

X-shape magnetic fields in galactic halos

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Cosmic rays and their interstellar medium environment
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Outline

- 1 Observational overview
- 2 Physical origin
- 3 Mathematical description
- 4 Our 4 models

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Observational tools

- Linear polarization of starlight & dust thermal emission

Due to *dust grains* → general ISM

☞ \vec{B}_\perp (orientation only)

- Zeeman splitting

Molecular & atomic *spectral lines* → neutral regions

☞ B_\parallel (strength & sign)

- Faraday rotation

Caused by *free electrons* → ionized regions

☞ B_\parallel (strength & sign)

- Synchrotron emission

Produced by *CR electrons* → general ISM

☞ \vec{B}_\perp (strength & orientation)

The Milky Way

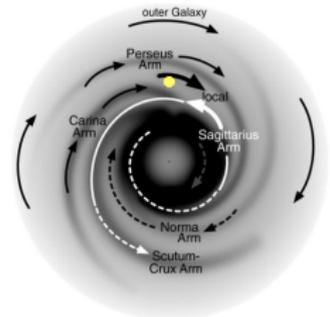
Magnetic field strength

- In general ISM
 - Near the Sun : $B_{\text{ord}} \sim 3 \mu\text{G}$ & $B_{\text{tot}} \sim 5 \mu\text{G}$
 - Global spatial distribution : $L_B \sim 12 \text{ kpc}$ & $H_B \sim 4.5 \text{ kpc}$
- In ionized regions
 - Near the Sun : $B_{\text{reg}} \simeq 1.5 \mu\text{G}$ & $B_{\text{fluct}} \sim 5 \mu\text{G}$
- In neutral regions
 - In atomic clouds : $B \sim \text{a few } \mu\text{G}$
 - In molecular clouds : $B \sim (10 - 3000) \mu\text{G}$

The Milky Way

Magnetic field direction

- Near the Sun
 - \vec{B}_{reg} is **horizontal** & **nearly azimuthal** ($p \simeq -7^\circ, -8^\circ$)
- In the Galactic disk
 - \vec{B}_{reg} is **horizontal** & **mostly azimuthal**
 - \vec{B}_{reg} **reverses direction** with decreasing radius
 - \vec{B}_{reg} is **symmetric in z**
 - \vec{B}_{reg} is neither pure **ASS** nor pure **BSS**
- In the Galactic halo
 - \vec{B}_{reg} has **horizontal** & **vertical** components
 - \vec{B}_{reg} is **anti-symmetric in z**



van Eck et al. (2011)

External spiral galaxies

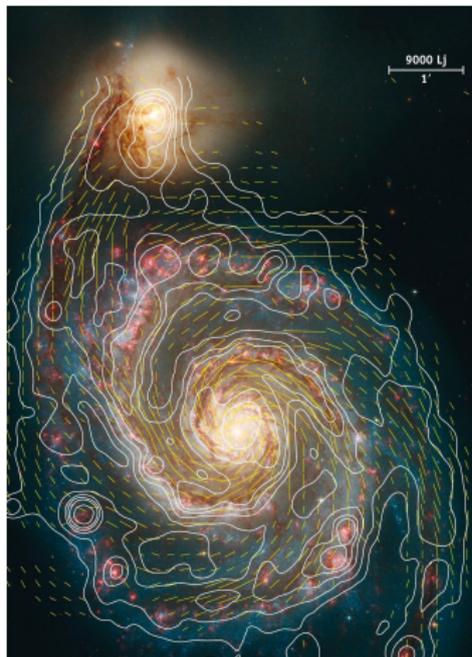
- In galactic disks
 - $B_{\text{ord}} \sim (1 - 5) \mu\text{G}$ & $B_{\text{tot}} \sim (5 - 20) \mu\text{G}$
 - \vec{B}_{ord} is **horizontal** & has a **spiral** structure

- In galactic halos
 - $B_{\text{tot}} \lesssim 10 \mu\text{G}$
 - \vec{B}_{ord} is **X-shaped**

Face-on spiral galaxy: M 51

Total intensity contours
+ apparent \vec{B} vectors
at λ 6 cm (5.0 GHz)
(100 m Effelsberg + VLA)

Optical image
(HST)

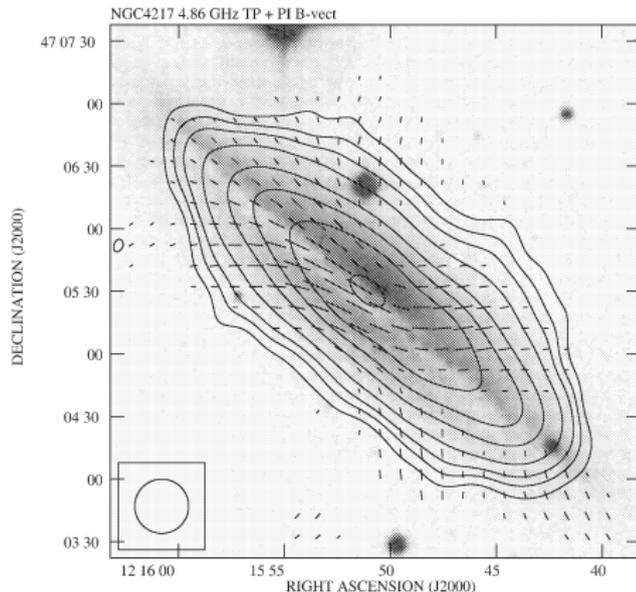


Fletcher et al. (2009)

NGC 4217

Total intensity contours
+ apparent \vec{B} vectors
at λ 6.2 cm (4.86 GHz)
(VLA)

Optical image
(DSS)

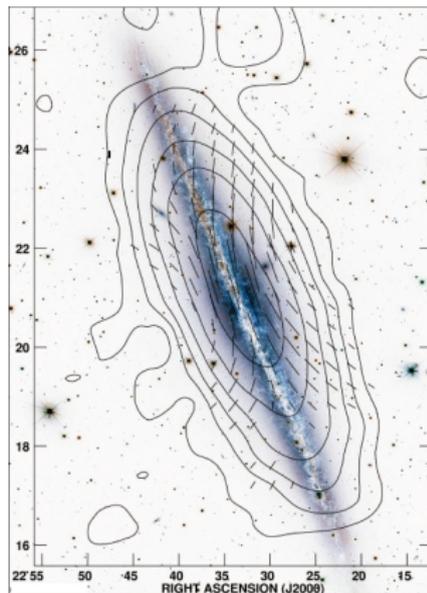


Soida (2004)

NGC 891

Total intensity contours
+ apparent \vec{B} vectors
at λ 3.6 cm (8.35 GHz)
(100m Effelsberg)

Optical image
(CFHT)



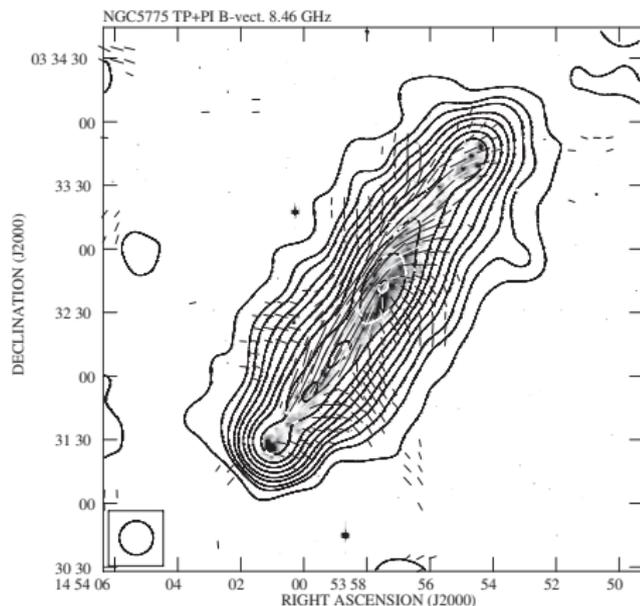
Krause (2009). © MPIfR Bonn & CFHT/Coelum

NGC 5775

Total intensity contours
+ apparent \vec{B} vectors
at λ 3.5 cm (8.46 GHz)
(VLA + 100 m Effelsberg)

H α image
(VLT)

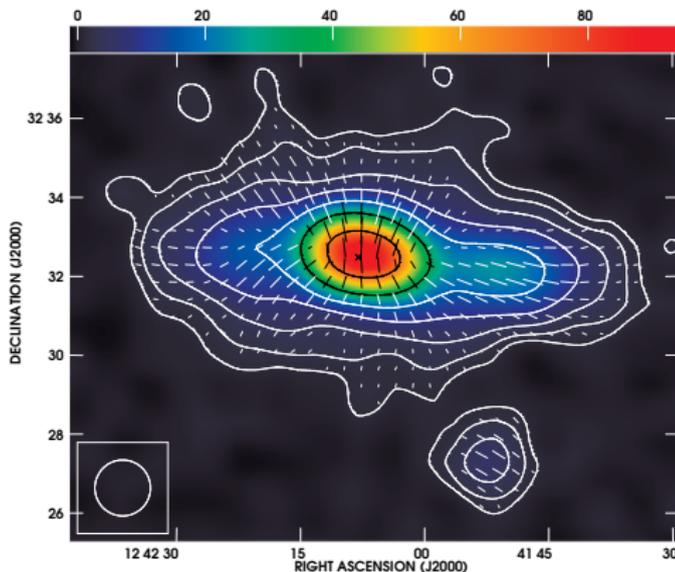
Tüllmann et al. (2000)



Soida et al. (2011)

NGC 4631

Total intensity contours
+ apparent \vec{B} vectors
at λ 3.6 cm (8.35 GHz)
(100 m Effelsberg)

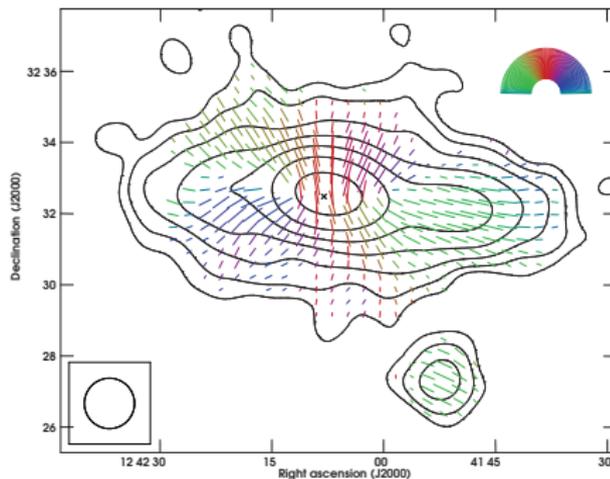


Mora & Krause (2013)

NGC 4631

Total intensity contours
at λ 3.6 cm (8.35 GHz)
(100 m Effelsberg)

+ intrinsic \vec{B} vectors
from λ 3.6 cm & λ 6.2 cm
(VLA + 100 m Effelsberg)



Mora & Krause (2013)

Outline

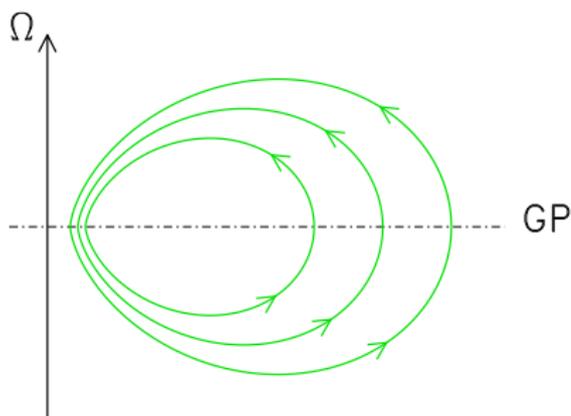
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Possible scenarios

- Large-scale regular magnetic field
 - ★ *Conventional dynamo* in the halo
 - ★ *Dynamo* in the halo + *large-scale wind* from the disk
or *outflow* from the central region
 - ★ *Dynamo* in the disk + *large-scale wind* from the disk
or *outflow* from the central region
- Small-scale anisotropic random magnetic field
 - ★ *Spiky wind* ☞ extremely elongated magnetic loops

Conventional dynamo in the halo

➡ Dipole-like magnetic field sheared out in the azimuthal direction



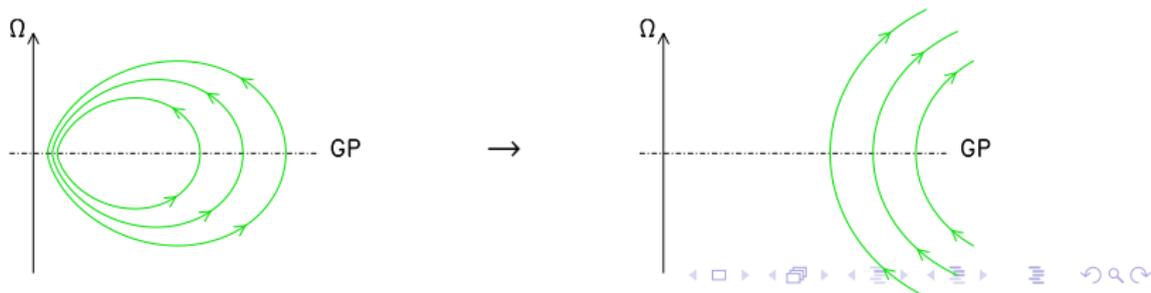
➡ Very different from an X-shape magnetic field

Halo dynamo + wind

- **Oblique wind** from the disk



- **Champagne flow** from the central region



Disk dynamo + wind

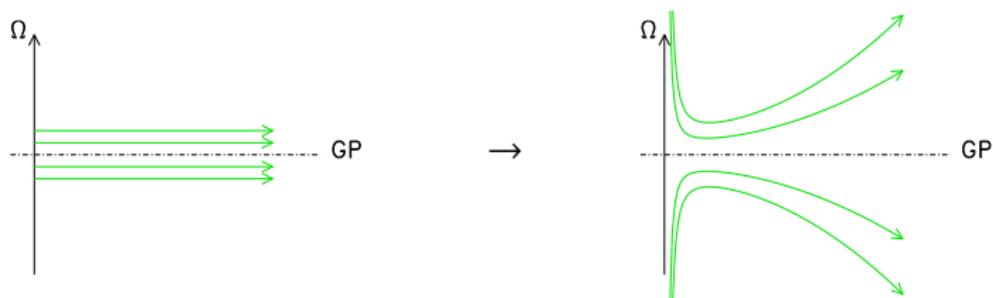
- Oblique wind from the disk



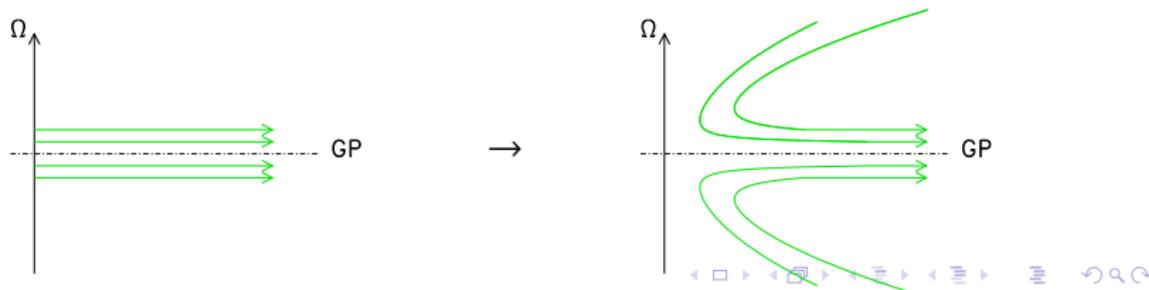
- ☞ Unphysical, because all field lines converge to the rotation axis
Mathematically, $B_r \rightarrow \infty$
- ☞ Must prevent field lines from reaching the rotation axis

Disk dynamo + wind

- **Oblique wind** from the disk + **bipolar jet** from the galactic center



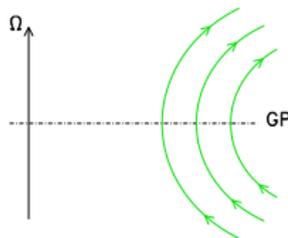
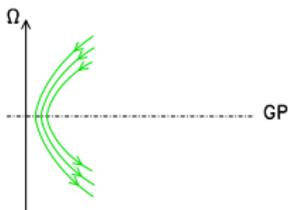
- **Champagne flow** from the central region



Vertical symmetry

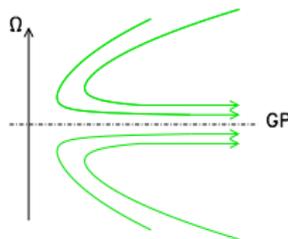
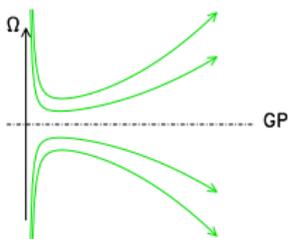
- Halo dynamo + wind

☞ \vec{B} is necessarily anti-symmetric in z



- Disk dynamo + wind

☞ \vec{B} is more likely symmetric in z (but could also be anti-symmetric)



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How to model X-shape magnetic fields ?

• Input

Consider a magnetic configuration

defined by a *network of field lines*

- ☞ - shape of field lines
- distribution of B_n on a reference surface

• Purpose

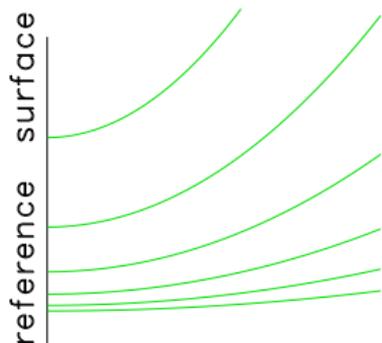
Derive an analytical expression

for the associated $\vec{B}(r, \varphi, z)$

• Method

Use the Euler potentials, α and β ,

defined such that $\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta$



Euler potentials

- Definition

2 scalar functions, α and β , such that

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta$$

- Advantages

★ \vec{B} is automatically *divergence-free*

$$\vec{\nabla} \cdot \vec{B} = 0$$

★ α and β are *constant along field lines*

$$\vec{B} \perp \vec{\nabla}\alpha \perp \text{surfaces of c}^{\text{st}} \alpha \quad \Rightarrow \quad \vec{B} \text{ tg surfaces of c}^{\text{st}} \alpha$$

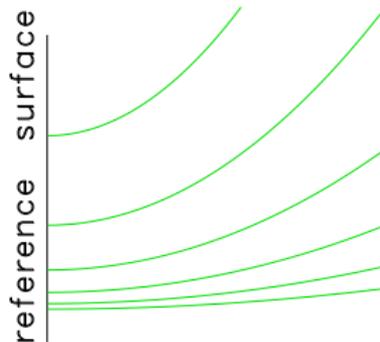
$$\vec{B} \perp \vec{\nabla}\beta \perp \text{surfaces of c}^{\text{st}} \beta \quad \Rightarrow \quad \vec{B} \text{ tg surfaces of c}^{\text{st}} \beta$$

★ Direct measure of *magnetic flux*

$$\vec{B} \cdot d\vec{S} = d\alpha d\beta$$

How to use the Euler potentials ?

- Consider a *network of field lines*
 - ☞ - shape of field lines
 - distribution of B_n on a reference surface
- Find 2 independent functions, α and β ,
 - with - α and β constant along field lines
 - $d\alpha d\beta = B_n dS$ on the reference surface
- Derive $\vec{B}(r, \varphi, z)$ using $\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta$



Poloidal, X-shape magnetic field

- Poloidal magnetic field

φ is cst along field lines

⇒ Take $\beta = \varphi$

- X-shape magnetic field

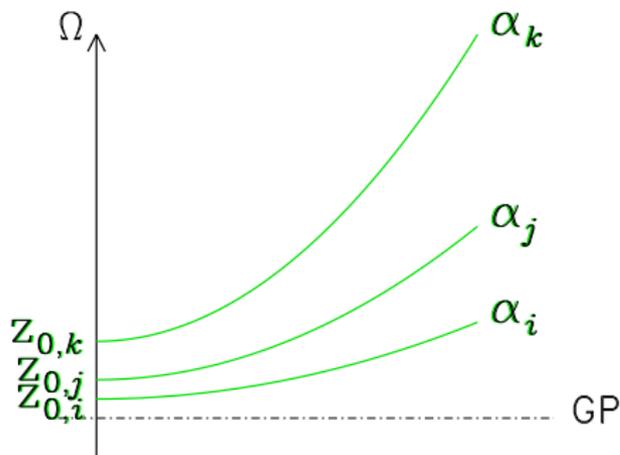
E.g., $z = z_0 (1 + ar^2)$

⇒ $z_0 = \frac{z}{1 + ar^2}$ is cst along field lines

⇒ Take $\alpha = \text{fc}(z_0)$

- Exponential decrease with z

Take $\alpha = \alpha_0 \exp\left(-\frac{|z_0|}{H}\right)$



Spiral, X-shape magnetic field

- Spiral magnetic field

E.g., $\varphi = \varphi_0 + f_\varphi(r, z)$

$\Rightarrow \varphi_0 = \varphi - f_\varphi(r, z)$ is c^{st} along field lines

\Rightarrow Take $\beta = \varphi_0$

- X-shape magnetic field

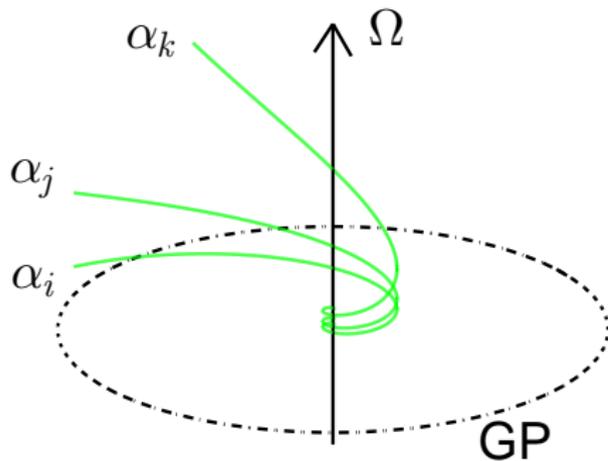
E.g., $z = z_0 (1 + ar^2)$

$\Rightarrow z_0 = \frac{z}{1 + ar^2}$ is c^{st} along field lines

\Rightarrow Take $\alpha = \text{fc}(z_0)$

- Exponential decrease with z

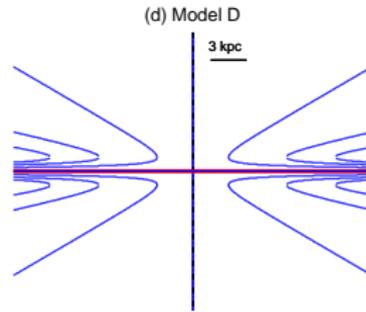
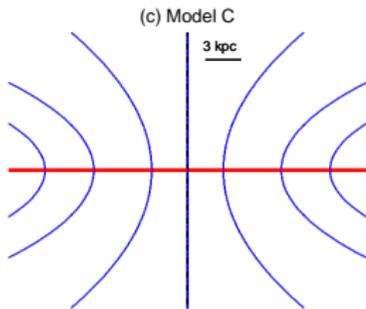
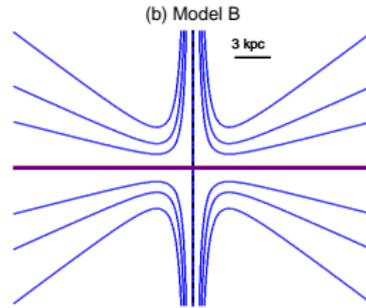
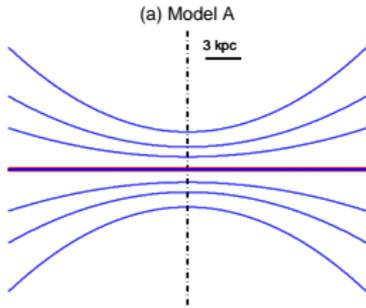
Take $\alpha = \alpha_0 \exp\left(-\frac{|z_0|}{H}\right)$



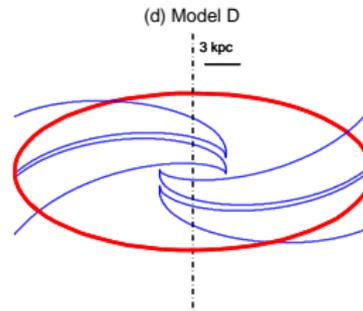
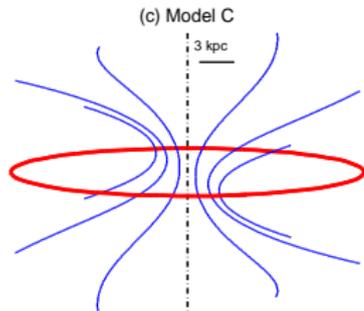
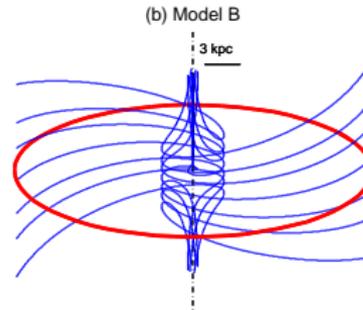
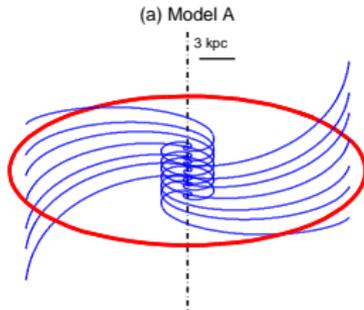
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Poloidal field lines

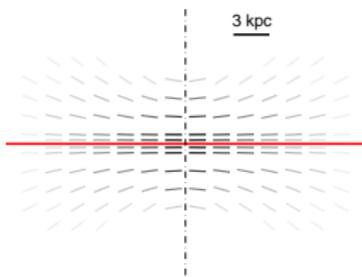


Full, spiraling field lines

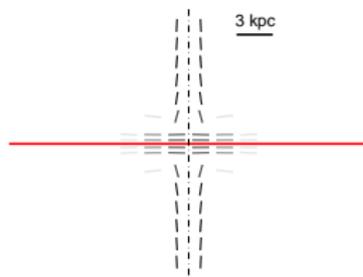


Synchrotron maps

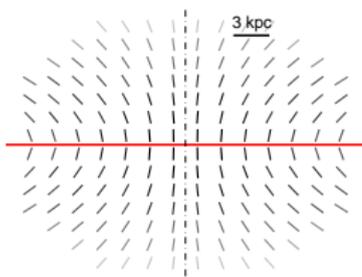
(a) Model A



(b) Model B



(c) Model C



(d) Model D

