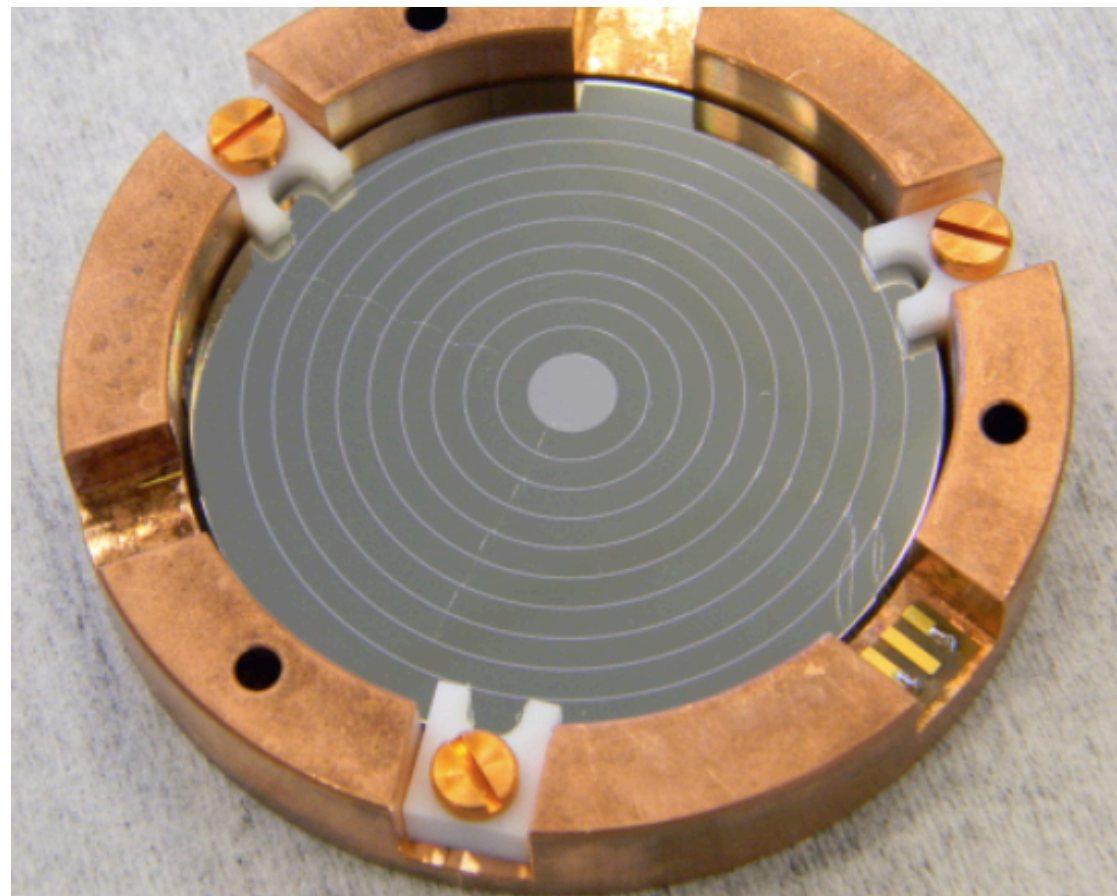


Neganov-Luke assisted Light Sensors

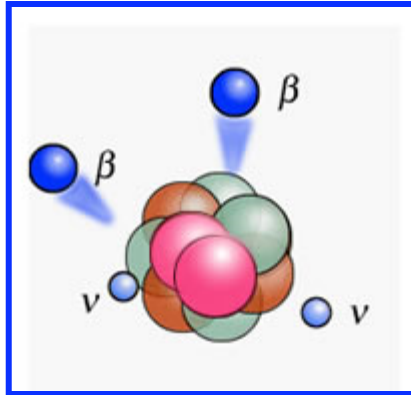
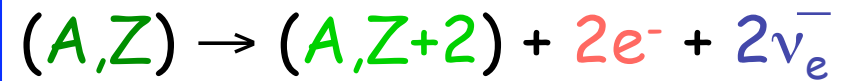


Emiliano Olivieri for the Detector R&D Group at
CSNSM Laboratory (Orsay Campus)

- ▶ Physics background and motivations: the double beta decay ($\beta\beta$)
- ▶ A bolometer in a nutshell...
- ▶ The Neganov-Luke Assisted light detector
 - Neganov-Trofimov-Luke effect
 - LukS working regimes
 - The first two prototypes
- ▶ Experimental set-up for a complete characterization
- ▶ Results and Perspectives
- ▶ Conclusions

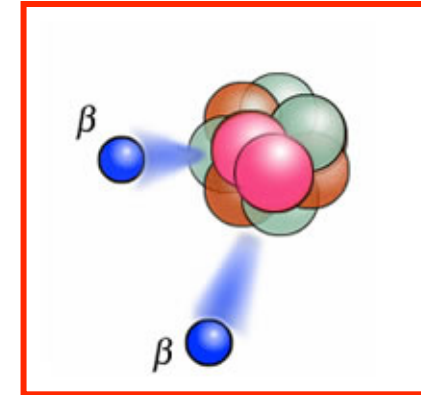
Rare event search: $0\nu\text{-}\beta\beta$ decay

Double Beta Decay ($2\nu\text{-}\beta\beta$)



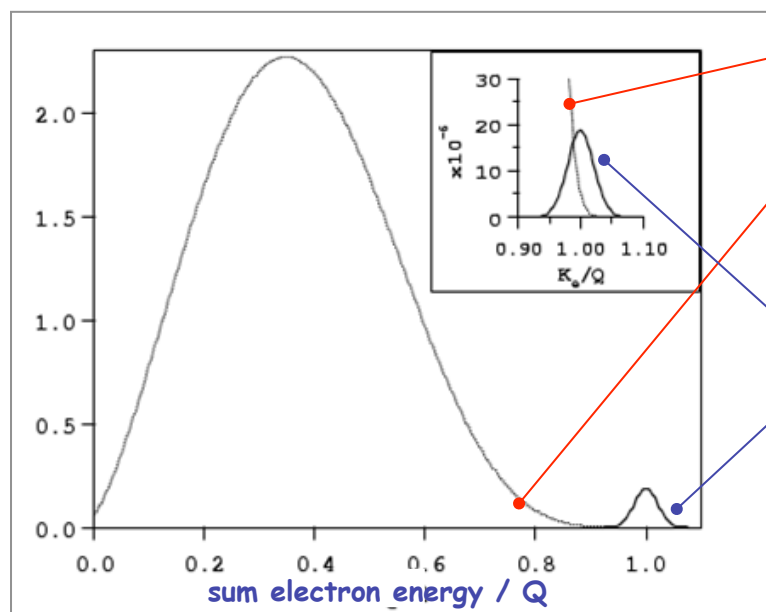
$$[T_{1/2}(2\nu)]^{-1} = G_{2\nu}(Q, Z) |M_{2\nu}|^2$$

Neutrino-less Double Beta Decay ($0\nu\text{-}\beta\beta$)



$$[T_{1/2}(0\nu)]^{-1} = G_{0\nu}(Q, Z) |M_{0\nu}|^2 m_{\beta\beta}^2$$

The **shape** of the **two electron sum energy spectrum** enables to distinguish among the two different discussed decay modes



2ν double beta decay continuum with maximum at $\sim 1/3 Q$

0ν double beta decay peak enlarged only by the detector energy resolution

↓

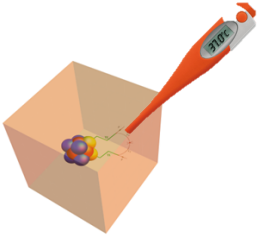
neutrinoless Double Beta Decay rate Phase space Nuclear matrix elements Effective Majorana mass

$$\frac{1}{\tau} = G(Q, Z) \cdot |M|^2 \cdot \langle m_{\beta\beta} \rangle^2$$

Absolute Neutrino Mass Scale

The $0\nu\text{-}\beta\beta$ candidates

Nucleus	I. A.	Q-value [keV]	Materials successfully tested as bolometers in crystalline form
^{76}Ge	7.8	2039	Ge
^{136}Xe	8.9	2479	NONE
^{130}Te	33.8	2527	TeO_2
^{116}Cd	7.5	2802	CdWO_4 , CdMoO_4
^{82}Se	9.2	2995	ZnSe
^{100}Mo	9.6	3034	PbMoO_4 , CaMoO_4 , SrMoO_4 , CdMoO_4 , ZnMoO_4 , Li_2MoO_4 , MgMoO_4
^{96}Zr	2.8	3350	ZrO_2
^{150}Nd	5.6	3367	NONE → many attempts
^{48}Ca	0.187	4270	CaF_2 , CaMoO_4



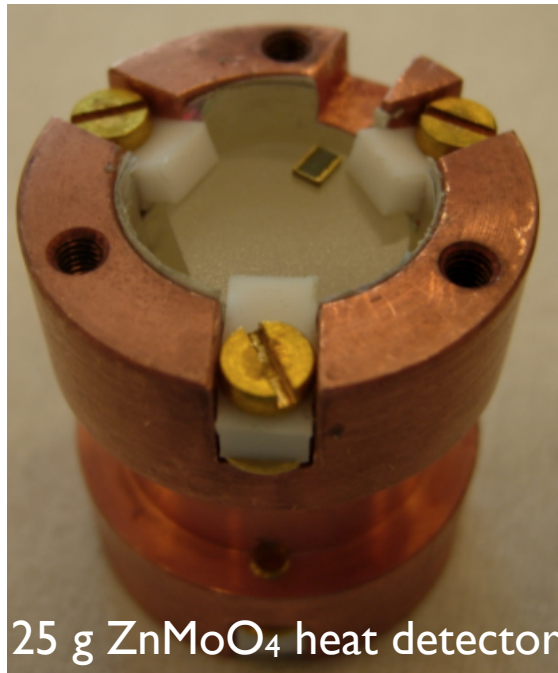
→ **Cuoricino, CUORE**
↓ **LUMINEU**
Orsay, Kiev, Novosibirsk, Como

Mass sensitivity: $\langle m_{\beta\beta} \rangle \propto \left(\frac{b\Delta E}{Mt_{\text{live}}} \right)^{\frac{1}{4}}$

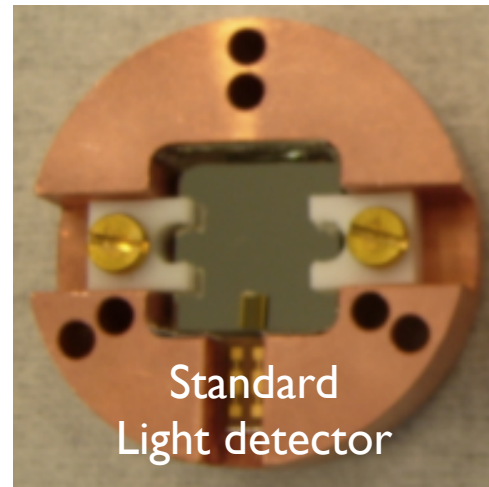
ΔE =energy res.
 b =bkg in counts/keV/kg/day
 M =effective mass
 t_{live} =life time

Next future $0\nu\text{-}\beta\beta$ experiment...

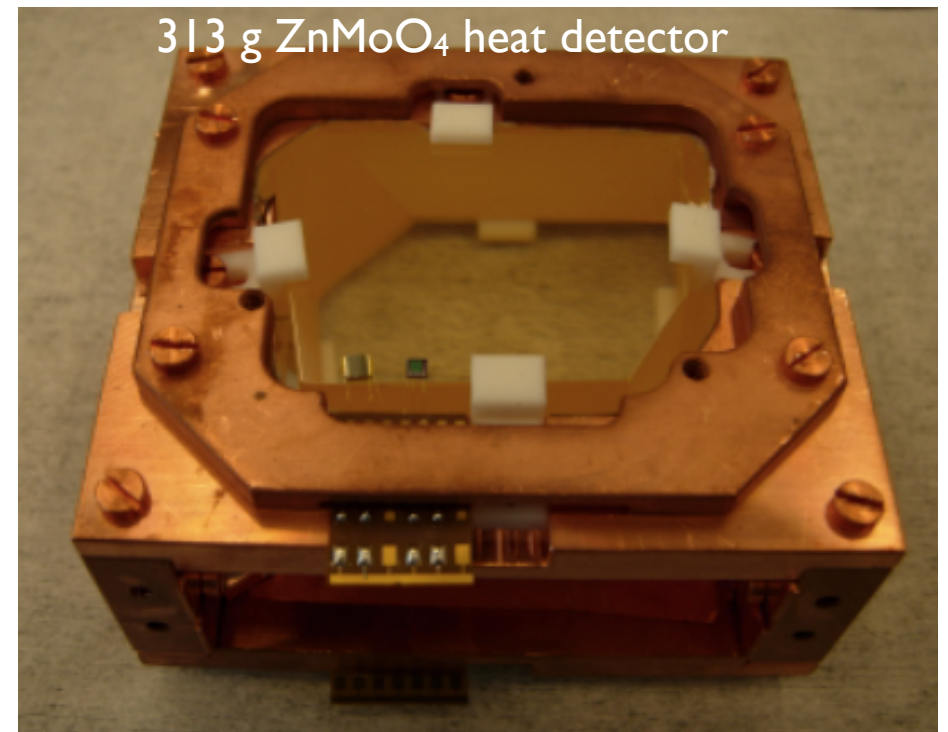
LUMINEU prototypes



25 g ZnMoO₄ heat detector



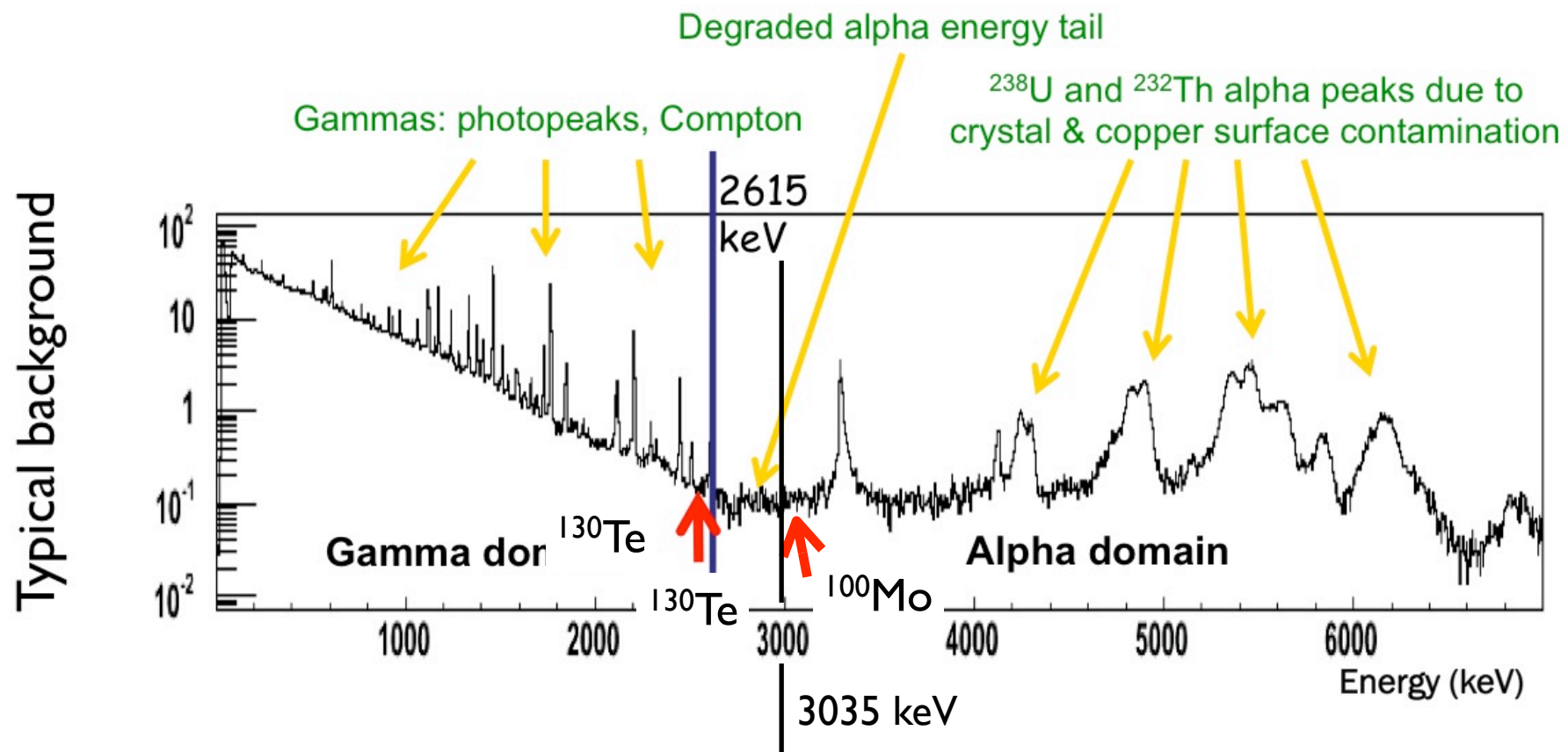
Standard
Light detector



313 g ZnMoO₄ heat detector

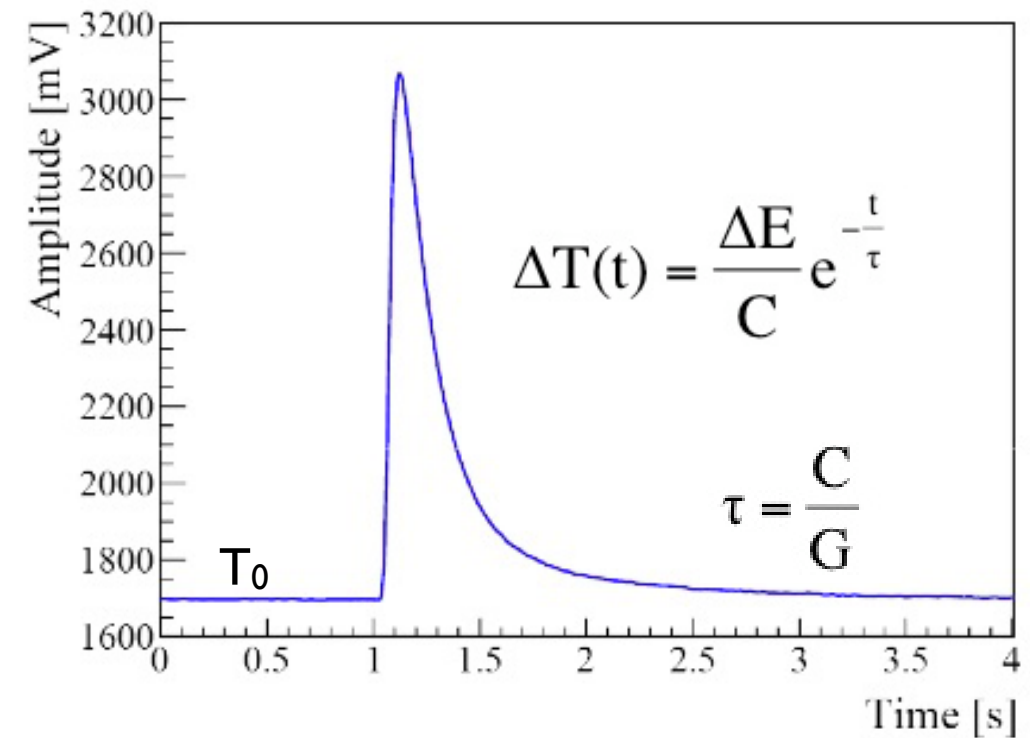
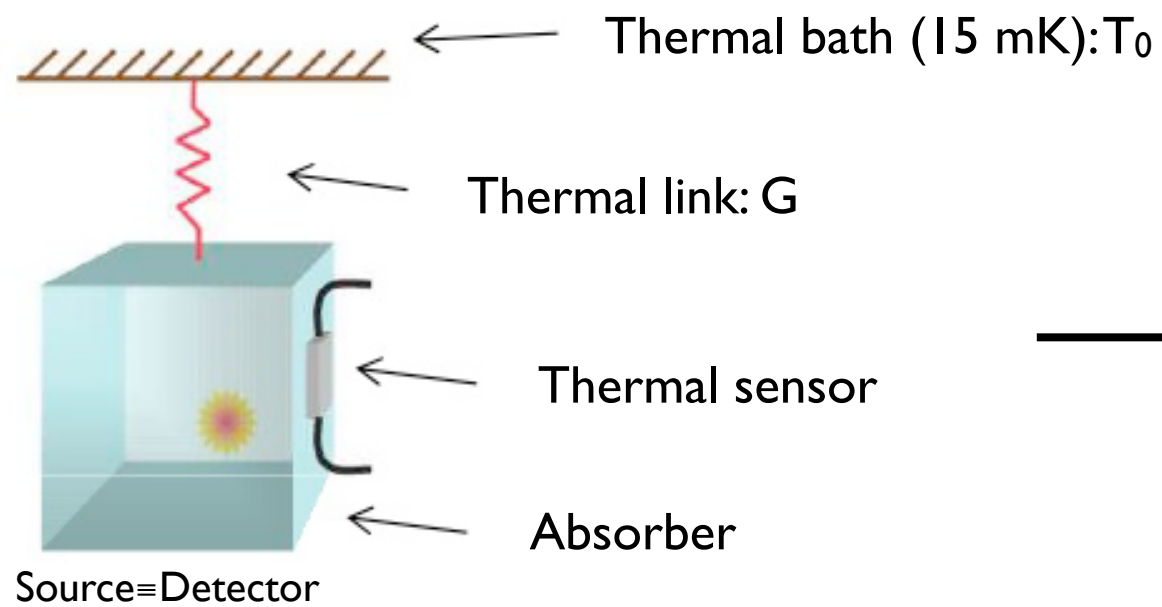
Actually running in LSM. Standard light detector equipped

Ready to run in LSM. LukS light detector equipped

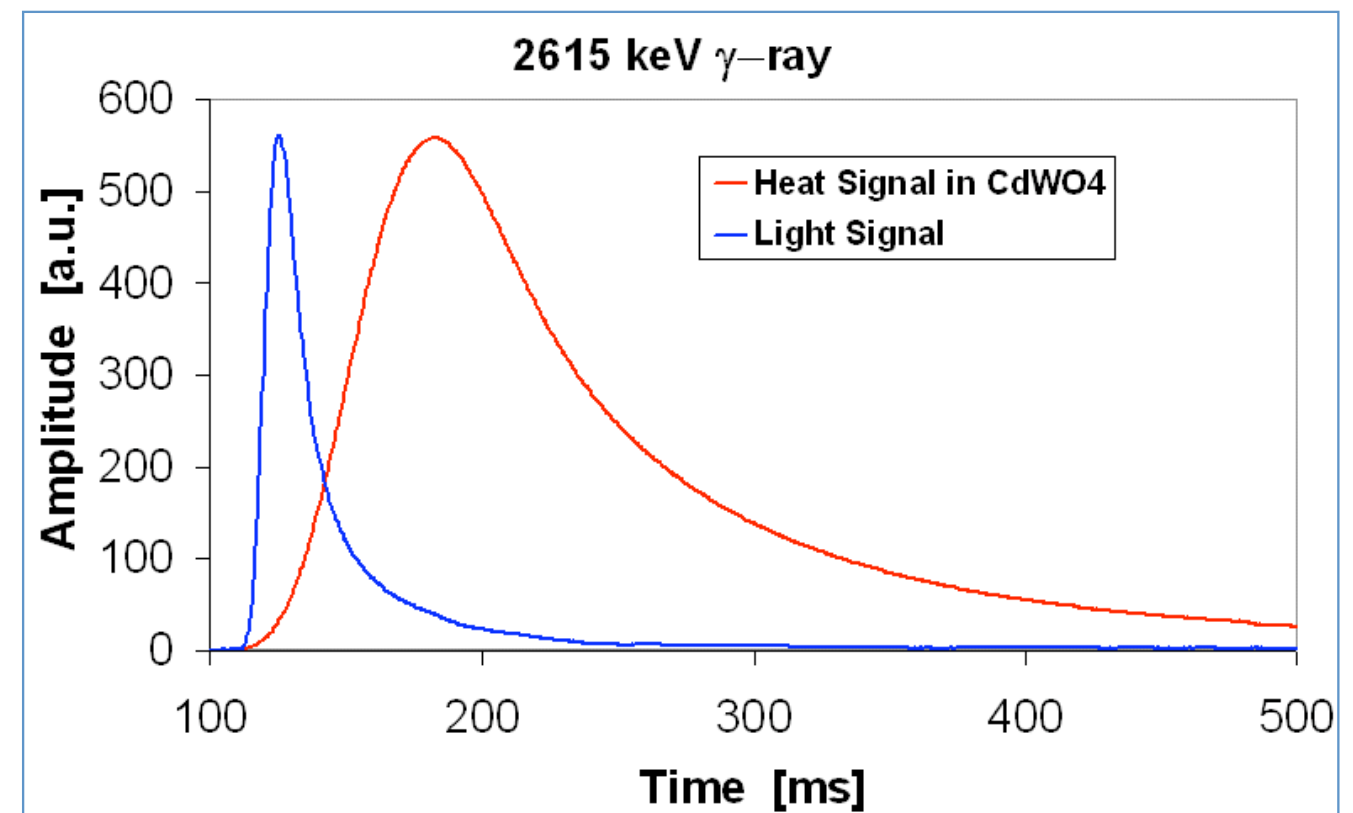
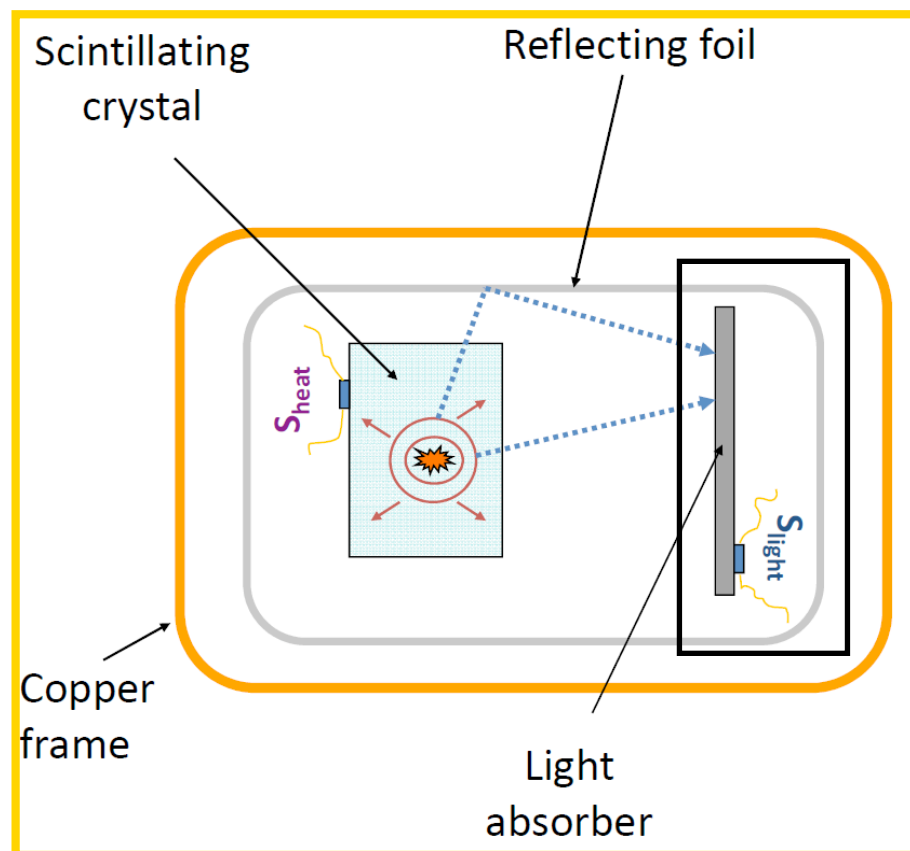


A bolometer in a nutshell

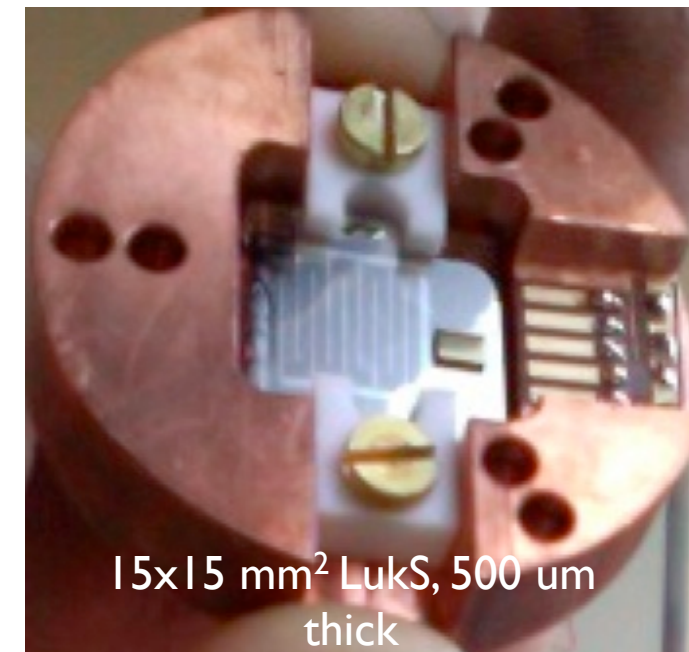
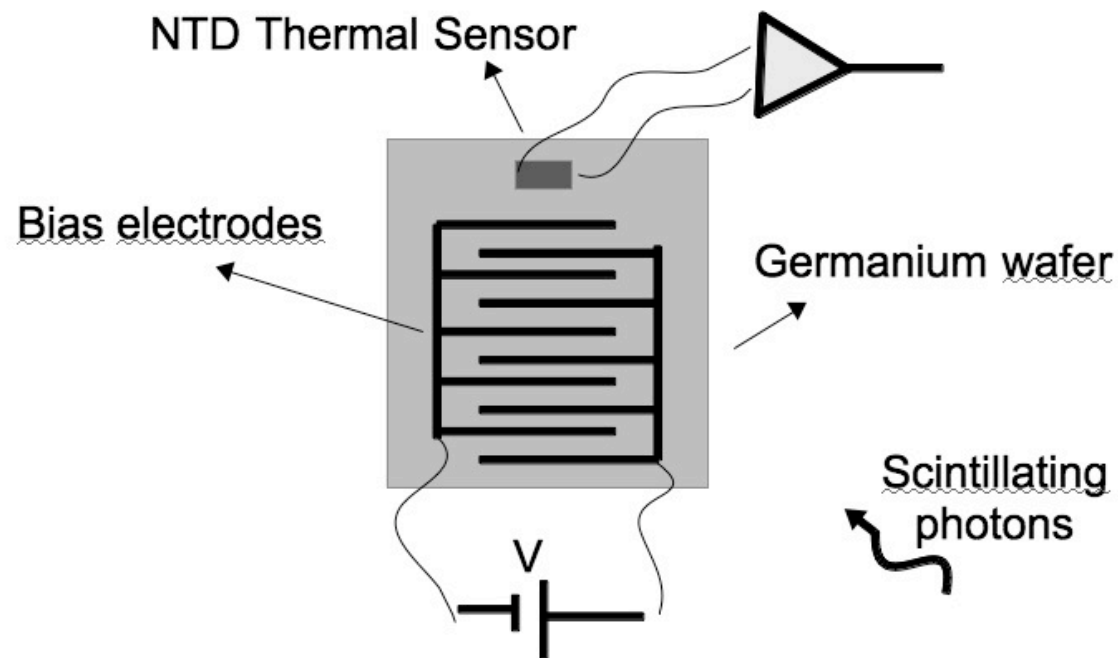
Heat



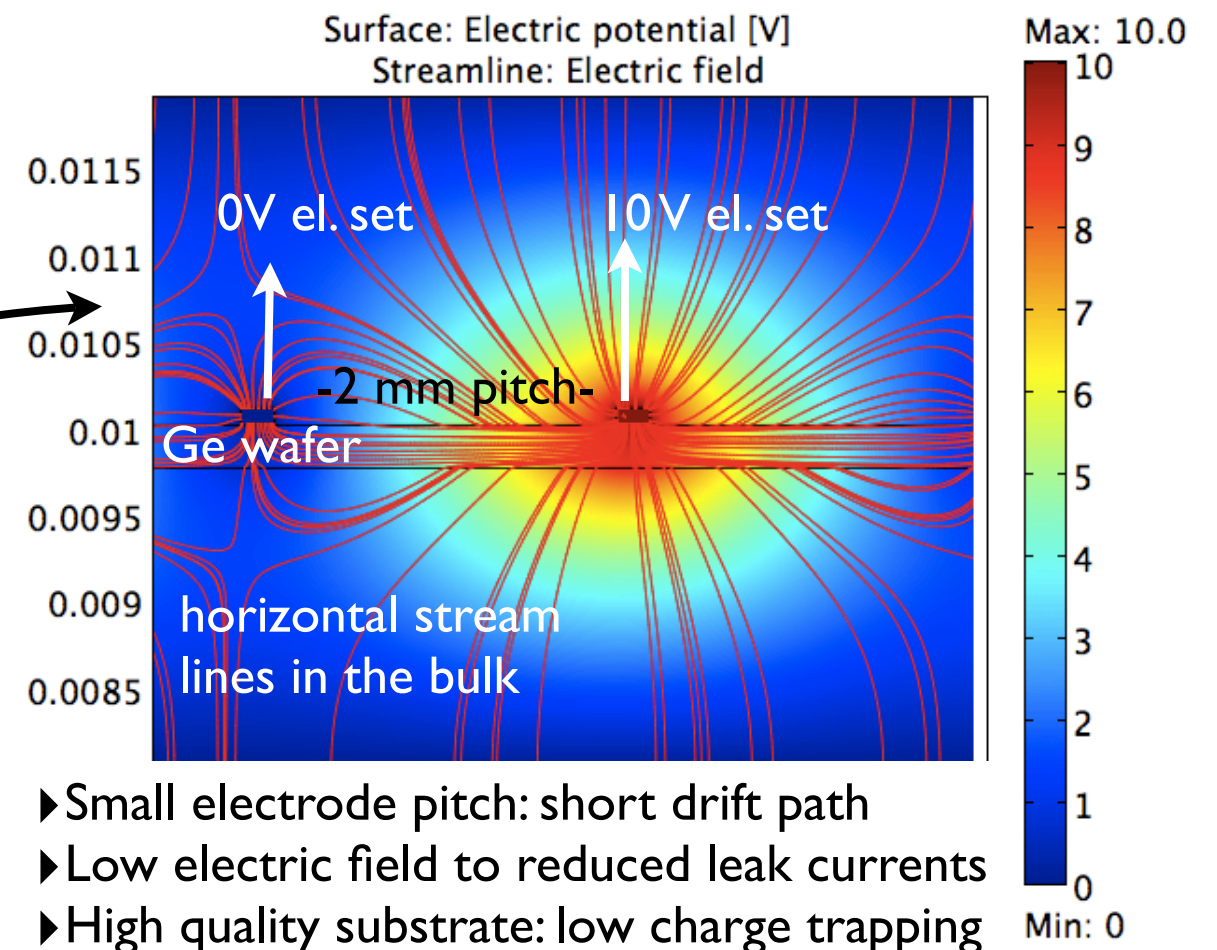
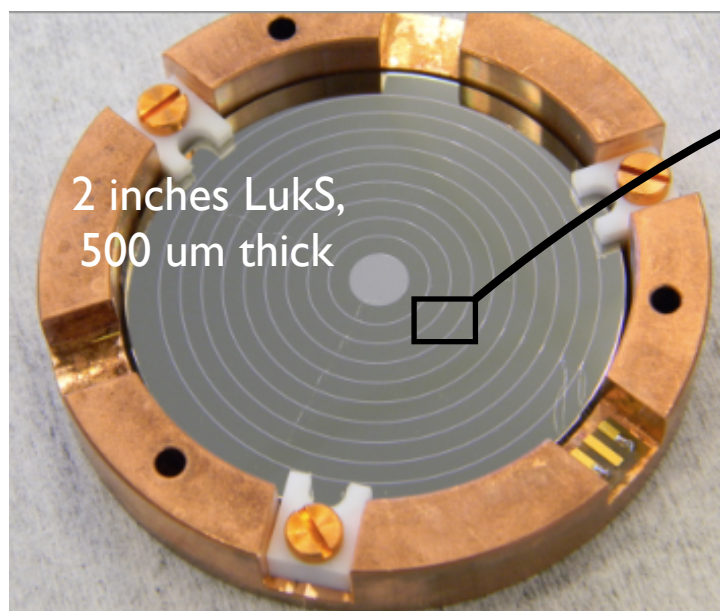
Heat and Light



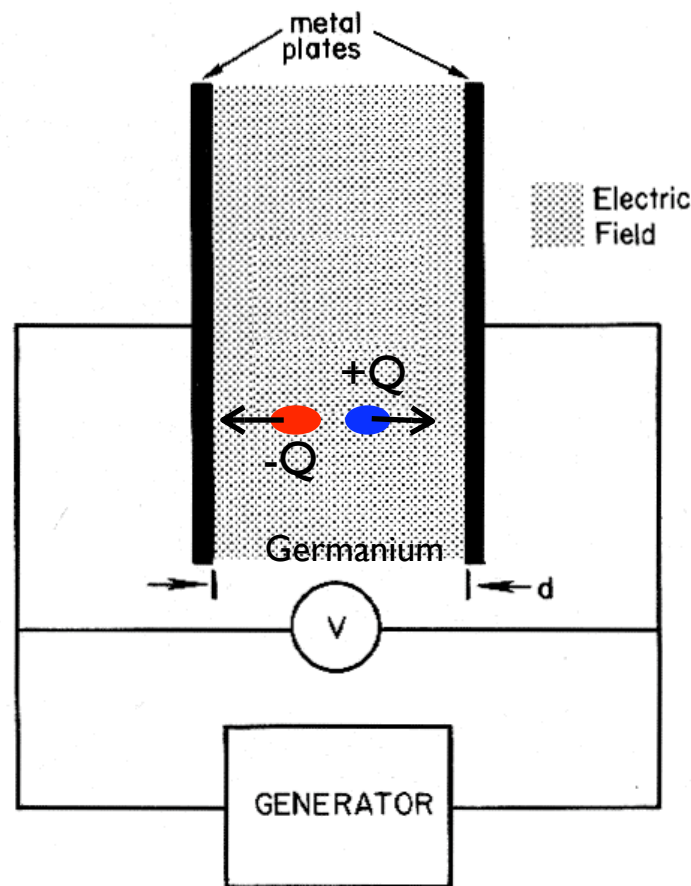
The LukS detector: generalities



Even and odd aluminum annular electrodes respectively connected by bonding aluminum wires. V_{bias} is applied within the two sets (Ex. 10V)

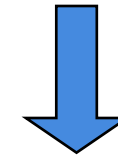


Neganov-Luke effect



- 1) Event creates e-h pairs
- 2) e-h pairs are drifted by E_{bias}
- 3) Phonon emission while e-h pairs drift. The external generator provide the energy

$$W_{\text{gen}} = Q \int_0^d \vec{E} \cdot d\vec{l}$$



$$E_{\text{heat}} = E_r (1 + qV_{\text{bias}}/\epsilon)$$

- 1) The over-heating is proportional to $N_{\text{e-h}} = E_r/\epsilon$, the number of pairs created by an event of E_r energy.
- 2) The quantum yield ϵ depends on the semiconductor (gap) and photon energy. For germanium in X region: $\epsilon = 3 \text{ eV}/(\text{e-h})$. Close to the gap, ϵ is much lower...
- 3) By increasing V the heat signal is mainly dominated by the charge creation \Rightarrow **beyond the thermodynamic fluctuation limit (Fano fluctuations)**

$$\langle (\Delta E)^2 \rangle = k_B \cdot T_0^2 \cdot C$$

Germanium quantum yield

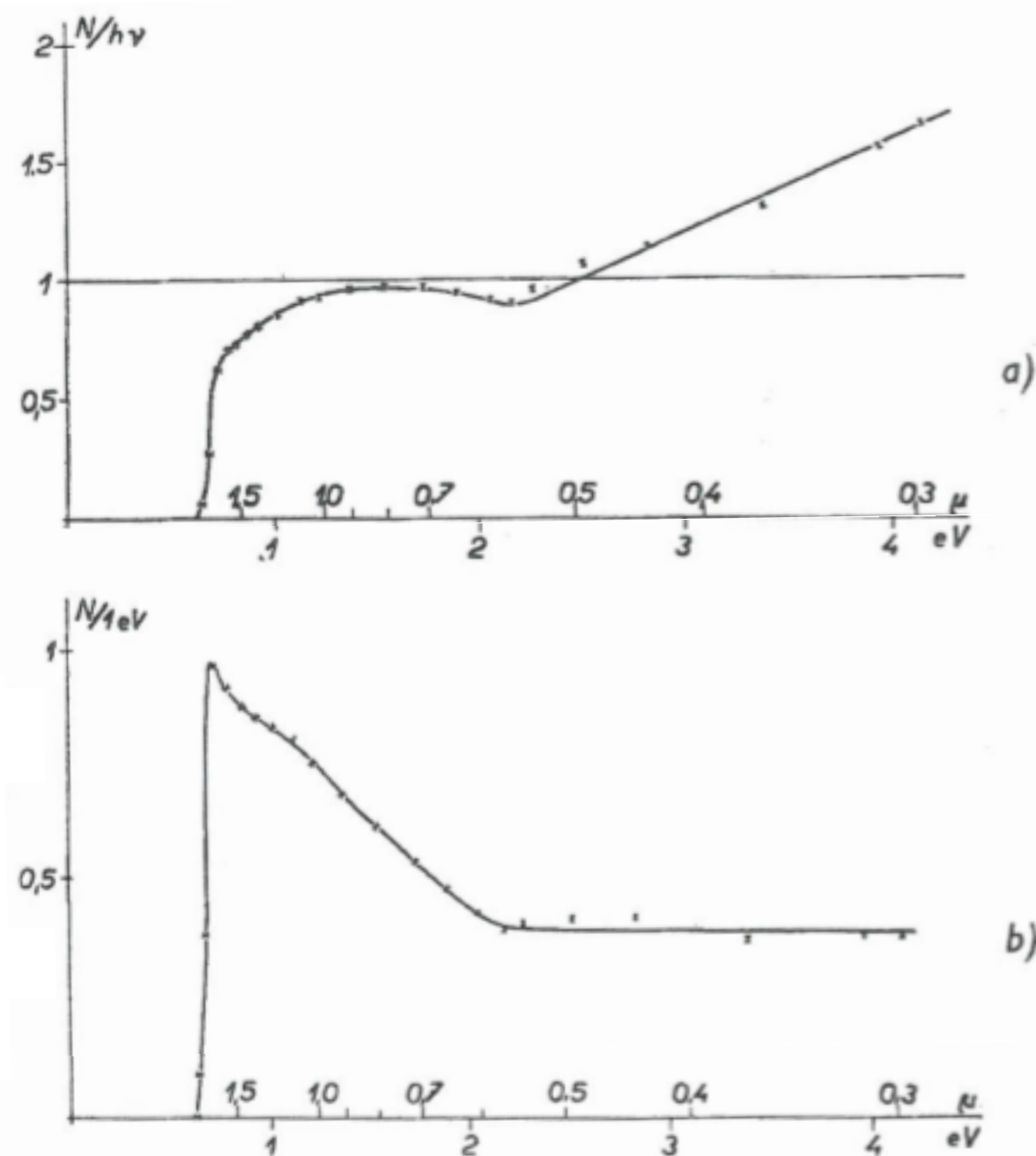
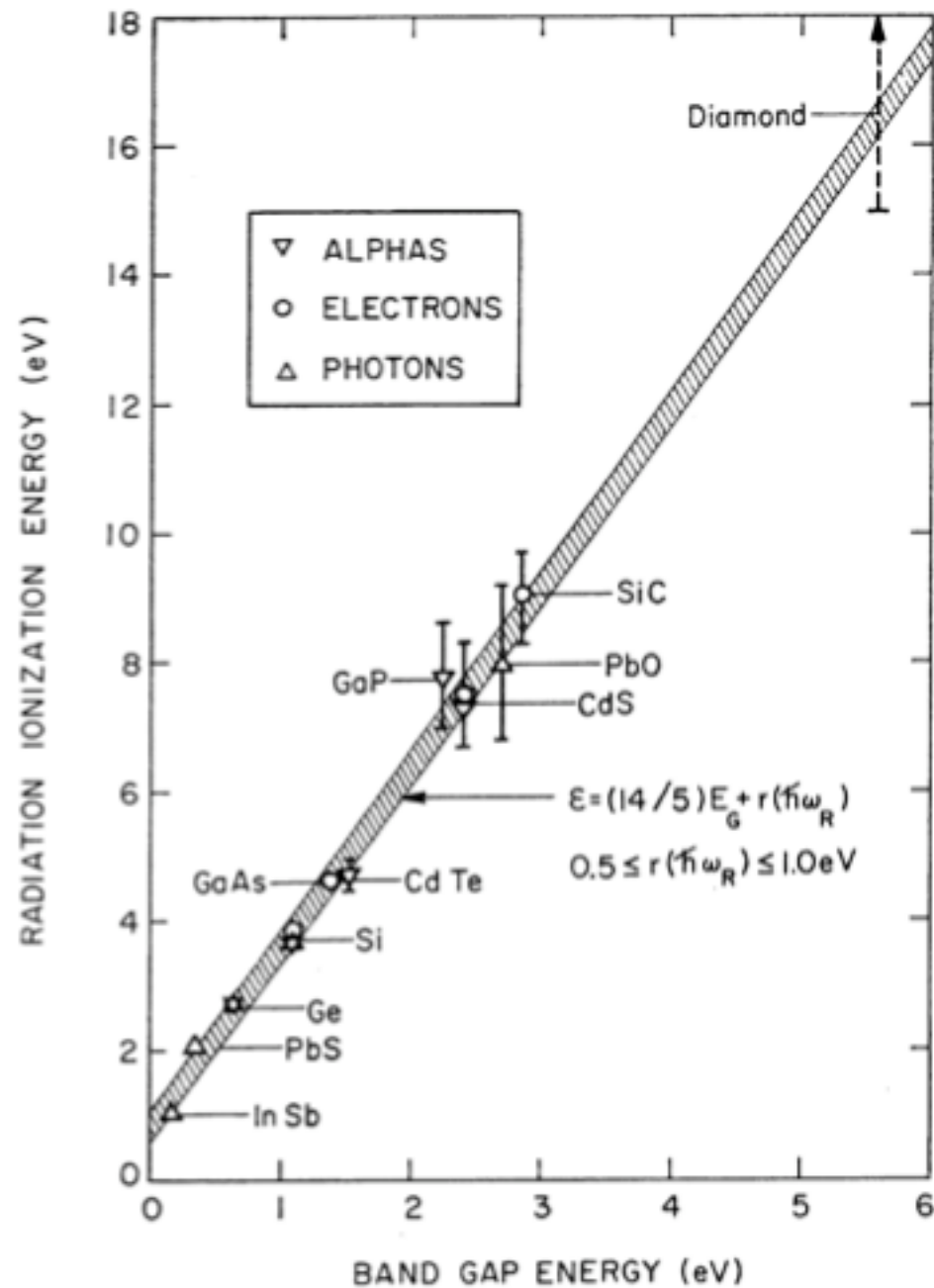


Fig. 1. a) Number of electron-hole pairs generated in germanium by the absorption of one photon as a function of the energy of the incident radiation.
b) Number of electron-hole pairs generated in germanium by the absorption of the total energy of one electronvolt as a function of the energy of the incident radiation.

Neganov-Luke sensor: the two regimes



Two bolometer working regimes:

Prompt Heat Regime

$$E_{\text{heat}} = E_r (1 + qV_{\text{bias}}/\epsilon)$$

Sensitive to thermo-dynamic fluctuation

Fano Regime

$$E_{\text{heat}} = E_r (1 + qV_{\text{bias}}/\epsilon)$$

Unresponsive to thermo-dynamic fluctuations but statistics of charge carriers

The resolution R (dE/E) is:
$$R = \sqrt{\frac{F\omega}{E}}$$

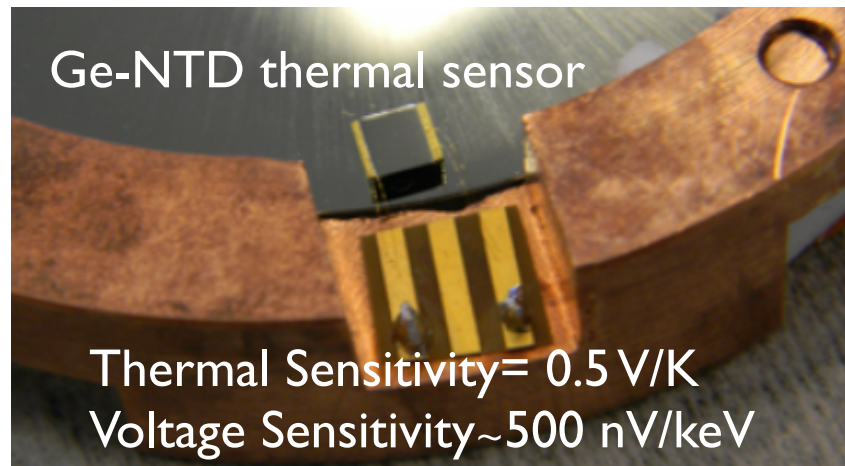
ω is the e-h creation energy. In germanium, it varies between 0.64 and 3 eV, depending on the E of the impinging particle. Fano factor F is material-dependent and in Ge is 0.13

We can increase the heat signal without any limit, increasing V (but dark currents)!

The 2-inches detector

The main questions:

- ▶ Baseline noise
- ▶ Energy resolution
- ▶ Gain under different excitation: X, IR Photons, Visible Photons

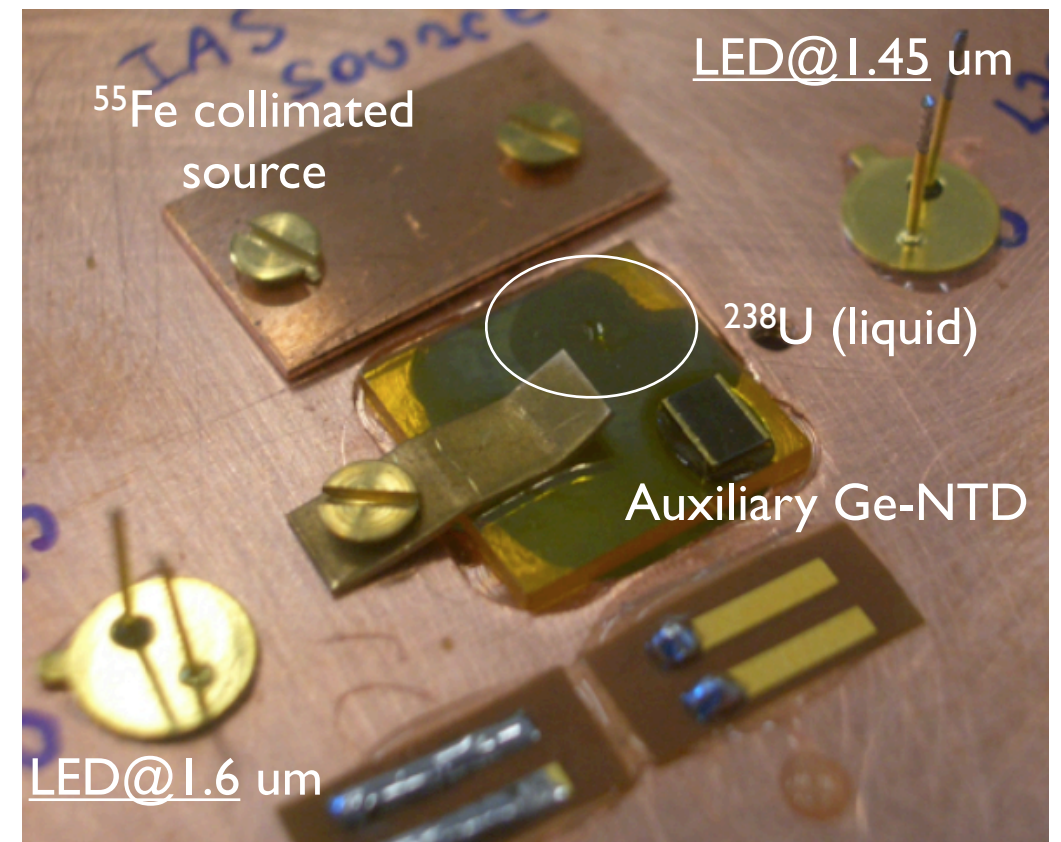
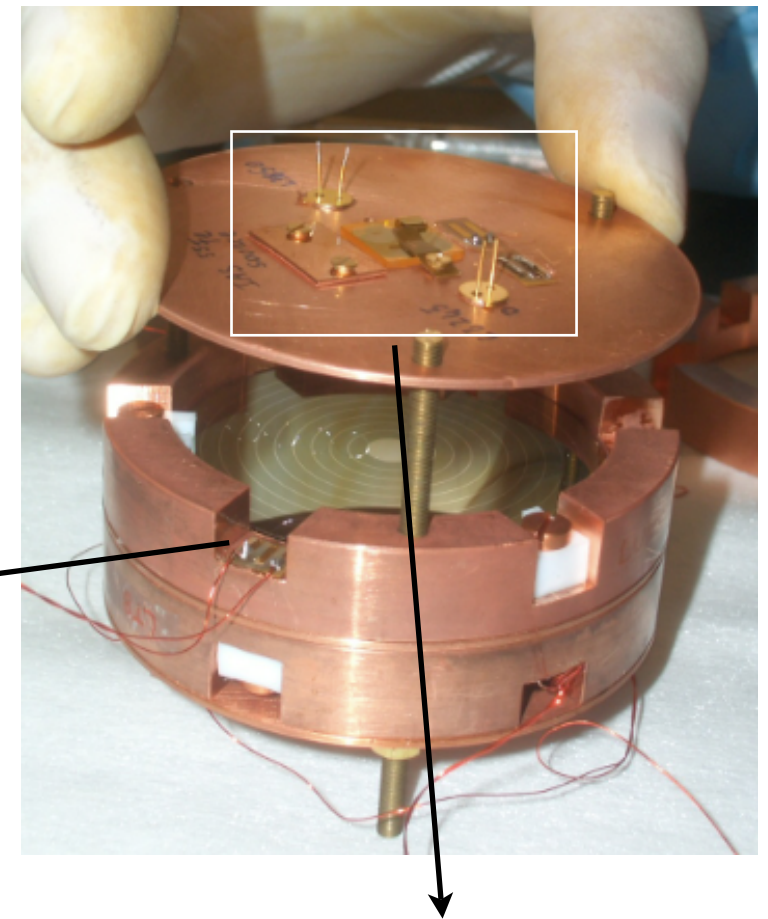


Light emitting source:

- ▶ ZnSe crystal which scintillates under ^{238}U liquid alpha source and provide almost 60 keV/alpha
- ▶ Tagging of alpha-induced light events via a second Ge-NTD sensor

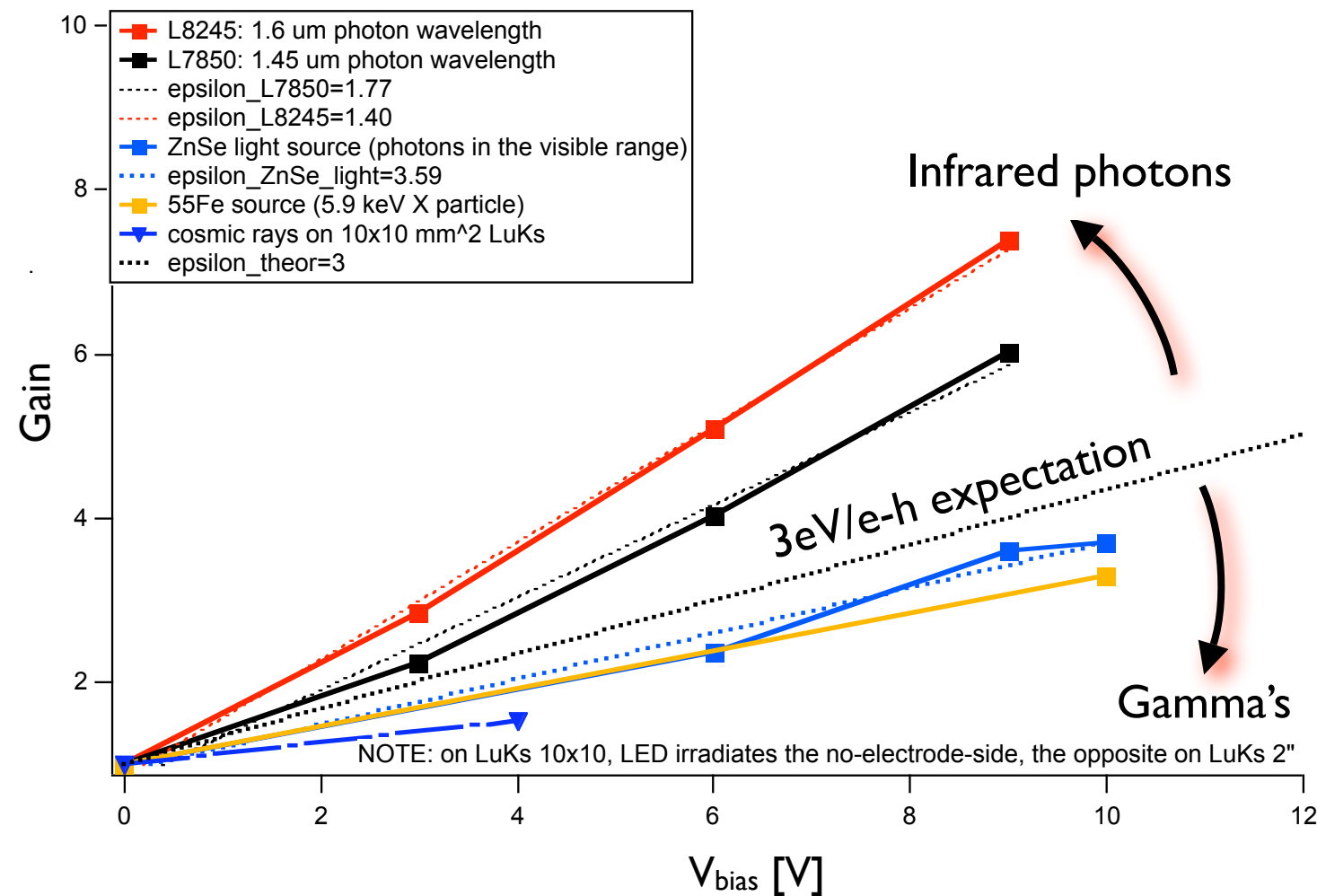
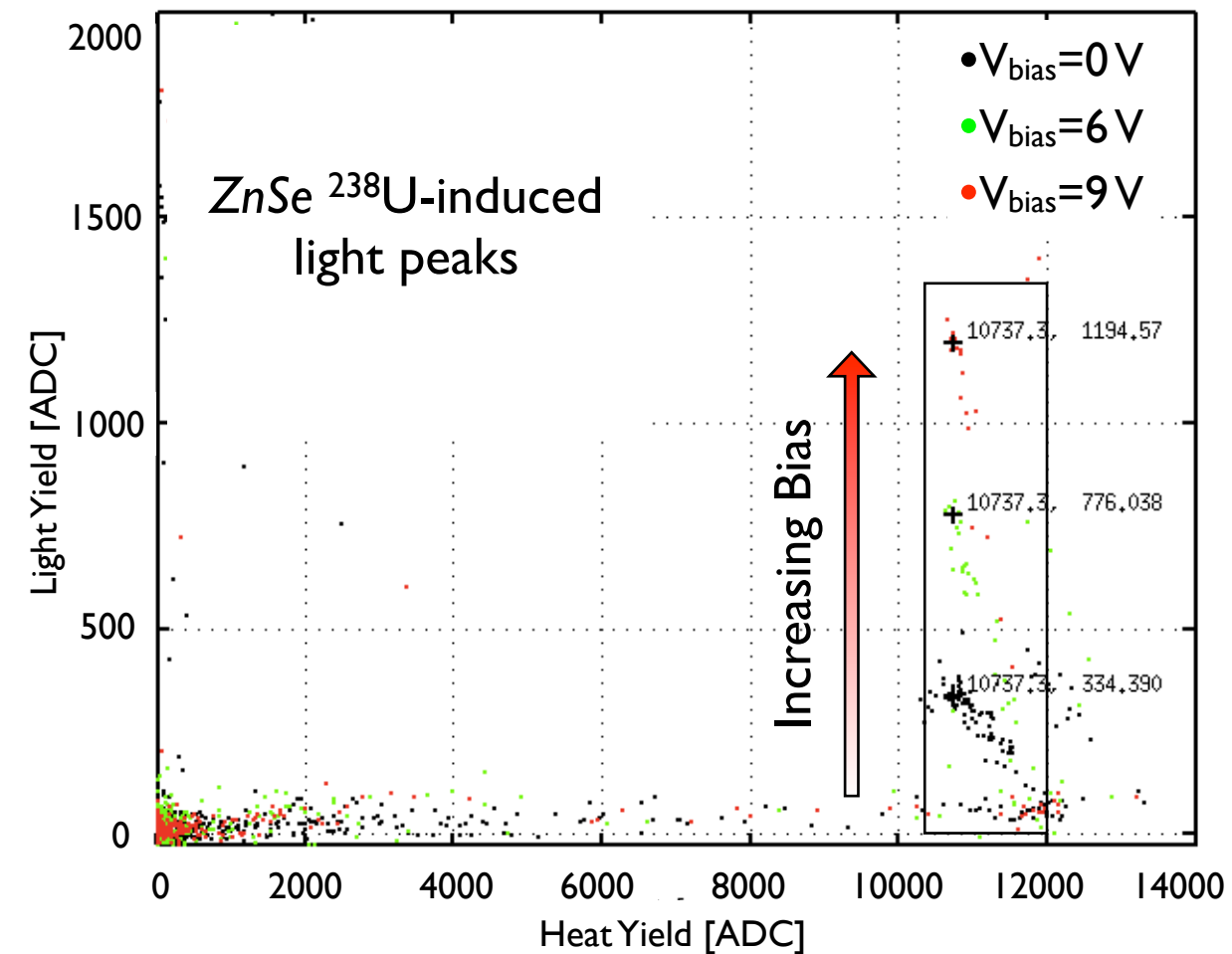
Absolute Calibration Source

- ▶ Collimated ^{55}Fe , which provides 5.9 and 6.4 keV X rays



Results: thermal sensitivity gain

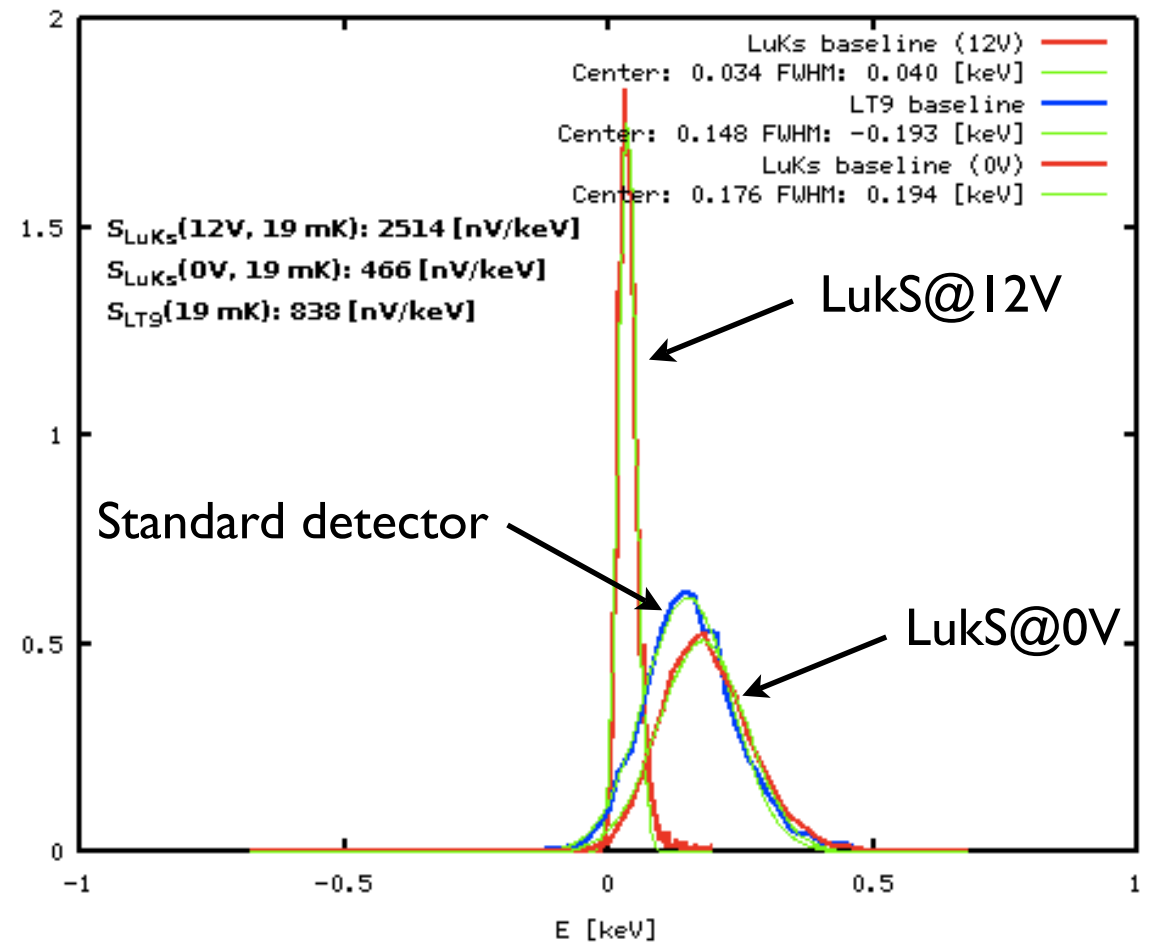
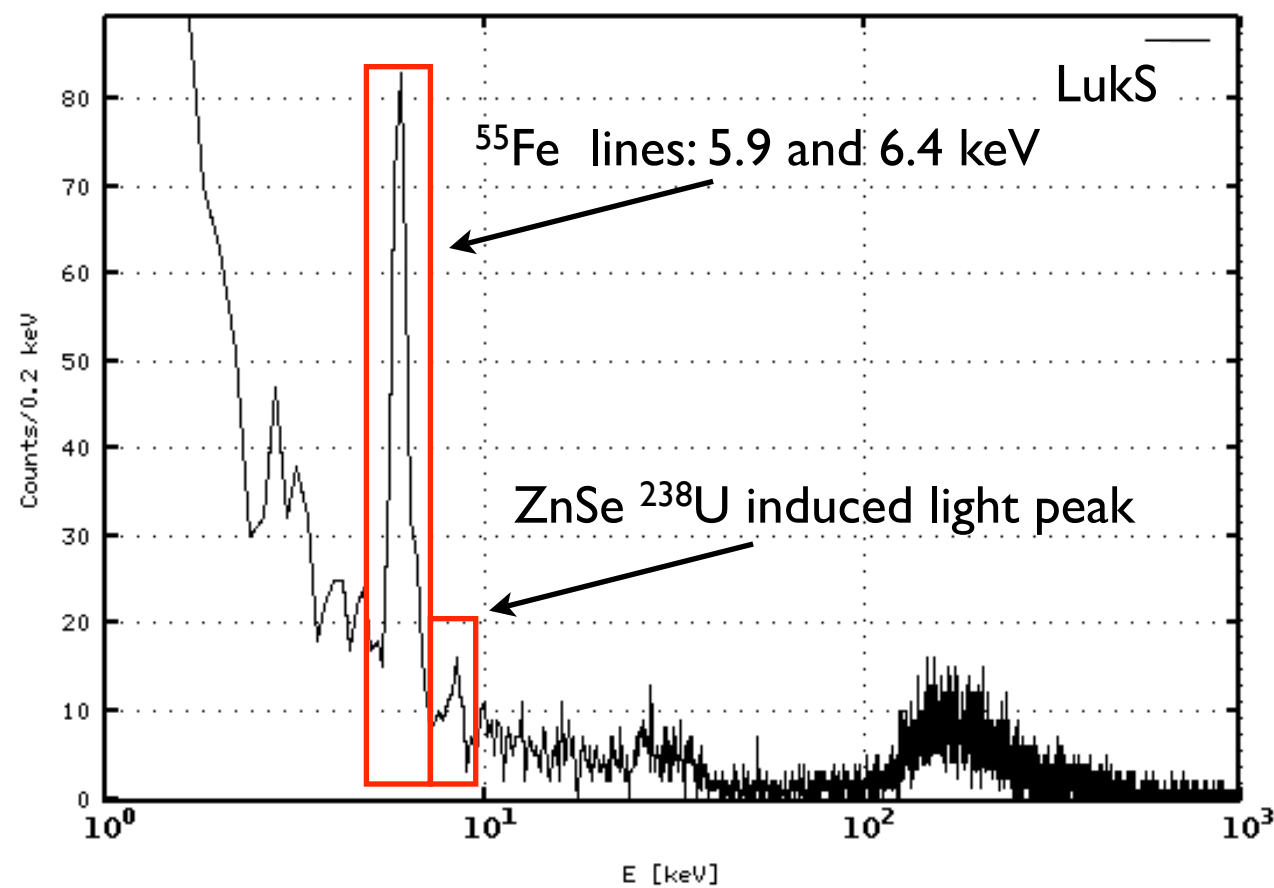
Heat and Light event-by-event scatterplot



Observations:

- ▶ Charge trapping at surfaces reduces the gain with respect to 3eV/e-h
- ▶ Gain depends on the particle/wavelength
- ▶ Leakage current appearance at high bias ($12\text{V} \rightarrow E_{\text{field}}=60\text{V/cm}$), which at the moment is the only limit to increase the gain!

Results: baseline noise reduction

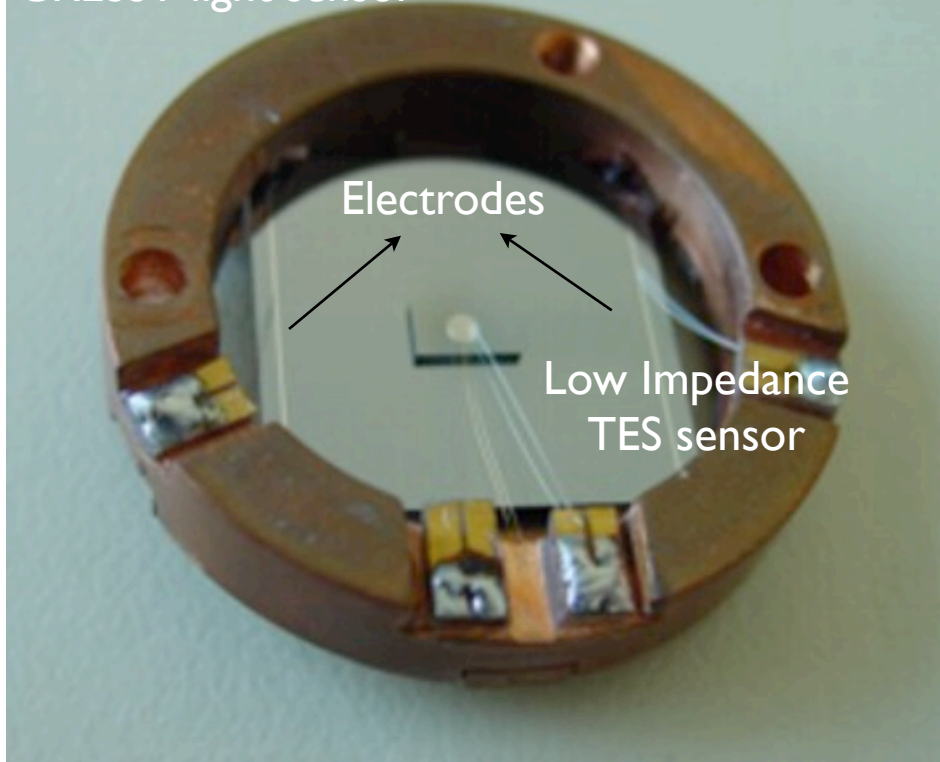


FWHM baseline noise: from 193 (0V bias) eV to 40 eV (12V bias)

- ▶ Neganov-Luke detector acts as a standard detector with zero bias \Rightarrow no specific heat added by the electrode deposition
- ▶ The baseline noise squeezes due to the sensitivity gain
- ▶ Baseline noise reduction well explained by the sensitivity gain due to Neganov-Luke \Rightarrow no excess noise injected by the electrode bias and no leakage current observed up to 12V bias!

The CRESST and LukS light detector comparison

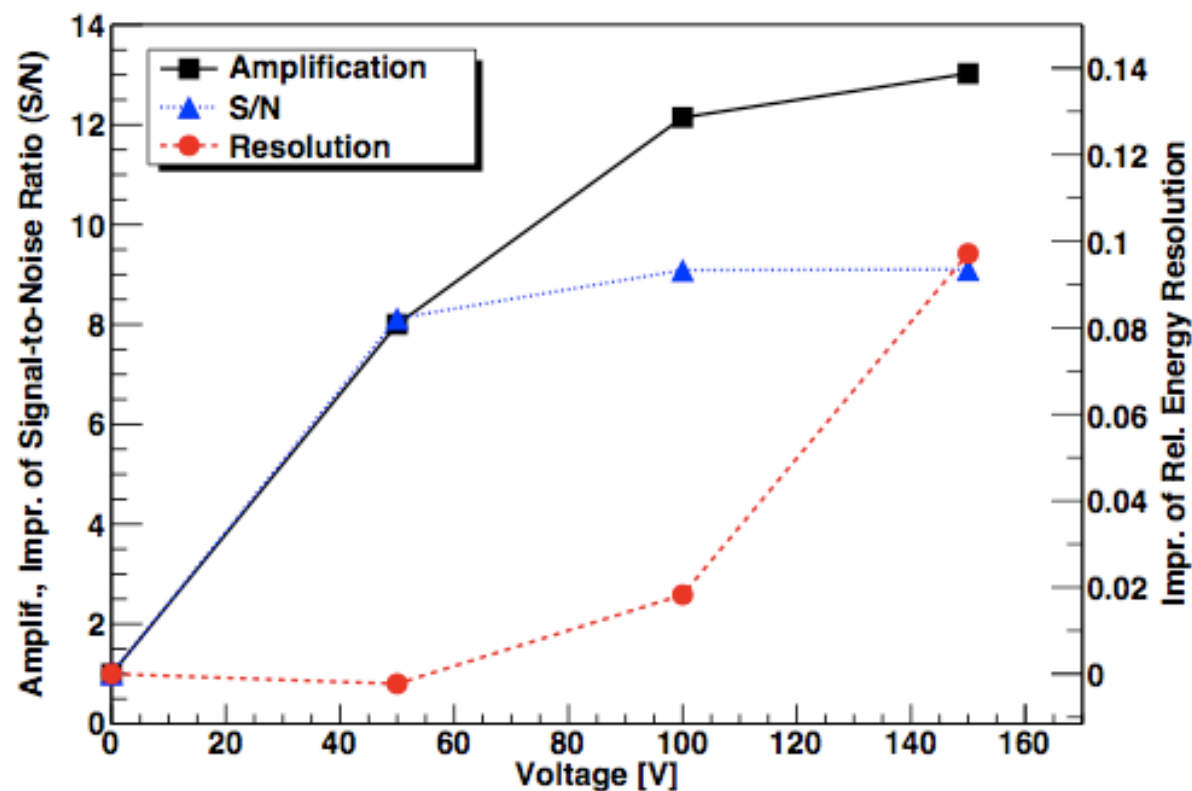
CRESST light sensor



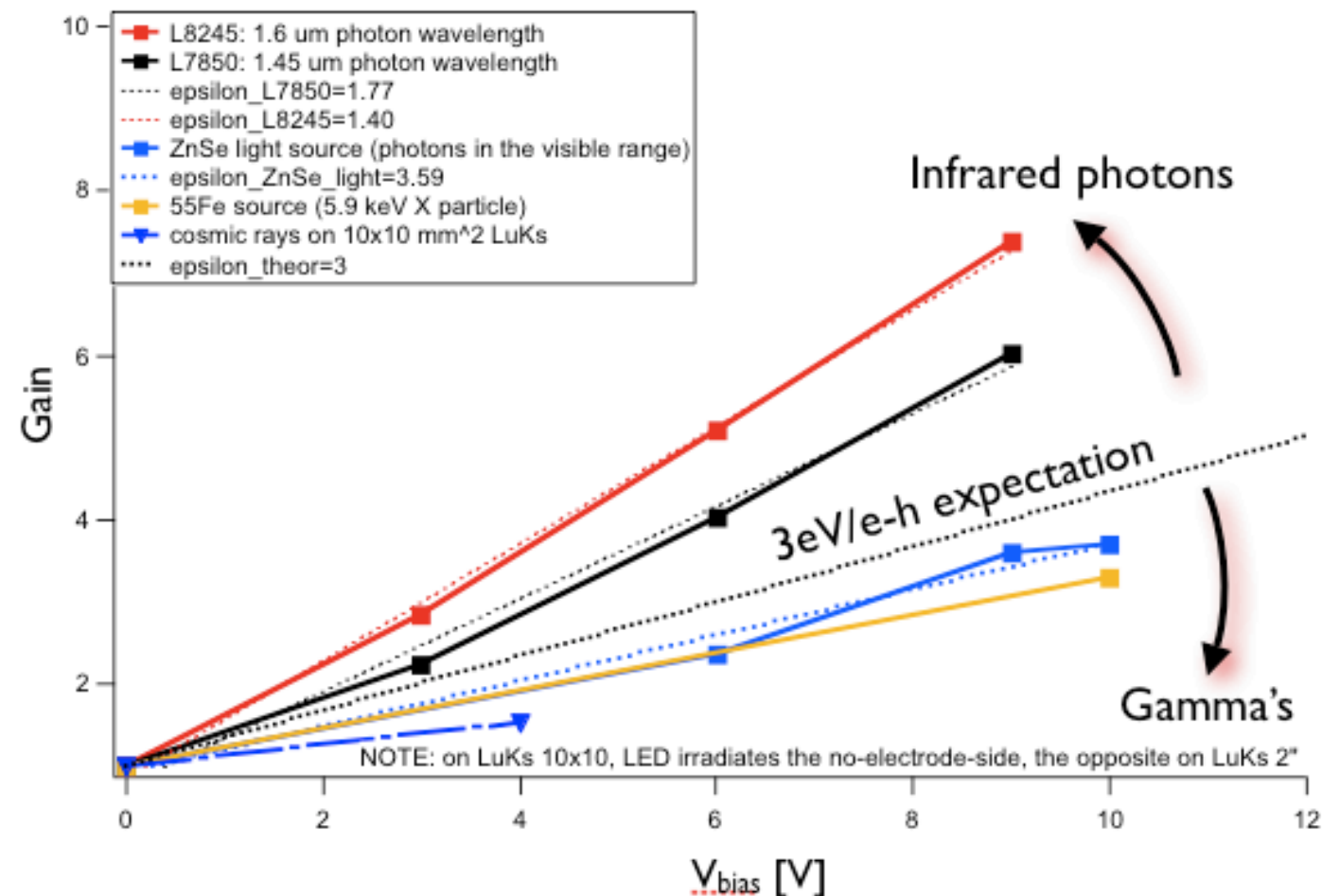
Composite detector

- Evaporation of the Ir-Au-TES onto a small substrate ($3 \times 5 \times 0.5 \text{ mm}^3$) using a shadow mask
- Thermometer is glued onto a Si substrate ($20 \times 20 \times 0.5 \text{ mm}^3$) (absorber)
- Al electrodes are directly evaporated onto the absorber

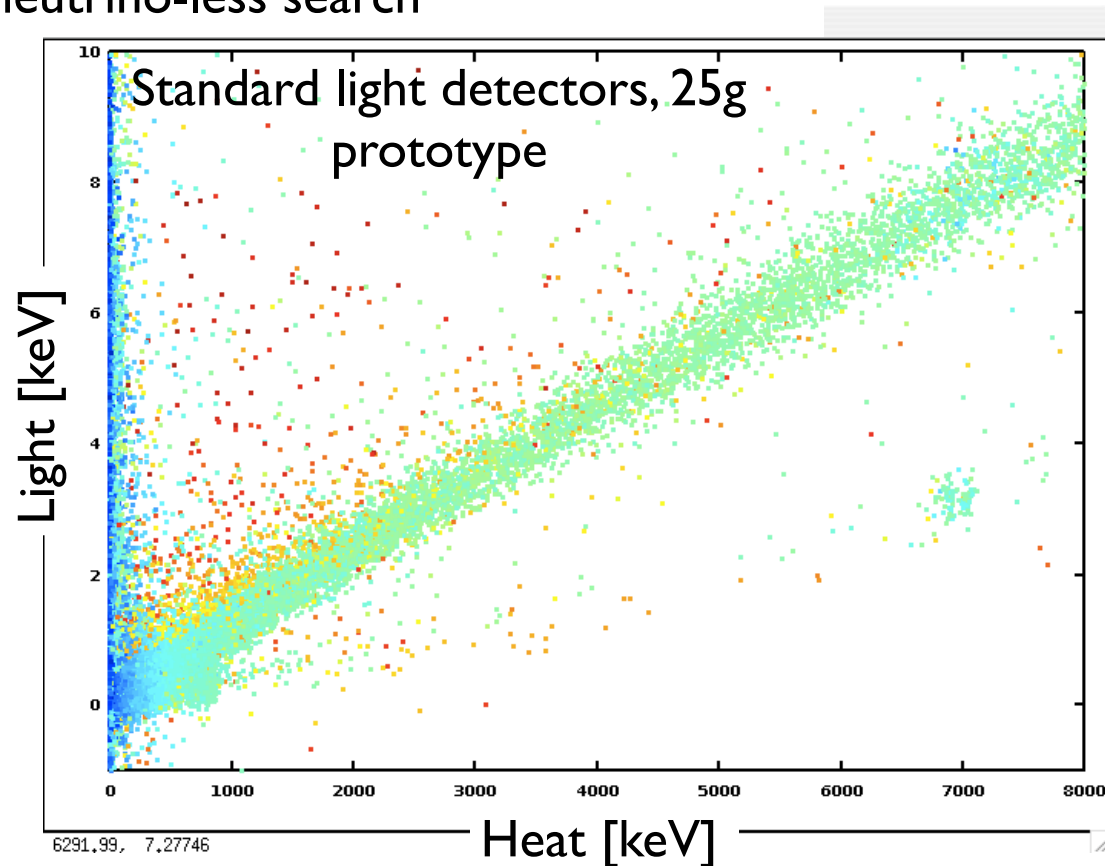
CRESST Luks



LUMINEU Luks



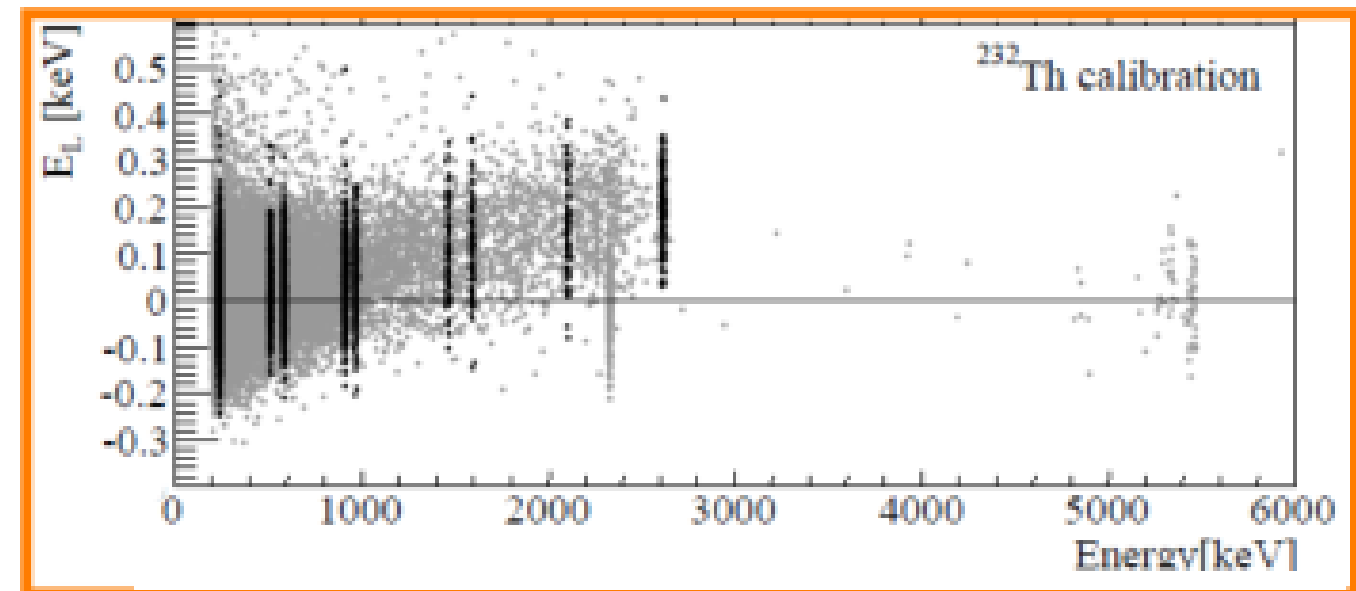
► ZnMoO_4 large mass experiment for neutrino-less search



When increasing the scintillating crystal, less light escapes the crystal

A much sensitive light detector is mandatory to scale the detector and LukS detector will fit

► Cherenkov light in TeO_2 massive bolometer



arXiv:1106.6286v1

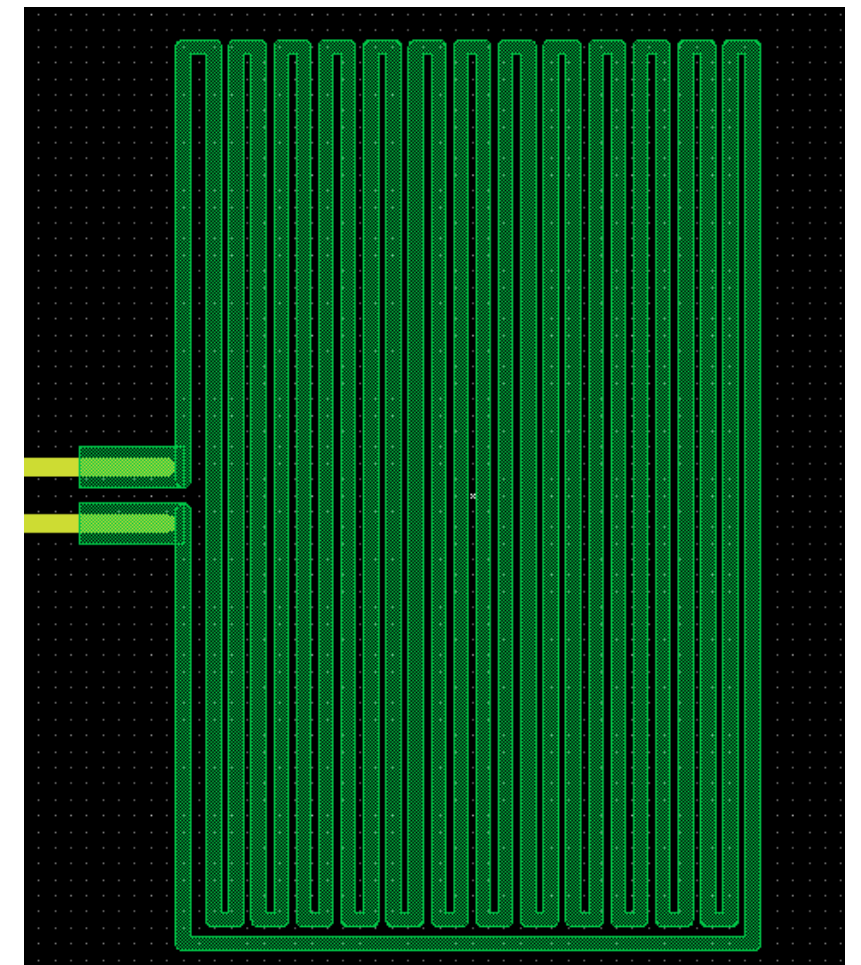
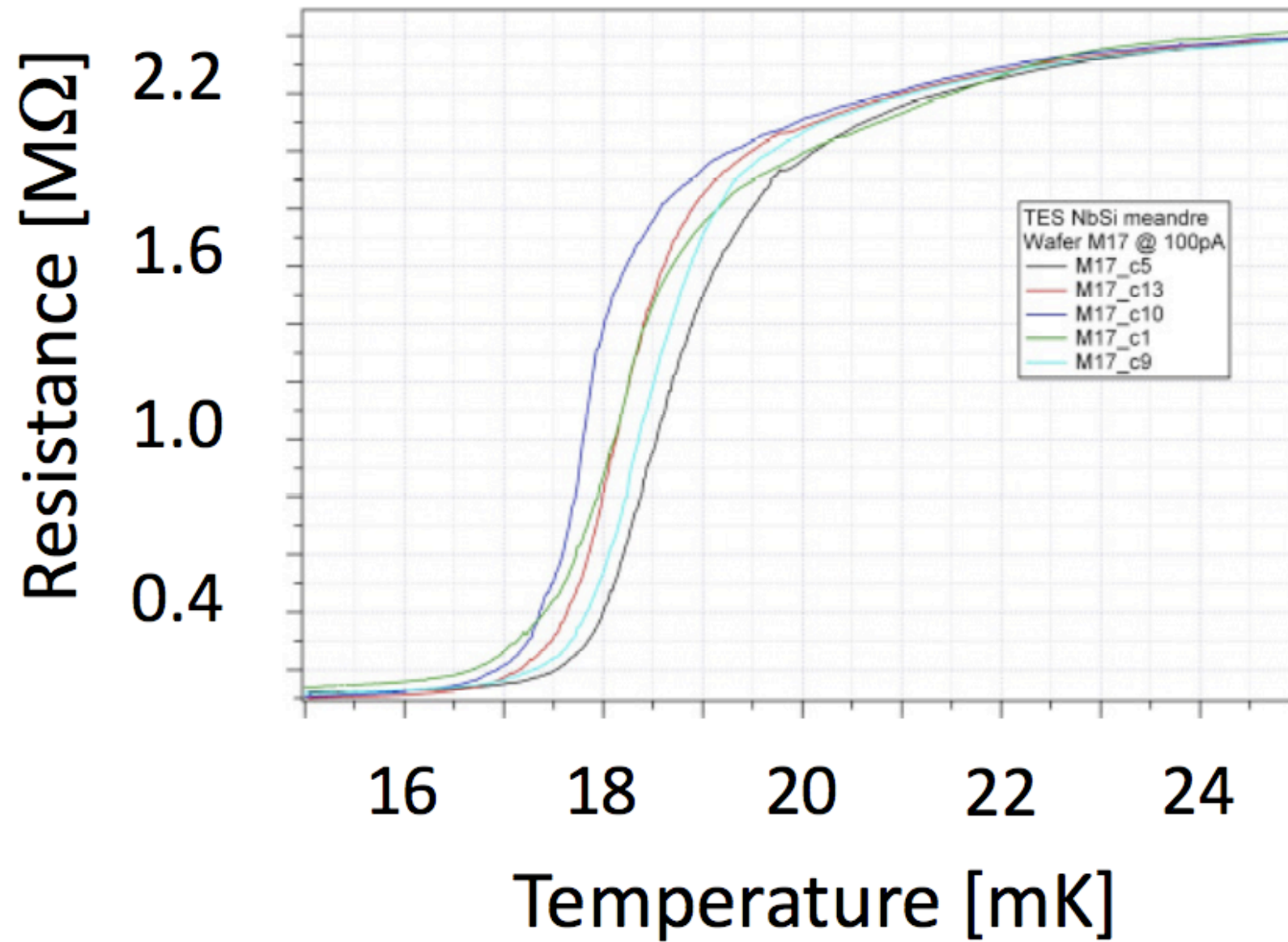
TeO_2 do not scintillate but Cherenkov!
There is no enough light resolution for event-by-event discrimination with standard light detector in TeO_2 !

LukS detectors will fit with this set-up

Replace/Re-processing all the standard light detector with Neganov-Luke assisted technologie

Next future LukS prototype

High impedance TES sensor, meander shaped.



- ✓ Different geometries of LukS detector have been successfully tested
- ✓ Gap between theoretical and experimental gain: still some charge trapping
- ✓ Sensitivity gain as high as 4 have been obtained
- ✓ Exportable technology on existing light detectors (evaporation of aluminum electrodes)
- ✓ New perspectives and possible updates for next generation
0V- 2β experiment

- ⊙ Decreasing the electrode pitch to closely obtain the theoretical gain
- ⊙ Increasing the V bias, with better control of the leakage current (we have hints...)
- ⊙ Testing the High Impedance NbSi TES sensors at the place of “standard” Ge-NTD sensors...

