



Photon-photon scattering in collisions of laser pulses

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• Vacuum polarisation introduction Heisenberg-Euler Lagrangian



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- Vacuum birefringence and dichroism in interacting lasers beams



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- Vacuum birefringence and dichroism in interacting lasers beams
- Elastic photon-photon scattering



- Vacuum birefringence and dichroism in interacting lasers beams
- Elastic photon-photon scattering
- Vacuum four-wave mixing



- Vacuum birefringence and dichroism in interacting lasers beams
- Elastic photon-photon scattering
- Vacuum four-wave mixing
- Summary





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D. Bernard (LPNHE) et al. EPJD **10**, 141-145 (2000)





$$\sigma_{\gamma\gamma}[\mathrm{cm}^2] = 7.3 \times 10^{-66} (\hbar\omega[\mathrm{eV}])^6$$
$$N_{\gamma} \approx 6 \times 10^{18} \mathcal{E}[\mathrm{J}] \hbar\omega[\mathrm{eV}]$$

D. Bernard (LPNHE) et al. EPJD **10**, 141-145 (2000)







$$\bigcirc = \bigwedge_{\sim} \bigvee_{\sim} \bigvee_{\sim} + \bigvee_{\sim} \bigvee_{\sim} \bigvee_{\sim} + \dots \quad a_0^2 = \frac{e^2}{m^2} \frac{k_{\mu} T^{\mu\nu} k_{\nu}}{(\varkappa k)^2} \quad a_0 \gg 1$$







"strong" field background k', ε' k, ε

perturbative "probe"





D. Bernard (LPNHE) et al. EPJD **10**, 141-145 (2000)





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"probe"





D. Bernard (LPNHE) et al. EPJD **10**, 141-145 (2000)





perturbative "probe"



$\begin{array}{l} \hbar\omega \ll mc^2 & (\lambda \ll \lambda_C) \\ \hat{A} \to A & & & & & & & & & \\ \end{array}$

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 $\alpha (E/E_{\rm cr})^2 \ll 1 \qquad E/E_{\rm cr} \lesssim 2 \times 10^{-4}$

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$$\mathcal{L} = \frac{1}{2} (E^2 - B^2) + c_1 (E^2 - B^2)^2 + c_2 (\mathbf{E} \cdot \mathbf{B})^2 \qquad (c_{1,2} \sim \frac{\alpha}{45\pi} \frac{1}{E_{\rm cr}^2})$$

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Maxwell H

$$\nabla^2 \mathbf{E} - \partial_t^2 \mathbf{E} = \partial_t \mathbf{J}_{\text{vac}} = \nabla \times \partial_t \mathbf{M}_{\text{vac}} + \partial_t^2 \mathbf{P}_{\text{vac}} - \nabla (\nabla \cdot \mathbf{P}_{\text{vac}})$$



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(c_{1,2} ~ \frac{\alpha}{45\pi} \frac{1}{E_{\rm cr}^2})

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Maxwell

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$$\mathbf{P}_{\text{vac}} = c_{1}'(E^{2} - B^{2})\mathbf{E} + c_{2}'(\mathbf{E} \cdot \mathbf{B})\mathbf{B}$$
$$\mathbf{M}_{\text{vac}} = -c_{1}'(E^{2} - B^{2})\mathbf{B} + c_{2}'(\mathbf{E} \cdot \mathbf{B})\mathbf{E}$$

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$$\mathbf{M}_{\text{vac}} = -c_{1}'(E^{2} - B^{2})\mathbf{B} + c_{2}'(\mathbf{E} \cdot \mathbf{B})\mathbf{E}$$
$$\mathbf{E}(\mathbf{x}, t) \approx \mathbf{E}^{(0)}(\mathbf{x}, t) + \int d^{3}x' \, \frac{\partial_{t}\mathbf{J}_{\text{vac}}(\mathbf{E}^{(0)}, \mathbf{x}', t_{\text{ret}})}{|\mathbf{x} - \mathbf{x}'|}$$

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(c_{1,2} ~ \frac{\alpha}{45\pi} \frac{1}{E_{\rm cr}^2})

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 $+ \mathbf{E}_{d}$

Maxwell

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 $\mathbf{P}_{\text{vac}} = c_1' (E^2 - B^2) \mathbf{E} + c_2' (\mathbf{E} \cdot \mathbf{B}) \mathbf{B}$



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$$I_s = 10^{24} \,\mathrm{Wcm}^{-2}$$
$$\omega_s = 1.5 \,\mathrm{eV}$$
$$\tau_s = 100 \,\mathrm{fs}$$
$$I_p = 10^{18} \,\mathrm{Wcm}^{-2}$$
$$\omega_p = 3 \,\mathrm{keV}$$
$$\tau_p = 100 \,\mathrm{fs}$$

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BK, A. Di Piazza and C. H. Keitel, PRA 82, 032114 (2010)

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R. Karplus and M. Neuman Phys. Rev. **80** 380-385 (1950)

Elastic photon-photon scattering



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R. Karplus and M. Neuman Phys. Rev. **80** 380-385 (1950)



Elastic photon-photon scattering



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Inelastic photon-photon scattering (four-wave mixing)




Photon-photon scattering

Elastic photon-photon scattering



R. Karplus and M. Neuman Phys. Rev. **80** 380-385 (1950)



Inelastic photon-photon scattering (four-wave mixing)





$$k_1^{\mu} + k_2^{\mu} + k_3^{\mu} + k_4^{\mu} = 0$$

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Strong laser:

 $150 \,\mathrm{PW}, 800 \,\mathrm{nm}, 30 \,\mathrm{fs}$ diffraction limited



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• Strong laser: $150 \, \mathrm{PW}, 800 \, \mathrm{nm}, 30 \, \mathrm{fs}$ diffraction limited

• Probe laser: 200 TW, 527 nm, 100 fsfocused to $300 \,\mu\text{m}$



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 Gaussian beams: 1st-order space, monochromatic



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- Gaussian beams: 1st-order space, monochromatic
- Parallel geometry



• Strong laser: 150 PW, 800 nm, 30 fs diffraction limited

• Probe laser: 200 TW, 527 nm, 100 fsfocused to $300 \,\mu\text{m}$

- Gaussian beams: 1st-order space, monochromatic
- Parallel geometry
- Required vacuum pressure at room temperature: $P \le 10^{-6}$ torr



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$$\frac{1}{2}\langle |\mathbf{E}_p + \mathbf{E}_d|^2 \rangle = I_p + I_{pd} + I_d$$

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Double-slit minima:

 $(m+1/2)\lambda_p = D\sin\theta$

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• "Detectable" vacuum signal: $I_d > 100(I_p + I_{pd})$



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- "Detectable" vacuum signal: $I_d > 100(I_p + I_{pd})$
- ~ 40 / 4 diffracted photons per shot



Photon-photon scattering scenario



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Photon-photon scattering scenario



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Photon-photon scattering scenario



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$$\mathbf{E}_{b} = \boldsymbol{\varepsilon}_{b} E_{b} f(x, y, z, t, \mathbf{w}_{b}, \omega_{b}) g(t - y)$$
$$\mathbf{E}_{a} = \boldsymbol{\varepsilon}_{a} E_{a} f(\tilde{x} - \tilde{x}_{0}, \tilde{y}, \tilde{z} - \tilde{z}_{0}, t - \Delta t, \mathbf{w}_{a}, \omega_{a}) g(t - \Delta t + \tilde{y})$$



Vulcan laser parameters:

- $\lambda_a = \lambda_b = 0.91 \,\mu\mathrm{m}$
- $\tau_a = \tau_b = 30 \, \mathrm{fs}$
- $P_a = P_b = 5 \,\mathrm{PW}$
- $w_a = 0.91 \,\mu\mathrm{m}$
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Vulcan laser parameters: $\lambda_a = \lambda_b = 0.91 \,\mu\mathrm{m}$ $\tau_a = \tau_b = 30 \, \mathrm{fs}$ $P_a = P_b = 5 \,\mathrm{PW}$ $w_a = 0.91 \,\mu{\rm m}$ $w_b = 100 \,\mu{\rm m}$ $\theta = 0.1$ $P_a = 2P_{\rm tot}/3$ $N_d(\lambda_b = 0.46\,\mu\mathrm{m}) \approx 4$ $N_d(\lambda_b = 0.23 \,\mu\mathrm{m}) \approx 16$







$$\omega = \omega_{a,1} + \omega_{a,2} + \omega_b$$
$$\omega \frac{y}{r} = \omega_{a,1} \cos \theta_{a,1} + \omega_{a,2} \cos \theta_{a,2} + \omega_b \cos \theta_b$$
$$\omega \frac{\rho}{r} = \omega_{a,1} \sin \theta_{a,1} + \omega_{a,2} \sin \theta_{a,2} + \omega_b \sin \theta_b$$

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 $\omega_{a,1}, \mathbf{k}_{a,1}$

 $\omega_{a,2}, \mathbf{k}_{a,2}$ —

 ω_b, \mathbf{k}_b —







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$$\omega \frac{\rho}{r} = \omega_{a,1} \sin \theta_{a,1} + \omega_{a,2} \sin \theta_{a,2} + \omega_b \sin \theta_b$$

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$$au_a = 1 \, \mathrm{fs}$$
 $au_b = 30 \, \mathrm{fs}$
 $N_d(\lambda_b = 0.46 \, \mu \mathrm{m}) \approx 1$
 $N_d(\lambda_b = 0.23 \, \mu \mathrm{m}) \approx 4$





 $\mathbf{E}_{a}(y+t) = \boldsymbol{\varepsilon}_{a} E_{a} \cos[\omega_{a}(y+t)] \operatorname{sech}[(y+t)/\tau_{a}]$ $\mathbf{E}_{b}(y-t) = \boldsymbol{\varepsilon}_{b} E_{b} \cos[\omega_{b}(y-t)] \operatorname{sech}[(y-t)/\tau_{b}]$



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BK and C. H. Keitel, NJP 14 103002 (2012)



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$$\mathbf{E}_{a}^{2}|\mathbf{E}_{b}| \sim \cos^{2}(\varphi_{a})\cos(\varphi_{b}) = \frac{1}{2}\left[\cos(\varphi_{b}) + \cos(2\varphi_{a})\cos(\varphi_{b})\right]$$

BK and C. H. Keitel, NJP 14 103002 (2012)

Summary






• Vacuum polarised by strong electromagnetic fields detectable with current and soon-to-be-available lasers

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- Single 10 PW beam split into two pulses sufficient for measuring elastic photon-photon scattering
- $$\begin{split} N_d(\lambda_p &= 0.23\,\mu\mathrm{m}) \approx 16\\ N_d(\lambda_p &= 0.46\,\mu\mathrm{m}) \approx 4 \end{split}$$

Summary



• Vacuum polarised by strong electromagnetic fields detectable with current and soon-to-be-available lasers

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- Single 10 PW beam split into two pulses sufficient for measuring elastic photon-photon scattering
- If strong pulse can be compressed to few-cycle, influence of inelastic photon-photon scattering also measurable

 $N_d(\lambda_p = 0.23 \,\mu\text{m}) \approx 16$ $N_d(\lambda_p = 0.46 \,\mu\text{m}) \approx 4$

 $N_d(\lambda_p = 0.23 \,\mu\text{m}) \approx 4$ $N_d(\lambda_p = 0.46 \,\mu\text{m}) \approx 1$ $\tau_a = 1 \,\text{fs} \ \tau_b = 30 \,\text{fs}$

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BK and C. H. Keitel, NJP **14** 103002 (2012) (see also NJP Highlights 2012)
BK, A. Di Piazza and C. H. Keitel, Nature Photon. **4**, 92 (2010)
BK, A. Di Piazza and C. H. Keitel, PRA **82**, 032114 (2010)



Current work

 $\nabla^2 \mathbf{E} - \partial_t^2 \mathbf{E} = \mathbf{J}_{\rm vac}(\mathbf{E})$







P. Boehl, BK and H. Ruhl (in preparation)

$$-\frac{\alpha}{(2\pi)^4} \int d^4x \operatorname{Tr} \langle x | \frac{1}{\gamma(p+eA)-m} e^{-ikx} \gamma^{\mu} \frac{1}{\gamma(p+eA)-m} e^{ik'x} \gamma^{\nu} | x \rangle$$

$$a_0 = \frac{eE\lambda_{\rm C}}{\hbar\omega},$$

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Envisaged scenario

