

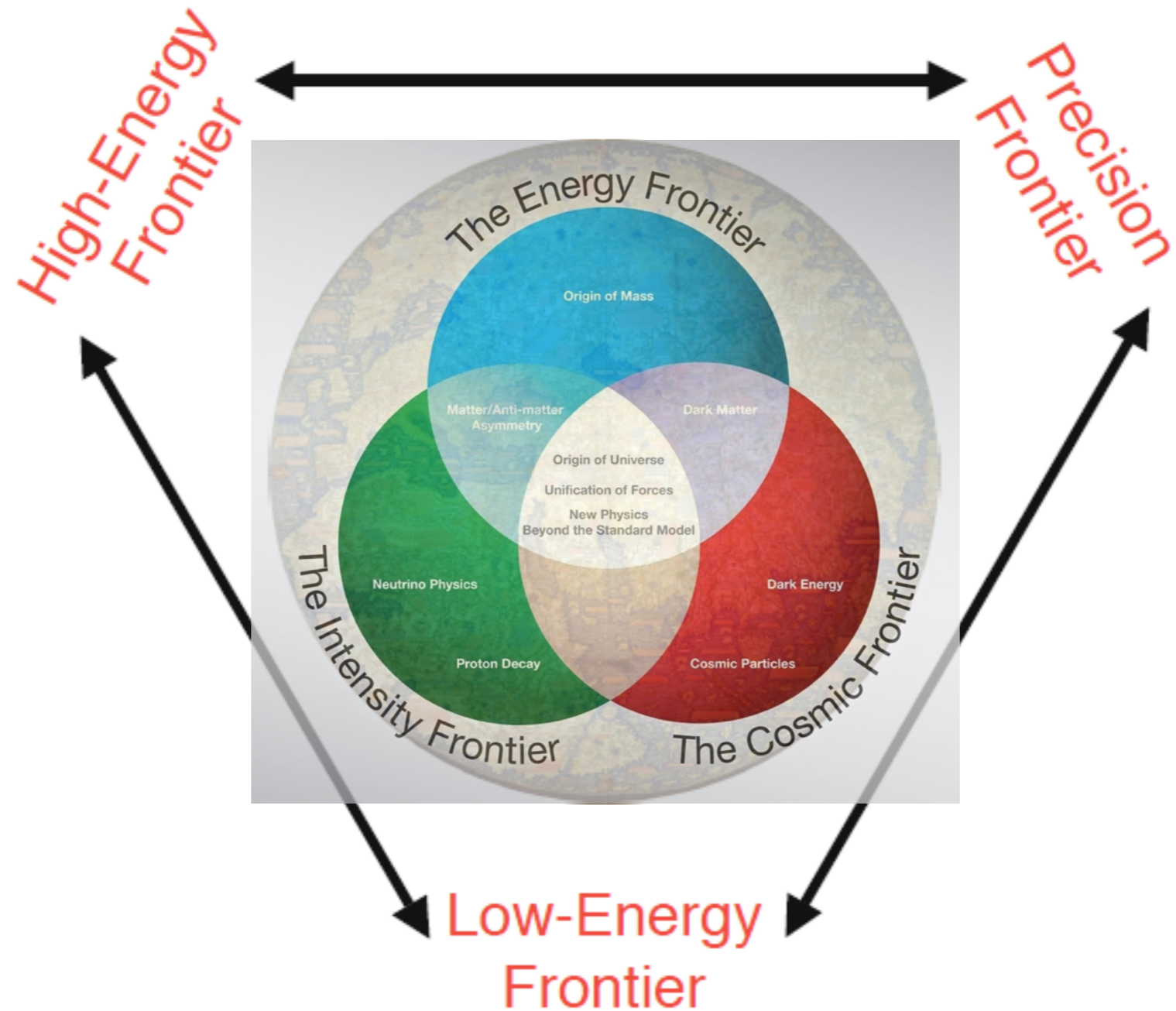
# NUCLEON STRUCTURE AT LOW-ENERGY / PRECISION FRONTIER

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**University of Mainz, Germany**



# Frontiers of subatomic physics



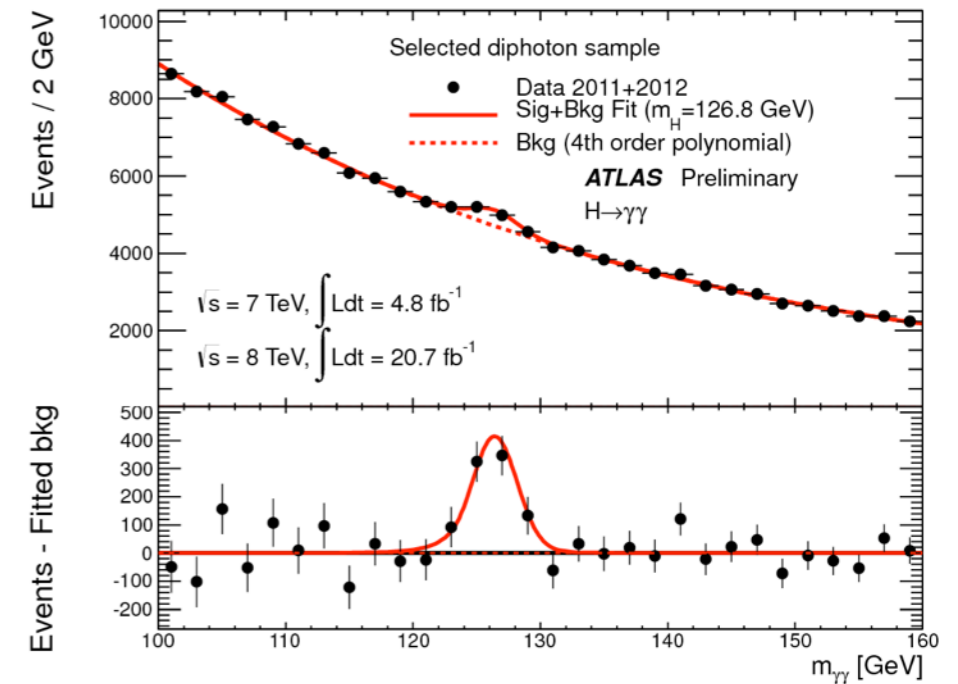
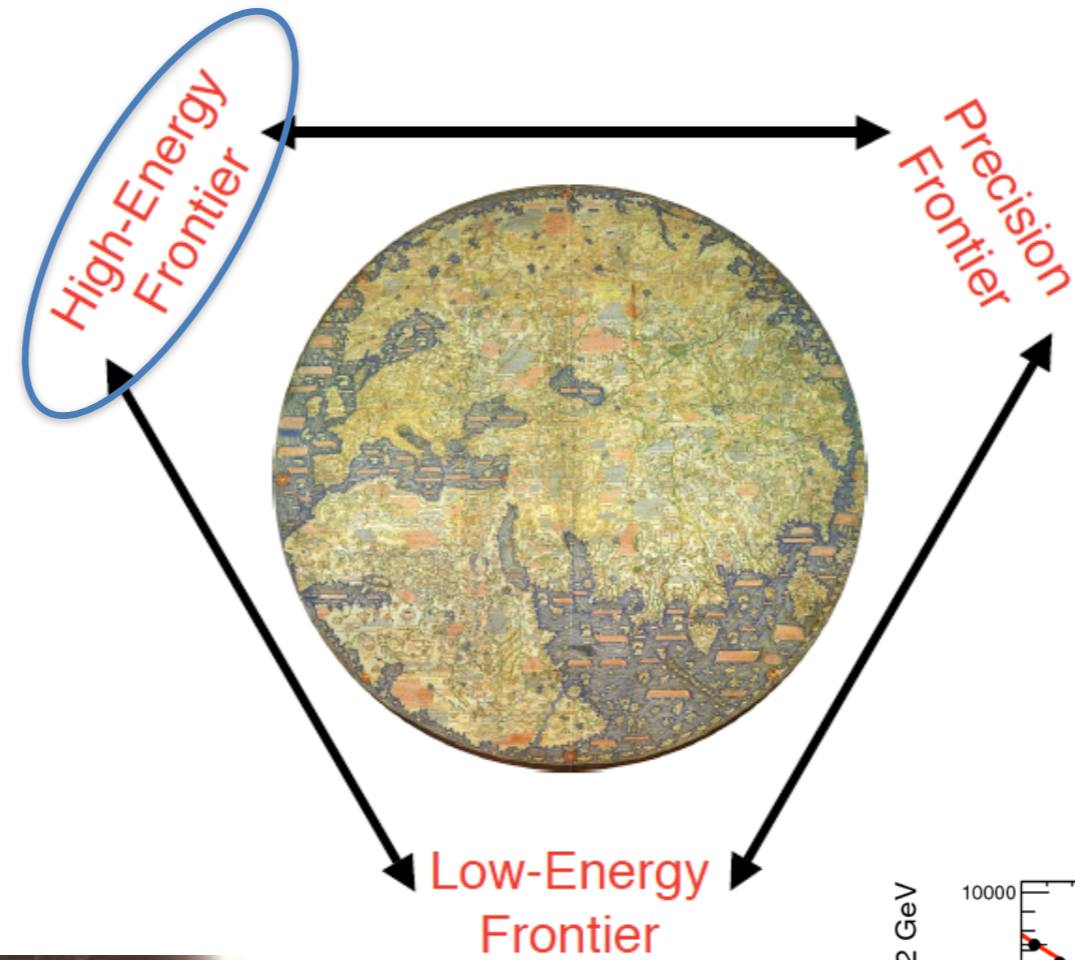


# High energy easy to identify

## LHC

Higgs boson discovery  
Production of new  
particles and stuff

Next: ILC

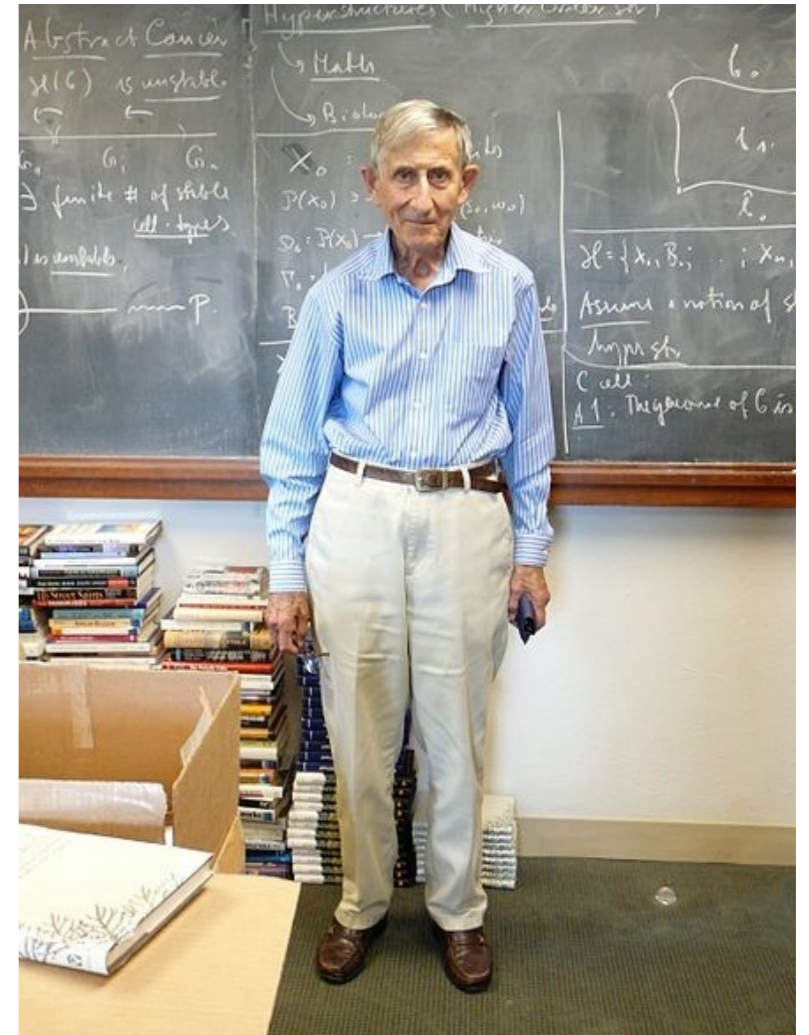


# Breaking through frontiers

**Freeman Dyson** on 16 discoveries awarded the Nobel Prize between 1945 and 2008:

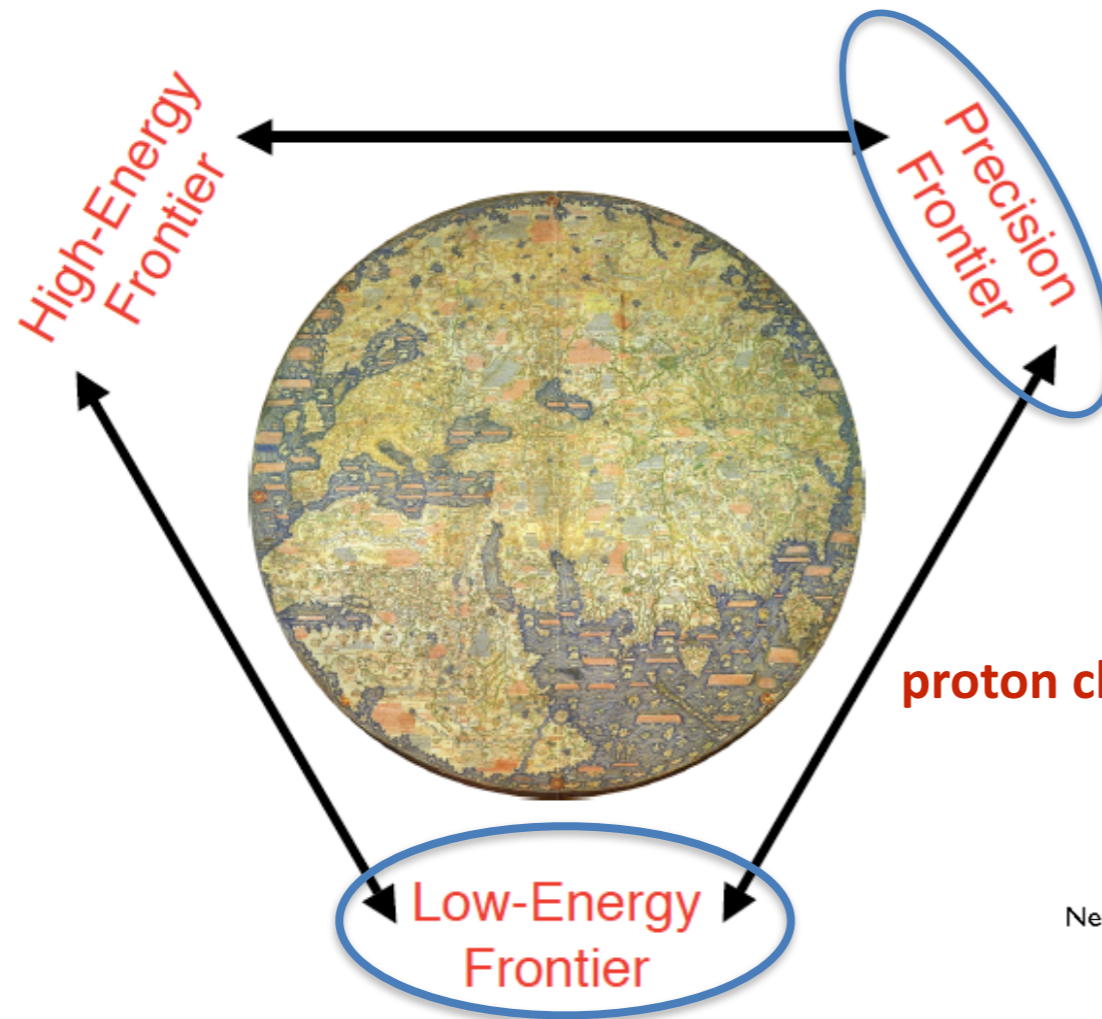
*“four discoveries on the energy frontier, four on the rarity frontier, eight on the accuracy frontier. Only a quarter of the discoveries were made on the energy frontier, while half of them were made on the accuracy frontier. **For making important discoveries, high accuracy was more useful than high energy.**”*

(Freeman Dyson, review of *The Lightness of Being*, F. Wilczek, *The New York Review of Books*, April 2009)





# Precision frontier and the stumbling stone



Testing SM at low energies, quantum loop corrections

$$\sin^2\theta_W$$

EDM

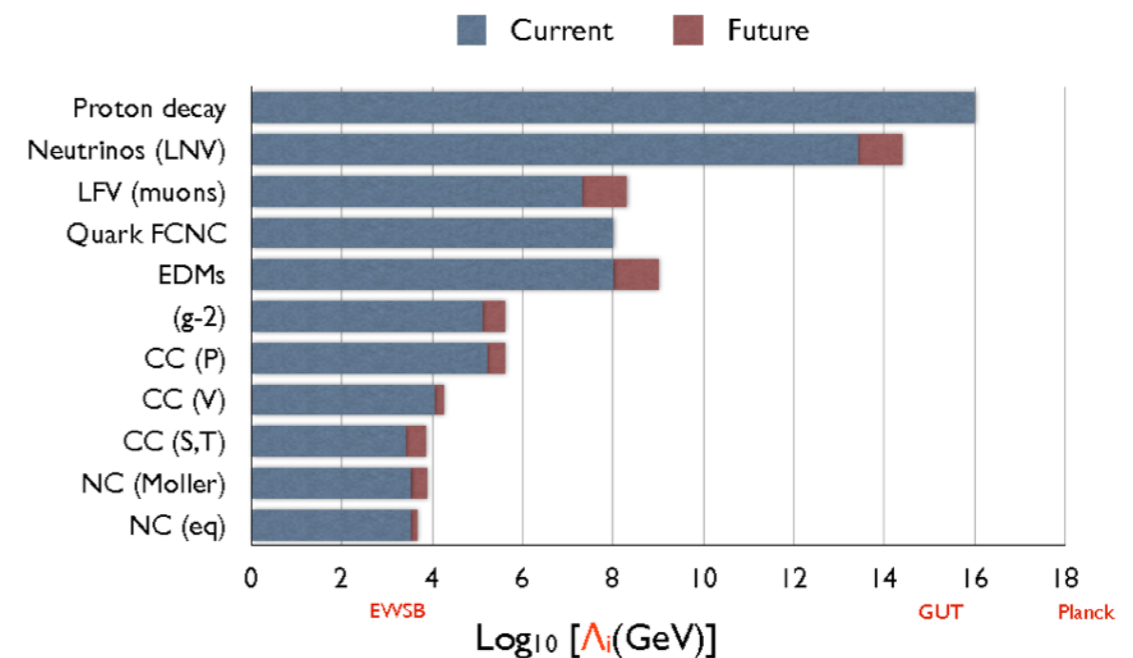
Flavor physics

Atomic tests

$$(g-2)_\mu$$

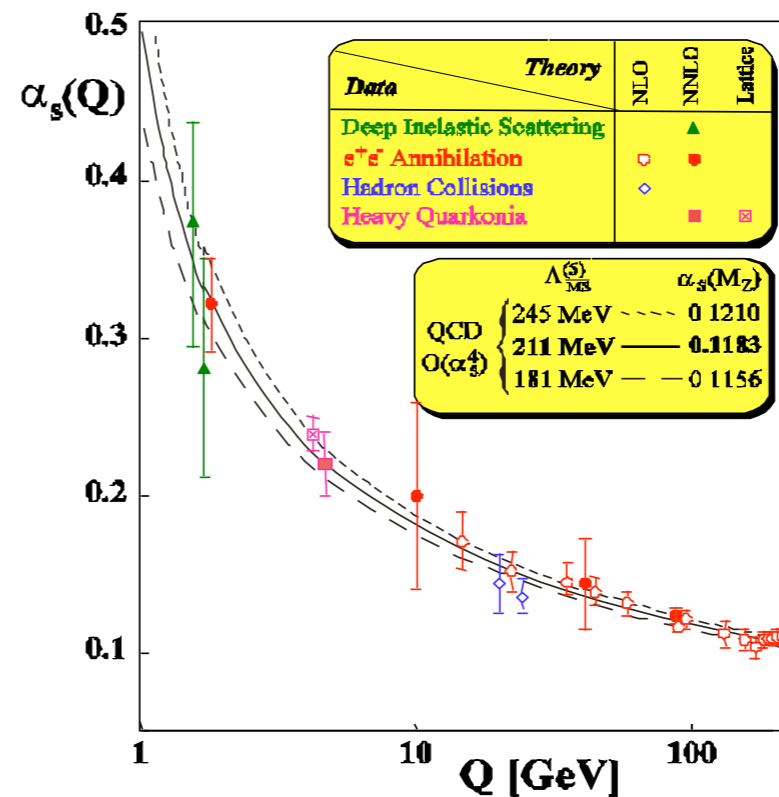
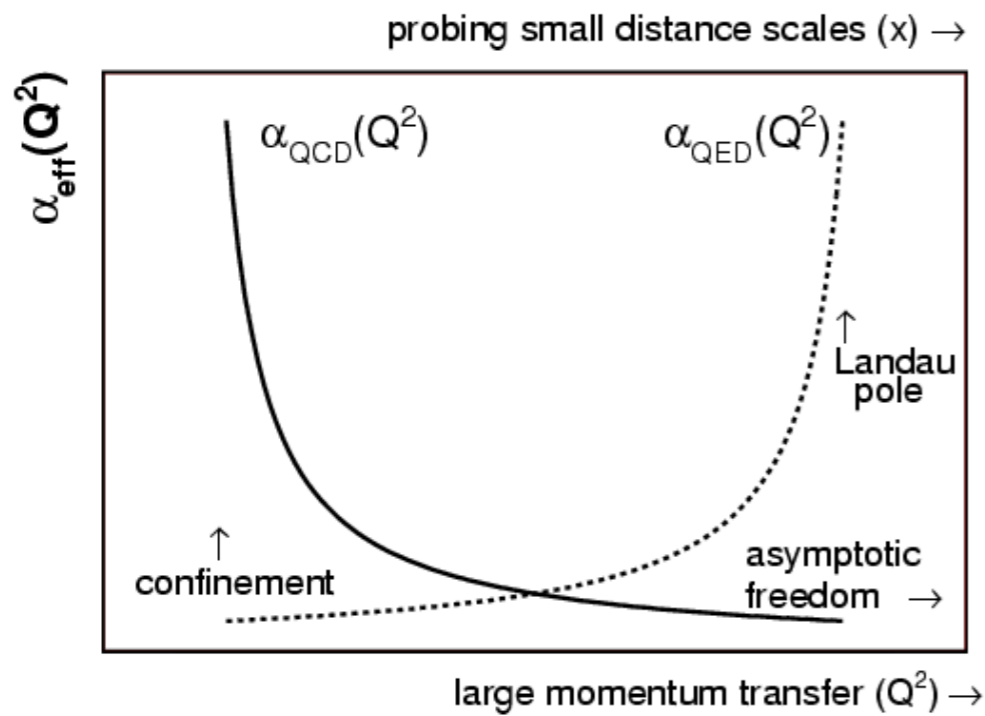
proton charge radius puzzle

**Low-Energy Strong Interaction:**  
non-perturbative QCD  
structure info (FFs, GPDs, PDFs)



[from Cigriliano & Ramsey-Musolf, Ann.Rev.Mod.Sci (2013) ]

# QCD coupling



For  $Q^2 \rightarrow \infty$ ,  $\alpha_s \rightarrow 0$  : asymptotic freedom

For  $Q \sim \Lambda_{QCD}$  non-perturbative phenomena:  
 color confinement,  
 spontaneous chiral symmetry breaking,  
 generation of nucleon mass, ...



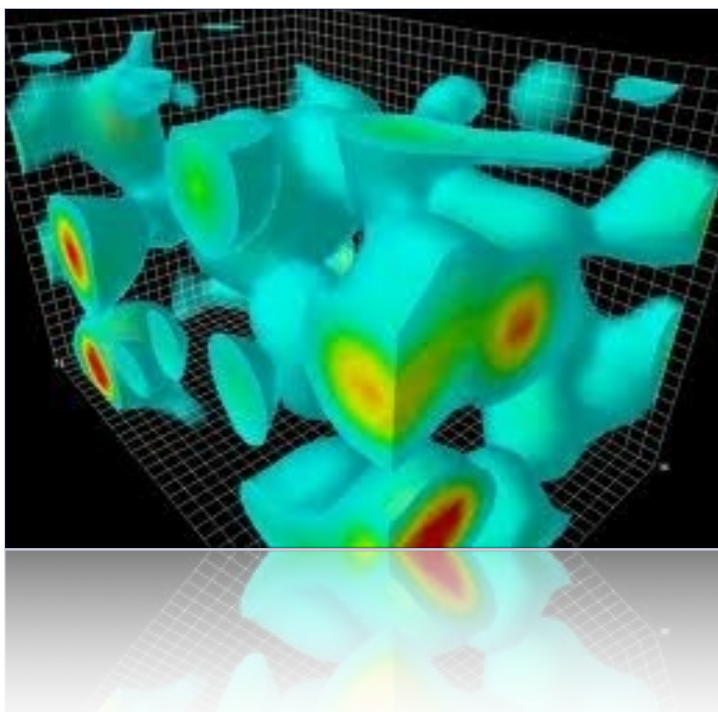


# QFTs of low-energy strong interaction

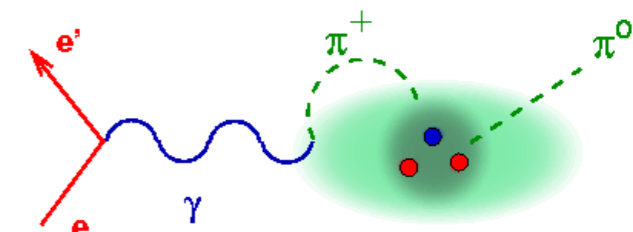
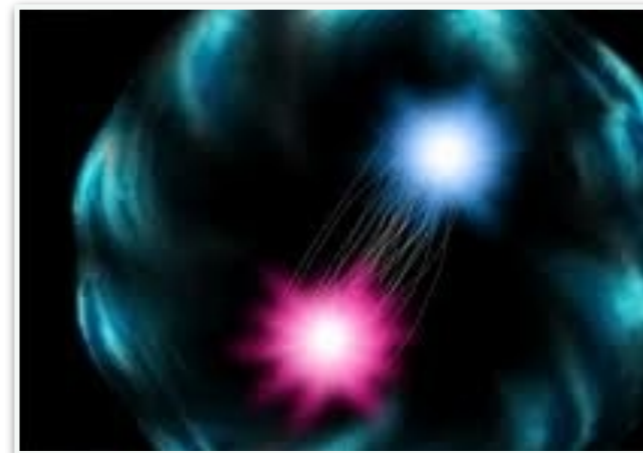
... turning the tumble stone into stepping stones



Lattice QCD



Chiral perturbation theory (ChPT),  
a.k.a Chiral Effective-Field Theory (ChEFT)



# ChPT basic facts

- ◆ **S. Weinberg, *Phenomenological Lagrangians*, Physica (1979):**  
aimed to obtain quantum corrections to PCAC (LETs + chiral symmetry),  
derived the **Effective Field Theory** framework
- ◆ **Gasser & Leutwyler (1984, 1985) worked out ChPT in the meson sector.**
- ◆ **‘Chiral’ and ‘Perturbative’ go together:**  
pions are Goldstone bosons of spontaneous ChSB,  
interaction goes with powers of energy, vanishes at  $E=0$  in the chiral limit.  
perturbative expansion in energy and pion mass (but not a series expansion!)
- ◆ **Most general Lagrangian (allowed by symmetries), hence infinitely many constants (LECs) parametrising the short-range physics.**
- ◆ **Predictive provided: *Hierarchy of scales and Naturalness***



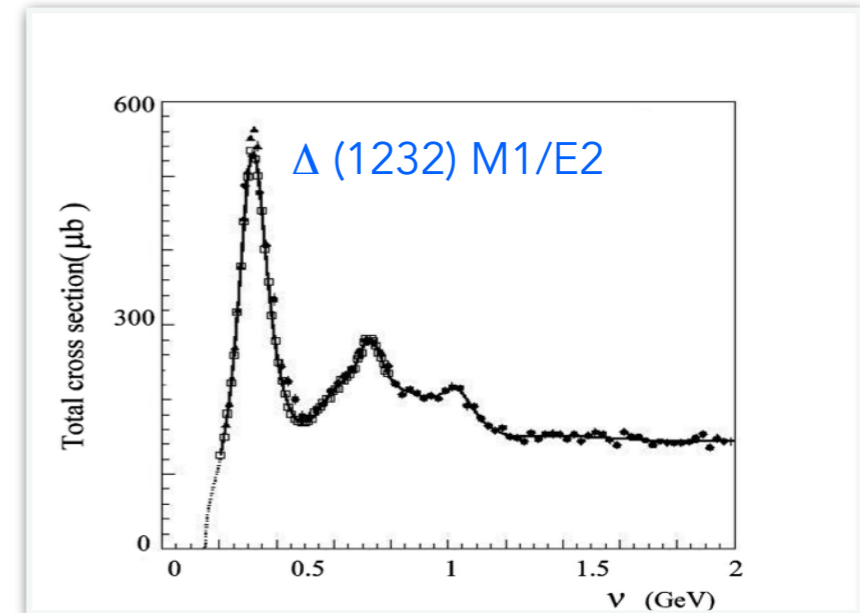
# Baryon ChPT

## Not just the pion cloud: Delta(1232) excitation

Jenkins & Manohar, PLB (1991)

Hemmert, Holstein, Kambor, JPhysG (1998)

V.P. & Phillips, PRC (2003)



$E$  (GeV)

1

$4\pi f_\pi$

$M_N$

$m_\rho$

0.3

$M_\Delta - M_N$

0.1

$m_\pi$

- The 1st nucleon excitation — Delta(1232) is within reach of chiral perturbation theory (293 MeV excitation energy is a light scale)
- Include into the chiral effective Lagrangian as explicit dof
- Power-counting for Delta contributions (SSE, “delta-counting”) depends on what chiral order is assigned to the excitation scale.

# Example: Nucleon mass

$$\mathcal{L} = \sum_k \mathcal{L}^{(k)}, \quad k = \# \text{ of pion derivatives and masses}$$

$$\begin{aligned} \mathcal{L}_{\pi N}^{(1)} &= \bar{N}(i\not{D} - \overset{\circ}{M}_N + \overset{\circ}{g}_A a_\mu \gamma^\mu \gamma_5)N \\ &= \bar{N}\left(i\not{\partial} - \overset{\circ}{M}_N + \frac{\overset{\circ}{g}_A}{2f_\pi} (\partial_\mu \pi) \gamma^\mu \gamma_5\right)N + O(\pi^2) \end{aligned}$$

$$\mathcal{L}_{\pi N}^{(2)} = 4\overset{\circ}{c}_{1N} m_\pi^2 \bar{N} N + \dots$$

**Power-counting:**

$$n = \sum_k k V_k + 4L - 2N_\pi - N_N$$

- $V_k$  # of vertices from  $\mathcal{L}^{(k)}$
- $L$  # of Loops
- $N_\pi$  # of internal pions
- $N_N$  # of internal nucleons

LECs

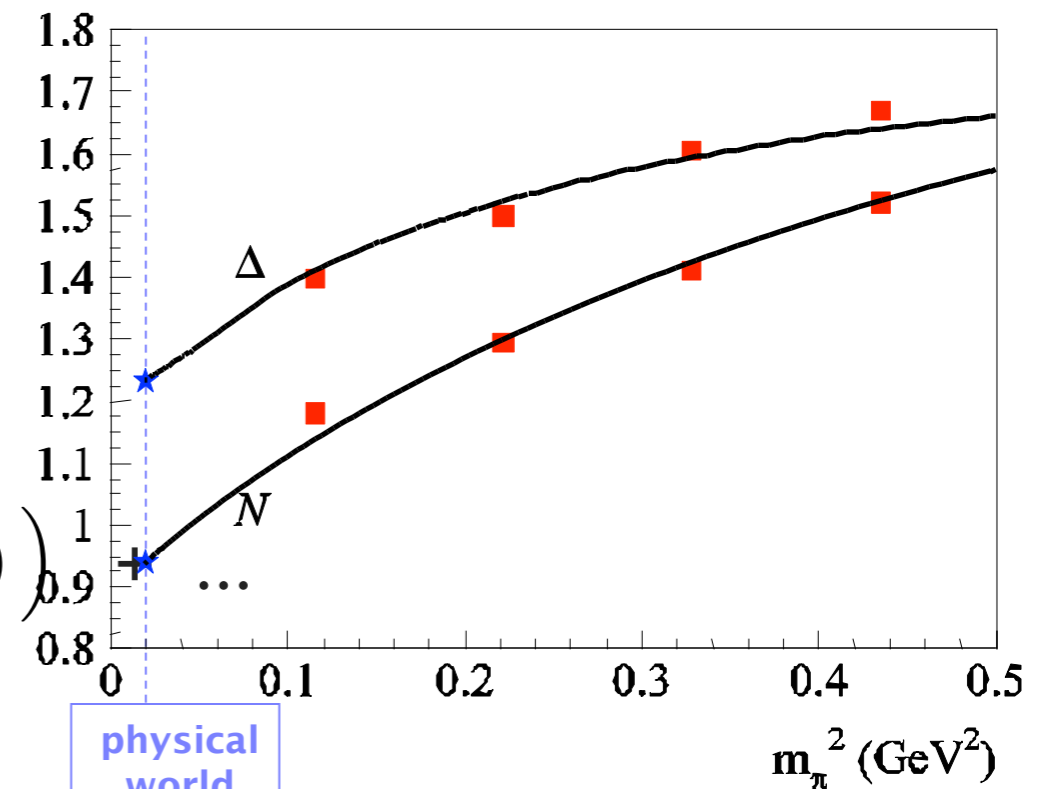
$$M_N = \overset{\circ}{M}_N - 4\overset{\circ}{c}_{1N} m_\pi^2 - \text{[Diagram: two nucleons connected by a dashed pion line]} + \dots$$

$O(p^3)$

$k = 1, V_k = 2, L = 1, N_\pi = 1, N_N = 1$

$$M_N = \overset{\circ}{M}_N - 4\overset{\circ}{c}_{1N} m_\pi^2 - \frac{g_A^2}{(4\pi f_\pi)^2} \left( \frac{3\pi}{2} m_\pi^3 + O(m_\pi^4) \right)$$

prediction of ChPT



[V.P. & Vanderhaeghen, PLB (2006)]



# Heavy-baryon ChPT? No!

LECs

$$M_N = \overset{\circ}{M}_N - 4 \overset{\circ}{c}_{1N} m_\pi^2 - \text{---} \overset{\frown}{\text{---}} \text{---} + \dots$$

$O(p^3)$

$$\text{---} \overset{\frown}{\text{---}} \text{---} = \frac{3g_A^2}{2(4\pi f_\pi)^2} \left\{ -M_N^3 L + M_N (1-L) m_\pi^2 - m_\pi^3 \left( \sqrt{1 - m_\pi^2/4M_N^2} \arccos \frac{m_\pi}{2M_N} + \frac{m_\pi^4}{4M_N} \ln \frac{m_\pi^2}{M_N^2} \right) \right\}$$

where  $L = \frac{1}{\epsilon} + \dots$  contains the UV-divergence, removed in MS-bar:  $L = 0$   
 remaining  $m_\pi^2$  "complicates life a lot" [GSS88]. **Violation of power counting?!!**

Gasser, Sainio & Svarc, NPB (1988); ...

Led to **Heavy-Baryon ChPT** [Jenkins & Manohar, PLB (1991)] which for a decade was considered as the only consistent formulation.

**Drawback: not working** — 1. removes  $m_\pi^2$  in dimreg but not in cutoff schemes, 2. demotes important contributions to "higher-orders"

**Fortunately, HB not needed:**  $m_\pi^2$  term removed by renormalization of the LEC.

Japaridze & Gegelia (1999), published in (2003)!

# Relevance to low-energy/precision frontiers

ChPT gives predictions\*, i.e. free-parameter free results, for:

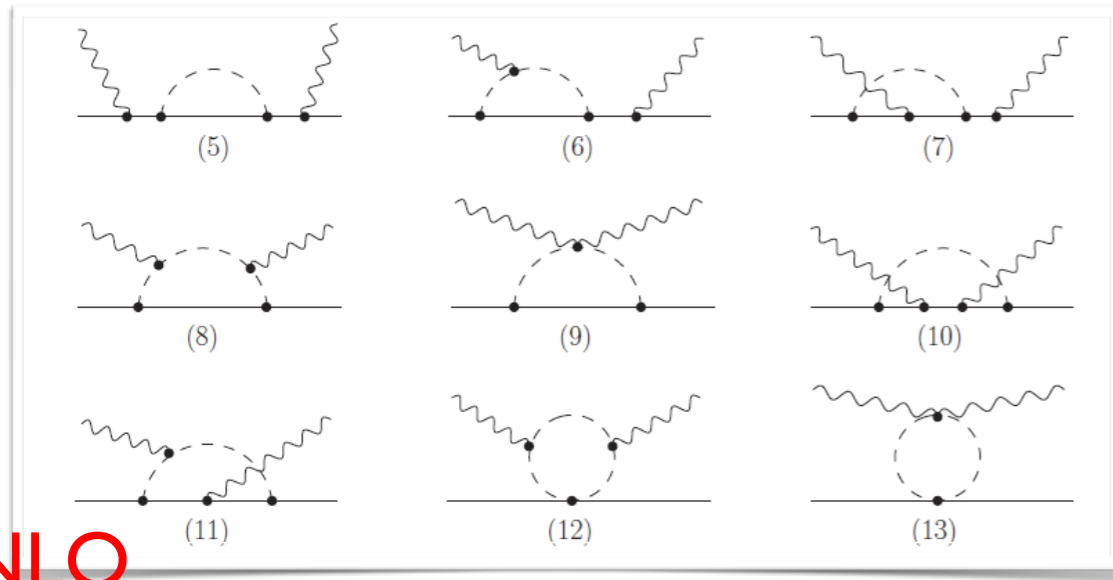
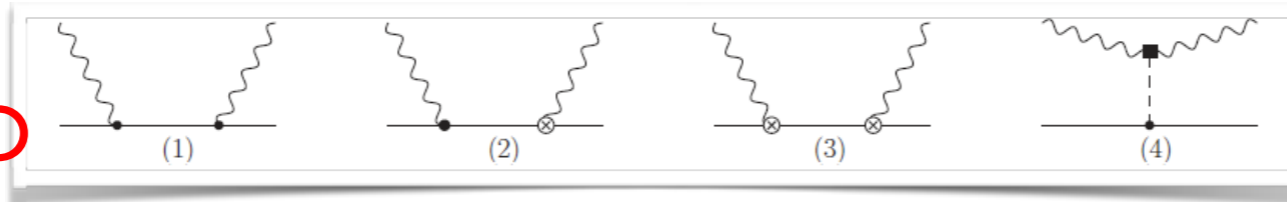
**1. Nucleon polarizabilities**

**2. Nucleon structure effects in hydrogen Lamb shift **beyond the charge radius****

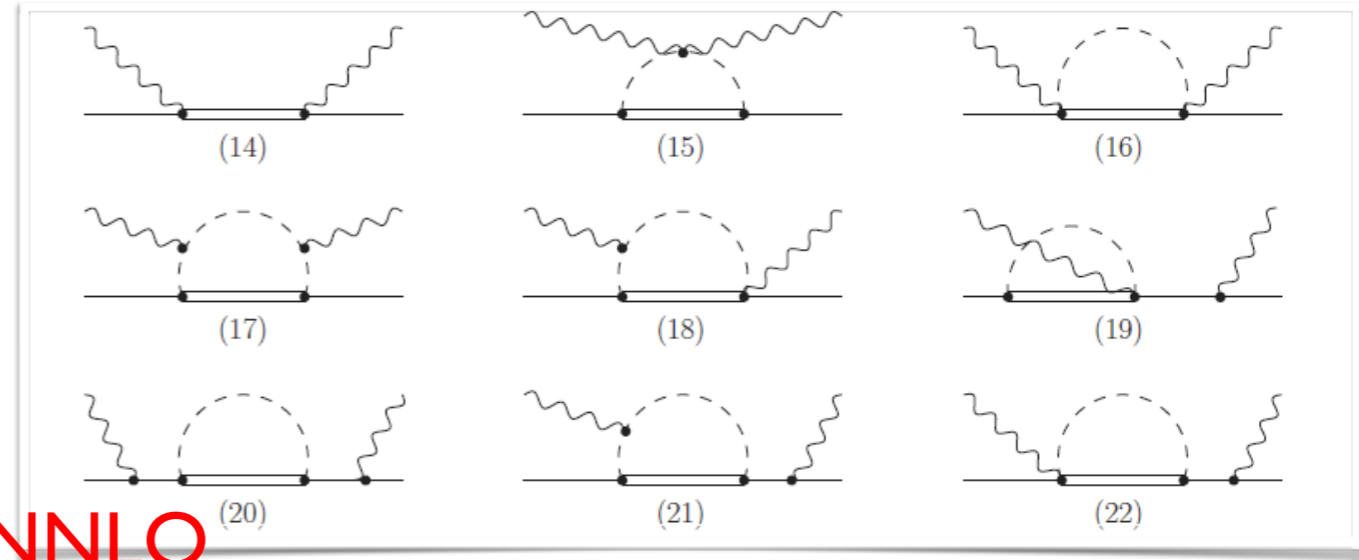
\*Predictions of HBChPT differ from BChPT

# ChPT of Compton scattering off protons

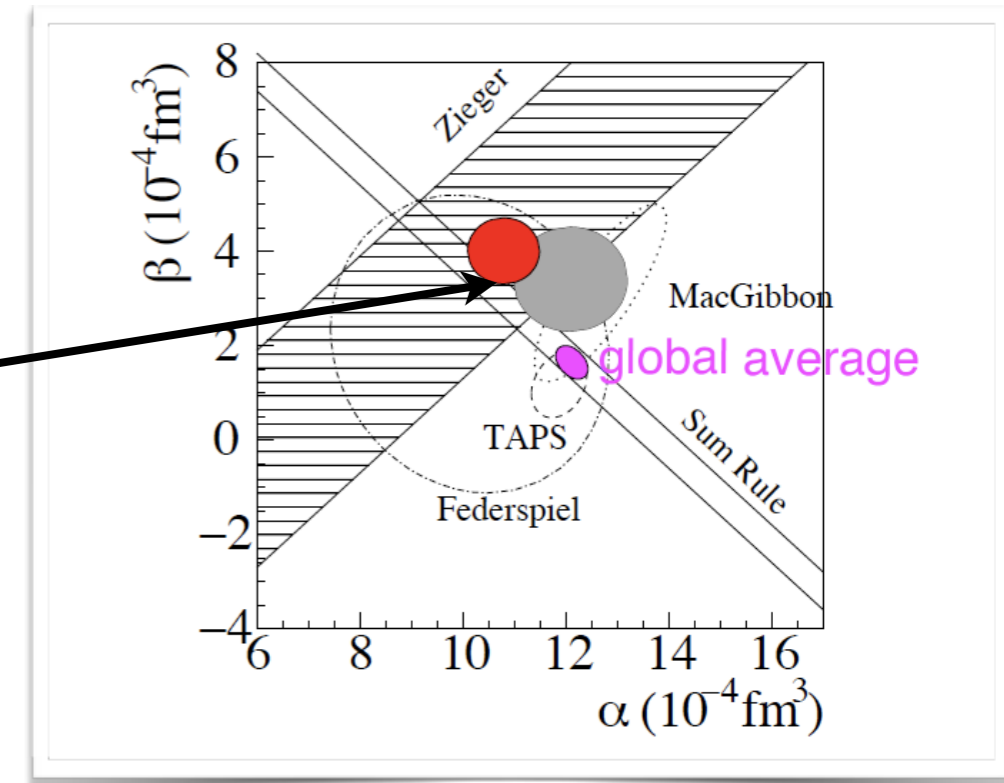
LO



NNLO



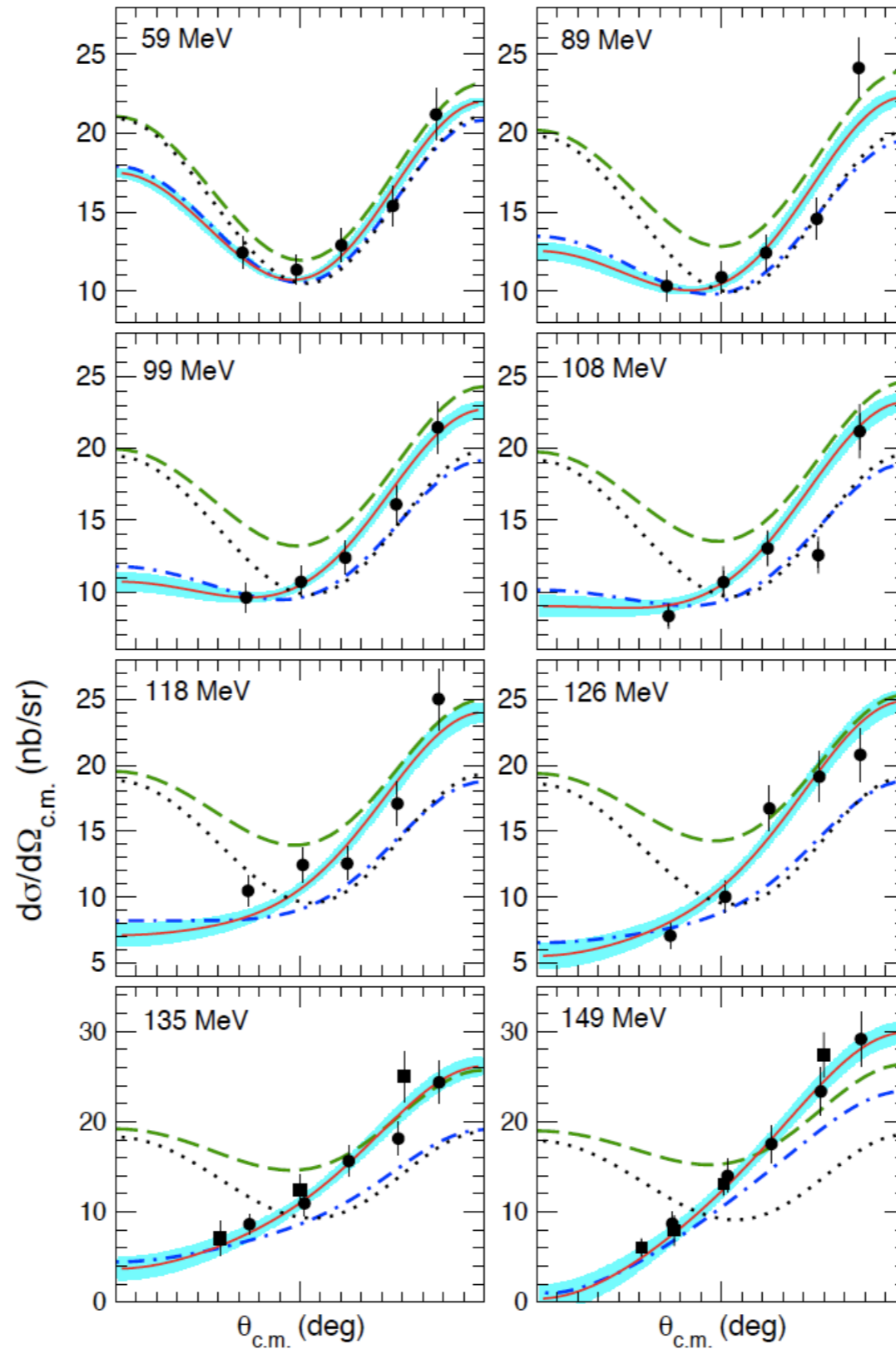
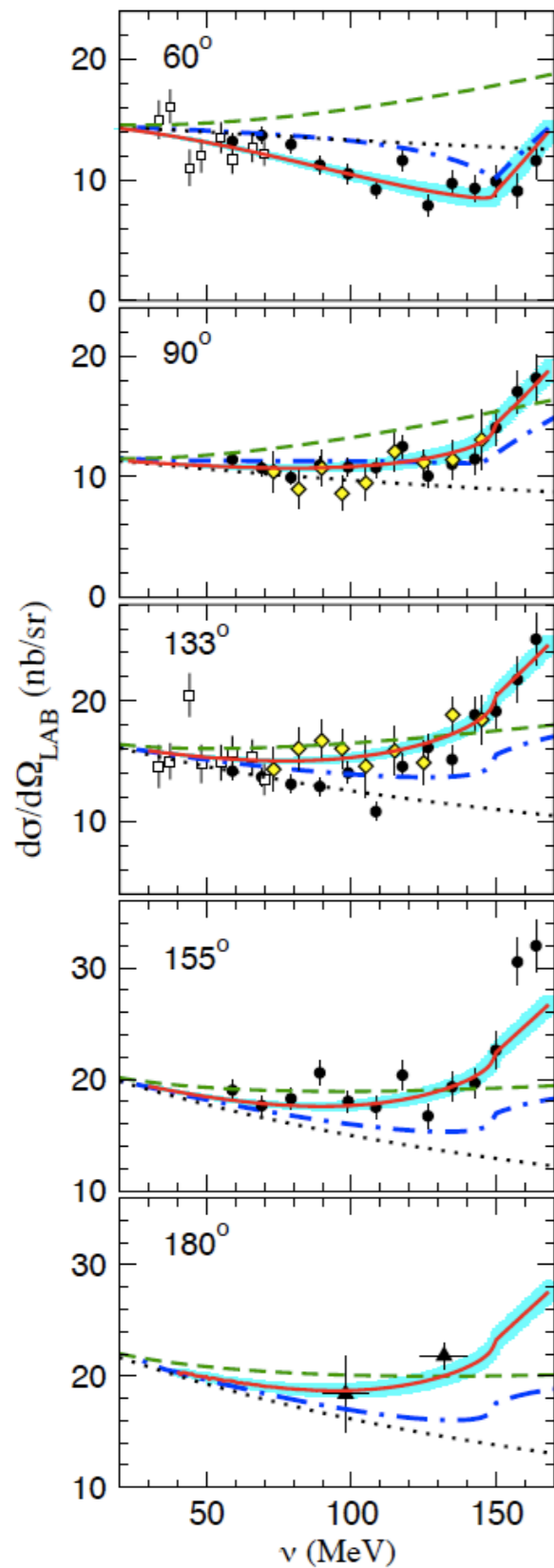
$\mathcal{O}(p^2)$	$\frac{e^2}{4\pi} = \frac{1}{137}, M_N = 938.3 \text{ MeV}, \hbar c = 197 \text{ MeV}\cdot\text{fm}$
$\mathcal{O}(p^3)$	$g_A = 1.267, f_\pi = 92.4 \text{ MeV}, m_\pi = 139 \text{ MeV}, m_{\pi^0} = 136 \text{ MeV}, \kappa_p = 1.79$
$\mathcal{O}(p^4/\Delta)$	$M_\Delta = 1232 \text{ MeV}, h_A = 2.85, g_M = 2.97, g_E = -1.0$
$\mathcal{O}(p^4)$	$\alpha_0, \beta_0 = \pm \frac{e^2}{4\pi M_N^3}$ <span style="margin-left: 20px;">size of the red blob</span>



Lensky & V.P., EPJC (2010)



# Unpolarized cross sections



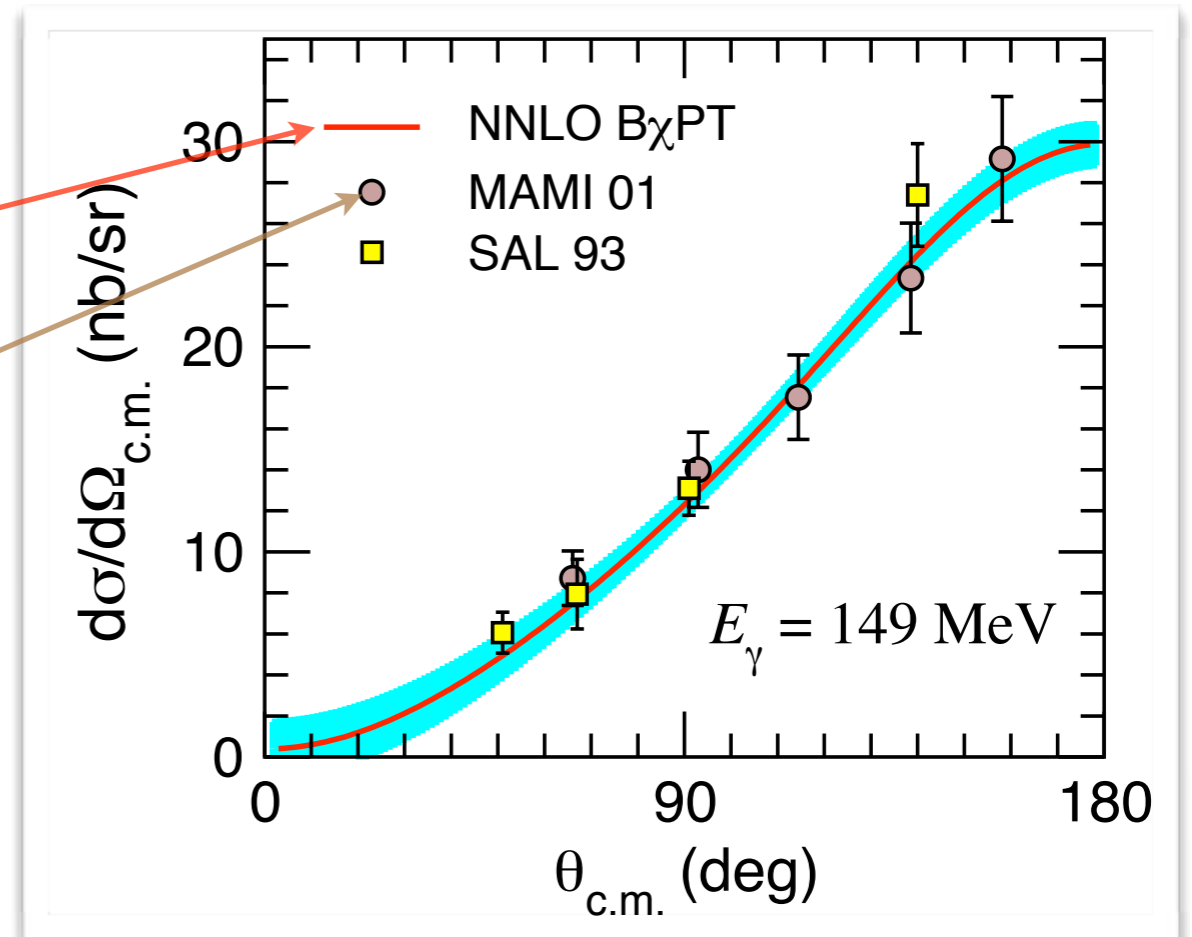
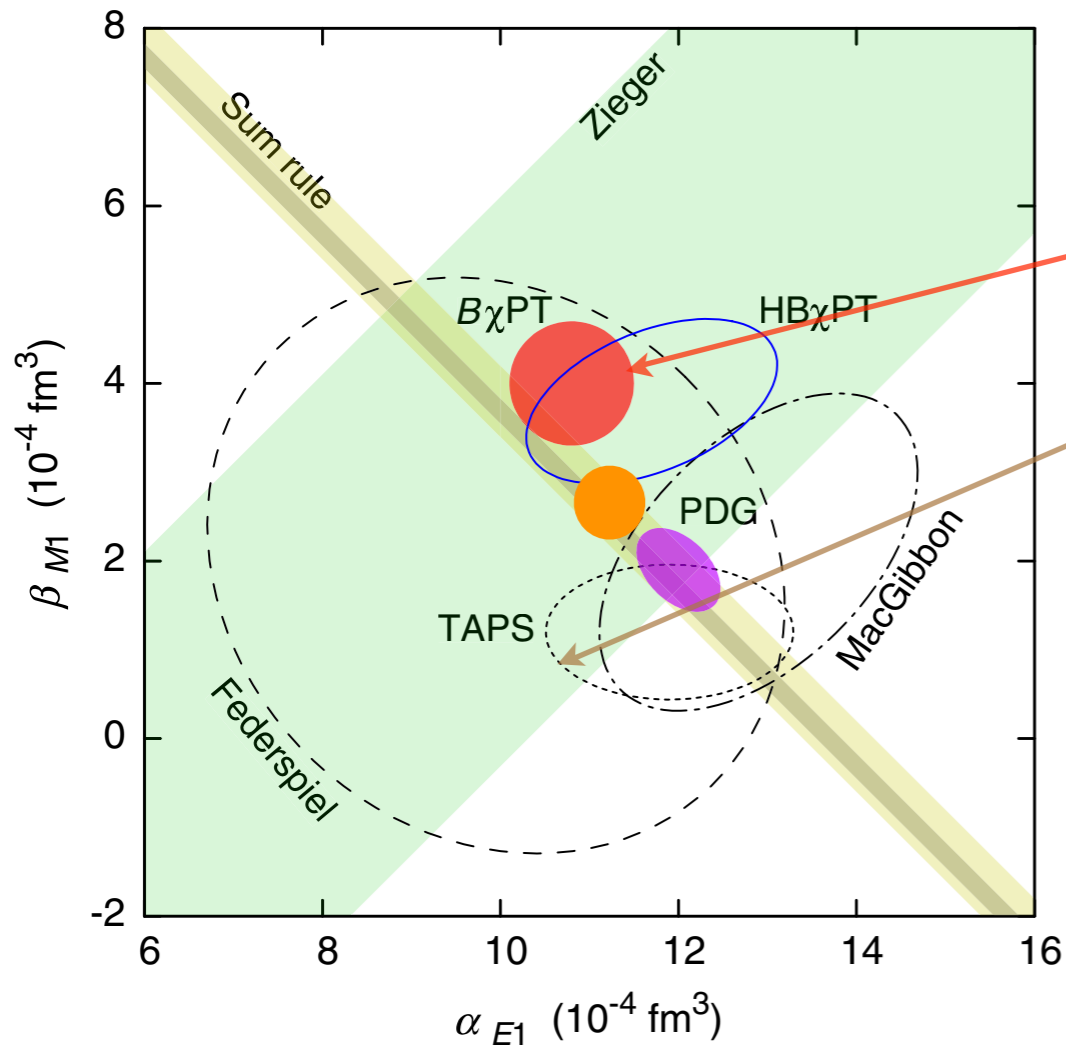
Data points:  
MAMI/TAPS  
(2001)  
SAL (1993)  
Illinois (1991)

Curves:

- ..... Klein-Nishina
- - - - Born + WZW
- . - . + p-qube
- Total NNLO

Lensky & V.P., EPJC (2010)

# Proton polarizabilities



$$\beta_{M1} = (1.9 \pm 0.5) \times 10^{-4} \text{ fm}^3 \text{ [PDG]}$$

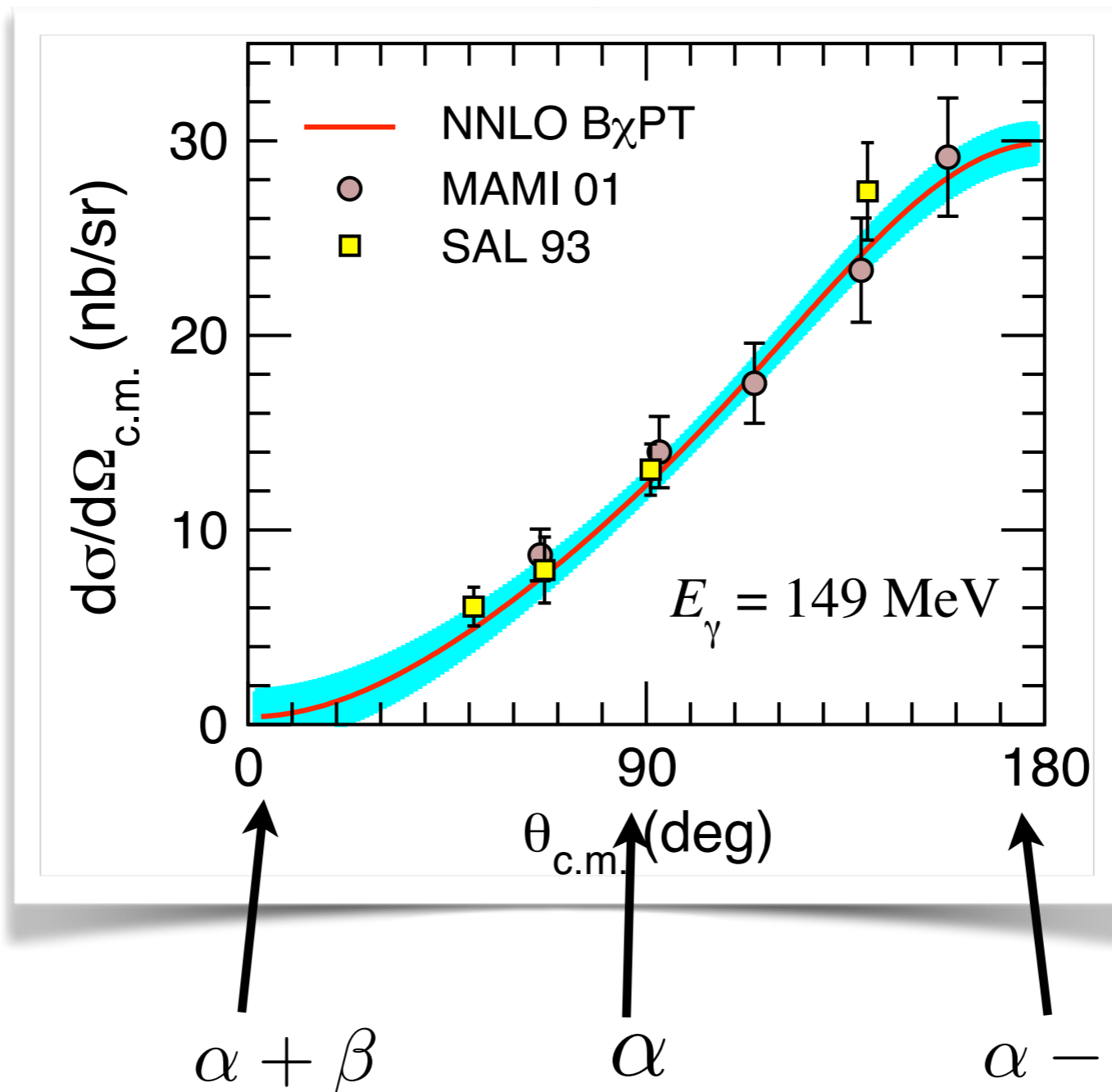
$$\beta_{M1} = (4.0 \pm 0.7) \times 10^{-4} \text{ fm}^3 \text{ [BChPT@NNLO]}$$

BChPT - Lensky & V.P., EPJC(2010)  
 HBChPT - Griesshammer, McGovern,  
 Phillips, EPJA (2013)

PDG adjusted values  
 from 2012 edition (**purple**) to  
 2013 on-line edition (**orange**)

# Extracting polarizabilities from angular dep.

$$\frac{d\sigma^{(\text{NB})}}{d\Omega} = -2\pi Z^2 \frac{\alpha}{M} \left(\frac{\nu'}{\nu}\right)^2 \nu\nu' [2\alpha_{E1} (1 + \cos^2 \theta) + 4\beta_{M1} \cos \theta] + O(\nu^4)$$

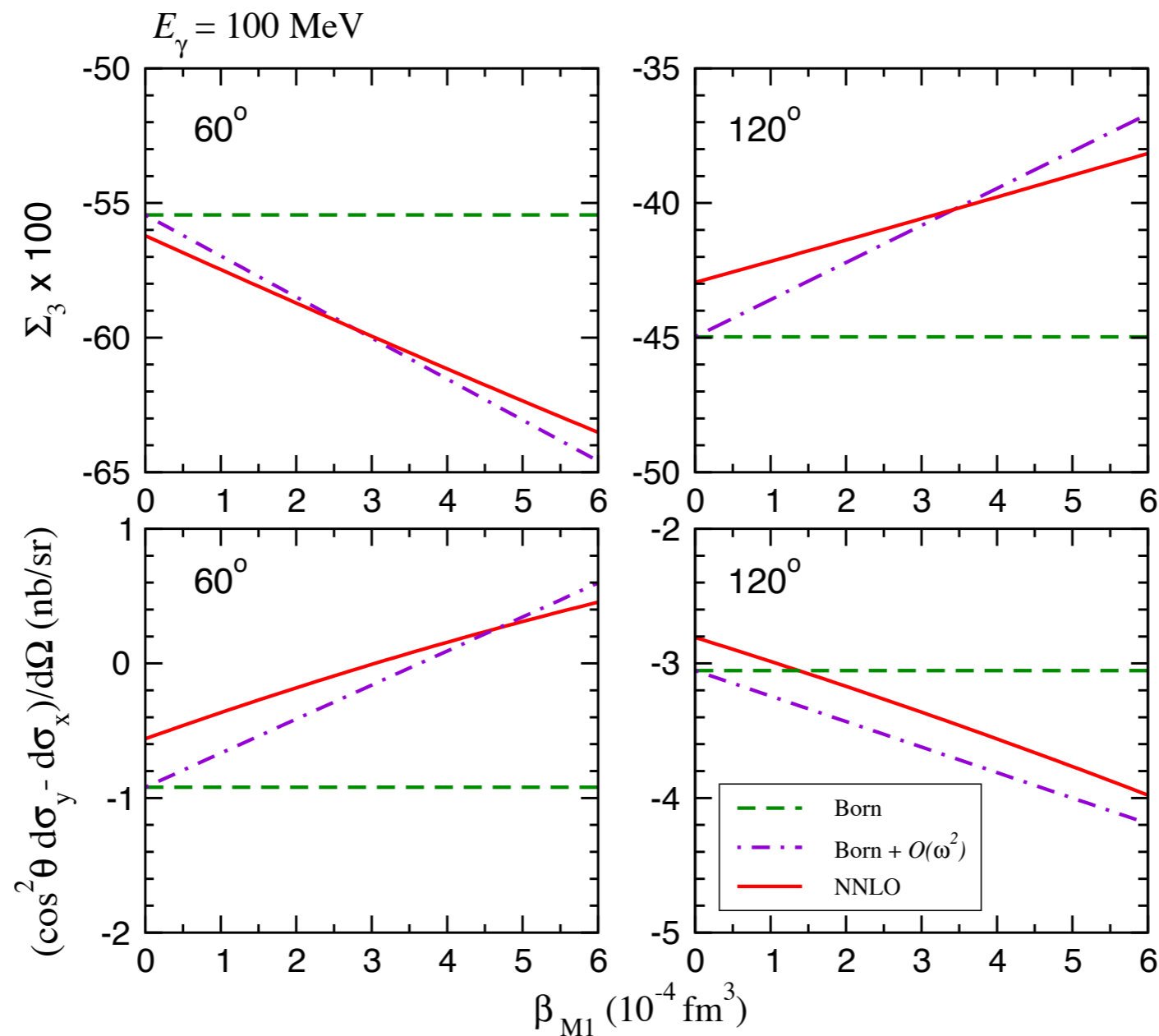




# From linear beam asymmetry

$$\Sigma_3 \equiv \frac{d\sigma_{||} - d\sigma_{\perp}}{d\sigma_{||} + d\sigma_{\perp}} \stackrel{\text{LEX}}{=} \Sigma_3^{(\text{Born})} - \frac{4\beta_{M1}}{Z^2\alpha_{em}} \frac{\cos\theta \sin^2\theta}{(1 + \cos^2\theta)^2} \omega^2 + O(\omega^4)$$

Krupina & V.P., PRL (2013)



# New Mainz data for Compton beam asymmetry

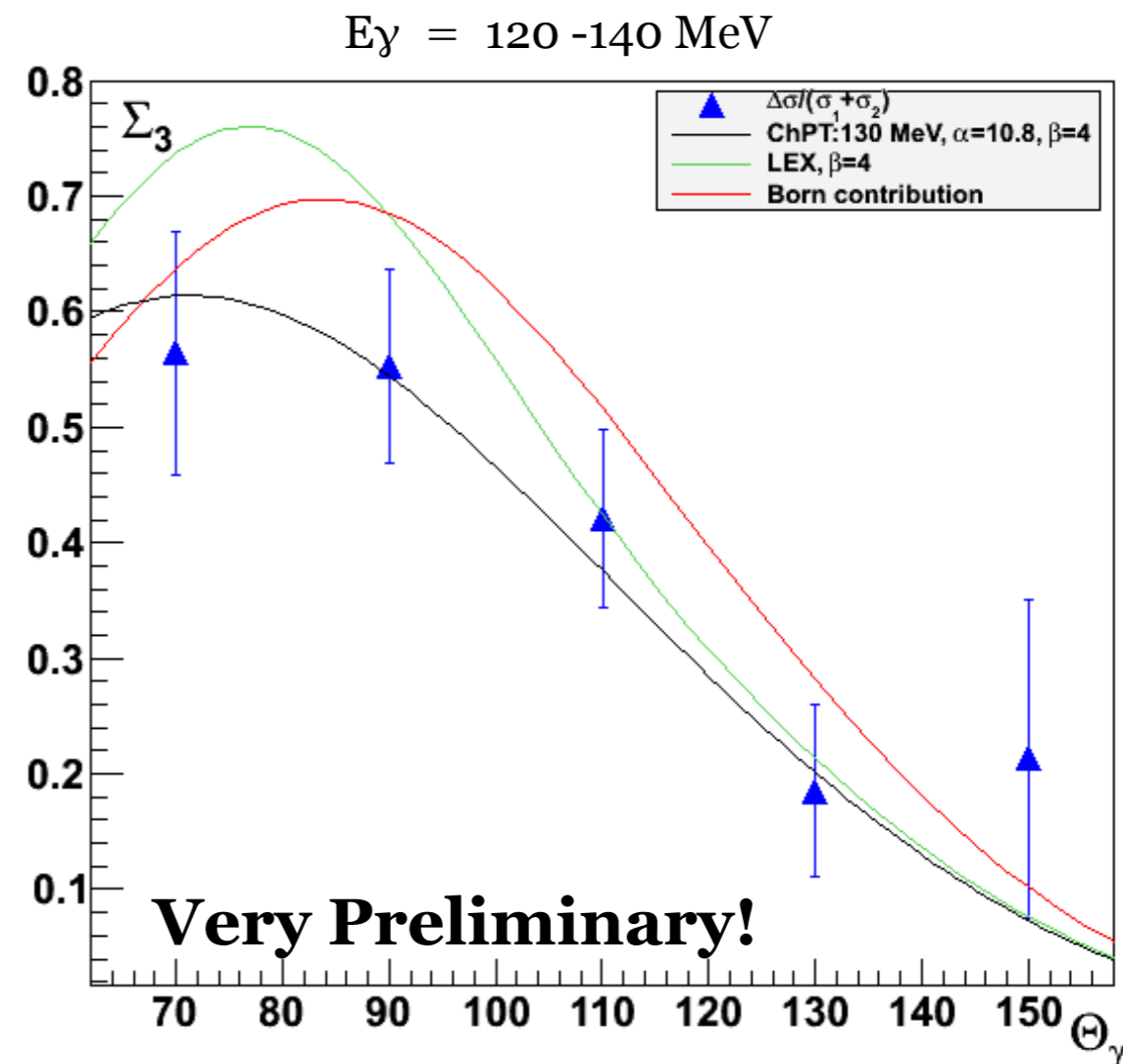
Data taken: 28.05. – 17.06.2013, 327 h

V. Sokhoyan, E. Downie et al.  
[A2 Coll.]

first data on this  
observable below pion  
production threshold!

**better precision needed!!**

## Beam asymmetry $\Sigma_3$ : Preliminary results



# Predictions of HBChPT vs BChPT

## HBChPT@LO

Bernard, Keiser, Meissner  
Int J Mod Phys (1995)

$$\alpha_p = \alpha_n = \frac{5\pi\alpha}{6m_\pi} \left( \frac{g_A}{4\pi f_\pi} \right)^2 = 12.2 \times 10^{-4} \text{ fm}^3,$$

$$\beta_p = \beta_n = \frac{\pi\alpha}{12m_\pi} \left( \frac{g_A}{4\pi f_\pi} \right)^2 = 1.2 \times 10^{-4} \text{ fm}^3,$$

paramagnetic

diamagnetic

$$\mu = m_\pi / M_N$$

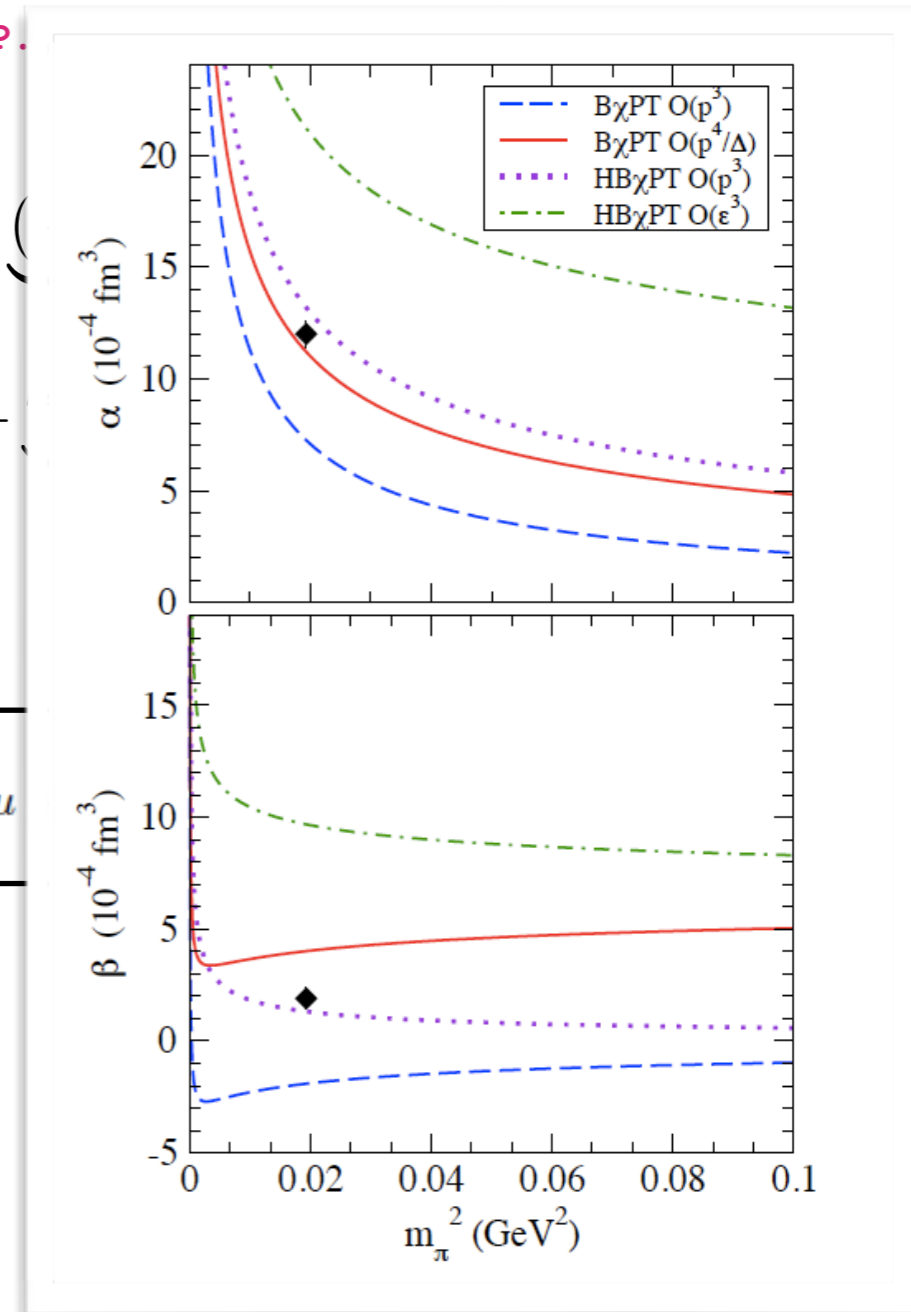
$$\beta = \frac{e^2 g_A^2}{192\pi^3 F^2 M_N} \left[ \frac{\pi}{4\mu} + 18 \log \mu + \frac{63}{2} - \frac{981\pi\mu}{32} - (100 \log \mu) \right]$$

## BChPT@NLO

Lensky & V.P.

$$\alpha = \underbrace{6.8}_{\mathcal{O}(p^3)} + \dots$$

$$\beta = \underbrace{-1.8}_{\mathcal{O}(p^3)} + \dots$$



Lattice QCD data expected soon

## HBChPT@NLO:

Griesshammer & Hemmert (2004)

Griesshammer, McGovern, Phillips (2012)

The Delta contribution is accompanied by “promoted” LECs, hence not predictive



# Discoveries relevant to modern precision frontier

## *The Nobel Prize in Physics 1955*

### **Willis Eugene Lamb**

"for his discoveries concerning the fine structure of the hydrogen spectrum"

### **Polykarp Kusch**

"for his precision determination of the magnetic moment of the electron"

## *The Nobel Prize in Physics 1961*

### **Robert Hofstadter**

"for his pioneering studies of electron scattering in atomic nuclei and for his thereby achieved discoveries concerning the structure of the nucleons"

# The proton radii puzzle

**SERGIO LEONE**



CLINT EASTWOOD

ELI WALLACH

LEE VAN CLEEF

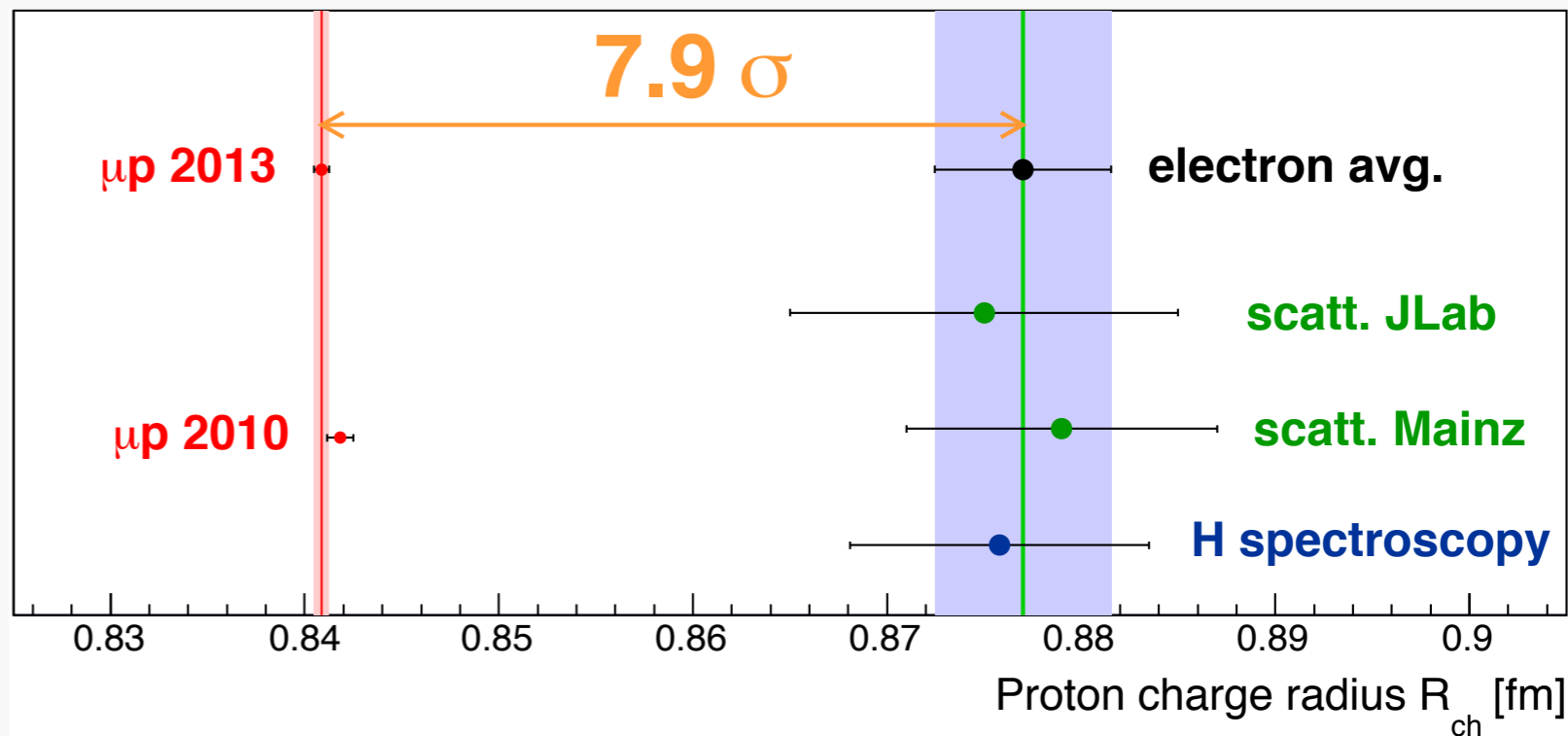
**THE UGLY**

**THE GOOD**

**AND THE BAD**

3 ways to the proton radius

e-p scattering  
 H precision laser spectroscopy  
 $\mu p$  laser spectroscopy



Pohl *et al.*, Nature 466, 213 (2010)

Antognini *et al.*, Science 339, 417 (2013)



# Proton size



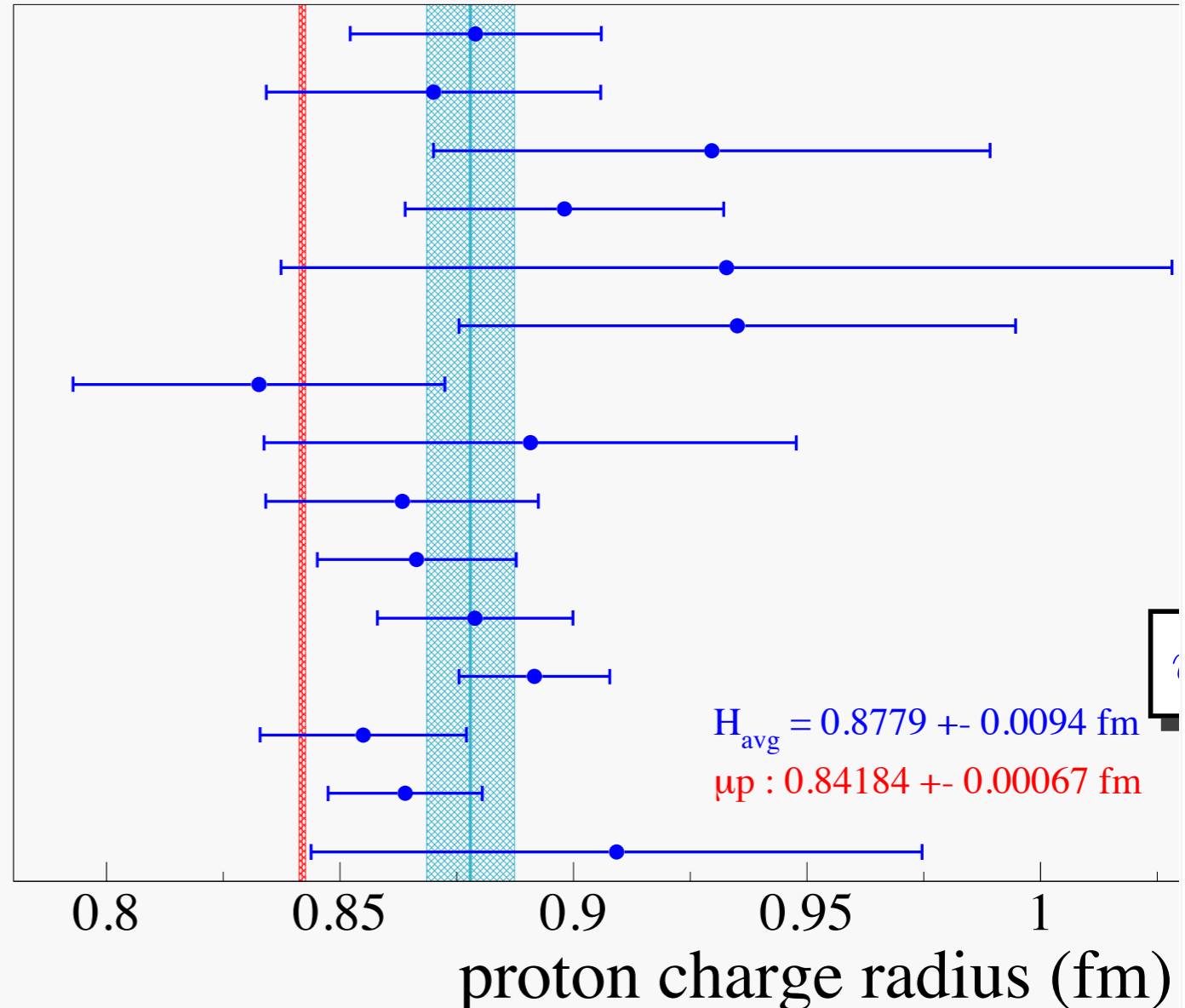
Pohl et al. (2010)  
Antognini et al. (2013)

$\mu\text{H}$ :  $r_E = 0.8409 \pm 0.0004 \text{ fm}$

$8\sigma$   
discrepancy

ep-data :  
CODATA 2010  $r_E = 0.8772 \pm 0.0046 \text{ fm}$   
Bernauer et al. (2010), (2013)  
Zhan et al. (2011)

- $2S_{1/2} - 2P_{1/2}$
- $2S_{1/2} - 2P_{3/2}$
- $2S_{1/2} - 2P_{1/2}$
- $1S-2S + 2S - 4S_{1/2}$
- $1S-2S + 2S - 4D_{5/2}$
- $1S-2S + 2S - 4P_{1/2}$
- $1S-2S + 2S - 4P_{3/2}$
- $1S-2S + 2S - 6S_{1/2}$
- $1S-2S + 2S - 6D_{5/2}$
- $1S-2S + 2S - 8S_{1/2}$
- $1S-2S + 2S - 8D_{3/2}$
- $1S-2S + 2S - 8D_{5/2}$
- $1S-2S + 2S - 12D_{3/2}$
- $1S-2S + 2S - 12D_{5/2}$
- $1S-2S + 1S - 3S_{1/2}$



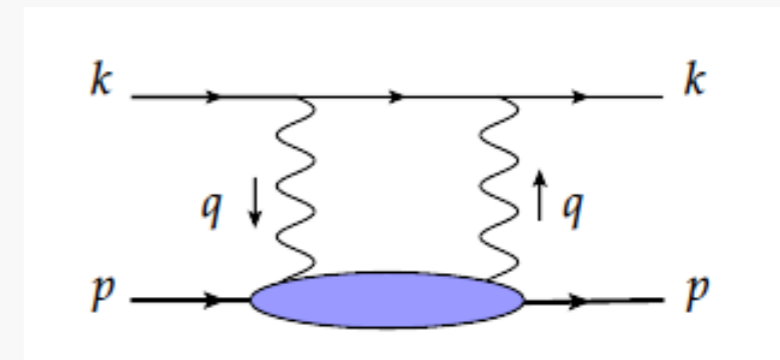
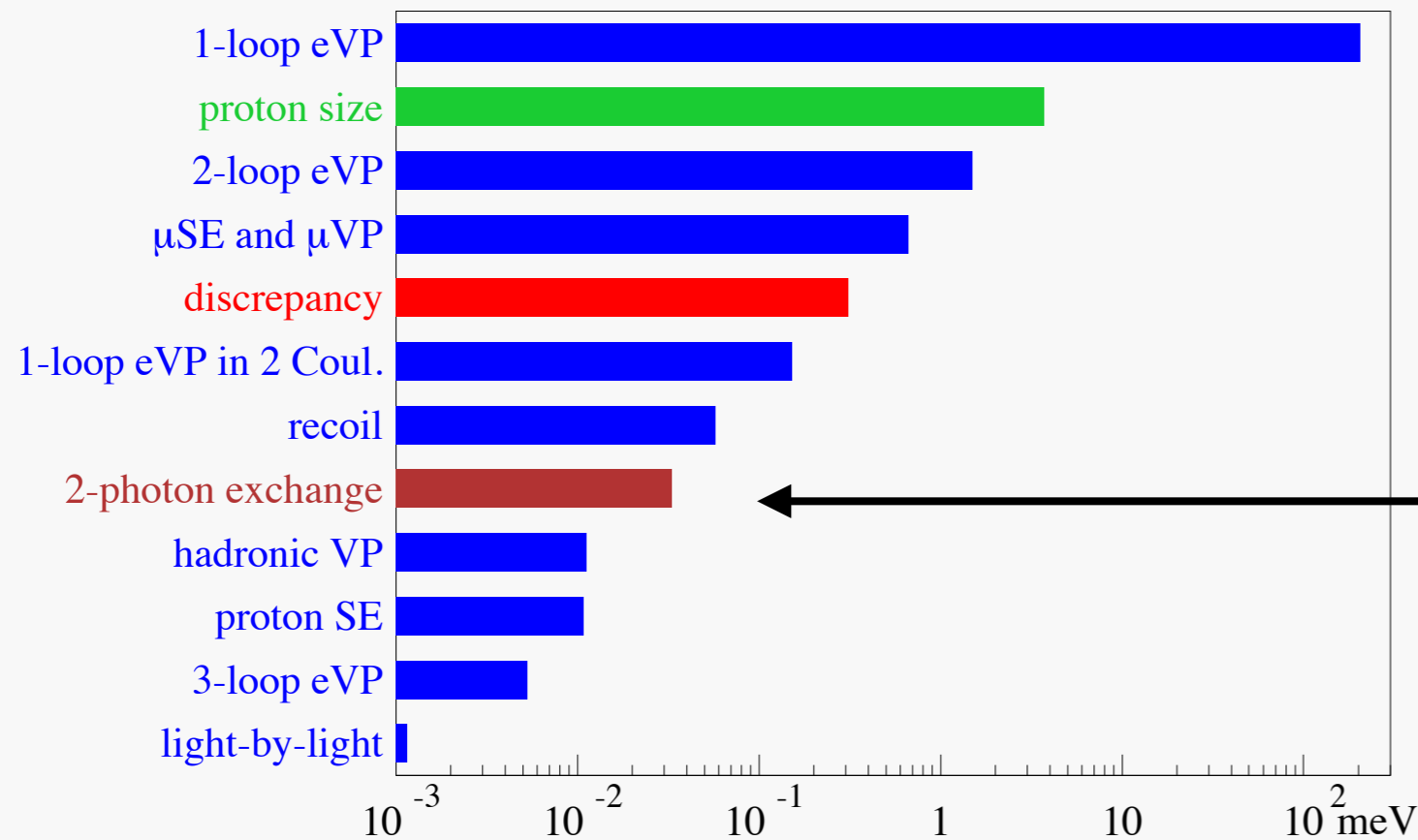
The discrepancy translates into  
**310 micro-eV deficit**  
in the Lamb shift of muonic hydrogen



# Theory of muonic hydrogen Lamb shift

Discrepancy = 0.31 meV  
 Theory uncertainty = 0.0025 meV  
 ⇒ 120 $\delta$ (theory) deviation?

$$\Delta E^{\text{th}} = 206.0668(25) - 5.2275(10) r_p^2 \text{ [meV]}$$



proposed to resolve the puzzle

De Rujula, PLB (2011)  
 Miller, PLB (2013)

Pachucki, PRA 60, 3593 (1999)

Borie, arXiv: 1103.1772-v6

Jentschura, Ann. Phys. 326, 500 (2011)

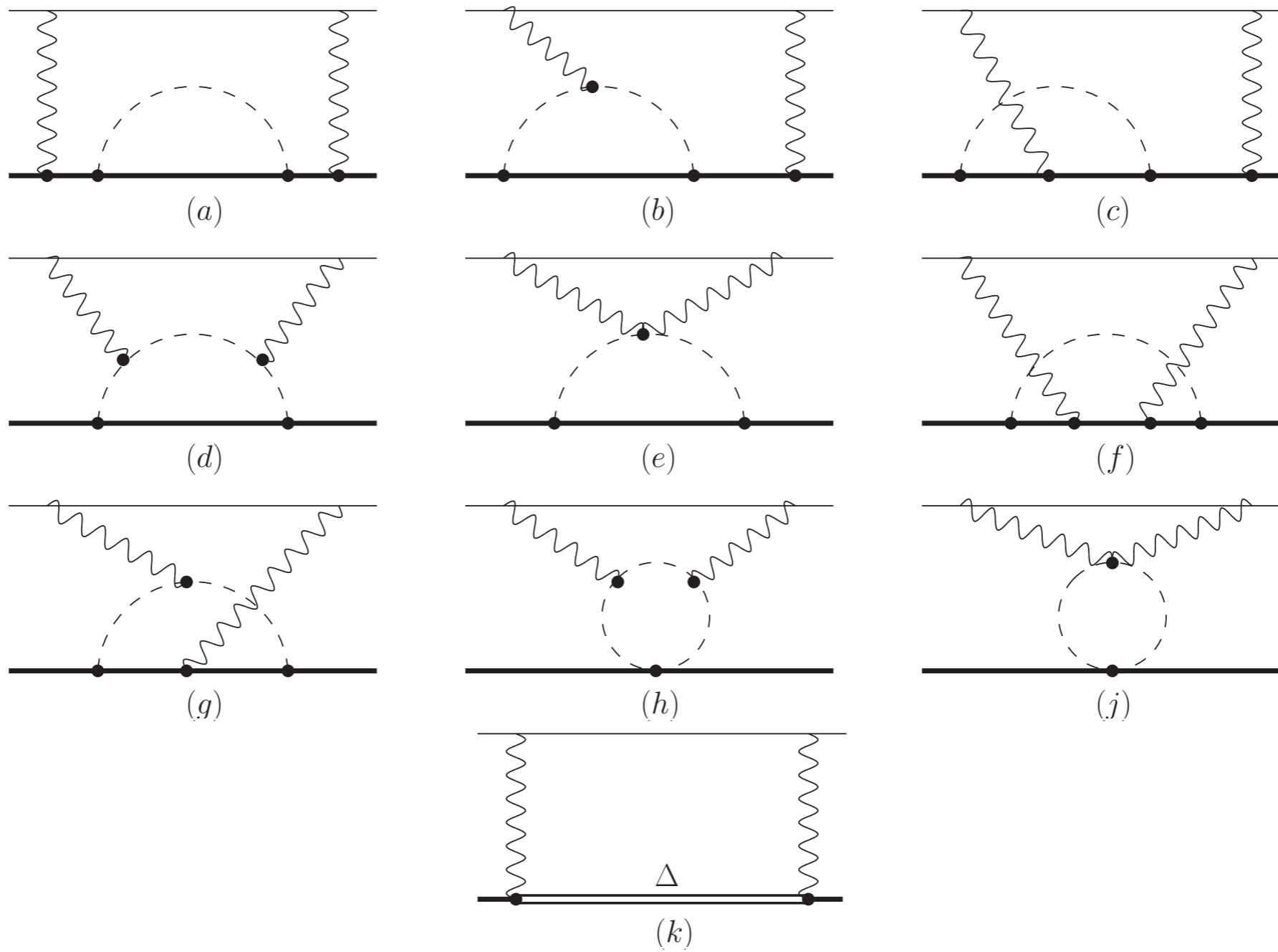
Karshenboim *et al.*, PRA 85, 032509 (2012)

calculable in ChPT

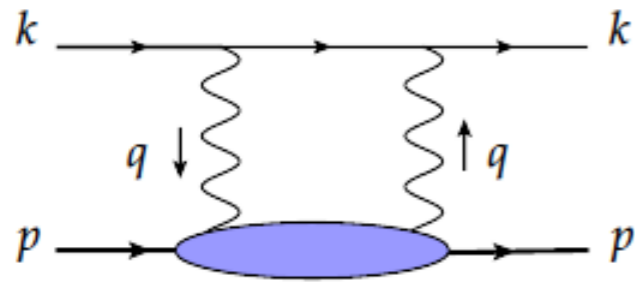
HB ChPT: Nevado & Pineda (2008)

BChPT: Alarcon, Lensky & V.P.

# Lamb shift in ChPT



# Lamb shift in terms of WVCS amplitudes



$$\Delta E_{nS}^{(\text{pol})} = -4\alpha_{em}\phi_n^2 \int_0^\infty \frac{dQ}{Q^2} w(Q^2/4m_\ell^2) \left[ T_2^{(\text{NB})}(0, Q^2) - T_1^{(\text{NB})}(0, Q^2) \right]$$

where unpolarized, forward Doubly-Virtual Compton scattering (VCS) amplitude:

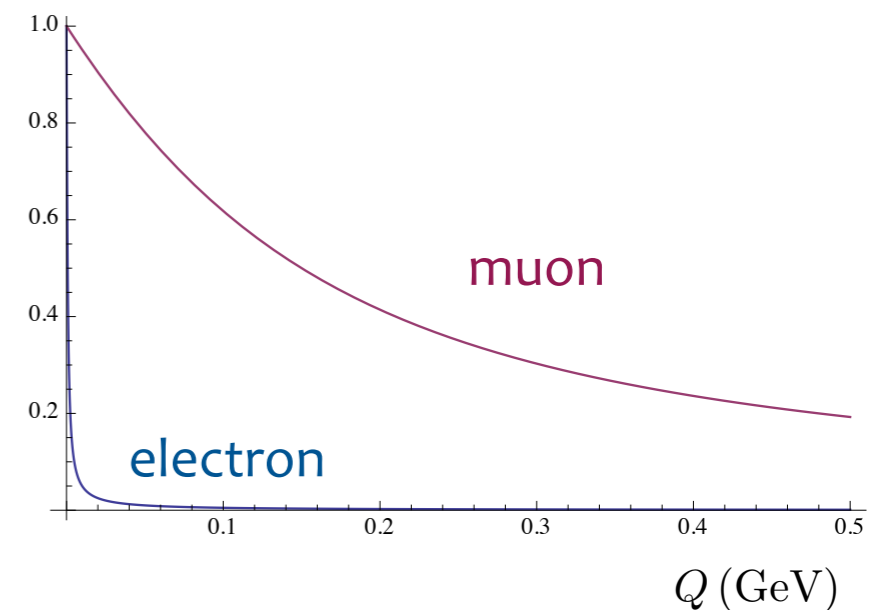
$$\begin{aligned} T^{\mu\nu}(p, q) &= \frac{i}{8\pi M} \int d^4x e^{iqx} \langle p | T j^\mu(x) j^\nu(0) | p \rangle \\ &= \left( -g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) T_1(\nu, Q^2) \\ &+ \frac{1}{M^2} \left( p^\mu - \frac{p \cdot q}{q^2} q^\mu \right) \left( p^\nu - \frac{p \cdot q}{q^2} q^\nu \right) T_2(\nu, Q^2) \end{aligned}$$

$$T_1^{(\text{NB})}(0, Q^2) \simeq Q^2 \beta_{M1}$$

$$T_2^{(\text{NB})}(0, Q^2) \simeq Q^2 (\alpha_{E1} + \beta_{M1}), \quad \text{for low } Q$$

$$\phi_n^2(0) = m_r^3 \alpha^3 / (\pi n^3)$$

$$w_\ell(Q) = \sqrt{1 + \frac{Q^2}{4m_\ell^2}} - \frac{Q}{2m_\ell}$$



# Proton polarizability effect in mu-H

Alarcon,  
Lensky,  
V.P.

HBChPT

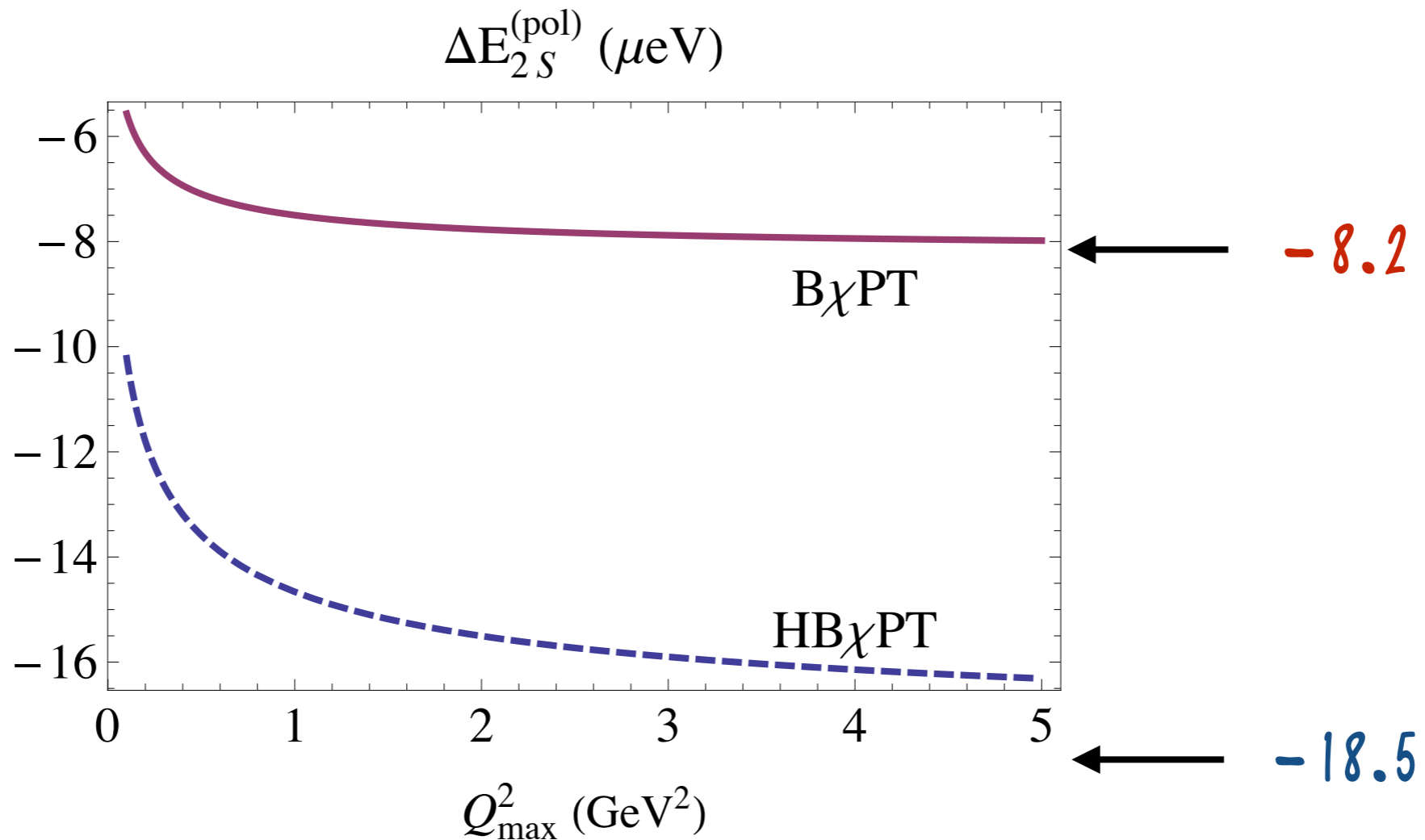
	Pachucki [9]	Marty- nenko [10]	Nevado & Pineda [11]	Carlson & Vanderhaeghen [12]	Birse & McGovern [13]	Gorchtein <i>et al.</i> [14]	LO-B $\chi$ PT [this work]
( $\mu\text{eV}$ )							
$\Delta E_{2S}^{(\text{subt})}$	1.8	2.3	—	5.3(1.9)	4.2(1.0)	3.3(4.6)	-3.0
$\Delta E_{2S}^{(\text{inel})}$	-13.9	-13.8	—	-12.7(5)	-12.7(5)*	-13.0(6)	-5.2
$\Delta E_{2S}^{(\text{pol})}$	-12(2)	-11.5	-18.5	-7.4(2.4)	-8.5	-9.7(5.6)	-8.2( $^{+1.2}_{-2.8}$ )

- [9] K. Pachucki, Phys. Rev. A **60**, 3593 (1999).  
 [10] A. P. Martynenko, Phys. Atom. Nucl. **69**, 1309 (2006).  
 [11] D. Nevado and A. Pineda, Phys. Rev. C **77**, 035202 (2008).  
 [12] C. E. Carlson and M. Vanderhaeghen, Phys. Rev. A **84**, 020102 (2011).  
 [13] M. C. Birse and J. A. McGovern, Eur. Phys. J. A **48**, 120 (2012).  
 [14] M. Gorchtein, F. J. Llanes-Estrada and A. P. Szczepaniak, Phys. Rev. A **87**, **052501** (2013).



# Polarizability effect in mu-H Lam shift

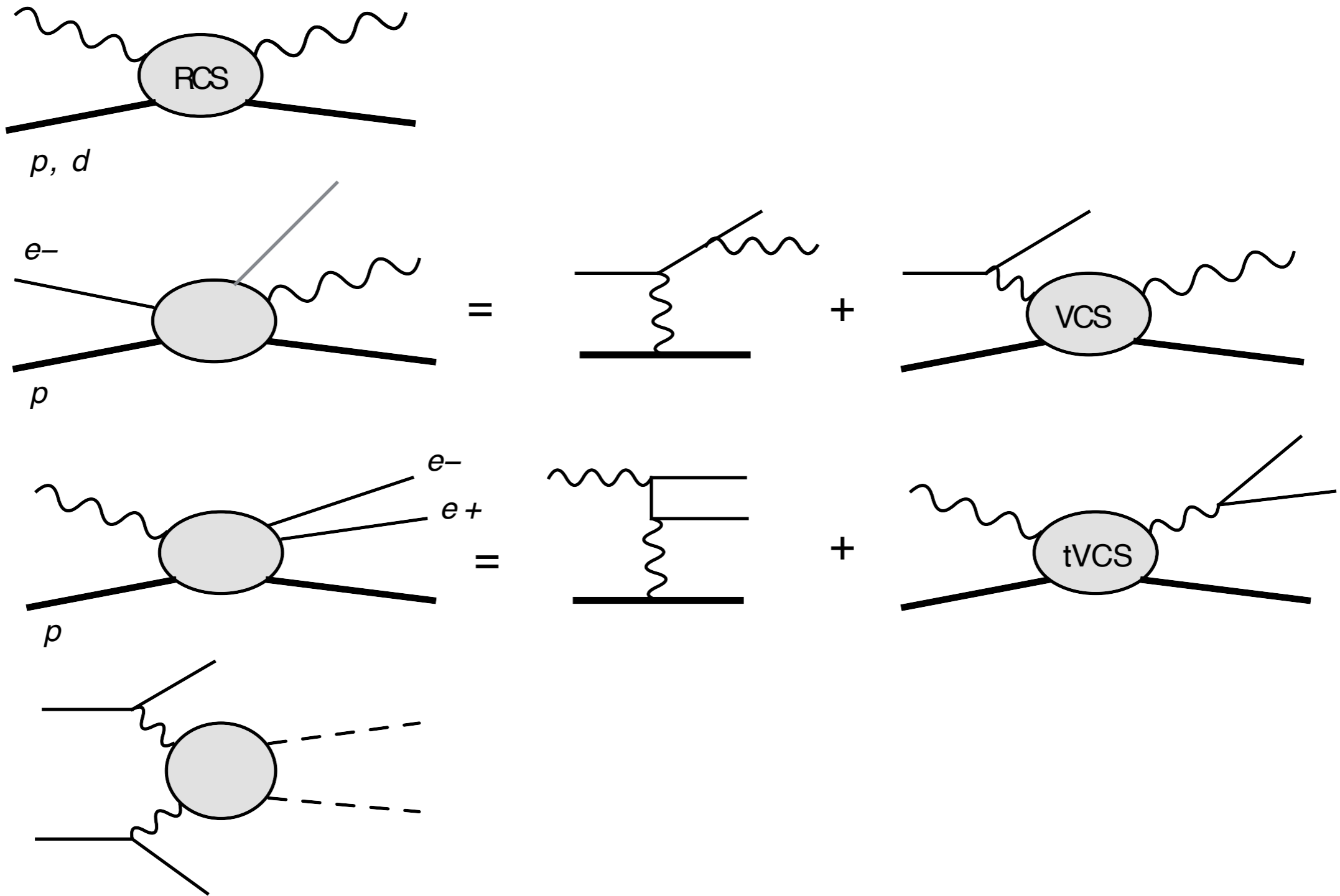
$$\Delta E_{nS}^{(\text{pol})} = -4\alpha_{em}\phi_n^2 \int_0^\infty \frac{dQ}{Q^2} w(Q^2/4m_\ell^2) \left[ T_2^{(\text{NB})}(0, Q^2) - T_1^{(\text{NB})}(0, Q^2) \right]$$



Heavy-Baryon and Baryon ChPT yield different predictions.. again this time for proton structure corrections to Lamb shift

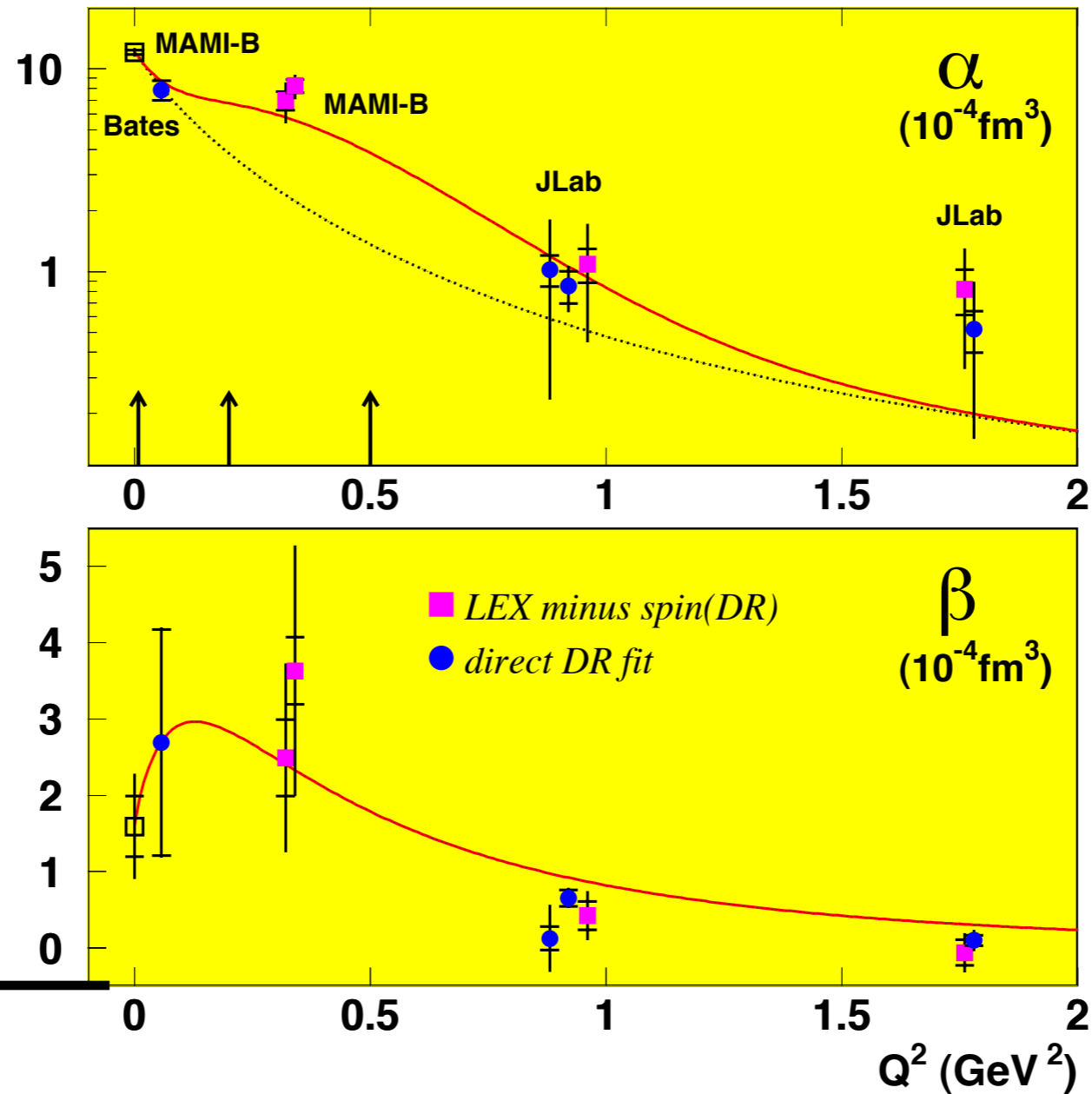
**but neither of them predicts the effect to be nearly enough to resolve the puzzle**

# More of two-photon processes



# Virtual Compton scattering (VCS)

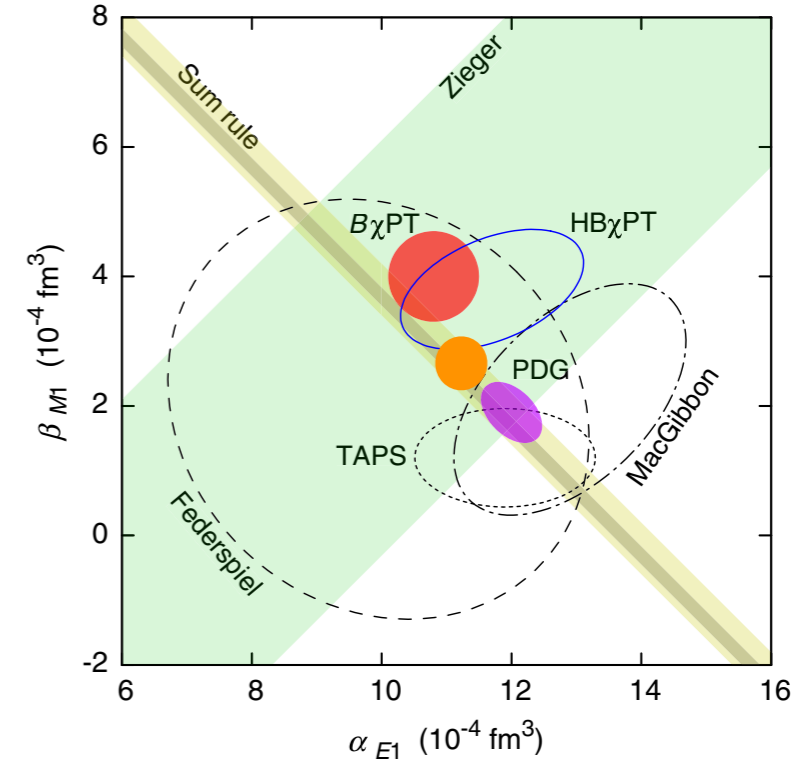
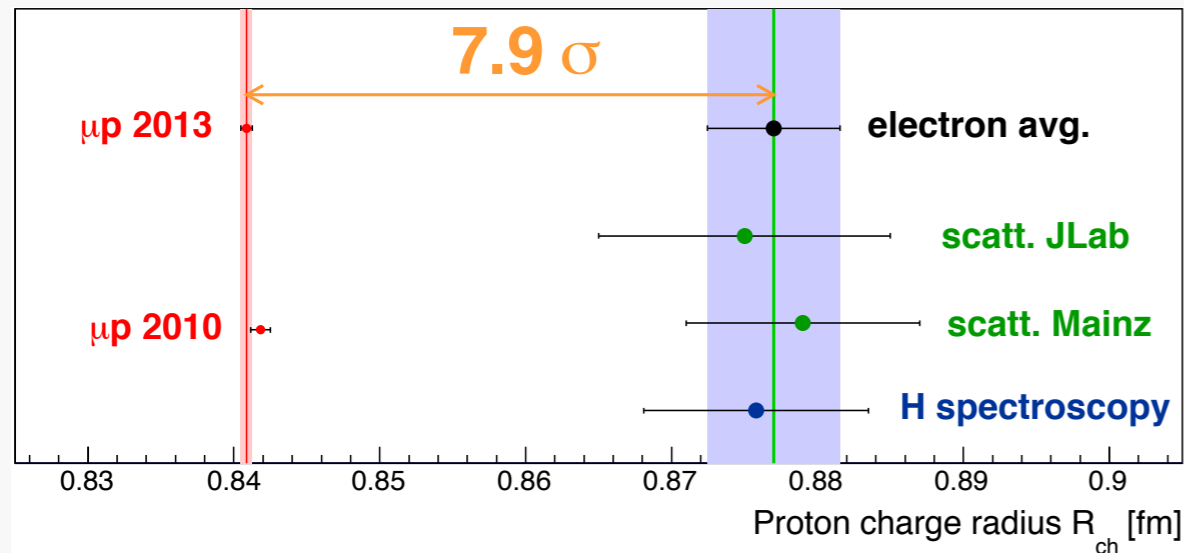
H. Fonvieille, H. Merkel et al.



"terra incognita"  
Time-like region  
( $\gamma N \rightarrow N e^+ e^-$ )

# Summary and outlook

- Nucleon structure on intersection of low-energy and precision frontiers: *proton charge radius, polarizabilities.*



Chiral PT predictions, tested in polarizabilities, *rule against scenarios* where the charge radius puzzle is explained by proton structure (beyond the radius itself)

- Stay tuned for Compton scattering (RCS, VCS, tVCS, VVCS) ongoing experiments at MAMI, HIGS, JLab and muon scattering at PSI !