## Phenomenology of a $U(2)^3$ flavour symmetry

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mainly based on: Barbieri,Buttazzo,S,Straub arXiv:1203.4218 and 1206.1327

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## Motivations and Conclusions

Flavour:excellent agreement between data and CKM pictureIn other words: $\Delta \mathcal{L} = \sum_{i} \frac{1}{\Lambda_{i}^{2}} \mathcal{O}_{i} \Rightarrow \Lambda_{i} \gtrsim 10^{3} \div 10^{4} \,\mathrm{TeV}$ Hierarchy problem: $m_{h} \approx \Lambda$ [A = highest scale h couples to]

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#### $U(2)^3 = \overline{U(2)_{Q_L} \times U(2)_{U_R} \times U(2)_{D_R}}$

 $\checkmark \xi \sim V_{CKM}^{2\div 4} \Rightarrow \Lambda \sim a \text{ few TeV}$  is OK with flavour bounds

✓ potentially rich phenomenology behind the corner

✓ other virtues...

Exact 
$$U(2)^3 \longrightarrow m_u = m_d = m_s = m_c = 0, V_{CKM} = 1$$

$$Y_u = y_t \left( \begin{array}{c|c} 0 & 0 \\ \hline 0 & 1 \end{array} \right) \qquad Y_d = y_b \left( \begin{array}{c|c} 0 & 0 \\ \hline 0 & 1 \end{array} \right)$$

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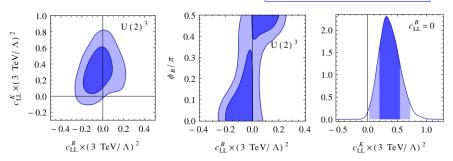
Assume: all flavour violation controlled by  $\Delta Y_{u,d}$ , V

i.e.  $\Delta \mathcal{L}$  built with bilinears like  $\mathbf{\bar{q}}_{L} \mathbf{V} \gamma_{\mu} q_{3L}$ ,  $\mathbf{\bar{q}}_{L} \Delta Y_{d} \mathbf{d}_{R}$ Example:  $\Delta \mathcal{L} \supset \frac{c_{L}^{B} e^{i\phi_{B}}}{\Lambda^{2}} (\mathbf{V}_{tb} \mathbf{V}_{ti}^{*})^{2} (\bar{d}_{L}^{i} \gamma_{\mu} b_{L})^{2}$ , i = d, s  $\begin{bmatrix} B_{d,s}^{0} - \bar{B}_{d,s}^{0} \end{bmatrix}$ 

### Bounds and New Effects

 $\Delta F = 2$ 

Can solve CKM fit tensions!



#### Messages

- Data consistent with  $\Delta \mathcal{L} = \sum_i \xi_i \frac{c_i}{\Lambda_i^2} \mathcal{O}_i$  and  $|c_i| = 0.2 \div 1$
- Larger effects than  $U(3)^3$  allowed

see also [Buras, Girrbach 2012]

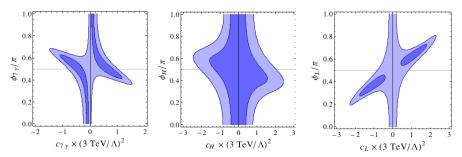
 $U(3)^3$ :  $Y_u, Y_d$  instead of  $\Delta Y_{u,d}, V$ 

 $c_I^B = c_I^K$  and  $\phi_B = 0$ 

### Bounds and New Effects

 $\Delta F = 1$ 

 $b o s(d) \ell ar{\ell}, \; b o s(d) 
u ar{
u}$ 



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#### How to know it is $U(2)^3$ ? (if some new physics signal seen)

- s d correlation in B decays (same as SM), with new phase
- no new phase in  $K \bar{K}$  mixings
- lots of specific correlations! see e.g. [Barbieri et al 2011, Straub 2013]

Observables to watch:  $S_{\phi K}$ ,  $S_{\psi \phi}$ ,  $|V_{ub}|$ ,  $b \to s(d) \ell \bar{\ell}, \nu \bar{\nu}$ , ...

#### $U(2)^3$ ideal for natural theories!

- SUSY with light stops (would also suppress EDMs)
- easy to satisfy collider and precision constraints
   e.g. in Composite Higgs models, see [Barbieri et al 2012]

# Back up

# Up sector within $U(2)^3$

Minimal  $U(2)^3$ : prediction of no detectable effects in

- Top FCNC [BR( $t \rightarrow c\gamma, cZ$ )]: below future LHC sensitivity
- CPV in  $D \overline{D}$  mixing  $[\phi_{12}]$ : below future LHCb sensitivity
- Direct CPV in D decay  $[A_{CP}^D(\pi\pi, KK)]$ : below per mille level

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Generic  $U(2)^3$  [add  $V_u \sim (1,2,1), V_d \sim (1,1,2)$ ]

- could explain  $\Delta A_{CP}^{exp}$
- respecting all current flavour and EDMs bounds

• keeping the same null predictions for  $\mathsf{BR}(t o c\gamma, cZ)$  and  $\phi_{12}$ 

## Generic $U(2)^3$ : bounds and new effects

$$\Delta Y_{u,d} = L(s_L^{u,d}) \cdot \Delta Y_{u,d}^{\text{diag}} \cdot R(s_R^{u,d}), \quad V = (0, \epsilon_L), \quad V_{u,d} = (0, \epsilon_R^{u,d})$$

 $s_L^{u,d}$ ,  $\epsilon_L$  fixed from tree level observables,  $s_R^{u,d}$ ,  $\epsilon_R^{u,d}$  bounded from above

