Hot QCD matter what has been learned at the LHC?





Jean-Paul Blaizot, IPhT-Saclay

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Símplícity often emerges in asymptotic situations

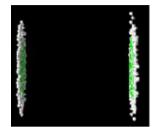
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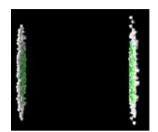
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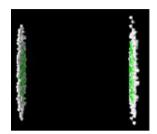
Símplícíty often emerges ín asymptotic sítuations Many phenomenologícal íssues (heavy íons are complex systems !)

colliding heavy nuclei



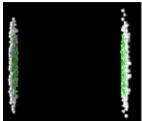


Initial conditions. Large Lorentz contraction. Nucleus wave function is mostly gluons.



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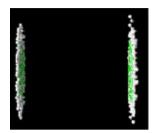
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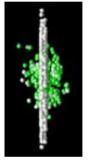


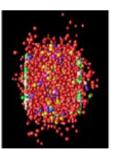


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Particle (entropy) production. Involves mostly 'small x' partons. One characteristic scale: saturation momentum Qs. Large initial fluctuations.

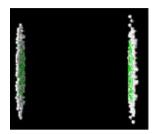


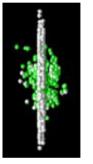




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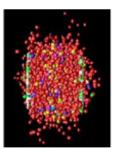




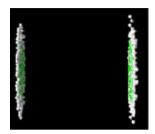


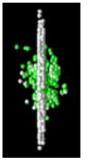
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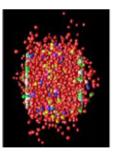
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Thermalization of produced partons. Quark-gluon plasma. Hydrodynamical expansion.



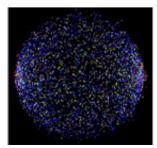


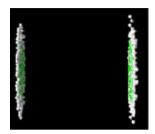


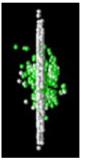
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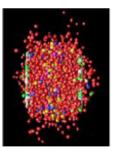
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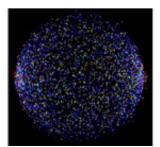




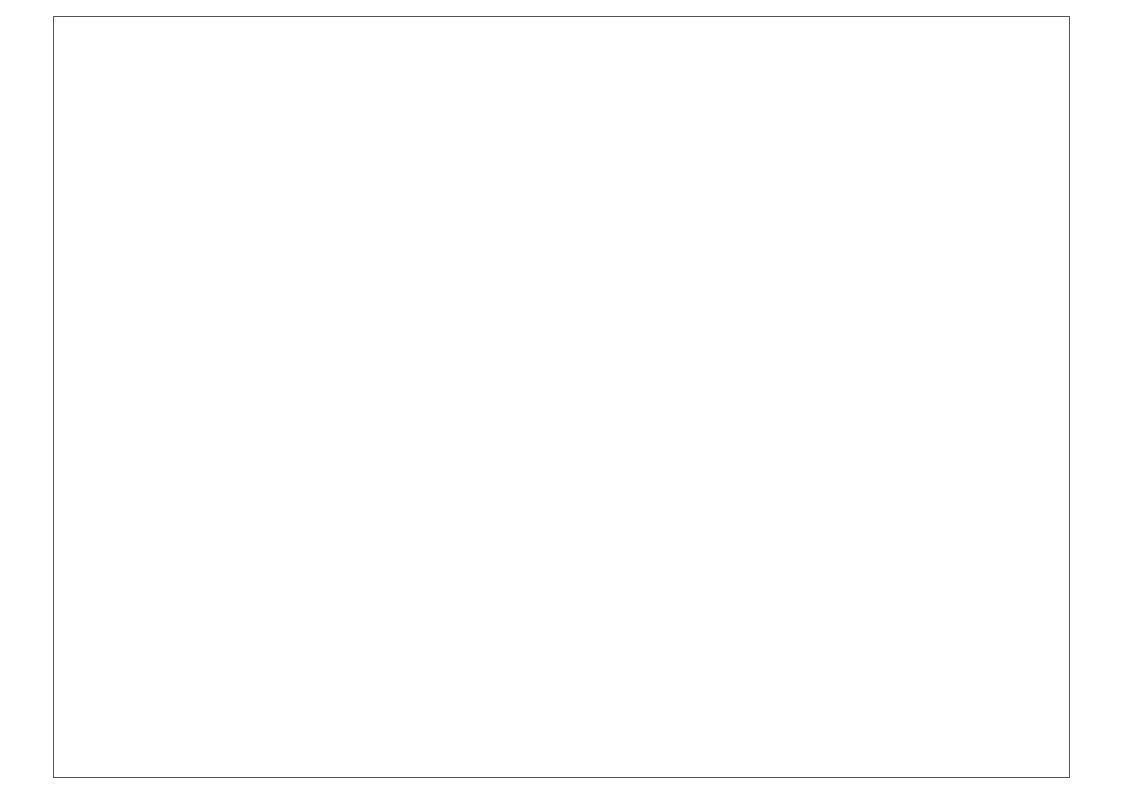
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Hadronization in apparent chemical equilibrium. Hadronic cascade till freeze-out. Measurements.





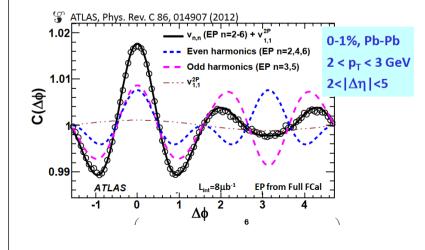
Pb+Pb @ sqrt(s) = 2.76 ATeV

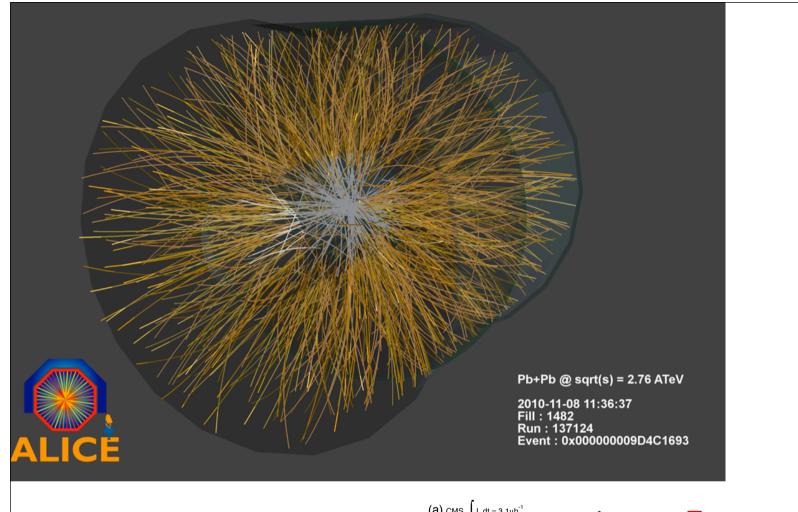
2010-11-08 11:36:37 Fill : 1482 Run : 137124 Event : 0x000000009D4C1693

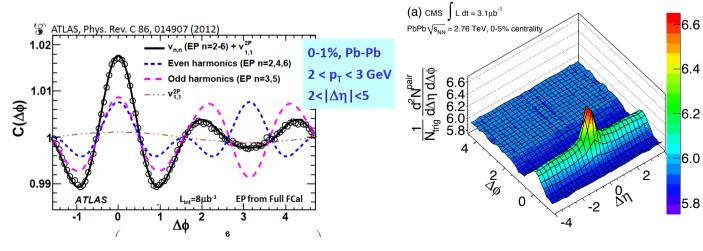


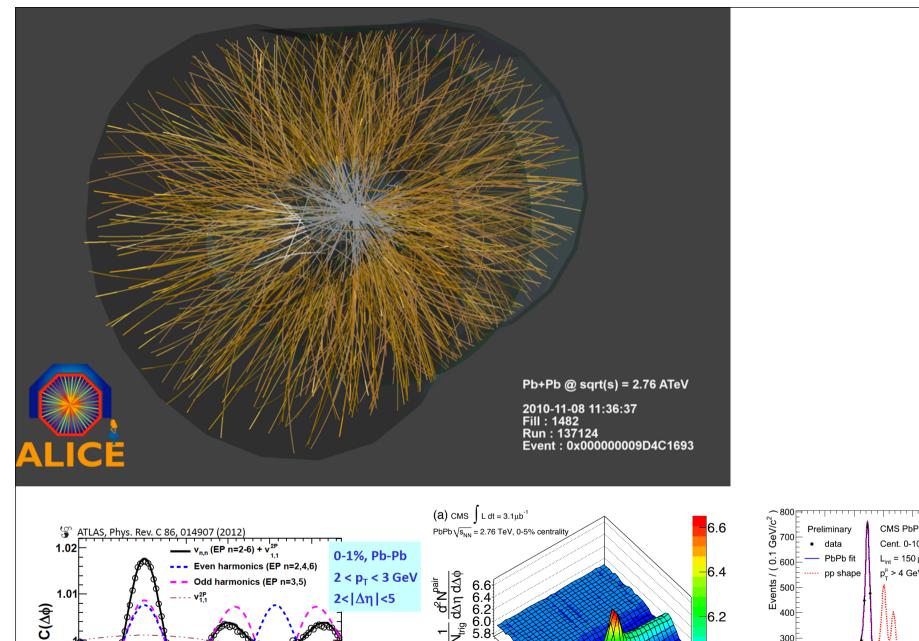
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0.99

ATLAS

0

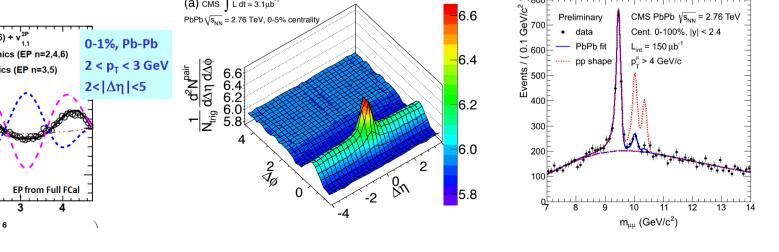
-1

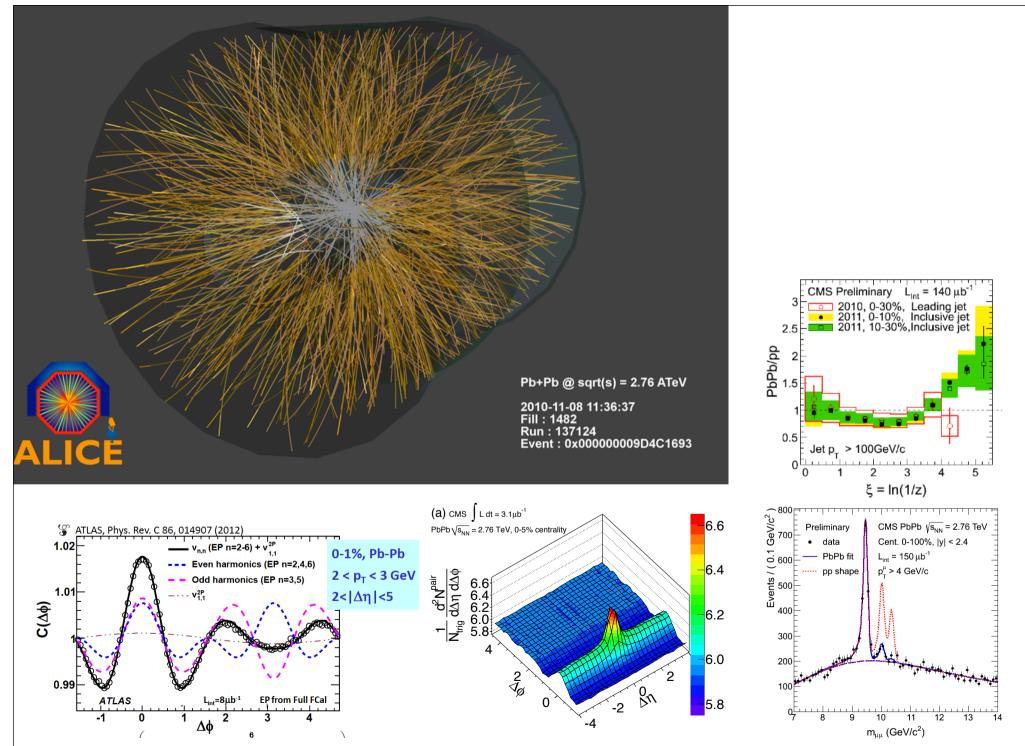
L_{int}=8µb⁻¹

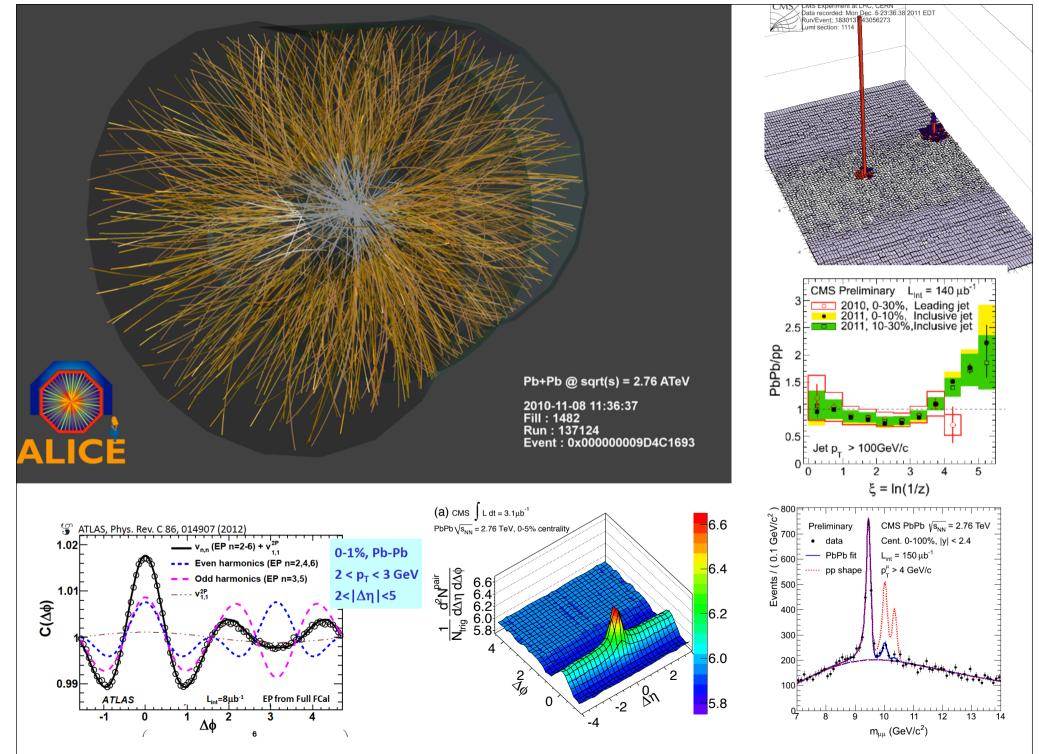
2

1

Δφ



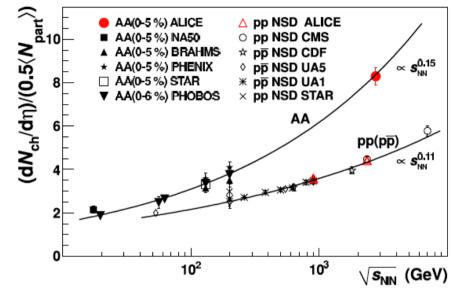




Moving backward in time Conditions are reached for the formation of a quark-gluon plasma

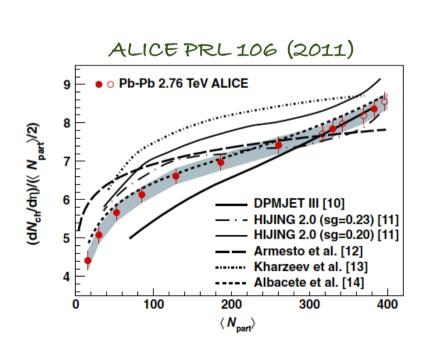
Matter at freeze-out is in chemical equilibrium

Counting particles



ALICE PRL 105 (2010)

Compatible with theoretical expectations, but large (theoretical) uncertainties remain...



2000

dN_{ch}/dη

1500

1000

ALICE

Busza [4]

HIJING 2.0 [5]

DPMJET III [6]

Levin et al. [8] Kharzeev et al. [9]

Kharzeev et al. [9]

Kharzeev et al. [10]

Armesto et al. [11] Eskola et al. [12] Bozek et al. [13] Sarkisyan et al. [14] Humanic [15]

Albacete [7]

The conditions for the formation of a quark-gluon plasma are reached in the early stages of the collisions

order of magnitude estimate

 $\leftarrow \tau_0 \longrightarrow$

$$\frac{dN_{ch}}{d\eta} \simeq 1600$$

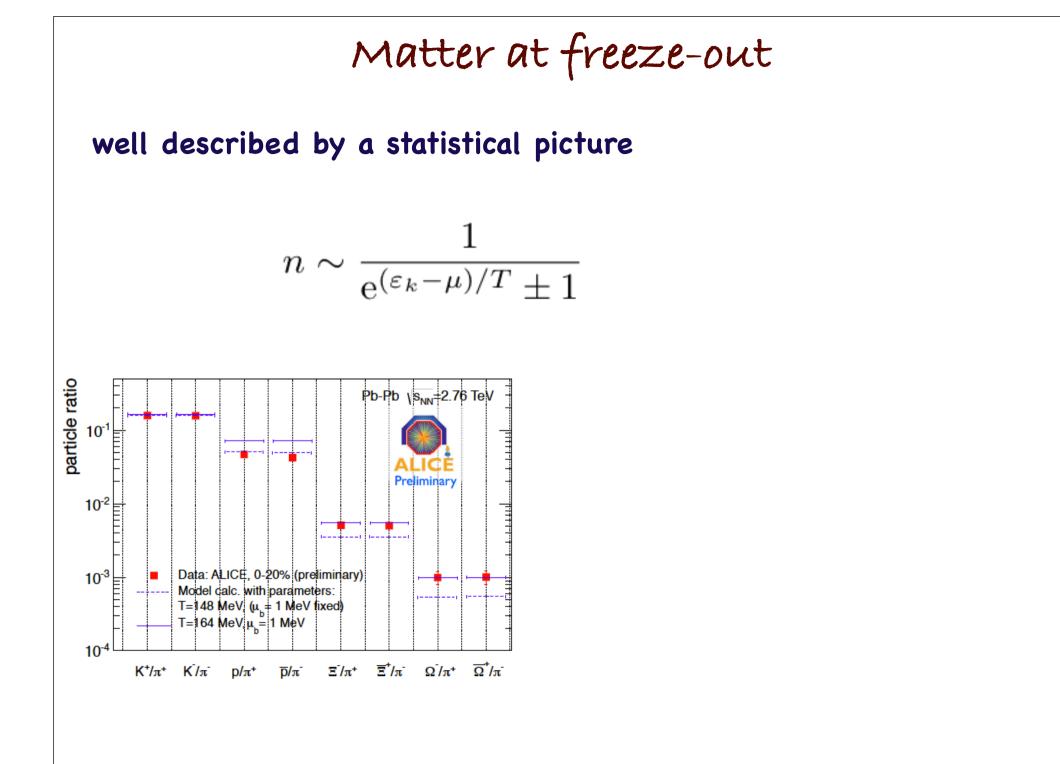
 $\epsilon \tau_0 \simeq 15 \text{GeV/fm}^2$

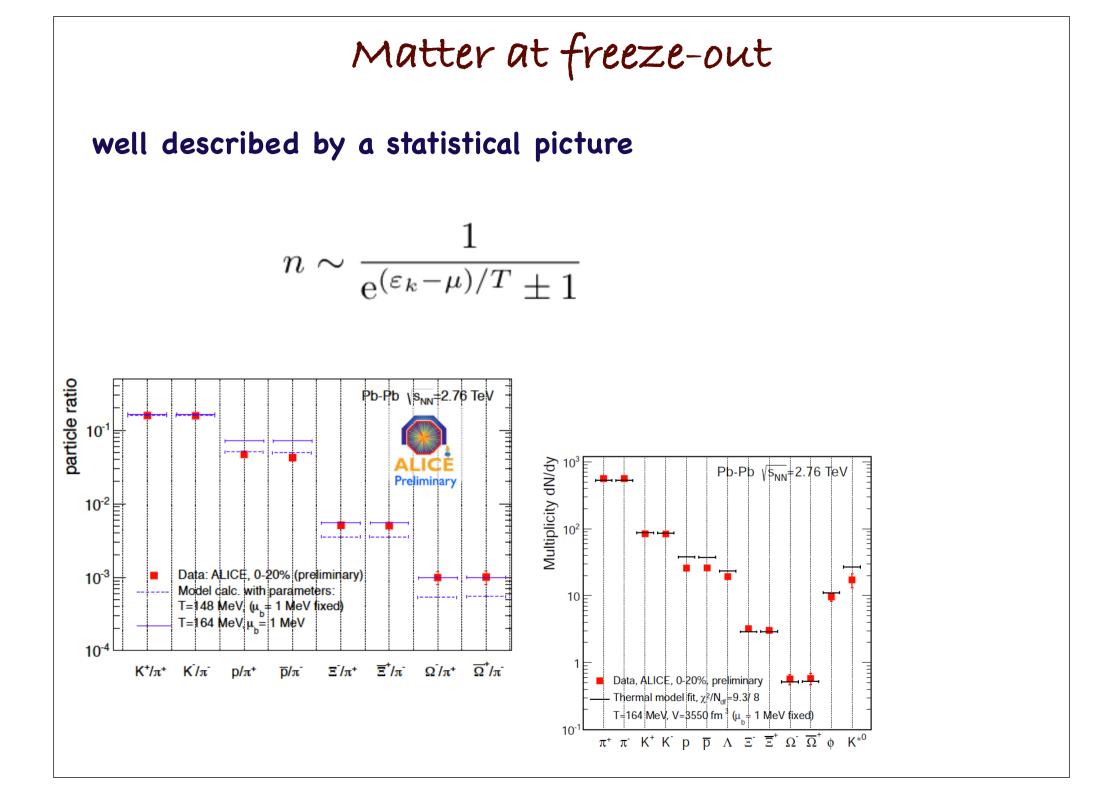
 $T_0 \simeq 300 \,\mathrm{MeV}$

Matter at freeze-out

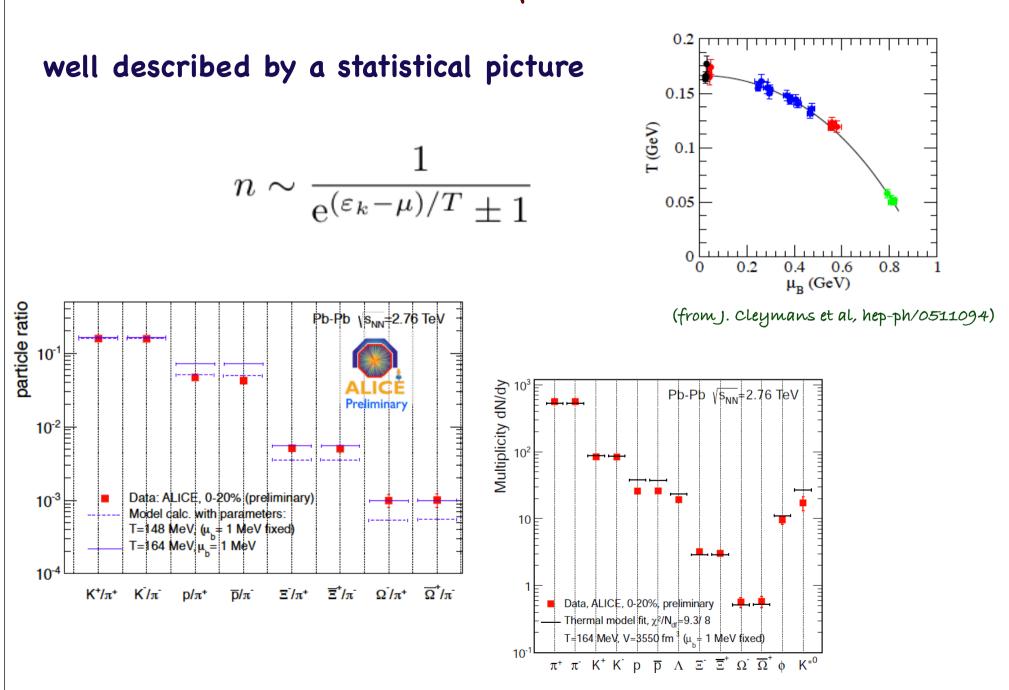
well described by a statistical picture

$$n \sim \frac{1}{\mathrm{e}^{(\varepsilon_k - \mu)/T} \pm 1}$$





Matter at freeze-out

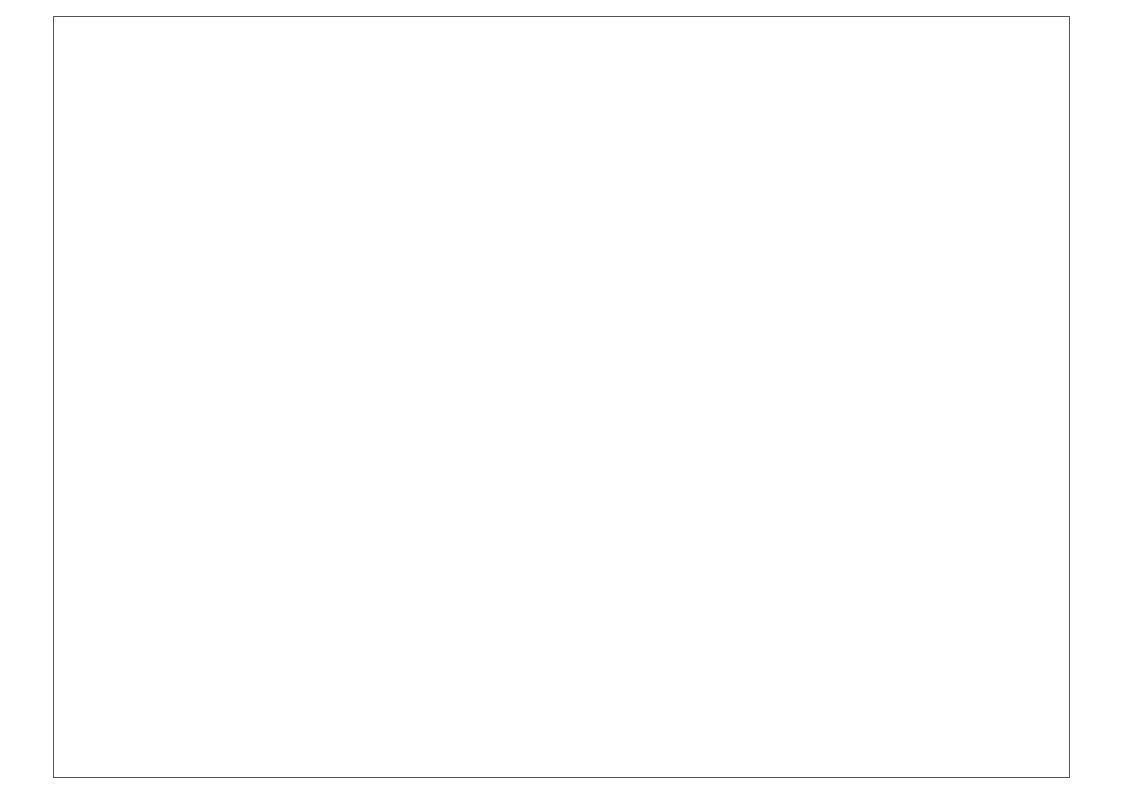


Moving backward in time

Matter flows like a fluid

The quark-gluon plasma as a nearly perfect fluid

Puzzles: viscosity, thermalization



Collective flow

Matter flows like a fluid and is well described by relatvistic hydrodynamics

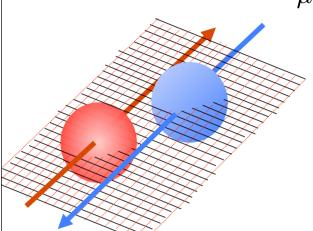
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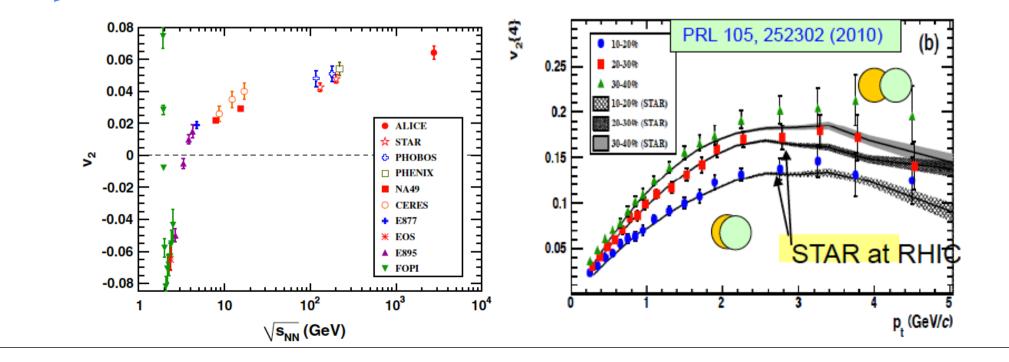


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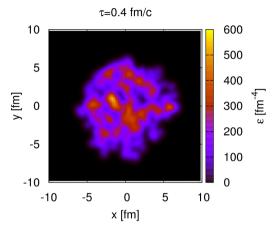
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Flow is best seen in azymuthal distributions of produced particles.

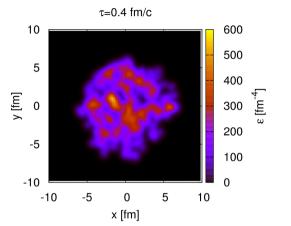


The flow is sensitive to initial density fluctuations



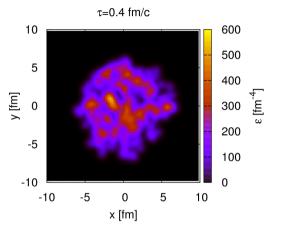
 $v_n \sim \epsilon_n$

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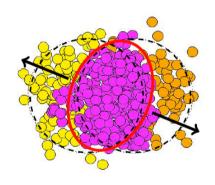


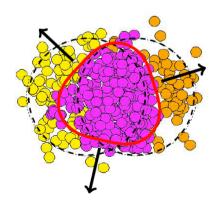
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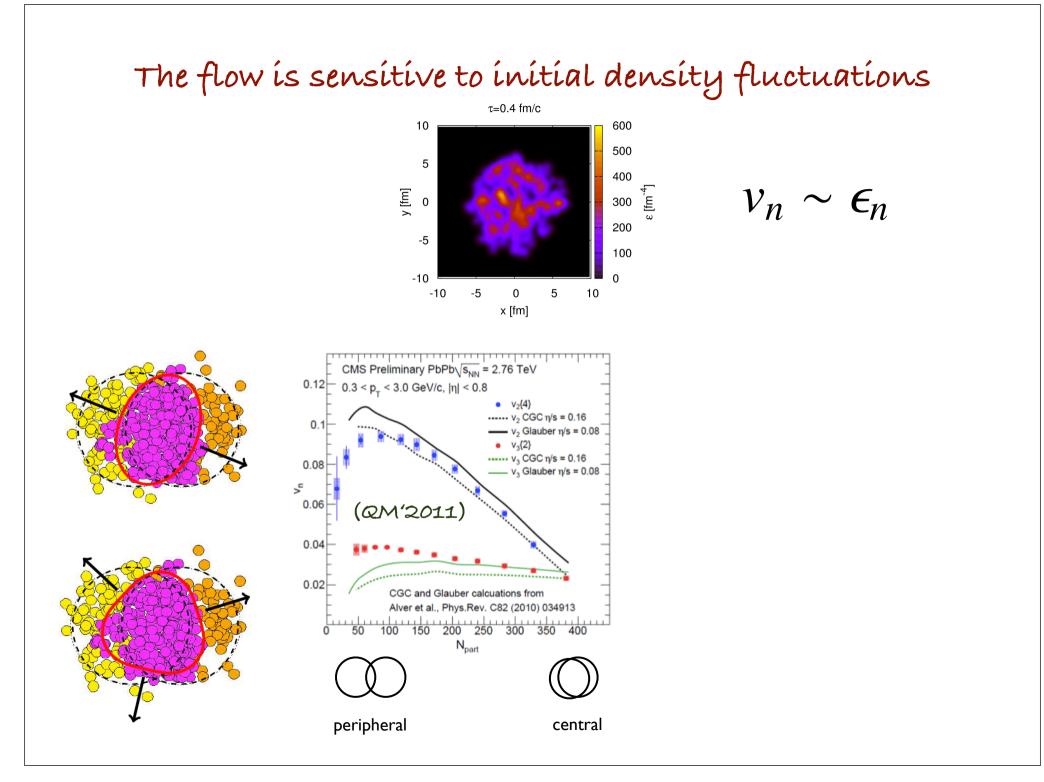
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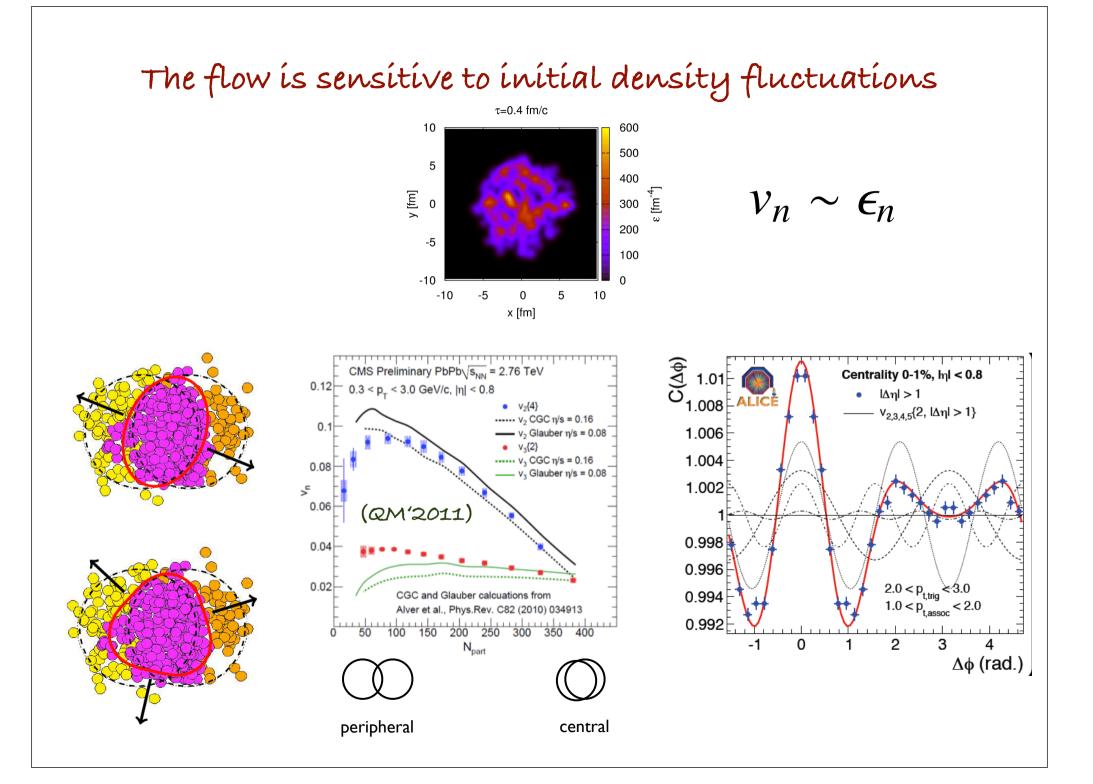


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- Rich flow pattern, sensitivity to initial conditions
- Sensitivity to the equation of state? (Pt, (I/S)dN/dy)

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Plasma: soft and hard modes, particles and fields. Long wavelength modes can remain strongly coupled....

THERMALIZATION

- How do we go from the intial nuclear wave-functions to the locally equilibrated fluid seen in experiments ?

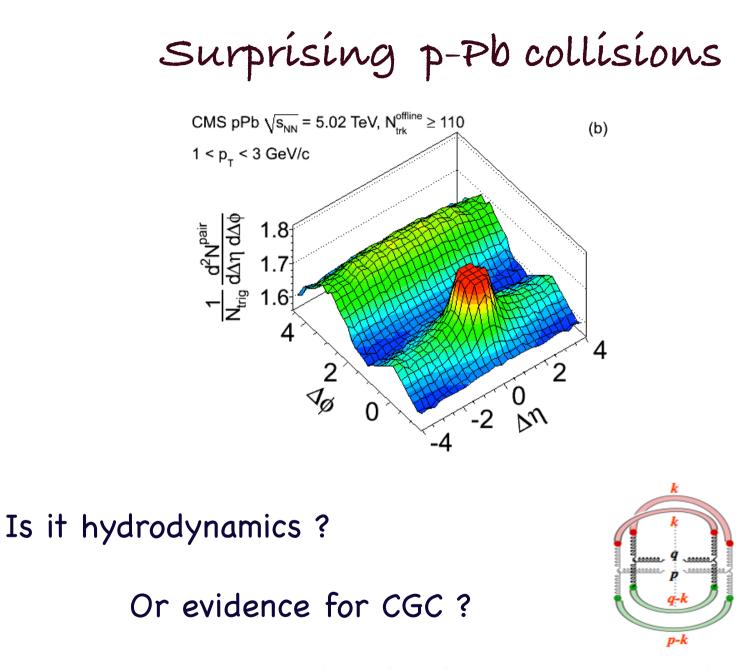
- What are the initial d.o.f.'s : partons ? color fields (CGC)? mixture of both ?

- Initial fields are typically unstable (e.g. if anisotric momentum distributions of particles). Instabilities provide 'fast' isotropization of momentum distributions

- Amplification of soft modes is a generic feature

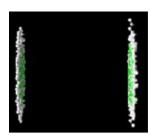
- CGC picture suggests an overpopulation of soft modes

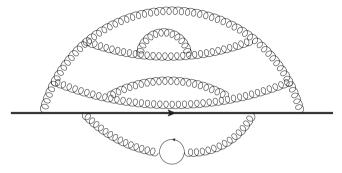
(for a summary see arXiv: 1203.2042)

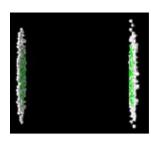


Dumitru, Dusling, Gelis, Jalilian-Marian, Lappi, Venugopalan : 1009.5295 Dusling, Venugopalan:1211.3701

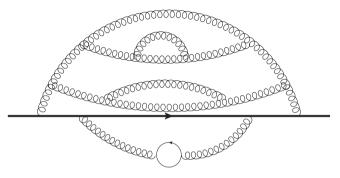
Moving backward in time Nucleí are made of densely packed gluons

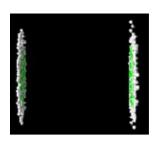




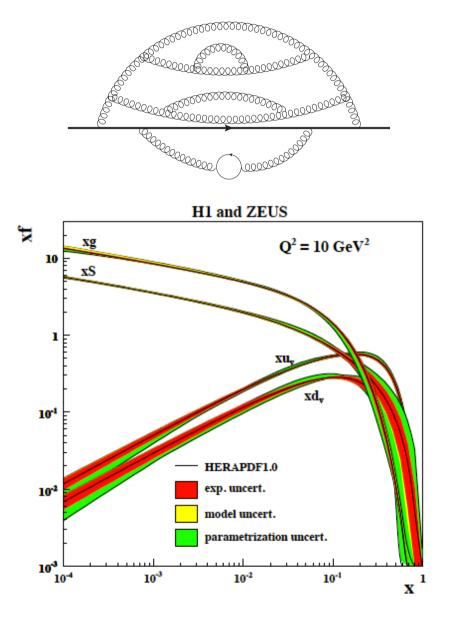


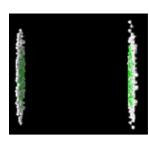
In a collision at high energy, one 'sees' mostly the gluons in the nuclei





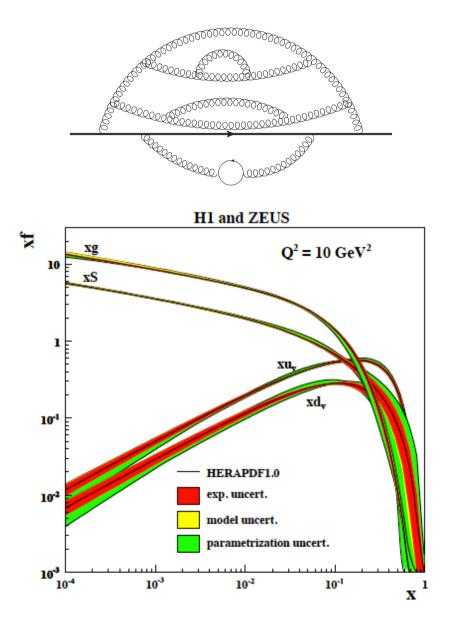
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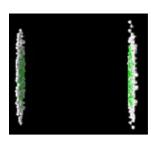




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Gluon density increases with energy (with decreasing x, increasing Q)



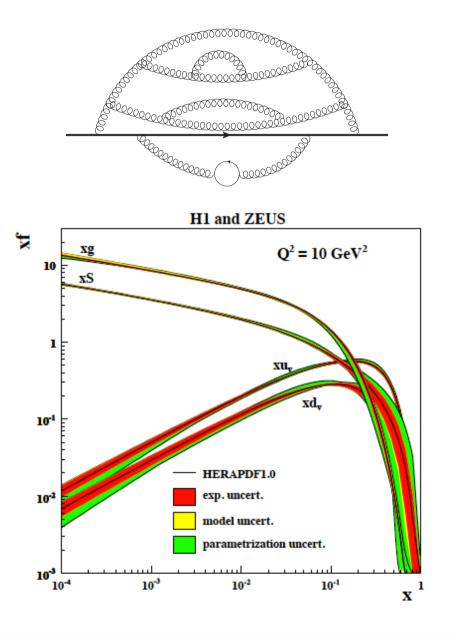


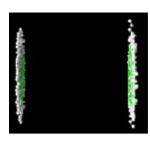
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Bulk of particle production ( p_{\rm T} \lesssim 2~{\rm GeV} )
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RHIC $(\sqrt{s} = 200 \text{ GeV})$ $x \sim 10^{-2}$ LHC $(\sqrt{s} = 5.5 \text{ TeV})$ $x \sim 4 \times 10^{-4}$





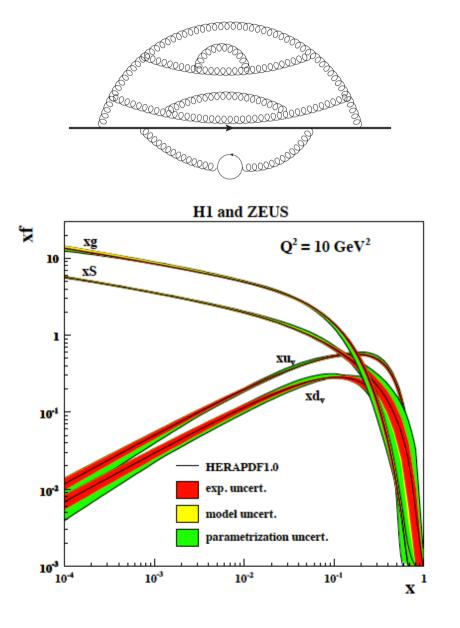
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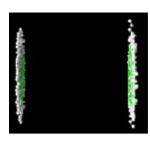
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Evolution equations describe the evolution with energy of relevant configurations (DGLAP, BFKL, JIMWLK...)





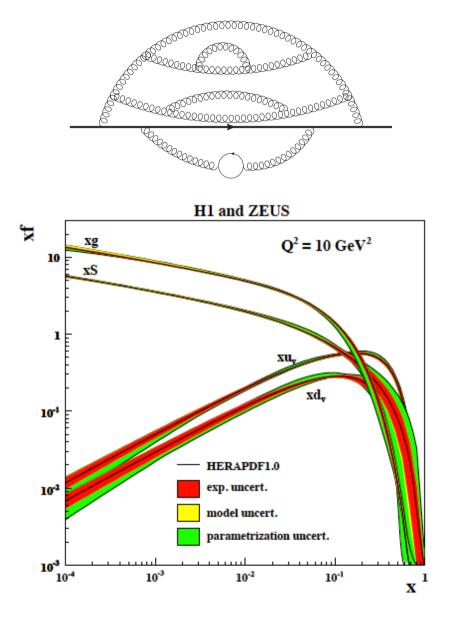
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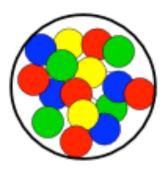
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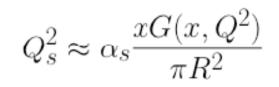
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The growth eventually saturates



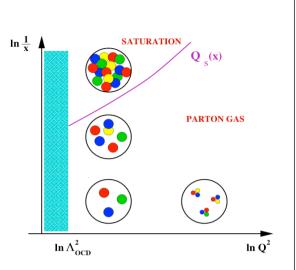
Saturation momentum





At saturation, occupation numbers are large

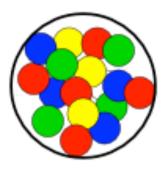
$$\frac{xG(x,Q^2)}{\pi R^2 Q_s^2} \sim \frac{1}{\alpha_s}$$



$$Q_s^2(x,A) \simeq Q_0^2 A^{1/3} \left(\frac{x_0}{x}\right)^{\lambda}$$

 $\lambda=0.2\div0.3$

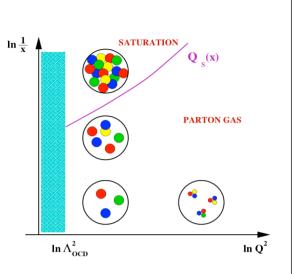
Saturation momentum



$$Q_s^2 \approx \alpha_s \frac{x G(x,Q^2)}{\pi R^2}$$

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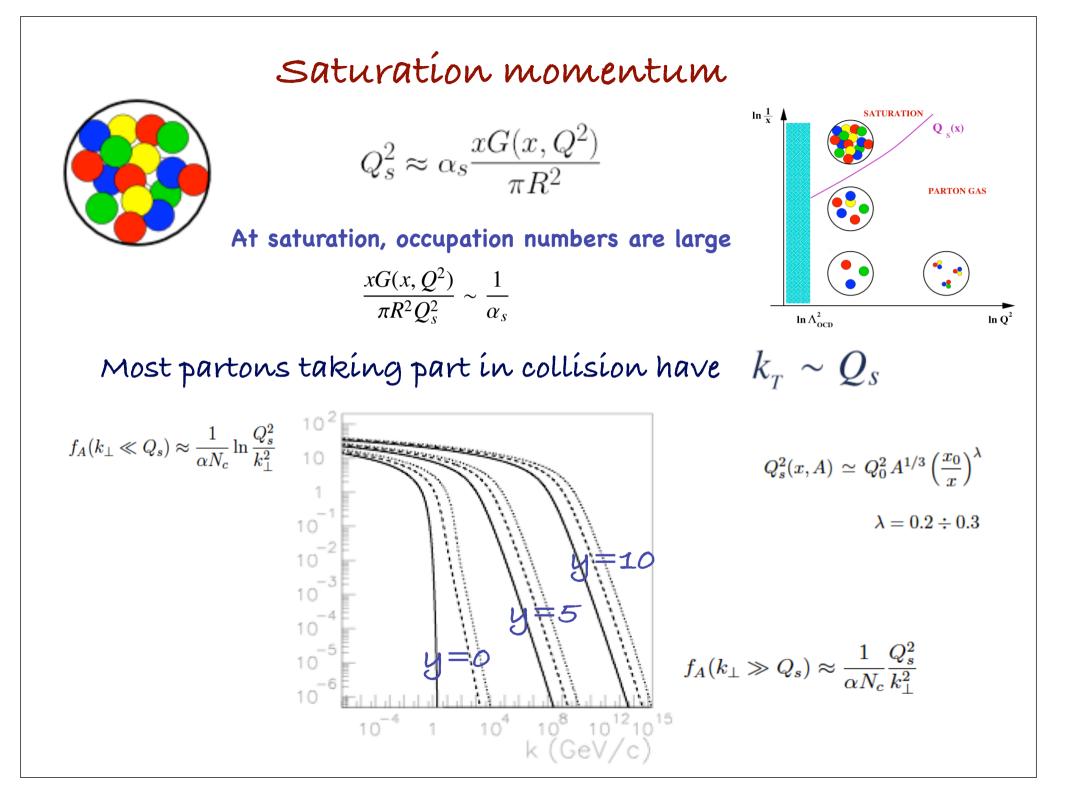
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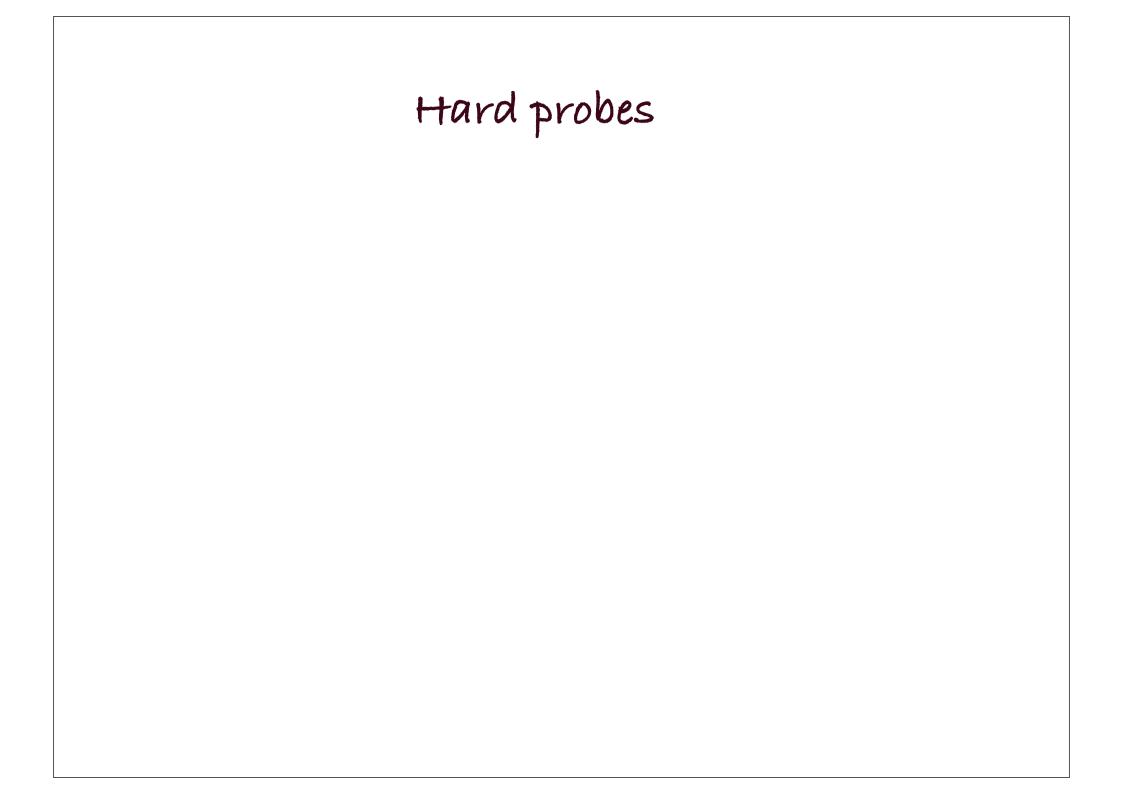
Most partons taking part in collision have $k_{ au} \sim Q_s$

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Moving backward in time Signals from the early stages Hard probes



Hard probes

Hard probes are produced on short space time scales, and their production rate can be calculated from pQCD

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Hard probes are like test particles. The study of their propagation provides much information about the medium in which they propagate.

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Examples of hard probes: heavy quarks, quarkonia, photons, Z and W, jets...

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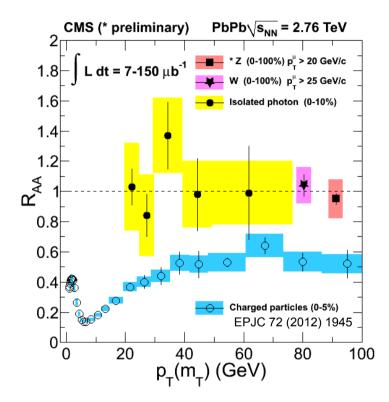
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Prospects for hard probes at the LHC are truly fascinating

hard processes are under control



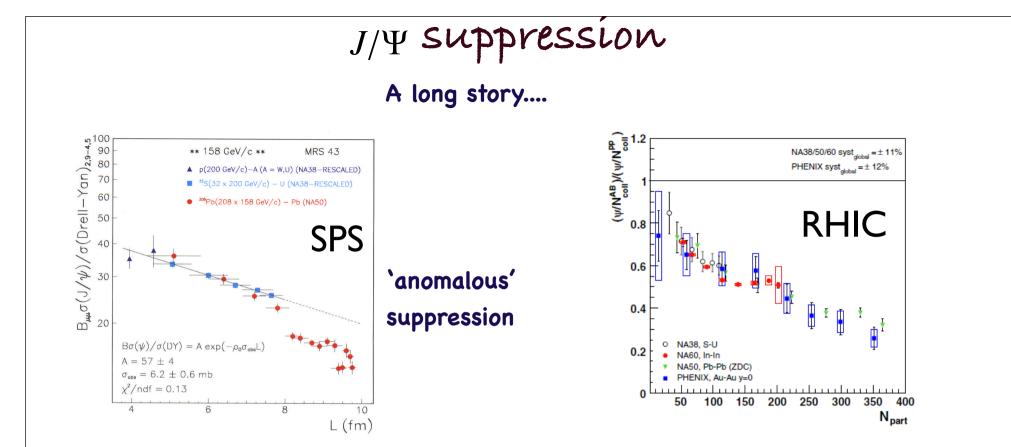
Hard processes are not affected by the nuclear environment, as expected.

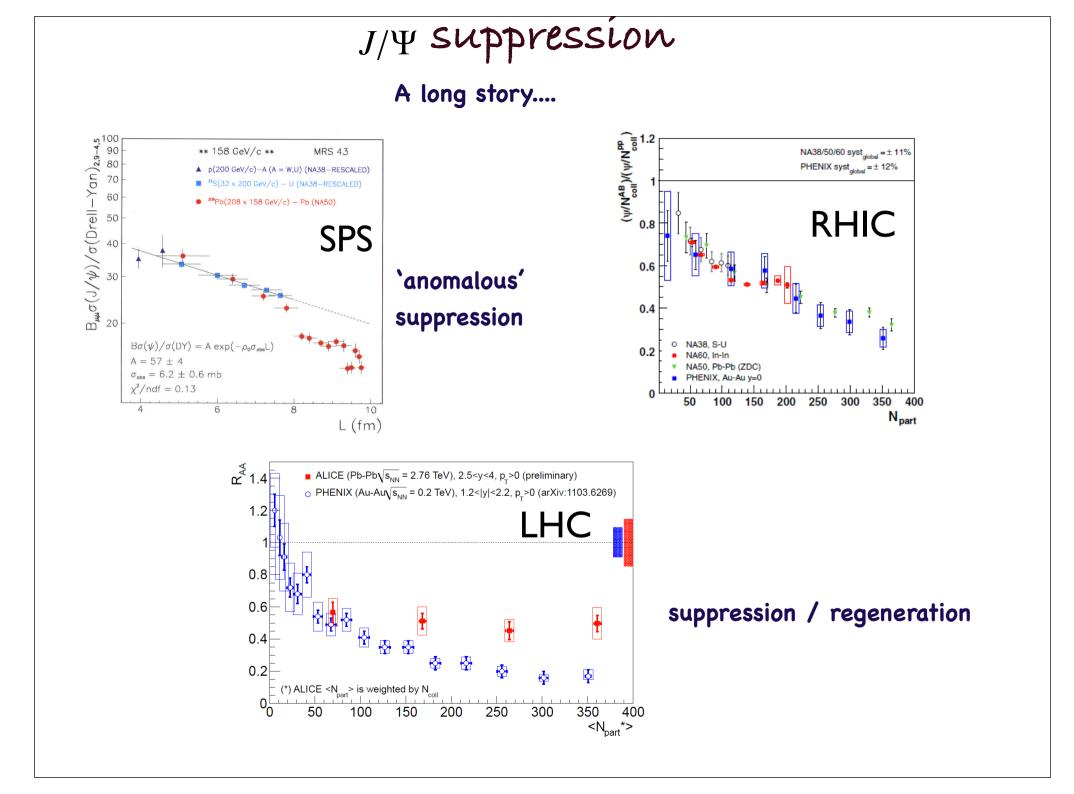


 J/Ψ suppression

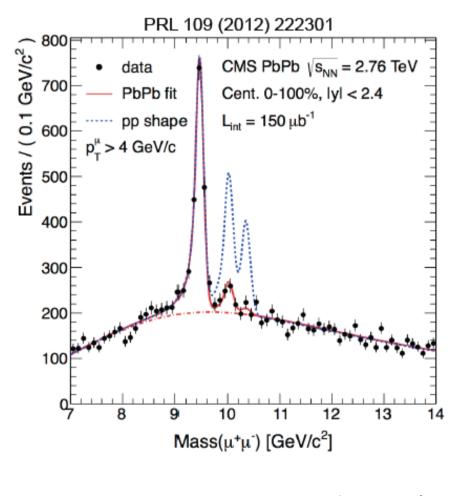
J/Ψ SUPPression A long story....

J/Ψ Suppression A long story.... $\begin{array}{cccc} \mathbb{B}_{\mu\mu}\sigma(J/\psi)/\sigma(\mathsf{Drell}-\mathsf{Yan})_{2;9-4,5}\\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \end{array} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \end{array} & \mathbb{C} \\ \end{array} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \end{array} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} & \mathbb{C} \\ \end{array} & \mathbb{C} \\ \end{array} & \mathbb{C} \\ \end{array} & \mathbb{C} \\ \end{array} & \mathbb{C} & \mathbb{C}$ ** 158 GeV/c ** **MRS 43** ▲ p(200 GeV/c) - A (A = W,U) (NA38 - RESCALED)■ ³²S(32 x 200 GeV/c) - U (NA38-RESCALED) • ²⁰⁸Pb(208 x 158 GeV/c) - Pb (NA50) SPS `anomalous' suppression $B\sigma(\psi)/\sigma(DY) = A \exp(-\rho_0 \sigma_{obs}L)$ $A = 57 \pm 4$ $\sigma_{\rm abs} = 6.2 \pm 0.6 \, {\rm mb}$ $\chi^{2}/ndf = 0.13$ 1 6 8 10 L (fm)

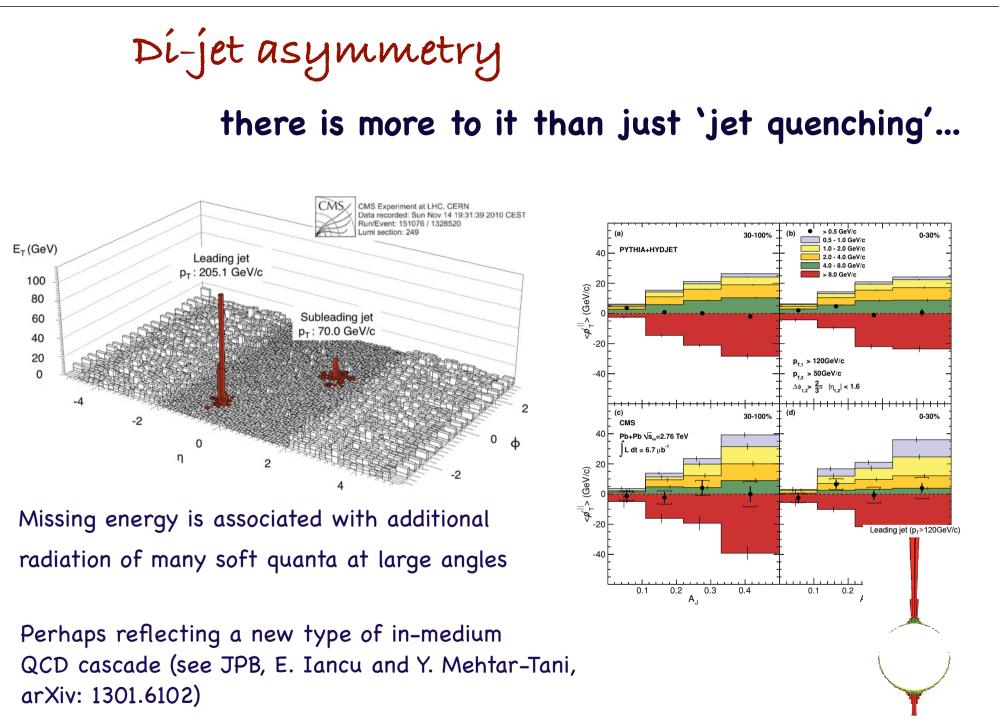




Ysuppression



excited states are more 'fragile'....



Subleading jet (p_T>30GeV/c)







A quark-gluon plasma is produced in ultra-relativistic heavy ion collisions, whose global properties do not seem to change much between RHIC and LHC (a liquid with low relative viscosity)



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The field has never been so exciting as now !