

Isospin breaking corrections to light hadron masses

Antonin Portelli

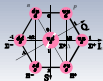
Budapest-Marseille-Wuppertal collaboration

14th of May 2012 – RPP 2012, Montpellier

Centre de Physique Théorique, Marseille, France

- 1 Motivations
- 2 Simulation setup & methodology
- 3 Preliminary results
- 4 Epilogue

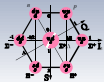
Motivations



Isospin symmetry breaking

Isospin symmetry is broken because :

	u	d
Mass [PDG 2012]	$2.5 \left(\begin{smallmatrix} +0.6 \\ -0.8 \end{smallmatrix} \right)$	$5.0 \left(\begin{smallmatrix} +0.7 \\ -0.9 \end{smallmatrix} \right)$
Charge	$\frac{2}{3}e$	$-\frac{1}{3}e$

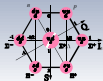


Isospin symmetry breaking

Isospin symmetry is broken because :

- up and down quark masses are different (strong breaking)

	u	d
Mass [PDG 2012]	$2.5 \left(\begin{smallmatrix} +0.6 \\ -0.8 \end{smallmatrix} \right)$	$5.0 \left(\begin{smallmatrix} +0.7 \\ -0.9 \end{smallmatrix} \right)$
Charge	$\frac{2}{3}e$	$-\frac{1}{3}e$

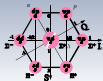


Isospin symmetry breaking

Isospin symmetry is broken because :

- up and down quark masses are different (**strong breaking**)
- up and down quark electric charges are different (**EM breaking**)

	u	d
Mass [PDG 2012]	$2.5 \left(\begin{smallmatrix} +0.6 \\ -0.8 \end{smallmatrix} \right)$	$5.0 \left(\begin{smallmatrix} +0.7 \\ -0.9 \end{smallmatrix} \right)$
Charge	$\frac{2}{3}e$	$-\frac{1}{3}e$



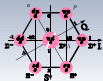
Isospin symmetry breaking

Isospin symmetry is broken because :

- up and down quark masses are different (**strong breaking**)
- up and down quark electric charges are different (**EM breaking**)

	u	d
Mass [PDG 2012]	$2.5 \left(\begin{smallmatrix} +0.6 \\ -0.8 \end{smallmatrix} \right)$	$5.0 \left(\begin{smallmatrix} +0.7 \\ -0.9 \end{smallmatrix} \right)$
Charge	$\frac{2}{3}e$	$-\frac{1}{3}e$

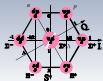
This breaking implies **mass splittings** in isospin multiplets.



Isospin breaking parameters

- EM breaking parameter :

fine-structure constant $\alpha \simeq 0.0073$



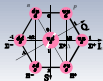
Isospin breaking parameters

- EM breaking parameter :

fine-structure constant $\alpha \simeq 0.0073$

- strong breaking parameter :

light quark mass splitting over a typical QCD scale $\frac{m_d - m_u}{\Lambda_{\text{QCD}}} \lesssim 0.01$



Isospin breaking parameters

- EM breaking parameter :

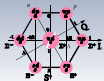
fine-structure constant $\alpha \simeq 0.0073$

- strong breaking parameter :

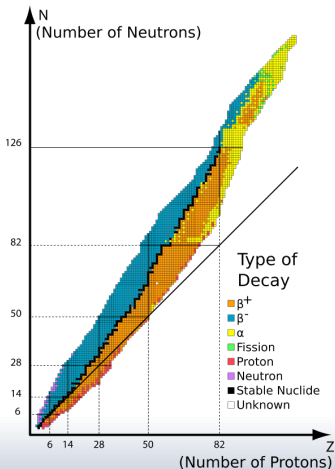
light quark mass splitting over a typical QCD scale $\frac{m_d - m_u}{\Lambda_{\text{QCD}}} \lesssim 0.01$

Isospin breaking effects

Sum of two little effects of the same order ($\sim 1\%$), eventually competing.



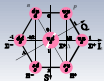
Nucleon mass splitting



Nucleon mass splitting is experimentally very well known :

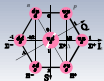
$$M_p - M_n = -1.29333214(43) \text{ MeV}$$

Although it is a 1‰ effect, it determines through β decay **the stable nuclides spectrum**.



Nucleon mass splitting

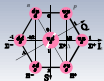
Ab-initio nucleon mass splitting prediction from QCD+QED **is still an open problem.**



Nucleon mass splitting

Ab-initio nucleon mass splitting prediction from QCD+QED **is still an open problem.**

Lattice QCD could give a way to solve numerically this problem.



Nucleon mass splitting

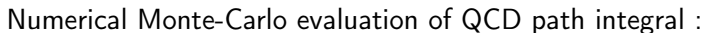
Ab-initio nucleon mass splitting prediction from QCD+QED **is still an open problem.**

Lattice QCD could give a way to solve numerically this problem.

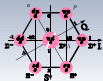
Predicting a 1‰ effect through lattice simulation is a **considerable computing challenge.**



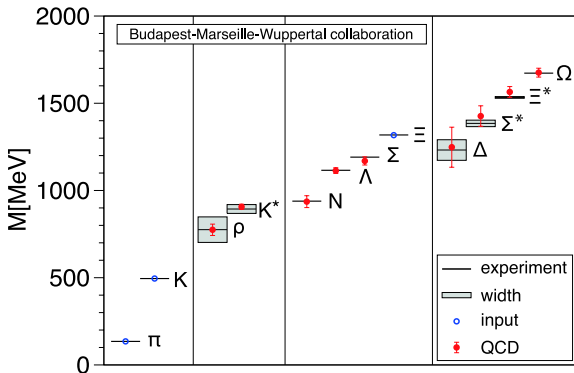
Lattice QCD



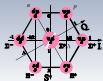
$$\langle O \rangle = \frac{1}{\mathcal{Z}} \int \mathcal{D}U_\mu O_{\text{Wick}}[D^{-1}] \det(D) \exp(-S_{\text{gauge}})$$



Light QCD spectrum solved

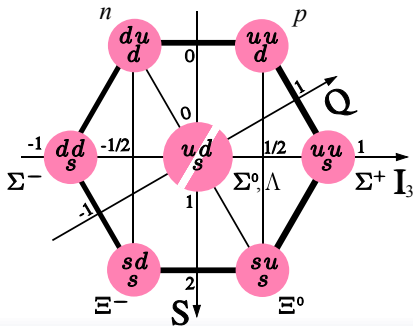


[BMWc 2008, Science, hep-lat/0906.3599]



Goal: octet baryon mass splittings

There are 3 stable baryon multiplets formed with u, d and s quarks :



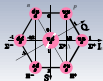
Mass splittings are experimentally known [PDG 2012] :

$$M_p - M_n = -1.29333214(43) \text{ MeV}$$

$$M_{\Sigma^+} - M_{\Sigma^-} = -8.08(08) \text{ MeV}$$

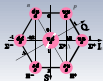
$$M_{\Xi^0} - M_{\Xi^-} = -6.85(21) \text{ MeV}$$

Simulation setup & methodology



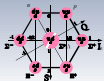
Actions

- $N_f = 2 + 1$ QCD simulations with Lüscher-Weisz gauge action, tree level $O(a)$ -improved Wilson fermions and two steps of HEX smearing.
[BMWc 2010, JHEP, hep-lat/1011.2711]



Actions

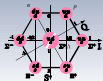
- $N_f = 2 + 1$ QCD simulations with Lüscher-Weisz gauge action, tree level $O(a)$ -improved Wilson fermions and two steps of HEX smearing.
[BMWc 2010, JHEP, hep-lat/1011.2711]
- Non-compact Maxwell action with Coulomb gauge fixing and zero mode subtraction.
[BMWc 2010, Lattice 2010, hep-lat/1011.4189]
[RPP 2011]



Actions

- $N_f = 2 + 1$ QCD simulations with Lüscher-Weisz gauge action, tree level $O(a)$ -improved Wilson fermions and two steps of HEX smearing.
[BMWc 2010, JHEP, hep-lat/1011.2711]
- Non-compact Maxwell action with Coulomb gauge fixing and zero mode subtraction.
[BMWc 2010, Lattice 2010, hep-lat/1011.4189]
[RPP 2011]
- At quark propagator computation time, we phase $SU(3)$ links using a generated electromagnetic field A_μ :

$$U_\mu \longmapsto \exp(iQ_q e A_\mu) U_\mu$$

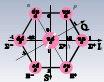


Actions

- $N_f = 2 + 1$ QCD simulations with Lüscher-Weisz gauge action, tree level $O(a)$ -improved Wilson fermions and two steps of HEX smearing.
[BMWc 2010, JHEP, hep-lat/1011.2711]
- Non-compact Maxwell action with Coulomb gauge fixing and zero mode subtraction.
[BMWc 2010, Lattice 2010, hep-lat/1011.4189]
[RPP 2011]
- At quark propagator computation time, we phase $SU(3)$ links using a generated electromagnetic field A_μ :

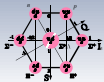
$$U_\mu \longmapsto \exp(iQ_q e A_\mu) U_\mu$$

- Thus, simulation is **quenched** in QED (neutral sea quarks).



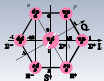
Gauge fields

- five lattice spacings: from 0.12 fm to 0.05 fm.



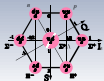
Gauge fields

- five lattice spacings: from 0.12 fm to 0.05 fm.
- 36 sea pion masses from 450 MeV to 128 MeV.



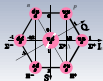
Gauge fields

- five lattice spacings: from 0.12 fm to 0.05 fm.
- 36 sea pion masses from 450 MeV to 128 MeV.
- Isospin symmetry: $m_u^s = m_d^s$.



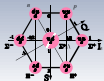
Gauge fields

- five lattice spacings: from 0.12 fm to 0.05 fm.
- 36 sea pion masses from 450 MeV to 128 MeV.
- Isospin symmetry: $m_u^s = m_d^s$.
- sea strange quark masses bracketing the physical value.



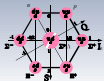
Gauge fields

- five lattice spacings: from 0.12 fm to 0.05 fm.
- 36 sea pion masses from 450 MeV to 128 MeV.
- Isospin symmetry: $m_u^s = m_d^s$.
- sea strange quark masses bracketing the physical value.
- 16 volumes from $(2 \text{ fm})^3$ to $(6 \text{ fm})^3$ with $M_\pi L > 4$ (negligible QCD finite volume effects).



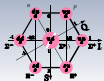
Gauge fields

- five lattice spacings: from 0.12 fm to 0.05 fm.
- 36 sea pion masses from 450 MeV to 128 MeV.
- Isospin symmetry: $m_u^s = m_d^s$.
- sea strange quark masses bracketing the physical value.
- 16 volumes from $(2 \text{ fm})^3$ to $(6 \text{ fm})^3$ with $M_\pi L > 4$ (negligible QCD finite volume effects).
- EM fields generated with α set to its physical value.



Observables

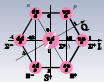
u, d, s quark propagators are computed on background gauge fields.



Observables

u, d, s quark propagators are computed on background gauge fields.

Valence u and s masses are set equal to sea masses.

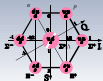


Observables

u, d, s quark propagators are computed on background gauge fields.

Valence u and s masses are set equal to sea masses.

Valence d quark mass is set in two different ways :



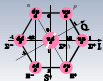
Observables

u, d, s quark propagators are computed on background gauge fields.

Valence u and s masses are set equal to sea masses.

Valence d quark mass is set in two different ways :

- equal to the sea d quark mass
(**mass isospin dataset**).



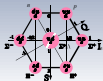
Observables

u, d, s quark propagators are computed on background gauge fields.

Valence u and s masses are set equal to sea masses.

Valence d quark mass is set in two different ways :

- equal to the sea d quark mass
(**mass isospin dataset**).
- $m_d^{\text{val.}} = m_d^{\text{sea}} + \varepsilon$ with ε near the physical value of $m_d - m_u$
(**physical dataset**).



Observables

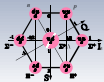
u, d, s quark propagators are computed on background gauge fields.

Valence u and s masses are set equal to sea masses.

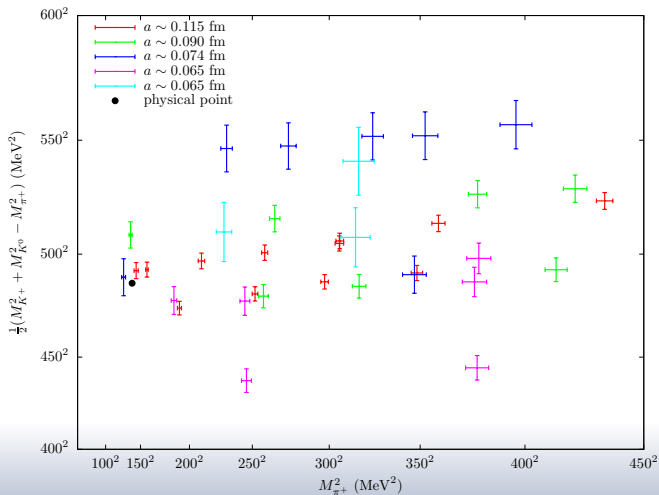
Valence d quark mass is set in two different ways :

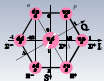
- equal to the sea d quark mass
(**mass isospin dataset**).
- $m_d^{\text{val.}} = m_d^{\text{sea}} + \varepsilon$ with ε near the physical value of $m_d - m_u$
(**physical dataset**).

Quark propagators are contracted to form hadron propagators.

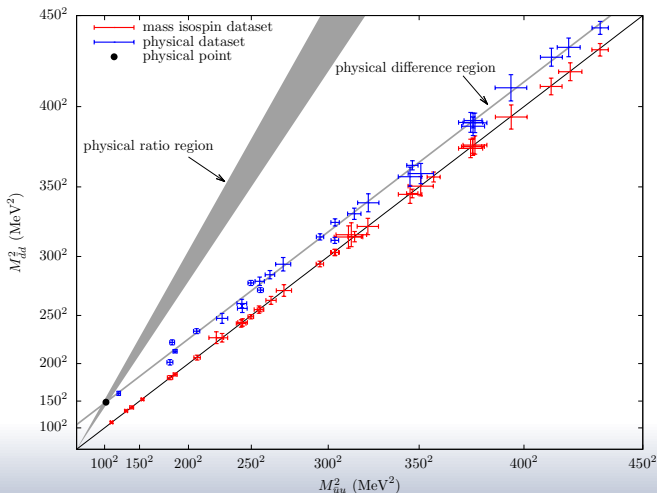


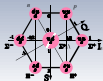
Simulation landscape





Simulation landscape



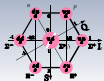


Propagator fits

Zero momentum propagators of two isospin partners X_1 and X_2 are fitted in a combined way with the model :

$$\begin{cases} G_{X_1}(t) = A_1 \exp[-(M_X + \Delta M_X) \frac{t}{2}] \\ G_{X_2}(t) = A_2 \exp[-(M_X - \Delta M_X) \frac{t}{2}] \end{cases}$$

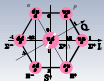
The parameter ΔM_X is the mass splitting $M_{X_1} - M_{X_2}$.



Interpolation model

ΔM_X can be expanded at NLO in isospin breaking parameters :

$$\Delta M_X = \alpha A_{\text{EM}} M_X + (m_u - m_d) A_{\text{QCD}} M_X$$

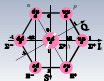


Interpolation model

ΔM_X can be expanded at NLO in isospin breaking parameters :

$$\Delta M_X = \alpha A_{\text{EM}} M_X + (m_u - m_d) A_{\text{QCD}} M_X$$

- α is already set to its physical value.

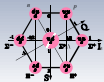


Interpolation model

ΔM_X can be expanded at NLO in isospin breaking parameters :

$$\Delta M_X = \alpha A_{\text{EM}} M_X + (m_u - m_d) A_{\text{QCD}} M_X$$

- α is already set to its physical value.
- $m_u - m_d$ can be fixed using $\Delta M_K^2 = M_{K^+}^2 - M_{K^0}^2$.

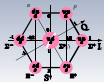


Interpolation model

ΔM_X can be expanded at NLO in isospin breaking parameters :

$$\Delta M_X = \alpha A_{\text{EM}} M_X + (m_u - m_d) A_{\text{QCD}} M_X$$

- α is already set to its physical value.
- $m_u - m_d$ can be fixed using $\Delta M_K^2 = M_{K^+}^2 - M_{K^0}^2$.
- $A_{\text{EM,QCD}}$ still depend on isospin symmetric parameters of QCD+QED : m_{ud}, m_s, L and a

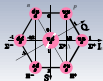


Interpolation model

ΔM_X can be expanded at NLO in isospin breaking parameters :

$$\Delta M_X = \alpha A_{\text{EM}} M_X + (m_u - m_d) A_{\text{QCD}} M_X$$

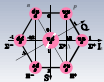
- α is already set to its physical value.
- $m_u - m_d$ can be fixed using $\Delta M_K^2 = M_{K^+}^2 - M_{K^0}^2$.
- $A_{\text{EM,QCD}}$ still depend on isospin symmetric parameters of QCD+QED : m_{ud}, m_s, L and a
- Due to our setup, distinction between sea and valence is a NNLO isospin breaking effect.



Interpolation model

$A_{\text{EM,QCD}}$ can then be expanded in the isospin symmetric parameters:

$$A_{\text{EM,QCD}} = A^\phi + P_{ud}(m_{ud} - m_{ud}^\phi) + P_s(m_s - m_s^\phi) + P_{\text{FV}}\frac{1}{L} + P_D a$$

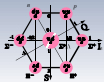


Interpolation model

$A_{\text{EM,QCD}}$ can then be expanded in the isospin symmetric parameters:

$$A_{\text{EM,QCD}} = A^\phi + P_{ud}(m_{ud} - m_{ud}^\phi) + P_s(m_s - m_s^\phi) + P_{\text{FV}}\frac{1}{L} + P_{Da}$$

- m_{ud} can be fixed using $M_{\pi^+}^2$.

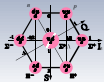


Interpolation model

$A_{\text{EM,QCD}}$ can then be expanded in the isospin symmetric parameters:

$$A_{\text{EM,QCD}} = A^\phi + P_{ud}(m_{ud} - m_{ud}^\phi) + P_s(m_s - m_s^\phi) + P_{\text{FV}}\frac{1}{L} + P_{Da}$$

- m_{ud} can be fixed using $M_{\pi^+}^2$.
- m_s can be fixed using $M_{K^+}^2 + M_{K^0}^2 - M_{\pi^+}^2$.

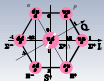


Interpolation model

$A_{\text{EM,QCD}}$ can then be expanded in the isospin symmetric parameters:

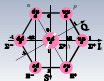
$$A_{\text{EM,QCD}} = A^\phi + P_{ud}(m_{ud} - m_{ud}^\phi) + P_s(m_s - m_s^\phi) + P_{\text{FV}}\frac{1}{L} + P_D a$$

- m_{ud} can be fixed using $M_{\pi^+}^2$.
- m_s can be fixed using $M_{K^+}^2 + M_{K^0}^2 - M_{\pi^+}^2$.
- Power finite volume corrections are expected with EM interaction.



Error estimation

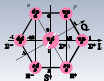
Each splitting is computed by choosing between:



Error estimation

Each splitting is computed by choosing between:

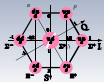
- 2 different propagator fit intervals per lattice spacing
(for a total of 32 possible choices)



Error estimation

Each splitting is computed by choosing between:

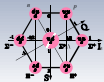
- 2 different propagator fit intervals per lattice spacing (for a total of 32 possible choices)
- 2 different M_{π^+} high cuts in the splitting fit (350 and 450 MeV)



Error estimation

Each splitting is computed by choosing between:

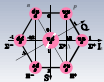
- 2 different propagator fit intervals per lattice spacing (for a total of 32 possible choices)
- 2 different M_{π^+} high cuts in the splitting fit (350 and 450 MeV)
- 2 different particles for lattice spacing determination (Ω^- and Ξ)



Error estimation

Each splitting is computed by choosing between:

- 2 different propagator fit intervals per lattice spacing (for a total of 32 possible choices)
- 2 different M_{π^+} high cuts in the splitting fit (350 and 450 MeV)
- 2 different particles for lattice spacing determination (Ω^- and Ξ)
- 2 different M_{π^+} high cuts in lattice spacing determination fit (350 and 450 MeV)

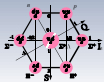


Error estimation

Each splitting is computed by choosing between:

- 2 different propagator fit intervals per lattice spacing (for a total of 32 possible choices)
- 2 different M_{π^+} high cuts in the splitting fit (350 and 450 MeV)
- 2 different particles for lattice spacing determination (Ω^- and Ξ)
- 2 different M_{π^+} high cuts in lattice spacing determination fit (350 and 450 MeV)

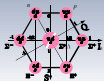
Total : 256 different analyses.



Error estimation

BMWc error evaluation :

Compute the histogram of the 256 fits with p-values as weights.

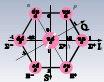


Error estimation

BMWc error evaluation :

Compute the histogram of the 256 fits with p-values as weights.

- statistical error: statistical error on the median

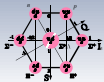


Error estimation

BMWc error evaluation :

Compute the histogram of the 256 fits with p-values as weights.

- statistical error: statistical error on the median
- systematic error: 1σ interval

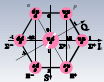


Error estimation

BMWc error evaluation :

Compute the histogram of the 256 fits with p-values as weights.

- statistical error: statistical error on the median
- systematic error: 1σ interval
- total error: quadratic sum of statistic and systematic errors



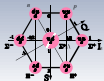
Error estimation

BMWc error evaluation :

Compute the histogram of the 256 fits with p-values as weights.

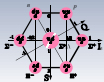
- statistical error: statistical error on the median
- systematic error: 1σ interval
- total error: quadratic sum of statistic and systematic errors

One additional source of systematics: **QED quenching**.



Error estimation

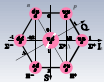
The sea EM contribution to a splitting is **not an isospin singlet**: only sea-valence interactions.



Error estimation

The sea EM contribution to a splitting is **not an isospin singlet**: only sea-valence interactions.

These interactions are proportional to the order e term of the fermionic determinant. With 3 sea flavors $\text{tr}(Q_s) = 0$, and this term is proportional to $\text{tr}(Q_s M_s) = \frac{1}{3}(m_{ud}^s - m_s^s)$.

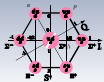


Error estimation

The sea EM contribution to a splitting is **not an isospin singlet**: only sea-valence interactions.

These interactions are proportional to the order e term of the fermionic determinant. With 3 sea flavors $\text{tr}(Q_s) = 0$, and this term is proportional to $\text{tr}(Q_s M_s) = \frac{1}{3}(m_{ud}^s - m_s^s)$.

EM quenching errors are $\text{SU}(3)$ suppressed.



Error estimation

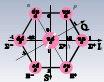
The sea EM contribution to a splitting is **not an isospin singlet**: only sea-valence interactions.

These interactions are proportional to the order e term of the fermionic determinant. With 3 sea flavors $\text{tr}(Q_s) = 0$, and this term is proportional to $\text{tr}(Q_s M_s) = \frac{1}{3}(m_{ud}^s - m_s^s)$.

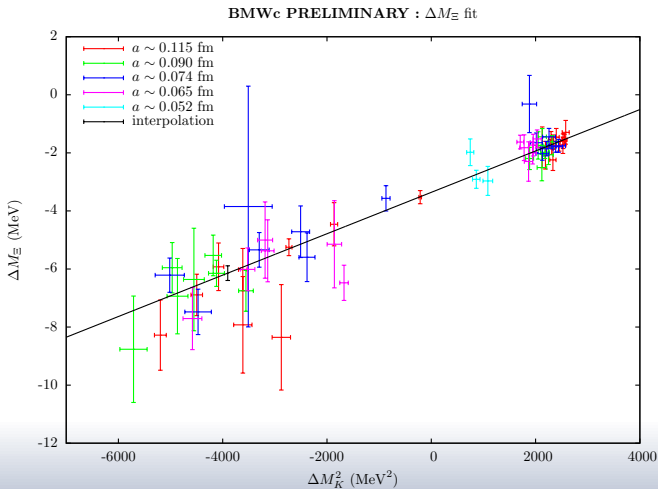
EM quenching errors are $\text{SU}(3)$ suppressed.

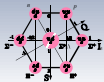
Typical relative error: $\frac{M_\Xi - M_N}{3M_N} \simeq 13\%$.

Preliminary results

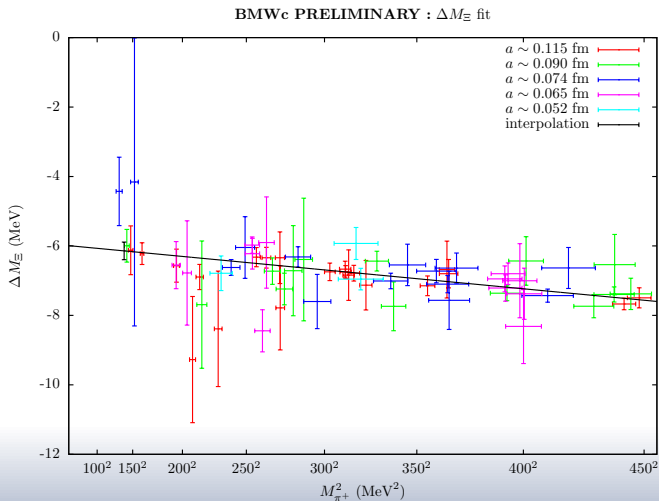


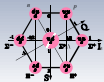
Ξ mass splitting



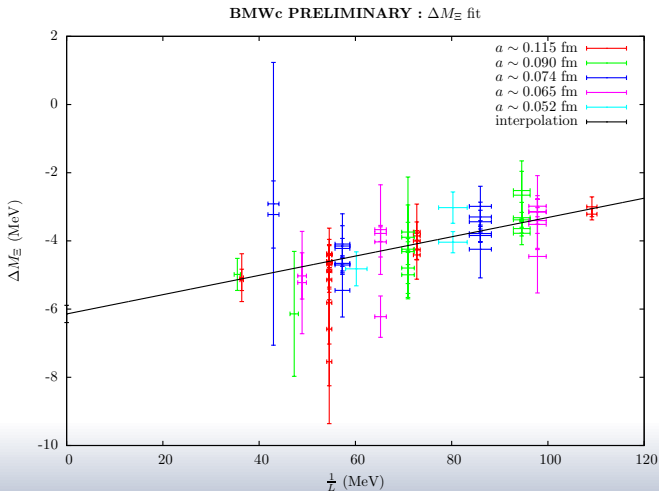


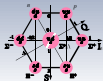
Ξ mass splitting



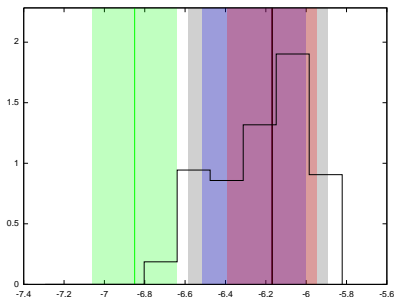


Ξ mass splitting



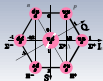


Ξ mass splitting

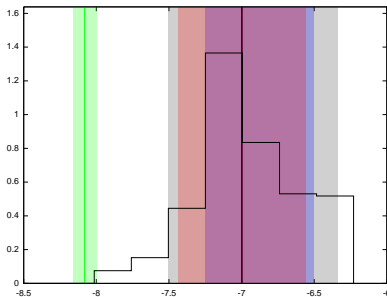


error: **systematic** **statistic** systematic + statistic **experimental**

$$M_{\Xi^0} - M_{\Xi^-} = -6.16(22)_{\text{stat.}} \left({}^{+16}_{-35} \right)_{\text{sys.}} (80)_{\text{quench.}} \left({}^{+85}_{-90} \right)_{\text{total}}$$

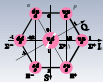


Σ mass splitting

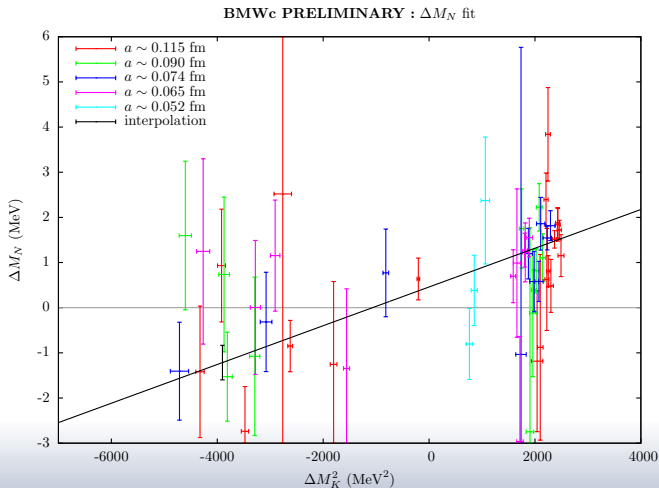


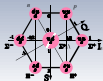
error: **systematic** **statistic** systematic + statistic **experimental**

$$M_{\Sigma^+} - M_{\Sigma^-} = -7.00(0.44)_{\text{stat.}} \left(\begin{smallmatrix} +0.49 \\ -0.25 \end{smallmatrix} \right)_{\text{sys.}} (0.91)_{\text{quench.}} \left(\begin{smallmatrix} +1.12 \\ -1.04 \end{smallmatrix} \right)_{\text{total}}$$

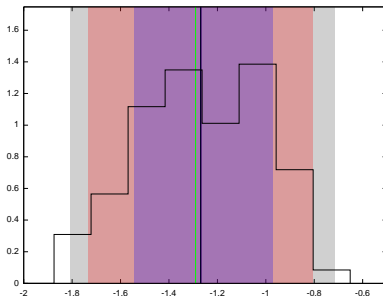


N mass splitting



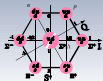


N mass splitting

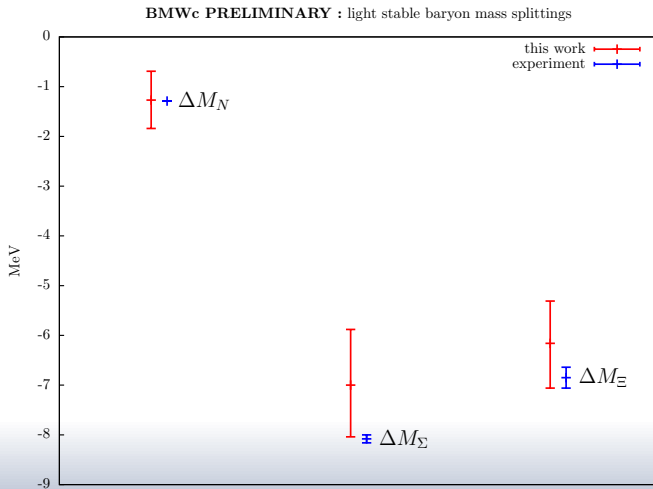


error: **systematic** **statistic** systematic + statistic **experimental**

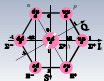
$$M_p - M_n = -1.27(46)_{\text{stat.}} \left(\begin{smallmatrix} +30 \\ -28 \end{smallmatrix} \right)_{\text{sys.}} (17)_{\text{quench.}} \left(\begin{smallmatrix} +58 \\ -57 \end{smallmatrix} \right)_{\text{total}}$$



Summary

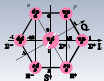


Epilogue



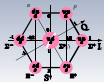
Conclusion

- The mass splittings of the baryon octet were computed with lattice QCD+(quenched)QED with masses down to the physical point.



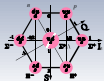
Conclusion

- The mass splittings of the baryon octet were computed with lattice QCD+(quenched)QED with masses down to the physical point.
- Statistical and systematic errors were estimated using usual BMWc strategy.



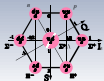
Conclusion

- The mass splittings of the baryon octet were computed with lattice QCD+(quenched)QED with masses down to the physical point.
- Statistical and systematic errors were estimated using usual BMWc strategy.
- Quenching errors were estimated using the $SU(3)$ suppression of the sea EM contribution to the splittings.



Conclusion

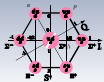
- The mass splittings of the baryon octet were computed with lattice QCD+(quenched)QED with masses down to the physical point.
- Statistical and systematic errors were estimated using usual BMWc strategy.
- Quenching errors were estimated using the $SU(3)$ suppression of the sea EM contribution to the splittings.
- With all sources of error taken into account, results are found compatibles with experiment.



Conclusion

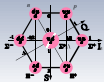
- The mass splittings of the baryon octet were computed with lattice QCD+(quenched)QED with masses down to the physical point.
- Statistical and systematic errors were estimated using usual BMWc strategy.
- Quenching errors were estimated using the $SU(3)$ suppression of the sea EM contribution to the splittings.
- With all sources of error taken into account, results are found compatibles with experiment.

Lattice QCD can have now **a significant sensitivity to 1‰ corrections** on hadronic energies.



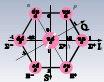
Perspectives

- Short term: isolate the **EM contribution to the baryon mass splittings** and investigate more the different sources of error on them.



Perspectives

- Short term: isolate the **EM contribution to the baryon mass splittings** and investigate more the different sources of error on them.
- Middle term: compute **light quark masses**.



Perspectives

- Short term: isolate the **EM contribution to the baryon mass splittings** and investigate more the different sources of error on them.
- Middle term: compute **light quark masses**.
- Long term: Explore ways to perform **full QCD+QED simulations**.

Thank you.

BMWc Collaboration

Budapest (Eötvös University)

S.D. Katz

Marseille (CPT)

J. Frison, L. Lellouch, A. Portelli, A. Sastre

Wuppertal (Bergische Universität)

S. Dürr, Z.I. Fodor, C. Hölbling, S. Krieg, T. Kurth, T. Lippert, C. Mc Neile,
K.K. Szabo