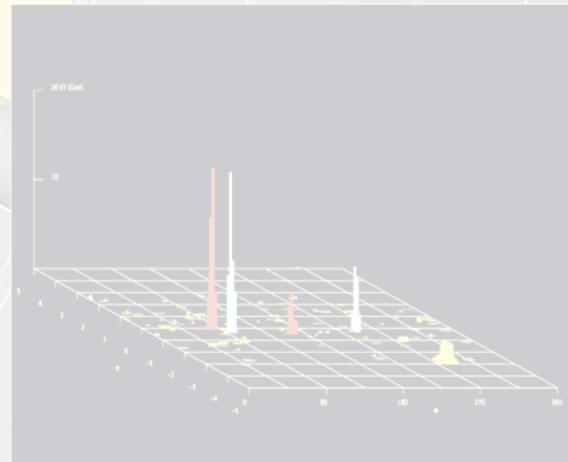
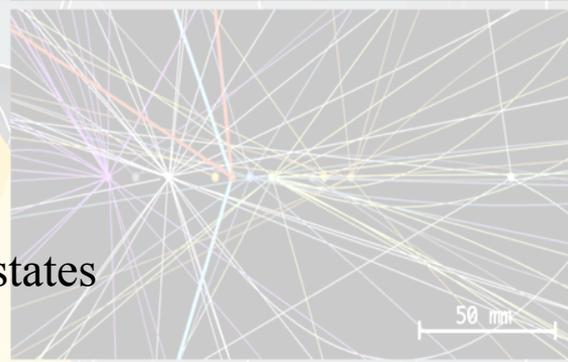
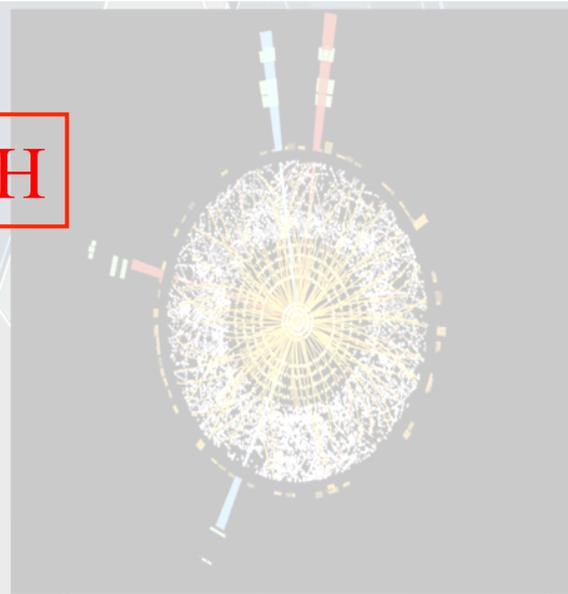


# Search for FCNC in $t \rightarrow cH$

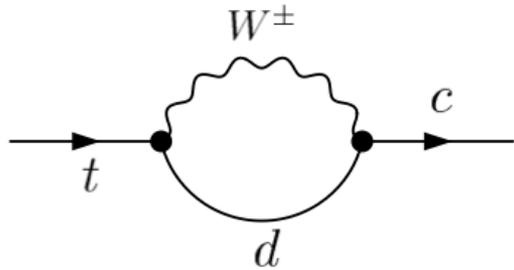
- Motivation
- Analyses
  - Di-photon analysis
  - Inclusive multi-lepton final states
- Conclusion



# Motivation

Why search for top decays to light up type quark + X,  $t \rightarrow cX$  ?

- ✓ Flavour changing neutral current (FCNC) involving light u, c quarks are highly suppressed in the standard model : **GIM mechanism**



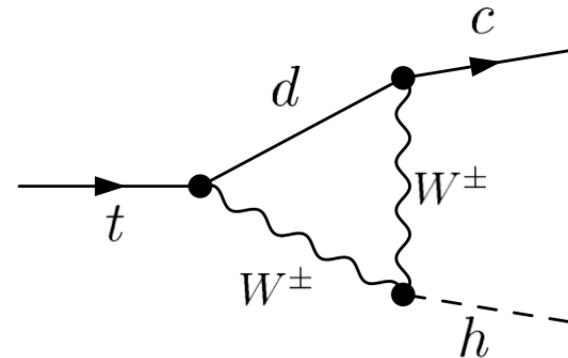
$$\Gamma(t \rightarrow cV) \propto |V_{cb}|^2 \alpha_{QED}^2 \alpha m_t \left(\frac{m_b}{m_W}\right)^4$$

+ Z,  $\gamma$ , (g) line attached to any (quark) line

Loop controlled by down quark  $\sim m_b$   
 $\Rightarrow$  much more suppressed than FCNC in the down sector (e.g.  $b \rightarrow s \gamma$ )

- $V = g : \alpha = \alpha_s, \text{Br} \sim 10^{-12}$
- $V = Z/\gamma : \alpha = \alpha_{QED}, \text{Br} \sim 10^{-14}-10^{-13}$

Even worth for the decay to Higgs boson  
 $\text{Br} \sim 10^{-15}$



Is this generic in most of popular BSM models ?

🦉 Additional heavy quarks, e.g. vector like quarks :

3x3 CKM matrix no longer unitary

⇒ GIM mechanism to suppress SM amplitudes relaxed

tree level flavour changing gauge and scalar interactions

E.g. for  $Q = 2/3$  up type singlet :

$$B(t \rightarrow cZ) \sim < 10^{-4}$$

$$B(t \rightarrow cH) \sim < 4 \cdot 10^{-5}$$

( $B(t \rightarrow c(\gamma/g))$  less enhanced)

(down type  $Q = -1/3$  singlet gives much less enhancements because of more constraints in the down quark sector)

🦄 **Two Higgs Doublet Models** : prejudice from *Natural Flavor Conservation*

Avoid fine tuning (small couplings) to get rid of FCNC by :

⇒ coupling a single doublet to fermions

**2HDM-I**

⇒ coupling up-type fermion to a doublet and down type to another one

**2HDM-II (SUSY)**

⇒ No tree level FCNC, but large enhancements still possible ( $B \sim 10^{-5}$ )  
especially in SUSY-QCD from tree level FC gluino couplings  
(more and more constrained by direct SUSY searches...)  
or in general 2HDM-II, at high  $\tan\beta$  (purely EW contribution)

However small couplings can be *natural* : NFC not really needed

e.g. Cheng-Sher scenario (inspired from Fritzsche ansatz) : mass matrix given by

$$M_{ij} = \Delta_{ij} \sqrt{m_i m_j} \quad \Delta_{ij} \sim O(1)$$

Leading to Yukawa couplings  $\lambda_{ij} \sim \frac{\sqrt{m_i m_j}}{v} \sim \frac{\sqrt{m_i m_j}}{2m_W} g$

⇒ **2HDM-III**

less well defined than the two others, tree level FCN scalar interactions  
but still no gauge FCNC (GIM intact)

For  $(i,j) = (t,c)$ ,  $\lambda \sim 10\% g$ , not a small coupling !

For light quarks, light masses involved ⇒ small couplings and constrains from flavour OK

## So, why search for top decays to light up type quark + X, $t \rightarrow cX$ ?

- ✓ LHC Run I : observation of a new (scalar) boson H with mass  $m_H \sim 125.5 \text{ GeV}/c^2$   
 $m_t > m_H + m_c \Rightarrow$  why not ? 😊
- ✓ FCNC are experimentally less constrained (single top production at LEP, Hera, Tevatron D-D mixing,  $t \rightarrow q Z/\gamma$  at Tevatron) than the ones in the down type sector
- ✓ LHC is a top factory, that can be used to search for such processes, e.g. in top decays

$$t \rightarrow q g/\gamma/Z/H$$

	SM expectation	Max expect. in some exotics (*)	Limits on Br, % (95% CL)
$t \rightarrow c(u) g$	$\sim 5 \cdot 10^{-12}$ ( $4 \cdot 10^{-14}$ )	$\sim 2 \cdot 10^{-4}$ ( $2 \cdot 10^{-4}$ )	Direct : very hard at LHC Search for single top strong production instead : $7.6 \cdot 10^{-5}$ ( $1.5 \cdot 10^{-5}$ ), ATLAS 8TeV, $14.2 \text{ fb}^{-1}$
$t \rightarrow c(u) \gamma$	$\sim 5 \cdot 10^{-14}$ ( $4 \cdot 10^{-16}$ )	$\sim 2 \cdot 10^{-6}$ ( $2 \cdot 10^{-6}$ )	hard
$t \rightarrow c(u) Z$	$\sim 10^{-14}$ ( $10^{-16}$ )	$\sim 10^{-4}$ ( $10^{-4}$ )	0.07 / 0.73 (8 TeV CMS $19.5 \text{ fb}^{-1}$ / 7 TeV ATLAS $2.1 \text{ fb}^{-1}$ )
$t \rightarrow c(u) H$	$\sim 3 \cdot 10^{-15}$ ( $2 \cdot 10^{-17}$ )	$\sim 10^{-3}$ ( $10^{-5}$ )	-

$\Rightarrow$  Any observation of such processes is a non ambiguous sign of new physics  
 Some models predict enhancement by several order of magnitude  
 (not to the % level though...), largest enhancement from **2HDM models** (especially type III)

\* J. A. Aguilar-Saavedra, hep-ph/0409342

## Playing with di-photon events (ATLAS) :

Search for FCNC in top decays  $t \rightarrow cH$ , followed by  $H \rightarrow \gamma\gamma$

Pair production of top quarks are considered :

one of the top decays to  $cH$ , the other decays to  $Wb$

$$pp \rightarrow t\bar{t} \rightarrow W^+ b H c \rightarrow f\bar{f}' b \gamma\gamma c + c.c.$$

For  $H = H_{SM}$  and  $B(t \rightarrow cH) = 1\%$ ,

expect  $N \sim 260$  produced events in the LHC RunI data set

A rather straightforward analysis : start from the standard  $H \rightarrow \gamma\gamma$  inclusive selection  
 $\rightarrow$  two high  $p_T$  (30/40 GeV/c) isolated photons (tight identification criteria)  
 $\rightarrow$  add cuts on jets and invariant masses to fully reconstruct the final state  
 (one W (lept. or hadr.), one Higgs boson, two tops)

Initial  $\gamma\gamma$  sample :

$\sim 24K$  (7 TeV) +  $\sim 119 K$  (8 TeV) events

with  $m_H \in [100, 160] \text{ GeV}/c^2$

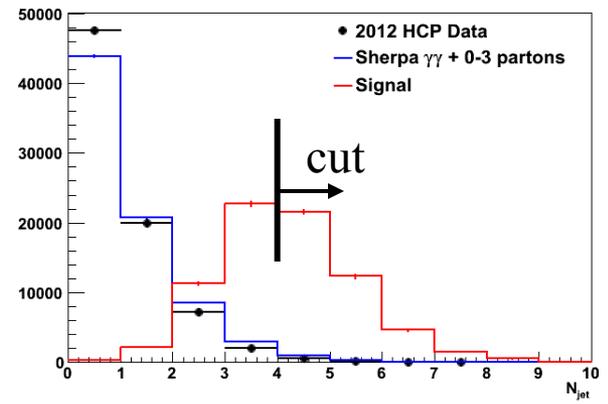
$\oplus$

- three jets : the **hadronic channel** (at 7 and 8 TeV)
- a lepton (e/ $\mu$ ), a neutrino and a b-jet : the **leptonic channel** (only at 8 TeV)

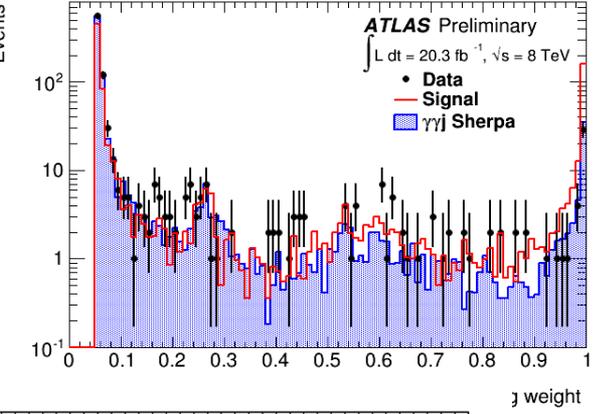
# Hadronic channel

At least 4 jets  
 At least one b-tag (70% b-efficiency)  
 No lepton  
 From the 4 leading jets :  
 $m_{\gamma\gamma j} \in [156, 191] \text{ GeV}/c^2$   
 $m_{\varphi\varphi j} \in [130, 210] \text{ GeV}/c^2$   
 (if 5<sup>th</sup> or 6<sup>th</sup> b-tagged, and 4<sup>th</sup> or 5<sup>th</sup> not, use it instead)

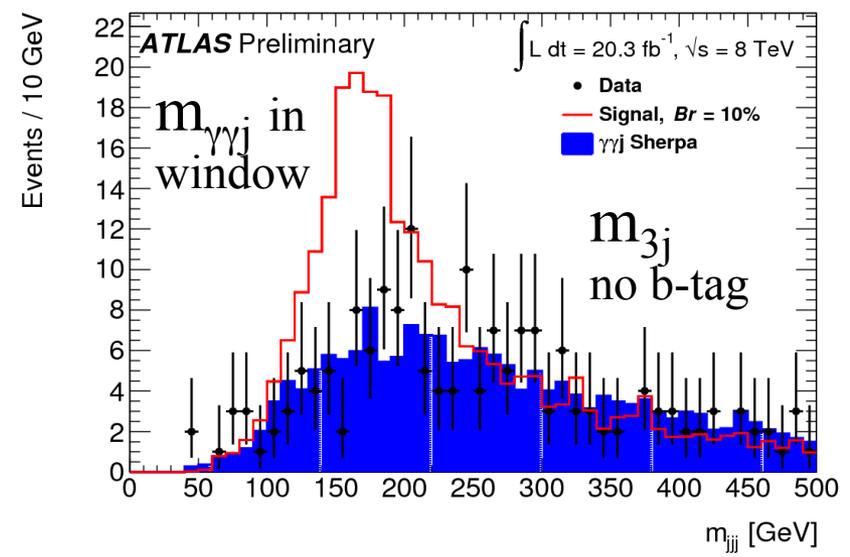
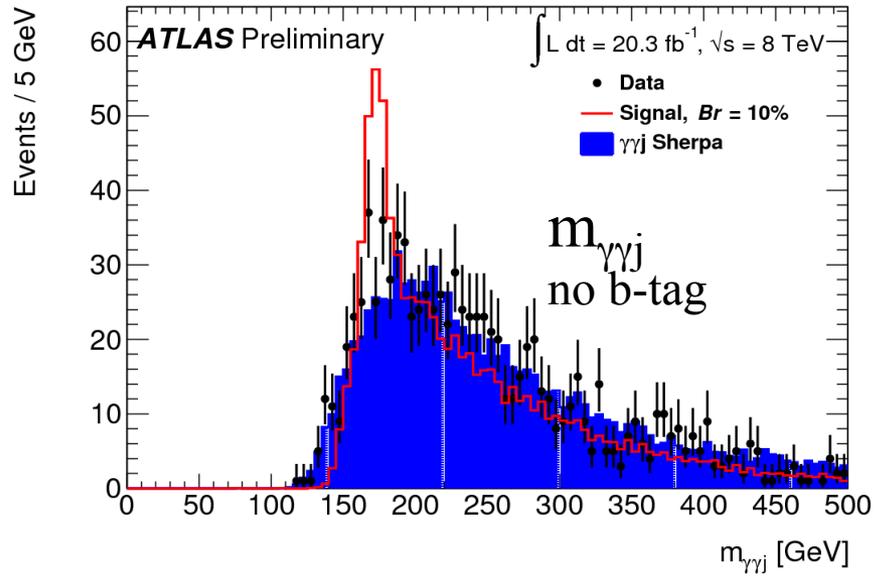
## Jet multiplicity



b-tag discriminant  
 $(N_{\text{jet}} > 3)$



Do not intend to predict the background from MC  
 but decent agreement between data and Sherpa MC shapes



# Leptonic channel

One lepton, at least 2 jets

At least one b-tag (70% b-efficiency)

No other lepton

$m_T(l, E_T^{\text{miss}}) > 30 \text{ GeV}/c^2$

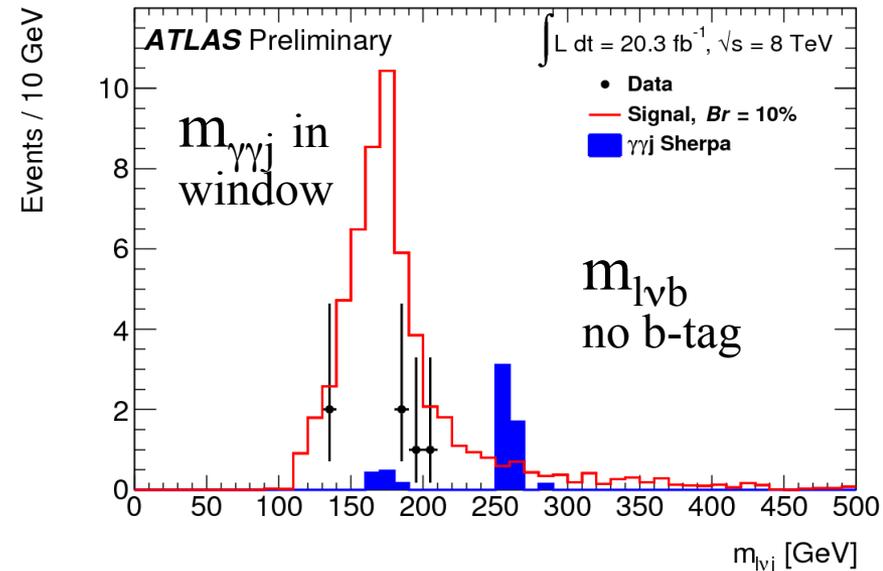
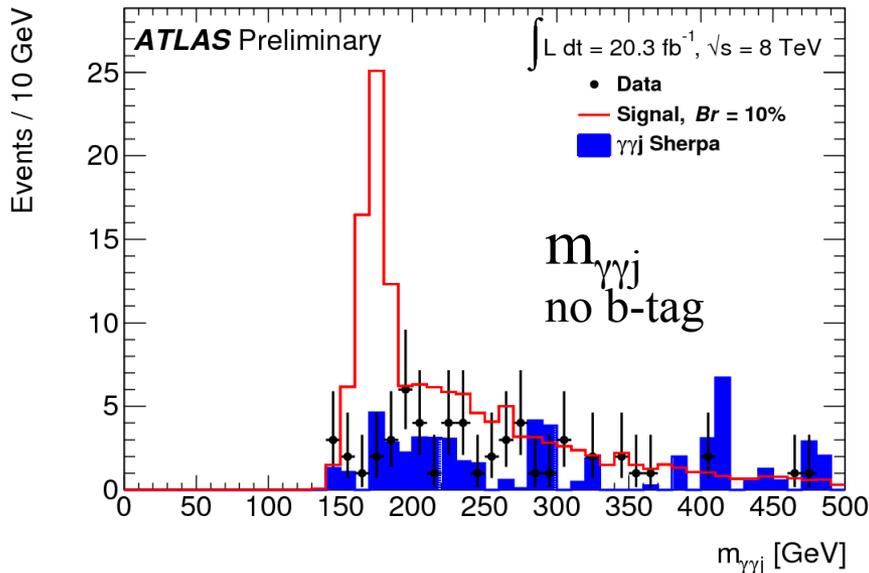
From the 2 leading jets :

$m_{\gamma\gamma j} \in [156, 191] \text{ GeV}/c^2$

$m_{l\nu j} \in [135, 205] \text{ GeV}/c^2$

( $p_z^{\nu}$  from W mass constrain)

(if 3<sup>rd</sup> or 4<sup>th</sup> b-tagged, and 2<sup>nd</sup> or 3<sup>rd</sup> not, use it instead)



## Results :

### 🦉 **Hadronic channel** (7+8 TeV combined) :

$$N_{\text{FCNC}}(\text{B}(t \rightarrow \text{cH}) = 1\%) = 10.9 \pm 0.8_{\text{theory}}$$

$$N_{\text{H}}^{\text{SM}} = 0.28 \pm 0.10_{\text{theory+lumi}}$$

$$N_{\text{obs}} = 50$$

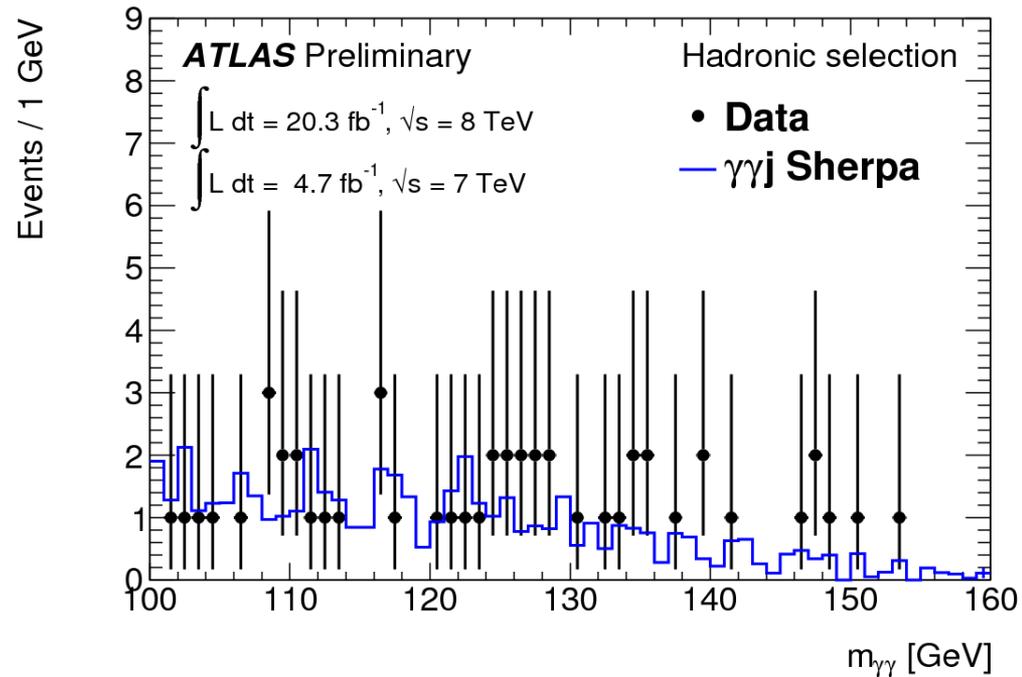
### 🦉 **Leptonic channel** (8 TeV) :

$$N_{\text{FCNC}}(\text{B}(t \rightarrow \text{cH}) = 1\%) = 2.9 \pm 0.2_{\text{theory}}$$

$$N_{\text{H}}^{\text{SM}} = 0.05 \pm 0.01_{\text{theory+lumi}}$$

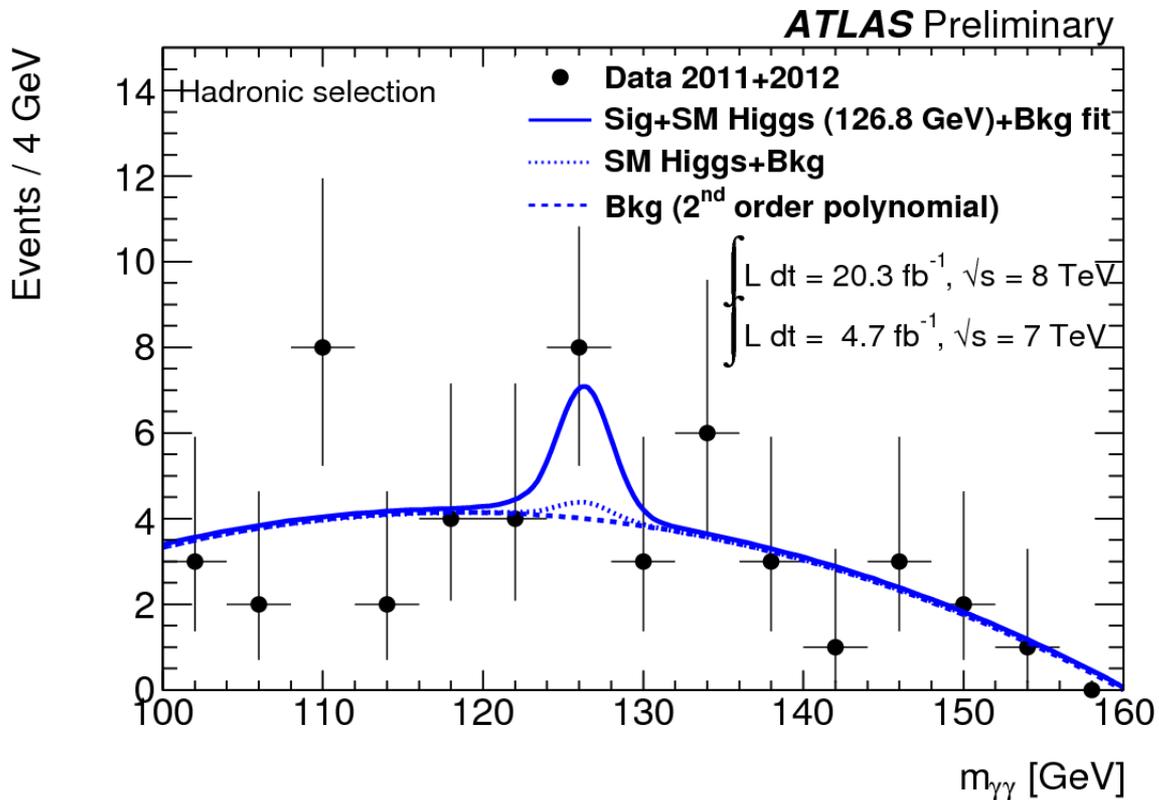
$$N_{\text{obs}} = 1, m_{\gamma\gamma} = 147 \text{ GeV}/c^2$$

- 🦉 **Systematic uncertainties** : dominated by photon ID and isolation,  
(increased w.r.t. standard H analysis to account for busier environment)  
+ jet energy scale, b-tagging and underlying event modelisation



## Statistical interpretation :

- Hadronic channel : enough events to determine background under the peak from fit to the  $m_{\gamma\gamma}$  distribution :  
bkg shape = 2<sup>nd</sup> order polynomial / signal shape : Crystal-Ball + wide Gaussian
- Leptonic channel, two bin shape analysis : constrain bkg expectation in signal region ( $[123,129]$  GeV/ $c^2$ ) from control region  $[100,123[U]129,160]$  GeV/ $c^2$



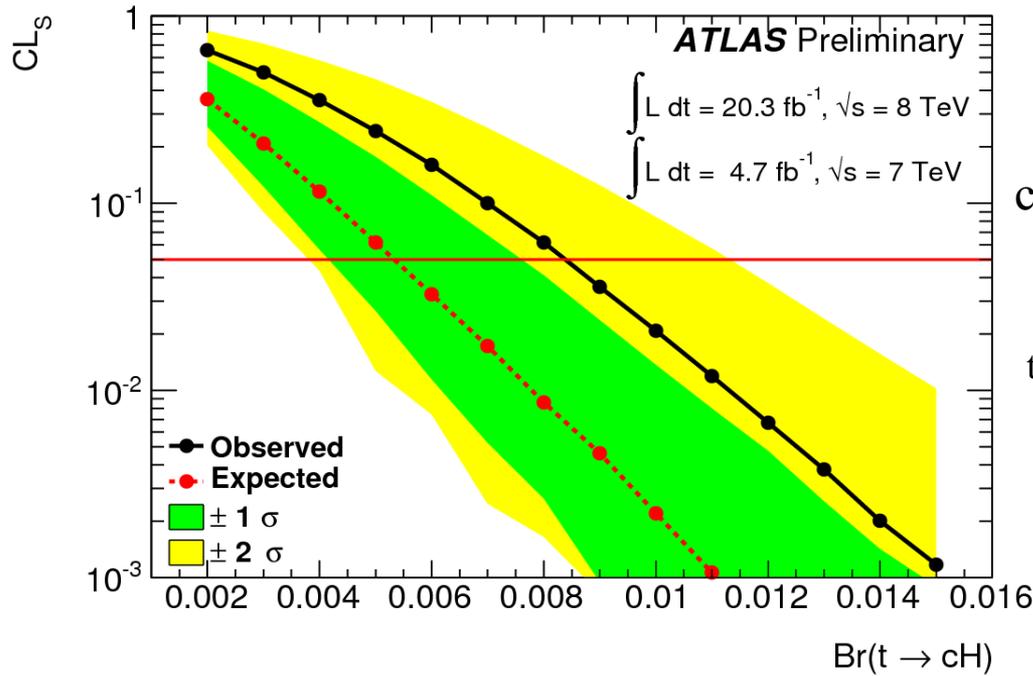
Fixing  $m_H = 126.8 \text{ GeV}/c^2$   
fitted branching ratio :

$$Br = 0.26^{+0.31}_{-0.26}\%$$

corresponding to a signal yield of

$$N_S = 3.7^{+4.4}_{-3.7}$$

No signal observed  $\Rightarrow$  limit



$B(t \rightarrow cH) < 0.83\% @ 95\% \text{ CL}$   
 (0.53% expected)

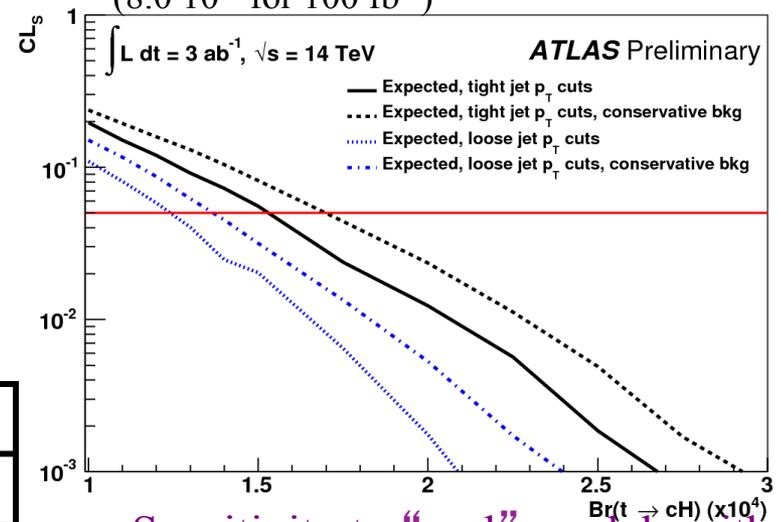
corresponding to a limit on the  $tcH$  coupling  
 $\lambda_{tcH} < 1.91 B^{0.5} \sim 0.17$  (0.14 expected)

to be compared to top quark Yukawa coupling

$$\lambda_{tcH} \sim 1$$

Might expect

$B(t \rightarrow cH) < 1.5 \cdot 10^{-4} \% @ 95\% \text{ CL}$   
 ( $8.0 \cdot 10^{-4}$  for  $100 \text{ fb}^{-1}$ )



Sensitivity to “real” models with  
 $H \rightarrow \gamma\gamma$  channel only !

Prospects at HL-LHC ( $14 \text{ TeV}/3\text{ab}^{-1}$ )

- ✓ signal cross-section ( $\sigma_{tt}$ ) : x 3.9
- ✓ non resonant bkg  $\sigma$  : x 2.2
- ✓ resonant bkg ( $ttH$ ) : x 4.7
- ✓ but needs tighter cuts (higher pile-up)

# events in  
 signal window

	FCNC	SM Higgs	Bkg
Hadronic	13	24	350
Leptonic	7	14	25

Re-interpreting multi lepton searches (CMS) :

- Use the weak boson and tau decays of H, in multi-lepton final states
    - more events ( $B(H \rightarrow WW^* \rightarrow e/\mu \nu e/\mu \nu) \sim 1\% + ZZ^* + \tau^+\tau^-$ )
    - more difficult to interpret than the di-photon channel in case of an excess (no full reconstruction)
- ⇒ better for limit setting...

3 leptons  
 + a number of categories defined from

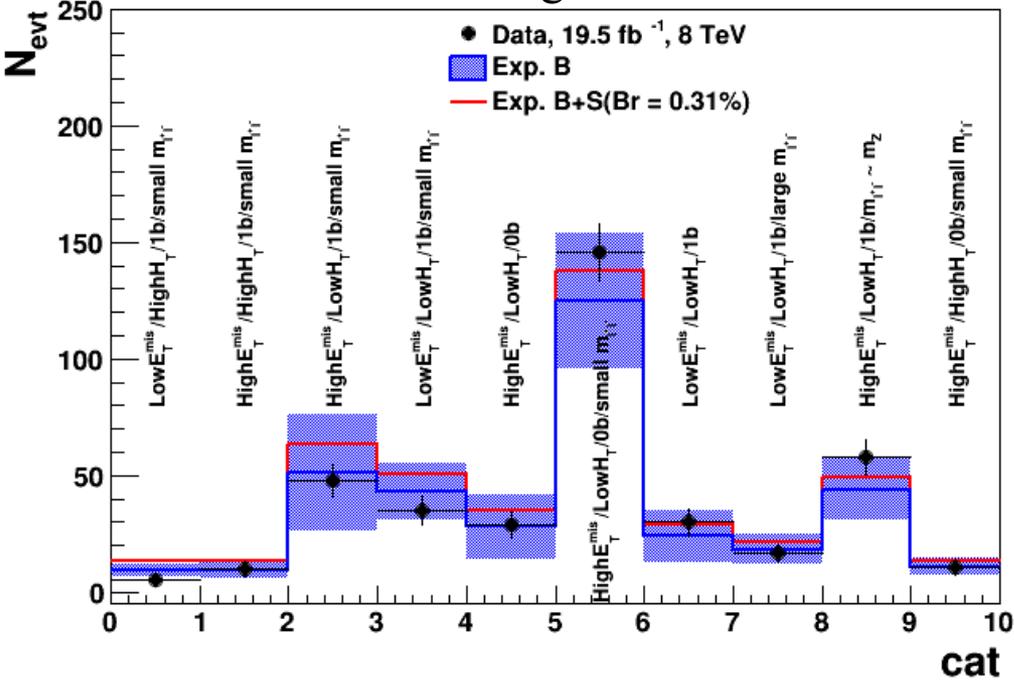
- ✓ number of b-tagged jets (0 / ≥1)
- ✓  $H_T$  (sum of all good jet  $p_T$ )
- ✓  $E_t^{mis}$
- ✓ the consistency of an Opposite Sign Same Flavour lepton pair with a Z

No excess seen ⇒ limit

$B(t \rightarrow cH) < 0.31\% @ 95\% CL$   
 (0.31% expected)

(essentially from  $H \rightarrow WW^*$  : if no  $ZZ^*/\tau^+\tau^-$  :  $B < 0.37\%$ )

Comparison Data / Bkg prediction for the 10 most sensitive categories :



corresponding to a limit on  $tcH$  coupling  
 $\lambda_{tcH} < 0.1$

## Conclusion

- SM yields for FCNC  $t \rightarrow cX$  process are extremely small (beyond reach of any experiment)
- Some not-so-unlikely models of BSM predict branching ratios within the reach of LHC : e.g. in 2HDM-III,  
 $B(t \rightarrow cH)^*$  up to  $10^{-3}$  can be obtained

⇒ spectacular signatures : Higgs feed down from top pair production

The sensitivity at LHC runII will be greatly improved thanks to the  $\sim 4x$  enhancement of the top pair production cross-section

- Some of these models have also direct impact on the Higgs boson phenomenology (e.g.  $H \rightarrow cc$  might be as large as  $H \rightarrow bb$ ,  
 $H \rightarrow gg / \tau^+\tau^-$  might be modified significantly)

\*  $t \rightarrow uH$  is also possible  
ATLAS search sensitive to this decay