





### Hadron electromagnetic form factors in Time-like region

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### Time-like electromagnetic processes of hadrons



- $\Gamma_{\mu}$  is parametrized in terms of 2S+1 FFs.
- Optical theorem requires imaginary part for  $q^2 > 4m_\pi^2$ 
  - Spin 0  $(\pi, K) \rightarrow$  one FF .
  - Spin 1/2 (nucleons)  $\rightarrow$  2 FFs :  $G_E$  and  $G_M$ .
    - $\bullet\,$  Differential cross section  $\rightarrow\,$  modulus of 2 FFs.
    - $\bullet\,$  polarization observables  $\rightarrow\,$  relative phase.
  - Spin 1 (Deuteron,  $\rho meson$ )  $\rightarrow$  3 FFs :  $G_C$ ,  $G_M$  and  $G_Q$ .

#### Time-like Form Factors : Data

• Proton  $(|G_M|; |G_E| = |G_M| \text{ or } |G_E| = 0)$ :

$$e^+ + e^- \leftrightarrow ar{p} + p$$



• Neutron  $(|G_M|)$ 

• 4 points measured by FENICE experiment (  $e^+ + e^- 
ightarrow ar{n} + n)$ 

- Deuteron and  $\rho$ -meson
  - no data in TL region.

#### Time-like Form Factors : Models



FFs are analytical functions of q<sup>2</sup>

- Most of FFs models are constructed to describe and Fit SL data .
- Time-like FFs ( $q^2$  positive) can be obtained by analytic continuation from the Space-like FFs ( $q^2$  negative)
- At high  $q^2$ , pQCD gives predictions for  $q^2$ -dependence.

 $\begin{array}{c} \mathbf{pQCD \ predictions \ at \ high \ q^2} \\ \mathrm{VMD \ parametrization \ of \ TL \ FFs} \\ \mathrm{Constraint \ on \ the \ VMD \ spin \ one \ FFs \ from \ BaBar} \end{array}$ 

### pQCD parametrization of proton FFs

# Space-like region $(Q^2 = -q^2)$



• quark counting rules :

$$G_{E,M}(Q^2)\sim 1/(Q^2)^2$$

• pQCD :

$$\mathcal{G}_{E,M}(Q^2) \sim rac{lpha_{s}^{2}(Q^2)}{(Q^2)^2}$$

• Asymptotic freedom, beta function :

$$\alpha_s = \frac{1}{b \ln(Q^2/\Lambda^2)}, \ b = \frac{11 - 2/3N_f}{4\pi}, \ Q^2 > \Lambda^2.$$

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pQCD parametrization of proton FFs

TIME-LIKE REGION

• Analytical extension :  $\ln(Q^2/\Lambda^2) \rightarrow \ln(q^2/\Lambda^2) - i\pi, \ q^2 > \Lambda^2.$ 

$$G_{E,M} = \frac{A}{(q^2)^2 [\ln(q^2/\Lambda^2) - i\pi]^2}, \ \Lambda = 0.3 \ GeV, \ A = 96.2 \ GeV^4$$
$$G_E(4M_p^2) = G_M(4M_p^2)$$



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Vector Meson Dominance formalism

- VMD model describes interaction of photon with hadron ( $\gamma HH$  vertex) through exchange of vectors mesons ( $\omega$ ,  $\phi$ ,  $\rho$ ).
- VMD parametrization:

$$F(Q^2) = \Sigma_\nu A_\nu \frac{m_\nu^2}{m_\nu^2 + Q^2}$$

- small number of parameters with physical meaning (masses, coupling constants,...).
- can be extended to TL region.

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## Spin 1/2

### Protons and neutrons

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### VMD parametrization of spin 1/2 FFs

### Space like

### lachello, Jackson and Lande (Physics letters 43 B 1973)

$$\begin{split} F_1^s(Q^2) \ &= \ \frac{g(Q^2)}{2} \left[ (1 - \beta_\omega - \beta_\phi) + \beta_\omega \frac{\mu_\omega^2}{\mu_\omega^2 + Q^2} + \beta_\phi \frac{\mu_\phi^2}{\mu_\phi^2 + Q^2} \right], \\ F_1^v(Q^2) \ &= \ \frac{g(Q^2)}{2} \left[ (1 - \beta_\rho) + \beta_\rho \frac{\mu_\rho^2 + 8\Gamma_\rho \mu_\pi / \pi}{(\mu_\rho^2 + Q^2) + (4\mu_\pi^2 + Q^2)\Gamma_\rho \alpha(Q^2) / \mu_\pi} \right], \\ F_2^s(Q^2) \ &= \ \frac{g(Q^2)}{2} \left[ (\mu_p + \mu_n - 1 - \alpha_\phi) \frac{\mu_\omega^2}{\mu_\omega^2 + Q^2} + \alpha_\phi \frac{\mu_\phi^2}{\mu_\phi^2 + Q^2} \right], \\ F_2^v(Q^2) \ &= \ \frac{g(Q^2)}{2} \left[ (\mu_p - \mu_n - 1) \frac{\mu_\rho^2 + 8\Gamma_\rho \mu_\pi / \pi}{(\mu_\rho^2 + Q^2) + (4\mu_\pi^2 + Q^2)\Gamma_\rho \alpha(Q^2) / \mu_\pi} \right]. \end{split}$$



### Time like

### lachello and Wan PRC 69 (2004)

$$g(q^2) = \frac{1}{(1 - \gamma e^{i\theta} q^2)^2} \qquad \qquad \frac{m_\rho^2}{m_\rho^2 - q^2} \to \frac{m_\rho^2 + 8\Gamma_\rho m_\pi / \pi}{m_\rho^2 - q^2 + (4m_\pi^2 - q^2)\Gamma_\rho \alpha(q^2)/m_\pi + i\Gamma_\rho 4m_\pi \beta(q^2)}$$

Hadron electromagnetic form factors in Time-like region

### VMD parametrization of spin 1/2 FFs

### Good agreement with both proton and neutron TL data

lachello and Wan, PRC 69 (2004)



## Spin 1

### Deuteron

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VMD parametrization of deuteron FFs

C. Adamuscin, G.I. Gakh and E.T-G. PRC 73, 045204 (2006)

• One hard core with surrounding meson cloud.

$$G_i(Q^2) = N_i g_i(Q^2) F_i(Q^2), \quad i = c, q, m$$

• Isoscalar vector mesons ( $\omega$ ,  $\phi$ )

$$F_i(Q^2) = 1 - \alpha_i - \beta_i + \alpha_i \frac{m_\omega^2}{m_\omega^2 + Q^2} + \beta_i \frac{m_\phi^2}{m_\phi^2 + Q^2}$$

Intrinsic structure

$$g_i(Q^2) = 1/[1+\gamma_iQ^2]^{\delta_i}$$

• *N<sub>i</sub>*'s : correspond to the normalisation static properties:

$$G_C(0) = 1, G_Q(0) = M_d^2 Q_d^2 = 25.83, G_M(0) = 1.714$$



VMD parametrization of deuteron FFs

C. Adamuscin, G.I. Gakh and E.T-G. PRC **73**, 045204 (2006)



Due to the quadratic dependence of FFs on the cross section and polarization observables in ed elastic scattering, there are 2 sets of experimental points for  $G_C$  and  $G_Q$ .

## Spin 1

### $\rho$ meson

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### Model for $\rho$ meson FFs

# C. Adamuscin, G.I. Gakh and E.T-G. Phys.Rev. C75 065202 (2007)

### Space-like region

#### **Time-like region**

$$G_{C}(q^{2}) = \frac{G_{C}(0)(A + Bq^{2})m_{C}^{2}}{(m_{C}^{2} - q^{2})^{2}} \qquad G_{C}(q^{2}) = \frac{G_{C}(0)(A + Bq^{2})m_{C}^{2}}{(m_{C}^{2} - q^{2} - im_{C}\Gamma_{C})^{2}}$$

$$G_{M}(q^{2}) = \frac{G_{M}(0)m_{M}^{4}}{(m_{M}^{2} - q^{2})^{2}} \qquad G_{M}(q^{2}) = \frac{G_{M}(0)m_{M}^{4}}{(m_{M}^{2} - q^{2} - im_{M}\Gamma_{M})^{2}}$$

$$G_{Q}(q^{2}) = \frac{G_{Q}(0)m_{M}^{4}}{(m_{Q}^{2} - q^{2})^{2}} \qquad G_{Q}(q^{2}) = \frac{G_{Q}(0)m_{M}^{4}}{(m_{Q}^{2} - q^{2} - im_{Q}\Gamma_{Q})^{2}}$$

- Parameters (A, B, m<sub>i</sub>) are fixed based on theoretical (covariant and light front) calculation in SL. PRC 55, 2043 (1997)
- A=1, B=0.33 from the node of  $G_C(q^2 = -3GeV^2)$

Helicity amplitudes measurements by BaBar

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PHYSICAL REVIEW D 78, 071103(R) (2008) Observation of  $e^+e^- \rightarrow \rho^+\rho^-$  near  $\sqrt{s} = 10.58$  GeV

$$egin{aligned} |F_{00}|^2 &: |F_{10}|^2 &: |F_{11}|^2 = 0.51 \pm 0.14(\textit{stat}) \pm 0.07(\textit{syst}) \ 0.10 \pm 0.04(\textit{stat}) \pm 0.01(\textit{syst}) \ 0.04 \pm 0.03(\textit{stat}) \pm 0.01(\textit{syst}) \end{aligned}$$

$$|F_{00}|^2 + 4|F_{10}|^2 + 2|F_{11}|^2 = 0.01$$

To fit the data, we need a model for the 3 FFs

Helicity amplitudes of  $\gamma^* \rightarrow \rho^+ \rho^-$ 

• Matrix element :

$$M = \epsilon \cdot (p_1 - p_2) [-G_1(q^2) U_1^* \cdot U_2^* + \frac{G_3(q^2)}{M^2} (U_1^* \cdot q U_2^* \cdot q - \frac{q^2}{2} U_1^* \cdot U_2^*)] \\ -G_2(q^2) (\epsilon \cdot U_1^* U_2^* \cdot q - \epsilon \cdot U_2^* U_1^* \cdot q),$$

 $\epsilon, \mathit{U}_1, \mathit{U}_2$  are the polarization vectors of  $\gamma^*, 
ho^+, 
ho^-$ 

helicity amplitudes:

$$F_{\lambda_1\lambda_2} = M_{\lambda_1\lambda_2}^{\lambda} = M(\epsilon \to \epsilon^{(\lambda)}, U_1 \to U_1^{(\lambda_1)}, U_2 \to U_2^{(\lambda_2)}), \ \lambda = \lambda_1 - \lambda_2$$

- spin one photon :  $F_{1-1} = F_{-11} = 0$
- symmetry properties :  $F_{-1-1} = F_{11}$ ,  $F_{10} = F_{01} = F_{-10} = F_{0-1}$

$$F_{00}, F_{10}, F_{11}$$

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Helicity amplitudes of  $\gamma^* \rightarrow \rho^+ \rho^-$ 

Helicity amplitudes are combination of the three complex FFs

$$\begin{split} F_{00} &= -\frac{\sqrt{q^2 - 4M^2}}{2M^2} [q^2(G_1 + G_2 + G_3) - 2M^2G_1], \\ F_{11} &= \sqrt{q^2 - 4M^2}(G_1 + 2\tau G_3), \\ F_{10} &= -\sqrt{\tau(q^2 - 4M^2)}G_2, \quad \tau = \frac{q^2}{4M^2}. \end{split}$$

Total cross section:

$$\sigma(e^+e^- o 
ho^+
ho^-) = rac{\pilpha^2eta^3}{3q^2}rac{1}{4M^2( au-1)}(|F_{00}|^2+4|F_{10}|^2+2|F_{11}|^2)$$

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pQCD predictions at high  $q^2$  VMD parametrization of TL FFs Constraint on the VMD spin one FFs from BaBar

### Fit the data

Ref.	$m_C$ (GeV)	$m_M$ (GeV)	$m_Q$ (GeV)
[3]	$1.34 \pm 2$	$1.42 \pm 0.5$	$1.51 \pm 0.1$
This work (I)	$1.05_{-0.09}^{+0.05}$	$1.28^{+0.06}_{-0.08}$	$0.97_{-0.01}^{+0.02}$
This work (II)	$0.77^{+0.05}_{-0.02}$	$1.28^{+0.06}_{-0.08}$	$1.12^{+0.05}_{-0.08}$



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Absolute value of  $\rho$ -meson FFs:  $|G_M|$ ,  $|G_C|$  and  $|G_Q|$ 





• Models based on vector meson dominance give useful parametrization in SL and TL regions for spin 1/2 hadrons.

• The VMD parametrization can reproduce the unique measurement on the spin one TL FFs given in terms of helicity amplitudes.

• PANDA experiment will provide a powerful test of the existent models of the nucleon FFs in a wide *q*<sup>2</sup> range.