

GDR – PH – QCD Workshop on Radiative Corrections in Annihilation and Scattering Experiments

Simulations with MC generators for $e^+e^-(X)$ processes







- Simulations with EKHARA^{*} will be presented
- Trick&tips on the implementation of EKHARA in the BES III code
- Examples from B factories: $e^+e^- \rightarrow e^+e^-X$
- PANDA: my first simulations for $p\overline{p} \rightarrow Y \rightarrow e^+e^-$
- Remarks

What is EKHARA



- EKHARA is a MC generator written for specific processes $e^+e^- \rightarrow e^+e^-X$, X = $\pi^0, \eta, \eta', \pi^+\pi^-$
- It is published and available online: prac.us.edu.pl/~ekhara/

Phys.Rev. D85 (2012) 094010, H.Czyz, S. Ivashyn

- It makes use of the double-octet model
- Fit makes use of the results from "real" data
- Very good simulations for projects willing to study form factors, cross section, in $e^+e^- \rightarrow e^+e^-X$ via $\gamma\gamma$ interactions.

π^{0} transition form factor



• The process under study here is: $e^+e^- \rightarrow e^+e^-\pi^0$ via $\gamma\gamma$ interactions



η,η' transition form factor





PUZZLING, compared to the case with π^{0} !

Double octet model: π^{0}





Still not a lot of information at very low Q², from experiments

Need to know the distribution of F(Q²)•Q² vs Q² in order to give important contribution to understanding LBL hadronic corrections to (g-2)

Theoretically, the most interesting range for the LBL hadronic corrections to $(g-2)_{\mu}$ is: $Q^2 \in [0 \div 2] \text{ GeV}^2$

FIG. 2 (color online). Transition form factor $\gamma^* \gamma \pi^0$ compared to the data. The Brodsky-Lepage [30] high- Q^2 limit (BL) is shown as a bold solid straight line at $2 \times f_{\pi} = 2 \times 0.0924$ GeV. The high- Q^2 limit in our 1-octet ansatz and and 2-octet ansatz are marked as (1) and (2), respectively.



• Double-octet model makes prediction for $F(q_1^2, q_2^2)$ with one lepton tagged, or both leptons tagged, or none of them. Tagging one of the 2 leptons makes the calculation of $F(q_1^2, q_2^2)$ easier, as it depends only on the transfer momentum of the **tagged lepton**(e.g. $q_1=Q$, so $q_2=0$). There are no free parameters in $F(Q^2, 0)$ in this model, except only one (the coupling constant Hv1)



FIG. 5 (color online). The cross section $d\sigma/dQ^2$ for the process $e^+e^- \rightarrow e^+e^-\pi^0$ compared to BABAR [1] (left) and CLEO [48] (right).

Double-octet model is used in these simulations

Some calculations as exercise:





 $\mathcal{M}[\gamma^*(q_1,\nu)\,\gamma^*(q_2,\beta)\to\mathcal{P}] = e^2\epsilon_{\mu\nu\alpha\beta}q_1^{\mu}q_2^{\alpha}F_{\gamma^*\gamma^*\mathcal{P}}(t_1,t_2)$

where $\epsilon_{\mu\nu\alpha\beta}$ is the totally antisymmetric Levi-Civita tensor

$$\begin{split} F_{\gamma^*\gamma^*\mathcal{P}}(t_1,t_2) &= F_{\gamma^*\gamma^*\mathcal{P}}(t_2,t_1) & \text{Bose symmetry of photons} \\ (q_1^2 &= t_1, \ q_2^2 &= t_2) \end{split}$$

$$F_{\gamma^*\gamma^*\pi^0}(t_1, t_2) = -\frac{N_c}{12\pi^2 f_\pi} + \sum_{i=1}^n \frac{4\sqrt{2}h_{V_i}f_{V_i}}{3f_\pi} t_1 \left(D_{\rho_i}(t_1) + D_{\omega_i}(t_1) \right) + \sum_{i=1}^n \frac{4\sqrt{2}h_{V_i}f_{V_i}}{3f_\pi} t_2 \left(D_{\rho_i}(t_2) + D_{\omega_i}(t_2) \right) - \sum_{i=1}^n \frac{4\sigma_{V_i}f_{V_i}^2}{3f_\pi} t_2 t_1 \left(D_{\rho_i}(t_2)D_{\omega_i}(t_1) + D_{\rho_i}(t_1)D_{\omega_i}(t_2) \right)$$

where n= number of vector meson resonance octets

 f_{π} = pion decay constant

Nc = number of quark color

 f_{ij} = coupling for vectors representation of the spin-1 fields for a fixed octet

 $D_V(Q^2) = [Q^2 - M_V^2 + i\sqrt{Q^2}\Gamma_{tot,V}(Q^2)]^{-1}$ is the vector meson propagator

Phys.Rev. D85 (2012) 094010

...let's continue our calculation!



$$\lim_{t_1 \to -\infty} F_{\gamma^* \gamma^* \mathcal{P}}(t_1, t_2) \Big|_{t_2 = const} = 0$$

In our case
$$t_2 = 0$$
. It implies:

$$\sqrt{2}h_{V_i}f_{V_i} - \sigma_{V_i}f_{V_i}^2 = 0, \quad i = 1, \dots, n$$

In the double-octet model used in these simulations, we have:

 $f_{V_1} = 0.20173(86)$ derived from: $\Gamma(\rho \to ee) = \frac{e^4 M_{\rho} f_{V_1}^2}{12\pi}$

Then it is possible to calculate: $h_{V_1} = 0.03121(14)$

$$h_{V_2} f_{V_2} = \frac{3}{32\pi^2 \sqrt{2}} - h_{V_1} f_{V_1} = 0.42(5) \times 10^{-3}$$

$$h_{V_3} f_{V_3} = \frac{3}{32\pi^2 \sqrt{2}} - h_{V_1} f_{V_1} - h_{V_2} f_{V_2} \approx 7.45 \times 10^{-3}$$

Phys.Rev. D85 (2012) 094010



For convenience, the authors of EKHARA define the slope of the transition form factor as:

$$a_{\mathcal{P}} = \frac{1}{F_{\gamma^*\gamma^*\mathcal{P}}(0,0)} \frac{dF_{\gamma^*\gamma^*\mathcal{P}}(t,0)}{dx} \Big|_{t=0} \qquad P = \text{pseudoscalar meson} \\ t = -Q^2 \\ x \equiv t/m_{\mathcal{P}}^2$$

$$a_{\pi} = \frac{16\sqrt{2}\pi^2 m_{\pi}^2}{N_C} \sum_{i=1}^n h_{V_i} f_{V_i} \left(\frac{1}{M_{\rho_i}^2} + \frac{1}{M_{\omega_i}^2} \right)$$
Coefficients evaluated in 2-octet model

TABLE III. Model prediction for the slope parameters $a_{\mathcal{P}}$ and two most recent experimental values. The "2 octets" column is calculated with the parameter values given by our global fit. The first error in experimental value is due to statistics and the second one is systematics.

a _π	1 octet 0.03003(1)	2 octets 0.02870(9)	Experiments		
			0.026(24) (48)	0.025(14) (26) sindrum:	
a_{η}	0.546(9)	0.521(2)	0.576(105) (39)	0.585(18) (13) NA62	
$a_{\eta'}$	1.384(3)	1.323(4)	•••	•••	

How does it work?



- In the source code flags are setup
- You can change those from a text file, depending on your analysis:
- 1. Tagging angle 2. Energy in c.m. 3. Min, max energy 4. Amplitude model (5 models are available) 5. Which final state $(\pi^0, \eta, \eta', \pi^+\pi^-)$ 6. Matrix (t, s, t+s)
- The source code is written in fortran, with quadrupole precision variables In order to run it, you need to install intel-fortran, which is not free for academic purposes

A free demo is available for a maximum period of 30 days. http://software.intel.com/en-us/intel-compilers



• Example of the EKHARA performance: energy c.m flag is set up to 3.77 GeV (Ψ ")

	$e+e- \rightarrow e+e-\gamma\gamma \rightarrow e+e-\pi^0$	e+e- →e+e-γγ →e+e-η	$e+e- \rightarrow e+e-\gamma\gamma \rightarrow e+e-\eta'$
EKHARA simulation	(nb)	(nb)	(nb)
Non tagged	(832.2± 2.9)x 10 ⁻³	(297.2 ± 1.0)x 10 ⁻³	(212.2 ± 1.1)x 10 ⁻³
Tagged e+ 21.6<θ<158.4	(6.672 ± 0.059)x 10 ⁻³	(5.240± 0.019)x 10 ⁻³	(6.776 ± 0.039)x 10 ⁻³
Double tagging	(2.020 ± 0.014)×10 ⁻⁴	$(1.451 \pm 0.010) \times 10^{-4}$	(3.613± 0.025)×10 ⁻⁴

No detector acceptance No radiative corrections are included

- Small cross section
- A drastic cross section reduction simulated when tagging one lepton This number is even smaller if 2 leptons are both tagged: measurement much cleaner, but high statistics needed in such a case



- BaBar could not check very low Q² values due to the trigger
- Simulations show that in BESIII it is possible
- Q² ∈[0÷ 2] GeV² is theoretically the best range to test hadronic LBL correction to (g-2).





Some examples: FF generated at E=3.77 GeV





No detector acceptance No radiative corrections are included

Single-tag mode

Test to verify the EKHARA interface





Figure 4: Studies performed with the MC generator EKHARA for the channel $e^+e^- \rightarrow e^+e^-\pi^0$, with tagged electrons: the momentum and energy of the tagged lepton is compared

EKHARA fortran / EKHARA in BES, true values: distributions must be aligned o 1

Some example: generated 4-momomentum distributions





Figure 8: Studies performed with the MC generator EKHARA for the channel $e^+e^- \rightarrow e^+e^-\pi^0$, with tagged electrons: the momentum and energy of the tagged lepton is shown



No detector acceptance No radiative corrections are included Single-tag mode



Figure 7: Distribution of the momentum transfer (Q²) performed with the MC generator EKHARA for the channel $e^+e^- \rightarrow e^+e^-\pi^0$, with tagged positrons, in both cases: a) EKHARA v2.1 standalone code; b) EKHARA v2.1 in BOSSv. 6.6.1, respectively for the non-tagged lepton (left) and tagged lepton (right).

Some example: azimuthal angular correlation 🧹





This analysis can be performed only if the momentum resolution is very good! 18

ISR corrections



- Radiative corrections are not implemented in EKHARA problems with numerical integration occur not easy to solve
- A MC generator where the rad. corr. are implemented is RABHAT (old fortran code) used from the Belle Collaboration for the analysis $e^+e^- \rightarrow e^+e^-\pi^0$
- The main background of this analysis is the VCS: $e^+e^- \rightarrow e^+e^-\gamma$
- A random photon can combine with one photon in e⁺e⁻γ and peak close to the π⁰ mass. Or a hard photon can be emitted from the electron and leads difference in Q².
 To reduce these sources of backgrounds it is possible to build *ad-hoc* variables (lesson learned by BaBar, PRD 80, 052002 (2009)): *the variables r* and *cosθ*_{EP}





Examples from BaBar





Examples from Belle



The Belle Collaboration chose the MC generator Rabhat to simulate events with a virtual Compton configuration

K. Tobimatsu and Y. Shimizu, Comp. Phys. Comm. 55, 337 (1989).

- Rabhat simulates the radiative Bhabha process with t-channel mass singularity
- The singularity occurs in a topology where the virtual photon, which is emitted in an extremely forward angle from one of the incident electron, is hard scattered off the other electron

Rabhat adopts a particular solution for integration variables and numerical integration



More on radiative corrections: MC generator BabaYaga@NLO



Approach from Pavia: MC generator Babayaga

- It can generate events for QED processes, mainly for e^+e^- colliders in the Q² range 1-10 GeV²
- It is a general MC generator for flavor factories for the processes: e^+e^- , $\mu^+\mu^-$, $\gamma\gamma$, $\pi^+\pi^-$.
- It is accommodated for LO, NLO and higher orders, which can have important impact on the cross section measurement
- For radiative corrections: QED Parton Shower model

http://www2.pv.infn.it/~hepcomplex/babayaga.html

DOCUMENTATION

BabaYaga@NLO Eur.Phys.J.C71:1680,2011

This version is a complete rewriting of the generator, which now could also simulate the production of a single dark matter massive photon decaying in e^+e^- , $\mu^+\mu^-$ pairs.

The code includes also non-log order alpha corrections, which were the main source of theoretical error in the previous releases of the code.

The theoretical error of the new release is estimated to be around 0.1%, for Standard Model processes.

What about PANDA?



- Anti-proton against a fix target: process of interest $p\overline{p} \rightarrow e^+e^-$
- Clearly, the calculations are different compared to e^+e^- colliders
- My attention has been to the decay of Y(4160) and Y(4260) to e^+e^-

 $p\overline{p} \to Y \to e^+e^-$

The goal of this study is:

- Study a rare process and determine the BR with precision
- Study interference between the Y states, decaying to a common final state

 $V \rightarrow e^+ e^-$



Vector state	BR(→e+e-)	Width (MeV))		
ψ(4040)	(1.07 ±0.16) ×10 ⁻⁵	80 ± 10	Vecto	rs: J ^{PC} = 1	
ψ(4160)	(8.1 ±0.9) ×10 ⁻⁶	103 ± 8	Expection	cted to decay	y to e⁺e⁻
Y(4260)		108 ± 12	Only 2	2 measured:	very low BR
Y(4360)		74 ± 18	Large	e width, <mark>prob</mark> a	<u>ably they interfere</u>
Y(4660)		48 ± 15	PANI	DA can do b	etter than
					similar Y(4160)/day
	4040	4160 4	260	4360	mass (MeV/c²)

Mitglied in der Helmholtz-Gemeinschaft.

* in high resolution mode

 $V \rightarrow e^+ e^-$



Vector sta	te BR(→e+e-) Width (Me	eV)		
ψ(4040) (1.07 ±0.16) ×1	$.0^{-5}$ 80 ± 10	Vec	ctors: $J^{PC} = 1^{-1}$	-
ψ(4160) (8.1 ±0.9) ×10	$^{-6}$ 103 ± 8	Exp	pected to deca	ay to e+e-
Y(4260)	108 ± 12	Onl	y 2 measured	1: very low BR
Y(4360)	74 ± 18	Lar	ge width, <mark>prol</mark>	bably they interfere
Y(4660)	48 ± 15	■ PA	NDA can do l	petter than
					similar Y(4160)/day
	4040	4160	4260	4360	mass (MeV/c²)

Due to Bremmstrahlung we expect a long tail on the left of the resonant state

* in high resolution mode



Remarks



I tried to use several MC generators in processes with e⁺e⁻(X) in the final state: EKHARA, PHOTOS. RABHAT: work in progress.

BABAYAGA @ NLO: I have still to try

All these MC generators are optimized for analysis at e⁺e⁻ colliders

Implementing EKHARA in a new framework was trial, but feasible: main problems due to the fortran version (ifort) and conversion of quadrupole precision variables to C++ Problems solved in BES using cfortran.h lib, and adding one subroutine to the original source code

Radiative corrections are still not implemented in EKHARA

Implementation of MC generators with correction N(N)L0 is tricky due to problems of numerical integration (Rabhat): divergences not easy to cancel
 Solution for numerical integration: BASES/SPRING (tested at SLAC, KEK, CERN) http://www.cpc.cs.qub.ac.uk/
 When using EKHARA in BES, specific variable were built to reduce VCS...

- Radiative corrections are important even in $p\overline{p}$ processes: what is here the estimation?
- If any algorithm is available, I can give help to introduce it in PANDA and make tests

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Thanks for your attention!