# Light hypernuclei

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- **O** Production of  $(n, n, \Lambda, \Lambda)$ 
  - Collaboration with Qiang Zhao (IHEP) and Qian Wang (Jülich)
  - Work presented at the Chinese-French meeting, April 8, 2014
  - Meeting on Strangeness, Frascati, May 14, 2014
  - Workshop on Clusters, Yokohama, May 2014
  - Preliminary results on
    - arXiv:1404.3473 [nucl-th]
    - arXiv: arXiv:1408.1323 [nucl-th] (May 23rd, 2014)

# Strangeness, hyperons

- Late 40s, 50s, new hadrons, pair produced
- Strangeness introduced, conserved by strong interactions,
- Violated in weak decays
- Generalized isospin, to account for new hadrons, SU(2)  $\rightarrow$  SU(3)
- Eightfold way (Gell-Mann & Ne'eman)  $\rightarrow$  quarks, see also ZWeig
- Successful prediction of the  $\Omega^-$  with strangeness -3;
- Hyperons = baryons with strangeness
- S = -1  $\Lambda$ ,  $\Sigma$  plus spin, radial and orbital excitations
- $S = -2 \equiv$  plus excitations
- $S = -3 \ \Omega^-$  plus excitations
  - Very early: interaction of the hyperons among themselves and with nucleons

# Light hypernuclei: survey

- NA not bound
- AA not bound, but might be attractive at short distances
- See, speculations on *H* dibaryon (Jaffe, ...)
- $^3_{\Lambda}\mathrm{H} = (\textit{n},\textit{p},\Lambda)$  bound slightly below  $^2\mathrm{H} + \Lambda$
- This probes the  $\Lambda N$  interaction of the Nijmegen + Japan group fitting the (rare) scattering data and constrained from *NN* by SU(3)<sub>F</sub>
- ${}^{5}_{\Lambda}$ He = ( $\alpha$ ,  $\Lambda$ ) bound, another example (more in recent reviews)
- Recent efforts on (n, n, Λ), with mixing of ΛN to ΣN included, found not bound (Gal et al., Hiyama et al., Valcarce et al.)

# Light double hypernuclei: survey

- Double hypernuclei also seen
- "Nagara" event  $^{6}_{\Lambda\Lambda}$ He seen
- Indicates a weak attraction for ΛΛ in medium
- Suppression of  $\Lambda\Lambda \leftrightarrow N\Xi \leftrightarrow \Lambda\Lambda$  due to Pauli principle
- Some discussion about (n, Λ, Λ)

### Borromean binding

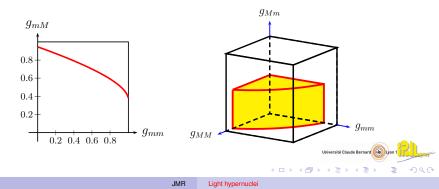
- Thomas (1935), Zhukov et al. (1993) milestones to stress the possibility of Borromean binding
- Bound state whose subsystems are unbound



- In quantum mechanics for bosons, a short-range attractive potential g V(r)
- Needs *g* > *g*<sub>2</sub> to bind (*m*, *m*)
- Needs  $g > g_3$  to bind (m, m, m)
- With  $g_3 < g_2$   $[g_3, g_2]$  Borromean windows

#### Borromean binding

- Example  ${}^{6}\text{He} = (\alpha, n, n)$  is stable
- While neither  $(\alpha, n)$  nor (n, n) are bound
- Rigorous limit  $2g_2/3 < g_3$  can be established
- And generalized to unequal masses and coupling, and to N > 3 body systems such as (mmM) and (mmMM)

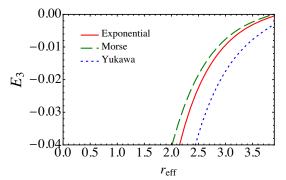


# Optimist view

- So, some hope is permitted, as the 'allowed" domain is large for (*m*, *m*, *M*, *M*)
- Including 4-body binding without 3-body binding
- As in atomic physics
  - Positronium hydride stable (Øre, 1951): (p, e<sup>+</sup>, e<sup>-</sup>, e<sup>-</sup>) stable though (p, e<sup>+</sup>, e<sup>-</sup>) is unstable
  - (M<sup>+</sup>, m<sup>+</sup>, M<sup>−</sup>, m<sup>−</sup>) stable for M ~ 2 m, though none of the 3-body subsystems such as (M<sup>+</sup>, m, m<sup>−</sup>) is stable (R., Phys. Rev. A 67, 034702)
- Will the magic of 4-body systems work once more?

### Borromean binding

- for a given scattering length a<sub>sc</sub>,
- Amount of Borromean binding depends on effective range
- Already stressed (Gibson + Lehman, Gal et al., ...)
- But the effect differs according to the sign of asc



- Figure:  $E_3$  vs.  $r_{\text{eff}}$  at fixed  $a_{\text{sc}}$
- Here for negative a<sub>sc</sub>

### Variational calculations

• Use simple central potentials (exp., Morse) reproducing  $a_{sc}$  and  $r_{eff}$ , as done by several groups,

• 3-body: 
$$\Psi = \sum_{i} \gamma_{i} \exp(-a_{i} x - b_{i} y - c_{i} z), \quad x = |r_{2} - r_{3}|, \dots$$

- For given {a<sub>i</sub>, b<sub>i</sub>, c<sub>i</sub>}, variational energy and {γ<sub>i</sub>} from a simple eigenvalue equation
- Non-linear parameters {a<sub>i</sub>, b<sub>i</sub>, c<sub>i</sub>} searched for numerically, as any triple inside a geometric series a, av, av<sup>2</sup>,.... So only a and v are varied (cf. Nakamura et al.)
- Cross-checked with Varga and Suzuki's stochastic variational method (SVM) based on Gaussians

$$\Psi = \sum_{i} \gamma_i \, \exp[-\sum_{k < \ell} a_i^{k\ell} \, (\boldsymbol{r}_\ell - \boldsymbol{r}_k)^2] \; ,$$

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SVM also used for 4-body

- N R Nijmegen–RIKEN series of fits with meson exchanges
- *CEFT Chiral Effective Field Theory* Which tend to give smaller *r*<sub>eff</sub> and thus more binding

#	( <i>np</i> ) <sub><i>l</i>=1</sub>	( <i>nn</i> )	$(np)_{I=0}$	$(\Lambda N)_{s=0}$	$(\Lambda N)_{s=1}$	(\\\)
N – R	-23.735 2.694	-16.51 2.85				-0.97 3.88
CEFT	-23.735 2.694			-2.90 2.65	-1.51 2.64	



#### Results

- Very delicate 3-body and 4-body calculations
- Warm-up without tuning of reff
  - With V = -g<sub>ij</sub> exp(-r), r in GeV, tuned to reproduce the scattering lengths (n, p), (n, p, Λ) OK, (n, n, Λ) unbound, (n, p, Λ, Λ) and (n, n, Λ, Λ) stable, but effective range too small
  - With  $V = -g_{ij} \exp(-0.2 r)$ , *r* in GeV, (*n*, *p*), (*n*, *p*,  $\Lambda$ ) OK, (*n*, *n*,  $\Lambda$ ) unbound, (*n*, *p*,  $\Lambda$ ,  $\Lambda$ ) bound and (*n*, *n*,  $\Lambda$ ,  $\Lambda$ ) unstable, but effective range too large,
- With tuning of  $r_{\text{eff}}$   $V = -g_{ij} \exp(-\mu_{ij} r)$ , or  $V = g_{ij} \exp[-2\alpha_{ij}(r-R)] - 2g_{ij} \exp[-\alpha_{ij}(r-R)]$  (Morse), and a reasonable *R*, one can adjust  $g_{ij}$  and  $\mu_{ij}$  or  $\alpha_{ij}$  to reproduce the scattering length and effective range for each pair, then  $(n, p), (n, p, \Lambda)$  OK,  $(n, n, \Lambda)$  unbound,  $(n, p, \Lambda, \Lambda)$  bound and  $(n, n, \Lambda, \Lambda)$  very weakly bound with CEFT,
- First indication for this neutral configuration which is fully
  Borromean

### Open questions: Three-body forces

- Probably very repulsive at high nuclear density: no room for additional u or d in a medium already saturated for u and d by two neighbouring nucleons
- Might be attractive for few-baryon systems
- Evolution along the years:
- In the early 70s, a hard core was assumed for any baryon–baryon interaction
- Now the core is softer, and is understood as coming from the Pauli principle for quarks
- Repulsive core e.g., for (u, u, d) + (u, d, d)
- No repulsion for (*u*, *d*, *s*) + (*u*, *d*, *s*)
- But onset of repulsion if more Λ are put together

### Open questions: coupled channels

- $\Lambda\Lambda \leftrightarrow \Xi N$  suppressed at high density, if medium saturated for Ns,
- May be not for a very dilute  $(\Lambda, \Lambda, n, n)$
- Comparison of ( $\Lambda NN$ ) with or without coupling to  $\Sigma NN$  when  $\Lambda N$  alone and  $\Lambda N + \Sigma N$  give the same  $a_{sc}$  and  $r_{eff}$
- Extra attraction if new spin-isospin configurations (Gibson, ...)
- For instance, leading s-wave decomposition schemes

$$\begin{split} &(n,p,\Lambda)_{I=0,s=3/2} = (np)_{0,1}\Lambda_{0,1/2} + (np)_{0,1}\Sigma_{1,1/2} \\ &(n,p,\Lambda)_{I=0,s=1/2} = (np)_{0,1}\Lambda_{0,1/2} + (np)_{0,1}\Sigma_{1,1/2} + (NN)_{1,0}\Sigma_{1,1/2} \end{split}$$

Similarly, for the example of the ΣΣ sector

 $\begin{aligned} (n, p, \Lambda\Lambda)_{I=0,s=1} &= (np)_{0,1}(\Lambda\Lambda)_{0,0} + (np)_{0,1}(\Sigma\Sigma)_{0,0} + (NN)_{1,0}(\Sigma\Sigma)_{1,1} \\ (n, p, \Lambda\Lambda)_{I=1,s=0} &= (nn)_{0,0}(\Lambda\Lambda_{0,0} + (nn)_{0,0}(\Sigma\Sigma)_{0,0} + (np)_{0,1}(\Sigma\Sigma)_{1,1} \end{aligned}$ 

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both can get some small downwards shift.

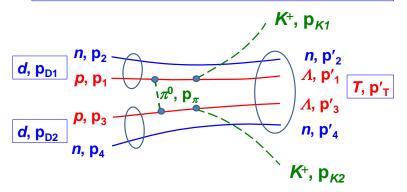
### Open questions: AA interaction

- Very short-range in the CEFT model, with  $r_{\rm eff} = 0.3$  fm.
- If mimicked by an exponential, one gets almost a delta function, this making the numerical estimate more delicate,
- Reminiscent of Jaffe's model for the  $H = (uuddss) = \Lambda\Lambda + \cdots$ , where the attraction is due the short-range part of the interquark force, the chromo-magnetic interaction.
- The existence of (n, n, Λ, Λ) would thus suggest a short-range, attractive ΛΛ interaction that is strongly suppressed in the nuclear medium
  - No medium effect for the dilute (n, n, Λ, Λ)
  - Suppression of  $\Lambda\Lambda \leftrightarrow N\Xi$  for heavier double hypernuclei

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### Production mechanism

Production mechanism for the tetra-baryon state ( $\Lambda\Lambda$ nn) in deuteron-deuteron collision

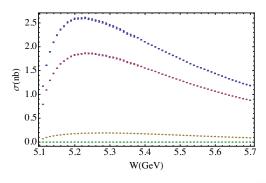


- The two K<sup>+</sup> production is via the elementary process,  $pp \rightarrow \Lambda\Lambda K^+K^+$ .
- The intermediate S11(1535) excitations dominate the threshold production.



### Production mechanism

- Rough estimate
- Total, and Born term, double *N*(1530), single *N*(1530) contributions



### Conclusions

#### A = 3

- A second  $(n, p, \Lambda)$ , with l = 0 and s = 3/2 cannot be excluded
- (n, p, Λ) with I = 1 or (n, n, Λ) hardly bound with 2-body forces only
- (Ν, Λ, Λ) debated
- Three-body forces to be studied more carefully

#### *A* = 4

- Calculations confirms the likely existence of a (n, p, Λ, Λ)
- New light hypernuclei expected such as (n, n, Λ, Λ)
- In a regime of extreme Borromean binding
- If ΛN and ΛΛ have small effective range, as suggested by effective field theories