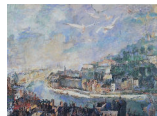


Multiquarks and molecules

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Content

- 1 $X(3872)$ and $X(3900)$
- 2 (n, n, Λ, Λ)

Work done in collaboration with Qiang Zhao (IHEP) and Qian Wang (Juelich),
and Emi Kou et al. for the first part.

Isospin partner of $X(3872)$ -1

$X(3872)$

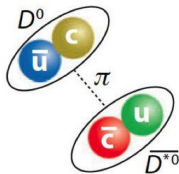
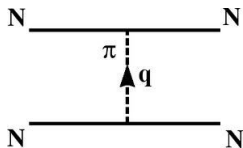
- $X(3872)$ seen in several experiments
- BELLE, BABAR, CDF, D0
- More recently at LHC
- Firmly established, $J^{PC} = 1^{++}$
- $M = 3871.68 \pm .17 \text{ MeV}$ (PDG) $\Gamma = 3.9^{+2.8+0.2}_{-1.4-1.1} \text{ MeV}$
- Decay mostly into $D\bar{D} + \text{c.c.}, D\bar{D}\pi$

$X(3872)$ Theory

- Predicted in a **molecular** model (Törnqvist, Voloshin, ...)
- Also studied in a **diquark** model (Maiani et al.)
- Several studies with QCD SR (Nielsen, Narison, S.-L. Zhu et al., Z.-G. Wang, ...)
- Firmly established
- Chromomagnetic model (Høgassen et al.)



Isospin partner of $X(3872)$ -2



- Favours one isospin
- Due to $\tau_1 \cdot \tau_2$ factor
- Historically, this was a problem in “bootstrap” theory, where mesons were mostly baryon-antibaryon
- Factor 3 better for $I = 0$
- Again, due to $\tau_1 \cdot \tau_2$ factor
- Favours $I = 0$ vs. $I = 1$

Isospin partner of $X(3872)$ -3

PHYSICAL REVIEW D 73, 054013 (2006)

Chromomagnetic mechanism for the $X(3872)$ resonance

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(Received 3 November 2005; published 17 March 2006)

The chromomagnetic interaction, with proper account for flavor-symmetry breaking, is shown to explain the mass and coupling properties of the $X(3872)$ resonance as a $J^{PC} = 1^{++}$ state consisting of a heavy quark-antiquark pair and a light one. It is crucial to introduce all the spin-color configurations compatible with these quantum numbers and diagonalize the chromomagnetic interaction in this basis. This approach thus differs from the molecular picture $D\bar{D}^*$ and from the diquark-anti-diquark picture.

annihilation term and $m_d - m_u = 3.5$ MeV, the “mostly $I = 1$ ” state lies 31 MeV above the “mostly $I = 0$ ” state.

symmetry

dimensions
of particle
physics

Chinese collider expands particle zoo

China's Beijing Electron-Positron Collider seems to be hosting a reunion; members of a poorly understood family of particles keep popping up in their data, which may help clarify the properties of this reclusive family.

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Isospin partner of $X(3872)$ -4

Observation of a Charged Charmoniumlike Structure in $e^+e^- \rightarrow \pi^+\pi^- J/\psi$ at $\sqrt{s} = 4.26$ GeV

BESIII Collaboration (M. Ablikim (Beijing, Inst. High Energy Phys.) et al.) [Afficher les 354 auteurs](#)

Mar 24, 2013 - 7 pages

to charmonium. A fit to the $\pi^\pm J/\psi$ invariant mass spectrum, neglecting interference, results in a mass of $(3899.0 \pm 3.6 \pm 4.9)$ MeV/c² and a width of $(46 \pm 10 \pm 20)$ MeV. Its production ratio is measured to be

Study of $e^+e^- \rightarrow \pi^+\pi^- J/\psi$ and Observation of a Charged Charmoniumlike State at Belle

Belle Collaboration (Z.Q. Liu (Beijing, Inst. High Energy Phys.) et al.) [Afficher les 187 auteurs](#)

Mar 30, 2013 - 7 pages

Isospin partner of $X(3872)$ -5

Chromomagnetic model

$$H = M + H_{\text{CM}} = \sum_i m_i - \sum_{i,j} C_{ij} \tilde{\lambda}_i \cdot \tilde{\lambda}_j \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j ,$$

- Gives $X(3872)$ with right mass and decay properties
- $I = 0$ and $I = 1$ degenerate
- **Annihilation term** predicted the right splitting
- X is mainly $(c\bar{c})_8(q\bar{q})_8$

Beauty partners of $X(3872)$ and $Z_c(3900)$

Physics Letters B 732 (2014) 97–100



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Isovector and hidden-beauty partners of the $X(3872)$

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$$H = M + H_{\text{CM}} + H_{\text{CE}} = \sum_i m_i - \sum_{i,j} C_{ij} \tilde{\lambda}_i \cdot \tilde{\lambda}_j \sigma_i \cdot \sigma_j - \sum_{i,j} A_{ij} \tilde{\lambda}_i \cdot \tilde{\lambda}_j ,$$

- One quadruplet near 10.62 GeV
- No possibility of two charged triplets with $J^{PC} = 1^{++}$ at about the same mass

Summary Part 1

- For a review of other approaches, see Q. Zhao at Beijing Chiral Workshop
- Near degeneracy of $I = 0$ and $I = 1$ natural in a quark picture
- Molecular model predicts a difference between $I = 0$ and $I = 1$
- Radiative decay seems crucial, in particular

$$\frac{X(3872) \rightarrow \psi(2s) + \gamma}{X(3872) \rightarrow J/\psi + \gamma} \gg 1$$

- See, the recent review by S.L. Zhu et al.
- See the recent better measurement by LHCb

Summary Part 1

X 3900 is a local speciality

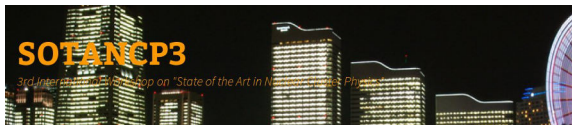
L'autorail X 3900



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Part 2: New light hypernuclei

- Much activity in the “molecular’ model,
- meson-meson, baryon–baryon, baryon–antibaryon, etc.
- stimulated by $X(3872)$ and other states in the heavy-quark sector
- also for nucleon-hyperon systems = **light hypernuclei**
- Qiang Zhao (IHEP) Qian Wang (Jülich) and JMR, to be posted and submitted next week
- and presented next month at the Conference on Nuclear Clusters (Yokohama)

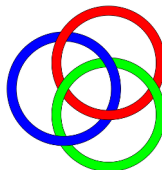


Light hypernuclei: survey

- $N\Lambda$ not bound
- $\Lambda\Lambda$ not bound, but might be attractive at short distances
- See, speculations on H dibaryon (Jaffe, ...)
- ${}^3_{\Lambda}\text{H} = (n, p, \Lambda)$ bound slightly below ${}^2\text{H} + \Lambda$
- This probes the ΛN interaction of the Nijmegen + Japan group fitting the (rare) scattering data and constrained from NN by $\text{SU}(3)_{\text{F}}$
- ${}^5_{\Lambda}\text{He} = (\alpha, \Lambda)$ bound, another example (More in Sakaguchi, 2009)
- **Double hypernuclei** also seen
- “Nagara” event ${}^6_{\Lambda\Lambda}\text{He}$ seen
- Indicates a weak attraction for $\Lambda\Lambda$ in medium
- Suppression of $\Lambda\Lambda \leftrightarrow N\Xi \leftrightarrow \Lambda\Lambda$ due to Pauli principle

Borromean binding

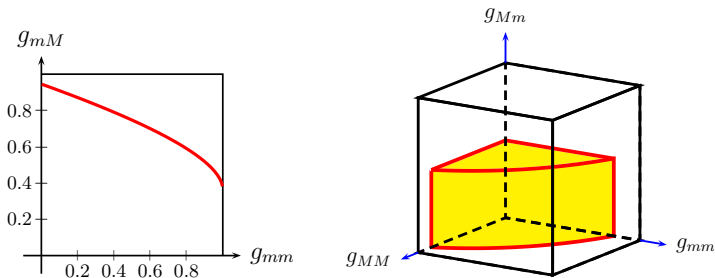
- Thomas (1935), Zhukov et al. (1993) milestones to stress the possibility of **Borromean binding**
- Bound state whose **subsystems are unbound**



- In quantum mechanics for bosons, a short-range attractive potential $g V(r)$
- Needs $g > g_2$ to bind (m, m)
- Needs $g > g_3$ to bind (m, m, m)
- With $g_3 < g_2$ $[g_3, g_2]$ Borromean windows

Borromean binding

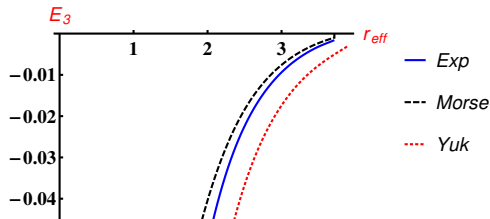
- Example ${}^6\text{He} = (\alpha, n, n)$ is stable
- while neither (α, n) nor (n, n) are bound
- Rigorous limit $2g_2/3 < g_3$ can be established
- And generalized to unequal masses and coupling, and to $N > 3$ body systems



So, some hope is permitted

Borromean binding

The amount of Borromean binding, for given scattering length, depends on the effective range a point sometimes overlooked.

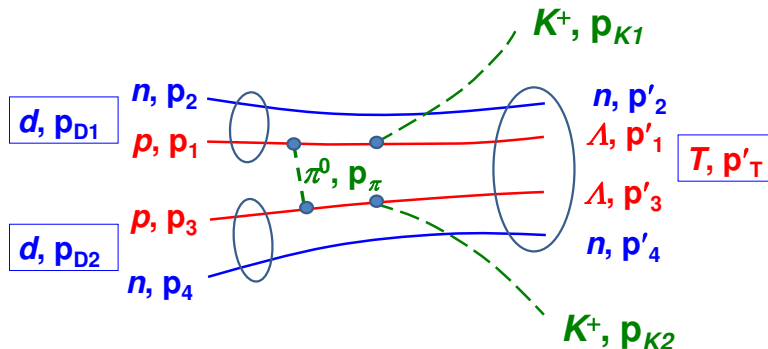


Results

- Very delicate 3-body and 4-body calculations
- With $V = -g_{ij} \exp(-r)$, r in GeV, tuned to reproduce the scattering lengths (n, p) , (n, p, Λ) OK, (n, n, Λ) unbound, (n, p, Λ, Λ) and (n, n, Λ, Λ) stable, but effective range too small
- With $V = -g_{ij} \exp(-0.2 r)$, r in GeV, (n, p) , (n, p, Λ) OK, (n, n, Λ) unbound, (n, p, Λ, Λ) bound and (n, n, Λ, Λ) unstable, but effective range too large,
- With $V = g_{ij} \exp[-2\alpha_{ij}(r - R)] - 2 g_{ij} \exp[-\alpha_{ij}(r - R)]$ (Morse), and a reasonable R , one can adjust g_{ij} and α_{ij} to reproduce the scattering length and effective range for each pair, then (n, p) , (n, p, Λ) OK, (n, n, Λ) unbound, (n, p, Λ, Λ) bound and (n, n, Λ, Λ) very weakly bound
- First indication for this neutral configuration
- Which is **fully Borromean**

Production mechanism

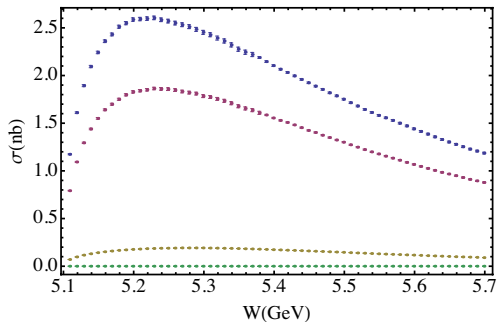
Production mechanism for the tetra-baryon state ($\Lambda\Lambda nn$) in deuteron-deuteron collision



- The two K^+ production is via the elementary process, $pp \rightarrow \Lambda\Lambda K^+ K^+$.
- The intermediate $S11(1535)$ excitations dominate the threshold production.

Production mechanism

- Preliminary estimate (still under checking)
- Total, and Born term, double $N(1530)$, single $N(1530)$ contributions



Conclusions

$Z_c(3900)$

- $Z_c(3900)$ expected in models based on quark dynamics
- Predictions for hidden beauty

Neutral tetrabaryon

- New light hypernuclei expected such as (n, n, Λ, Λ)
- In a regime of extreme Borromean binding