

Neutrino Oscillations (or Physics beyond the Standard Model of elementary particles)

Marcos Dracos

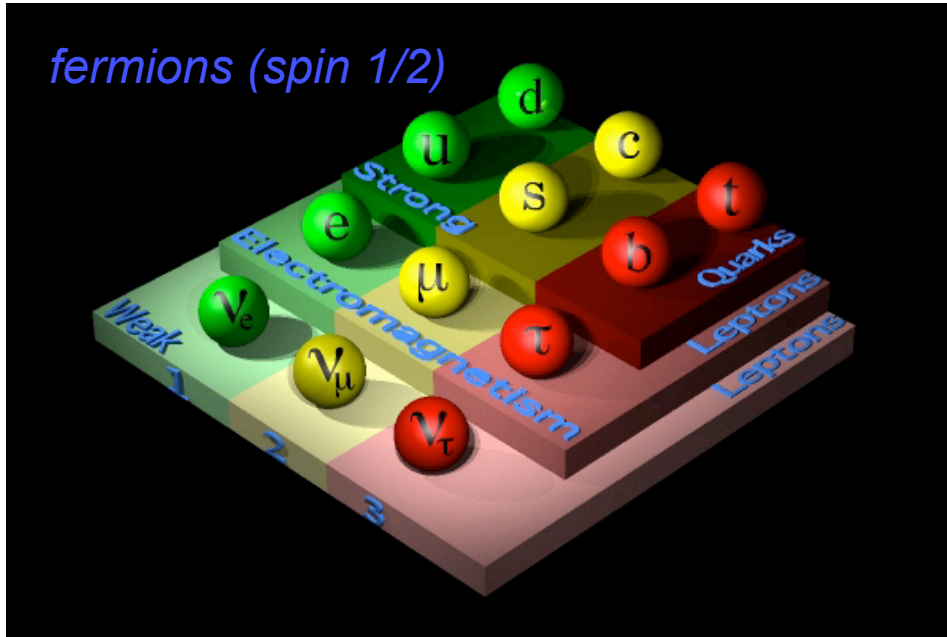
Institut Pluridisciplinaire Hubert Curien

Strasbourg

[\(marcos.dracos@in2p3.fr\)](mailto:marcos.dracos@in2p3.fr)

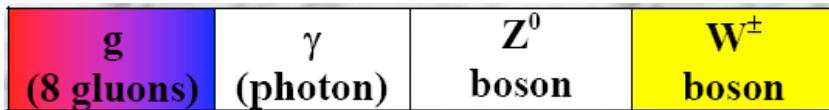


But, what a neutrino is?



- elementary particles,
- neutral (electrical charge=0)
- interacting only through weak interaction,
- they have massive and charged partners,
- massless up to recently (this is what was assumed by the Standard Model of elementary particles)?

+bosons carrying the interactions (spin integer)



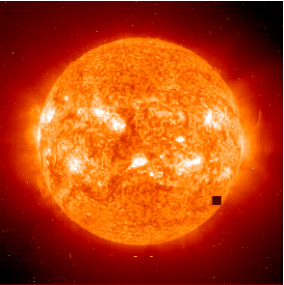
strong

electromagnetic

weak

+ the fermions anti-particles...

Where can we find neutrinos?



- Solar Neutrinos : $2 \cdot 10^{38}$ ν /s \rightarrow 40 billions ν /s/cm² on the earth \rightarrow 400000 billions ν /s/human (<20 MeV).
- Universe :
 - Big-Bang: 330 ν /cm³ (0.0004 eV \rightarrow 2000 km/s if $m_\nu = 10$ eV/c²).
 - Stars : 0.000006 ν /cm³.
 - Supernovae : 0.0002 ν /cm³.
- Earth radioactivity : 50 billions ν /s human.
- Nuclear Plants: 10-100 billions ν /s/human (1-10 MeV).
- Human body: 340 millions ν /day (20 mg of ⁴⁰K, β decay).

Neutrino Oscillations

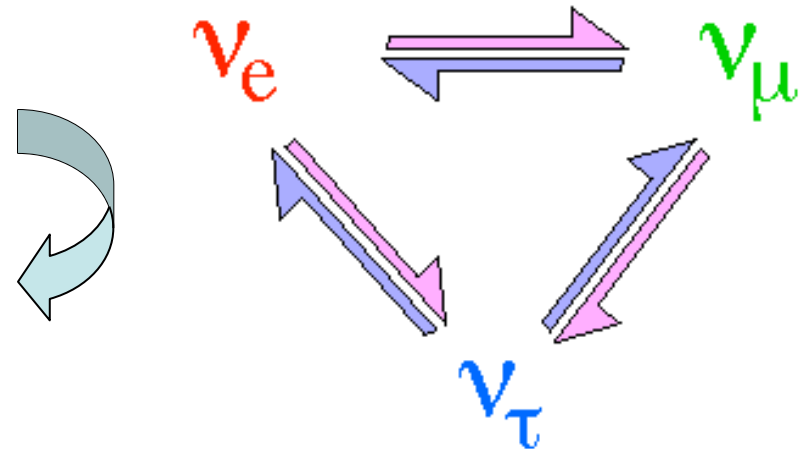
- The Standard Model doesn't predict any mass for the neutrinos.
- According to Quantum Mechanics, if neutrinos have a non-zero mass they can "oscillate" (change family during traveling).
- Why? Because their mass eigen states could not coincide with their flavour eigen states (or interaction eigen states).

States participating to the weak interaction
(flavour eigen states)

ν_e ν_μ ν_τ

States with masses well defined
(mass eigen states)

ν_1 ν_2 ν_3



Lepton Mixing and Quantum Mechanics

- The “known” neutrinos are combinations of mass eigen state neutrinos, for example, for electron neutrinos:

$$|\nu_e\rangle = U_{e1}|\nu_1\rangle + U_{e2}|\nu_2\rangle + U_{e3}|\nu_3\rangle$$

- For all neutrinos, we can write:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \underbrace{\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix}}_{\text{unitary mixing matrix}} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Maki-Nakagawa-Sakata matrix



- change of basis,*
- U is the transformation operator (matrix, unitarity),*
- the hypothetical states ν_1, ν_2, ν_3 have unique masses and are neutrino fundamental eigen states.*

Rotations between states

The most popular representation

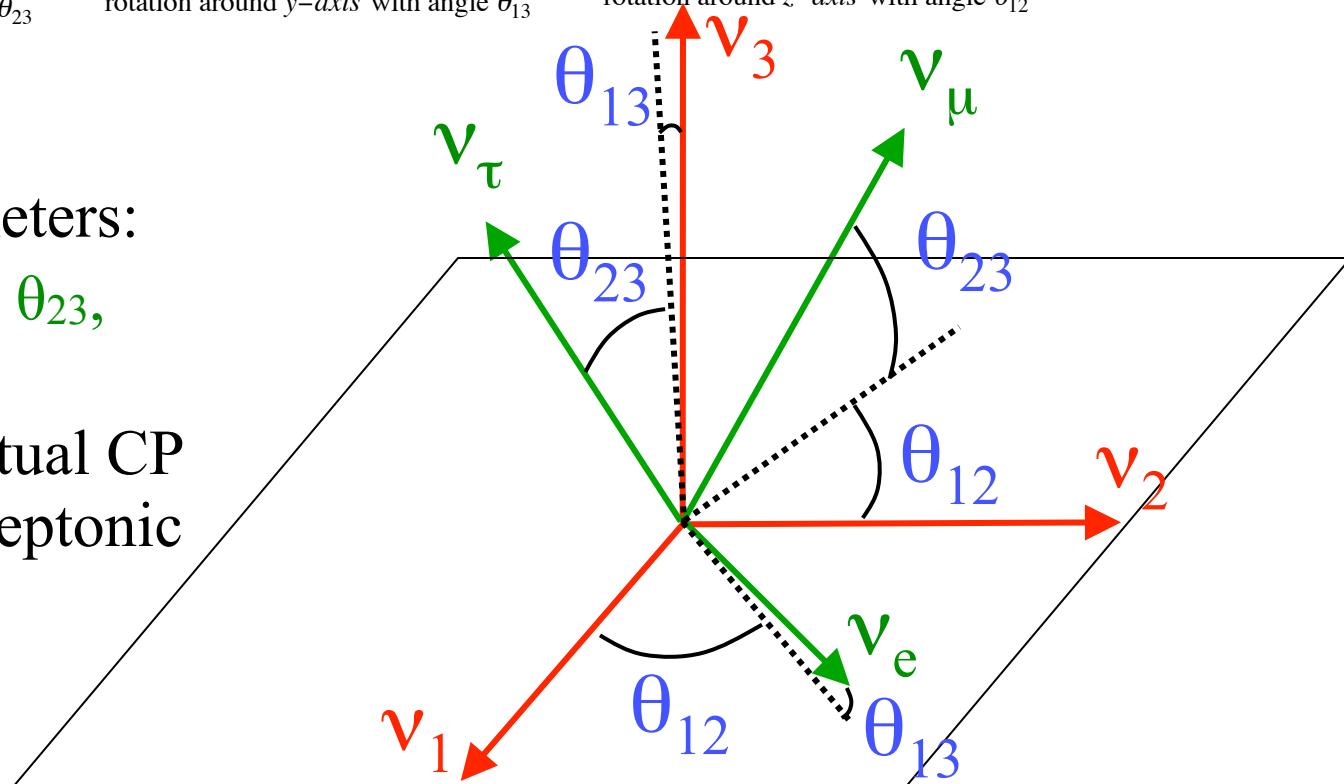
with $c_{ij} = \cos\theta_{ij}$ and $s_{ij} = \sin\theta_{ij}$

$$U_{\alpha i} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix}$$

$$= \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{rotation around } x\text{-axis with angle } \theta_{23}} \cdot \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix}}_{\text{rotation around } y\text{-axis with angle } \theta_{13}} \cdot \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{rotation around } z\text{-axis with angle } \theta_{12}}$$

Only 4 free parameters:

- 3 mixing angles: θ_{23} , θ_{12} , θ_{13}
- 1 phase for eventual CP violation in the leptonic sector: δ_{CP}



Transition Probability

$$|\nu_j(t)\rangle = e^{-iHt/\hbar} |\nu_j(0)\rangle$$

$$\Rightarrow P_{\nu_\alpha \rightarrow \nu_\beta} = \left| \delta_{\alpha\beta} + \sum_{k=2}^3 U_{\beta k} U_{\alpha k}^* \left[e^{-i \frac{\Delta m_{k1}^2 L}{2E}} - 1 \right] \right|^2$$

with: $\Delta m_{kj}^2 \equiv m_k^2 - m_j^2$ (the time has been replaced by the distance)



the transition probabilities do not depend on the particle masses but on the squared mass differences

Finally, the transition probabilities depend on the mixing matrix elements, the 2 squared mass differences and on the parameter L/E .

No transitions if:

$$m_\nu = 0 \text{ or}$$

$$\Delta m = 0 \text{ or}$$

$$\Delta m_{k1}^2 L / E \ll 1$$

oscillation parameters:

$\Delta m_{13}, \Delta m_{23}, \Delta m_{12}$ (only 2 are free)

$\theta_{13}, \theta_{12}, \theta_{23}$

Approximations according present measurements

$$\Delta m_{32}^2 = 2.45 \pm 0.09 \times 10^{-3} \text{ eV}^2$$

$$\Delta m_{21}^2 = 7.59_{-0.18}^{+0.20} \times 10^{-5} \text{ eV}^2$$

$$\sin^2 \theta_{23} = 0.51 \pm 0.06 \quad (\sim 45^\circ)$$

$$\sin^2 \theta_{12} = 0.312_{-0.015}^{+0.017} \quad (\sim 34^\circ)$$

$$\sin^2 \theta_{13} \leq 0.046 \quad (< 12.4^\circ)$$

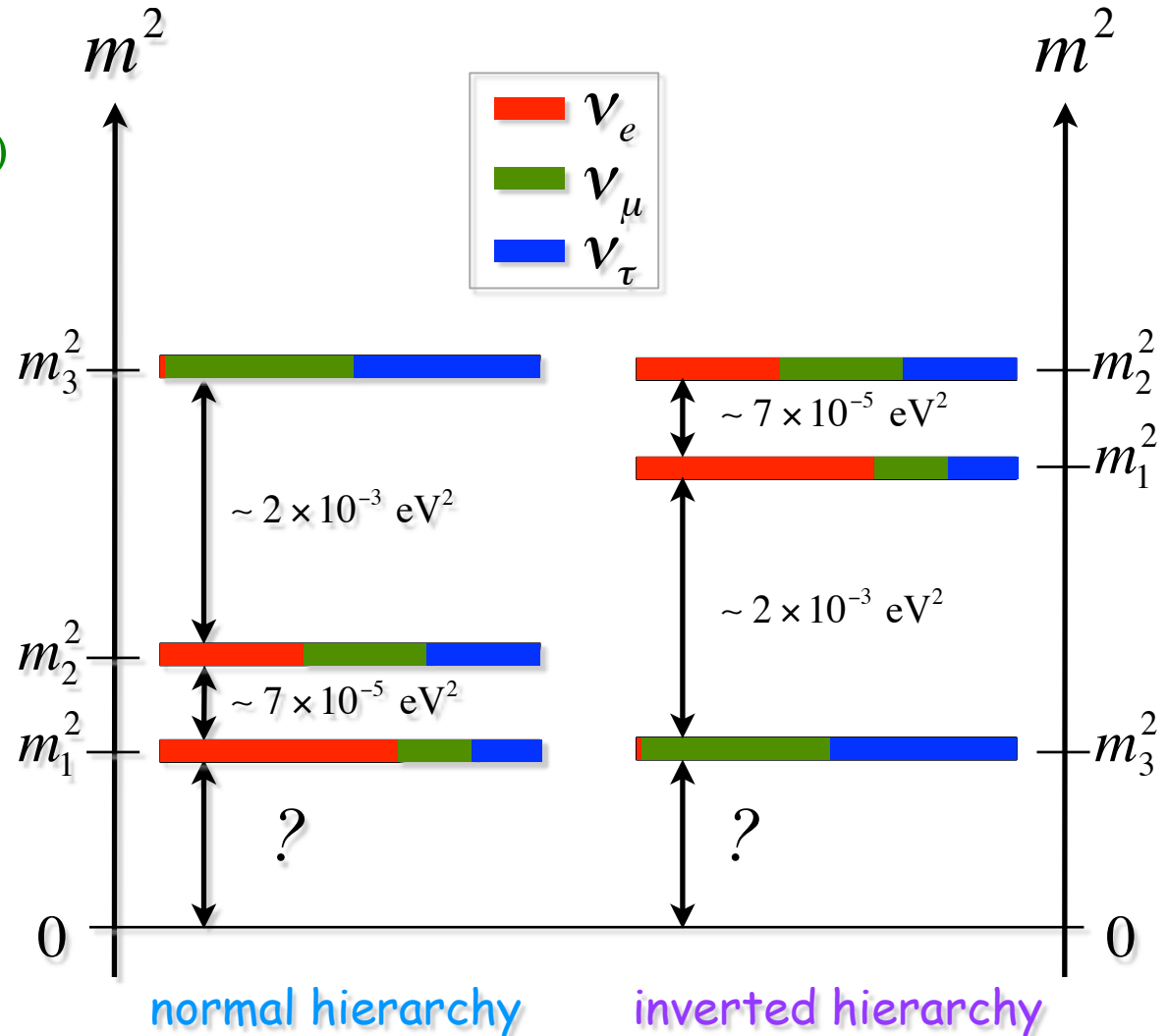
up to recently

We can then write:

$$\Delta m_{13} \approx \Delta m_{23} \equiv \Delta m$$

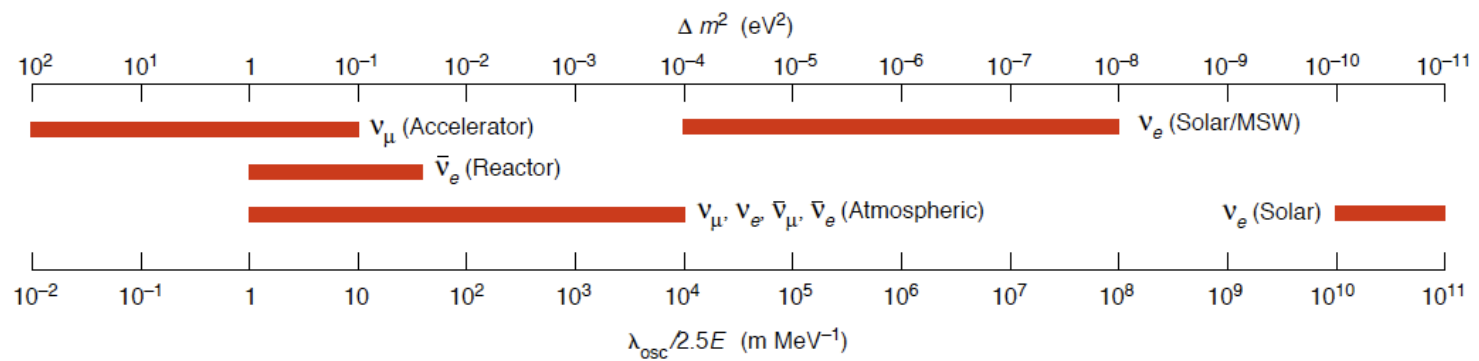
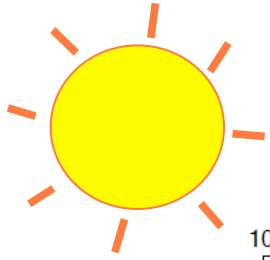
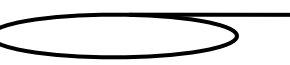
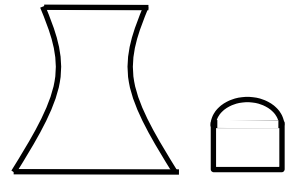
$$\Delta m_{12} \equiv \delta m$$

$$\Delta m \gg \delta m$$



L/E classification of the experiments

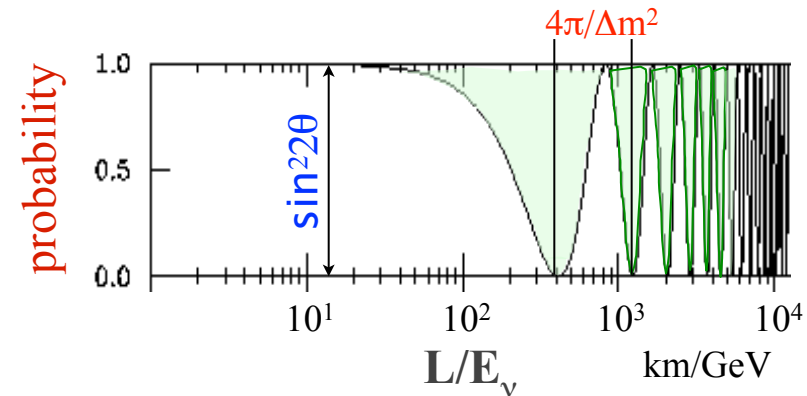
Experiment	L(m)	E (MeV)	L/E	Δm^2 (eV ²) (sensitivity)
SBL Reactors	10 ²	1	10 ²	10 ⁻²
LBL Reactors	10 ³	1	10 ³	10 ⁻³
SBL Accelerators	10 ³	10 ³	1	1
LBL Accelerators	10 ⁶	10 ³	10 ³	10 ⁻³
Atmospheric	10 ⁷	10 ³	10 ⁴	10 ⁻⁴
Solar	10 ¹¹	1	10 ¹¹	10 ⁻¹¹



Appearance/disappearance experiments

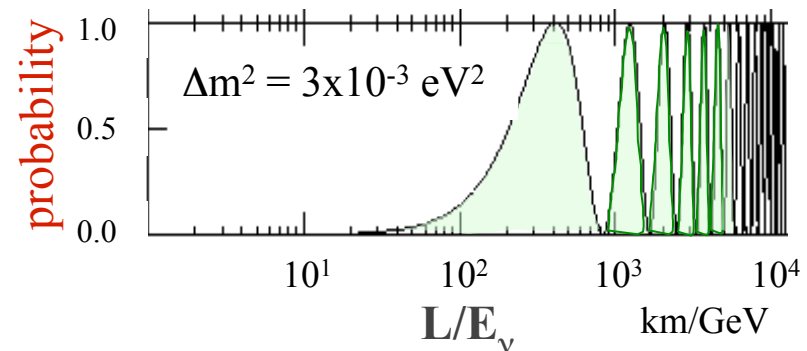
In disappearance experiments we count how many initial neutrinos ν_α survive after traveling over a distance L (*case of 2 flavours*):

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - \sin^2 2\theta \sin^2 \left(\frac{1.27 \Delta m^2 L}{E_\nu} \right)$$



In appearance experiments we look for neutrinos ν_β in a ν_α neutrino beam:

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \left(\frac{1.27 \Delta m^2 L}{E_\nu} \right)$$



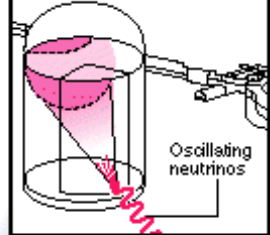
Atmospheric Neutrinos and oscillation confirmation

SuperK(first results in 1998)

Discovering Mass

The farther neutrinos travel, the more time they have to oscillate. By comparing the ratio of flavors of neutrinos coming "up" through the Earth to those coming from overhead, physicists determined that neutrinos oscillate, which neutrinos can only do if they have mass.

SUPER KAMIOKANDE DETECTOR



Oscillating neutrinos

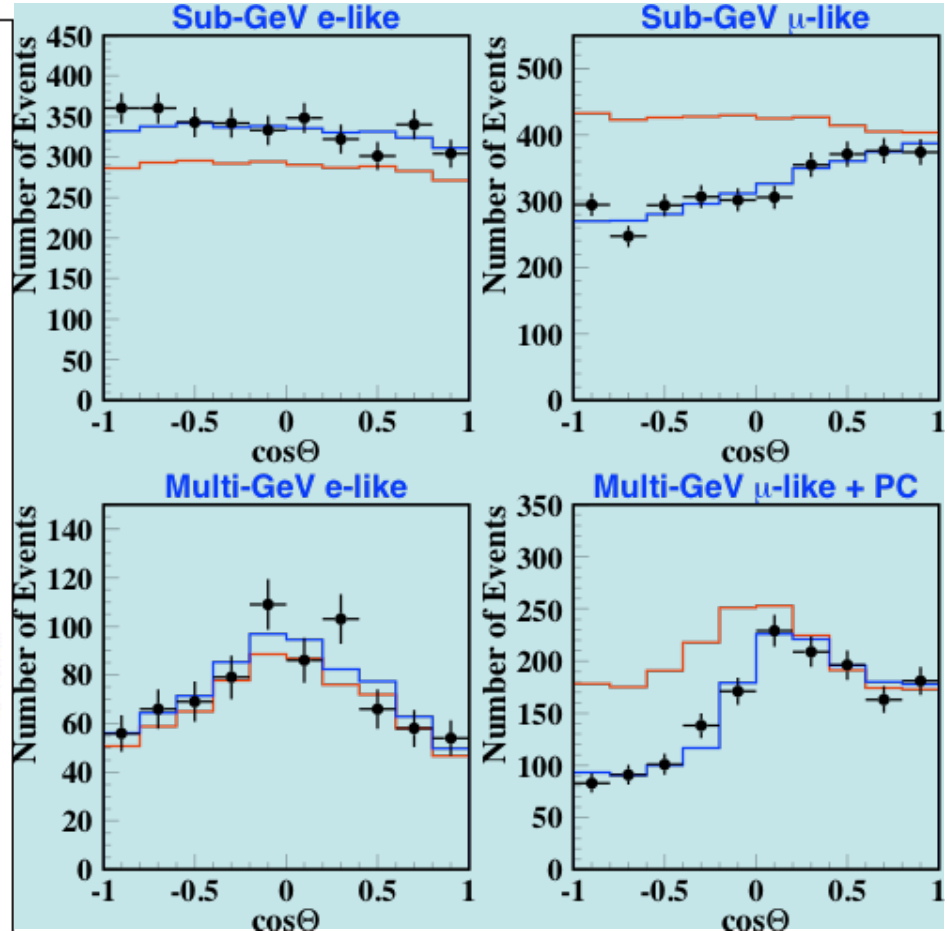
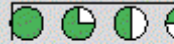
A cosmic ray (usually a proton) from space

3 A neutrino strikes another elementary particle in the detector tank. The interaction is recorded and analyzed by scientists to identify both the flavor of the neutrino and its flight path.

2 Neutrinos continue on the trajectory and begin to oscillate as they pass through the earth

1 The cosmic ray hits the earth's atmosphere, making a spray of secondary particles, some of which decay into neutrinos

One cycle of an oscillating neutrino as it passes through earth



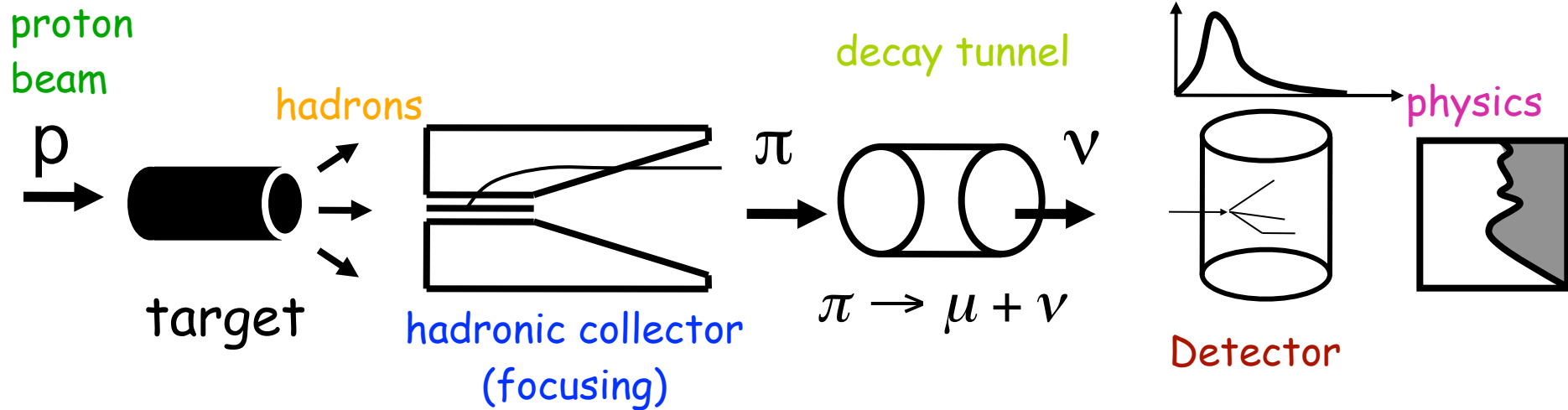
$\sim 13000\text{km}$

$\sim 500\text{km}$

$\sim 15\text{km}$

— no oscillation
— best fit

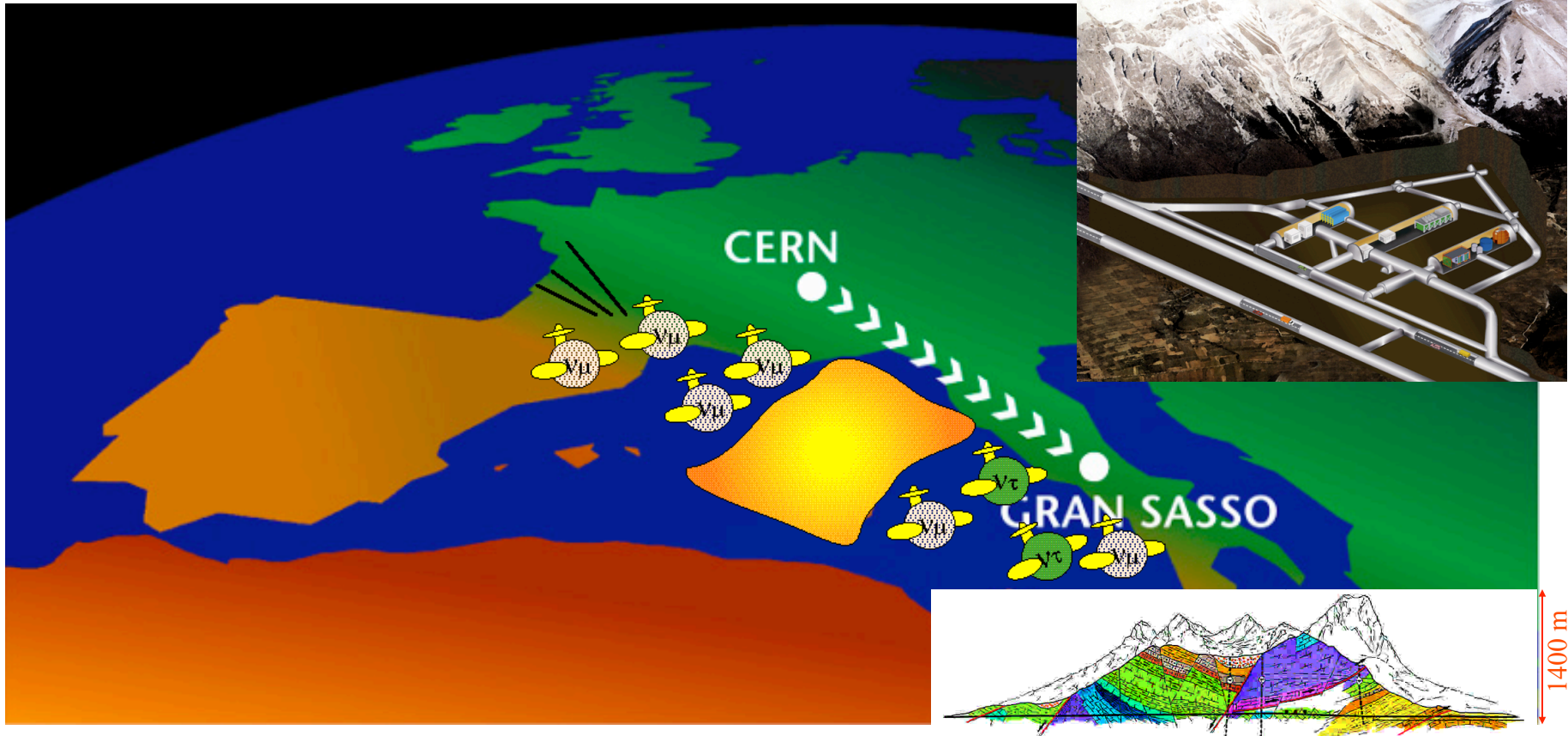
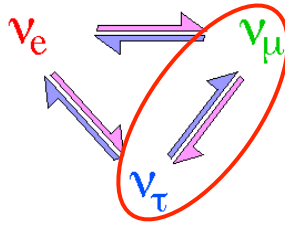
How can we produce a neutrino beam?



- adjustable proton energy,
- adjustable distance between the production point and the detector,
- neutrino type choice,
- but as usually, the neutrino energy is not well defined...

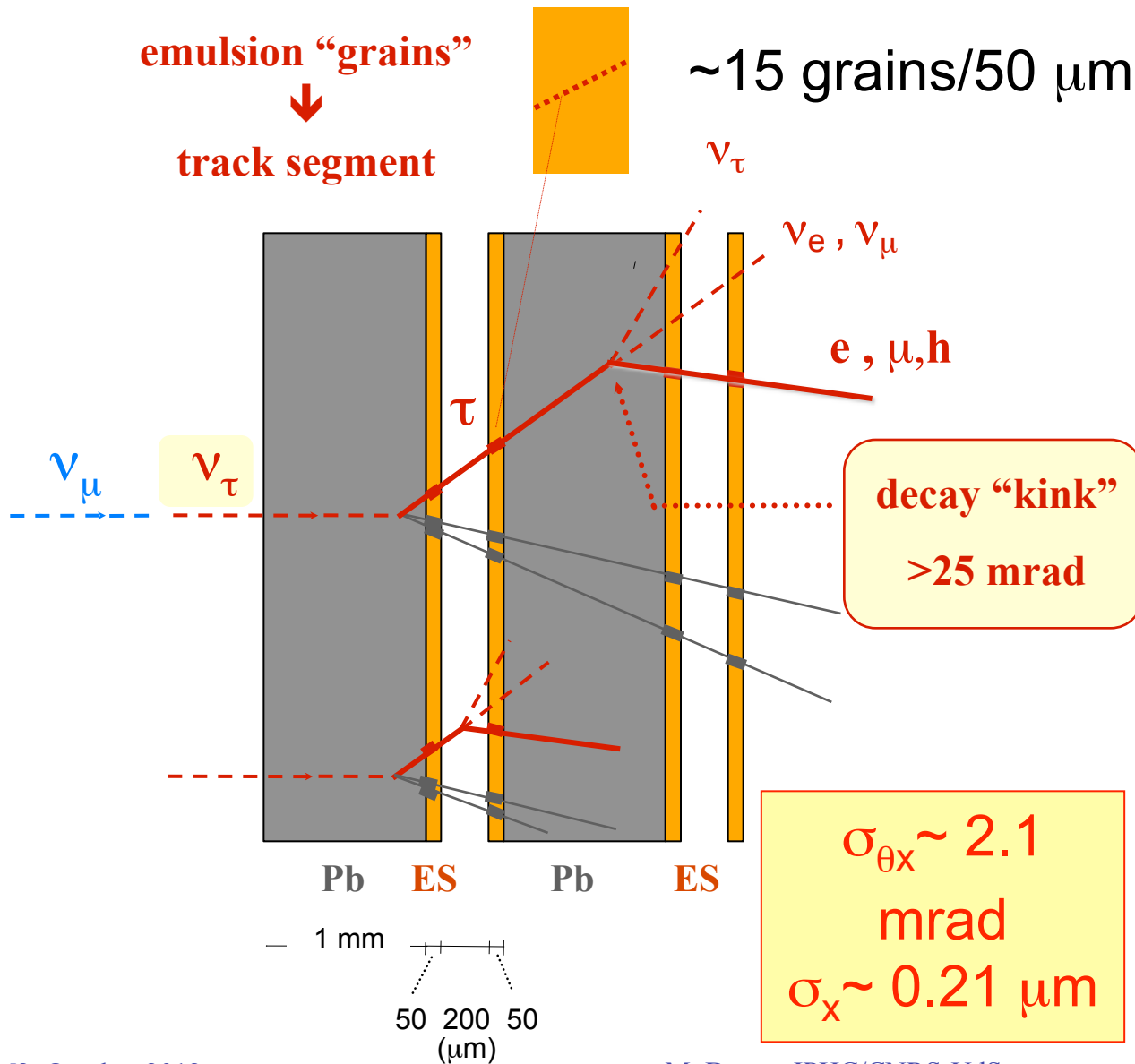
Long Base Line Appearance Experiment

up to now we have assumed
this oscillation valid:



CNGS ν_μ beam CERN-Gran Sasso 732 km (look for ν_τ in a pure ν_μ beam)

ν_τ detection in OPERA



A very high spatial resolution is necessary (do not forget that large surfaces have to be covered)

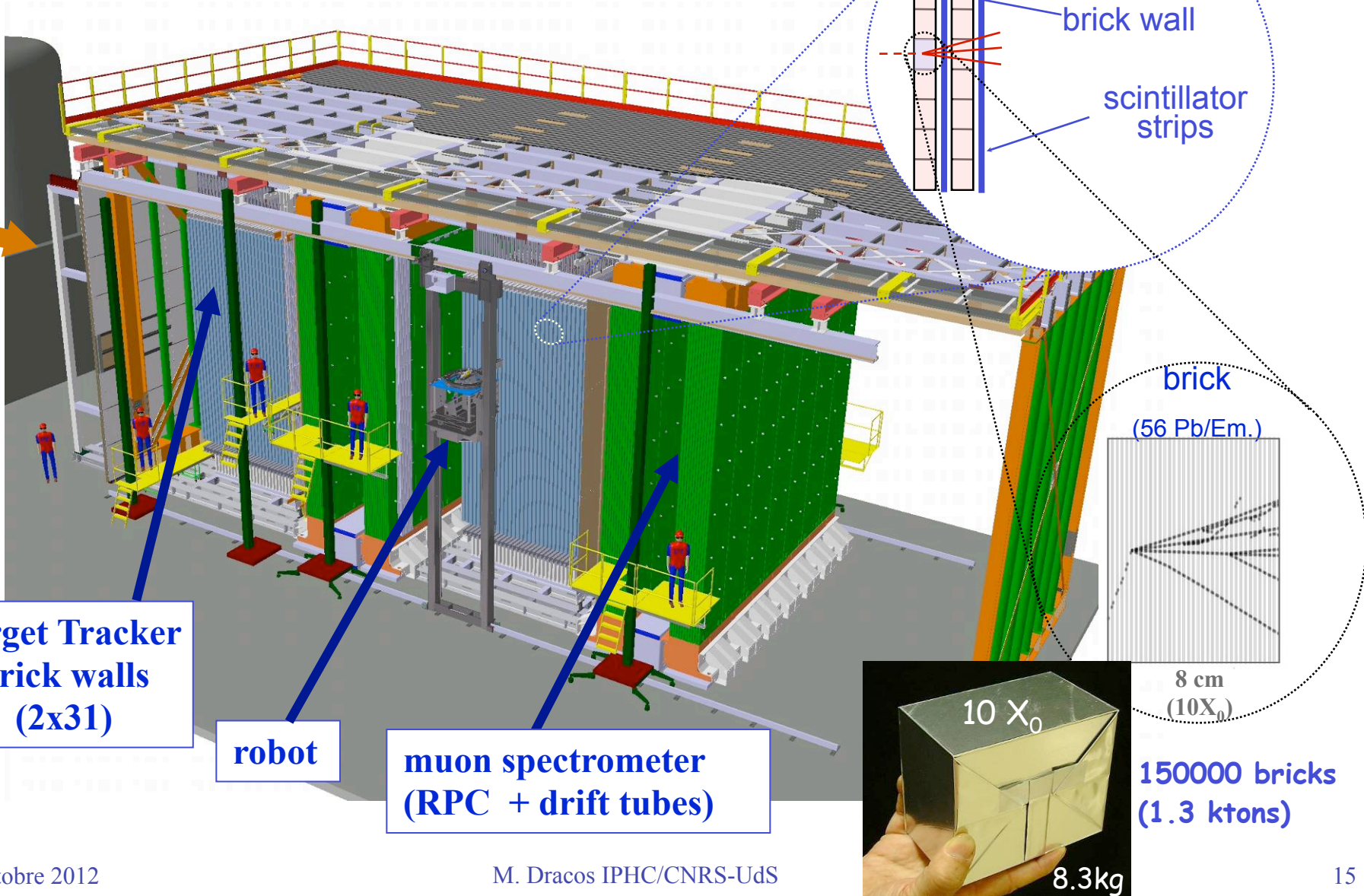


Nuclear Emulsions (photographic)



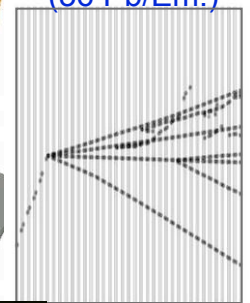


The OPERA Detector



“target” wall
 brick wall
 scintillator strips

brick
 (56 Pb/Em.)

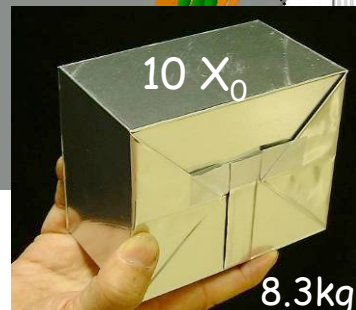


8 cm
 (10X₀)

Target Tracker
 brick walls
 (2x31)

robot

muon spectrometer
 (RPC + drift tubes)



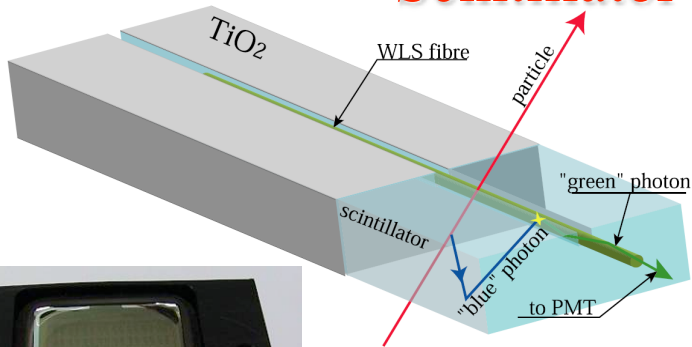
150000 bricks
 (1.3 ktons)

8.3kg

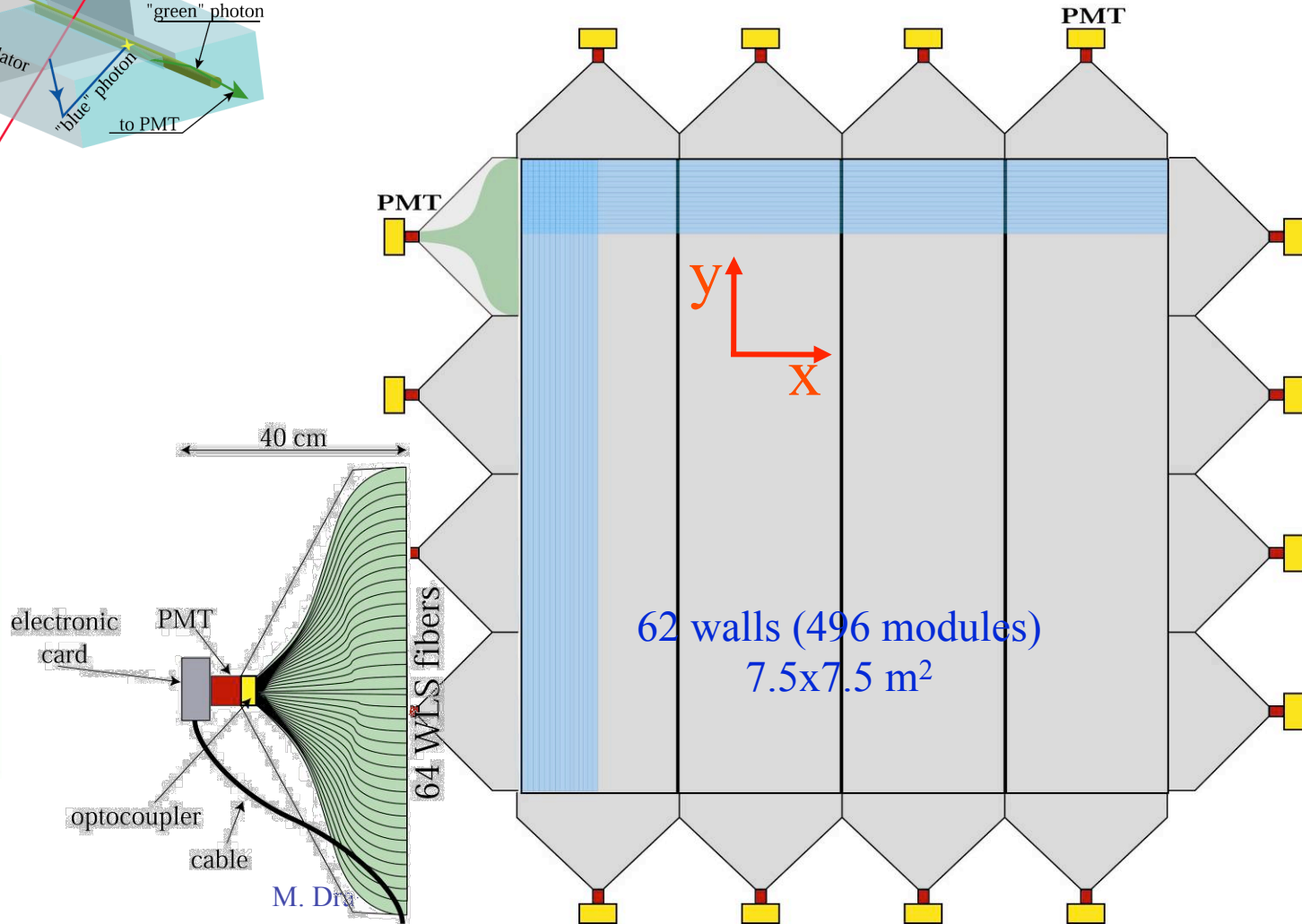
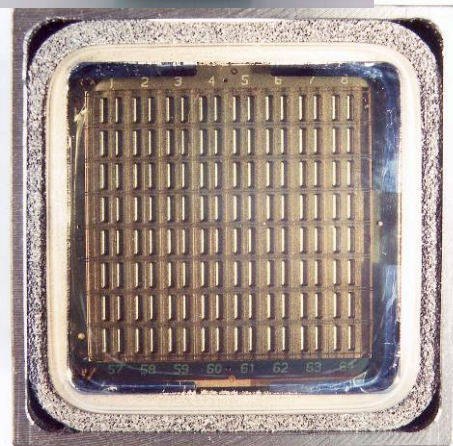
Target Tracker



Scintillator polystyrene strips (plastic)



(Bern, Brussels, Dubna, IPHC, LAL, Neuchâtel)



electronic card
PMT
optocoupler
cable

40 cm

64 WLS fibers

62 walls (496 modules)
 $7.5 \times 7.5 \text{ m}^2$

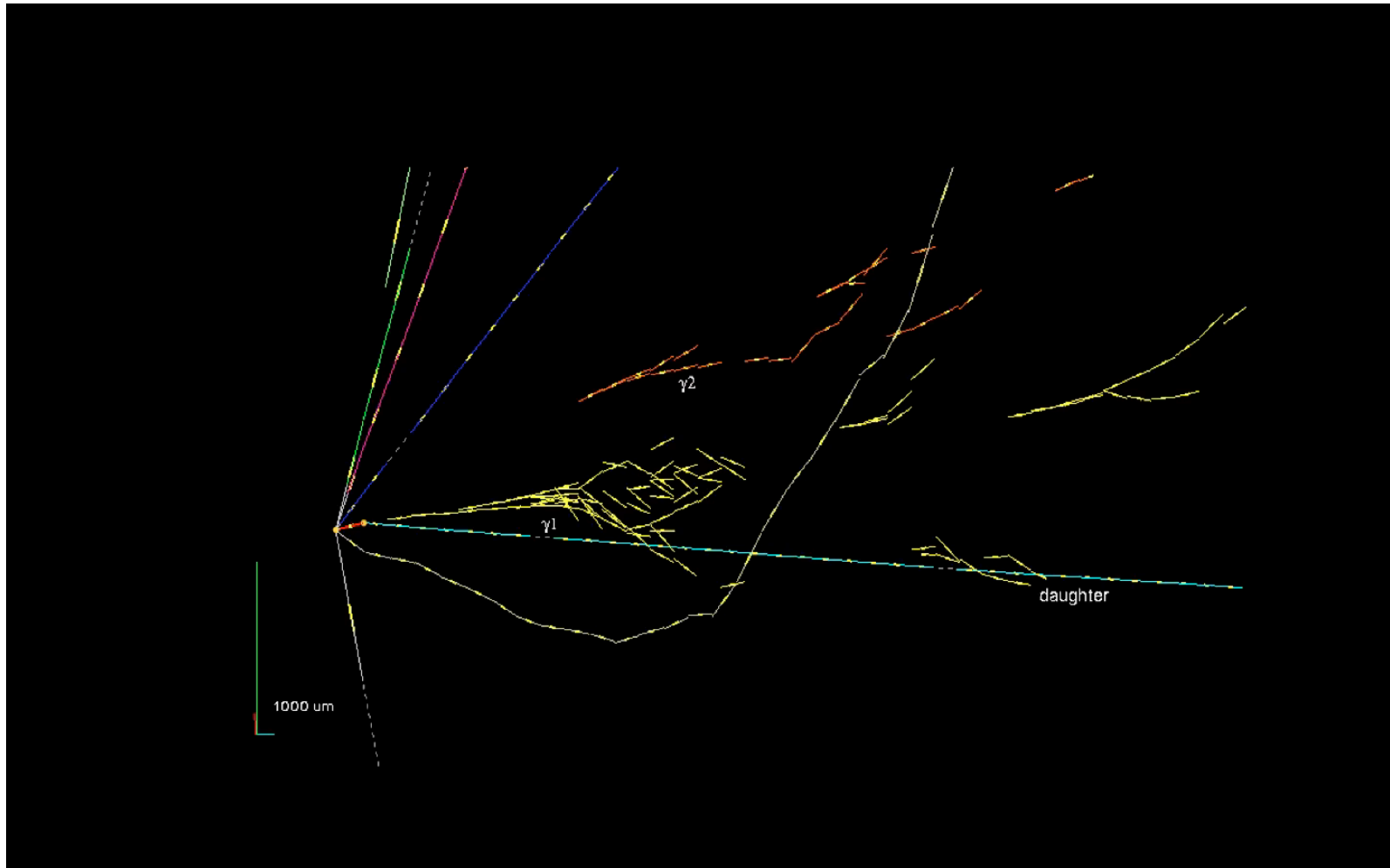
M. Dr

Hamamatsu MA-PMT
(64 channels) $3 \times 3 \text{ cm}^2$

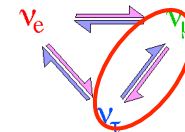
M2, Octobre 2012

The first ν_τ candidate event

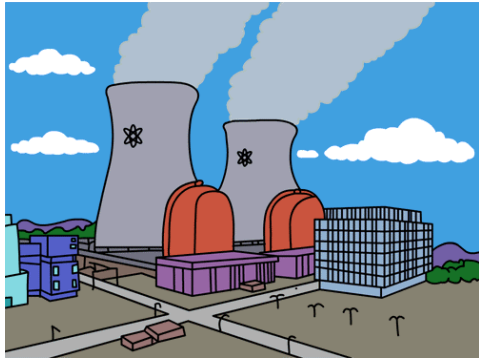
Phys. Lett. B 691 (2010) 138



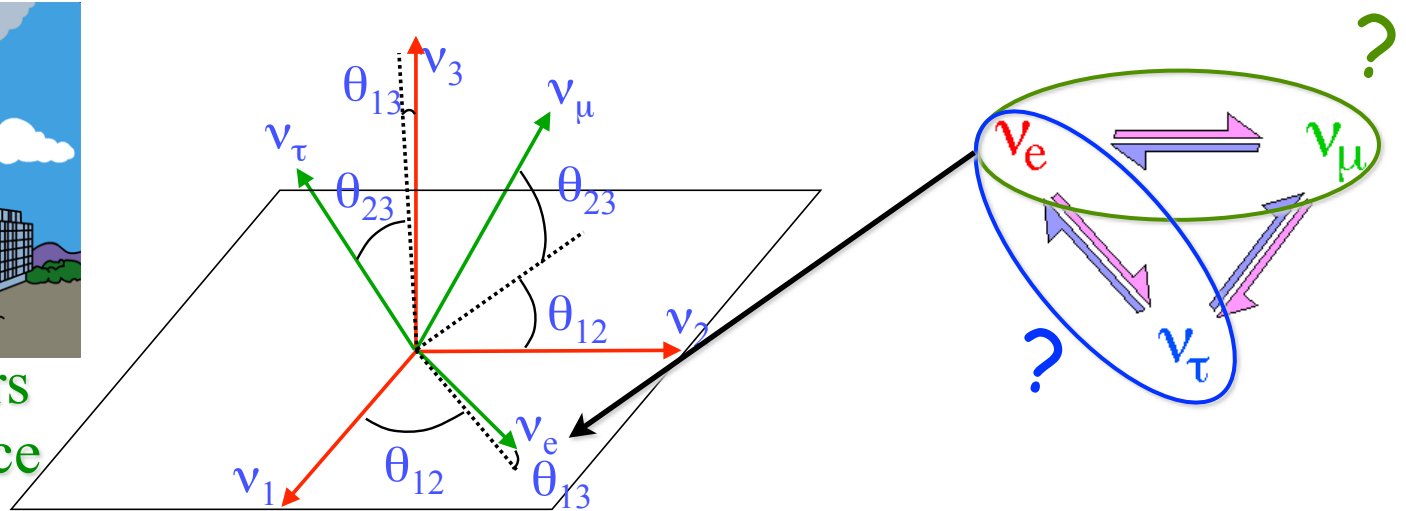
candidate: $\tau \rightarrow \underbrace{\pi^- \pi^0}_{\rho^-} \bar{\nu}_\tau$



The θ_{13} hunting



Nuclear Reactors
as neutrino source



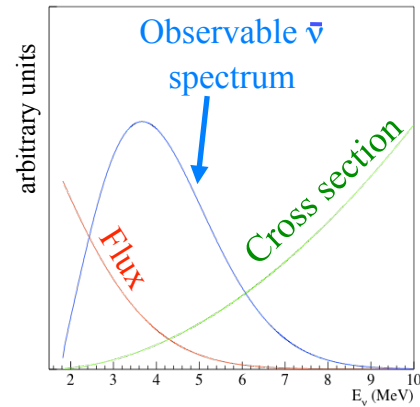
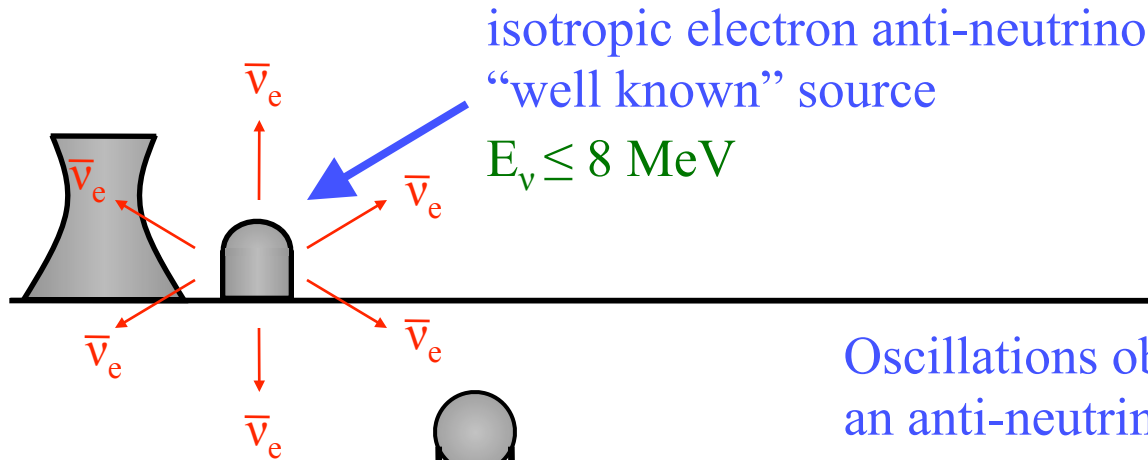
$$U = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{rotation around } x\text{-axis with angle } \theta_{23}} \cdot \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix}}_{\text{rotation around } y\text{-axis with angle } \theta_{13}} \cdot \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{rotation around } z\text{-axis with angle } \theta_{12}}$$

atmospheric,
accelerators

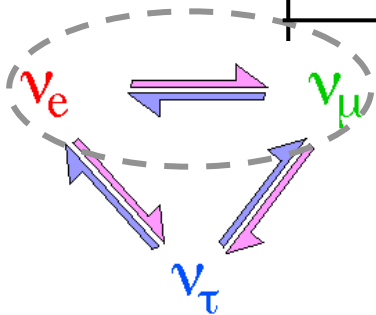
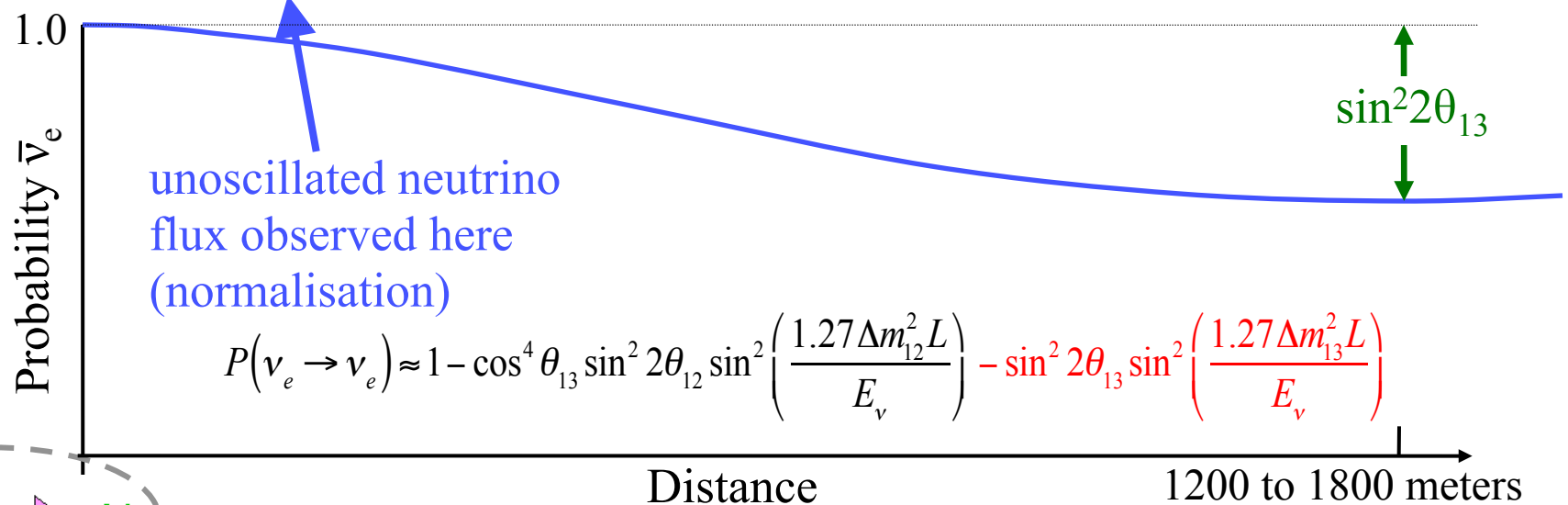
if $\theta_{13}=0$ we could not measure
CP violation

solar,
reactors

$\sin^2 2\theta_{13}$ and the reactor experiments

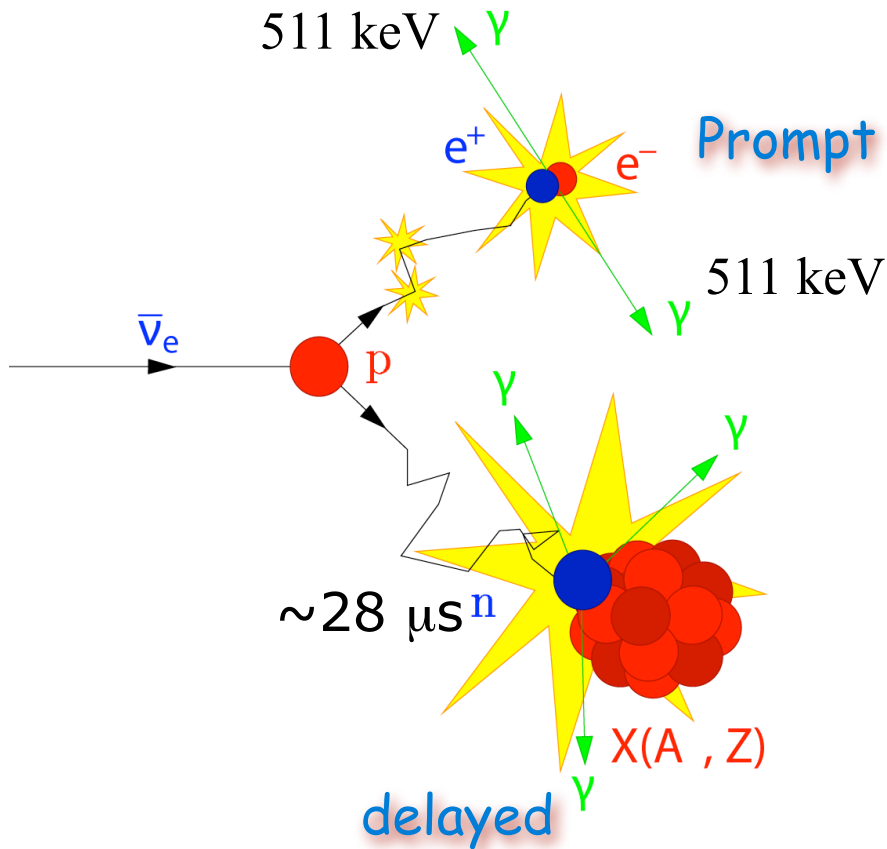
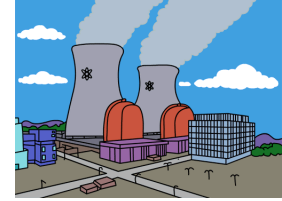


Oscillations observed as an anti-neutrino deficit



$$P_{surv.}(E_\nu, d_{far}) = \frac{N_{p,near}}{N_{p,far}} \times \frac{\epsilon_{near}}{\epsilon_{far}} \times \left(\frac{d_{far}}{d_{near}} \right)^2 \times \frac{N_{far}}{N_{near}}$$

Inverse β decay and detection mode



In a pure scintillator the neutron will be captured by hydrogen:

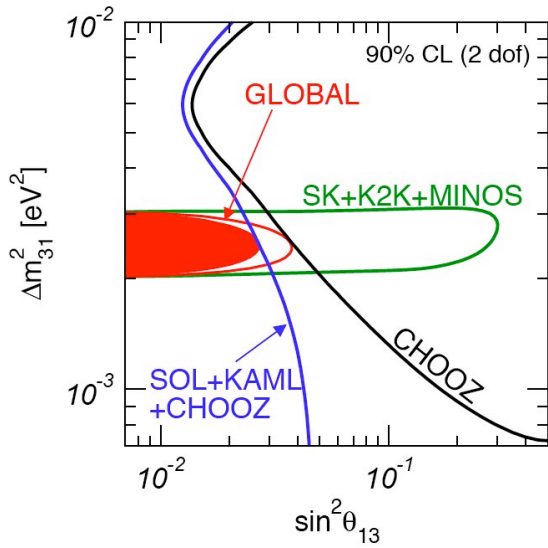
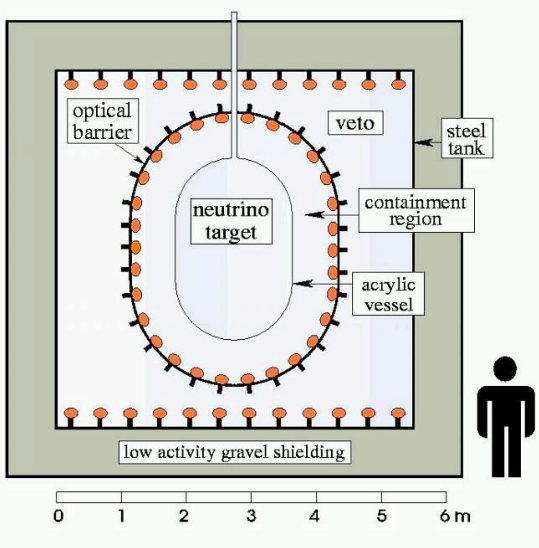
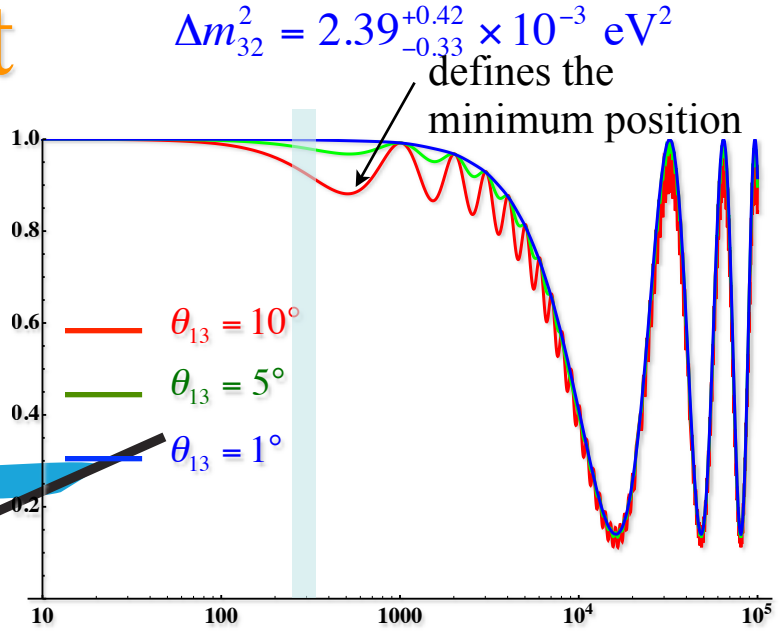
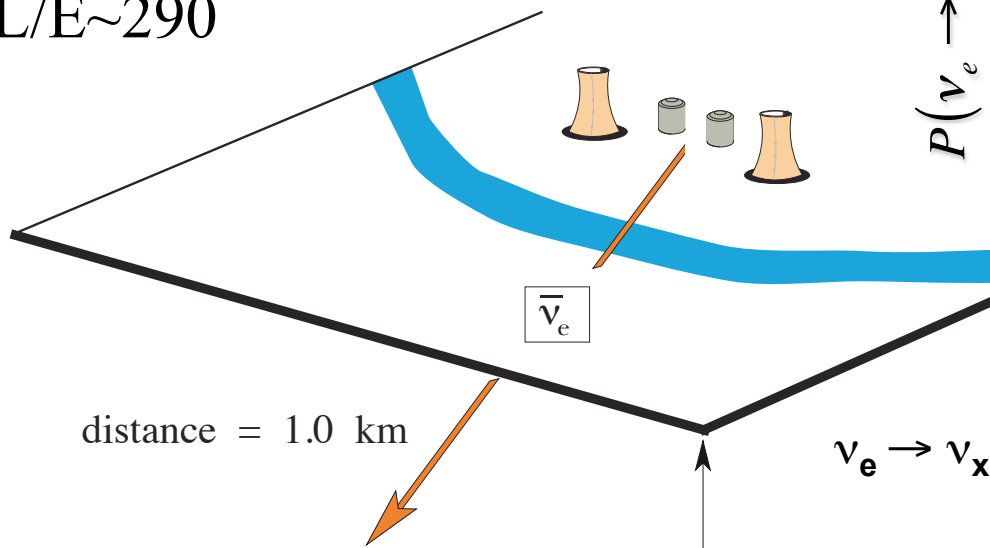


Very often the scintillator is doped with gadolinium that increase the capture probability and liberates more γ 's:



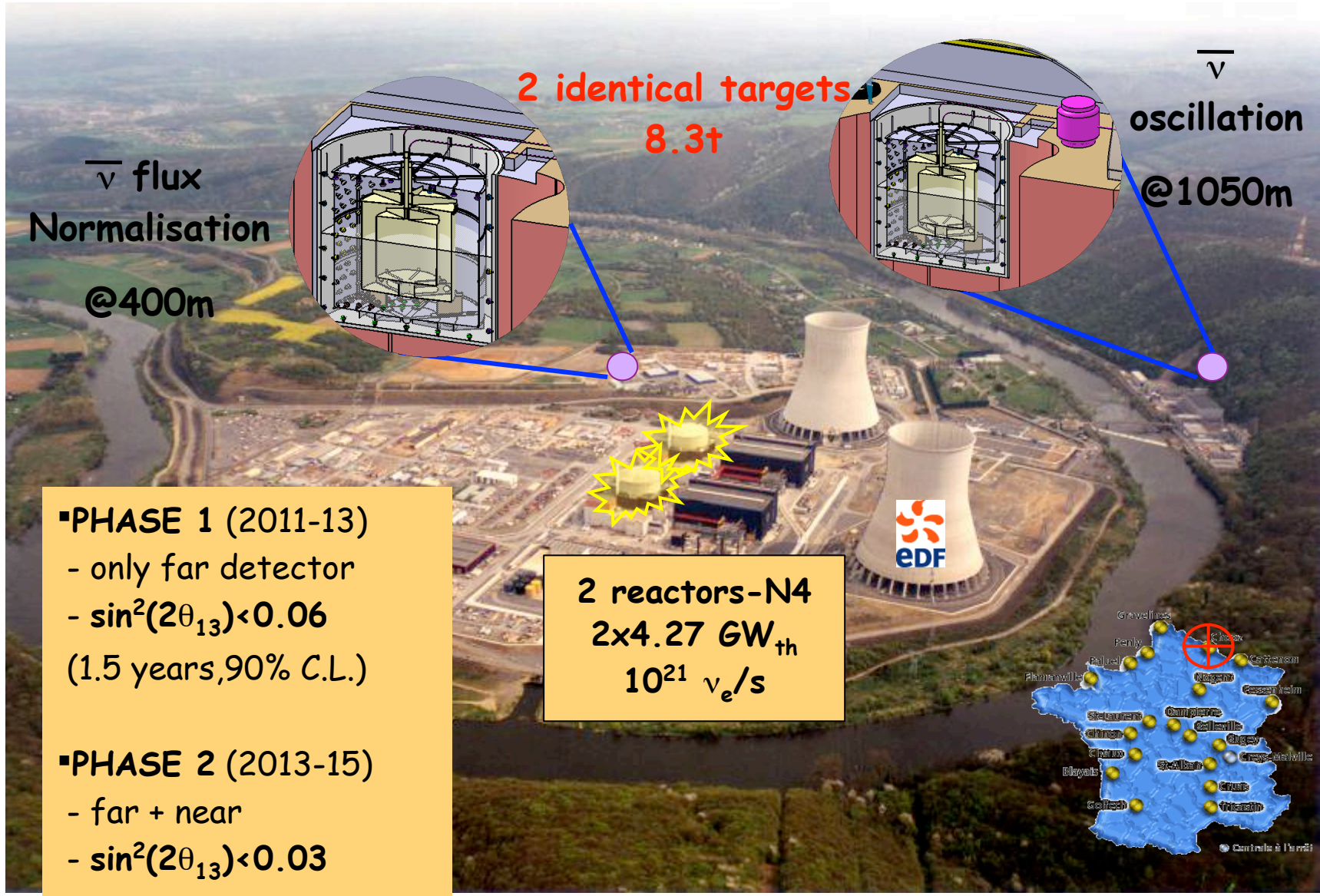
The CHOOZ experiment

L/E ~ 290



$\sin^2(2\theta_{13}) < 0.12 - 0.2$
(90% C.L.)

Double Chooz site



2 identical targets
8.3t

$\bar{\nu}$ flux
Normalisation
@400m

$\bar{\nu}$
oscillation
@1050m

- **PHASE 1 (2011-13)**
 - only far detector
 - $\sin^2(2\theta_{13}) < 0.06$
 - (1.5 years, 90% C.L.)
- **PHASE 2 (2013-15)**
 - far + near
 - $\sin^2(2\theta_{13}) < 0.03$
 - (3.5 years, 90% C.L.)

2 reactors-N4
2x4.27 GW_{th}
10²¹ ν_e/s



DC detector



$$\bar{\nu}_e + p \rightarrow e^+ + n$$



Target ν : 10,3 m³

80% C₁₂H₂₆ + 20% PXE + PPO + Bis-MSB
+0.1% Gd

γ Catcher : 22,6 m³

80% C₁₂H₂₆ + 20% PXE + PPO + Bis-MSB

Non scintillating Buffer : 114 m³

Mineral oil

Buffer vessel & 390 10" PMTs :

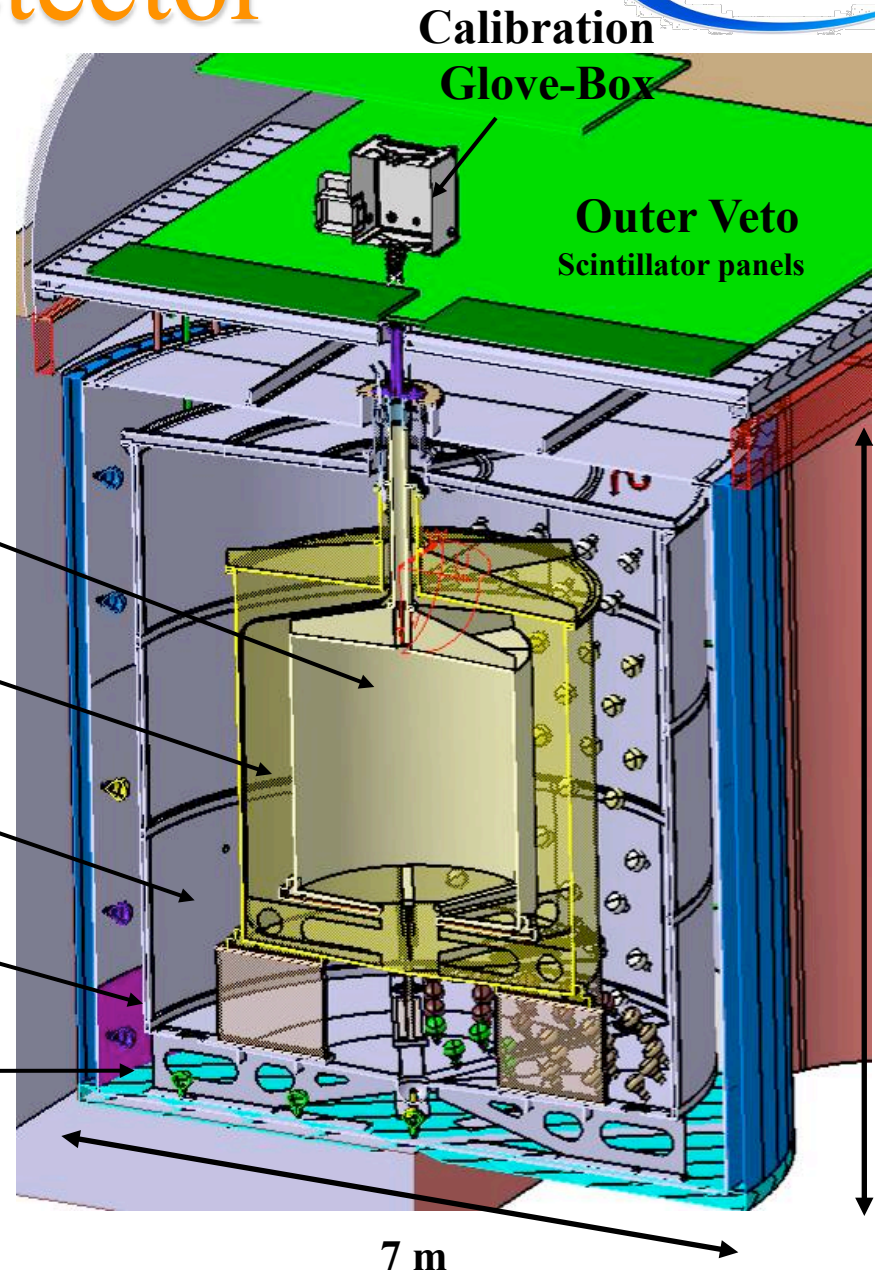
Stainless steel 3 mm

Inner Muon Veto : 90 m³

Mineral oil + 78 8" PMTs

Steel Shielding :

15 cm steel, All around



7 m

7 m

Background (key element)



Accidental:

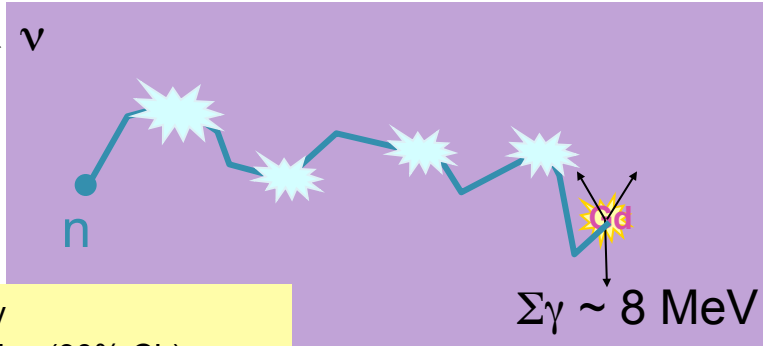
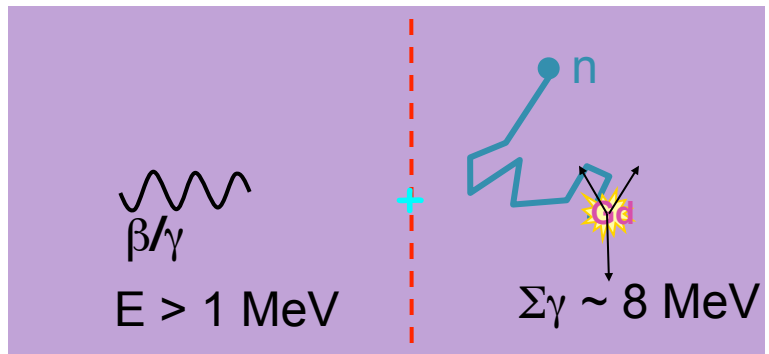
gamma or beta events with $E > 1 \text{ MeV}$
 +
 neutron capture by Gd, $E \sim 8 \text{ MeV}$

*radiopurity of detector components,
 shielding against external radiation sources,
 the "single"s rate can be measured online*

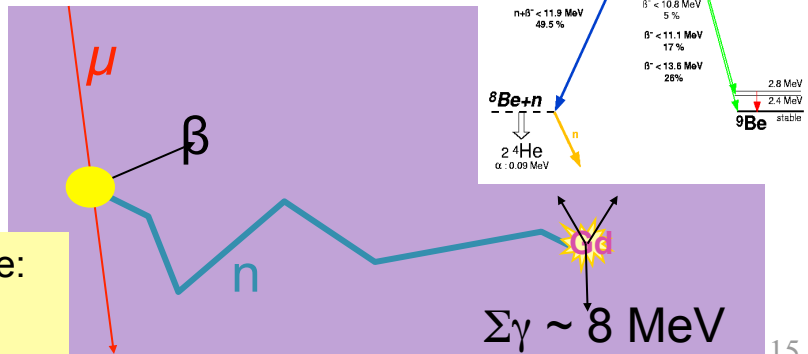
Correlated:

produced by muons and their secondaries
 •fast neutrons (produced by μ in the surrounding rock)
 •beta-neutron cascades (${}^9\text{Li}$, ${}^8\text{He}$): produced by the μ or n interactions with ${}^{12}\text{C}$
 mean lifetime $\sim (0.1 - 1) \text{ s}$

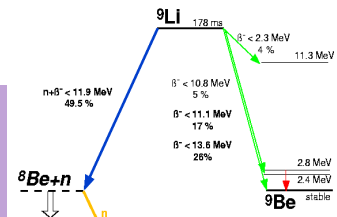
*shielding against cosmic rays,
 active veto to recognise μ and n ,
 measurement of the background with the reactors off (if and when possible...).*



Chooz: $\sim 1/\text{day}$
 •far: $N_b < 0.6/\text{day}$ (90% CL)
 •near: $N_b \sim 3.3/\text{day}$ (90% CL)



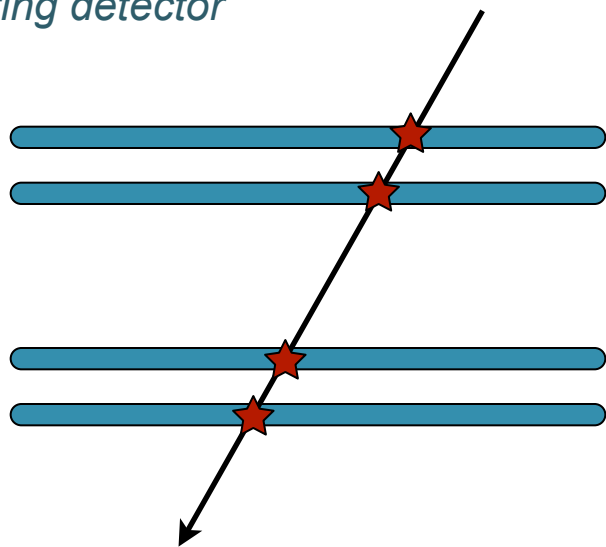
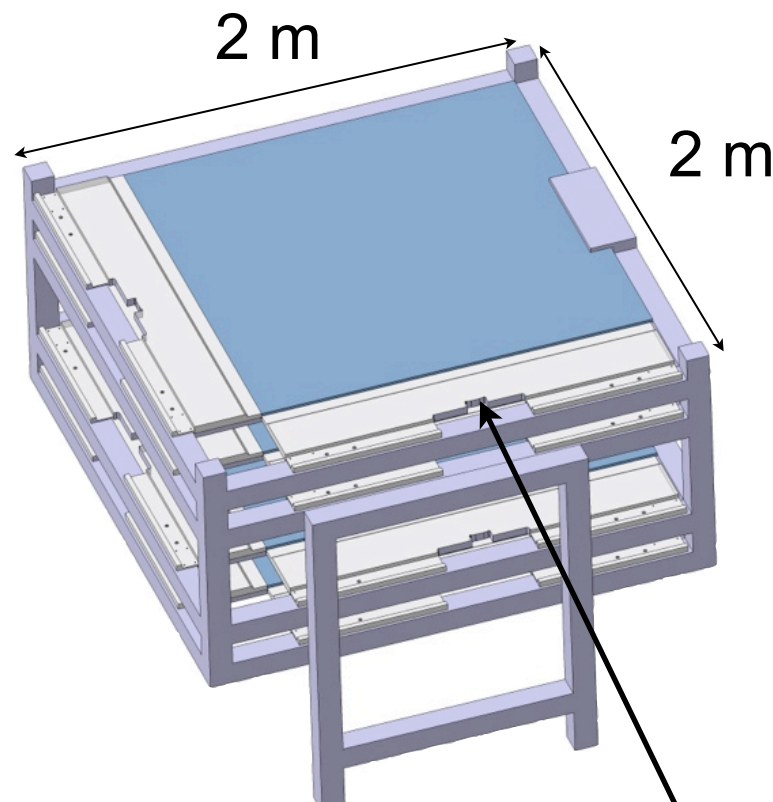
expected rate:
 •far: 1.4/day
 •near: 9/day



Muon telescope



- *In the near detector the cosmic background will be more important due to the low depth*
- *Better cosmic ray simulation needed*
- *Use a muon telescope to well measure the the muon spacial distribution and introduce it in the simulation*
- *4 layers of plastic scintillators (OPERA)*
- *rotating detector*

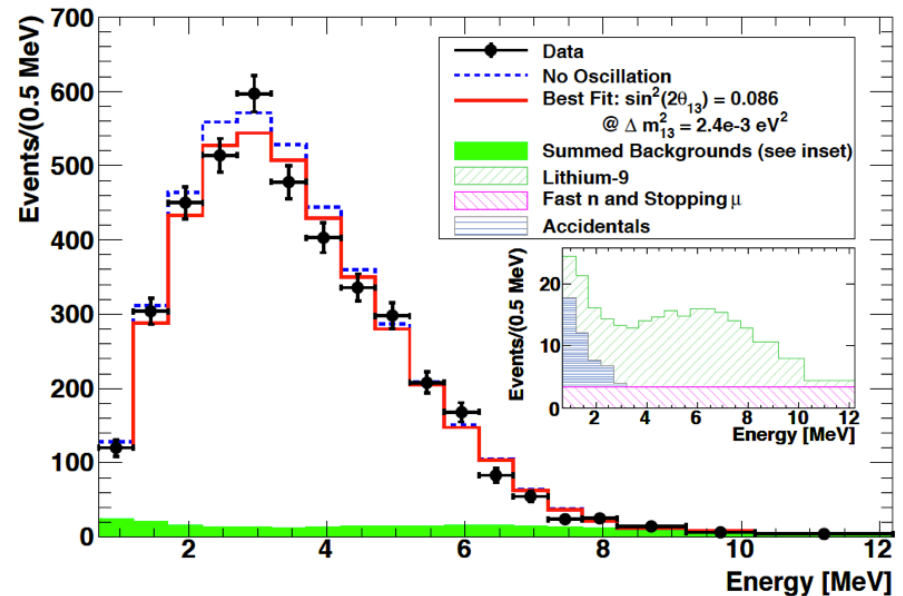
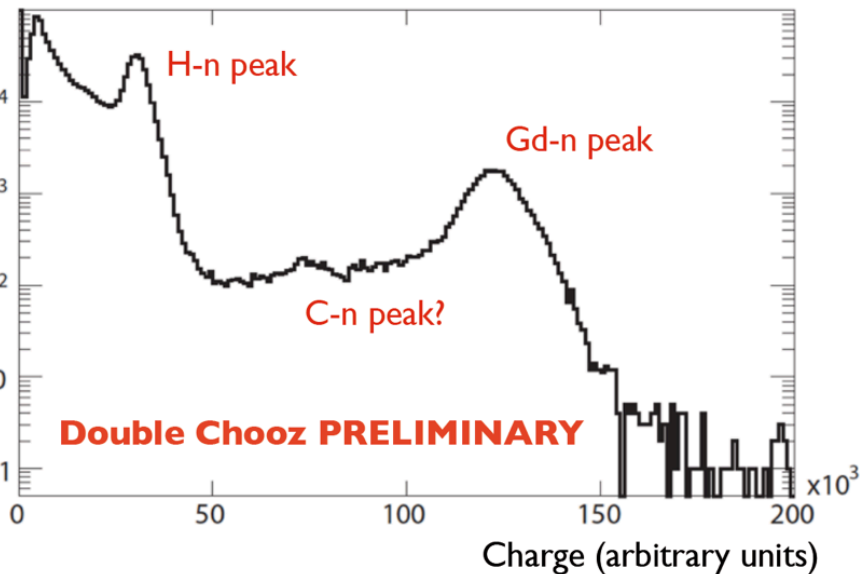


to be constructed during the training

PMT



First DC results (Nov. 2011)



- Rate + Shape Analysis:

- ◆ $\sin^2(2\theta_{13}) = 0.085 \pm 0.029(\text{stat}) \pm 0.042(\text{syst})$

- Rate Only:

- ◆ $\sin^2(2\theta_{13}) = 0.093 \pm 0.029(\text{stat}) \pm 0.073(\text{syst})$



Future Projects on Neutrino Oscillations

(essentially to discover CP violation
and measure the mass hierarchy)

Neutrino Related European Projects

Two FP7 projects:

- **EUROv** (<http://www.euronu.org/>)

- Design Study for new neutrino facilities in Europe
- Project started: 1st September 2008
- Duration: 4 years – completion August 2012



- **LAGUNA-LBNO** (Large Apparatus for Grand Unification and Neutrino Astrophysics, <http://laguna.ethz.ch:8080/Plone>)

- Design Study for large underground laboratories for astroparticle and neutrino studies
- Project started: 2011
- Duration: 3 years – completion in 2014

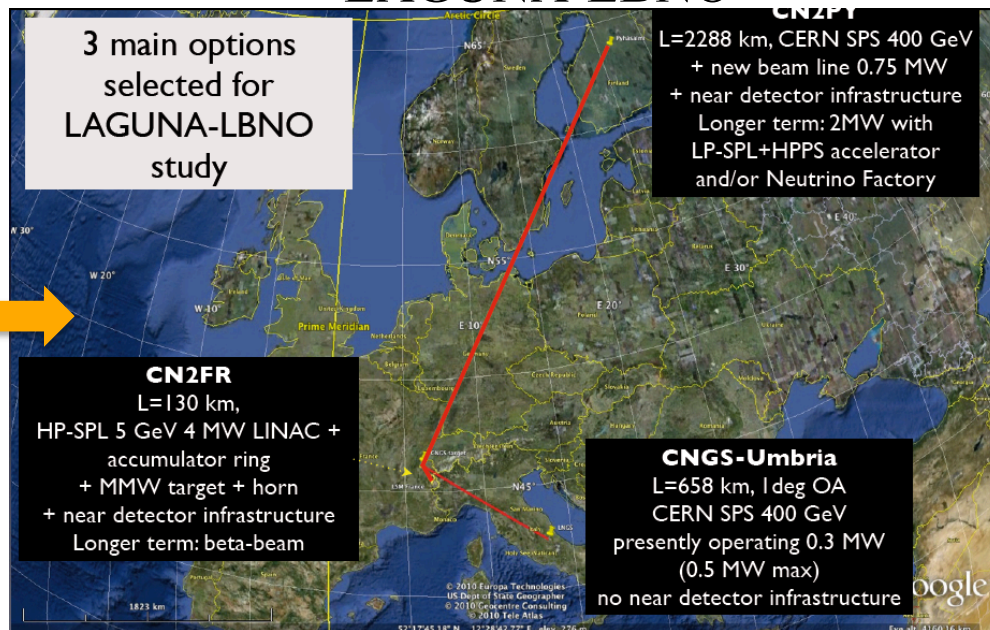


All sites and detection techniques under consideration by LAGUNA and possible neutrino beams from CERN

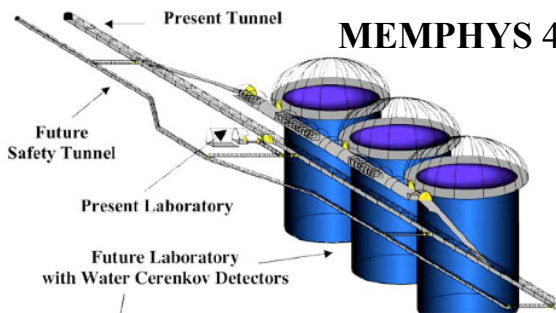
LAGUNA



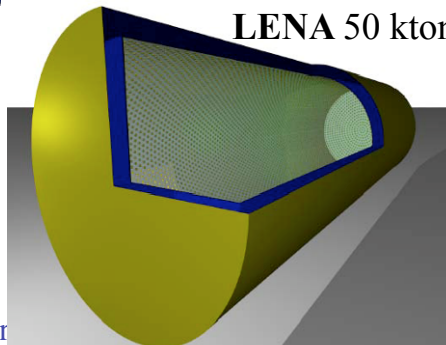
LAGUNA-LBNO



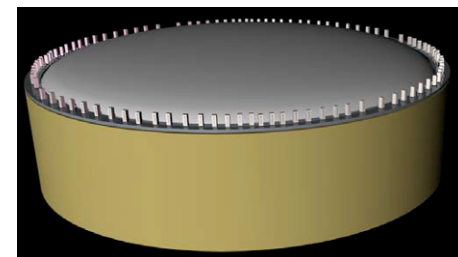
MEMPHYS 450 Ktons



LENA 50 ktons



GLACIER 100 ktons



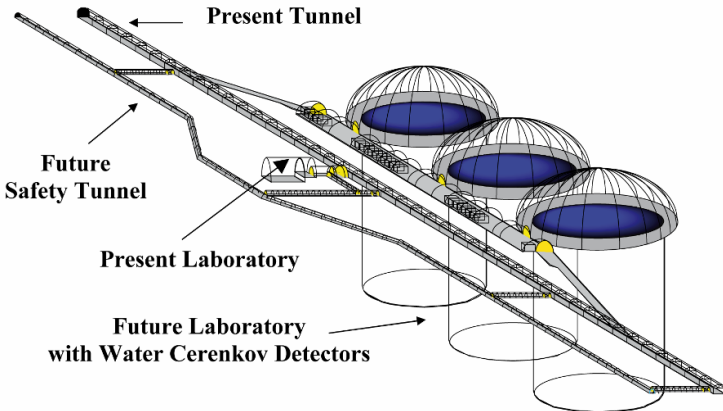
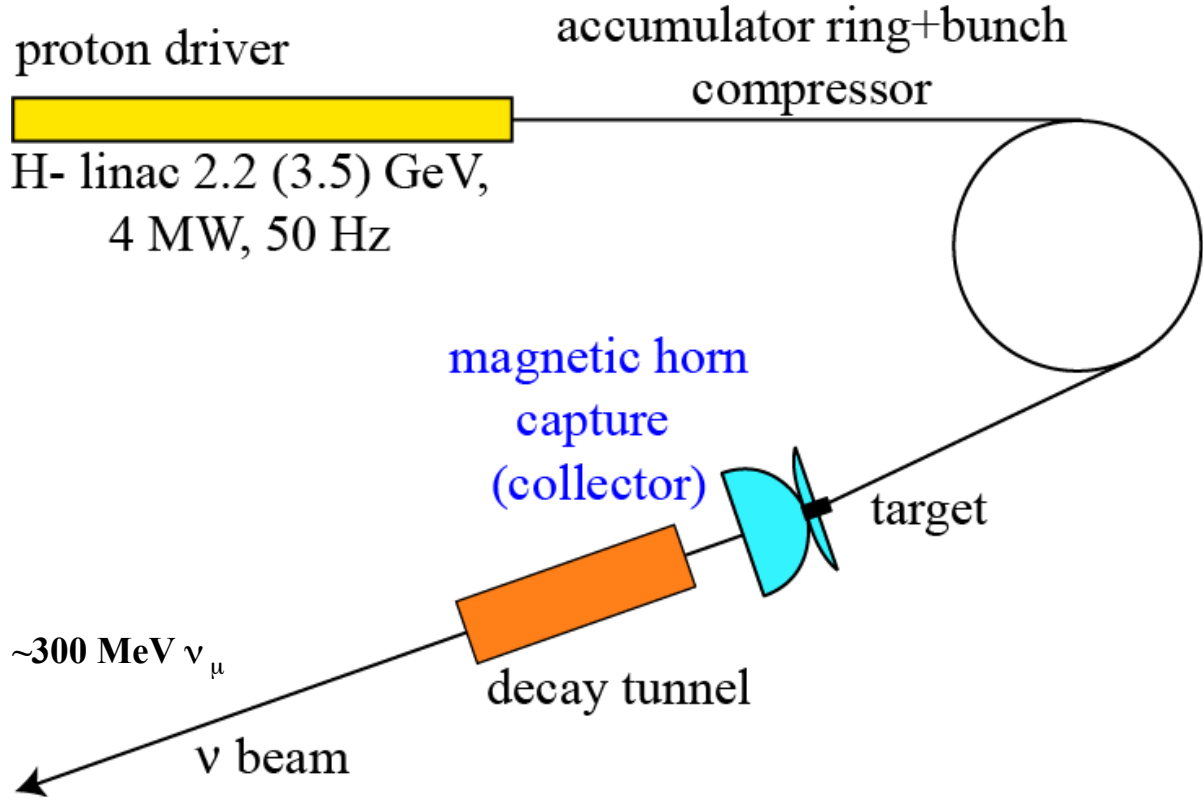


Super-Beams

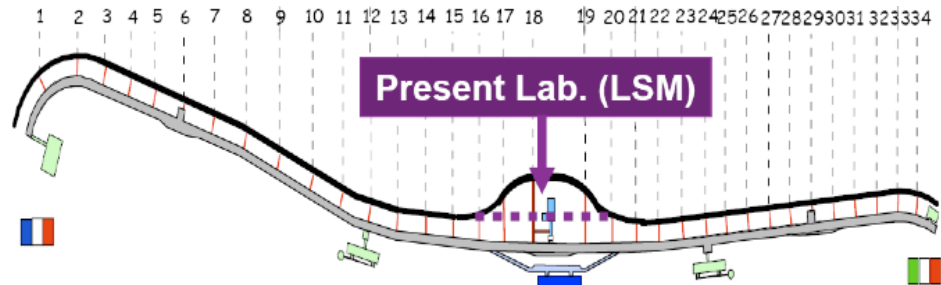
CERN-SPL



MEMPHYS detector (Water Cherenkov)

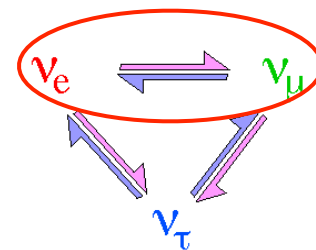
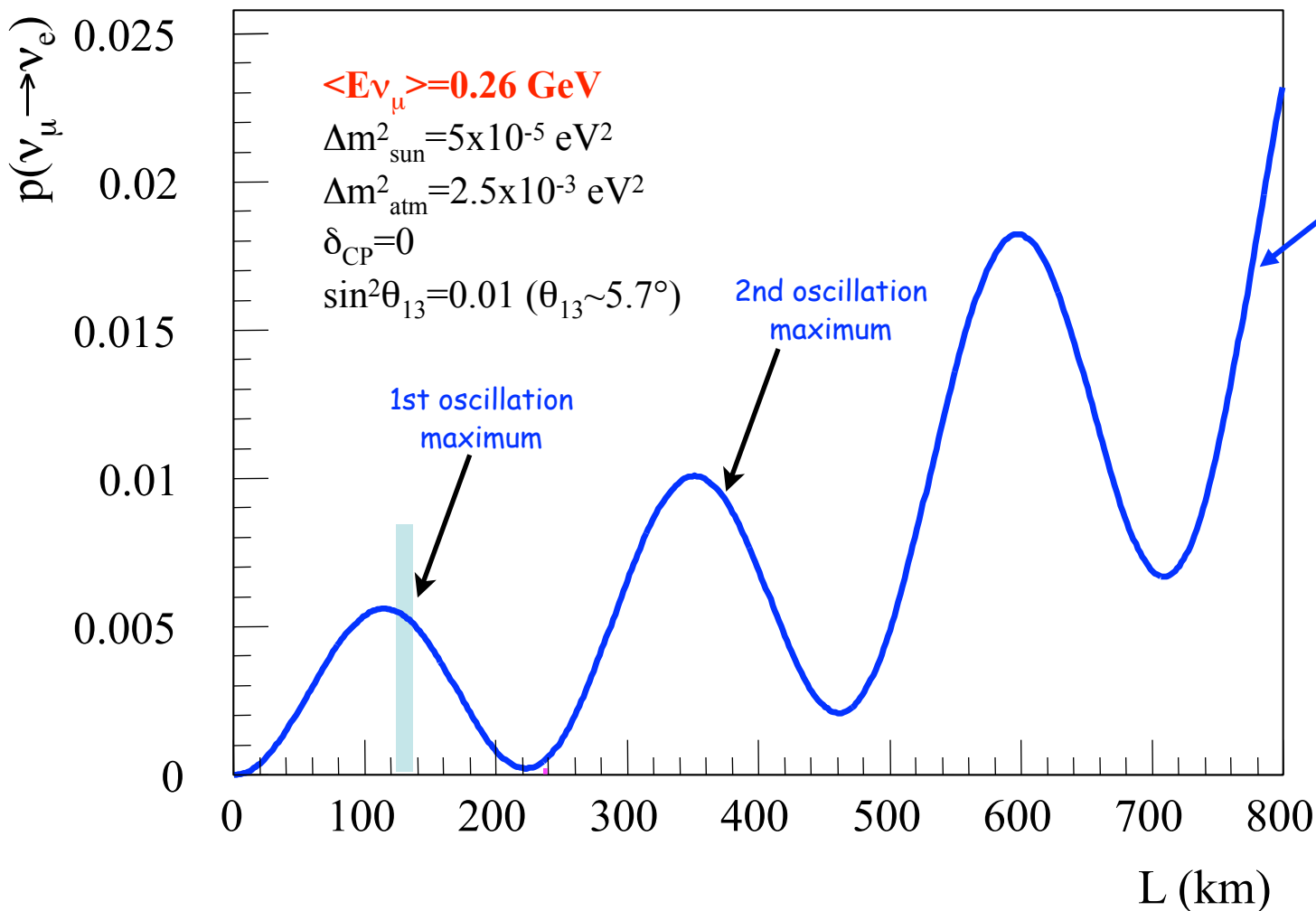


- 3 shafts= 440 kT fid., 1 shaft= 4xSK
- 30% coverage by 81k PMT 12" per shaft



Why this Super-Beam?

$\ell = 130$ km (distance CERN-Fréjus)



without too many approximations

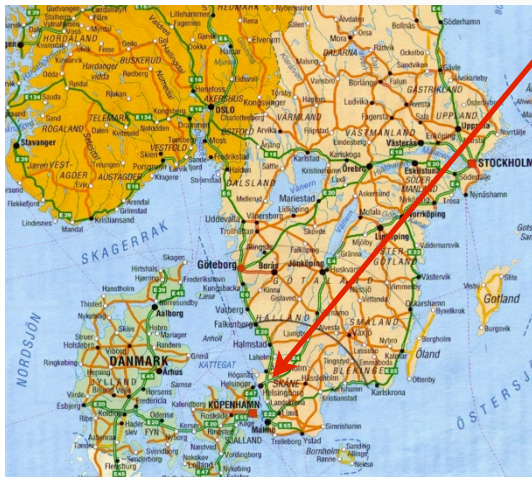
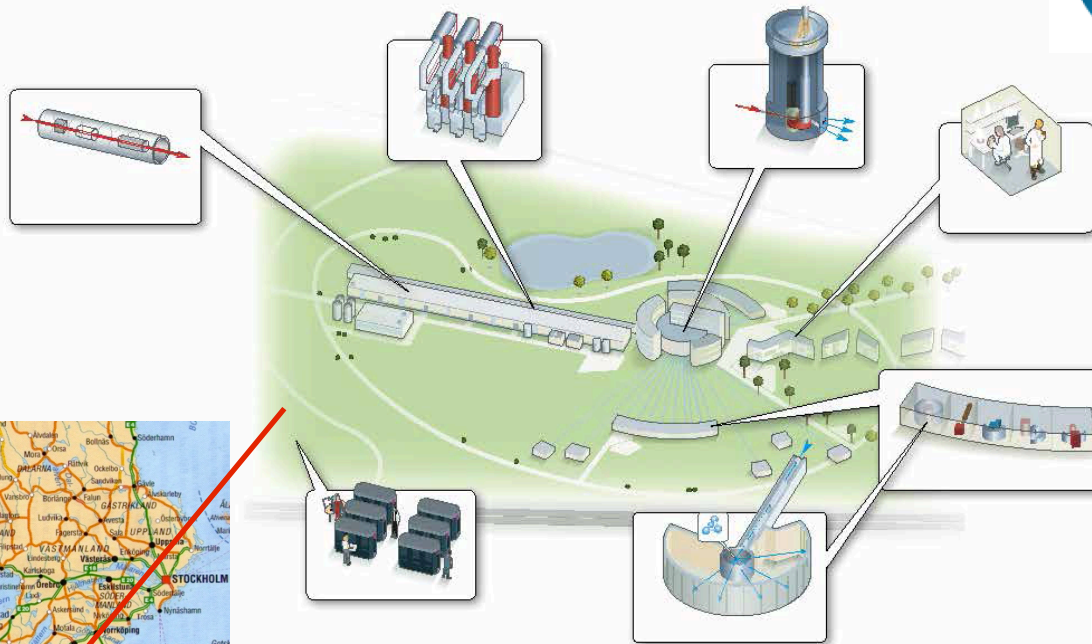
neutrinos of ~ 300 MeV necessary for this particular projet, C2F (CERN to Fréjus)

European Spallation Source (Lund)

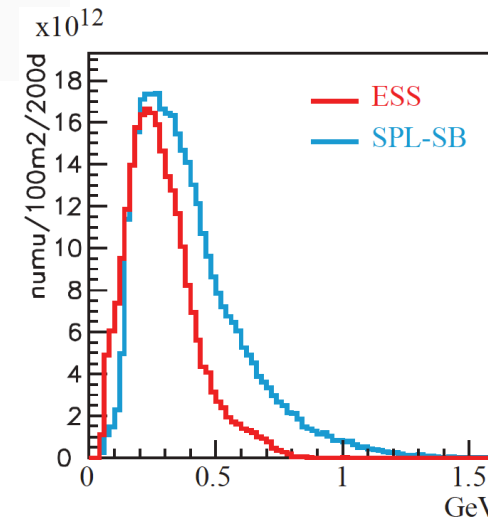


EUROPEAN
SPALLATION
SOURCE

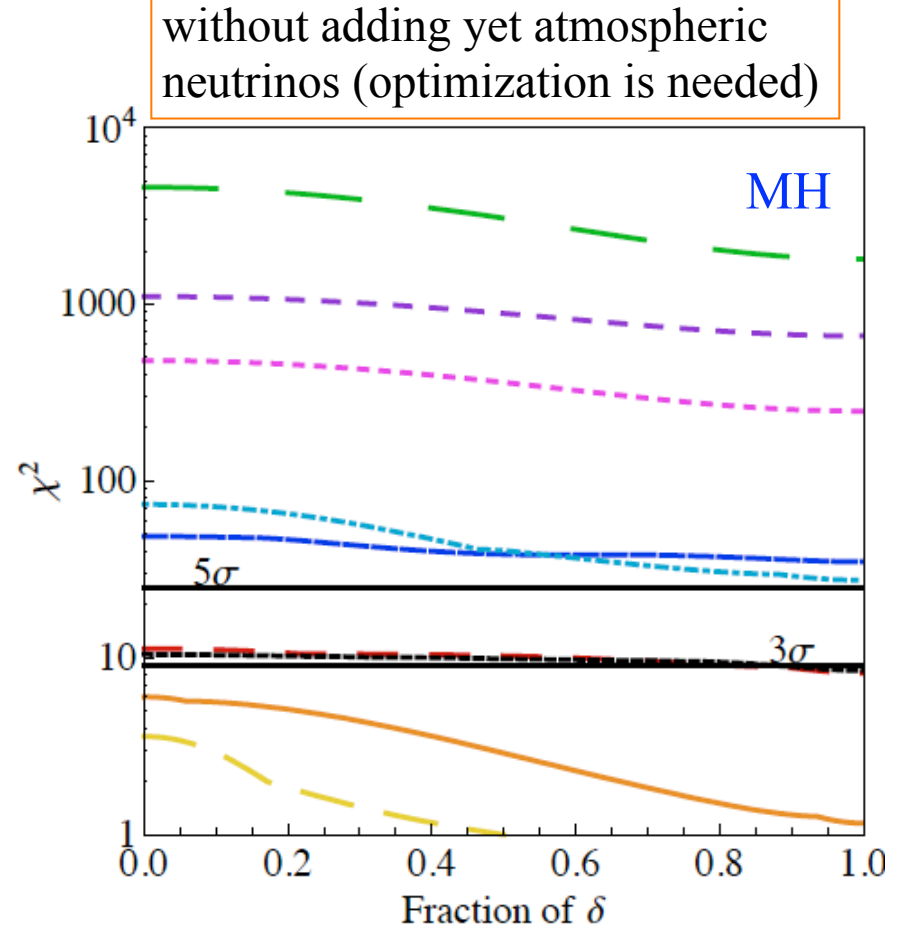
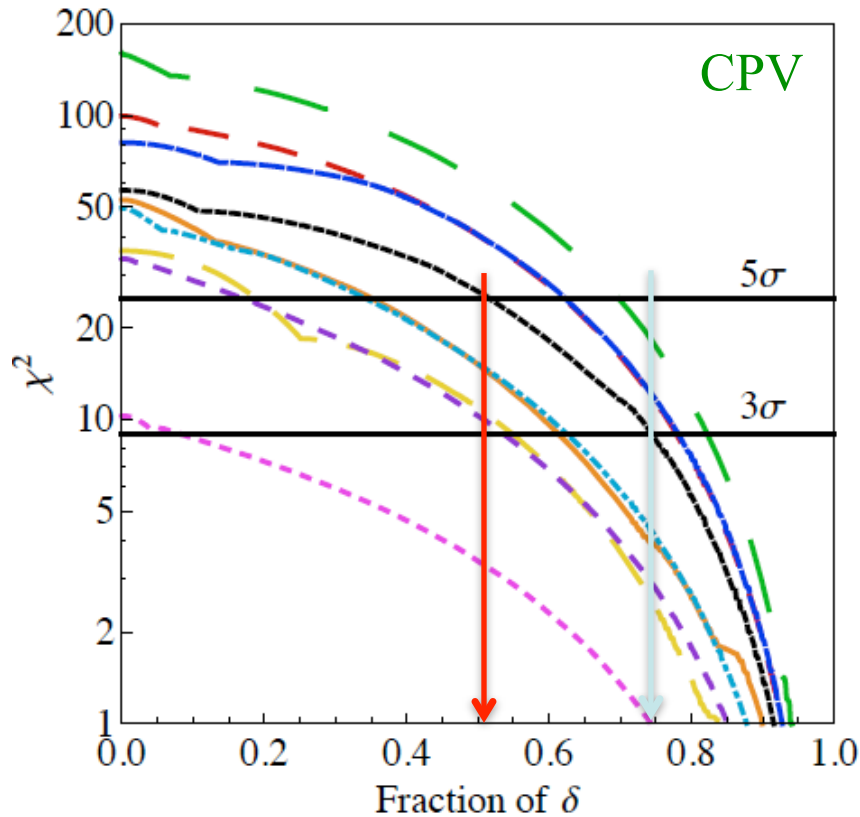
- First beam ~2020



- The ESS will be a copious source of spallation neutrons but also of **neutrinos**
- 5 MW proton beam
- 2.5 GeV protons
- ~300 MeV neutrinos



Physics Performance



EUROv parameters without any particular optimization for ESS

- 1 Mton WC detector (440 kton fiducial), 5% syst.
- 2.5 GeV protons
- 5 MW proton beam



Conclusions

- Deux stages sur Double Chooz
 - Etude du bruit de fond induit par les neutrons rapides.
 - Construction d'un télescope à muons cosmiques pour étudier le bruit de fond du détecteur proche.
- Un sujet de thèse
 - Etude des performances de futurs faisceaux neutrinos.

Dec. 1930: A desperate remedy

Letter of W. Pauli, 4 December 1930

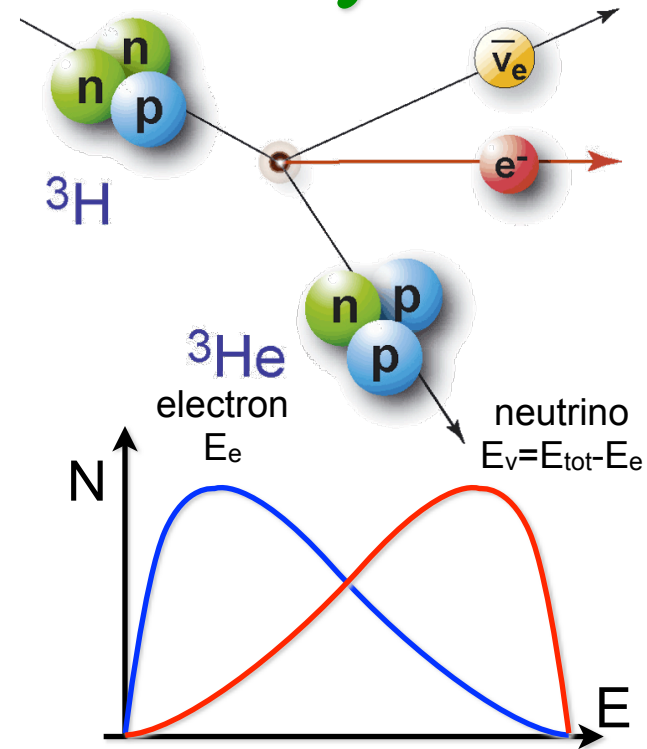
Dear Radioactive Ladies and Gentlemen,

As the bearer of these lines, to whom I graciously ask you to listen, will explain to you in more detail, how because of the "wrong" statistics of the N and Li^6 nuclei and the continuous beta spectrum, I have hit upon a **desperate remedy** to save the "exchange theorem" of statistics and the law of conservation of energy. Namely, the possibility that **there could exist in the nuclei electrically neutral particles**, that I wish to call **neutrons**, which have **spin 1/2** and obey the exclusion principle and which further differ from light quanta in that they do not travel with the velocity of light. The mass of the neutrons should be of the same order of magnitude as the electron mass and in any event not larger than 0.01 proton masses. The continuous beta spectrum would then become understandable by the assumption that in beta decay a neutron is emitted in addition to the electron such that the sum of the energies of the neutron and the electron is constant...

I agree that my remedy could seem incredible because one should have seen those neutrons very earlier if they really exist. But only the one who dare can win and the difficult situation, due to the continuous structure of the beta spectrum, is lighted by a remark of my honoured predecessor, Mr Debye, who told me recently in Bruxelles: "Oh, It's well better not to think to this at all, like new taxes". From now on, every solution to the issue must be discussed. Thus, dear radioactive people, look and judge. Unfortunately, I cannot appear in Tübingen personally since I am indispensable here in Zurich because of a ball on the night of 6/7 December. With my best regards to you, and also to Mr Back.

Your humble servant

W. Pauli



“I have done something very bad today by proposing a particle that cannot be detected. It is something no theorist should ever do.”

W. Pauli



How neutrinos propagate through the time?

According to Quantum Mechanics:

$$\left| \nu_j(t) \right\rangle = e^{-iHt/\hbar} \left| \nu_j(0) \right\rangle \quad \text{(H: Hamiltonian)}$$

Solutions of Schrödinger equation

For 3 neutrinos with a well defined mass and energy:

Schrödinger equation:

$$i \frac{d}{dt} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} = H \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

for the mass eigen states

$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = H_f \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

for the flavour eigen states

$$\text{with: } H_f = U H U^\dagger$$

following the ν history...

1933 : First estimation of the neutrino interaction cross-section (interaction probability) by **Hans Bethe** and **Rudolf Peierls**

$$\sigma_{\nu N} \approx 10^{-10} \sigma_{eN} \quad (N \text{ for nucleon), very very weak cross-section!!!}$$

$$\sigma_{\bar{\nu}p} \approx 10^{-43} \text{ cm}^2$$

$$\lambda = \frac{1}{N_A \rho \sigma}$$

mean free path

mean free path in lead
for **3 MeV** neutrino

$$\lambda(\text{Pb}) \approx \frac{1}{6 \times 10^{23} (\text{nucleons / g})(7.9 \text{ g / cm}^3) \times 10^{-43} \text{ cm}^2}$$

~ 4 light-years!



to stop a neutrino a lot of lead is needed
or many neutrinos...

The beginning of a long neutrino hunting which lasted 26 years...

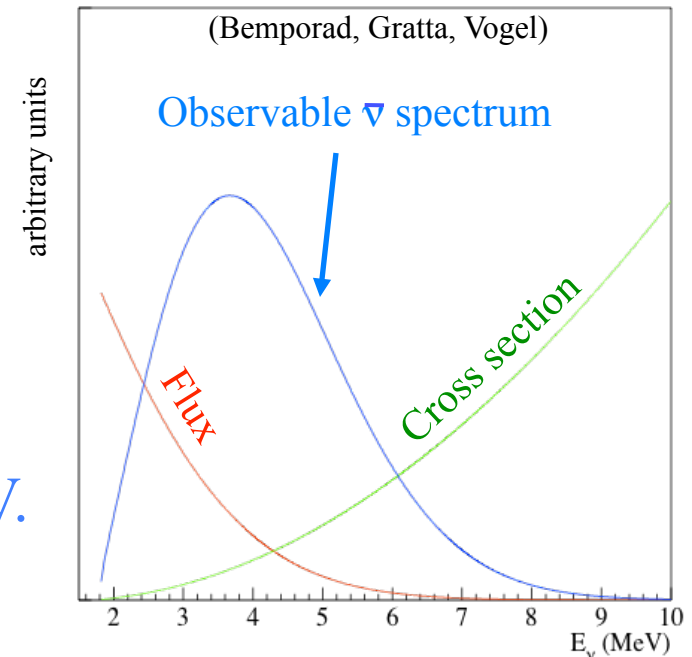
(Pauli: "I bet a case of champagne that nobody would ever detect the neutrino")

Nuclear Reactors as neutrino source

- Nuclear reactors are a very intense electron anti-neutrino source (β decay of neutron rich fission fragments).
- Each fission release an energy of ~ 200 MeV and generates ~ 6 electron anti-neutrinos. For a typical commercial reactor (3 GW thermal energy):

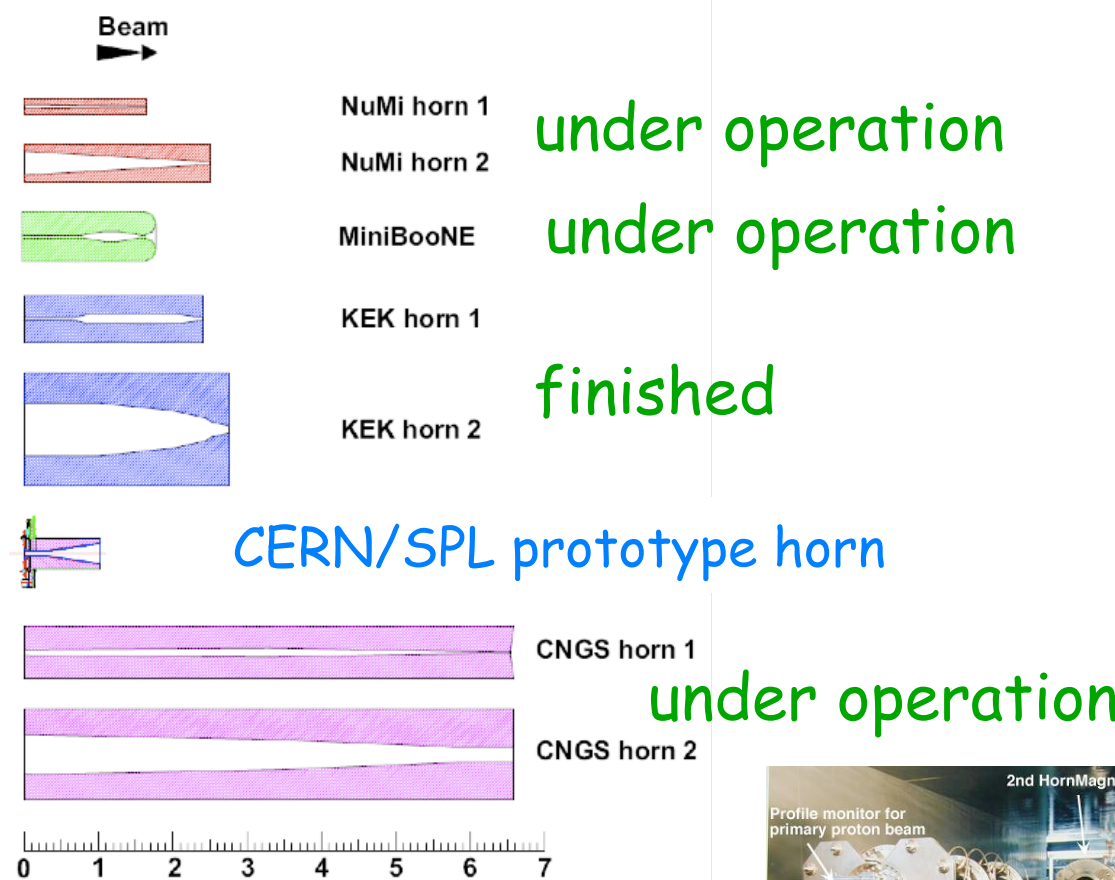
$$3 \text{ GW} \approx 2 \times 10^{21} \text{ MeV/s} \rightarrow 6 \times 10^{20} \bar{\nu}_e/\text{s}$$

- Observable neutrino energy spectrum = **neutrino flux** * **cross section**.
- The spectrum has a maximum at ~ 3.7 MeV.



Present Collectors

Experiment	Current	Rep. Rate	Pulses per time period
<i>Numi</i> (120 GeV)	200 kA	0.5 Hz	6 Mpulses 1 year
<i>MiniBoone</i> (8 GeV)	170 kA	5 Hz	11 Mpulses 1 year
<i>K2K</i> (12 GeV)	250 kA	0.5 Hz	11 Mpulses 1 year
<i>Super-Beam</i> (3.5 GeV)	300 kA	50 Hz	200 Mpulses 6 weeks
<i>CNGS</i> (400 GeV)	150 kA	2 pulses/ 6 sec	42 Mpulses 4 year

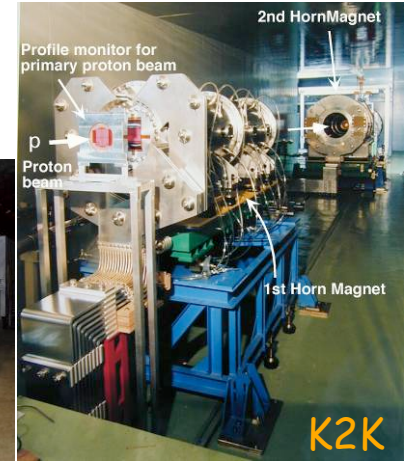
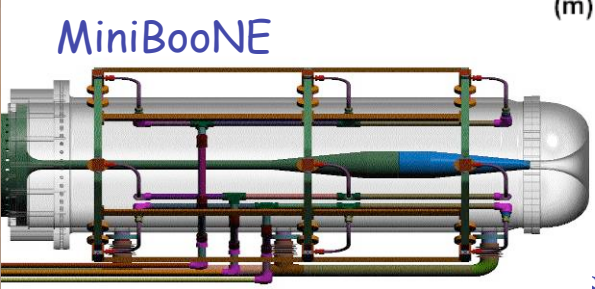
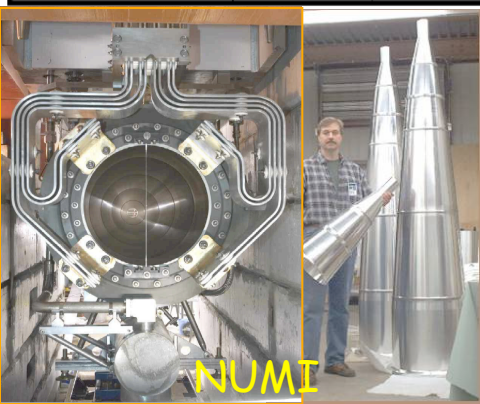


under operation

under operation

finished

under operation

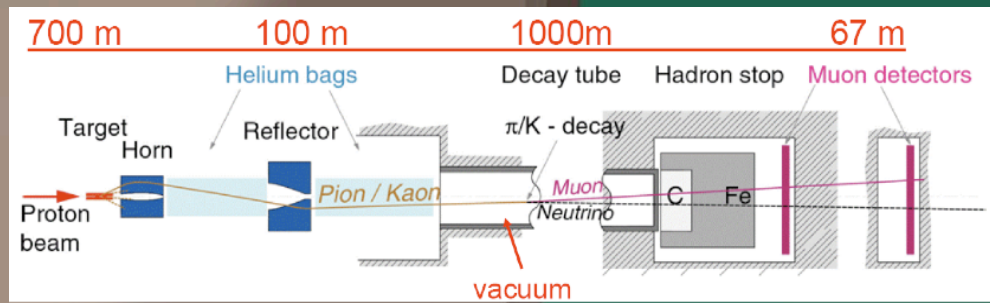
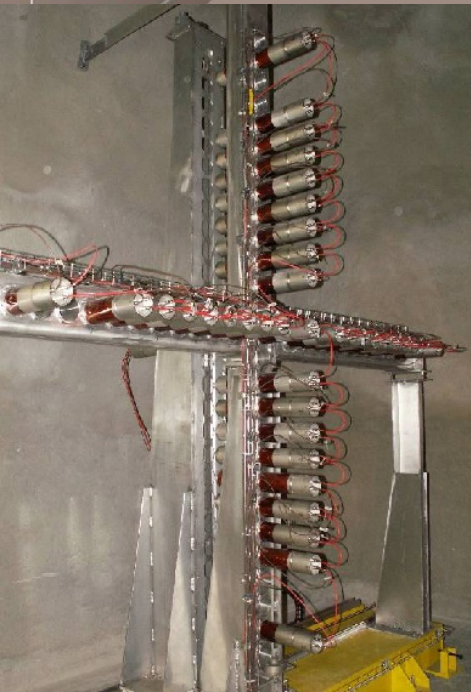
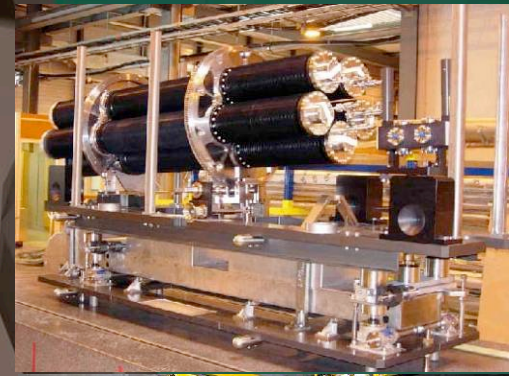
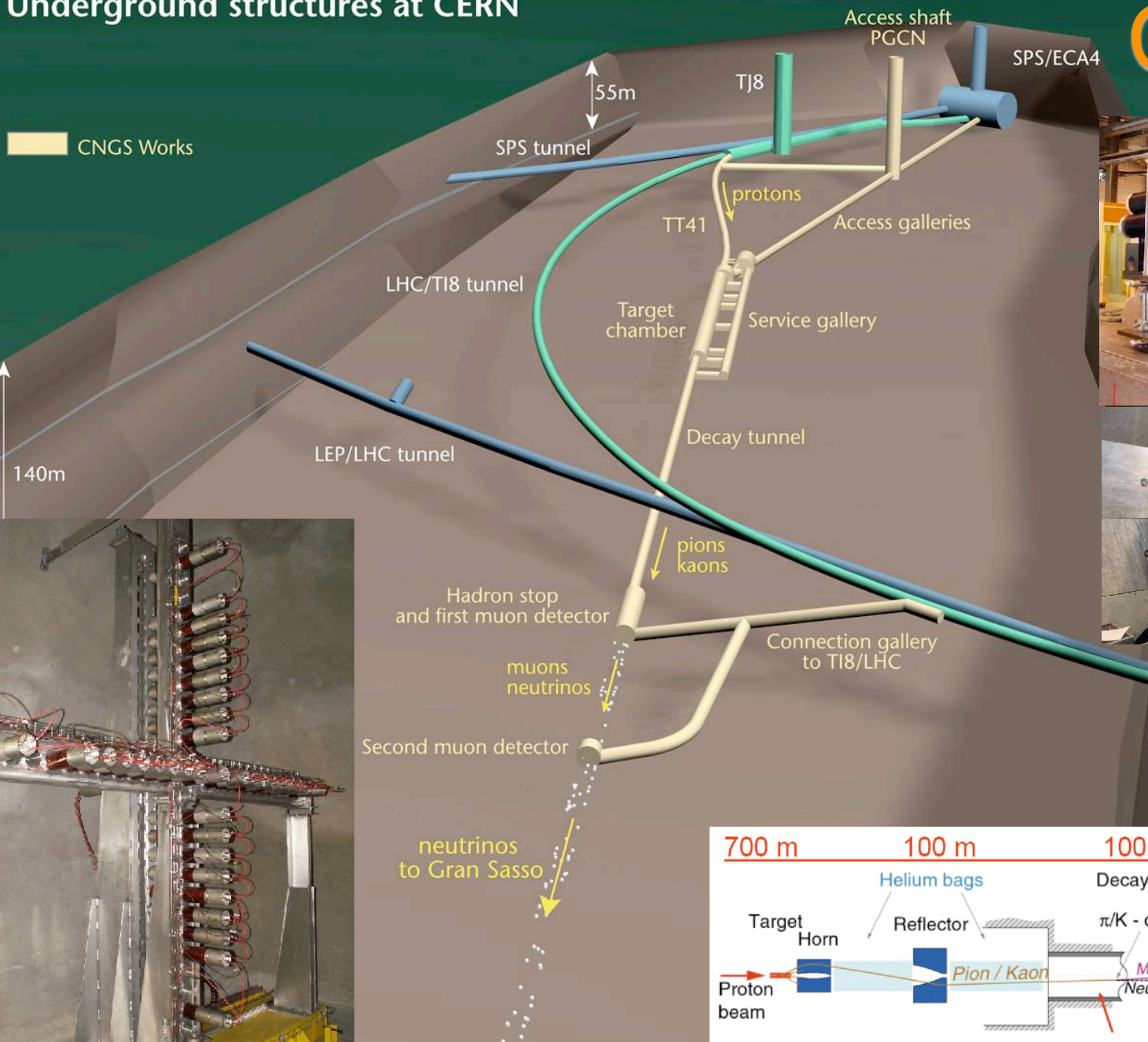


CERN NEUTRINOS TO GRAN SASSO

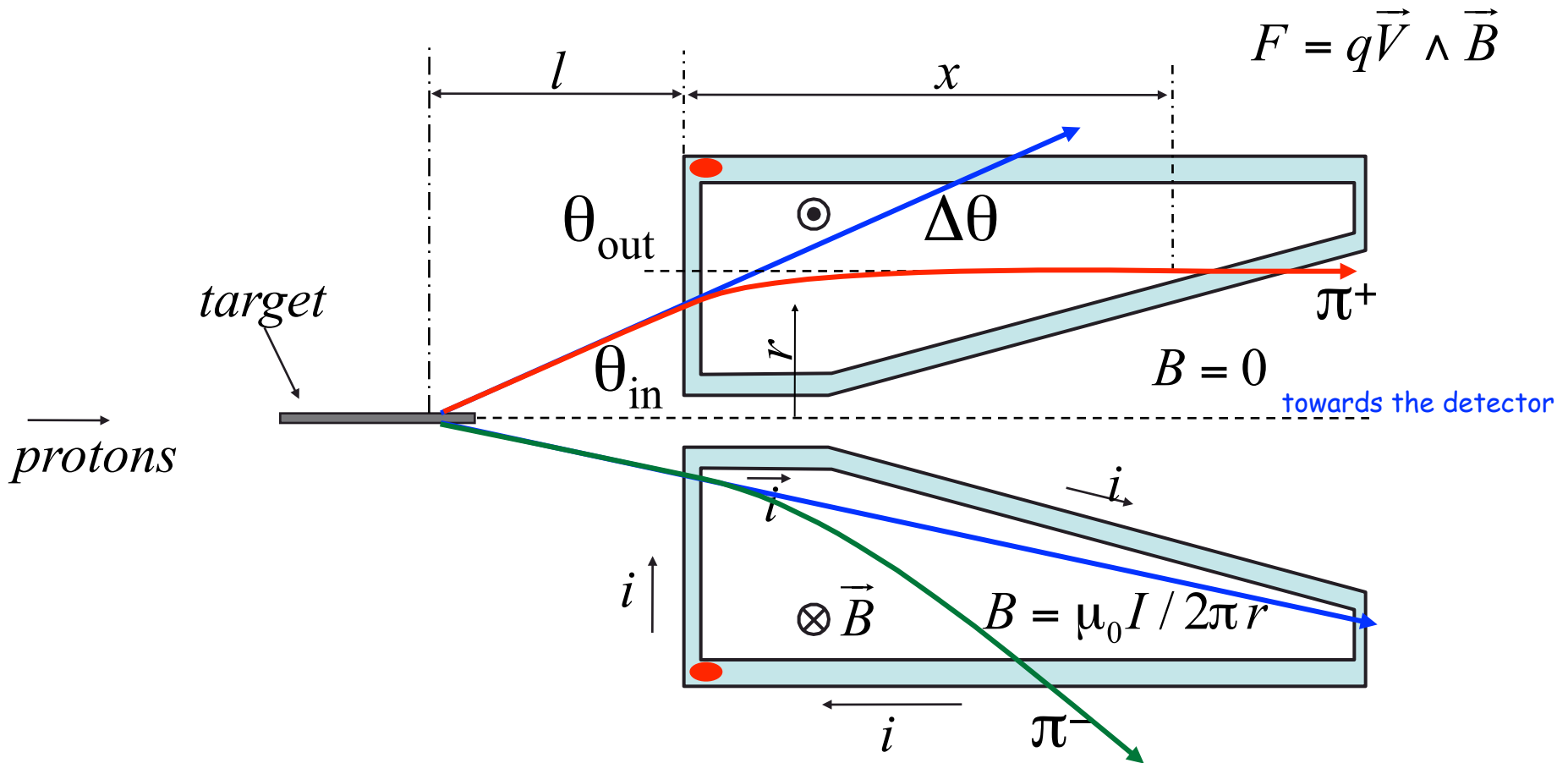
Underground structures at CERN

CNGS

■ CNGS Works

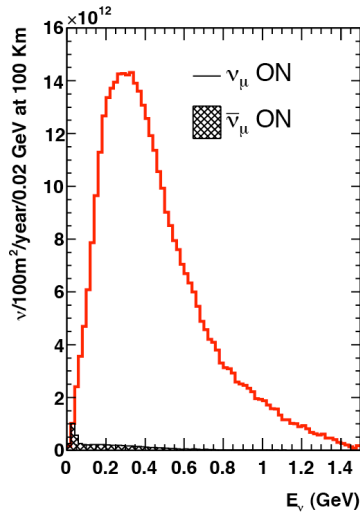
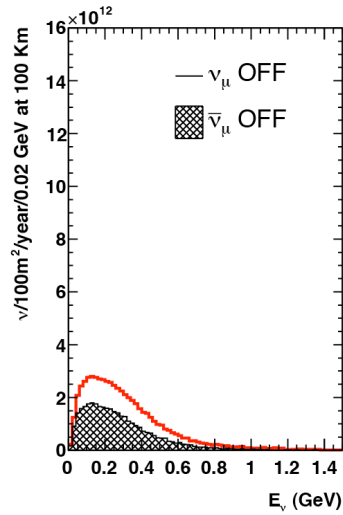


Hadronic Collector

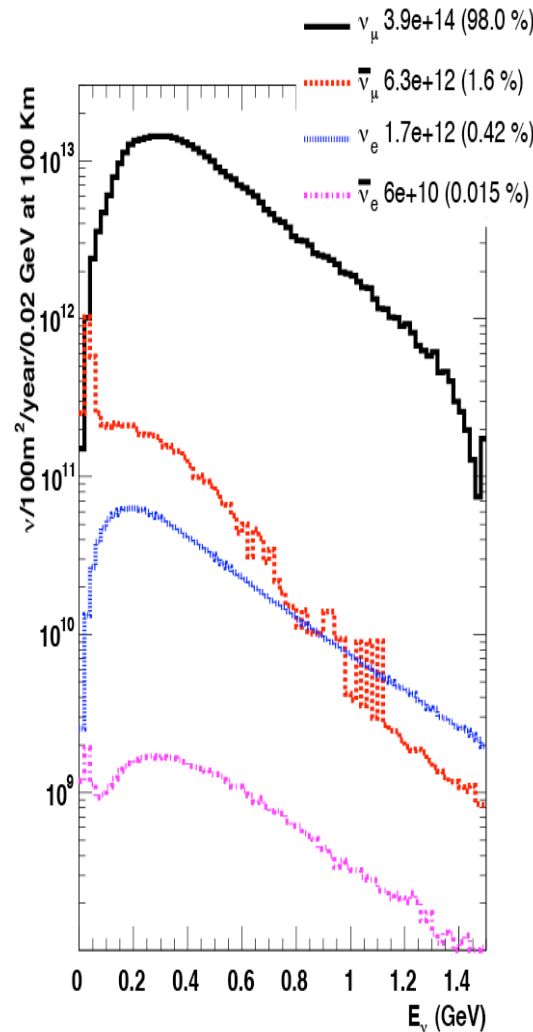


 horn shape (conical)

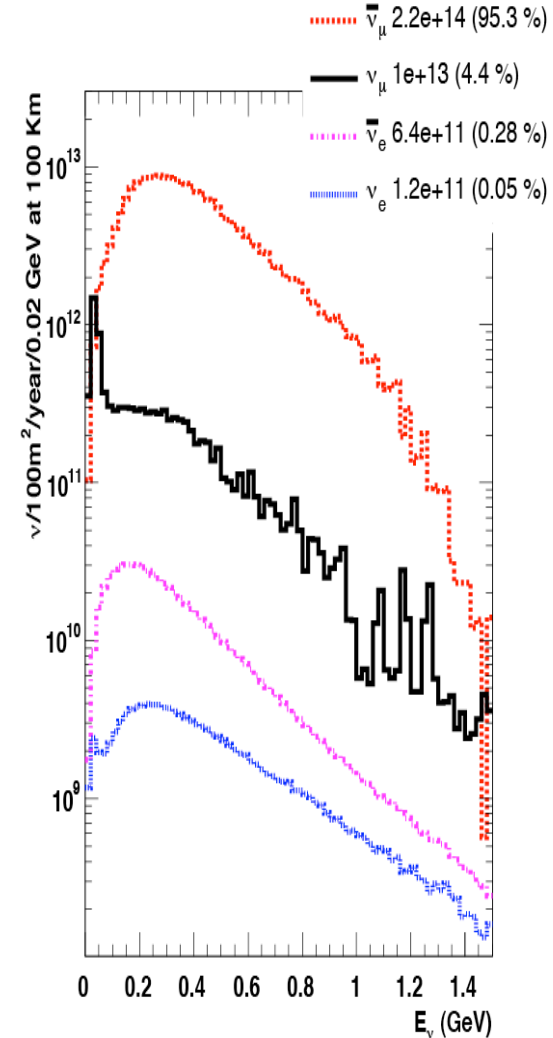
Neutrino Spectra



horn on/off



neutrinos



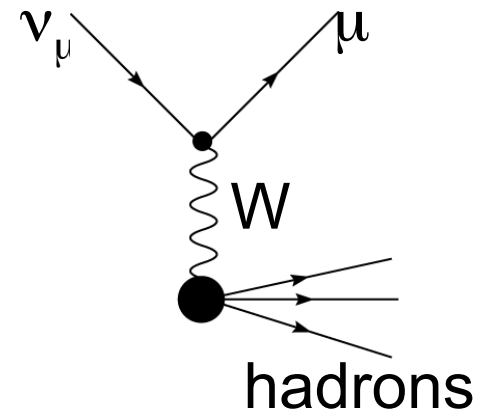
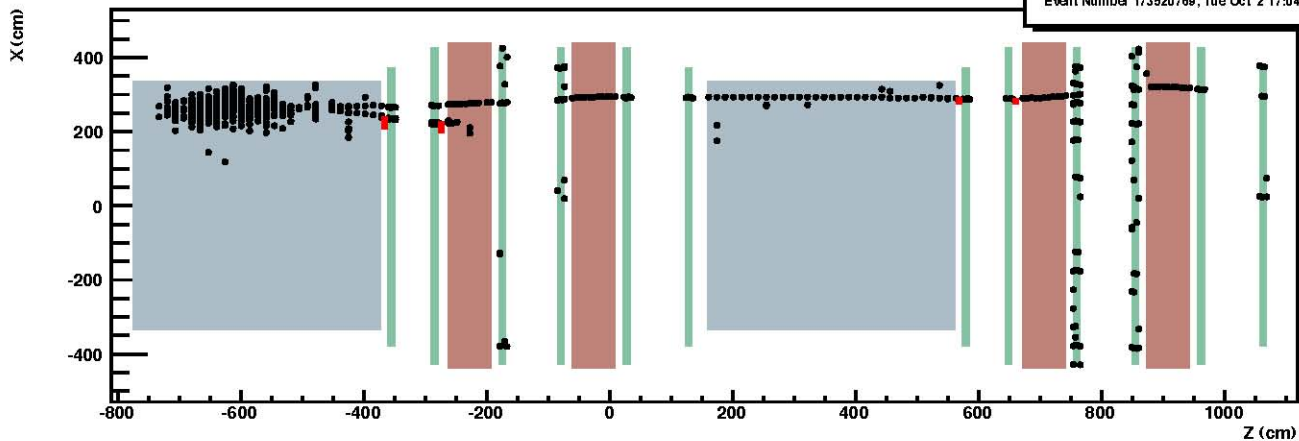
anti-neutrinos

First OPERA events

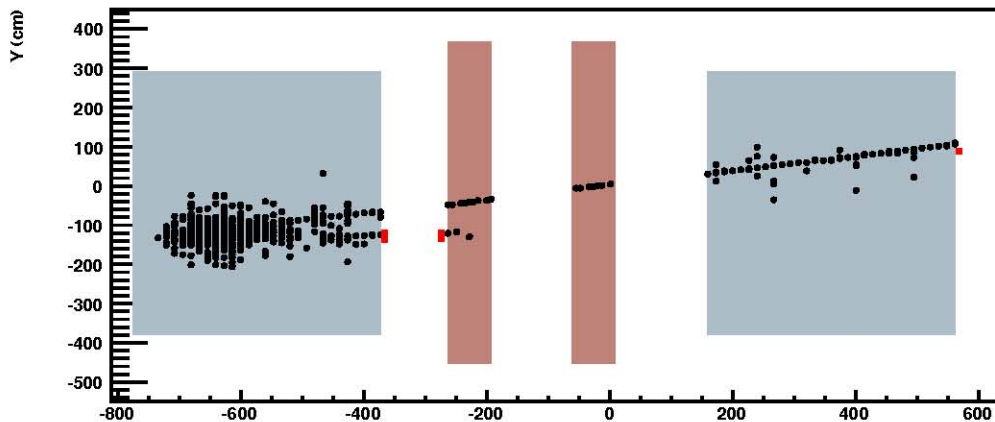


2007

TOP VIEW (horizontal projection)

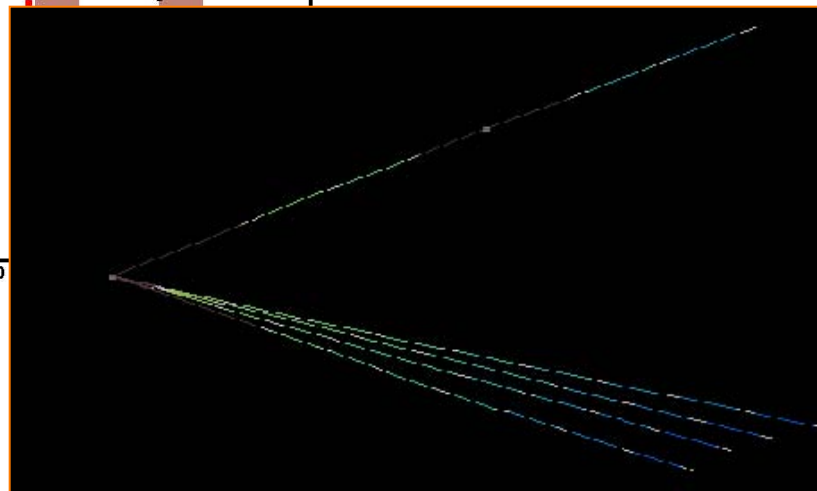


SIDE VIEW (Vertical projection)



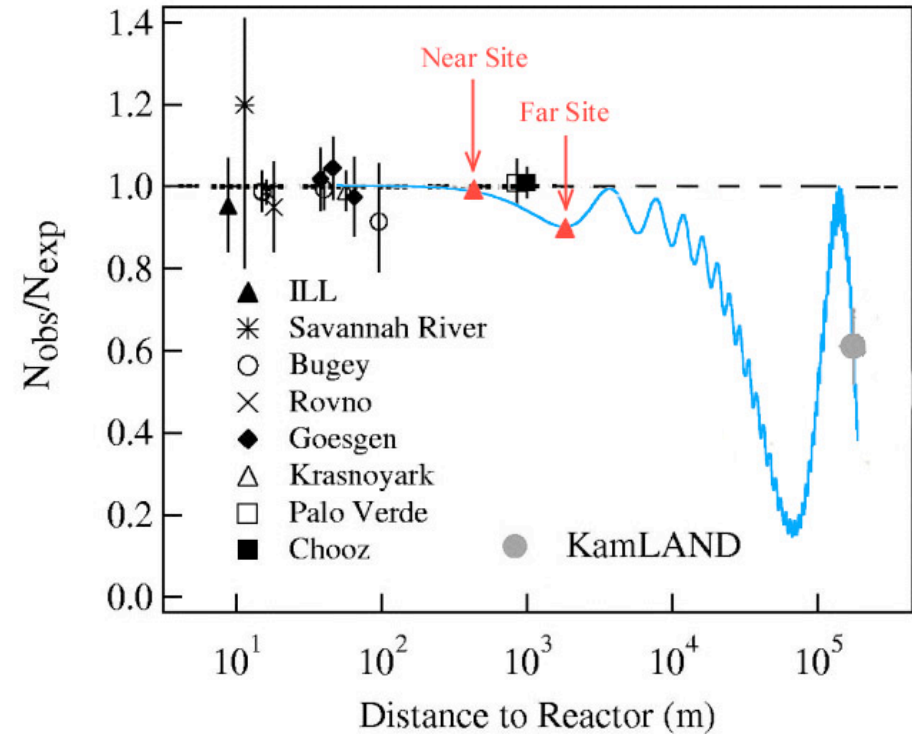
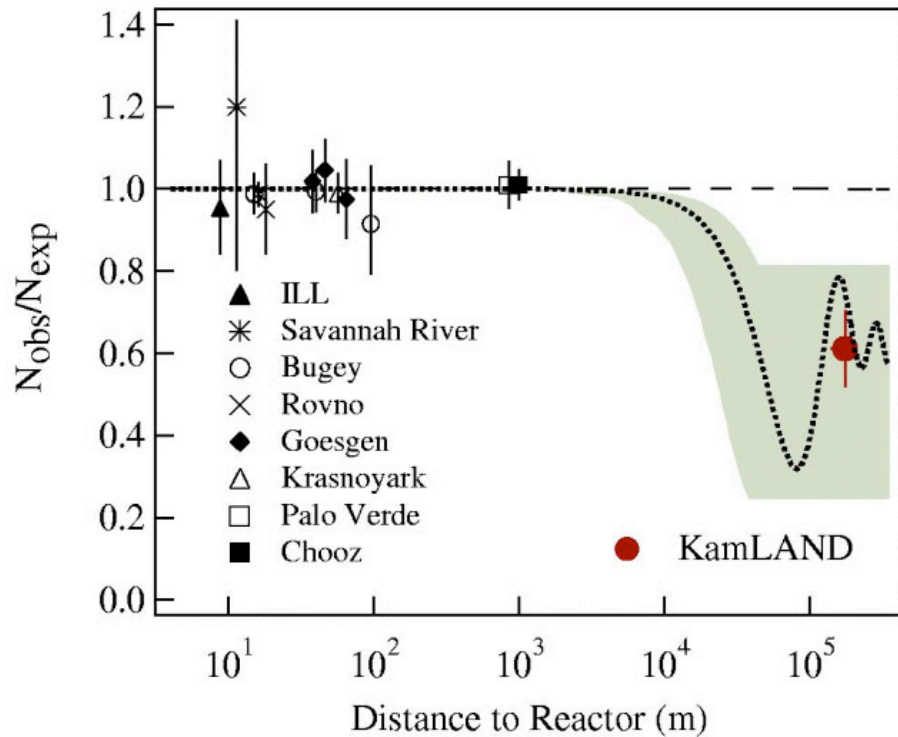
hadrons
charged current

in the emulsions



The θ_{13} hunting

disappearance of electron neutrinos

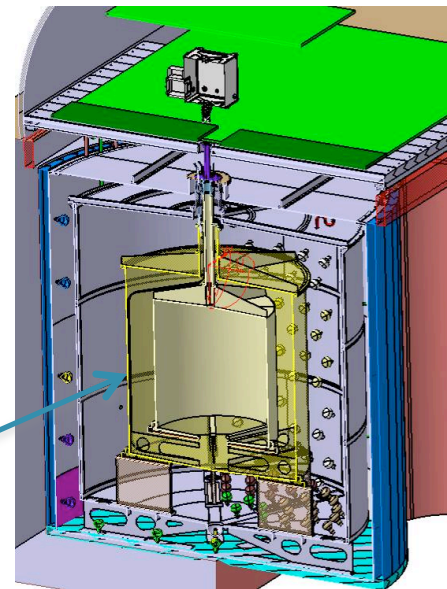
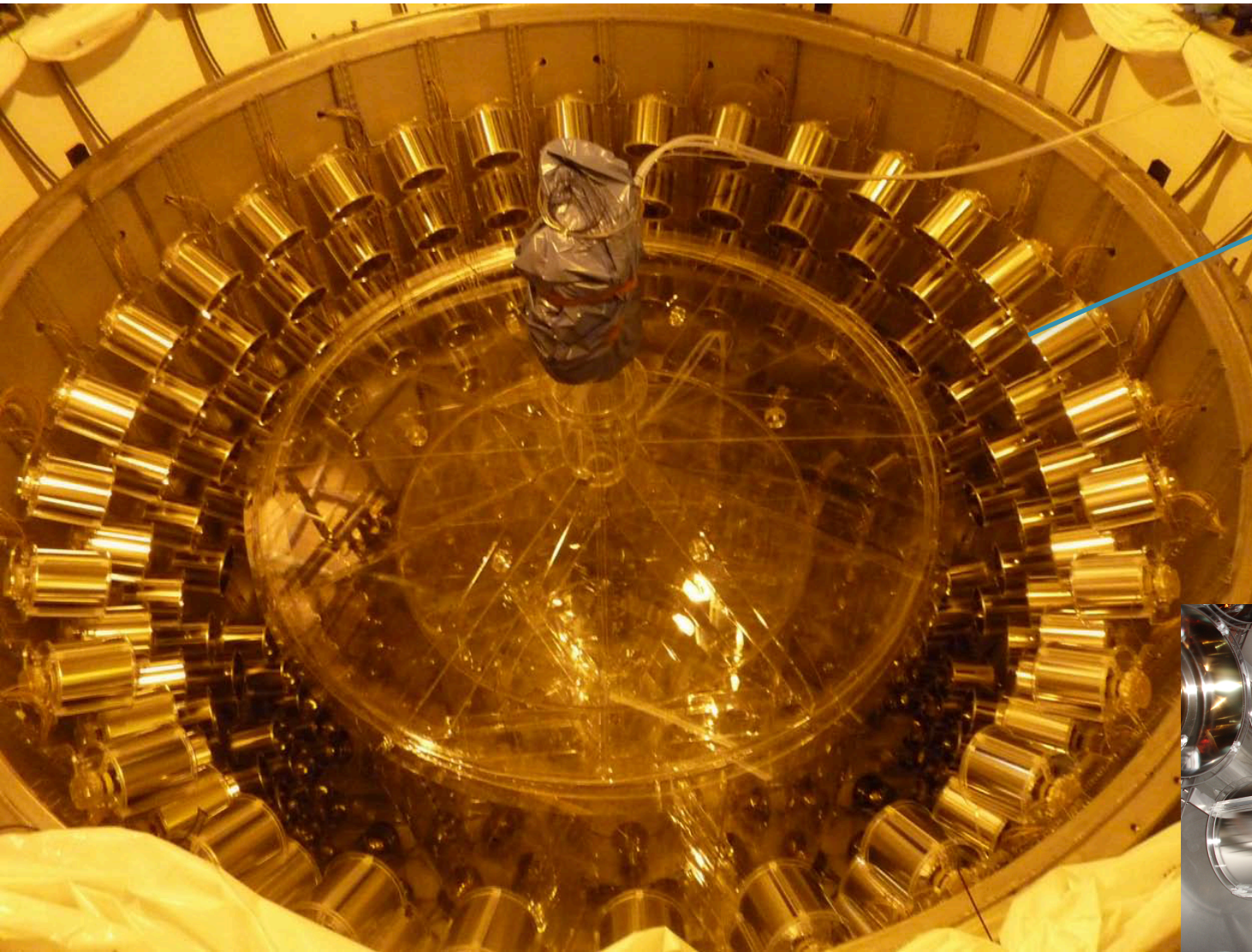


Actually, we have almost neglected θ_{13} on this figure

For $\theta_{13} \sim 10^\circ$

Far DC Detector construction

Acrylic Gamma Catcher
installation (May 2010)



Project Comparison

Daya Bay
(China)



Double Chooz
(France)



RENO
(South Korea)

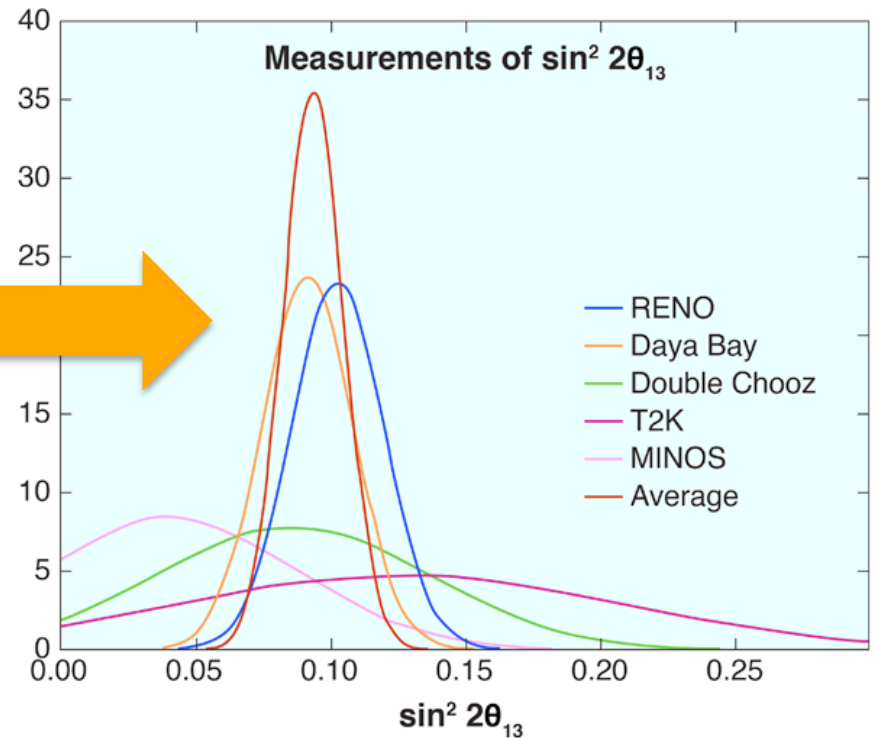
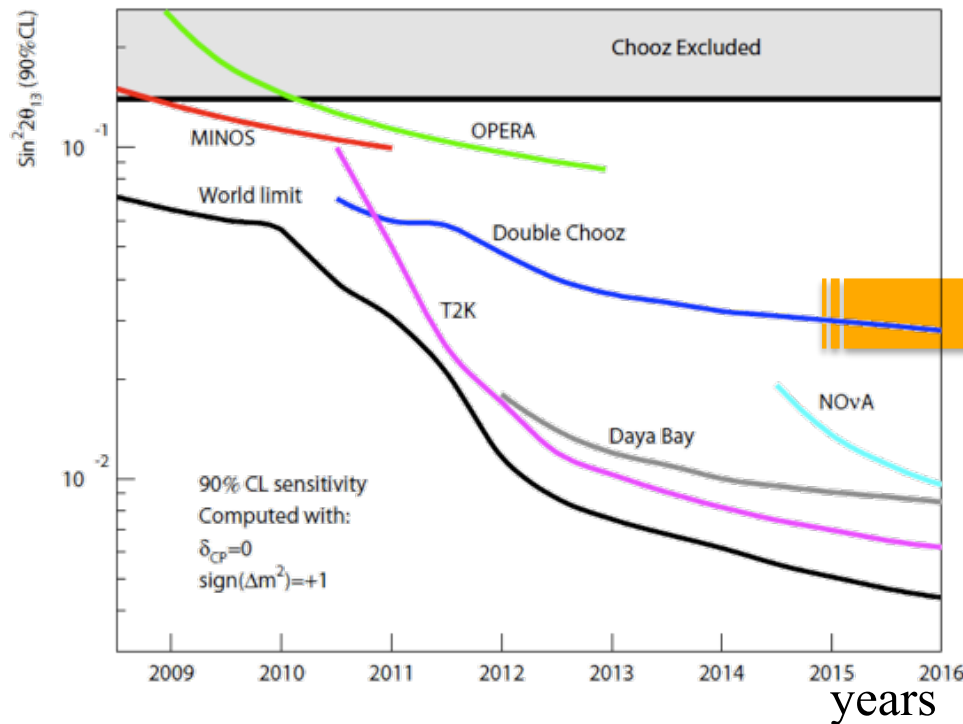


	Luminosity in 3 years (ton·GW·y)	Overburden near/far (mwe)	Expected sensitivity	Start of data taking
Daya Bay	4200	270/950	<0.01	August 2011
Double Chooz	210	80/300	0.02~0.03	April 2011
RENO	740	90/440	~0.02	August 2011

Results on θ_{13}

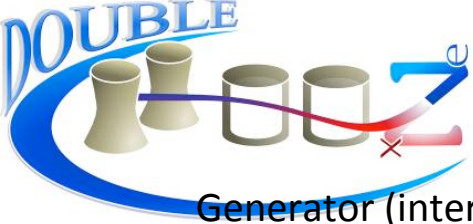
expectations

measurements



$$\sin^2 2\theta_{13} = 0.089 \pm 0.010(\text{stat}) \pm 0.005(\text{syst})$$

- θ_{13} is large!
- Proposed facilities to be readjusted...

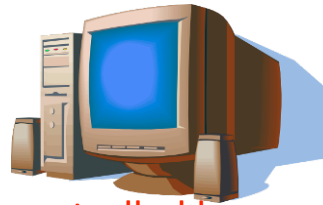
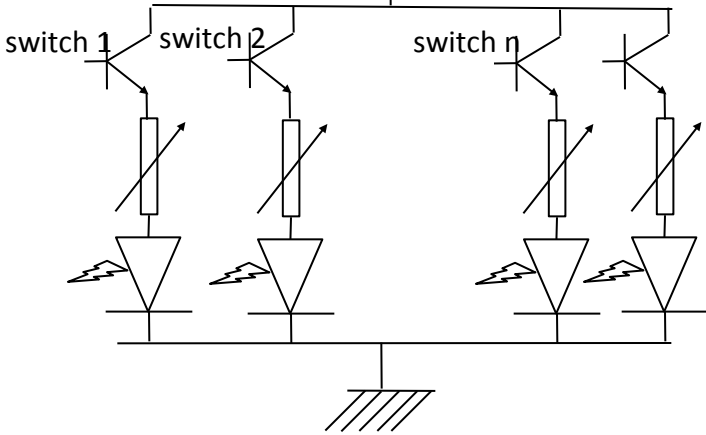


IV Calibration system

Generator (internal or external trigger, delay, remote control)

send to the DAQ a 96 bit word to give the flashed LED pattern

quartz fibre (going to the detector)



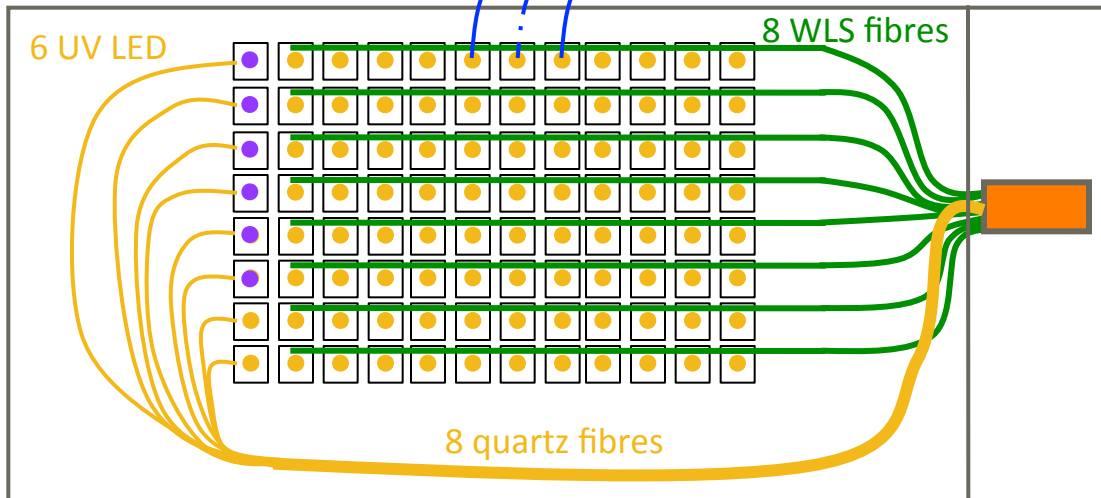
controlled by LabView (GPIB or ethernet)

WLS fibre



diameter: 600 μm
(tot. diam.: 1040 μm)

to the detector



LED (475 nm or 365 nm)

monitoring photodetector (signal to be sent in a "normal channel, FADC, if possible use the same HT than the other IV PMT's)

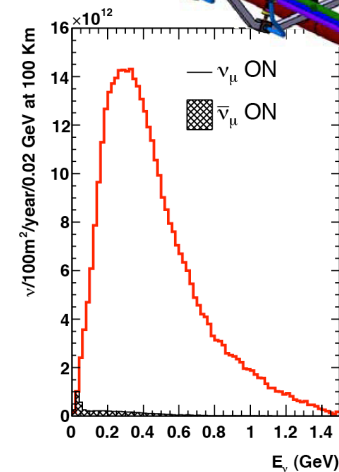
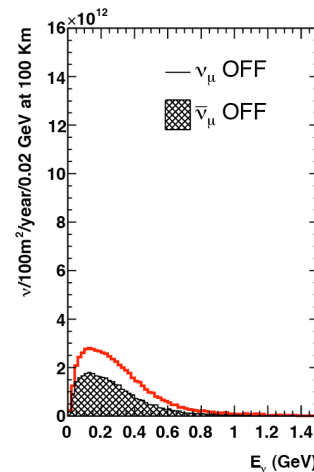
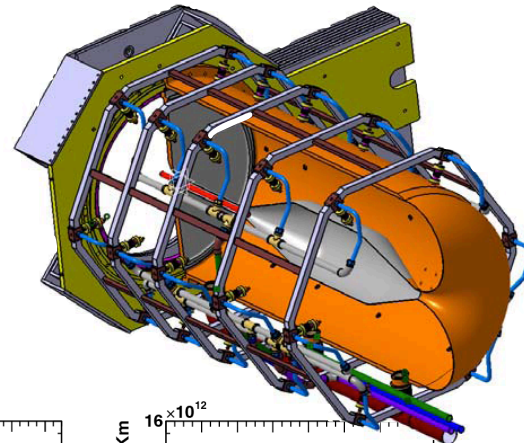
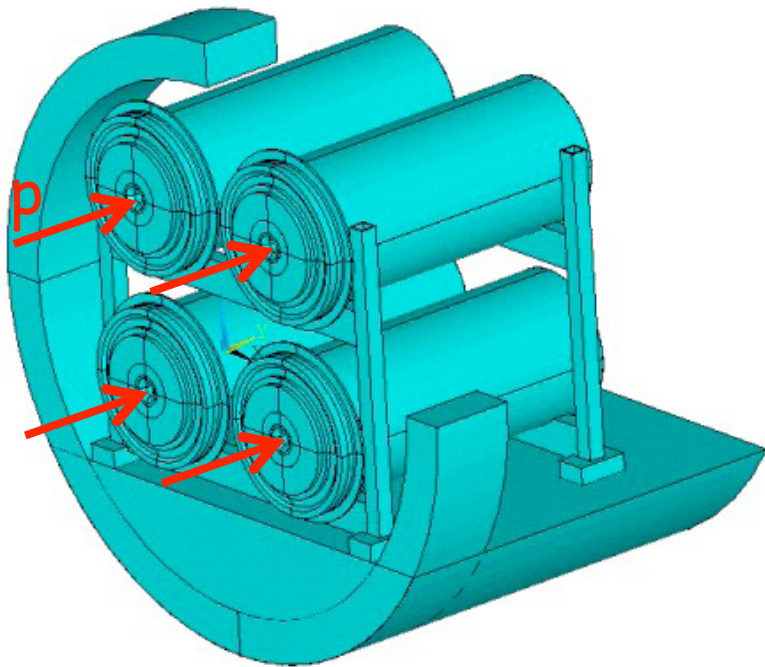
How to mitigate the power effect

4 target/horn system
(4x4 m²) with single
decay tunnel (~30 m)



more expensive but more reliable system

back to solid targets able to
afford up to ~1.5 MW proton
beam (get rid of liquid Hg)



Possible Detector Locations



- Many mines (active or not) are available in Sweden
- What is the optimal position for CPV
- How this project could help for MH?

$$\Delta m_{\text{sun}}^2 = 7.6 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{\text{atm}}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

$$\theta_{23} = 44.4^\circ$$

$$\theta_{12} = 34.4^\circ$$

$$\theta_{13} = 9.3^\circ$$

$$\delta_{\text{CP}} = 0$$

$$E \sim 290 \text{ MeV}$$

