

Noncommutative Geometry in the LHC-Era

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Overview

- 1 NCG: Basic Ideas
- 2 The Standard Model
- 3 Beyond the Standard Model
- 4 Conclusions

Overview

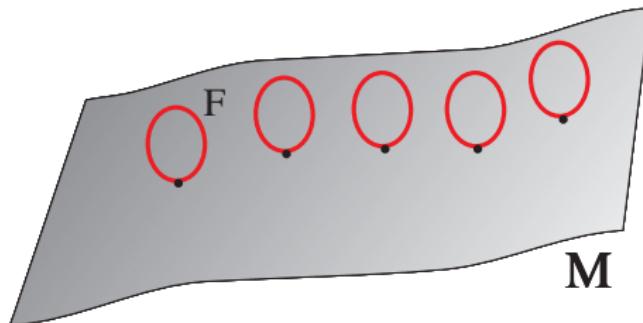
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Analogy: Noncomm. geometry \leftrightarrow Kaluza-Klein space

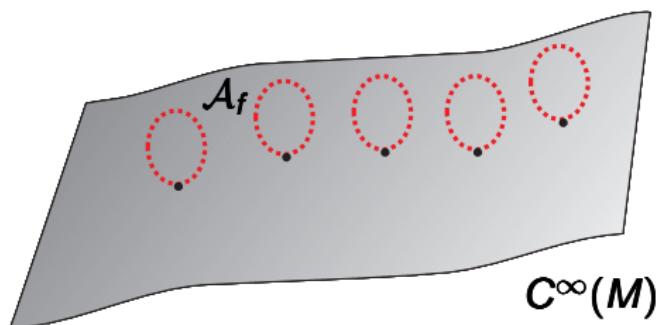


Idea:

$M \rightarrow C^\infty(M)$, $F \rightarrow$ some "finite or inner space",

differential geometry \rightarrow spectral triple

Noncommutative Geometry (A. Connes '90s)

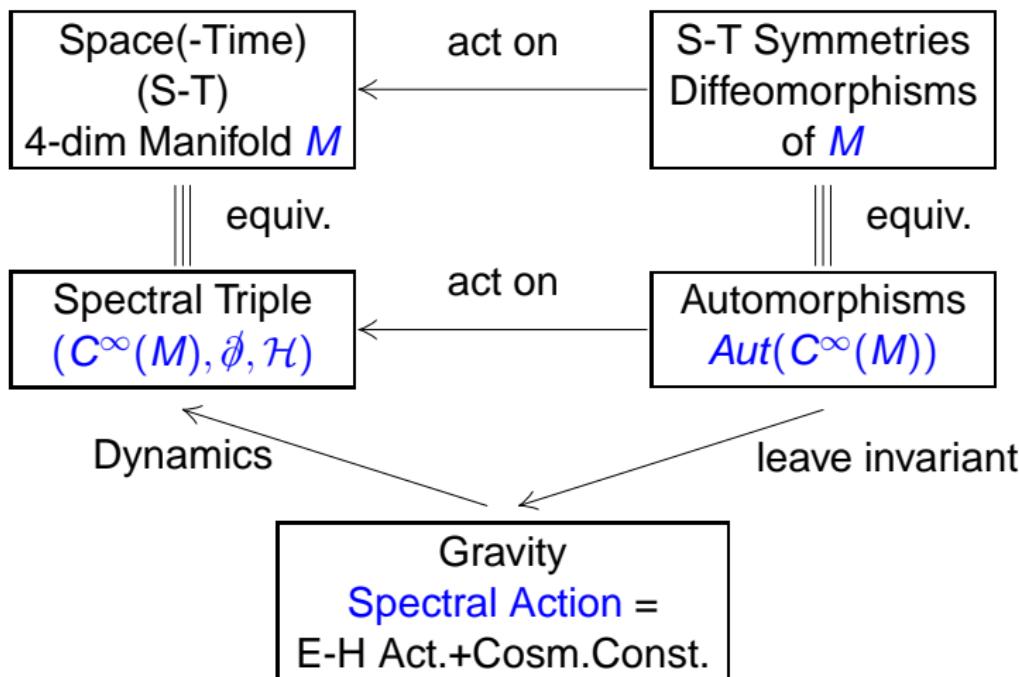


Replacing manifolds by algebras

extra dimension: $F \rightarrow \mathcal{A}_f = M_1(\mathbb{K}) \oplus M_2(\mathbb{K}) \oplus \dots$

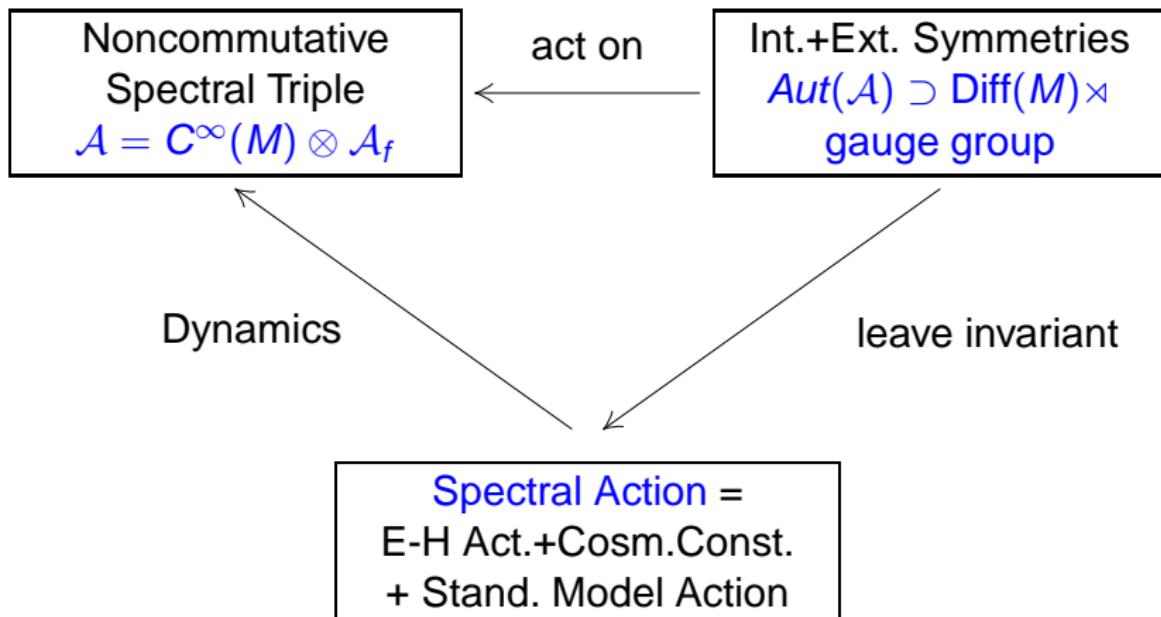
Kaluza-Klein space: $M \times F \rightarrow \mathcal{A} = C^\infty(M) \otimes \mathcal{A}_f$

Euclidean space(-time)!



Noncommutative Standard Model

(A.Chamseddine, A.Connes '96):

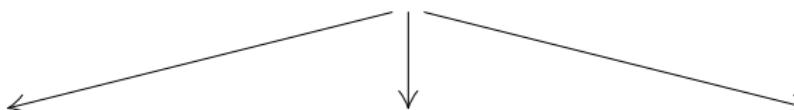


Noncommutative Geometry in a nutshell:

- finite dimensions treated the same as “ordinary” space
- Dirac operator plays the rôle of the metric
- gauge fields and scalar fields are interpreted as “internal” gravitational fields

The Dirac operator plays a multiple role:

$$D = \not{\partial} \otimes 1_f + \not{A} + \gamma_5 \otimes \Phi$$



Scalar & Gauge
Bosons

Particle Dynamics,
Ferm. Mass Matrix

Metric of M ,
Internal Metric

Generalising to noncommutative internal space:

$$C^\infty(\mathcal{M}) \rightarrow C^\infty(\mathcal{M}) \otimes \mathcal{A}_f$$

generalise to Chamseddine-Connes-Dirac operator:

$$\not{D} \rightarrow \not{D} = \not{\partial} \otimes 1_f + \not{A} + \gamma_5 \otimes \Phi$$

The geometric setup imposes constraints:

- mathematical axioms of Noncomm. Geometry
→ Restrictions on particle content & poss. interactions
- symmetries of finite space
→ determines gauge group
- representation of matrix algebra
→ representation of non-abelian gauge sub-group
- Dirac operator → allowed mass terms / scalar fields

General requirements for Particle Models:

- requirement for internal space
→ Standard Model is a sub-model
- no harmful (Yang-Mills) anomalies
→ representation of abelian gauge sub-group
- Action from **Spectral Action**
→ high energy effective action + **constraints**
- low energy physics by renormalisation group flow

Experimental requirements at m_Z :

- $g_1(m_Z) = 0.3575$, $g_2(m_Z) = 0.6514$, $g_3(m_Z) = 1.221$
- $m_{W^\pm} = 80.39 \text{ GeV}$
- $m_{top} = 172.9 \pm 1.5 \text{ GeV}$
- $m_{SMS} = 125.5 \pm 1.1 \text{ GeV}$

The Spectral Action (A. Connes, A. Chamseddine 1996)

$$(\Psi, \mathcal{D}\Psi) + S_{\mathcal{D}}(\Lambda) \quad \text{with } \Psi \in \mathcal{H}$$

- $(\Psi, \mathcal{D}\Psi)$ = fermionic action
 - includes Yukawa couplings
 - & fermion–gauge boson interactions
 - + constraints at Λ
- $S_{\mathcal{D}}(\Lambda)$ = # eigenvalues of \mathcal{D} up to cut-off $\pm\Lambda$
 - \equiv Einstein-Hilbert action + Cosm. Const.
 - + bosonic action
 - + constraints at Λ
- constraints \Rightarrow fewer free parameters

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Constraints on the SM parameters at the cut-off Λ :

$$\frac{5}{3} g_1(\Lambda)^2 = g_2(\Lambda)^2 = g_3(\Lambda)^2 = \frac{Y_2(\Lambda)^2}{H(\Lambda)} \frac{\lambda(\Lambda)}{24} = \frac{1}{4} Y_2(\Lambda)$$

- g_1, g_2, g_3 : $U(1)_Y, SU(2)_w, SU(3)_c$ gauge couplings
- λ : quartic SMS coupling
- Y_2 : trace of the Yukawa matrix squared
- H : trace of the Yukawa matrix to the fourth power

Consequences from the SM constraints:

Input:

- Big Desert
- $g_1(m_Z) = 0.3575$, $g_2(m_Z) = 0.6514$, $g_3(m_Z) = 1.221$
- renormalisation group equations
- ($m_{top} = 171.2 \pm 2.1$ GeV)

Output:

- $g_2^2(\Lambda) = g_3^2(\Lambda)$ at $\Lambda = 1.1 \times 10^{17}$ GeV
- $m_{top} < 190$ GeV (follows from constraints alone)
- no 4th SM generation

Excluded by Tevatron & LHC since:

- $m_{SMS} \neq 168.3 \pm 2.5$ GeV
- $\frac{5}{3} g_1(\Lambda)^2 \neq g_2(\Lambda)^2$

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Excluded models:

Extensions of the Standard Model:

- AC -model (C.S. '05):
new particles (A and C) with opposite hypercharge
dark matter as bound AC -states (Fargion,Khlopov,C.S. '05)
- θ -model (C.S. '07): new particles with $SU_c(D)$ -colour
- Vector-Doublet Model (Squellari, C.S. '07):
new $SU_w(2)$ -vector doublets

Problem for models: • $m_{SMS} \geq 170$ GeV

- constraints on g_1 , g_2 , g_3 at Λ .

SM + $U(1)_X$ scalar field + new fermions (C.S. '10, C.S. t.a.):

- SM as a sub-model: comme il faut!
- gauge group: $U(1)_Y \times SU(2)_W \times SU(3)_C \times U(1)_X$
- new fermions: $X_\ell^1, X_r^2 : (0, 1, 1, +1)$, $X_r^1, X_\ell^2 : (0, 1, 1, 0)$
 $V_\ell^c, V_r^c : (-1/6, 1, \bar{3}, +1/2)$
 $V_\ell^w, V_r^w : (0, \bar{2}, 1, +1/2)$
in each SM-generation
- new scalar: $\varphi : (0, 1, 1, +1)$
- $\mathcal{L}_{scalar} = -\mu_1^2 |H|^2 - \mu_2^2 |\varphi|^2 + \frac{\lambda_1}{6} |H|^4 + \frac{\lambda_2}{6} |\varphi|^4 + \frac{\lambda_3}{3} |H|^2 |\varphi|^2$
- $\mathcal{L}_{ferm} = g_{\nu, X^1} \bar{\nu}_r \varphi X_\ell^1 + g_{X^1} \bar{X}_\ell^1 \varphi X_r^1 + g_{X^2} \bar{X}_\ell^2 \varphi X_r^2$
 $+ \bar{V}_\ell^c m_c V_r^c + \bar{V}_\ell^w m_w V_r^w + h.c.$
- $\mathcal{L}_{gauge} = 1/g_4^2 F_X^{\mu\nu} F_{X,\mu\nu}$
- $U(1)_Y \times SU(2)_W \times SU(3)_C \times U(1)_X \rightarrow U(1)_{el.} \times SU(3)_C$

The constraints at Λ :

- $g_2(\Lambda) = g_3(\Lambda) = \sqrt{\frac{7}{6}} g_1(\Lambda) = \sqrt{\frac{3}{2}} g_4(\Lambda)$
- $\lambda_1 = 36 \frac{H}{Y_2^2} g_2^2, \quad \lambda_2 = 36 \frac{H_X + \text{tr}(g_{\nu,X^1}^2 g_{X^1}^2)}{Y_X^2} g_2^2$
- $\lambda_3 = 36 \frac{\text{tr}(g_\nu^2)}{Y_2} \frac{\text{tr}(g_{\nu,X^1}^2)}{Y_X} g_2^2, \quad Y_2 = Y_X = 6 g_2$
- $Y_X := \text{tr}(g_{\nu,X^1}^2) + \text{tr}(g_{X^1}^2) + \text{tr}(g_{X^2}^2)$
- $H_X := \text{tr}(g_{\nu,X^1}^4) + \text{tr}(g_{X^1}^4) + \text{tr}(g_{X^2}^4)$

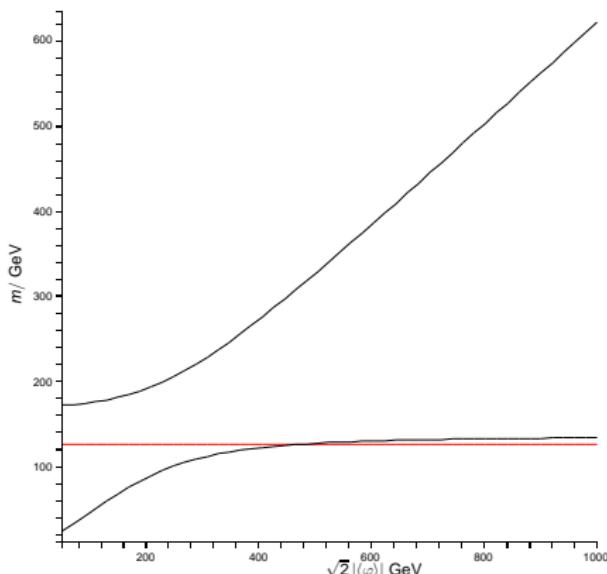
Example: one point in parameter space at Λ

- $Y_2 \approx 3g_{top} + g_{\nu_\tau}$
- $Y_X \approx g_{\nu,X}^2, \quad H_X \approx g_{\nu,X}^4, \quad \text{tr}(g_{\nu,X^1}^2 g_{X^1}^2) \ll 1$
- $(m_w)_{ij} \approx \Lambda, \quad (m_c)_{ij} \approx 10^{14} \text{ GeV}$

Results for 1-loop renormalisation groups:

(threshold effects have been neglected)

- Constraints
 $\Rightarrow \Lambda \approx 2 \times 10^{18} \text{ GeV}$
- $m_{top} \approx 172.9 \pm 1.5 \text{ GeV}$
- $m_{\varphi_1, \text{SMS}} \approx 125 \pm 1.1 \text{ GeV}$
- $m_{\varphi_2} \approx 379 \pm 93 \text{ GeV}$
- $m_{Z_X} \approx 105 \pm 25 \text{ GeV}$
- $g_4(m_Z) \approx 0.35$
- $m_{X^1, X^2} \lesssim 50 \text{ GeV}$
- free parameter: $|\langle \varphi \rangle|$
- φ_2 -sector could be much heavier



Mass EVs of scalar fields

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Questions & to-do-list

- Is the SM + scalar model compatible with LHC-data?
- Does the SM + scalar model contain viable dark matter candidates?
- Explore parameter space (g_{ν, X^1} , $g_{X^1}^2$, $g_{X^2}^2$, m_w , m_c)
- Extend renormalisation group analysis to n -loop, $n \geq 2$
- Classify models / geometries beyond the Standard Model
- Renorm. group flow for all couplings in the spectral action
(Exact Renorm Groups, M. Reuter et al. ?)
- Spectral triples with Lorentzian signature
(A. Rennie, M. Paschke, R. Verch,...)

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