
solid H	0.088	1	26	260
liquid H	0.068	1	20	200
liquid D	0.16	2	24	240
Be	1.85	9	62	620
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solid H	0.088	1	11	11
liquid H	0.068	1	8	8
liquid D	0.16	2	10	10
Be	1.85	9	25	25
Cu	8.96	64	17	17
W	19.1	185	13	13
Pb	11.35	207	7	7

Table 2: Instantaneous and yearly luminosities obtained with an extracted beam of  $2 \times 10^5 \text{ Pb}/\text{s}$  with a momentum per nucleon of 2.76 TeV for various 1cm thick targets

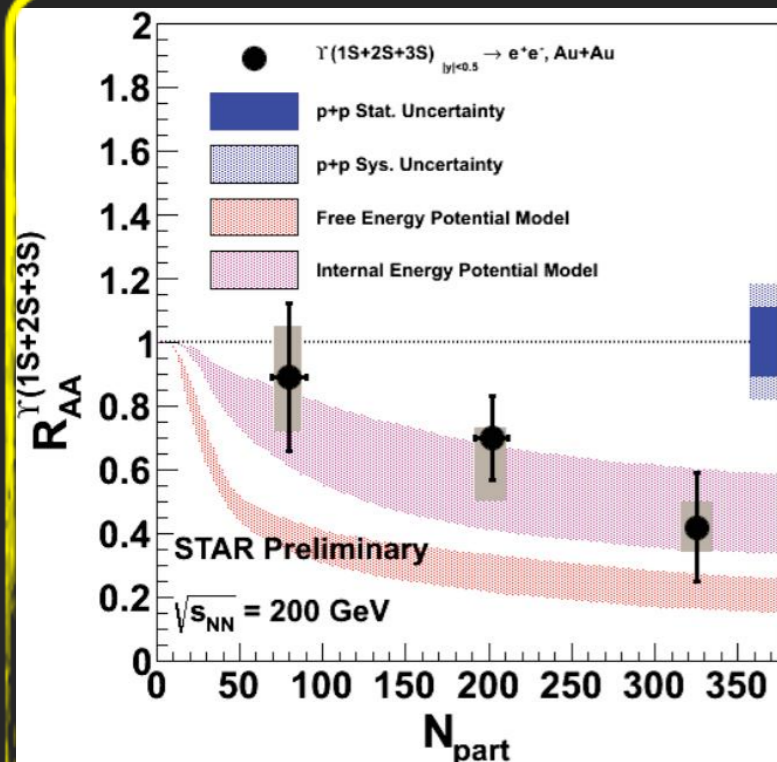
extracted beam  $N_{\text{beam}} = 5 \cdot 10^8 \text{ p}^+/\text{s}$   
9 months running / year  $\Leftrightarrow 10^7 \text{ s}$

extracted beam  $N_{\text{beam}} = 2 \cdot 10^5 \text{ Pb}/\text{s}$   
1 month running / year  $\Leftrightarrow 10^6 \text{ s}$

Planned luminosity for PHENIX :

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  - @ 200 GeV run15AuAu 2.8  $\text{pb}^{-1}$  ( 0.13  $\text{nb}^{-1}$  @ 62 GeV)
- Nominal LHC luminosity PbPb 0.5  $\text{nb}^{-1}$

# Bottomonium studies: from RHIC to AFTER



[ Bielcik for STAR, HP2013 ]

Today :

- ▶ inclusive  $\Upsilon$   $R_{AA}$  vs centrality
- ▶ the most central point is compatible with a complete melting of 3S and a very strong suppression of 2S, with  $T_{\text{initial}} \sim 430$  MeV in this model

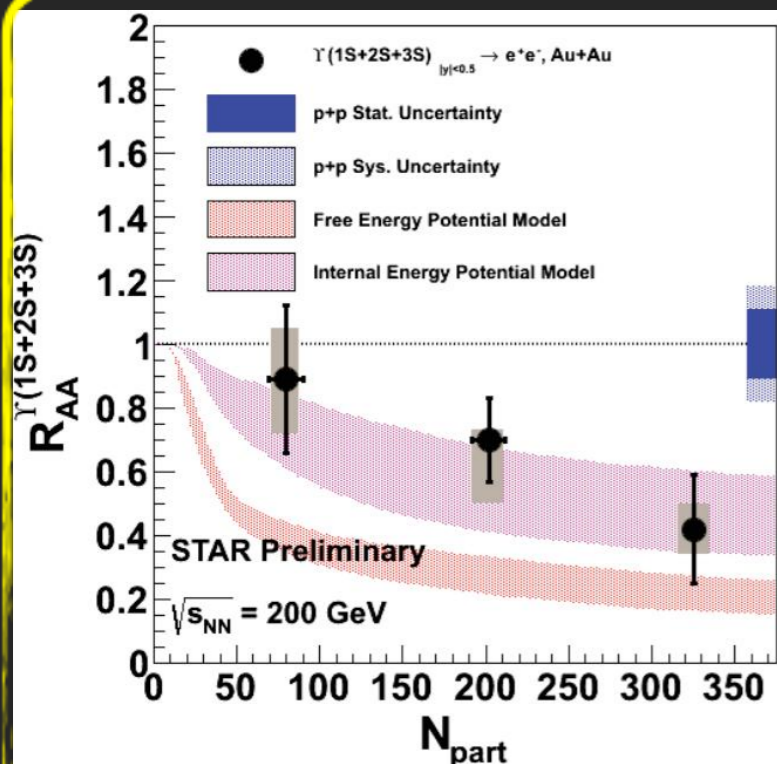
From thermal photon  $p_T$  spectra :

$$T_{\text{avg}} = 221 \pm 19 \text{ (stat)} \pm 19 \text{ (syst)}$$

MeV (0-20% AuAu)

[ PHENIX, PRL. 104 (2010) 132301 ]

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[ PHENIX, PRL. 104 (2010) 132301 ]

decompose this model into each state



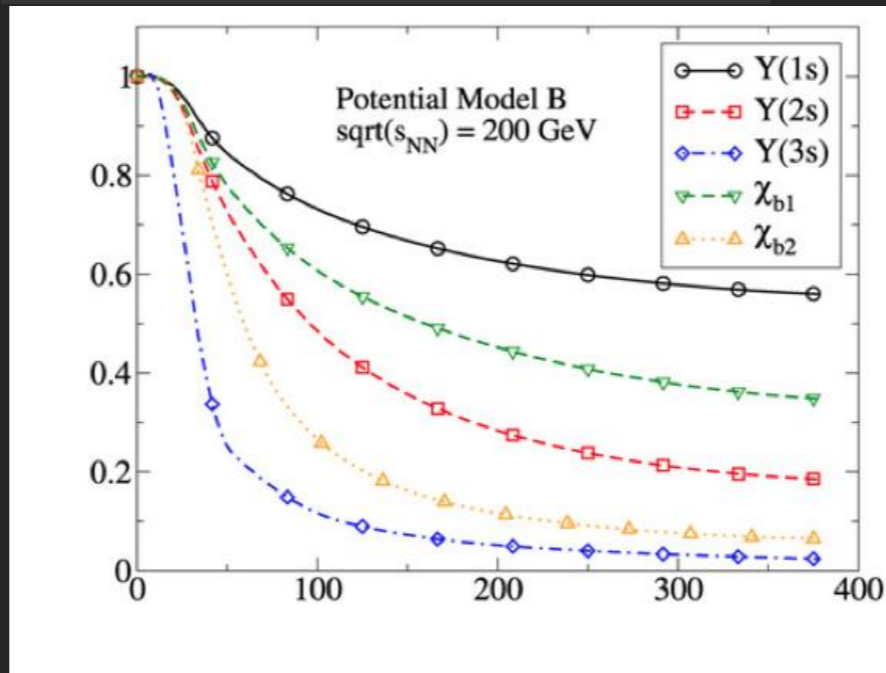
need more stat in AA + very good resolution

reminder

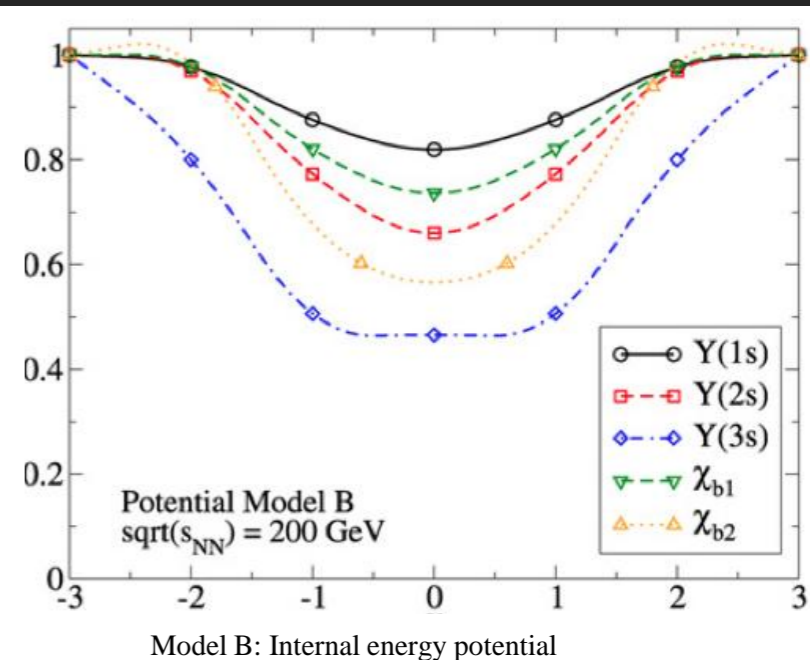
STAR :  $\sim 200$   $\Upsilon$

CMS :  $\sim 1k$   $\Upsilon$

[ Strickland et al., NPA 879 (2012) 25-58 ]



The dream measurements :

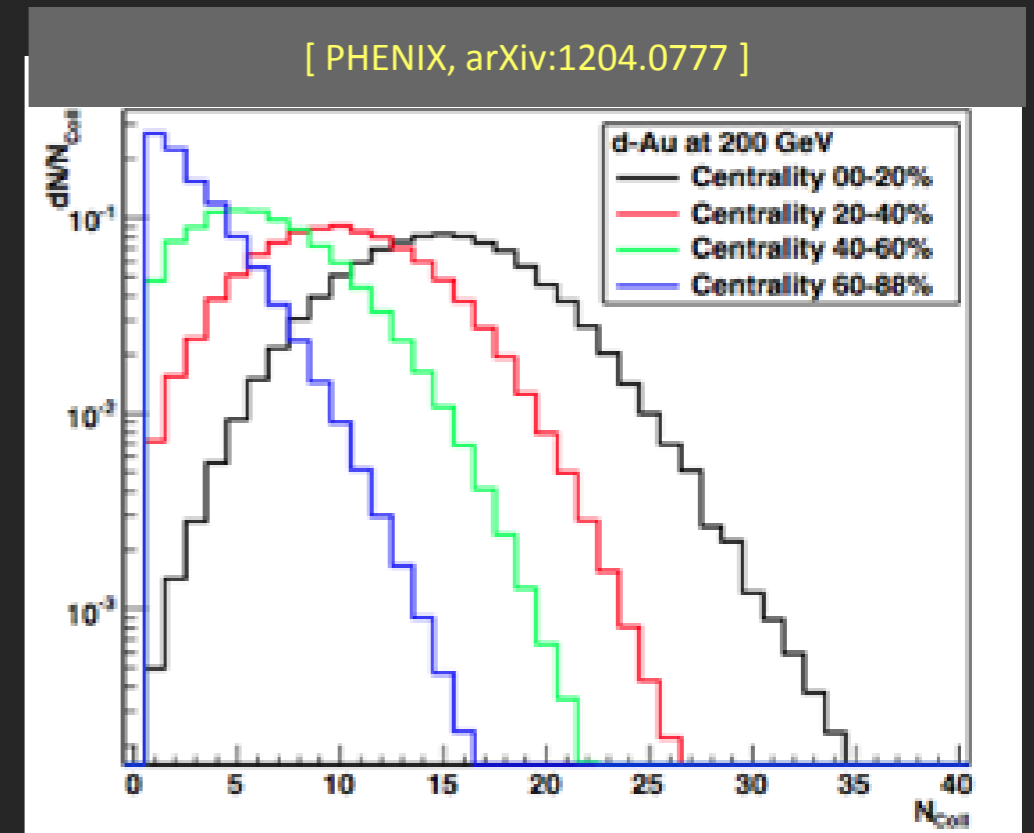


# High statistics $pA$ studies with AFTER: reference for nuclear effects & nPDF *per se*

$pA$

PbPb

- A dependence thanks to target **versatility**  
< $N_{\text{coll}}$ > dependence vs.  $A$  dependence (à la NA50, NA60)

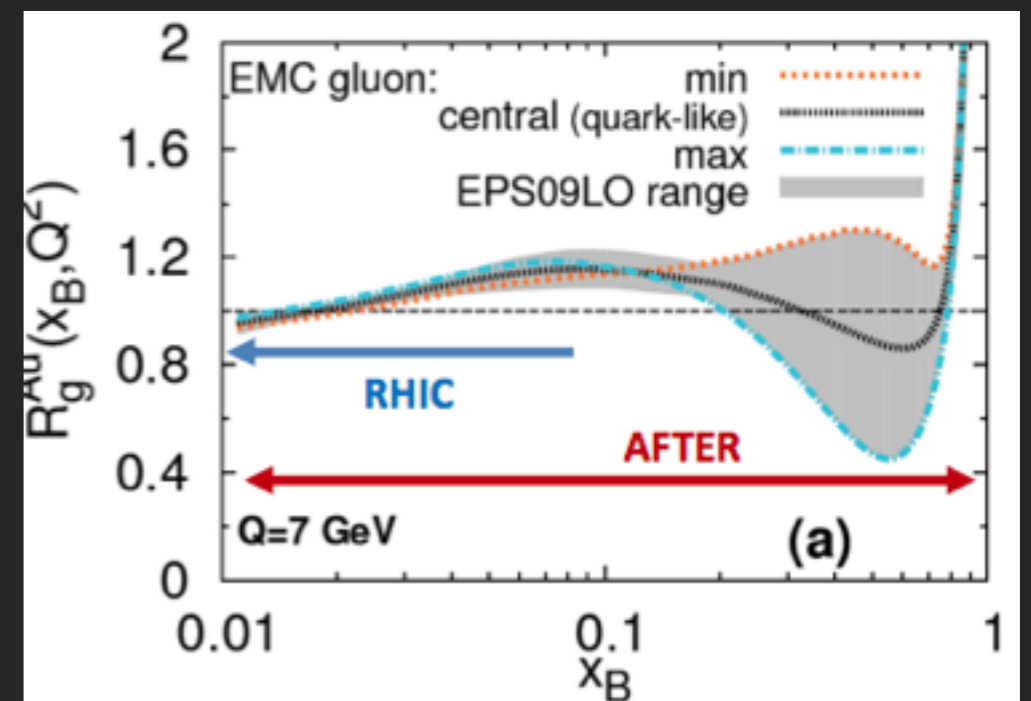
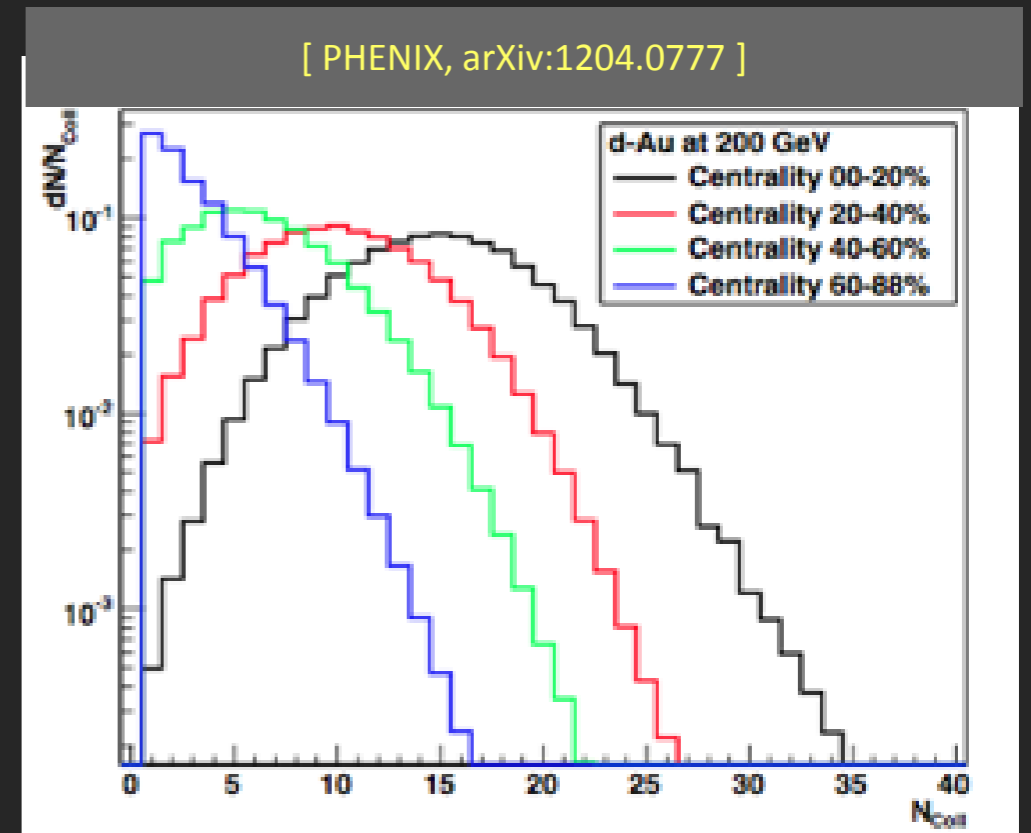




# High statistics $pA$ studies with AFTER: reference for nuclear effects & nPDF *per se*

$pA$   
PbPb

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- $\langle N_{\text{coll}} \rangle$  dependence vs.  $A$  dependence (à la NA50, NA60)
- nuclear PDF from intermediate to high  $x$  : antishadowing , EMC region , Fermi motion
- Gluon nPDF extraction using quarkonia (+ correlations), isolated photons, photon-jet correlation

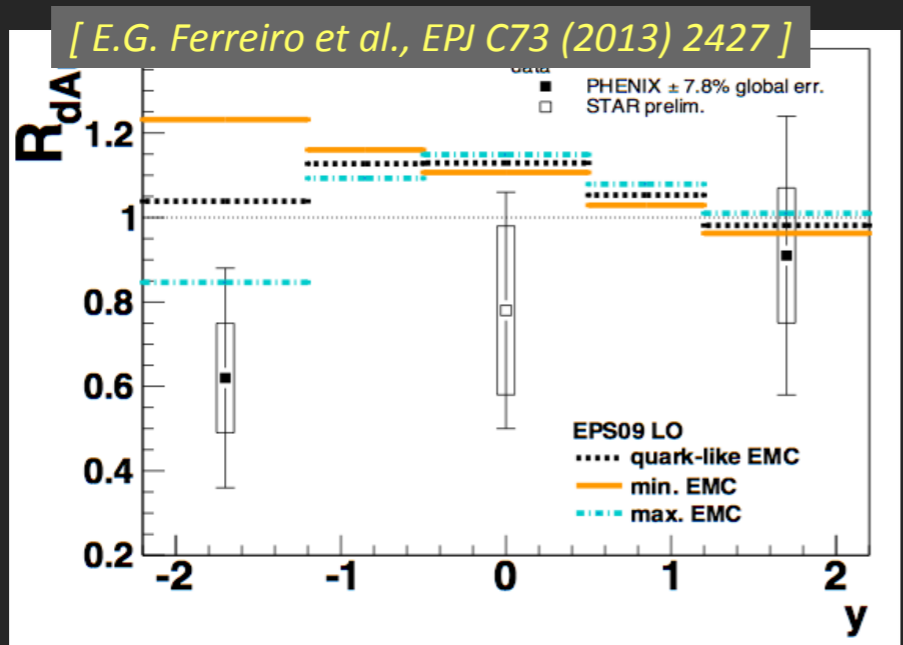
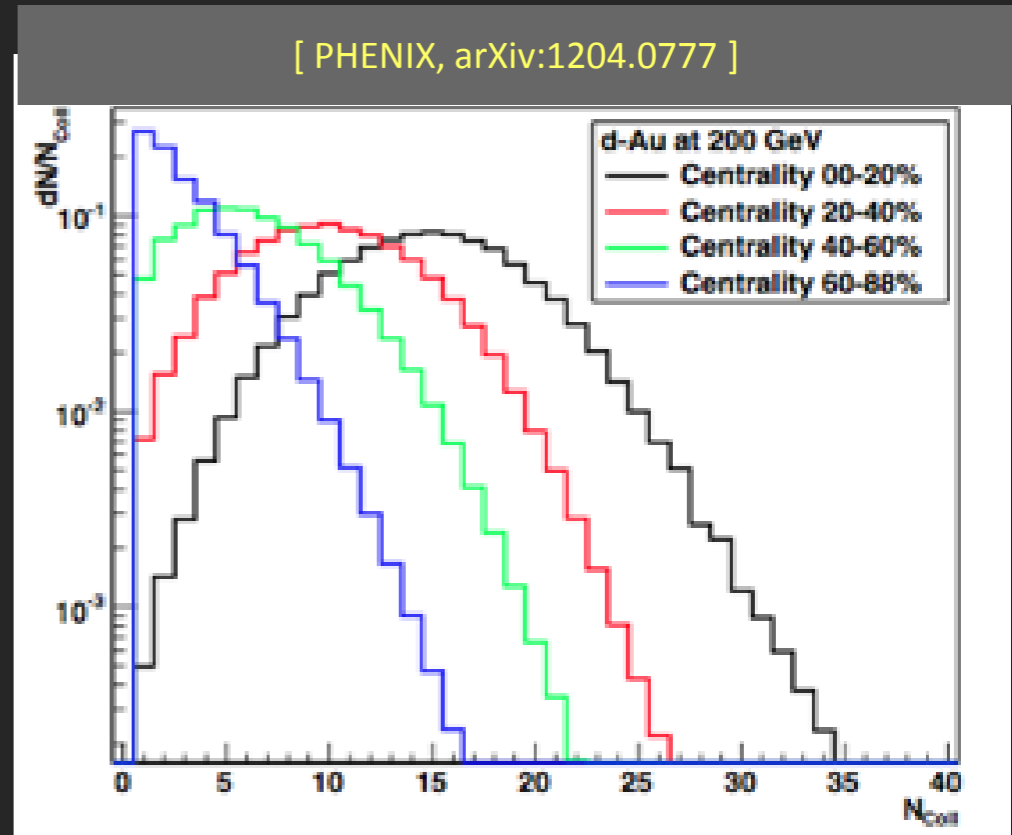


nuclear modification of  $g$  PDF in Au

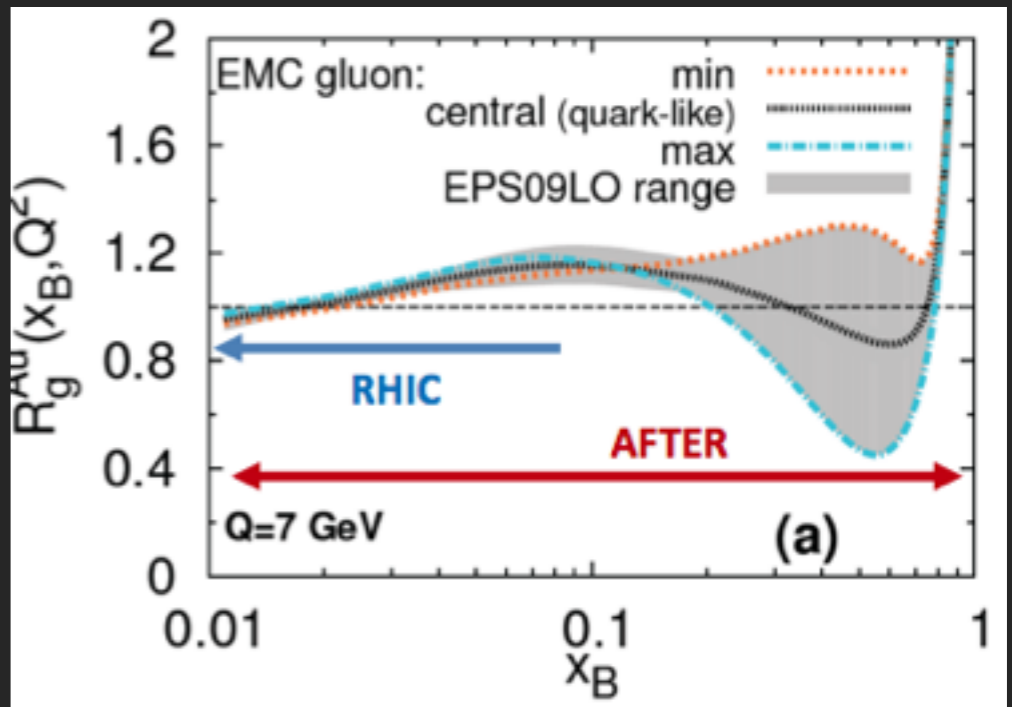
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- Gluon nPDF extraction using quarkonia (+ correlations), isolated photons, photon-jet correlation
- Strongly limited at RHIC



$Y$  in dAu @ 200 GeV,  
STAR and PHENIX



nuclear modification of  $g$  PDF in Au

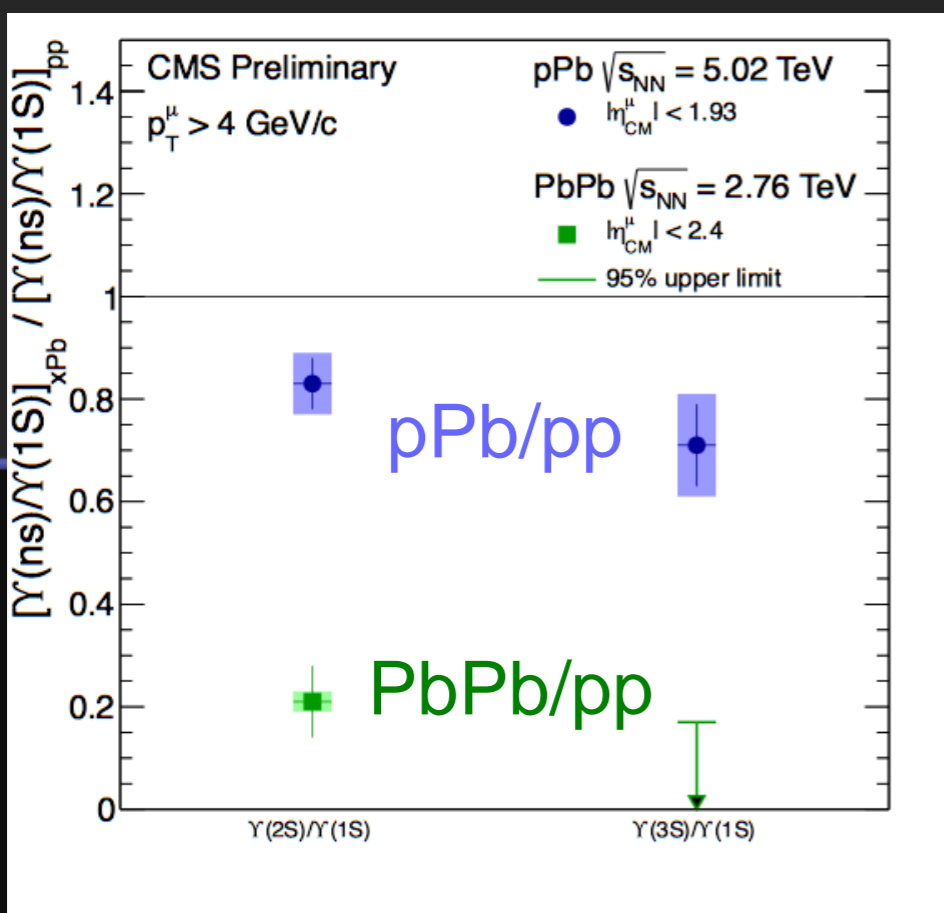
# Bottomonium : a cleaner QGP probe ?

- 3 states (2S & 3S not too fragile)
- Better applicability of pQCD w.r.t.  $J/\psi$
- in the QGP : negligible regeneration effects

BUT

# Bottomonium : a cleaner QGP probe ?

pPb vs. pp: excited states suppressed more than the ground state in pPb compared to pp collisions (significance  $< 3\sigma$  ?)



[ Benhabib for CMS, HP2013 ]

- 3 states (2S & 3S not too fragile)
- Better applicability of pQCD w.r.t. J/ $\psi$
- in the QGP : negligible regeneration effects

BUT

Cold effects (i.e. not QGP) :

- ◆ non-trivial effects seen in  $pA$  collisions
- ◆ need more studies and high stat  $pA$  measurements

➔ This is where AFTER cannot be challenged

# Summary and outlooks



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- **RHIC experiments cannot** resolve the 3 states and are limited by the luminosity (stronger limitation at 62 GeV)
- Measurement of  $\chi_b$  states not required, since we could use all  $3Y(nS)$  states, but would certainly add very interesting pieces of information.

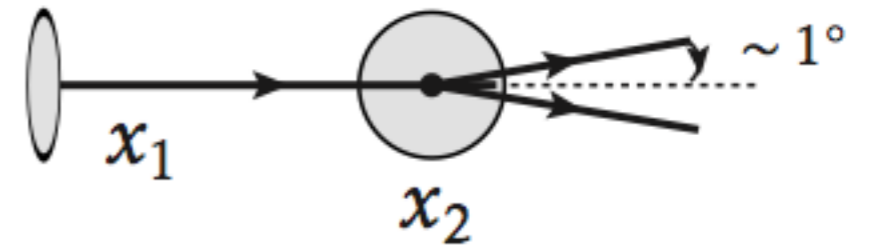
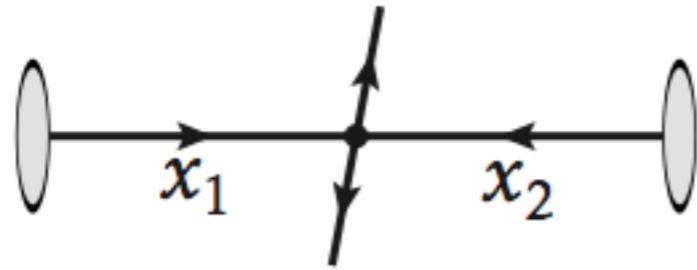
# SPARE SLIDES

# Backward physics

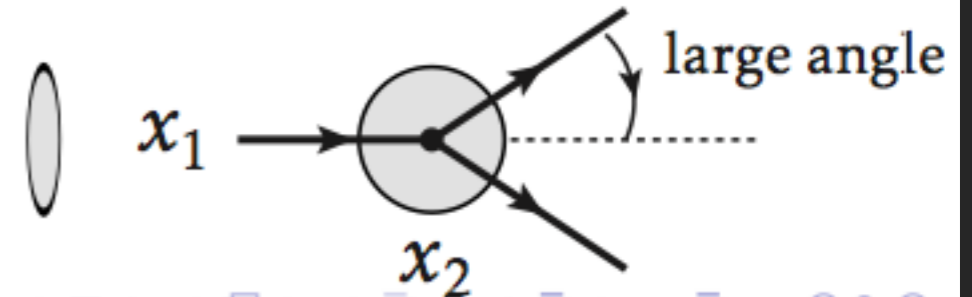
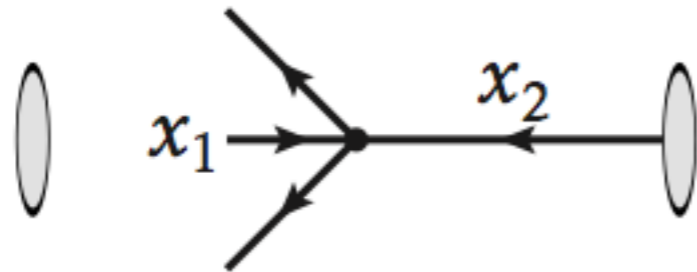
Hadron center-of-mass system

Target rest frame

$x_1 \simeq x_2$

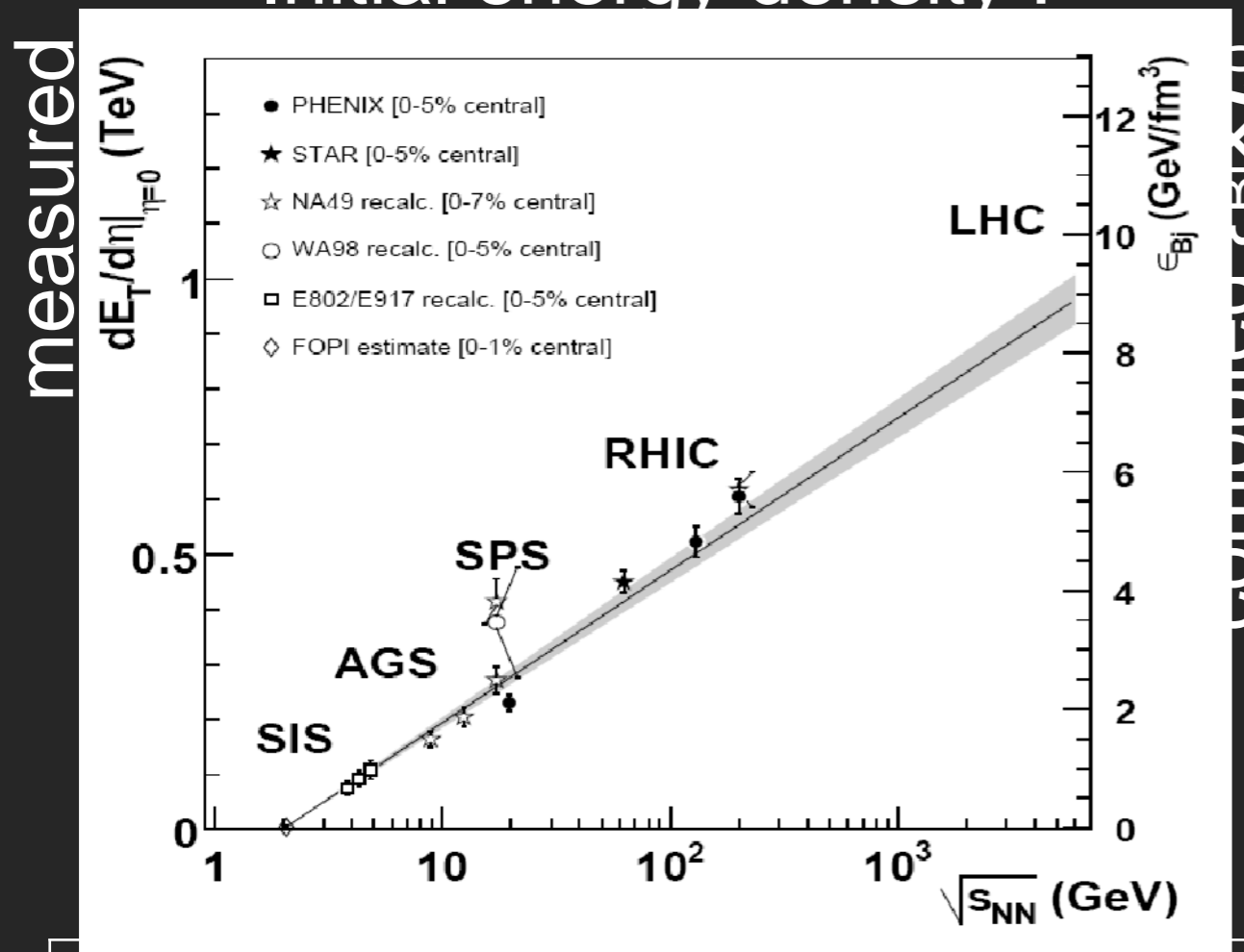


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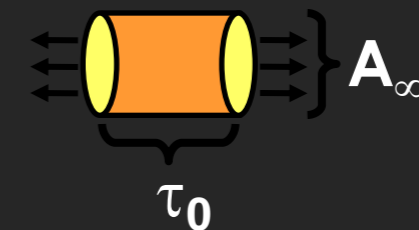


# Energy density in heavy ion collisions

Initial energy density :



Longitudinal QGP expansion :



Bjorken formula :

$$\epsilon_{Bj} = \frac{dE_t}{dy} \frac{1}{A_{\perp} \tau_0}$$

facility	collision species	$\sqrt{s_{NN}}$ (GeV)	$\epsilon_{Bj} \times \tau_0$ (GeV/fm <sup>-3</sup> . fm/c)
AGS (BNL)	Au+Au	5	1,5
SPS (CERN)	Pb+Pb	17	3,9
RHIC (BNL)	Au+Au	200	5,5
LHC (CERN)	Pb+Pb	5500	10