

# Fixed-Target Dimuon Experiments at Fermilab with Proton Beams

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ECT\* Workshop on “Physics at A Fixed Target ExpeRiment  
(AFTER) using the LHC beams”

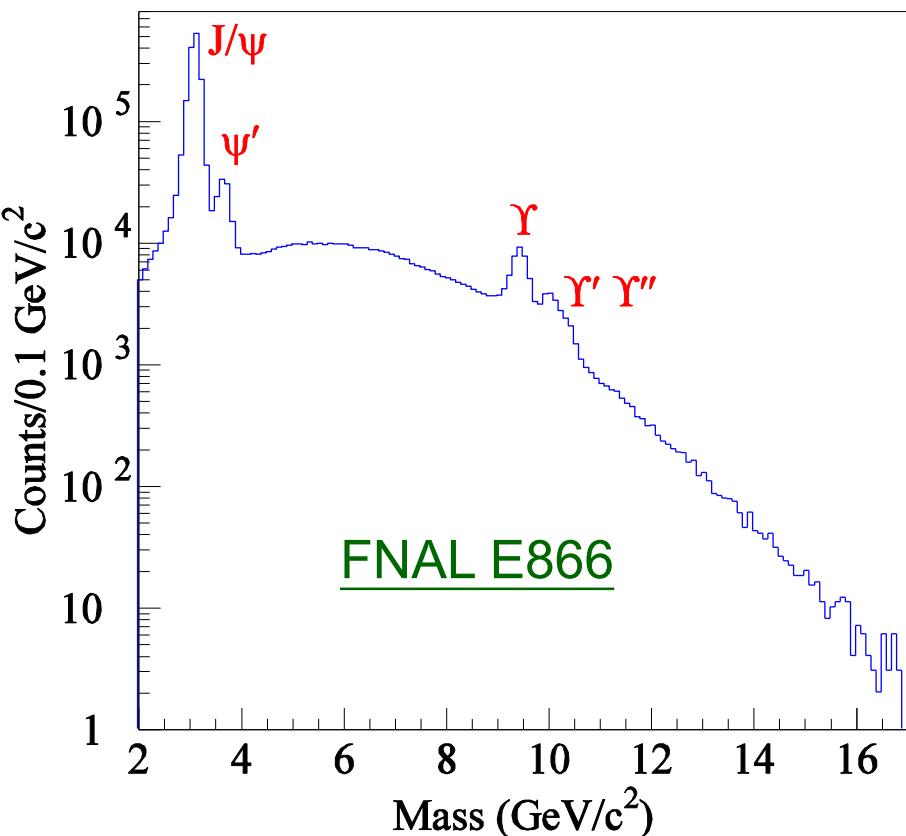
Trento, Feb. 3-13, 2013

## Outline

- Highlights from the Fermilab dimuon experiments
- Physics opportunities at AFTER with dilepton experiments

# Dilepton production in hadron-hadron collision

$p + p(d) \rightarrow \mu^+ \mu^- x$  at 800 GeV/c



Two components in the dilepton mass spectrum

(a) Continuum: Drell-Yan process

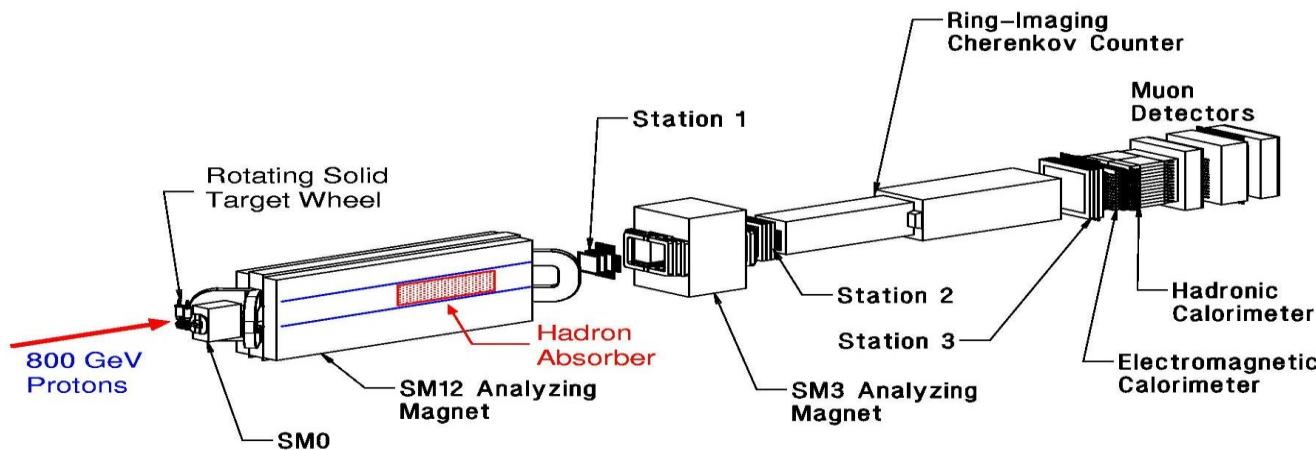
- Electromagnetic process
- Quark - antiquark annihilation

(b) Vector mesons: J/ $\psi$ ,  $\Upsilon$

- Strong interaction
- Gluon - gluon fusion  
(quark - antiquark annihilation)

# Fermilab Dimuon Spectrometer

(E605 / 772 / 789 / 866 / 906)



1) Fermilab E772 (proposed in 1986 and completed in 1988)

"Nuclear Dependence of Drell-Yan and Quarkonium Production"

2) Fermilab E789 (proposed in 1989 and completed in 1991)

"Search for Two-Body Decays of Heavy Quark Mesons"

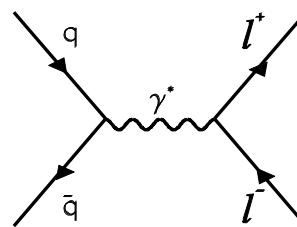
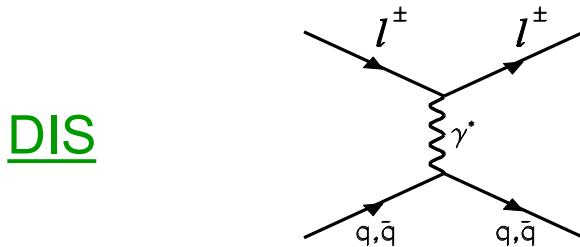
3) Fermilab E866 (proposed in 1993 and completed in 1996)

"Determination of  $\bar{d} / \bar{u}$  Ratio of the Proton via Drell-Yan"

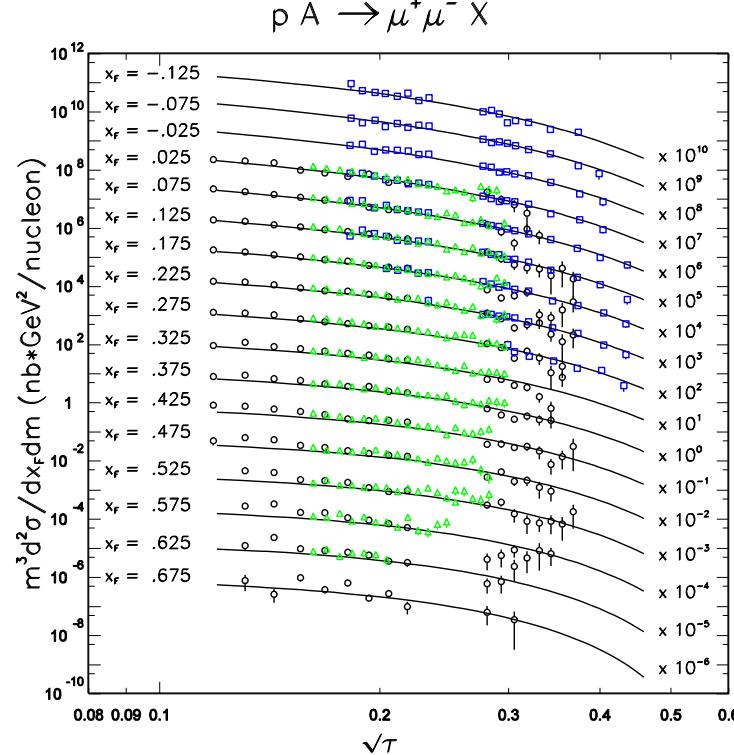
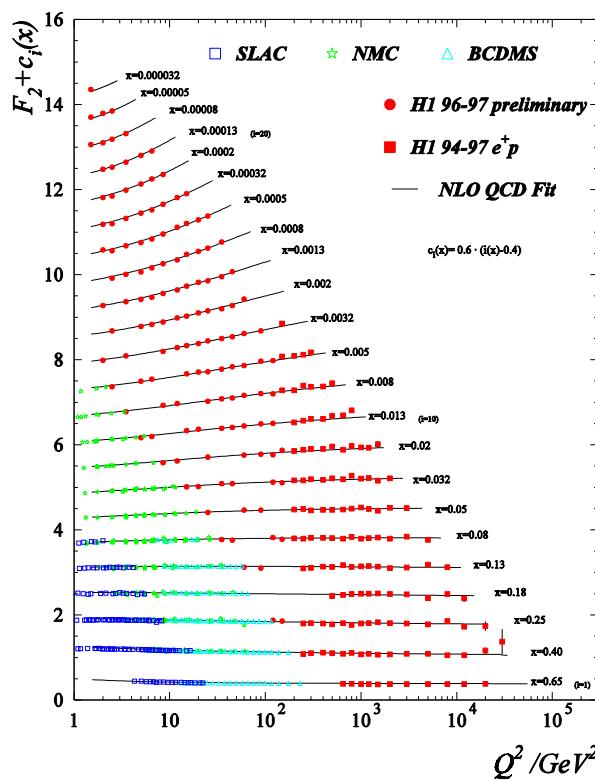
4) Fermilab E906 (proposed in 1999, run started in 2012)

"Drell-Yan with the FNAL Main Injector"

# Complementarity between DIS and Drell-Yan



Drell-Yan

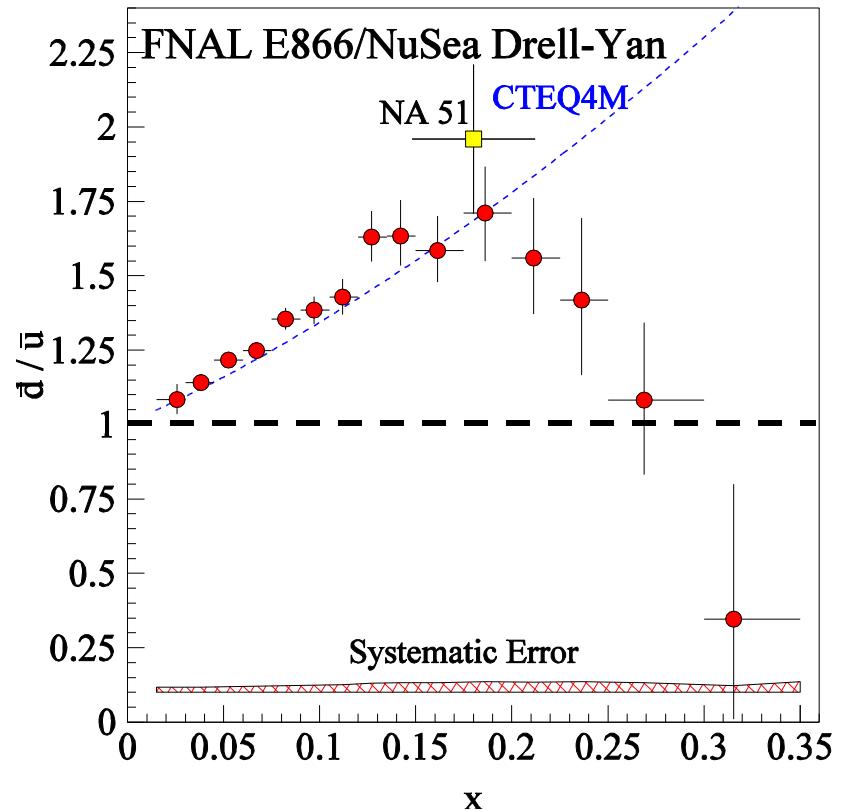
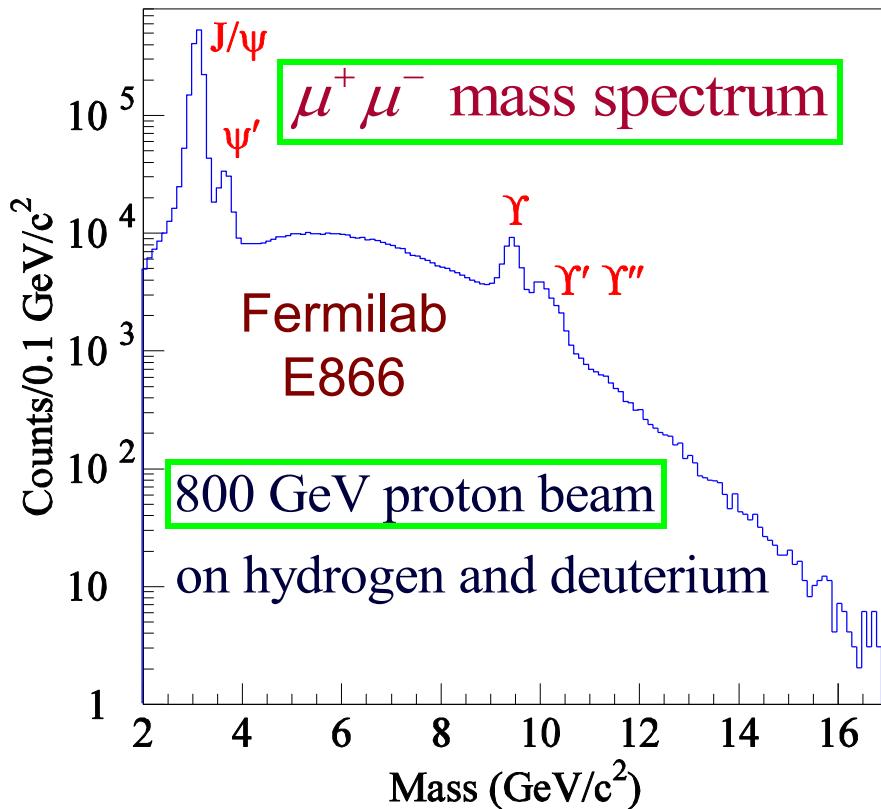


McGaughey,  
Moss, JCP,  
Ann.Rev.Nucl.  
Part. Sci. 49  
(1999) 217

Both DIS and Drell-Yan process are tools to probe the quark and antiquark structure in hadrons (factorization, universality)

# $\bar{d}/\bar{u}$ flavor asymmetry from Drell-Yan

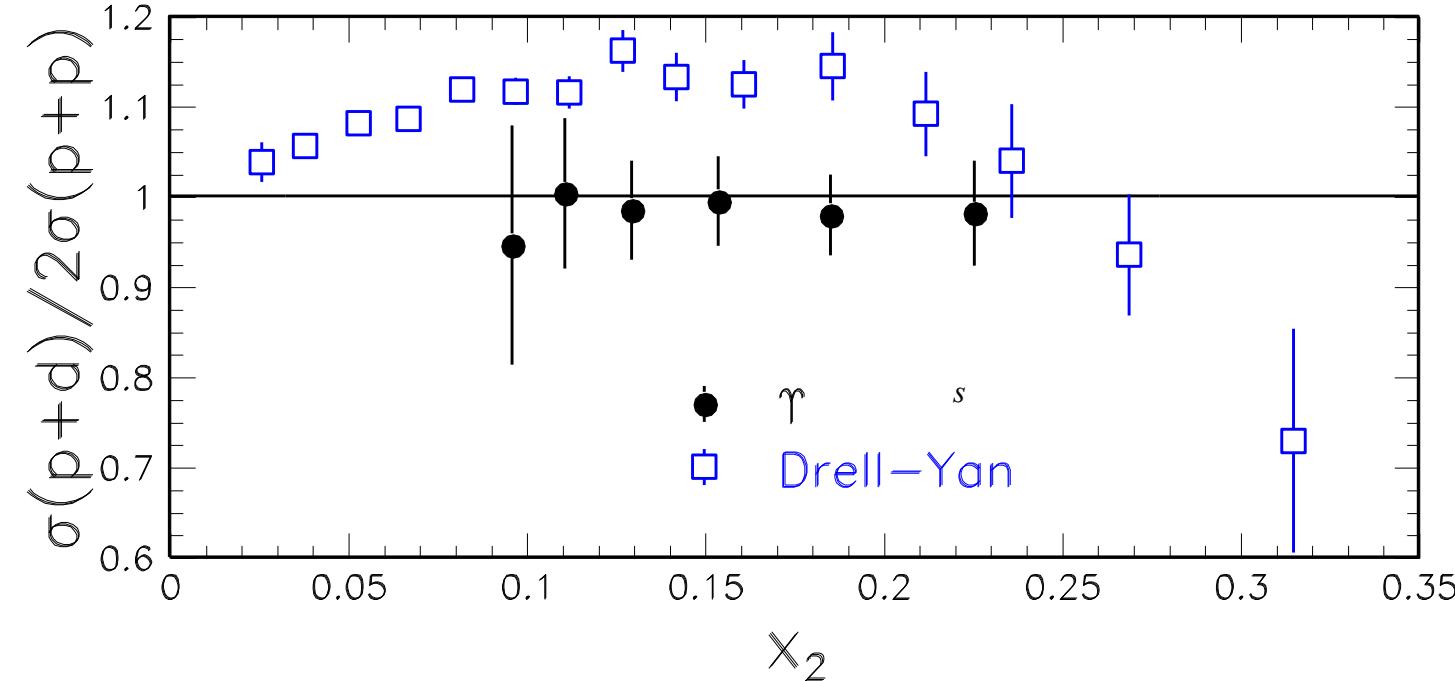
$$\left( \frac{d^2\sigma}{dx_1 dx_2} \right)_{D.Y.} = \frac{4\pi\alpha^2}{9sx_1x_2} \sum_a e_a^2 [q_a(x_1)\bar{q}_a(x_2) + \bar{q}_a(x_1)q_a(x_2)]$$



at  $x_1 > x_2$  : Drell-Yan:  $\sigma^{pd} / 2\sigma^{pp} \sim \frac{1}{2}(1 + \bar{d}(x_2)/\bar{u}(x_2))$

# Gluon distributions in proton versus neutron?

E866 data:  $\sigma(p+d \rightarrow \Upsilon X) / 2\sigma(p+p \rightarrow \Upsilon X)$



Lingyan Zhu et al.,  
PRL, 100 (2008)  
062301 (arXiv:  
0710.2344)

$$Drell-Yan: \sigma^{pd} / 2\sigma^{pp} \simeq [1 + \bar{d}(x) / \bar{u}(x)] / 2$$

$$J/\Psi, \Upsilon: \sigma^{pd} / 2\sigma^{pp} \simeq [1 + g_n(x) / g_p(x)] / 2$$

If gluon distributions in proton and neutron are different, then charge-symmetry is violated at the partonic level

# Charge Symmetry Violation in PDF?

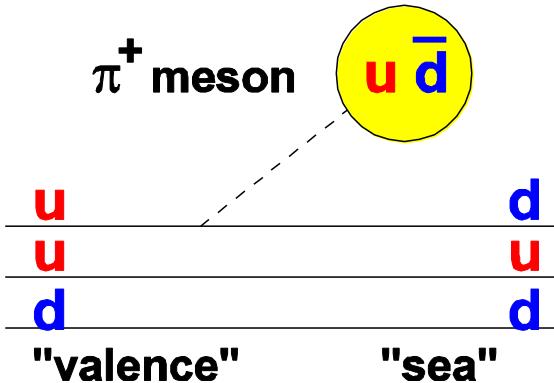
- Charge symmetry : rotation around isospin axis  $I_y$  by  $180^\circ$   
 $p \leftrightarrow n$     $u \leftrightarrow d$     $\bar{u} \leftrightarrow \bar{d}$    ( $u_p = d_n$ ;  $\bar{u}_p = \bar{d}_n$  etc.)
- Charge symmetry in PDF is generally assumed, but not tested  
A comparison of  $F_2^u$  with  $F_2^d$  shows upper limit of CS breaking of  $\square 5 - 10\%$  (Londergan and Thomas, hep-ph/0407247)
- The NuTeV anomaly in the Weinberg angle could be partly explained by charge symmetry violation (Londergan and Thomas, PL B558 (2003) 132)

See recent review of Charge Symmetry violation:

Londergan, Peng, and Thomas (Rev. Mod. Phy. 82 (2010) 2009)

# Origins of $\bar{u}(x) \neq \bar{d}(x)$ ?

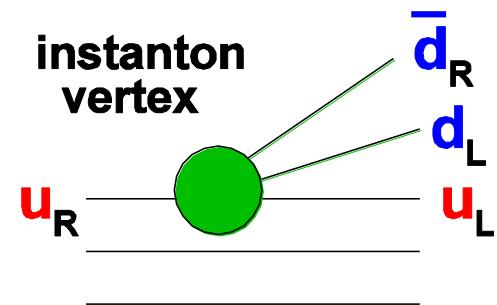
Meson Cloud Models



Chiral-Quark Soliton Model

- nucleon = chiral soliton
- expand in  $1/N_c$
- Quark degrees of freedom in a pion mean-field

Instantons



**Theory:** Thomas, Miller, Kumano, Ma, Londergan, Henley, Speth, Hwang, Melnitchouk, Liu, Cheng/Li, etc.

(For reviews, see Speth and Thomas (1997), Kumano (hep-ph/9702367 ), Garvey and Peng (nucl-ex/0109010))

These models also have implications on

- asymmetry between  $s(x)$  and  $\bar{s}(x)$
- flavor structure of the polarized sea

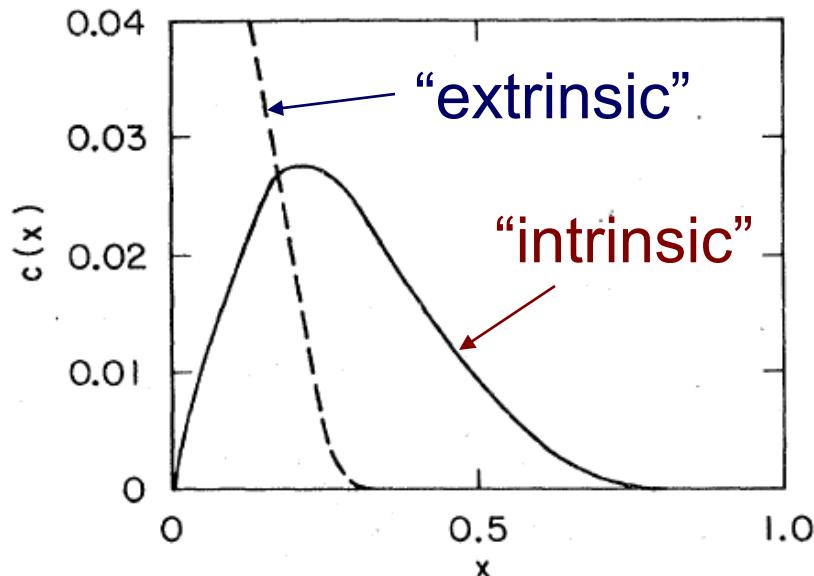
Meson cloud has significant contributions to sea-quark distributions

# Implications on the “intrinsic” quark sea

In 1980, Brodsky, Hoyer, Peterson, Sakai (BHPS) suggested the existence of “intrinsic” charm

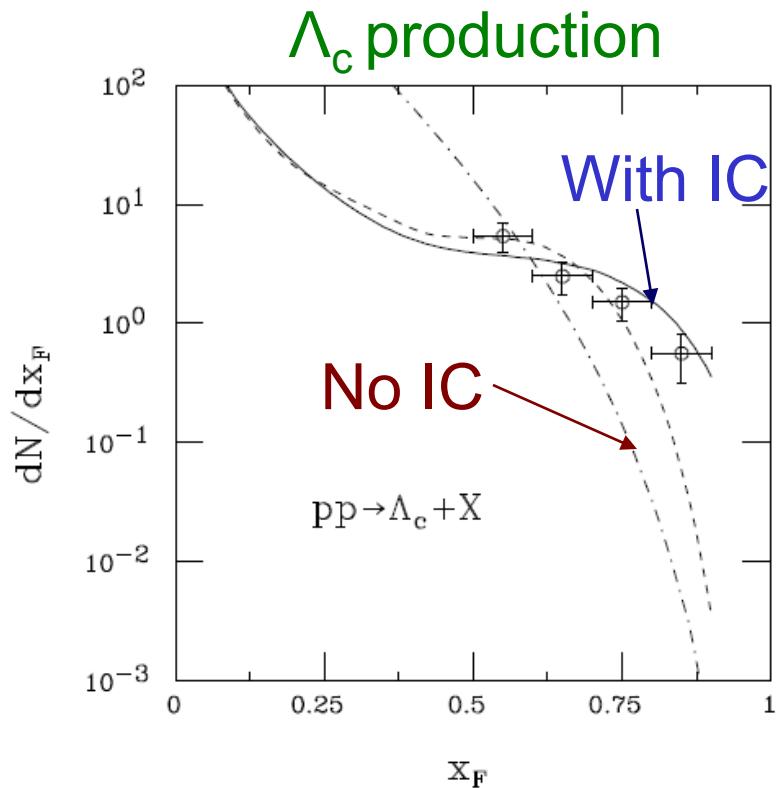
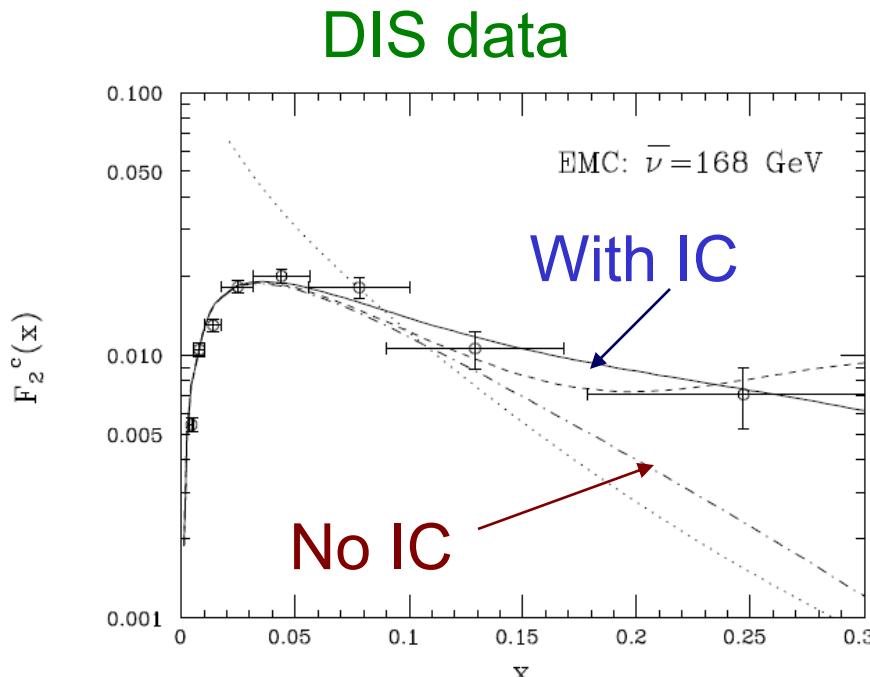
$$|p\rangle = P_{3q} |uud\rangle + P_{5q} |uudQ\bar{Q}\rangle + \dots$$

The “intrinsic”-charm from  $|uudc\bar{c}\rangle$  is “valence”-like and peak at large  $x$  unlike the “extrinsic” sea ( $g \rightarrow c\bar{c}$ )



The  $|uudc\bar{c}\rangle$  intrinsic-charm can lead to large contribution to charm production at large  $x$

# Evidence for the “intrinsic” charm (IC)



Gunion and Vogt (hep-ph/9706252)

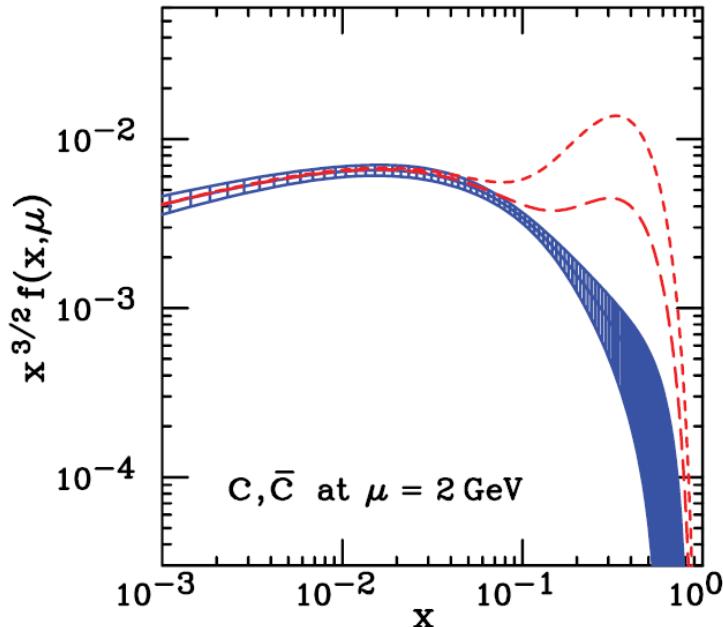
(Evidence is subjected to the uncertainties of charmed-quark parametrization in the PDF)

# A global fit by CTEQ to extract intrinsic-charm

PHYSICAL REVIEW D 75, 054029 (2007)

## Charm parton content of the nucleon

J. Pumplin,<sup>1,\*</sup> H. L. Lai,<sup>1,2,3</sup> and W. K. Tung<sup>1,2</sup>



Blue band corresponds to CTEQ6  
best fit, including uncertainty

Red curves include intrinsic charm of  
1% and 3% ( $\chi^2$  changes only slightly)

We find that the range of IC is constrained to be from zero (no IC) to a level 2–3 times larger than previous model estimates. The behaviors of typical charm distributions within this range are described, and their implications for hadron collider phenomenology are briefly discussed.

No conclusive evidence for intrinsic-charm

# Search for the lighter “intrinsic” quark sea

$$| p \rangle = P_{3q} | uud \rangle + P_{5q} | uudQ\bar{Q} \rangle + \dots$$

No conclusive experimental evidence  
for intrinsic-charm so far

Are there experimental evidences for the intrinsic  
 $| uudu\bar{u} \rangle$ ,  $| uudd\bar{d} \rangle$ ,  $| uuds\bar{s} \rangle$  5-quark states ?

$$P_{5q} \sim 1/m_Q^2$$

The 5-quark states for lighter  
quarks have larger probabilities!

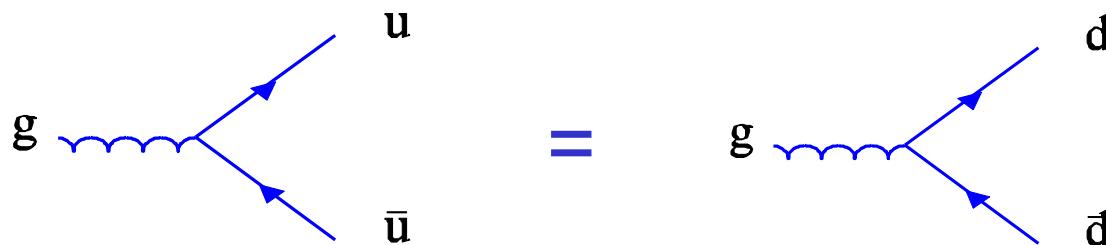
# How to separate the “intrinsic sea” from the “extrinsic sea”?

- Select experimental observables which have no contributions from the “extrinsic sea”
- “Intrinsic sea” and “extrinsic sea” are expected to have different  $x$ -distributions
  - Intrinsic sea is “valence-like” and is more abundant at larger  $x$
  - Extrinsic sea is more abundant at smaller  $x$

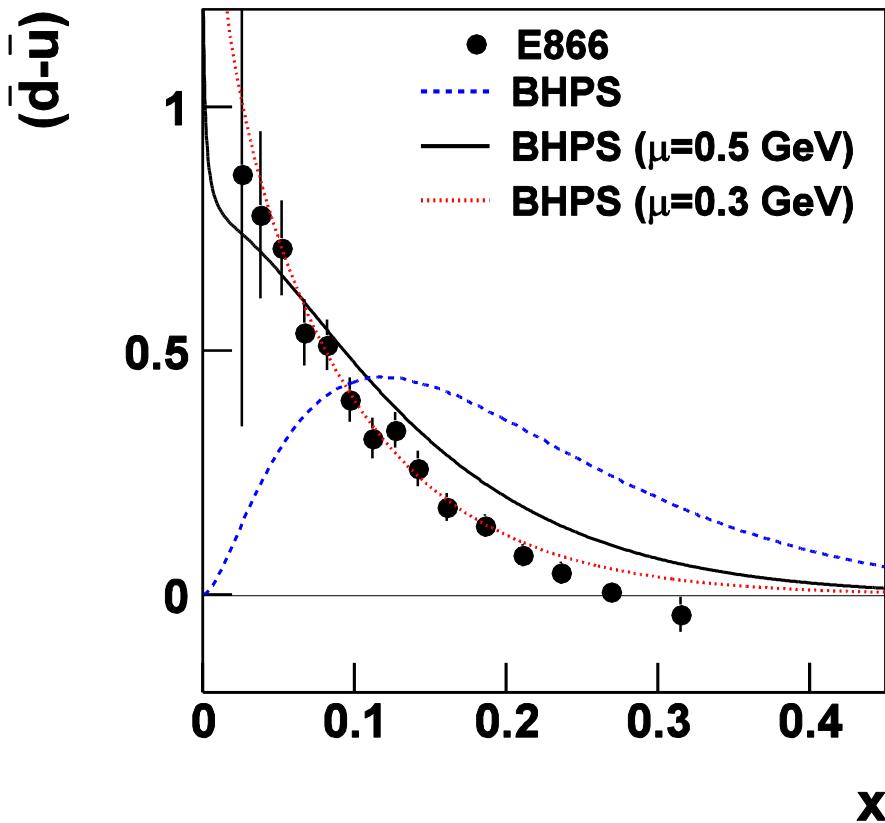
# How to separate the “intrinsic sea” from the “extrinsic sea”?

- Select experimental observables which have no contributions from the “extrinsic sea”

$\bar{d} - \bar{u}$  has no contribution from extrinsic sea ( $g \rightarrow \bar{q}q$ )  
and is sensitive to "intrinsic sea" only



# Comparison between the $\bar{d}(x) - \bar{u}(x)$ data with the intrinsic 5- $q$ model



The data are in good agreement with the 5-q model after evolution from the initial scale  $\mu$  to  $Q^2=54$  GeV $^2$

The difference in the two 5-quark components can also be determined

(W. Chang and JCP , PRL 106, 252002 (2011))

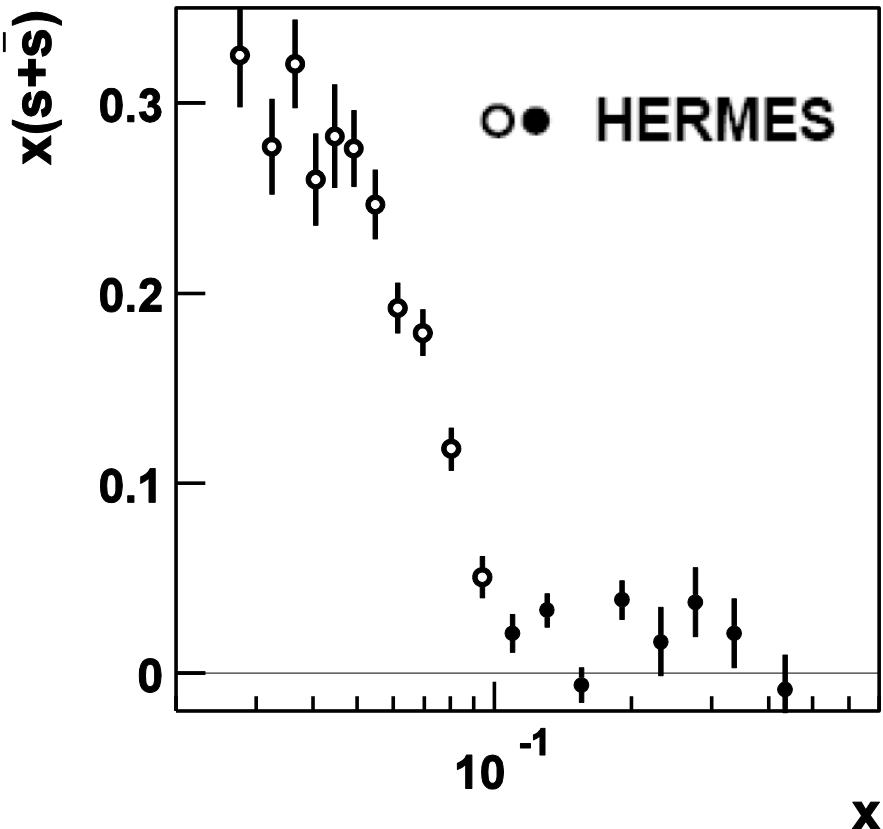
$$P_5^{uudd\bar{d}} - P_5^{uudu\bar{u}} = 0.118$$

# How to separate the “intrinsic sea” from the “extrinsic sea”?

- “Intrinsic sea” and “extrinsic sea” are expected to have different  $x$ -distributions
  - Intrinsic sea is “valence-like” and is more abundant at larger  $x$
  - Extrinsic sea is more abundant at smaller  $x$

An example is the  $s(x) + \bar{s}(x)$  distribution

# Comparison between the $s(x) + \bar{s}(x)$ data with the intrinsic 5- $q$ model

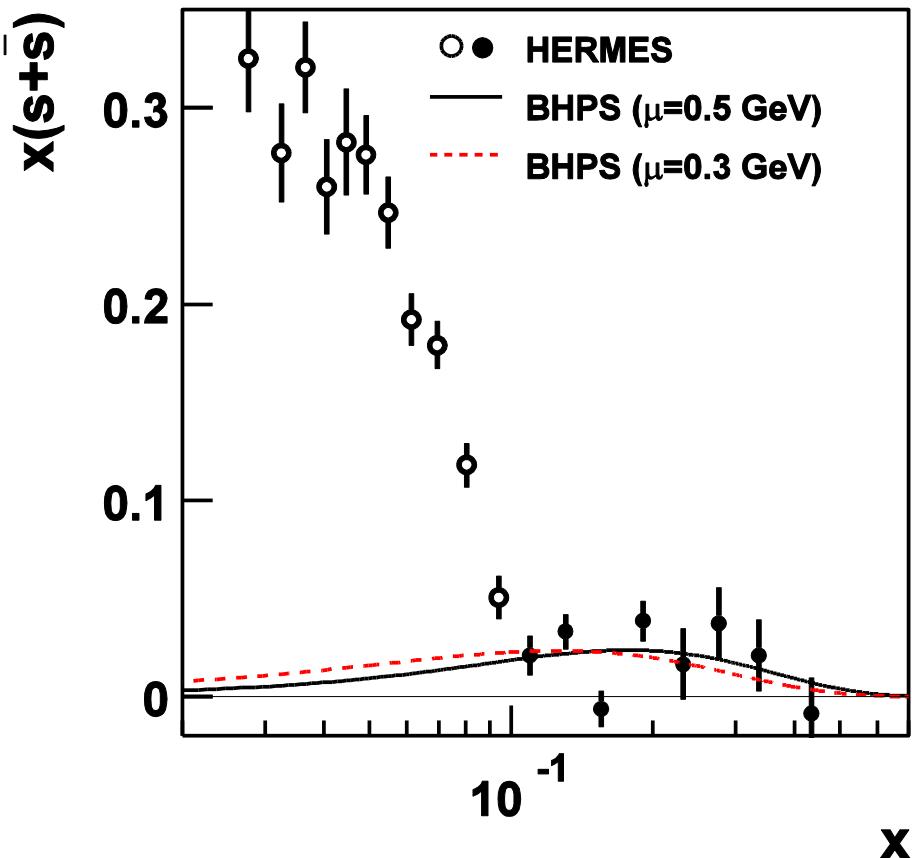


$s(x) + \bar{s}(x)$  from HERMES kaon  
SIDIS data at  $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$

The data appear to consist  
of two different components  
(intrinsic and extrinsic?)

HERMES collaboration, Phys. Lett.  
B666, 446 (2008)

# Comparison between the $s(x) + \bar{s}(x)$ data with the intrinsic 5- $q$ model



$s(x) + \bar{s}(x)$  from HERMES kaon SIDIS data at  $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$

Assume  $x > 0.1$  data are dominate by intrinsic sea (and  $x < 0.1$  are from QCD sea)

This allows the extraction of the intrinsic sea for strange quarks

(W. Chang and JCP, PL B704, 197(2011))

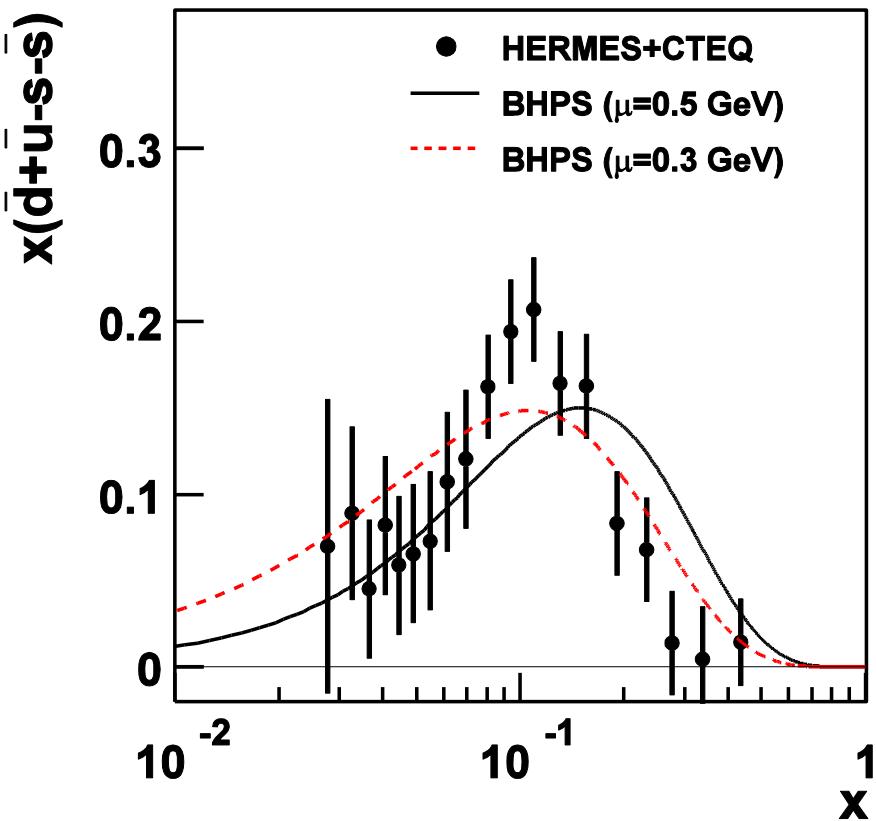
$$P_5^{uud\bar{s}\bar{s}} = 0.024$$

# How to separate the “intrinsic sea” from the “extrinsic sea”?

- Select experimental observables which have no contributions from the “extrinsic sea”

$\bar{d} + \bar{u} - s - \bar{s}$  has no contribution from extrinsic sea ( $g \rightarrow \bar{q}q$ )  
and is sensitive to "intrinsic sea" only

# Comparison between the $\bar{u}(x) + \bar{d}(x)$ - $s(x) - \bar{s}(x)$ data with the intrinsic 5- $q$ model



$\bar{d}(x) + \bar{u}(x)$  from CTEQ6.6  
 $s(x) + \bar{s}(x)$  from HERMES

$$\bar{u} + \bar{d} - s - \bar{s}$$

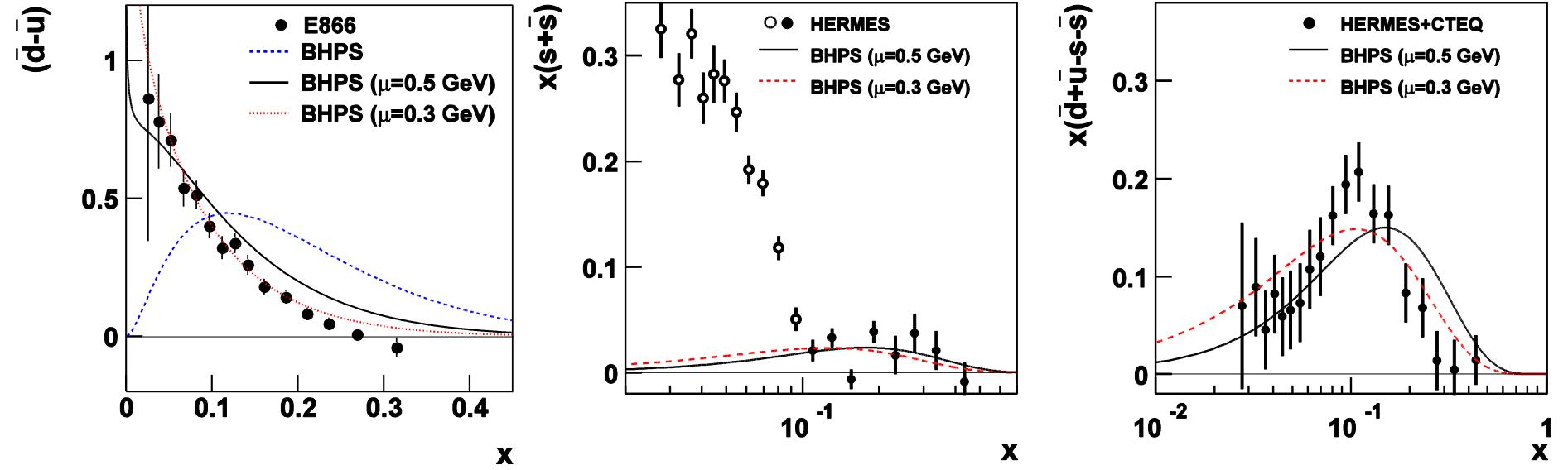
$$\sim P_5^{uudu\bar{u}} + P_5^{uudd\bar{d}} - 2P_5^{uuds\bar{s}}$$

(not sensitive to extrinsic sea)

(W. Chang and JCP, PL B704, 197(2011))

$$P_5^{uudu\bar{u}} + P_5^{uudd\bar{d}} - 2P_5^{uuds\bar{s}} = 0.314$$

# Extraction of the various five-quark components for light quarks



$$P_5^{uudd\bar{d}} - P_5^{uudu\bar{u}} = 0.118$$

$$P_5^{uud\bar{s}\bar{s}} = 0.024$$

$$P_5^{uudu\bar{u}} + P_5^{uudd\bar{d}} - 2P_5^{uud\bar{s}\bar{s}} = 0.314$$

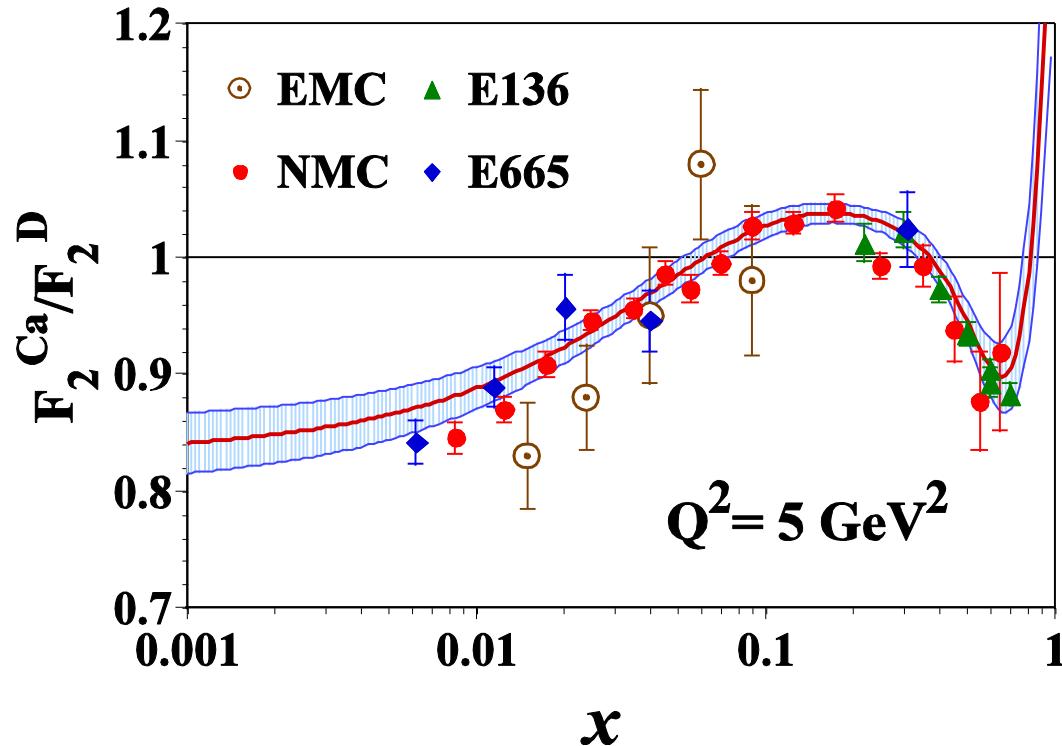
$$P_5^{uudd\bar{d}} = 0.240; \quad P_5^{uudu\bar{u}} = 0.122; \quad P_5^{uud\bar{s}\bar{s}} = 0.024$$

# Future Prospect (Experimental)

- Kaon production in SIDIS at COMPASS and 12 GeV Jlab/EIC for more information on s and s-bar?
- Kaon-induced Drell-Yan at COMPAS for probing s and s-bar?
- Open-charm and open-beauty production at forward rapidity at LHC to search for intrinsic charm and beauty.
- $J/\psi$  and Upsilon production at forward  $x_F$  to search for intrinsic charm and beauty at AFTER?

# Modification of Parton Distributions in Nuclei

EMC effect observed in DIS

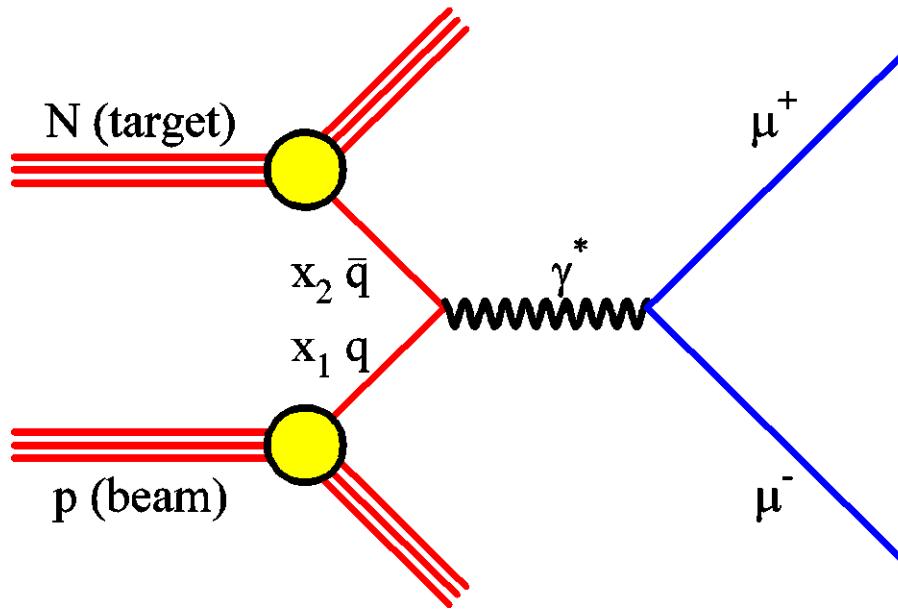


(Ann. Rev. Nucl. Part. Phys., Geesaman, Sato and Thomas)

$F_2$  contains contributions from quarks and antiquarks

How are the antiquark distributions modified in nuclei?

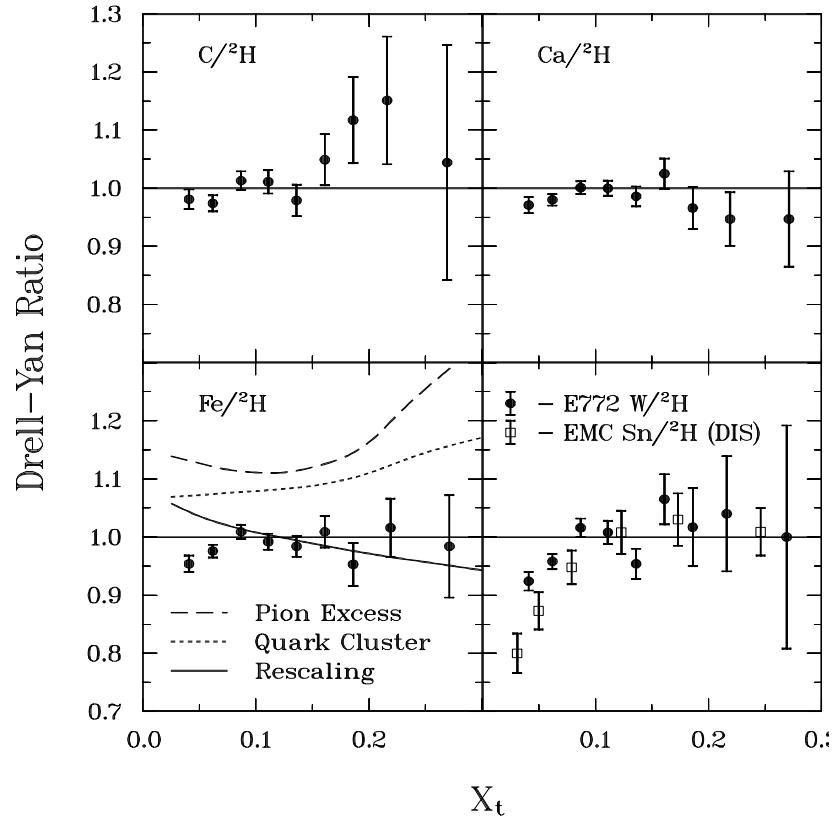
# Drell-Yan on nuclear targets



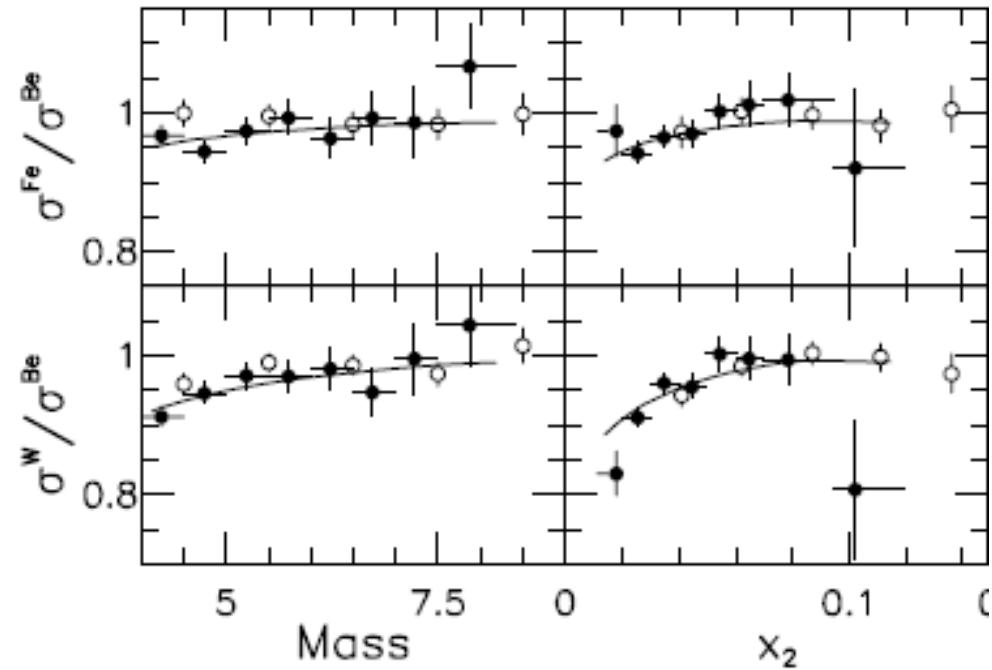
$$\frac{\sigma^{pA}}{\sigma^{pd}} \approx \frac{\bar{u}_A(x)}{\bar{u}_N(x)}$$

The  $x$ -dependence of  $\bar{u}_A(x)/\bar{u}_N(x)$  can be directly measured

# Drell-Yan on nuclear targets



PRL 64 (1990) 2479

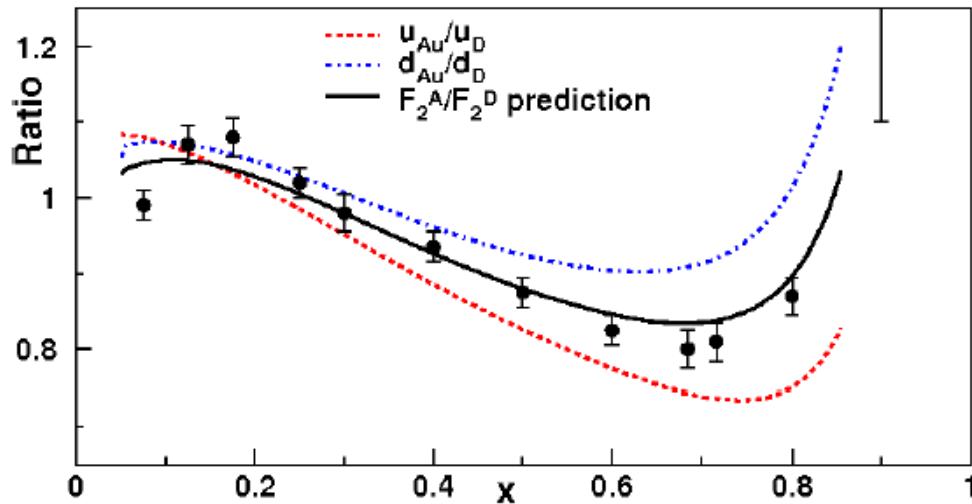


PRL 83 (1999) 2304

No evidence for enhancement of antiquark in nuclei !?

LHC (AFTER) can explore antiquark shadowing/saturation in nuclei at small  $x$

# Flavor dependence of the EMC effects ?



Isovector mean-field generated in  $Z \neq N$  nuclei  
can modify nucleon's  $u$  and  $d$  PDFs in nuclei

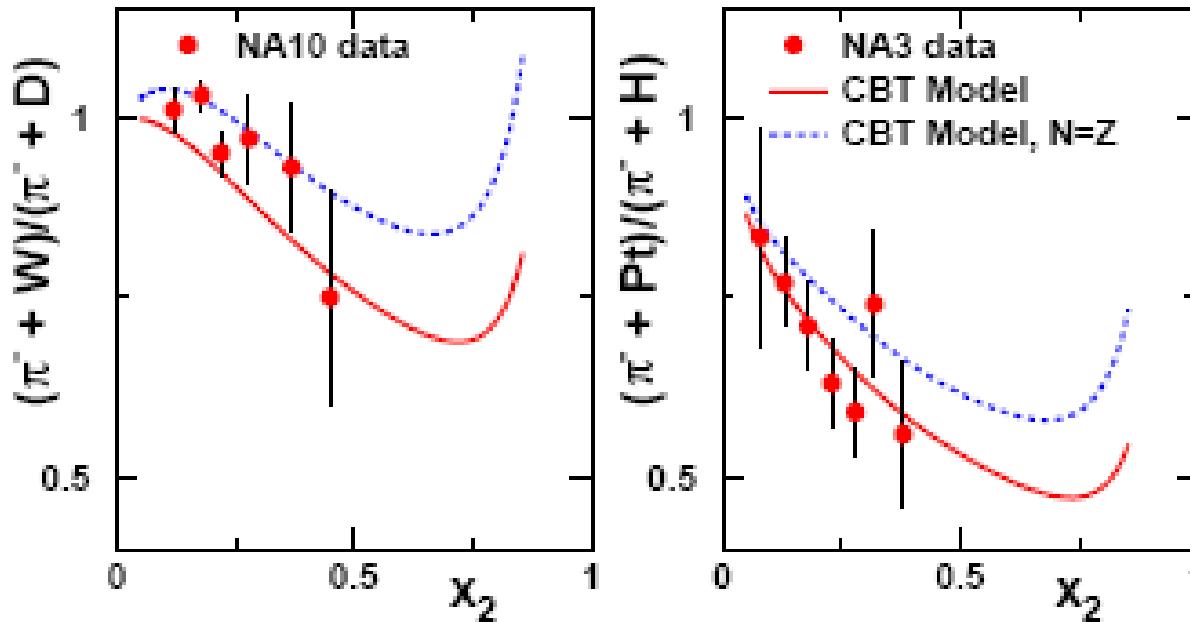
Cloet, Bentz, and Thomas, arXiv:0901.355  
(see also Kumano et al.)

How can one check this prediction?

- SIDIS (Semi-inclusive DIS) and PVDIS (Parity-violating DIS)
- Pion-induced Drell-Yan

# Pion-induced Drell-Yan and the flavor-dependent EMC effect

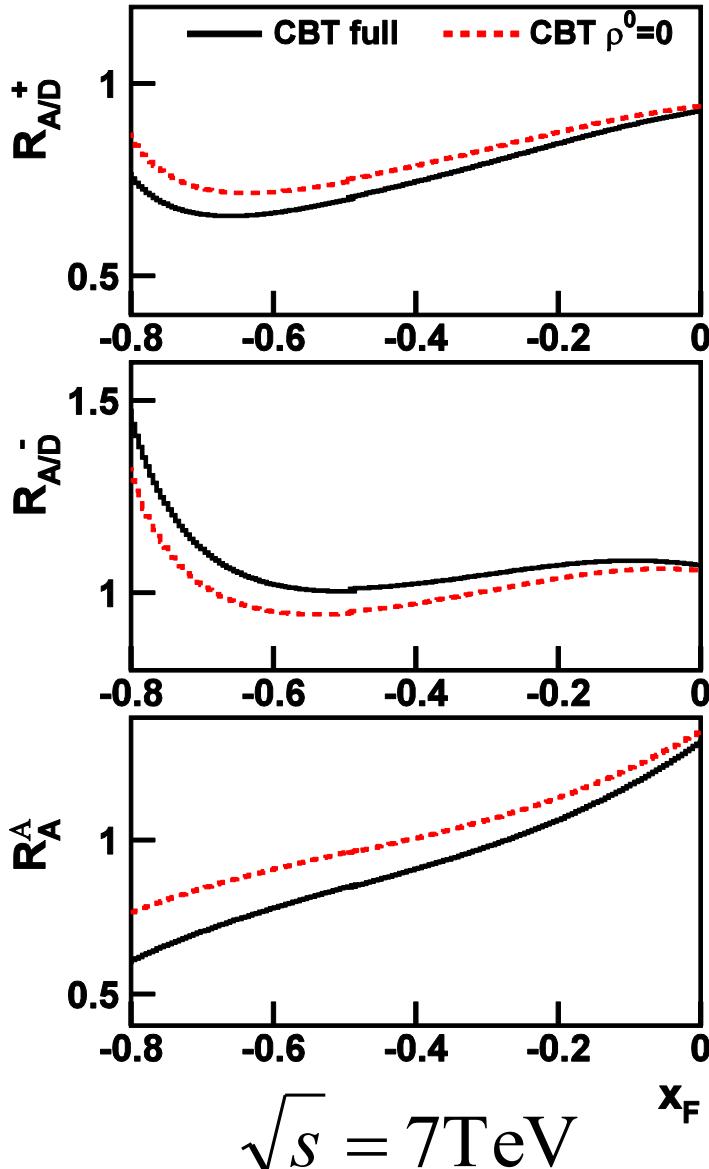
$$\frac{\sigma^{DY}(\pi^- + A)}{\sigma^{DY}(\pi^- + D)} \approx \frac{u_A(x)}{u_D(x)}$$



**Red (blue) curves correspond to flavor-dependent (independent) EMC**

(D. Dutta, JCP, Cloet, Gaskell, arXiv: 1007.3916)

# W-production at LHC and the flavor-dependent EMC effect



$$R_{A/D}^+ \equiv \frac{d\sigma(p+A \rightarrow W^+ + x)}{d\sigma(p+D \rightarrow W^+ + x)}$$

$$\approx \frac{u_A(x_2)}{u_D(x_2)}$$

$$R_{A/D}^- \equiv \frac{d\sigma(p+A \rightarrow W^- + x)}{d\sigma(p+D \rightarrow W^- + x)}$$

$$\approx \frac{d_A(x_2)}{d_D(x_2)}$$

(Chang, Cloet, Dutta,  
JCP, 1109.3108)

$$R_A^\pm \equiv \frac{d\sigma(p+A \rightarrow W^\pm + x)}{d\sigma(p+A \rightarrow W^\mp + x)}$$

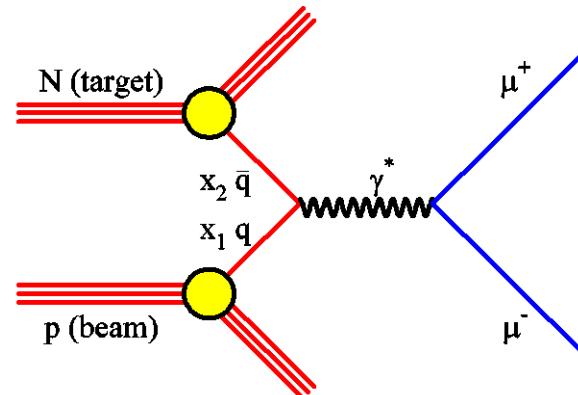
$$\approx \frac{\bar{d}_p(x_1)}{\bar{u}_p(x_1)} \frac{u_A(x_2)}{d_A(x_2)}$$

Pb beam at AFTER is ideal for exploring the negative  $x_F$  region

# Can one measure $\bar{d} / \bar{u}$ in nuclei?

The Drell-Yan Process

$$pN \rightarrow \mu^+ \mu^- X$$



Drell-Yan ratios  
for p-A / p-d :

$$\frac{2\sigma^{pA}}{A\sigma^{pd}} \approx \frac{2}{A} \frac{Z + N \bar{d} / \bar{u}}{1 + \bar{d} / \bar{u}}$$

Assuming  $\bar{d} / \bar{u} = 1.5$  for the nucleons at  $x=0.15$ , then the above ratios are:

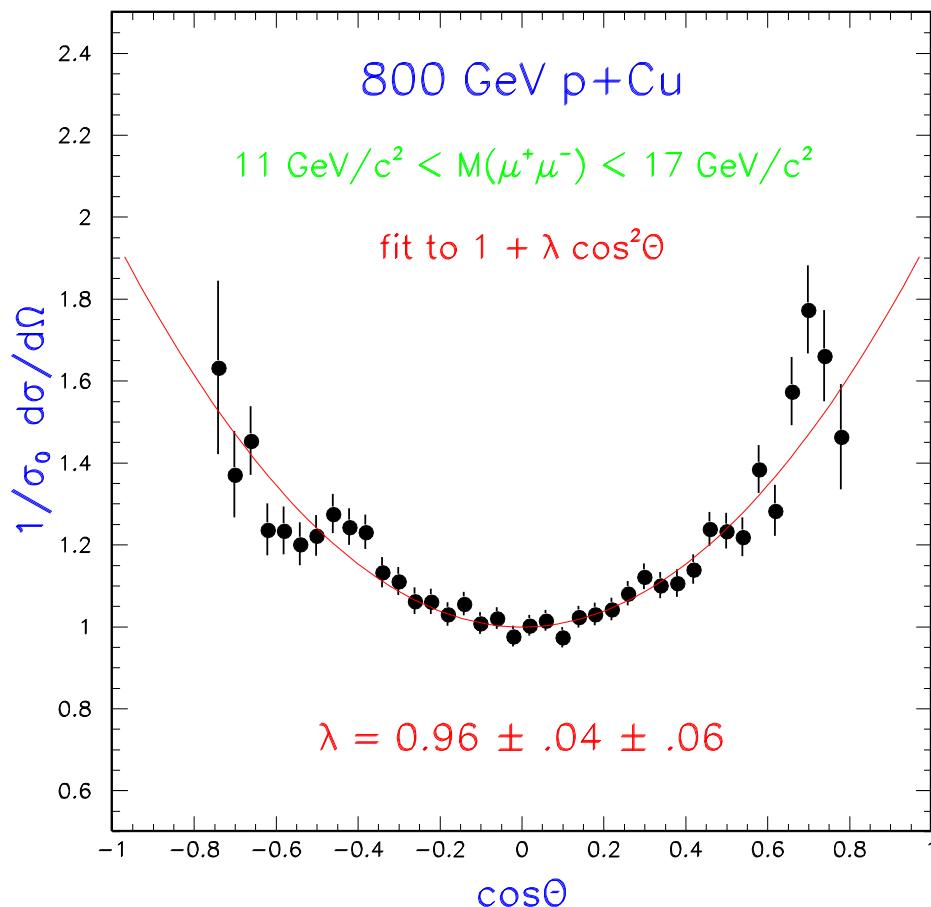
1.0 for  ${}^{40}\text{Ca}$ , 1.042 for  ${}^{208}\text{Pb}$

Could probably be measured at AFTER

# Dilepton angular distribution

Decay Angular Distribution of “naïve” Drell-Yan:

$$\frac{d\sigma}{d\Omega} = \sigma_0(1 + \cos^2 \theta)$$



Data from  
Fermilab E772

# Polarization of J/Ψ and Υ

- Decay angular distribution in the quarkonium rest frame

$$\frac{d\sigma}{d\Omega} \sim 1 + \lambda \cos^2 \theta$$

- \* Transverse :  $\sigma_T$  ; helicity:  $\pm 1$ ;  $\lambda = 1$
- \* Longitudinal :  $\sigma_L$  ; helicity: 0;  $\lambda = -1$
- \* Unpolarized :  $\sigma_T = 2\sigma_L$  ; helicity: 0,  $\pm 1$ ;  $\lambda = 0$
- $\lambda = \frac{\sigma_T - 2\sigma_L}{\sigma_T + 2\sigma_L} = (1 - 2\sigma_L / \sigma_T) / (1 + 2\sigma_L / \sigma_T)$
- $\sigma_L / \sigma_T$  depends on the color - spin states of the Q $\bar{Q}$  pair :

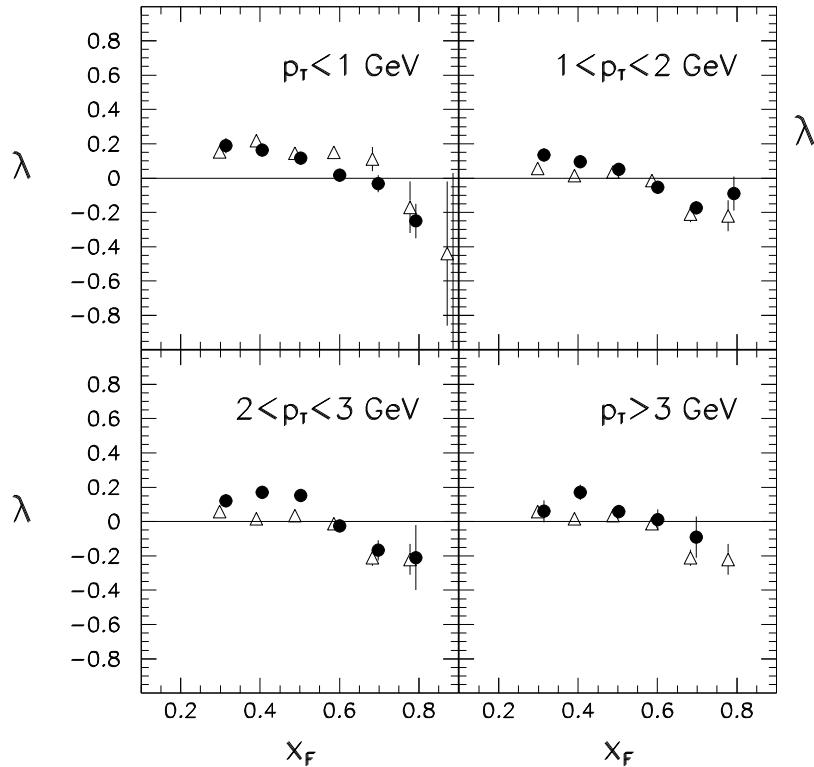
State:	$^3S_1^{(1)}$	$^1S_0^{(8)}$	$^3P_J^{(8)}$	$^3S_1^{(8)}$
$\sigma_L / \sigma_T :$	1/3.4	1/2	1/6	0/1
- Polarization of Q $\bar{Q}$  is sensitive to the production mechanism

# Polarization of J/ $\Psi$ in p + Cu Collision

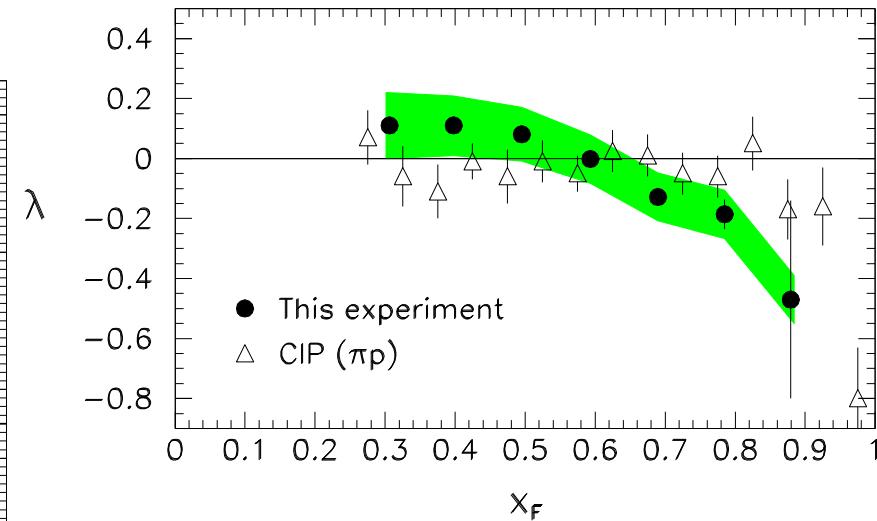
$$d\sigma/d\Omega \sim 1 + \lambda \cos^2\theta$$

( $\lambda=1$ : transversely polarized,  $\lambda = -1$ : longitudinally polarized  $\lambda = 0$ , unpolarized)

E866 data



hep-ex/030801



$\lambda$  is small, but nonzero

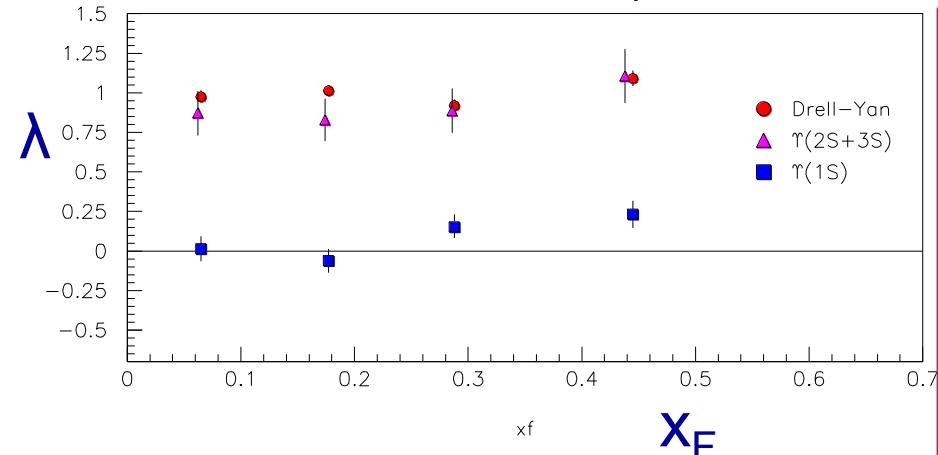
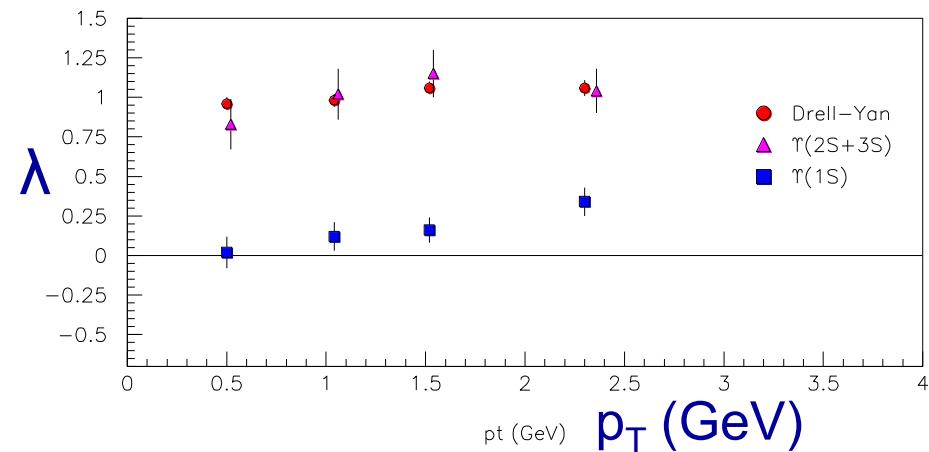
$\lambda$  becomes negative at large  $x_F$

No strong  $p_T$  dependence for  $\lambda$

# Polarization of Y(1S), Y(2S+3S)

$p + Cu \rightarrow Y + x$  (E866 beam-dump data)

$\lambda$  for D-Y, Y(1S), Y(2S+3S)



Brown et al. PRL 86, 2529 (2001)

D-Y is transversely polarized

Y(1S) is slightly polarized  
(like J/ $\Psi$ )

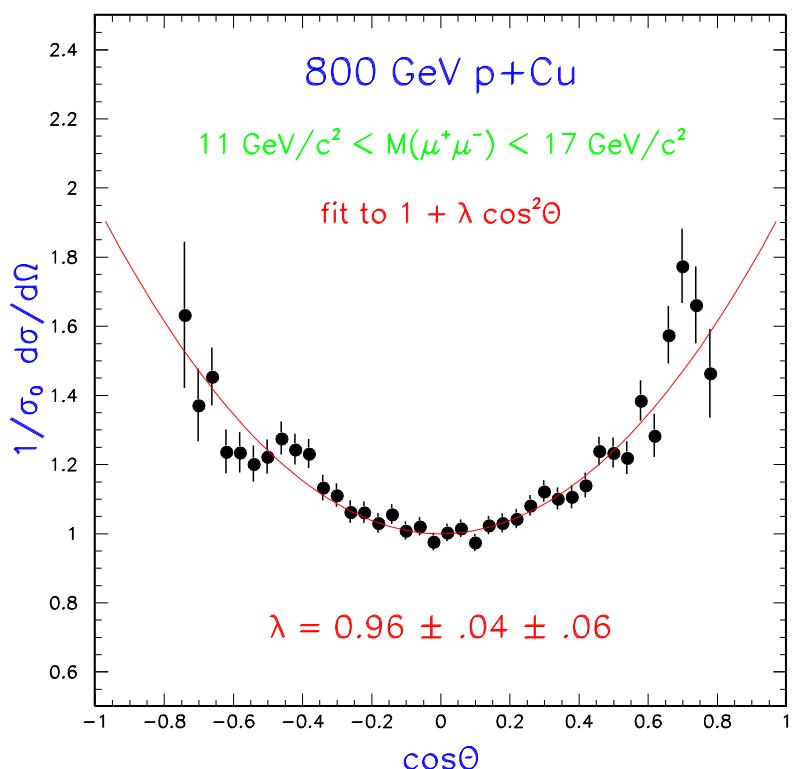
Y(2S+3S) is transversely  
polarized!

Measurement of Upsilon  
polarization in  $p+p$  and  $p+d$   
at AFTER would be very  
interesting (nuclear  
dependence?)

# Parity violation in Drell-Yan?

Forward-backward asymmetry in decay angular distribution of Drell-Yan:

$$\frac{d\sigma}{d\Omega} = \sigma_0 (1 + \boxed{a \cos \theta} + b \cos^2 \theta)$$



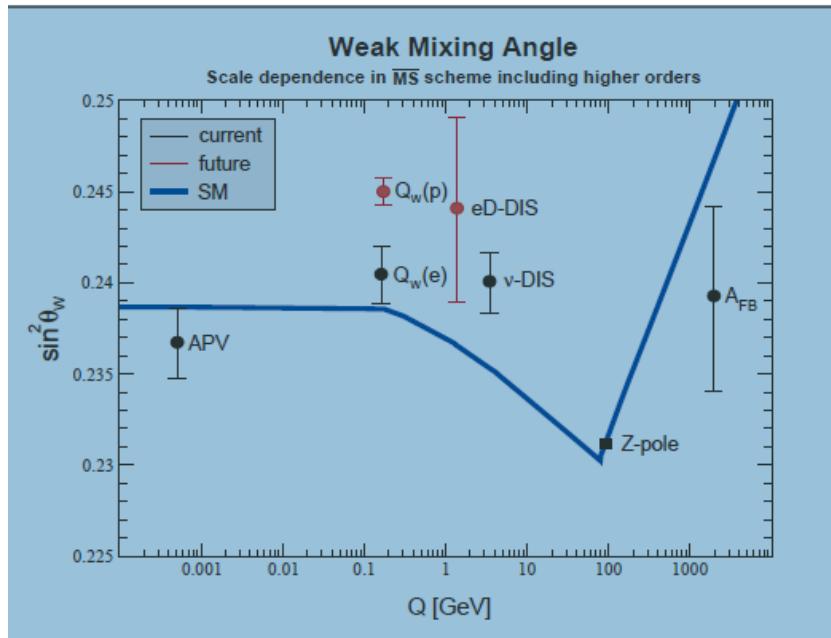
Interference between  $\gamma$  and  $Z^0$  diagrams can lead to non-zero parity-violating  $\cos \theta$  term

Parity violation in D-Y can probe Weinberg angles at low  $Q^2$

# Parity violation in Drell-Yan?

Forward-backward asymmetry in decay angular distribution of Drell-Yan:

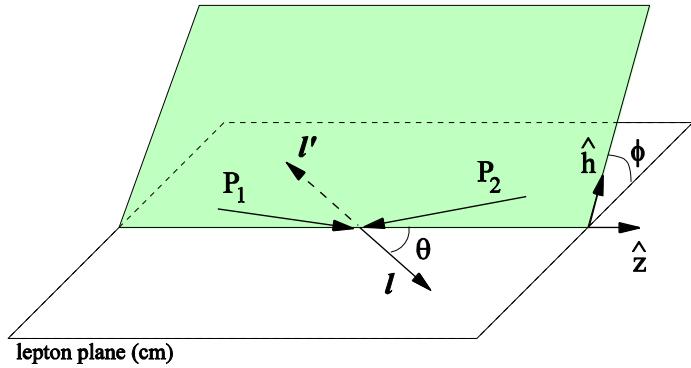
$$\frac{d\sigma}{d\Omega} = \sigma_0(1 + [a \cos \theta] + b \cos^2 \theta)$$



Interference between  $\gamma$  and  $Z^0$  diagrams can lead to non - zero parity - violating  $\cos \theta$  term

Parity violation in D-Y can probe Weinberg angles at low  $Q^2$  (Never been studied yet)

# Drell-Yan decay angular distributions



$\Theta$  and  $\Phi$  are the decay polar and azimuthal angles of the  $\mu^+$  in the dilepton rest-frame

**Collins-Soper frame**

A general expression for Drell-Yan decay angular distributions:

$$\left(\frac{1}{\sigma}\right)\left(\frac{d\sigma}{d\Omega}\right) = \left[\frac{3}{4\pi}\right] \left[ 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi \right]$$

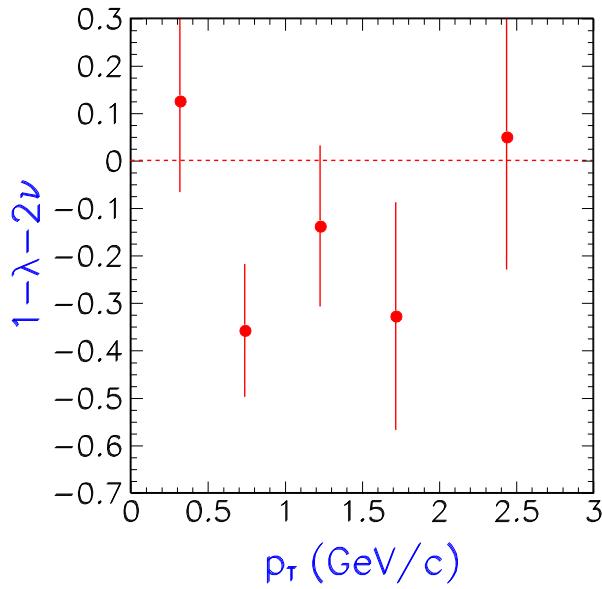
$\lambda$  can differ from 1, but should satisfy  $1 - \lambda = 2\nu$  (Lam-Tung)

- Reflect the spin-1/2 nature of quarks  
(analog of the Callan-Gross relation in DIS)
- Insensitive to QCD - corrections

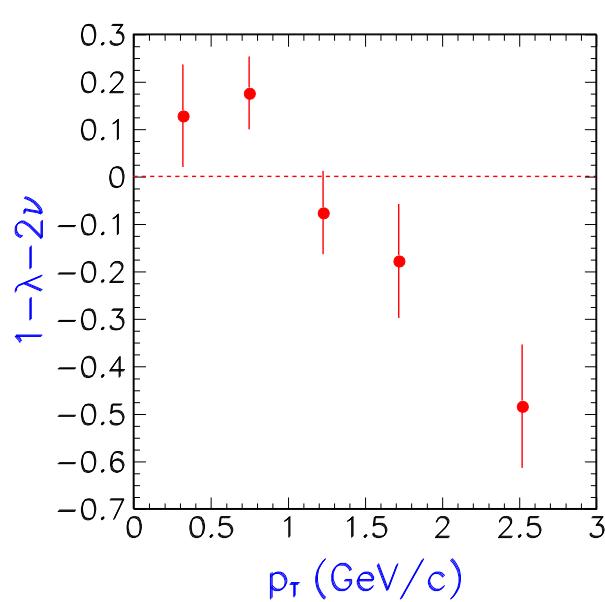
# Decay angular distributions in pion-induced Drell-Yan

## Is the Lam-Tung relation violated?

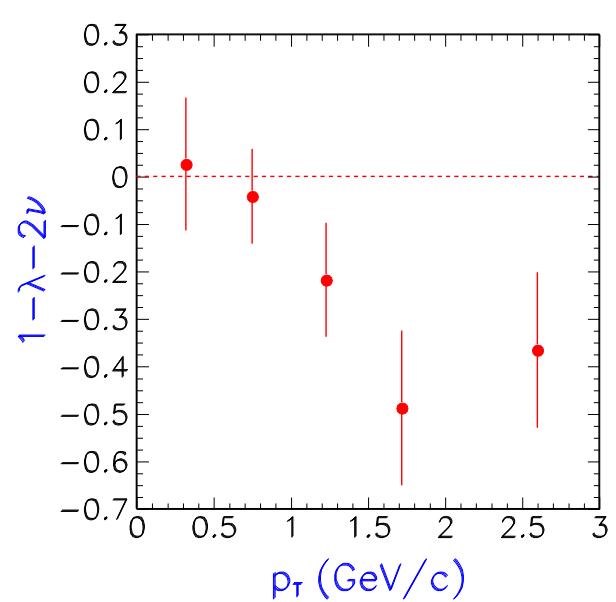
140 GeV/c



194 GeV/c



286 GeV/c



Data from NA10 (Z. Phys. 37 (1988) 545)

Violation of the Lam-Tung relation suggests

new mechanisms with non-perturbative origin

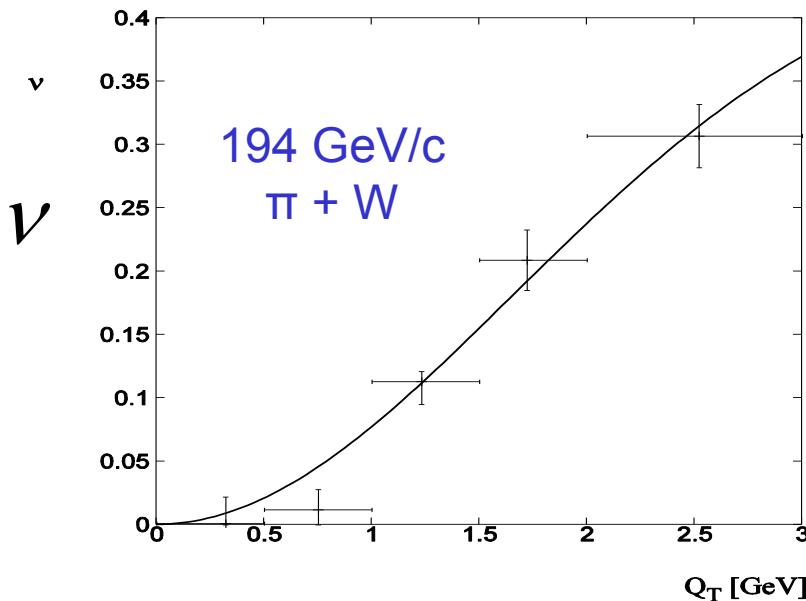
- $q - \bar{q}$  spin correlation in QCD color field (Nachtmann et al.)

# Boer-Mulders function $h_1^\perp$



- $h_1^\perp$  represents a correlation between quark's  $k_T$  and transverse spin in an unpolarized hadron
- $h_1^\perp$  is a time-reversal odd, chiral-odd TMD parton distribution
- $h_1^\perp$  can lead to an azimuthal  $\cos(2\phi)$  dependence in Drell-Yan

$$\left(\frac{1}{\sigma}\right)\left(\frac{d\sigma}{d\Omega}\right) = \left[\frac{3}{4\pi}\right] \left[1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi\right]$$



Boer, PRD 60 (1999) 014012

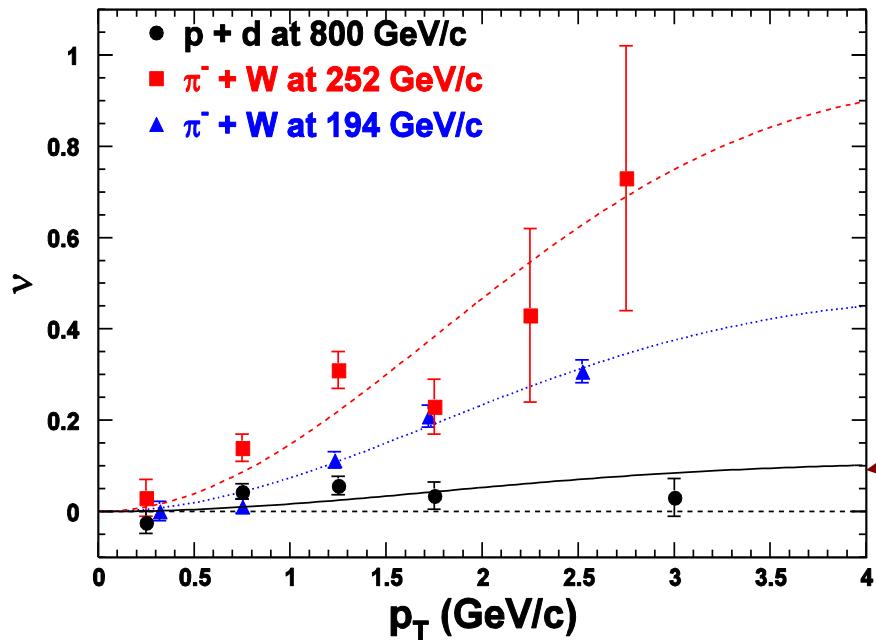
- Observation of large  $\cos(2\Phi)$  dependence in Drell-Yan with pion beam

- $\nu \propto h_1^\perp(x_q)h_1^\perp(x_{\bar{q}})$

- B-M functions have same signs for pion and nucleon

# Azimuthal $\cos 2\Phi$ Distribution in p+p and p+d Drell-Yan

E866 Collab., Lingyan Zhu et al.,  
PRL 99 (2007) 082301; PRL 102 (2009) 182001



Small  $\nu$  is observed  
for p+d and p+p D-Y

With Boer-Mulders function  $h_1^\perp$ :

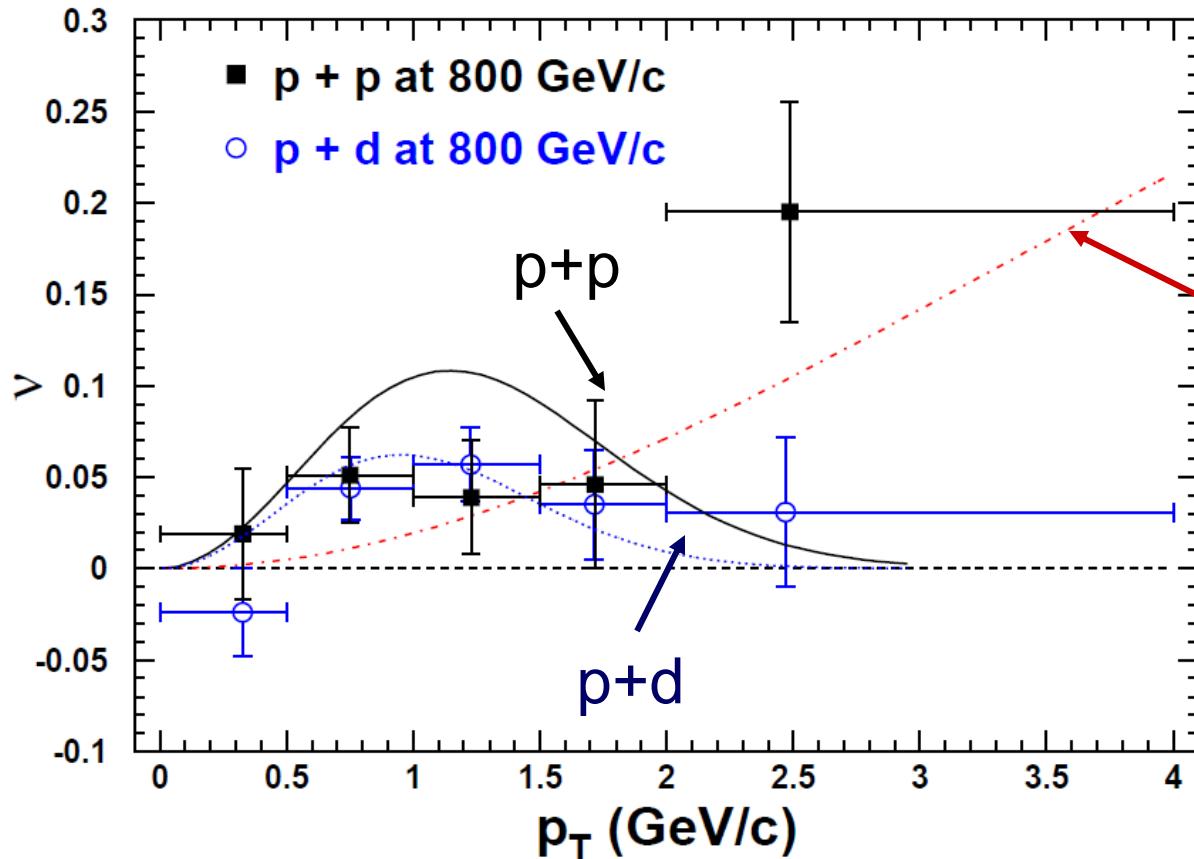
$$\nu(\pi^- W \rightarrow \mu^+ \mu^- X) \sim [\text{valence } h_1^\perp(\pi)] * [\text{valence } h_1^\perp(p)]$$

$$\nu(p d \rightarrow \mu^+ \mu^- X) \sim [\text{valence } h_1^\perp(p)] * [\text{sea } h_1^\perp(p)]$$

Sea-quark BM functions are much smaller than valence quarks

# Results on $\cos 2\Phi$ Distribution in p+p Drell-Yan

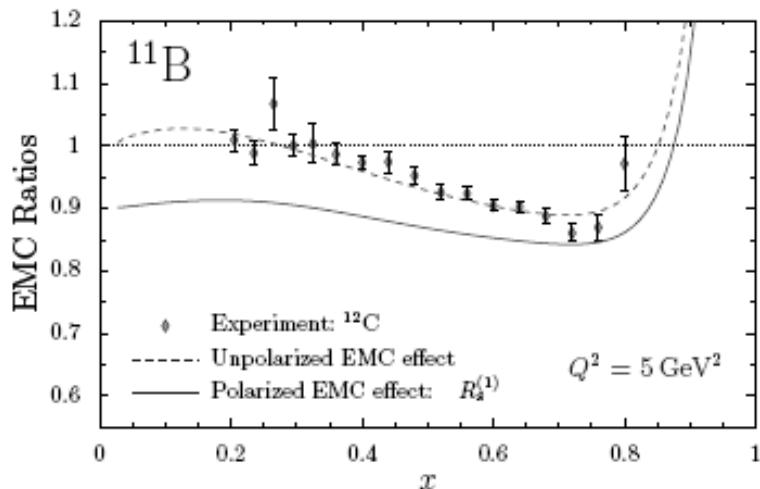
L. Zhu et al., PRL 102 (2009) 182001



QCD  
(Boer, Vogelsang;  
Berger, Qiu,  
Rodriguez-Pedraza)

Data at larger  $p_T$  (from AFTER)  
would be very interesting

# Nuclear modification of spin-dependent PDF?



EMC effect  
for  $g_1(x)$

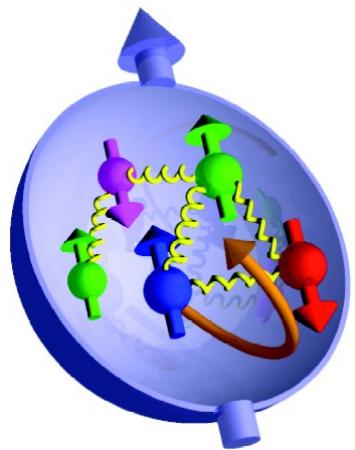
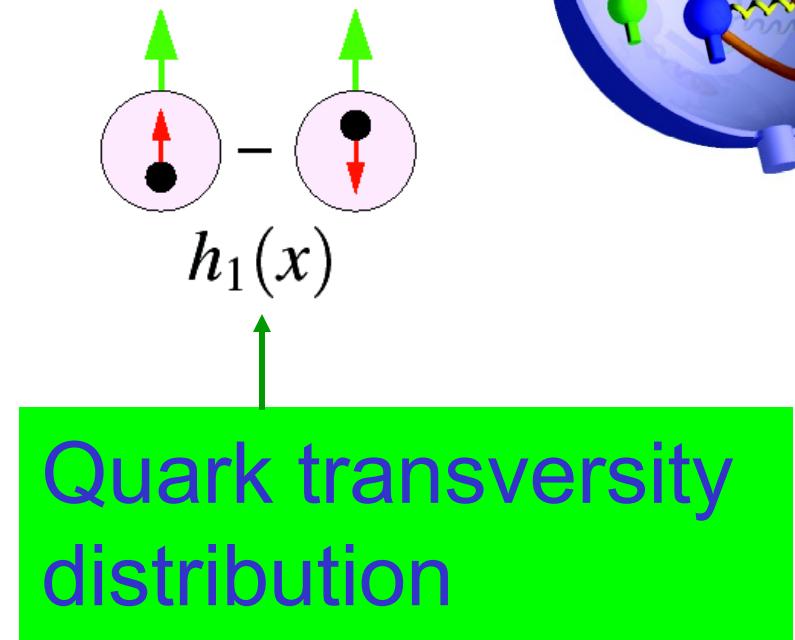
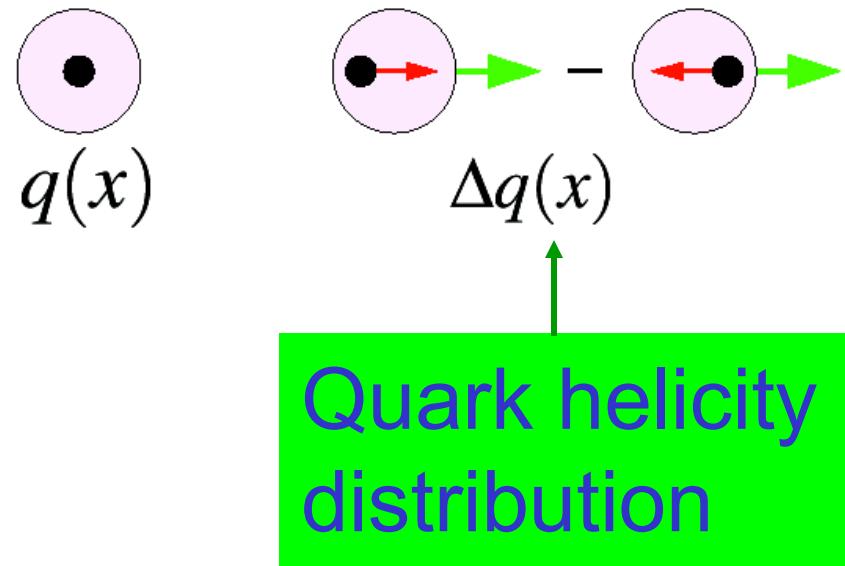
Benz, Cloet, Thomas  
et al., arXiv:0711.0392

Remains to be tested by experiments

Measure the nuclear modification of Boer-Mulders  
functions with Drell-Yan ?  
(only unpolarized Drell-Yan is required)

# Polarized Drell-Yan with polarized proton beam?

- Polarized Drell-Yan experiments have never been done before
- Provide unique information on the quark (antiquark) spin

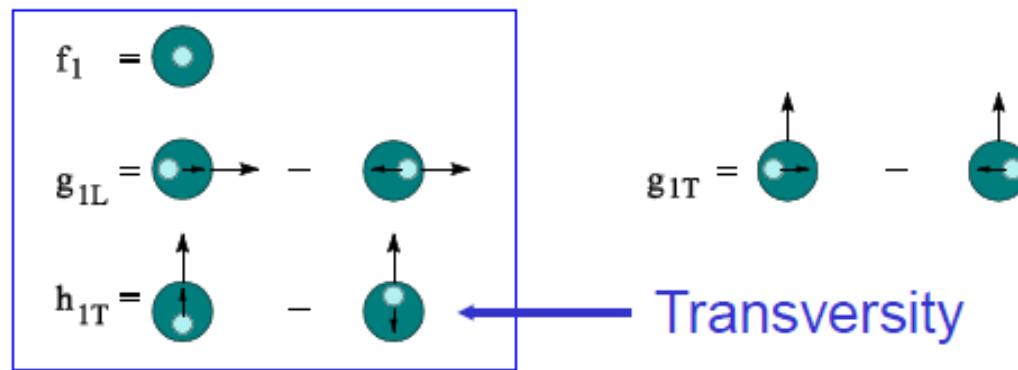


# Drell-Yan and Transverse Momentum Dependent (TMD) Quark Distributions

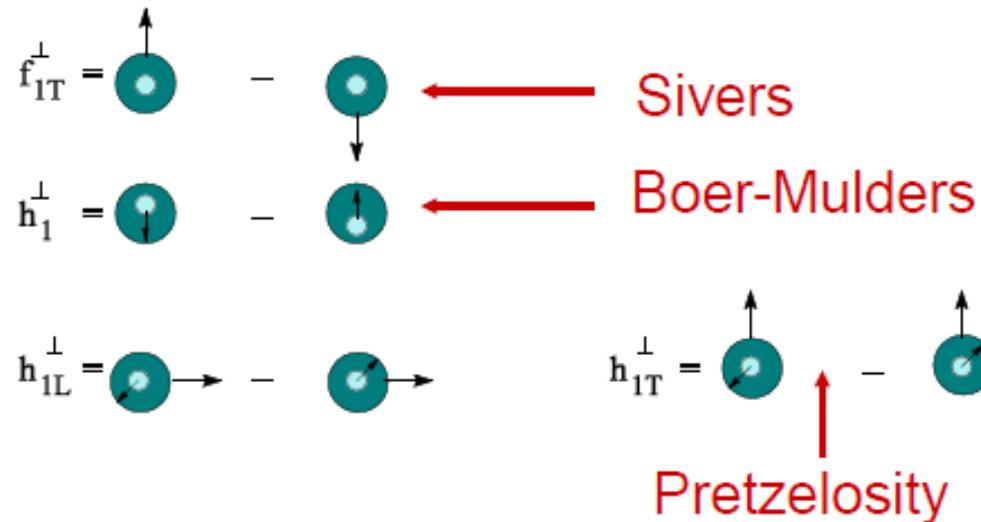
Leading-Twist Quark Distributions

( A total of eight distributions)

Three would remain after  $k_\perp$  integration



The other five are transverse momentum ( $k_\perp$ ) dependent (TMD)



# Transversity and Transverse Momentum Dependent PDFs are also probed in Drell-Yan

a) Boer-Mulders functions:

- Unpolarized Drell-Yan:  $d\sigma_{DY} \propto h_1^\perp(x_q)h_1^\perp(x_{\bar{q}})\cos(2\phi)$

b) Sivers functions:

- Single transverse spin asymmetry in polarized Drell-Yan:

$$A_N^{DY} \propto f_{1T}^\perp(x_q)f_{\bar{q}}(x_{\bar{q}})$$

c) Transversity distributions:

- Double transverse spin asymmetry in polarized Drell-Yan:

$$A_{TT}^{DY} \propto h_1(x_q)h_1(x_{\bar{q}})$$

- Drell-Yan does not require knowledge of the fragmentation functions
- T-odd TMDs are predicted to change sign from DIS to DY  
(Boer-Mulders and Sivers functions)

Remains to be tested experimentally!

# Outstanding questions to be addressed by future Drell-Yan experiments

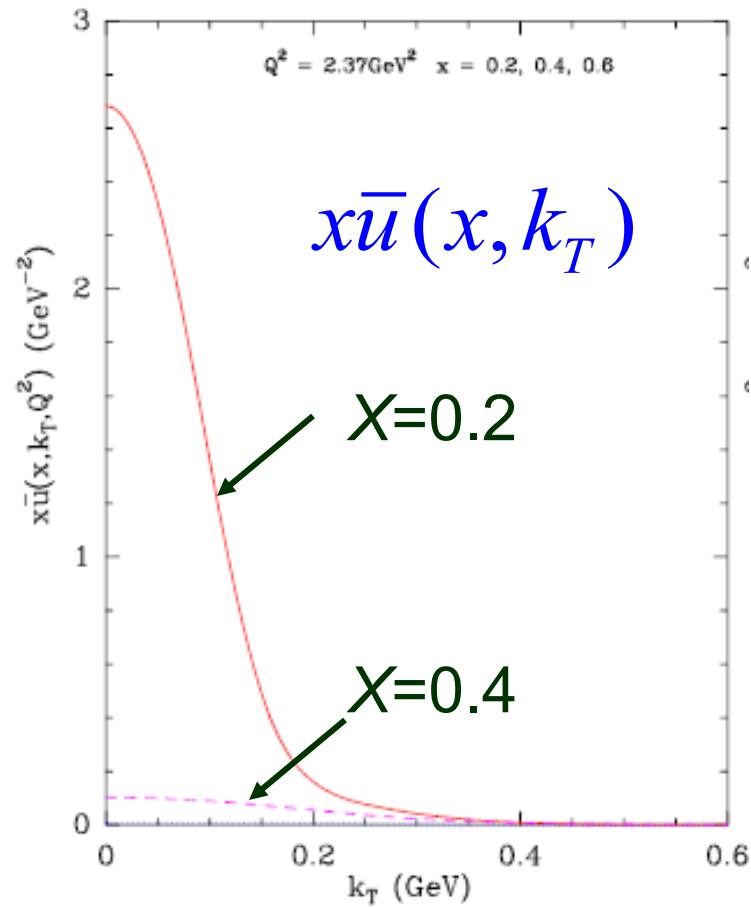
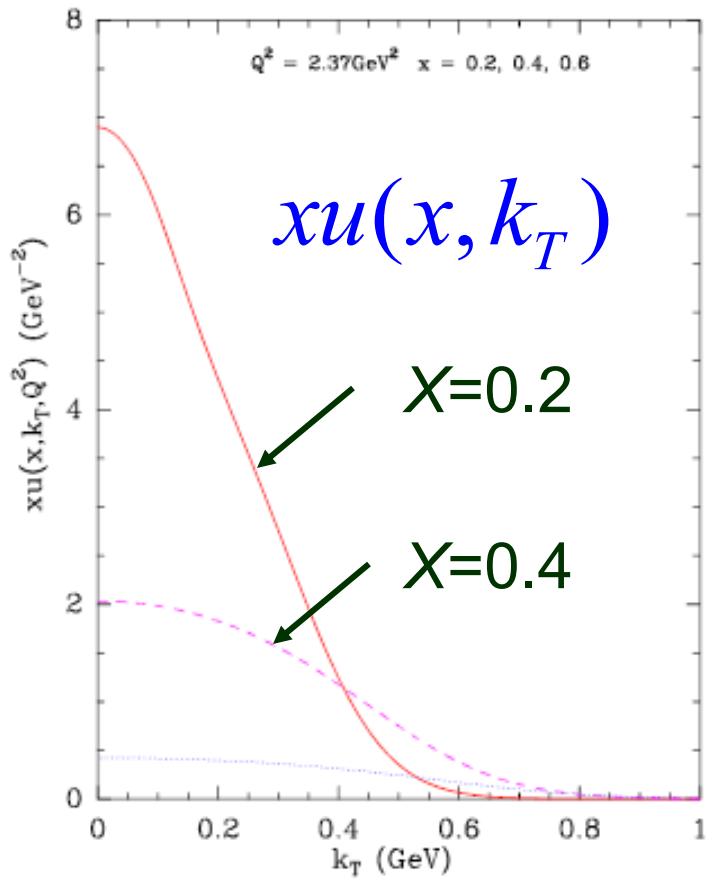
- Does Sivers function change sign between DIS and Drell-Yan?
- Does Boer-Mulders function change sign between DIS and Drell-Yan?
- Are all Boer-Mulders functions alike (proton versus pion Boer-Mulders functions)
- Flavor dependence of TMD functions
- Independent measurement of transversity with Drell-Yan

Can be studied at COMPASS, RHIC, FAIR, LHC, JPARC, JINR, etc

# What do we know about the quark and gluon intrinsic transverse momentum distributions?

- Does the quark  $k_T$  distribution depend on  $x$ ?
- Do valence quarks and sea quarks have different  $k_T$  distributions?
- Do  $u$  and  $d$  quarks have the same  $k_T$  distribution?
- Do nucleons and mesons have different quark  $k_T$  distribution?
- Do gluons have  $k_T$  distribution different from quarks?
- Important for extracting the TMD parton distributions
- Interesting physics in its own right

# Flavor and $x$ -dependent $k_T$ -distributions? (Bourrely, Buccella, Soffer, arXiv:1008.5322)

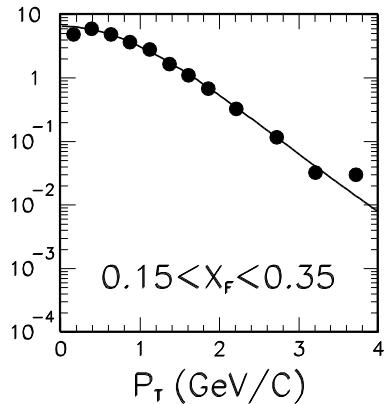
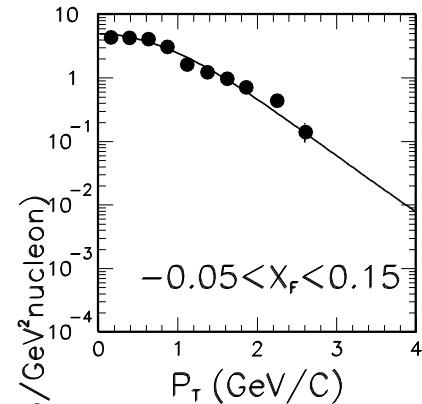


- $\langle k_T \rangle$  increases when  $x$  increases
- $\langle k_T \rangle$  for sea quarks is smaller than for valence quarks

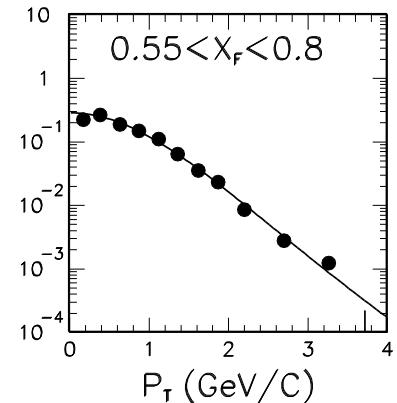
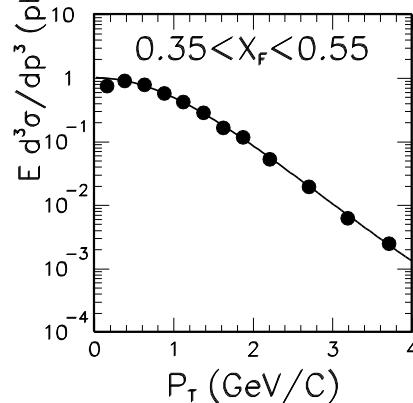
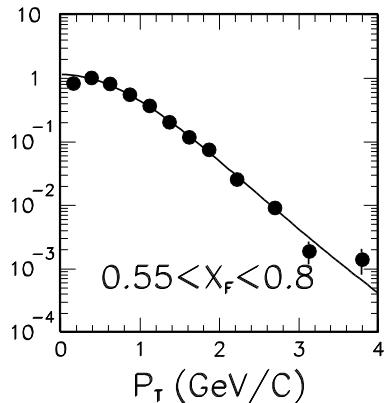
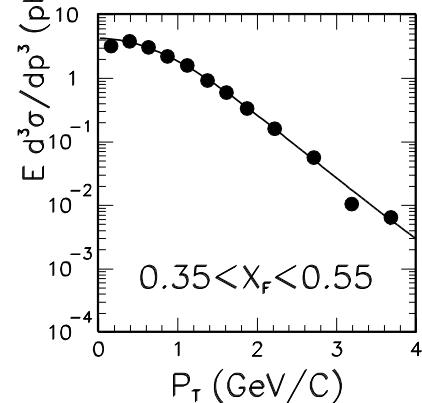
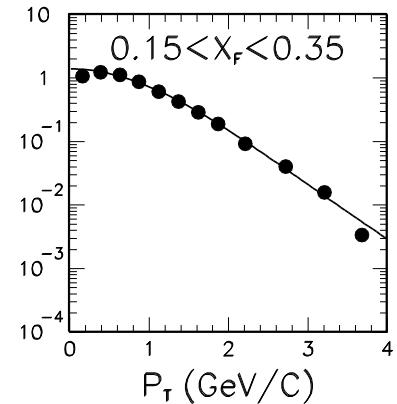
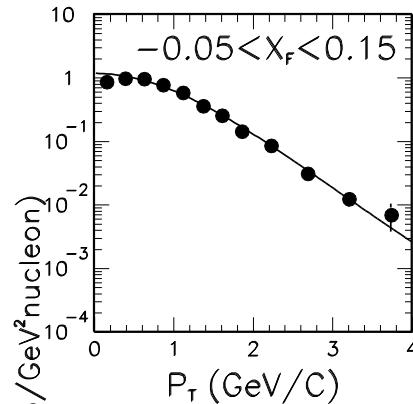
# Test of possible $x$ -dependent $k_T$ -distributions

E866 p+d D-Y data (800 GeV beam)

$5.2 < M < 6.2 \text{ GeV}$



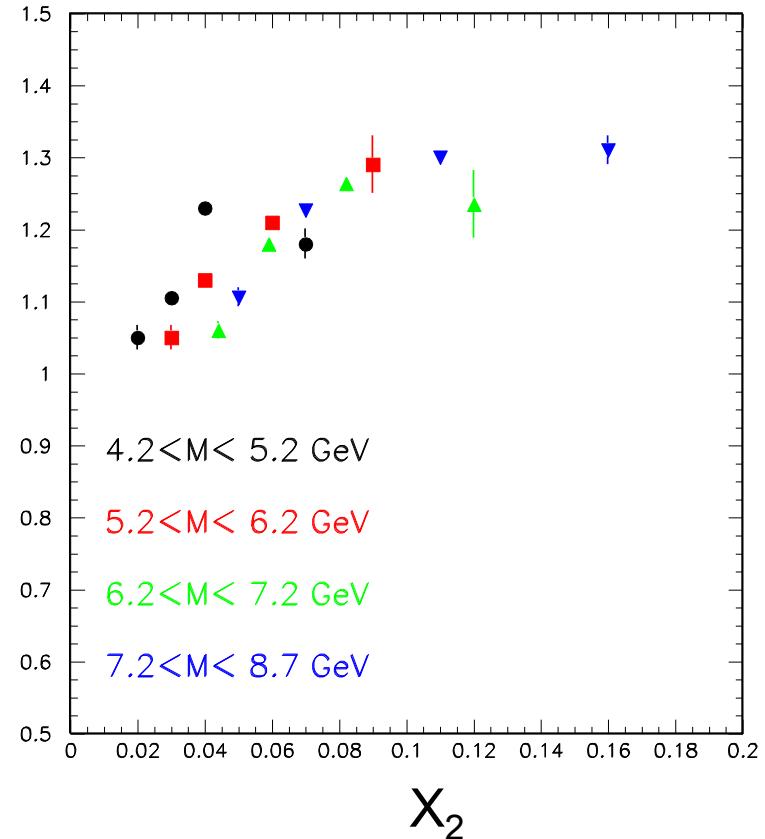
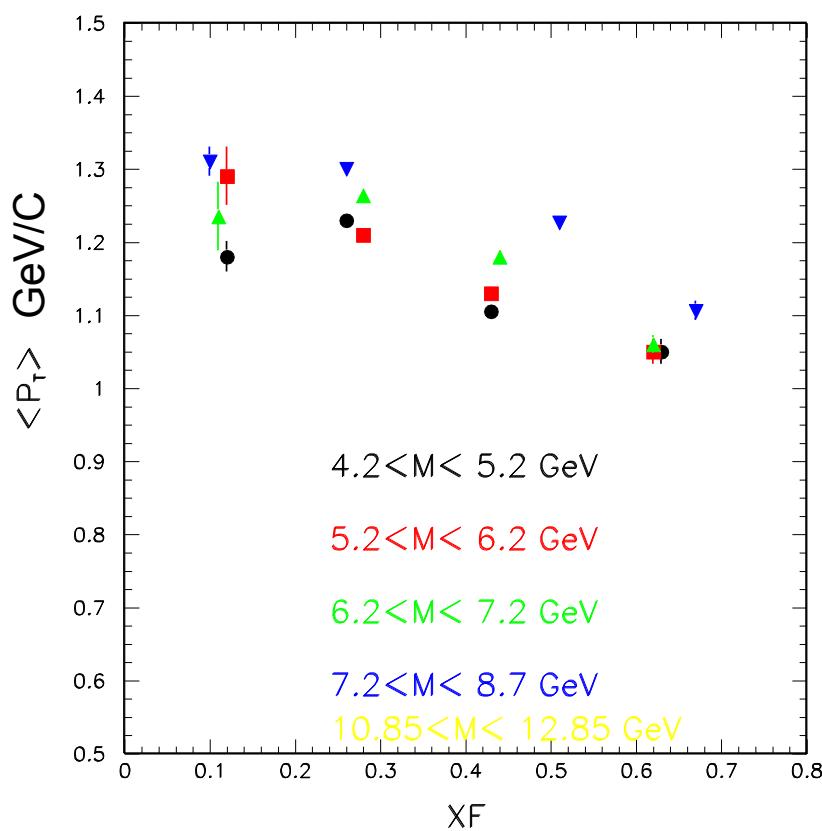
$7.2 < M < 8.7 \text{ GeV}$



Data from thesis of J. Webb

# Possible $x$ -dependent $k_T$ -distributions

E866 p+d D-Y data (800 GeV beam)



$\langle p_T \rangle$  scale with  $x_2$ ?

Analysis is ongoing (A. Ghalsasi, E. McClellan, JCP)

Would be very interesting to have new data at AFTER

# Summary

- The Drell-Yan process is a powerful experimental tool complimentary to the DIS for exploring quark structures in nucleons and nuclei.
- A rich physics program in Drell-Yan and quarkonium production can be pursued at AFTER.

There must be physics AFTER Higgs