### **DM with supersymmetric triplets**

CA, V.Martín-Lozano and G.Nardini, arXiv:1403.6434

**Chiara Arina** 



### GDR Terascale@Palaiseau June 4<sup>th</sup> 2014

# Why Triplet extension of the MSSM?

- Little hierarchy problem within the MSSM to accommodate  $m_h$ = 126 GeV
- Extension of the MSSM with SINGLETS or TRIPLETS mitigates the problem

- Higgs is SM-like except in the loop-suppressed processes
- Triplets are particularly interesting for that



# The Triplet extension of the MSSM

$$\Sigma = \begin{pmatrix} \xi^0 / \sqrt{2} & -\xi_2^+ \\ \xi_1^- & -\xi^0 / \sqrt{2} \end{pmatrix} \qquad Y = 0 \ SU(2)_L \text{-triplet superfield}$$

$$W_{\text{TMSSM}} = W_{\text{MSSM}} + \lambda H_1 \cdot \Sigma H_2 + \frac{1}{2} \mu_{\Sigma} \operatorname{Tr} \Sigma^2$$

$$\begin{split} \mathcal{L}_{\text{TMSSM}_{\text{SB}}} &= \mathcal{L}_{\text{MSSM}_{\text{SB}}} + \frac{m_4^2}{4} \operatorname{Tr}(\Sigma^{\dagger} \Sigma) \\ &+ \left[ \frac{B_{\Sigma}}{2} \operatorname{Tr}(\Sigma^2) + \frac{\lambda A_{\lambda}}{4} H_1 \cdot \Sigma H_2 + \text{h.c.} \right] \end{split}$$

A.Delgado, G.Nardini and M.Quiros '12, '13

TMSSM considered in a definite configuration to introduce the minimal amount of new parameters and satisfy a SMlike Higgs

Scalar triplet vev is constrained by electroweak parameters

 $\langle \xi^0 
angle \lesssim 4 \, {
m GeV}$ 

$$|A_{\lambda}|,\,|\mu|\,,|\mu_{\Sigma}|\lesssim 10^{-2}rac{m_{\Sigma}^2+\lambda^2v^2/2}{\lambda v}$$
EW scale

• Scalar triplet vev is constrained by electroweak parameters  $\langle \xi^0 
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m GeV}$ 

$$egin{aligned} A_{\lambda}(, \left|\mu
ight|, \left|\mu_{\Sigma}
ight| \lesssim 10^{-2}rac{m_{\Sigma}^{2}+\lambda^{2}v^{2}/2}{\lambda v} \end{aligned}$$
 EW scale  $\int M_{\Sigma}=5\,\mathrm{TeV} \end{aligned}$ 

• Scalar triplet vev is constrained by electroweak parameters  $\langle \xi^0 \rangle \lesssim 4 \, \text{GeV}$ 



After electroweak symmetry breaking the CP-even Higgs mass matrix is

$$\mathcal{M}_{h,H}^2 = \begin{pmatrix} m_A^2 \cos^2\beta + m_Z^2 \sin^2\beta & (\lambda^2 v^2 - m_A^2 - m_Z^2) \sin\beta\cos\beta \\ (\lambda^2 v^2 - m_A^2 - m_Z^2) \sin\beta\cos\beta & m_A^2 \sin^2\beta + m_Z^2 \cos^2\beta \end{pmatrix}$$

• Scalar triplet vev is constrained by electroweak parameters  $\langle \xi^0 \rangle \leq 4 \, \text{GeV}$ 



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• Decoupling limit 
$$m_A \to \infty$$
  
1.  $m_{h,tree}^2 = m_Z^2 \cos^2 2\beta + \frac{\lambda^2}{2} v^2 \sin^2 2\beta$   
2. Higgs is SM-like except in loop-induced processes

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## **Electroweak sector**

additional degrees of freedom in the electroweakino sector • Triplet

$$\mathcal{M}_{\widetilde{\chi}^0}^{tree} = egin{pmatrix} M_1 & 0 & -rac{1}{2}g_1v_1 & rac{1}{2}g_1v_2 & 0 \ 0 & M_2 & rac{1}{2}g_2v_1 & -rac{1}{2}g_2v_2 & 0 \ -rac{1}{2}g_1v_1 & rac{1}{2}g_2v_1 & 0 & -\mu & -rac{1}{2}v_2\lambda \ rac{1}{2}g_1v_1 & -rac{1}{2}g_2v_2 & -\mu & 0 & -rac{1}{2}v_1\lambda \ 0 & 0 & -rac{1}{2}v_2\lambda & -rac{1}{2}v_1\lambda & \mu_\Sigma \end{pmatrix}$$

Neutralino sector relevant for Higgs invisible decay width and for DM

$$\mathcal{M}_{\widetilde{\chi}^{\pm}}^{tree} = egin{pmatrix} M_2 & g_2 v \sineta & 0 \ g_2 v \coseta & \mu & -\lambda v \sineta \ 0 & \lambda v \coseta & \mu_\Sigma \end{pmatrix}$$

Chargino sector relevant for  $\begin{cases} h \to \gamma \gamma \\ h \to Z \gamma \end{cases}$ 

## **Higgs signatures**



$$\begin{split} R_{\gamma\gamma} &= \left| 1 + \frac{A_{\widetilde{\chi}_{1,2,3}}^{\gamma\gamma}}{A_W^{\gamma\gamma} + A_t^{\gamma\gamma}} \right|^2 \\ A_{\widetilde{\chi}_{1,2,3}}^{\gamma\gamma} &= \sum_{i=1}^3 \frac{2M_W}{\sqrt{2} m_{\widetilde{\chi}_i^\pm}} (g_{h\widetilde{\chi}_i^+\widetilde{\chi}_i^-}^L + g_{h\widetilde{\chi}_i^+\widetilde{\chi}_i^-}^R) A_{1/2} (\tau_{\widetilde{\chi}_i^\pm})^2 \\ A_W^{\gamma\gamma} &= -8.3 \qquad A_t^{\gamma\gamma} = 1.9 \end{split}$$

## **Higgs signatures**



$$\begin{split} R_{\gamma\gamma} &= \left| 1 + \frac{A_{\widetilde{\chi}_{1,2,3}}^{\gamma\gamma}}{A_W^{\gamma\gamma} + A_t^{\gamma\gamma}} \right|^2 \\ A_{\widetilde{\chi}_{1,2,3}}^{\gamma\gamma} &= \sum_{i=1}^3 \frac{2M_W}{\sqrt{2} m_{\widetilde{\chi}_i^{\pm}}} (g_{h\widetilde{\chi}_i^+\widetilde{\chi}_i^-}^L + g_{h\widetilde{\chi}_i^+\widetilde{\chi}_i^-}^R) A_{1/2}(\tau_{\widetilde{\chi}_i^{\pm}}) \\ A_W^{\gamma\gamma} &= -8.3 \qquad A_t^{\gamma\gamma} = 1.9 \end{split}$$

## **Higgs signatures**



$$R_{Z\gamma} = \left| 1 + \frac{A_{\widetilde{\chi}_{1,2,3}^{\pm}}^{Z\gamma}}{A_W^{Z\gamma} + A_t^{Z\gamma}} \right|^2$$

$$\begin{split} A_{\widetilde{\chi}_{1,2,3}}^{Z\gamma} &= \sum_{j,k=1}^{3} \frac{g_2 \, m_{\widetilde{\chi}_{j}^{\pm}}}{g_1 \, m_Z} \, f\Big(m_{\widetilde{\chi}_{j}^{\pm}}, m_{\widetilde{\chi}_{k}^{\pm}}, m_{\widetilde{\chi}_{k}^{\pm}}\Big) \, (g_{h\widetilde{\chi}_{j}^{+}\widetilde{\chi}_{i}^{-}}^{L} + g_{h\widetilde{\chi}_{j}^{+}\widetilde{\chi}_{i}^{-}}^{R}) (g_{Z\widetilde{\chi}_{j}^{+}\widetilde{\chi}_{i}^{-}}^{L} + g_{Z\widetilde{\chi}_{j}^{+}\widetilde{\chi}_{i}^{-}}^{R}) \\ A_W^{Z\gamma} &= -12 \qquad A_t^{Z\gamma} = 0.6 \end{split}$$

# Set up of the analysis

#### SUSY Model = TMSSM

SARAH

# SPheno

(masses computed at full 1-loop + Higgs has 2 loop corrections)

#### SPheno, CPSuperH

**Higgs Physics** 

## Sampling method and free parameters

#### $\{\theta_i\} = \{M_1, M_2, M_3, \widetilde{m}, \tan\beta, \mu, \lambda, \mu_{\Sigma}\}$

Sampling with the algorithm MultiNest

$$p(\theta_i|d) \propto \mathcal{L}(d|\theta_i)\pi(\theta_i)$$

- Nested sampling
- Sampling scale as n instead of n<sup>2</sup> as for a random scan

NS parameters	Prior range
$\log_{10}(M_1/\text{GeV}), \log_{10}(\mu_{\Sigma}/\text{GeV})$	$1 \rightarrow 3$
$\log_{10}(\mu/{\rm GeV}), \log_{10}(M_2/{\rm GeV})$	$2 \rightarrow 3$
$\widetilde{m}/{ m TeV}$	0.63  ightarrow 2
$\log_{10}(\tan\beta)$	$0 \rightarrow 1$
$\lambda$	0.5  ightarrow 1.2

$$m_{\Sigma} = 5 \,{
m TeV} \;, \quad A_t = A_b = 0 \;, \quad M_3 = 1.4 \,{
m TeV} \;, \quad m_A = 1.5 \,{
m TeV}$$

# Likelihood and priors



# Running the machinery for Higgs physics ...

















# Set up of the analysis

#### SUSY Model = TMSSM

SARAH

# SPheno

(masses computed at full 1-loop + Higgs has 2 loop corrections)

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**Higgs Physics** 



# Running the machinery with DM constraints ...



Type	Observable	Measurement/Limit
Collider data	$m_h$	$125.85 \pm 0.4 \text{ GeV} (\exp) \pm 3 \text{ GeV} (\text{theo})$
	$\Gamma(Z \to \widetilde{\chi}_1^0 \widetilde{\chi}_1^0)$	$< 2 { m MeV}$
	$m_{\widetilde{t}_1}$	> 650  GeV (LHC 90%  CL)
	$m_{\widetilde{\mathbf{v}}^+}$	> 101  GeV (LEP  95%  CL)
_DM data	$\Omega_{ m DM}h^2$	$0.1186 \pm 0.0031 \text{ (exp) } \pm 20\% \text{ (theo)}$
	$\sigma^{SI}_{ m Xe}$	LUX (90% CL)



Neutralino (with triplino component)

# Direct detection constraint on spin-independent scattering



## **Neutralino as DM**



Complementarity between DM direct detection and colliders

# **Neutralino spin-dependent interaction**



- Xenon100 more sensitive to scattering on neutron
- Spin-dependent scattering on proton is one order of magnitude lower
- Large uncertainties on the nuclear form factors

### **Composition of the neutralino**



800

700

700 800

# Impact of DM on Higgs pheno



## Conclusions

#### Triplet extension of the MSSM has some adgvantages

- (a) Reduces fine-tuning in the Higgs mass (new contribution at tree level)
- (b) Higgs phenomenology is SM-like
- (c) Only loop-induced processes deviate from SM behavior (extra charginos)

 $\Gamma(h \to \gamma \gamma) < 60\%$  $\Gamma(h \to \gamma Z) < 40\%$ 

loop-induced processes are correlated

#### If the lightest neutralino is the DM candidate

(a) Compatible with DM constraints on the Higgs pole or co-annihilation regions

(b) Strong impact on the loop-induces processes because of DD constraints

$$\Gamma(h \to \gamma \gamma), \Gamma(h \to \gamma Z) < 20\%$$

# Back up slides

## More on the Higgs sector

Minimization conditions for electroweak symmetry breaking:

$$\begin{split} m_3^2 &= m_A^2 \sin\beta\cos\beta \ , \\ m_Z^2 &= \frac{m_2^2 - m_1^2}{\cos 2\beta} - m_A^2 + \lambda^2 v^2/2 \ , \\ m_A^2 &= m_1^2 + m_2^2 + 2|\mu|^2 + \lambda^2 v^2/2 \ , \\ m_H^\pm &= m_A^2 + m_W^2 + \lambda^2 v^2/2 \ , \end{split}$$

$$\begin{aligned} \frac{1}{2}(h_2, h_1, x) \,\widehat{\mathcal{M}}^2 \begin{pmatrix} h_2 \\ h_1 \\ x \end{pmatrix} & H_i^0 = v_i + (h_i + i\chi_i)/\sqrt{2} & x = \operatorname{Re}\xi^0/\sqrt{2} \\ & & & \\ \widehat{\mathcal{M}}^2 = \begin{pmatrix} \mathcal{M}^2 & \cdot \\ & \ddots & m_{\Sigma}^2 + \frac{\lambda^2}{2}v^2 \end{pmatrix} & \hat{\mu} = \max\{|\mu|, |\mu_{\Sigma}|, |A_{\lambda}|\} \end{aligned}$$

$$\mathcal{M}_0^2 = \left( egin{array}{c} m_A^2 \cos^2eta + m_Z^2 \sin^2eta & (\lambda^2 v^2 - m_A^2 - m_Z^2)\sineta\coseta \ (\lambda^2 v^2 - m_A^2 - m_Z^2)\sineta\coseta & m_A^2\sin^2eta + m_Z^2\cos^2eta \end{array} 
ight)$$

## More on the Higgs sector

Higgs mass has full 2 loop corrections from RGEs resummation + full one loop mass spectrum + corrections of the order

 $\mathcal{O}(h_t^2 g_3^2) \ \mathcal{O}(h_t^4)$ 

$$\Gamma(Z \to \tilde{\chi}_1^0 \tilde{\chi}_1^0) = \frac{1}{12\pi} \frac{G_F}{\sqrt{2}} m_Z^3 \left( 1 - \frac{4m_{\tilde{\chi}_1^0}^2}{m_Z^2} \right)^{3/2} \left( |N_{13}|^2 - |N_{14}|^2 \right)^2$$

$$\Gamma(h \to \tilde{\chi}_1^0 \tilde{\chi}_1^0) = \frac{G_F m_W^2}{2\sqrt{2}\pi} m_h \left( 1 - \frac{4m_{\tilde{\chi}_1^0}^2}{m_h^2} \right)^{3/2} g_{h\chi_1^0\chi_1^0}^2$$

$$g_{h\chi_1^0\chi_1^0} = (N_{12} - rac{g_1}{g_2}N_{11})(\sineta N_{14} - \coseta N_{13}) + rac{\lambda}{g_2}N_{15}(N_{14}\sineta + N_{13}\coseta)$$

 $m_{\Sigma}^2 \equiv m_4^2 + \mu_{\Sigma}^2 + B_{\Sigma} \mu_{\Sigma}$ 

### **Details on diphoton etc..**



$$g_{Z\chi_{j}^{+}\chi_{k}^{-}}^{L} = -\left(V_{i1}V_{j1}^{*} + \frac{1}{2}V_{i2}V_{j2}^{*} + V_{i3}V_{j3}^{*} - \delta_{ij}s_{W}^{2}\right)$$
$$g_{Z\chi_{j}^{+}\chi_{k}^{-}}^{R} = -\left(U_{i1}U_{j1}^{*} + \frac{1}{2}U_{i2}U_{j2}^{*} + U_{i3}U_{j3}^{*} - \delta_{ij}s_{W}^{2}\right)$$

$$g_{h\tilde{\chi}_{i}^{+}\tilde{\chi}_{j}^{-}}^{L} = \frac{1}{\sqrt{2}} \left[ \left( U_{j1}V_{i2} - \frac{\lambda}{g_{2}}U_{j2}V_{i3} \right) \sin\beta + \left( U_{j2}V_{i1} + \frac{\lambda}{g_{2}}U_{j3}V_{i2} \right) \cos\beta \right]$$
$$g_{h\tilde{\chi}_{i}^{+}\tilde{\chi}_{j}^{-}}^{R} = \frac{1}{\sqrt{2}} \left[ \left( U_{i1}V_{j2} - \frac{\lambda}{g_{2}}U_{i2}V_{j3} \right) \sin\beta + \left( U_{i2}V_{j1} + \frac{\lambda}{g_{2}}U_{i3}V_{j2} \right) \cos\beta \right]$$





# Higgs mass



# Neutralino as DM in the Higgs pole



# Neutralino as DM in the Higgs pole



# Well tempered neutralino



## Well tempered neutralino



## **Neutralino as DM and Higgs**

≺ <sup>1.2</sup>



 $M_2$  [GeV]

9

 $\tan\beta$ 

10

 $R_{\gamma\gamma}$ 

#### Sampling method and free parameters

PARAMETER SPACE with 7 free parameters

$$\{ heta_i\} = \{M_1, M_2, M_3, \widetilde{m}, \taneta, \mu, \lambda, \mu_{\Sigma}\}$$

Sampling with the algorithm MultiNest

- Nested sampling
- Sampling scale as n instead of n<sup>2</sup> as for a random scan
- Based on Bayes theorem



from the pdf distribution, no statistical meaning