

Outline

- **DM observables** calculated by micrOMEGAs:
 - Relic density
 - Direct detection
 - Indirect detection (*Green function method*)
 - Sun/Earth captured DM neutrino signal
- **Dark matter models**
 - MSSM, NMSSM, CPVMSS (*SUSY*)
 - IDM, Z3M, LHM, RHNM,
 - **any user model**
- **Particle decays and cross sections**
 - Higgs decays with $\Upsilon\Upsilon$, gg and virtual Z/W channels
- **Easy interface with other packages:**
 - SuSpect, SoftSUSY, SPHENO, ..., NMSSMTools, CPsuperH (spectra)
 - HIGGSBOUNGS, SUPERISO, (observables)

Guiding principles of micrOMEGAs

- Human make mistakes - computer not
 - Automation
- Several groups are developing specialized codes
 - Interface routines
- Users might want to improve one aspect
 - Modularity
- We do not know what DM is made of
 - Possibility to include different DM candidates
- Models are often complex with huge parameter space
 - Speed of execution
- Ready made, stand-alone package for the non-expert
 - User friendly interface and manual

Matrix elements in micrOMEGAS

MicrOMEGAS uses **CalcHEP** program

arXiv 1207.6082(hep_ph)

for generation of matrix elements. Because **CalcHEP** is designed to work with generic model, **micrOMEGAS** is able to calculate DM related observables for generic model as well.

Model of particle interaction in CalcHEP is presented by 5 tables.

1) Variables

Inert Doublet Model

Variables

| Name | Value | > Comment | < |
|------|---------|-----------------------------------|---|
| EE | 0.31333 | Electromagnetic coupling constant | |
| SW | 0.474 | sin of the Weinberg angle | |
| MZ | 91.187 | Mass of Z | |
| MHX | 111 | Mass of Inert Doublet Higgs | |
| MH3 | 222 | Mass of CP-odd Higgs | |
| MHC | 333 | Mass of charged Higgs | |
| LaL | 0.01 | Coupling in Inert Sector | |

3) Constrained parameters of the model.

Inert Doublet

Constraints

| Name | > Expression |
|------|--------------------------------|
| CW | sqrt(1-SW^2) |
| MW | MZ*CW |
| Mb | MbEff(Q) |
| Mc | McEff(Q) |
| mu2 | MHX^2 - laL*(2*MW/EE*SW)^2 |
| la3 | 2*(MHC^2 - mu2)/(2*MW/EE*SW)^2 |
| la5 | (MHX^2 - MH3^2)/(2*MW/EE*SW)^2 |

%Local!

4) External functions and their prototypes.

Libraries

% - comment

extern type func(type X) % C-prototype

Lines of other type are passed to linker

Particle list

Inert Doublet Model

Particles

| Full Name | P | aP | number | spin2 | mass | width | color | aux |
|---------------|-----|-----|--------|-------|------|-------|-------|-----|
| photon | A | A | 22 | 2 | 0 | 0 | 1 | G |
| Z boson | Z | Z | 23 | 2 | MZ | wZ | 1 | G |
| gluon | G | G | 21 | 2 | 0 | 0 | 8 | G |
| W boson | W+ | W- | 24 | 2 | MW | wW | 1 | G |
| neutrino | n1 | N1 | 12 | 1 | 0 | 0 | 1 | L |
| electron | e1 | E1 | 11 | 1 | 0 | 0 | 1 | |
| mu-neutrino | n2 | N2 | 14 | 1 | 0 | 0 | 1 | L |
| muon | e2 | E2 | 13 | 1 | Mm | 0 | 1 | |
| tau-neutrino | n3 | N3 | 16 | 1 | 0 | 0 | 1 | L |
| tau-lepton | e3 | E3 | 15 | 1 | Mt | 0 | 1 | |
| u-quark | u | U | 2 | 1 | 0 | 0 | 3 | |
| d-quark | d | D | 1 | 1 | 0 | 0 | 3 | |
| c-quark | c | C | 4 | 1 | Mc | 0 | 3 | |
| s-quark | s | S | 3 | 1 | Ms | 0 | 3 | |
| t-quark | t | T | 6 | 1 | Mtop | wtop | 3 | |
| b-quark | b | B | 5 | 1 | Mb | 0 | 3 | |
| Higgs | h | h | 25 | 0 | Mh | !wh | 1 | |
| odd Higgs | ~H3 | ~H3 | 36 | 0 | MH3 | !wH3 | 1 | |
| Charged Higgs | ~H+ | ~H- | 37 | 0 | MHC | !wHC | 1 | 6 |
| second Higgs | ~X | ~X | 35 | 0 | MHX | !wHX | 1 | |

4) List of Feynman rules

Inert Doublet

Lagrangian

| P1 | P2 | P3 | P4 | Factor | $ dLagrangian/dA(p1)dA(p2)dA(p3)$ |
|-----------|-----------|-----------|----------|------------------|-----------------------------------|
| A | W+ | W- | | -EE | $ m3.p2*m1.m2-m1.p2*m2.m3- \dots$ |
| B | b | A | | EE/3 | $ G(m3)$ |
| B | b | G | | GG | $ G(m3)$ |
| B | b | Z | | $EE/(12*CW*SW)$ | $ 3*G(m3)*(1-G5)-4*SW^2*G(m3)$ |
| W+ | $\sim H-$ | $\sim H3$ | | $-EE/(2*SW)$ | $ m1.p2-m1.p3$ |
| W+ | $\sim H-$ | $\sim X$ | | $i*EE/(2*SW)$ | $ m1.p2-m1.p3$ |
| W- | $\sim H+$ | $\sim H3$ | | $-EE/(2*SW)$ | $ m1.p3-m1.p2$ |
| W- | $\sim H+$ | $\sim X$ | | $-i*EE/(2*SW)$ | $ m1.p3-m1.p2$ |
| Z | $\sim H+$ | $\sim H-$ | | $EE/(2*CW*SW)$ | $ B00003*m1.p3+B00005*m1.p2$ |
| Z | $\sim H3$ | $\sim X$ | | $i*EE/(2*CW*SW)$ | $ m1.p2-m1.p3$ |
| $\sim H3$ | $\sim H3$ | $\sim X$ | $\sim X$ | $-2*\ln a2$ | $ 1$ |
| $\sim X$ | $\sim X$ | $\sim X$ | $\sim X$ | $-6*\ln a2$ | $ 1$ |

LanHEP package is able to generate model files with Feynman rules based on Lagrangian.

Semenov, A.

LanHEP - a package for automatic generation of Feynman rules from the Lagrangian.

arXiv: 1005.1909 [hep-ph]

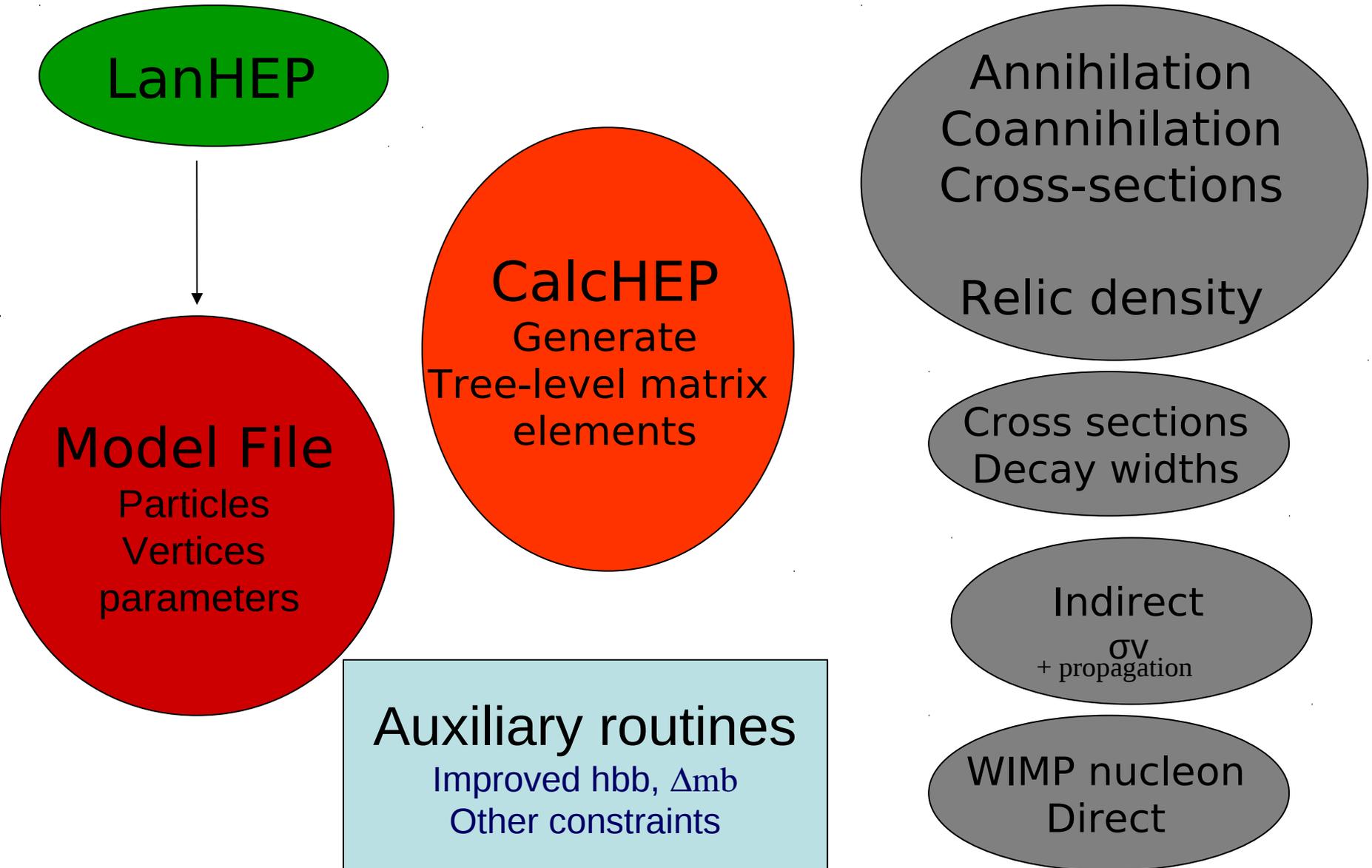
Matrix elements codes are **not** generated in advance.
When some matrix element needs, it is

- 1) generated by means of CalcHEP,
- 2) compiled as a shared library,
- 3) linked to the main code,
- 4) included in the list of attached matrix elements,
- 5) stored on the disk.

Last two options need for subsequent call of the same matrix element. Such scheme was initially realized for co-annihilation processes, but now is used for all matrix elements in micrOMEGAs.

LanHEP package is able to generate model files with Feynman rules based on Lagrangin.

micrOMEGAs



LanHEP

Model File

Particles
Vertices
parameters

CalcHEP

Generate
Tree-level matrix
elements

Auxiliary routines

Improved hbb, Δ_{mb}
Other constraints

Annihilation
Coannihilation
Cross-sections

Relic density

Cross sections
Decay widths

Indirect
 σv
+ propagation

WIMP nucleon
Direct

MicrOMEGAs_3 :new features

- Neutrino signal of DM captured by Sun/Earth
- Generalization of relic density computation
 - 3-body final state with virtual W/Z
 - Z₃ discrete symmetry
 - Asymmetric dark matter
 - g_{eff}(T), h_{eff}(T) input (SM thermodynamics)

Particle decays

- Virtual W/Z channels (automatically)
- Loop-induced Higgs decay (semi-automatically)

Loop induces 2γ and γZ signals of DM annihilation

- for SUSY-like models (semi-automatically)

SLHA interface with BSM related packages

DM capture in Sun

- DM particles captured by Sun/Earth, concentrate in center and annihilate into SM
- Lead to neutrino flux, can be observed at Earth (SuperKamiokande, IceCube)
- Shape of neutrino flux depends on dominant DM annihilation channel
- Signal determined by cross section for DM scattering on nuclei --related to DD
- Capture rate

$$C_\chi = \frac{\rho_\chi}{m_\chi} \int_0^\infty du f_1(u) \int 4\pi r^2 dr \sum_A \sigma_{\chi A N_A}(r) \frac{\beta_A}{\alpha_A} \left(e^{-\alpha_A u^2} - e^{-\alpha_A (u^2 + v_{esc}^2(r))/\beta_A} \right)$$

$$\alpha_A = \frac{1}{3} m_\chi m_A R_A^2, \quad \beta_A = \frac{(m_\chi + m_A)^2}{4m_\chi m_A}$$

- Annihilation

$$A_x = \langle \sigma v \rangle \frac{V_2}{V_1^2}$$

- Number of DM particles

$$\begin{aligned} \dot{N}_x &= C_x - A_{xx} N_x^2 - A_{x\bar{x}} N_x N_{\bar{x}} - EN_x, \\ \dot{N}_{\bar{x}} &= C_{\bar{x}} - A_{x\bar{x}} N_x N_{\bar{x}} - A_{\bar{x}\bar{x}} N_{\bar{x}}^2 - EN_{\bar{x}}, \end{aligned}$$

- When capture/annihilation is large, no evaporation, self-conj. particle, equilibrium is reached and annih. rate determined by capture rate

$$\Gamma_x = A_{xx} N_x^2 = C_x/2.$$

- In general solve equation for number density numerically and obtain ν flux at Earth

$$\frac{d\phi_\nu}{dE_\nu} = \frac{1}{4\pi d^2} \left(\Gamma_{xx} Br_{\nu\nu} \frac{dN_{\nu\nu}}{dE} + \Gamma_{x\bar{x}} \sum_f Br_{f\bar{f}} \frac{dN_f}{dE} \right)$$

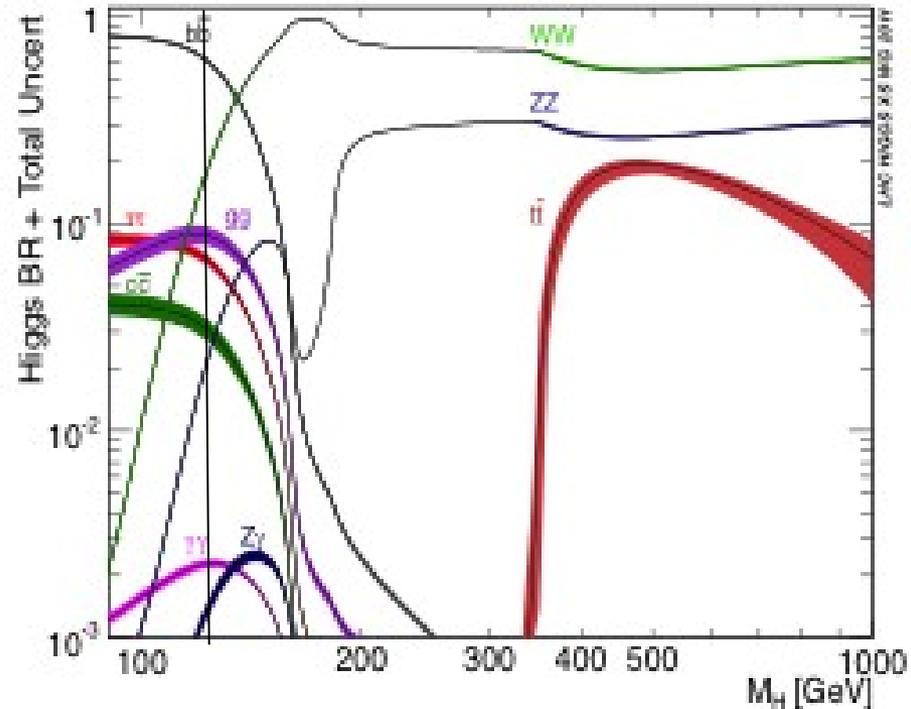
- neutrino spectrum originating from different SM annihilation channels and including oscillation available in
 - M. Cirelli, hep-ph/0506298

Neutrino-muon conversions in rocks is calculated according to Erkoca, Reno, Sarcevic 2009

- Upward and contained and muon fluxes are computed

Higgs decays

- Higgs at 126 GeV at LHC
- Decay modes :
bb, WW*, ZZ*, gg, $\gamma\gamma$
- Include
 - 3 body decays
 - loop-induced decays
- Experimental results give signal strength



$$R_{gg}^{h_i}(X) \equiv \frac{\Gamma(h_i \rightarrow gg) \text{BR}(h_i \rightarrow X)}{\Gamma(h_{\text{SM}} \rightarrow gg) \text{BR}(h_{\text{SM}} \rightarrow X)},$$

$$R_{\text{VBF}}^{h_i}(X) \equiv \frac{\Gamma(h_i \rightarrow WW) \text{BR}(h_i \rightarrow X)}{\Gamma(h_{\text{SM}} \rightarrow WW) \text{BR}(h_{\text{SM}} \rightarrow X)}$$

Loop induced Higgs decays

- Introduce effective operators

- $\lambda h F_{\mu\nu} F^{\mu\nu}$ (CP-even) $\lambda' h F_{\mu\nu} F^{\mu\nu}$ (CP-odd)

$$\mathcal{L} = g_{h\psi\psi} \bar{\psi}\psi h + i g'_{h\psi\psi} \bar{\psi}\gamma_5\psi h + g_{h\phi\phi} M_\phi h \phi\phi + g_{hVV} M_V h V_\mu V^\mu$$

$$\lambda = \frac{\alpha}{8\pi} \left[g_{h\nu\nu} f_\psi^c q_\phi^2 \frac{1}{M_\phi} A_{1/2}\left(\frac{M_h^2}{4M_\phi^2}\right) - g_{h\nu\nu} f_V^c q_V^2 \frac{1}{2M_V} A_1\left(\frac{M_h^2}{4M_V^2}\right) + g_{h\phi\phi} f_\phi^c q_\phi^2 \frac{1}{2M_\phi} A_0\left(\frac{M_h^2}{4M_\phi^2}\right) \right]$$

- Add QCD corrections

- e.g. for gluons

$$\lambda = -R \left[A_{htt}^{LO} C_t + (A_{hbb}^{LO} + A_{hcc}^{LO}) C_q + \sum_{\tilde{q}} A_{h\tilde{q}\tilde{q}}^{LO} C_{\tilde{q}} \right]$$

micrOMEGAS packages contains Passarino-Veltman functions needed for $H \rightarrow gg$, $H \rightarrow \gamma\gamma$ and tabulated functions for QCD corrections for $H \rightarrow \gamma\gamma$ amplitudes.

At the level of LanHEP package it was implemented an option to extract couplings of Higgs- $\rightarrow X, X$ interaction and convert them to Hgg , $H\gamma\gamma$ vertexes off effective Lagrangian.

A good agreement with HDECAY for MSSM Higgses was obtained, although some special SUSY effects can not be included in this way.

Files with BSM/SM coupling rates is **HIGGBOUNDS** format is generated.

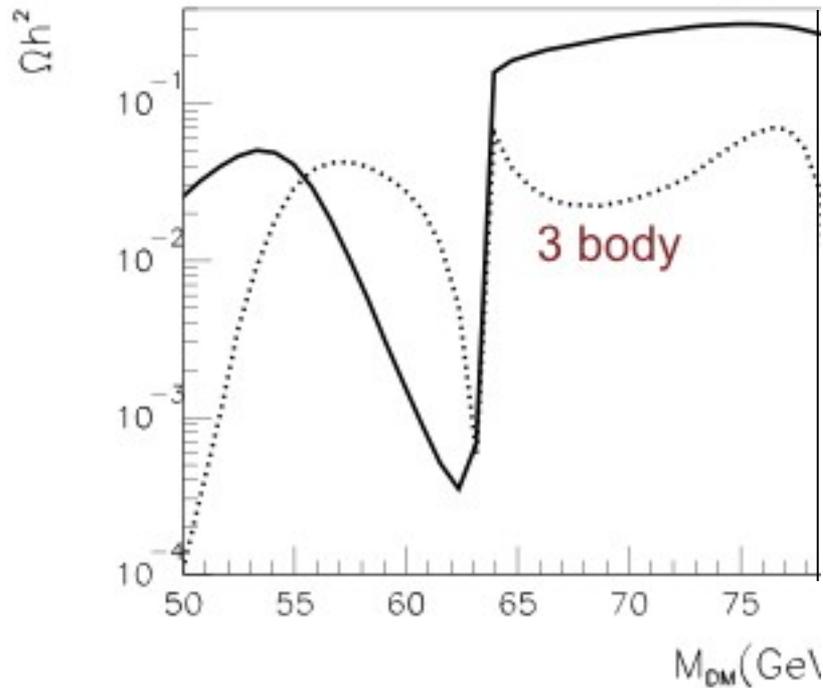
Generalization of relic density calculation

Three-body processes

- Annihilation into 3-body final state can be as large. Example is IDM where annihilation $DM, DM \rightarrow$ fermions is suppressed and processes with virtual W are important before WW threshold (Yaguna, arXiv: 1003.2730)
- About 5% correction in MSSM (close to threshold)
- **Problems:**
 - I) As we know from Higgs decay calculation that 4-body kinematics needs. But it is very slow.
 - a) calculate $Dm, Dm \rightarrow W^+, e, \nu$
 - b) assume Breit-Wigner structure of 4-body reaction and calculate for it K -factor 3-body/4-body
 - II) $Dm, Dm \rightarrow W^+, e, \nu$ contains not resonance diagrams, but resonance ones are not gauge invariant.

Select diagrams with at least one virtual gauge boson

Example : MSSM



- bino/higgsino LSP, $\mu=150\text{GeV}$, $M_2=2M_1$

Semi-annihilation

- Discrete remnant of some broken gauge group, in general does not have to be Z_2 - consider Z_N
- Impact for dark matter :
 - New processes
 - **semi-annihilation** : processes involving different number of “odd particles” $xx \rightarrow x^* \text{ SM}$
 - T. Hambye, 0811.0172, T. Hambye, M. Tytgat, 0907.1007
 - More than one DM candidate
 - **Assisted freeze-out**/DM conversion : interaction between particles from different dark sectors
 - » $X_1 X_1 \leftrightarrow X_2 X_2$

The Z_3 case

- Number density (x : dark sector X: SM)

$$\frac{dn}{dt} = -v\sigma^{xx^* \rightarrow XX} (n^2 - \bar{n}^2) - \frac{1}{2}v\sigma^{xx \rightarrow x^* X} (n^2 - n\bar{n}) - 3Hn.$$

$$\sigma_v \equiv v\sigma^{xx^* \rightarrow XX} + \frac{1}{2}v\sigma^{xx \rightarrow x^* X} \quad \text{and} \quad \alpha = \frac{1}{2} \frac{\sigma^{xx \rightarrow x^* X}}{\sigma_v}$$

$$3H \frac{dY}{ds} = \sigma_v (Y^2 - \alpha Y\bar{Y} - (1 - \alpha)\bar{Y}^2).$$

- Modified equation solved numerically. At low temperatures standard freeze-out picture is expected ($Y=Y_{\text{eq}}+\Delta Y$)

$$3H \frac{d\bar{Y}}{ds} = \sigma_v \bar{Y} \Delta Y (2 - \alpha)$$

Asymmetric DM

- The case where DM is not self-conjugate (e.g. Dirac fermion, complex scalar)
- $Y^+(Y^-)$: abundance of DM particle(anti-)

$$\frac{dY^\pm}{ds} = \frac{2 \langle \sigma v \rangle}{3H} (Y^+ Y^- - Y_{eq}^+ Y_{eq}^-)$$

- $\Delta Y = Y^+ - Y^-$ is constant

- Define $Y = 2(Y^+ Y^-)^{1/2}$

$$\frac{dY}{ds} \equiv \frac{\langle \sigma v \rangle}{3H} (Y^2 - Y_{eq}^2) \sqrt{1 + \left(\frac{\Delta Y}{Y}\right)^2}$$

- Equation is solved numerically. No freeze-out for large ΔY

- Relic density

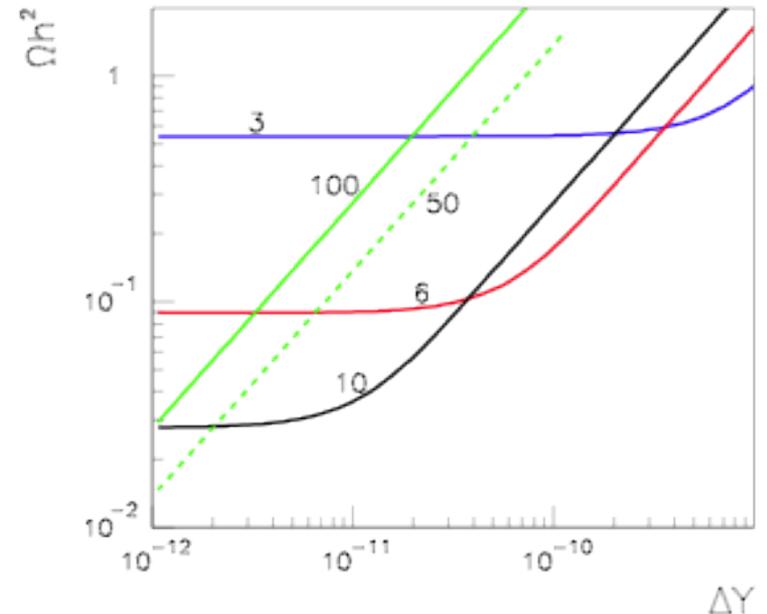
$$\Omega h^2 = \frac{8\pi}{3H_{100}^2} \frac{m_\chi}{M_{\text{Planck}}} \frac{\sqrt{Y_0^2 + \Delta Y^2}}{80}$$

- For each specie

$$\Omega_{\pm} h^2 = \frac{\Omega h^2}{1 + e^{\mp \delta_{DM}}}$$

- DM asymmetry always increase relic abundance
- Example : neutrino DM
- ΔY global parameter - taken into account for DD (always compute DM-nucleon and antiDM-nucleon) and indirect detection

$$Q = \frac{1}{2} \langle \sigma v \rangle \frac{\rho_\chi \rho_{\bar{\chi}}}{m_\chi^2} \frac{dN_a}{dE}$$



Monochromatic gamma-rays

- Monochromatic gamma rays ($\gamma\gamma, \gamma Z$) are loop-induced (suppressed) BUT lead to very distinctive signal
 - C. Weniger, 1204.2797, T. Bringmann et al, 1203.1312 (signal in Fermi-LAT?)

In generic models only Higgs contribution are known.



NEW

- Available in micrOMEGAs for SUSY models (MSSM, NMSSM, CPVMSSM) and Z3M
 - Computed with SloopS, a code for computation of one-loop processes in the SM, MSSM and some extensions
 - F. Boudjema, A. Semenov, D. Temes, hep-ph/0507127
 - G. Chalons, A. Semenov, arXiv:1110.2064

Can be extended for any model.

Conclusion

- To understand the nature of dark matter clearly need information and cross checks from cosmology, direct and indirect detection as well as from collider physics
- micrOMEGAs is tool to perform these analyses in a generic model