

Moduli stabilization and mixed gauge-gravity mediation

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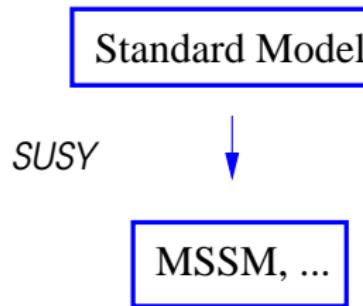
E. Dudas, Y. Mambrini, S. Pokorski, A.R.
JHEP 04 (2008) 015 [arXiv:0711.4934]

E. Dudas, Y. Mambrini, S. Pokorski, A.R. *work in progress*

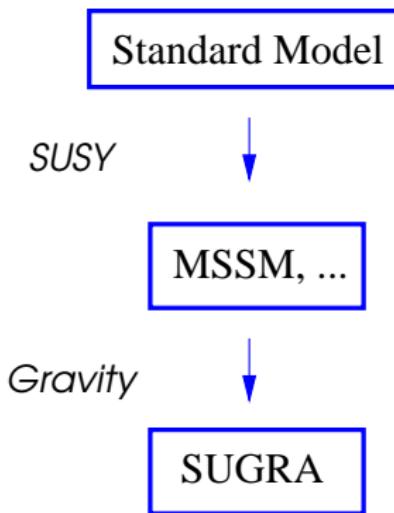
Motivations and Ingredients

Standard Model

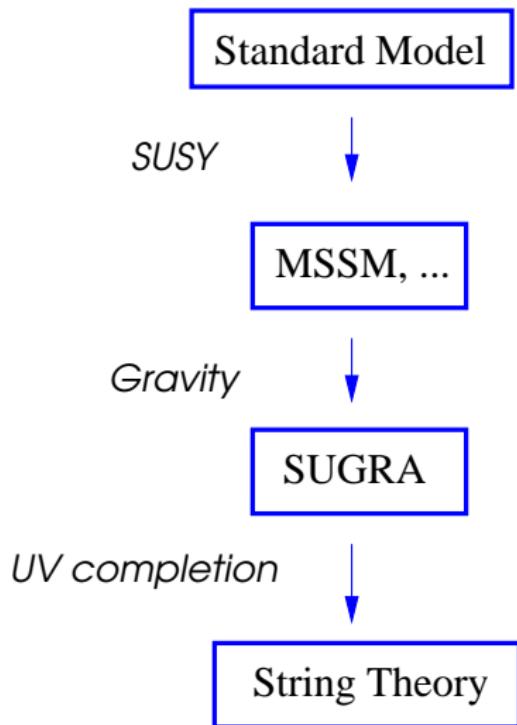
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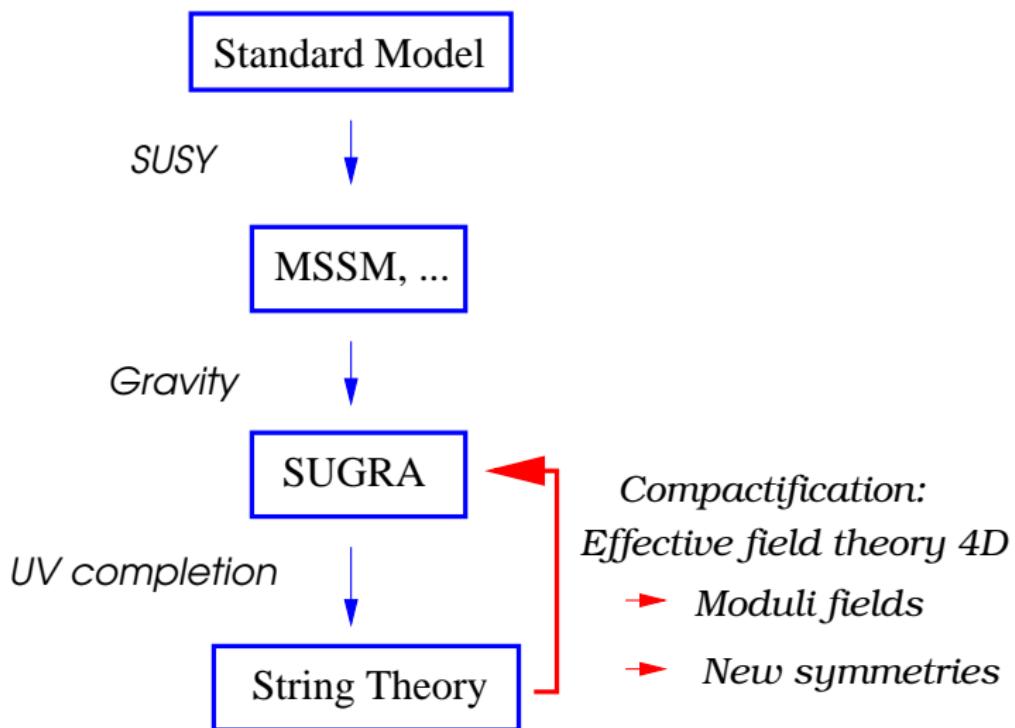
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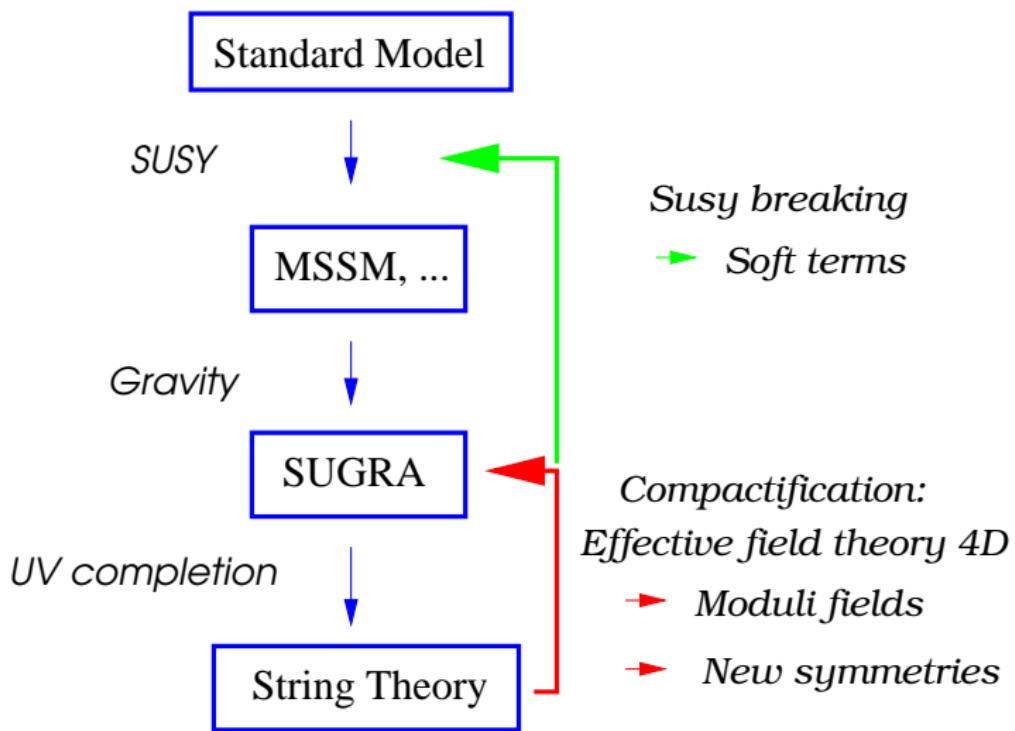
Motivations and Ingredients



Motivations and Ingredients



Motivations and Ingredients



Motivations and Ingredients

Moduli

- Fields parametrizing "flat directions": there's no potential for them
- Their vev's are related to physical parameters
→ they have to be taken into account and **stabilized**

New Symmetries

- New $U(1)_X$ gauge symmetries appear very naturally
- New charged fields Φ_i in the spectrum → interplay in the game

Moduli problem: Main Ideas of KKLT model

Kachru, Kallosh, Linde, Trivedi, hep-th/0301240

General Framework

Type IIB 10D compactified on orientifold \rightarrow SUGRA in 4D

Three logical "steps"

- 1- Turn on magnetic fields Giddings, Kachru, Polchinski, hep-th/0301240
 \hookrightarrow All moduli stabilized, except "T" (overall volume)
- 2- Provide nonperturbative potential for T
 \hookrightarrow AdS vacuum ($\Lambda_c < 0$)
- 3- Uplift the vacuum energy adding $\overline{D3}$ -branes [KKLT]
 \hookrightarrow Susy explicitly broken, not standard SUGRA

From the MSSM point of view

SUSY breaking in a new hidden sector

\hookrightarrow soft terms for the MSSM masses & couplings

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Other possibilities: stay in SUGRA & change step 3

Main ingredients for SUSY theories

- real **Kahler potential**: $K \rightarrow$ kinetic terms
- holomorphic **superpotential**: $W \rightarrow$ interaction terms

Auxiliary fields

- Scalar multiplets $\rightarrow F$
- Vector multiplets $\rightarrow D$

E.O.M.: F, D functions of the scalar fields depending on $K & W$

Scalar potential in SUGRA

$$V \sim \Sigma (|F - terms|^2 + |D - terms|^2) - 3|W|^2$$

SUSY preserved if $F_i = D_i = 0 \rightarrow \Lambda_c \sim -3|W|^2$

Other possibilities: stay in SUGRA & change step 3

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Scalar potential in SUGRA

$$V \sim \Sigma (|F - \text{terms}|^2 + |D - \text{terms}|^2) - 3|W|^2$$

Idea: F_i and/or $D_i \neq 0 \rightarrow$ uplift and spontaneous SUSY breaking

Adding an extra anomalous $U(1)_X$

T charged under $U(1)_X$

- it transforms as a Goldstone boson
- mass to $U(1)_X$ gauge boson
- coupling depends on T → anomalies → G.-S. mechanism
- gauge invariance + SUSY + other fields
 - ↪ Fayet-Iliopoulos term depending on T
 - ↪ constraints for terms in the superpotential

More in details

$$K = -3 \ln(T + \bar{T}) + |\Phi_+|^2 + |\Phi_-|^2 + K_{MSSM}$$

$$W = W_0 + e^{-bT} \Phi_-^q + m \Phi_+ \Phi_-$$

$$\hookrightarrow V = e^K (|F_T|^2 + |F_+|^2 + |F_-|^2 - 3|W|^2) + g_T^2 D^2$$

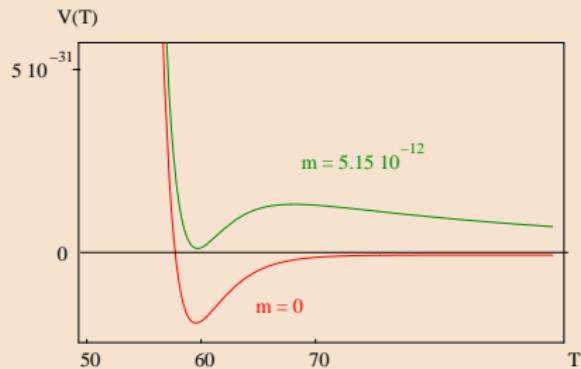
FI mechanism → spontaneous SUSY Breaking: $D, F_i \neq 0$

First Phenomenological Results

Uplift

Parameters in the superpotential fixed by some phenomenological requirements:

- $m_{3/2} \sim \frac{W}{M_p^2} \Rightarrow$ fix W_0
- Imposing $\Lambda_{cosm} = 0$
 \Rightarrow fix m



Soft terms by gravitational effects:

$$m_{1/2} \sim \frac{F_T}{T} \quad (\text{if } T \text{ couples to SM vector fields})$$
$$m_0 \sim m_{3/2}$$

Usually F_T is small....here:

$$m_{1/2} \lesssim m_0 \sim m_{3/2} \sim \text{TeV} \quad D \sim (100 \text{ TeV})^2$$

Anomalies & Gauge Messengers

If T couples to the SM gauge fields, since it is charged under $U(1)_X$, its shift generates mixed $U(1)_X - G_a^2$ anomalies

$$(G_a = \text{SU}(3), \text{SU}(2)_L, \text{U}(1)_Y)$$

$$\int d^2\theta \ T \ W_a^\alpha \ W_{\alpha a} \rightarrow Re[T] \ F_a^{\mu\nu} F_{\mu\nu a} + i \text{Im}[T] \ F_a^{\mu\nu} \tilde{F}_{\mu\nu a}$$

Possibilities :

- MSSM fields charged under $U(1)_X \Rightarrow$
problems: large masses for squarks and sleptons $\sim \sqrt{D}$,
tachyonic directions,...
- MSSM fields uncharged but additional fields M, \tilde{M}

M, \tilde{M} "Messengers" fields introduced in Gauge Mediated
Supersymmetry Breaking

Gauge Mediation & Soft masses

Giudice,Rattazzi hep-ph/9801271

GMSB: Messengers M, \tilde{M} couple to Susy Breaking hidden sector (superfield S)

$$\lambda \int d^4\theta S M \tilde{M}$$



Gauge Mediation & Soft masses

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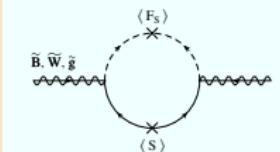
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→ mediation of Susy Breaking at perturbative level

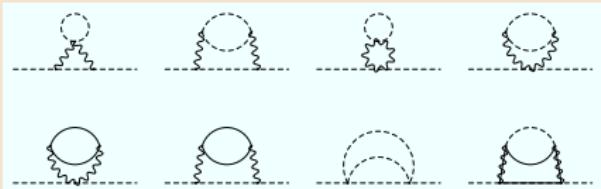
Gaugino masses

$$m_{1/2}^{GMSB} \sim \frac{g_a^2}{8\pi^2} \frac{F_S}{S}$$



Scalar masses

$$(m_0^{GMSB})^2 \sim \frac{g_a^4}{128\pi^4} \left(\frac{F_S}{S} \right)^2$$



Gauge Mediation & Uplift + D-term

Scalar soft masses

Poppitz, Trivedi, hep-ph/9703246

$$\text{STr } \mathcal{M}_{\text{Mess}}^2 = \text{Tr } m_{\text{scal}}^{\text{mess}}{}^2 - \text{Tr } m_{\text{ferm}}^{\text{mess}}{}^2 \neq 0$$

⇒ different GMSB result for scalars

In this model...

- Anomalies argument: M, \tilde{M} charged "+" under $U(1)_X$

↪ $\text{STr } \mathcal{M}_{\text{Mess}}^2 \sim D$ (D-term)

↪ Coupling with SUSY breaking sector: $\lambda \int d^4\theta \Phi_- M \tilde{M}$

True GMSB contribution for scalar soft masses

$$(m_0^{\text{GMSB}})^2 = \frac{g_a^4}{128\pi^4} \left\{ \left(\frac{F_-}{\Phi_-} \right)^2 + D \left[2 - \ln \left(\frac{\Lambda_{UV}^2}{m_{\text{Mess}}^2} \right) \right] \right\}$$

Mixed Mediation

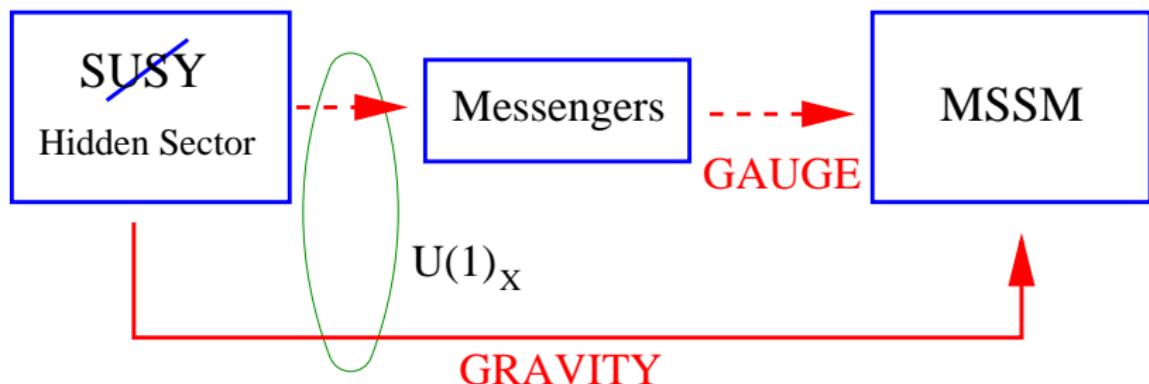
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Nontrivial interplay between $(m_0^2)_{\text{grav.}}$ and $(m_0^{GMSB})^2$:

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 - ↪ Here: $(m_0^{GMSB})^2 < 0$ and $|(m_0^{GMSB})^2| \sim m_{3/2}^2 \sim (\text{TeV})^2$

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- $(m_0^{GMSB})^2$ depends on couplings between scalars and gauge vectors

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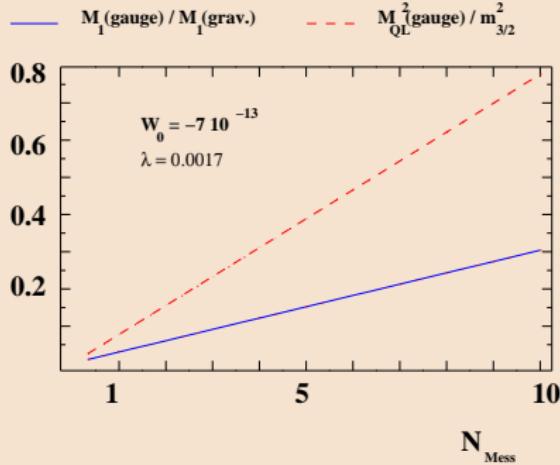
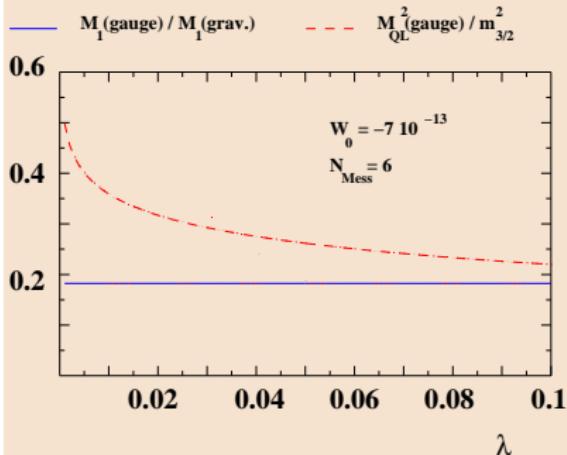
\Rightarrow Squarks can be lighter than sleptons at high energy
 \hookrightarrow measurable effects at low energy after RG flow

Gravity vs. Gauge

Playing with λ and N_{Mess}

$$(m_0^2) = (m_0^2)_{\text{grav.}} + N_{\text{Mess}} (m_0^{\text{GMSB}})^2$$

$$m_{1/2} = (m_{1/2})_{\text{grav.}} + N_{\text{Mess}} (m_{1/2}^{\text{GMSB}})$$



Some numbers

	A	B
W_0	$-7 \cdot 10^{-13}$	$-4.3 \cdot 10^{-13}$
m	$7.3 \cdot 10^{-12}$	$4.5 \cdot 10^{-12}$
a	1	1
b	0.3	0.5
q	1	1
$\tan \beta$	30	15
t	98.3	59.4



μ (GeV)	810	1070
$B\mu$ (GeV) 2	(400) 2	(870) 2
$m_{\chi_1^0}$	110	140
$m_{\chi_1^+}$	220	290
$m_{\tilde{g}}$	760	950
m_h	120	120
m_A	2220	3290
$m_{\tilde{t}_1}$	1380	1770
$m_{\tilde{t}_2}$	1920	2610
$m_{\tilde{c}_1}, m_{\tilde{u}_1}$	2580	3300
$m_{\tilde{b}_1}$	1910	2610
$m_{\tilde{b}_2}$	2310	3230
$m_{\tilde{s}_1}, m_{\tilde{d}_1}$	2580	3300
$m_{\tilde{\tau}_1}$	2290	3200
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$m_{\tilde{\mu}_1}, m_{\tilde{e}_1}$	2550	3270

	A+GMSB	B+GMSB
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$\tan \beta$	30	15
t	97.3	59.4
λ	$1.7 \cdot 10^{-3}$	$1.1 \cdot 10^{-3}$
N_{Mess}	6	6
μ (GeV)	186	216
$B\mu$ (GeV) 2	(330) 2	(730) 2
$m_{\chi_1^0}$	120	150
$m_{\chi_1^+}$	160	200
$m_{\tilde{g}}$	850	1060
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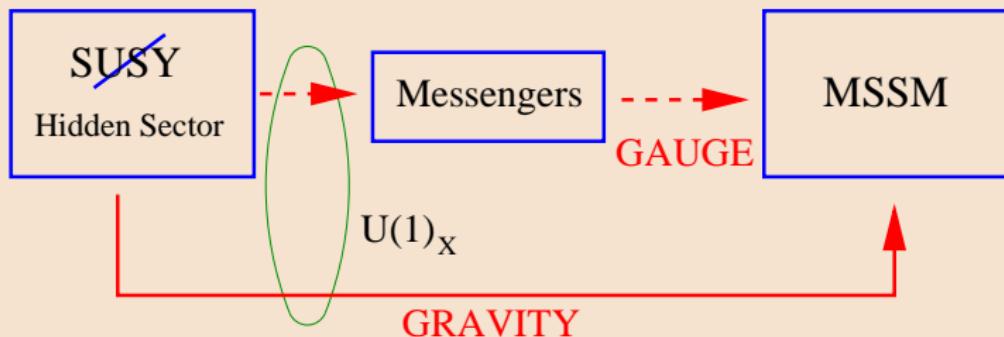
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Hybrid Model

Gauge - Gravity mixed mediation

Idea: Generalize the previous stringy motivated example



Generic scenario with **both** gauge and gravity contributions to the soft terms at the GUT scale

→ Suitable **parametrization** and **scan** on the parameters space

Parametrization

Soft terms for the masses

$$(m_0^2)_i = M_S^2 \left[1 + N_{\text{mess}} \tilde{\alpha}_i (1 - \tilde{\beta}) \right]$$

$$(m_{1/2})_a = f M_S$$

Parametrization

Soft terms for the masses

$$(m_0^2)_i = \textcolor{red}{M_S^2} [1 + N_{\text{mess}} \tilde{\alpha}_i (1 - \tilde{\beta})]$$

$$(m_{1/2})_a = f \textcolor{red}{M_S}$$

- M_S → scale of the **gravity** mediation

Parametrization

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- $\tilde{\beta}$ → effects of $\text{STr } \mathcal{M}_{\text{Mess}}^2 \neq 0$
- f → both gravity & gauge contributions universal

Parametrization... $\tilde{\alpha}_i$

$$m_{Q_L}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[\frac{4}{3}n_q\alpha_3^2 + \frac{3}{4}n_l\alpha_2^2 + \frac{3}{5}\left(\frac{1}{6}\right)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$m_{U_R}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[\frac{4}{3}n_q\alpha_3^2 + \frac{3}{5}\left(\frac{-2}{3}\right)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$m_{D_R}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[\frac{4}{3}n_q\alpha_3^2 + \frac{3}{5}\left(\frac{1}{3}\right)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$m_{E_L}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[+\frac{3}{4}n_l\alpha_2^2 + \frac{3}{5}\left(\frac{-1}{2}\right)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$m_{E_R}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[\frac{3}{5}(1)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$m_{H_u}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[\frac{3}{4}n_l\alpha_2^2 + \frac{3}{5}\left(\frac{1}{2}\right)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$m_{H_d}^2 = M_S^2 + (1 - \tilde{\beta})\tilde{\alpha} \left[\frac{3}{4}n_l\alpha_2^2 + \frac{3}{5}\left(\frac{-1}{2}\right)^2 \frac{2n_q + 3n_l}{5}\alpha_1^2 \right] M_S^2,$$

$$(m_{1/2})_a = M_S f = M_{1/2}$$

Sample spectra

SUSPECT + micrOMEGAs2.0

	A	B	C
M_S	1.3 TeV	1.3 TeV	1.3 TeV
$M_{1/2}$	500	500	500
$\tilde{\alpha}$	0	1	1
$\tilde{\beta}$	0	0	1.85
$\tan \beta$	5	5	5
μ	750	1070	258
M_1	212	213	209
$m_{\chi_1^0}$	211	211	186
$m_{\chi_1^+}$	400	408	430
$m_{\tilde{g}}$	1222	1251	1175
m_h	113.2	114.1	112.4
m_A	1530	1920	1095
$m_{\tilde{t}_1}$	1066	1260	844
$m_{\tilde{t}_2}$	1403	1780	1008
$m_{\tilde{c}_1}, m_{\tilde{u}_1}$	1612	2052	1119
$m_{\tilde{b}_1}$	1388	1771	954
$m_{\tilde{b}_2}$	1588	1891	1282
$m_{\tilde{s}_1}, m_{\tilde{d}_1}$	1613	2053	1122
$m_{\tilde{\tau}_1}$	1275	1448	1046
$m_{\tilde{\tau}_2}$	1302	1550	1106
$m_{\tilde{\mu}_1}, m_{\tilde{e}_1}$	1303	1551	1047
Ωh^2	15.1	29.4	0.105

Sample spectra

SUSPECT + micrOMEGAs2.0

	A	B	C
M_S	1.3 TeV	1.3 TeV	1.3 TeV
$M_{1/2}$	500	500	500
$\tilde{\alpha}$	0	1	1
$\tilde{\beta}$	0	0	1.85
$\tan \beta$	5	5	5
μ	750	1070	258
M_1	212	213	209
$m_{\chi_1^0}$	211	211	186
$m_{\chi_1^+}$	400	408	430
$m_{\tilde{g}}$	1222	1251	1175
m_h	113.2	114.1	112.4
m_A	1530	1920	1095
$m_{\tilde{t}_1}$	1066	1260	844
$m_{\tilde{t}_2}$	1403	1780	1008
$m_{\tilde{c}_1}, m_{\tilde{u}_1}$	1612	2052	1119
$m_{\tilde{b}_1}$	1388	1771	954
$m_{\tilde{b}_2}$	1588	1891	1282
$m_{\tilde{s}_1}, m_{\tilde{d}_1}$	1613	2053	1122
$m_{\tilde{\tau}_1}$	1275	1448	1046
$m_{\tilde{\tau}_2}$	1302	1550	1106
$m_{\tilde{\mu}_1}, m_{\tilde{e}_1}$	1303	1551	1047
Ωh^2	15.1	29.4	0.105

- Squarks more sensitive than sleptons to GMSB negative contribution

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$$\mu^2 \approx -m_{H_2}^2 - \frac{1}{2} M_Z^2$$

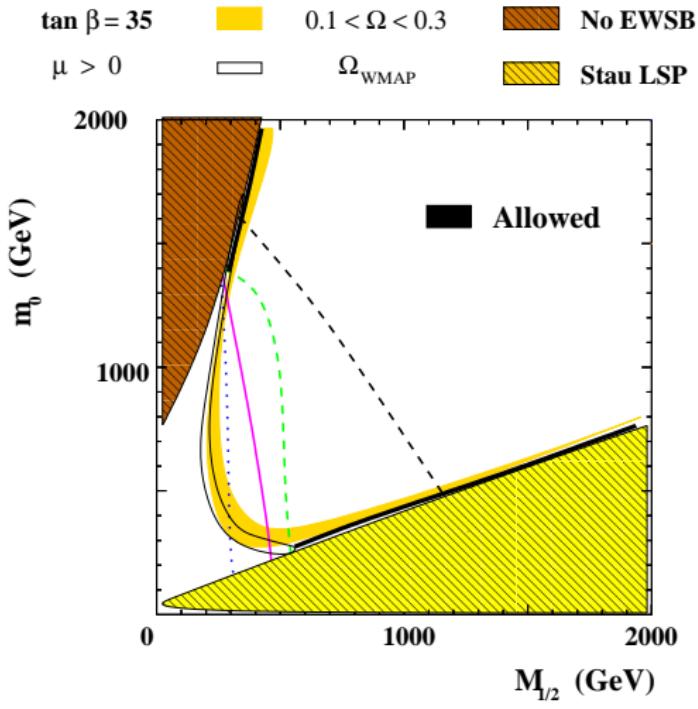
$$\frac{\partial m_{H_2}^2}{\partial \log \mu} \approx 6y_t^2(m_{H_2}^2 + m_{U_3}^2 + m_{Q_3}^2 + A_t^2)$$

Initial conditions + RG flow effects:

- Lighter squarks
- lower $|m_{H_2}^2|$ value
- lower μ value
- Lightest neutralino is generally a mixed bino-higgsino state
- Suitable dark matter relic density value allowed

Examples of scan

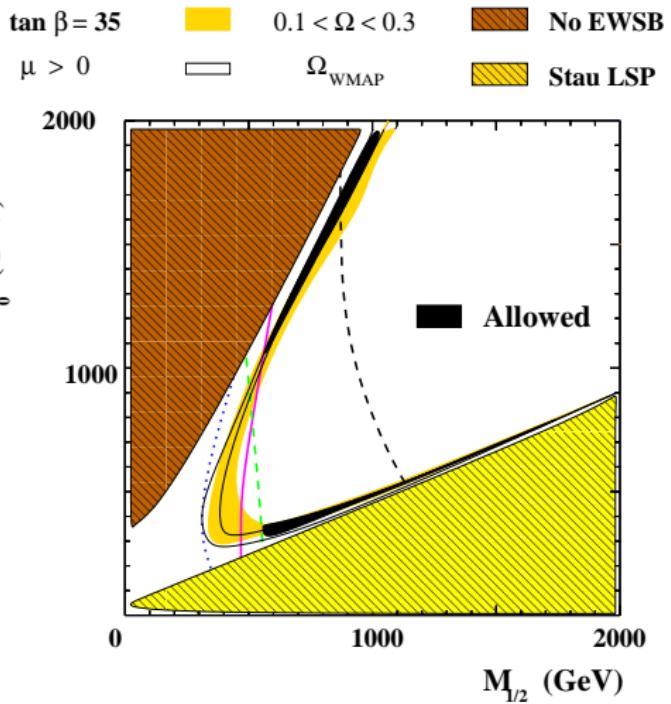
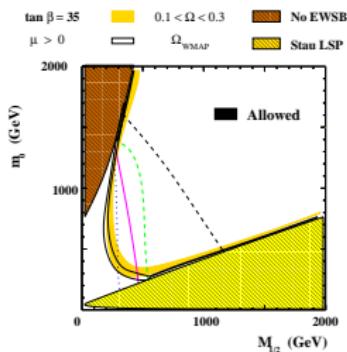
$$N_{mess} = 5 \quad \tilde{\beta} = 5$$



$$\tilde{\alpha} = 0 \quad (\text{mSUGRA})$$

Examples of scan

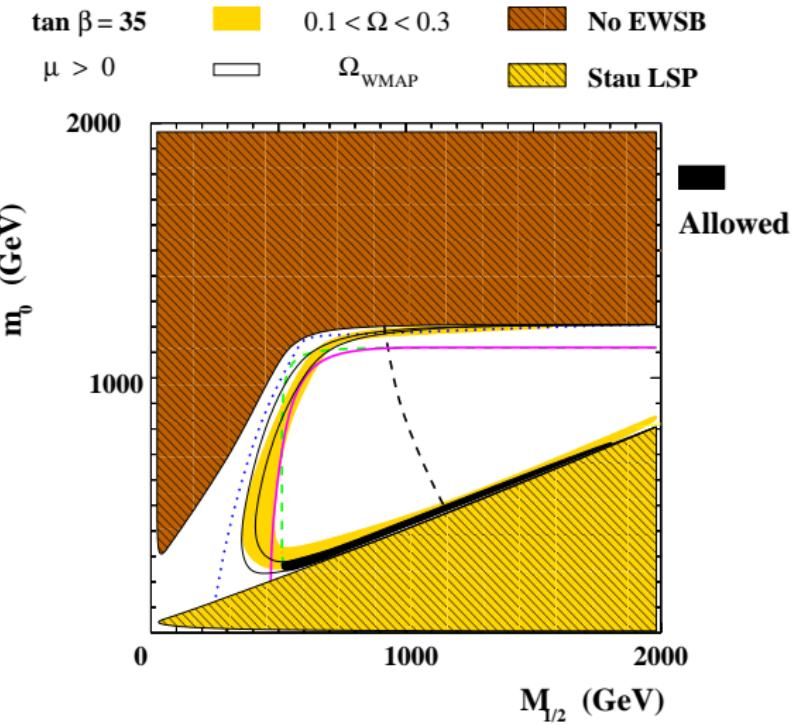
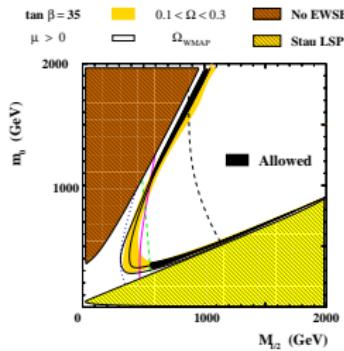
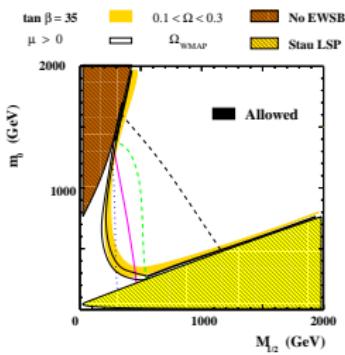
$$N_{mess} = 5 \quad \tilde{\beta} = 5$$



$$\tilde{\alpha} = 1$$

Examples of scan

$$N_{mess} = 5 \quad \tilde{\beta} = 5$$



$$\tilde{\alpha} = 1.1$$

Conclusions & Outlook

Results

- High energy:
 - Example of a **non-decoupled** uplift: modulus contribution F_T to SUSY Breaking is not negligible
 - Natural mixing between **gauge & gravity** mediation mechanisms
- Low energy:
 - Gaugino masses **comparable** to the scalars and gravitino masses
 - **Spectrum "compressed"**: squarks more sensitive to GMSB

Work in progress

- Top-down: Can F_T be increased?
- Bottom-up: Systematic study of the phenomenology of the **hybrid** model