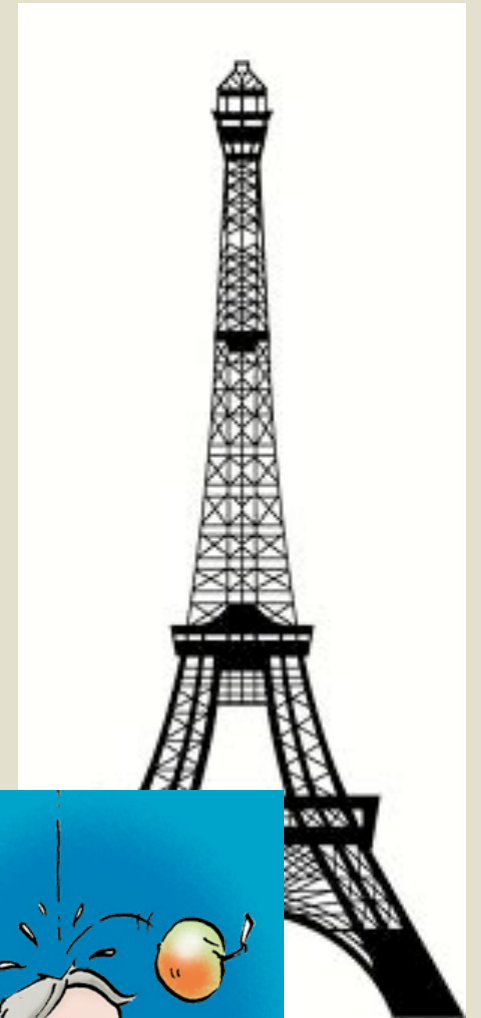


Understanding Gravity on Cosmic Scales

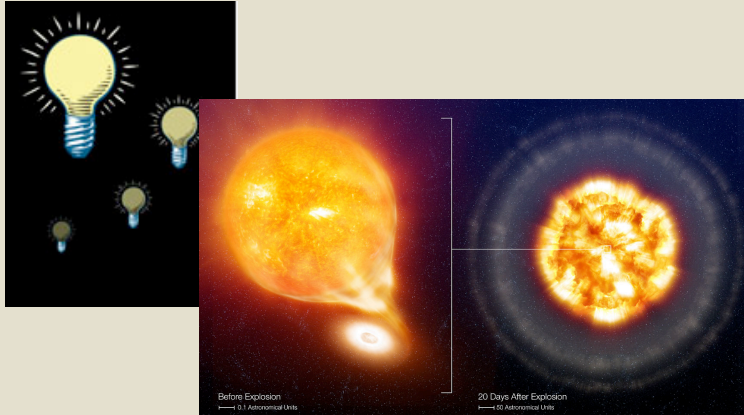
Rachel Bean
Cornell University

Work in collaboration with
Istvan Laszlo (Cornell)
Eva-Marie Mueller (Cornell)
Scott Watson (Syracuse)
Donnacha Kirk & Sarah Bridle (UCL)

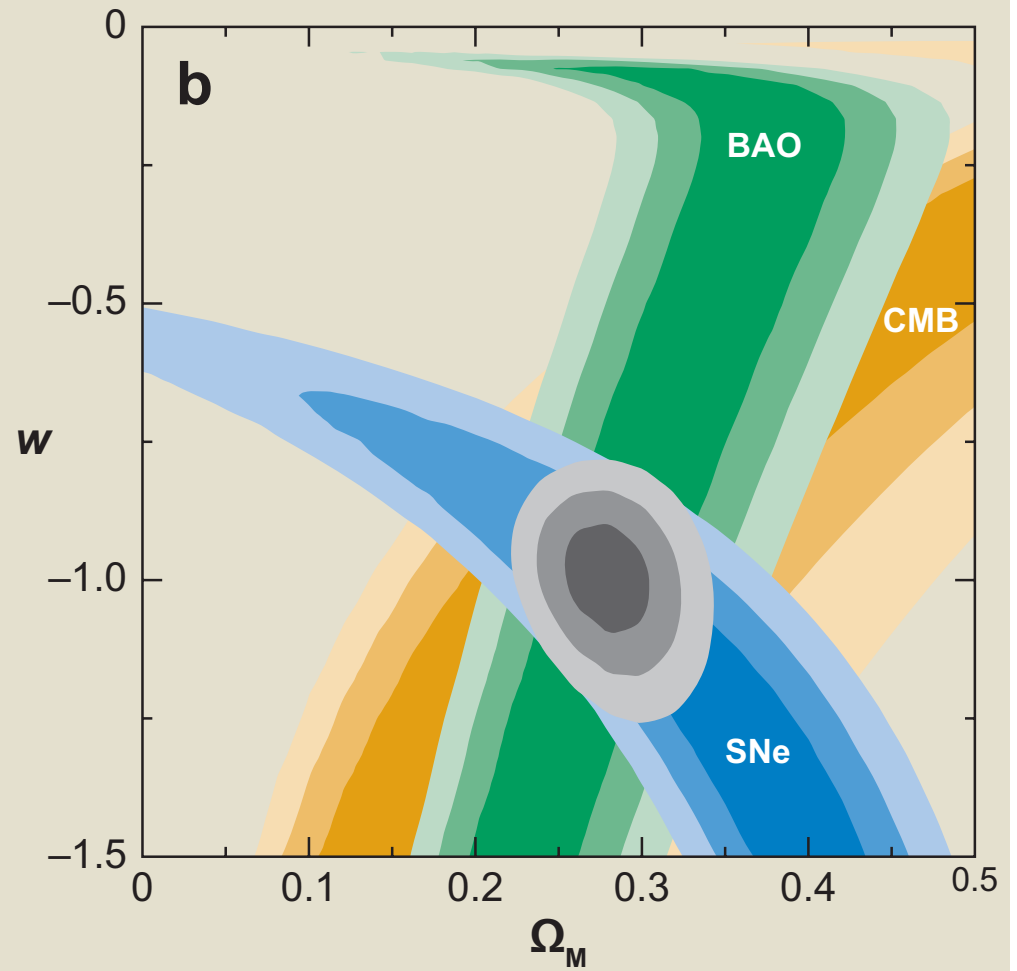
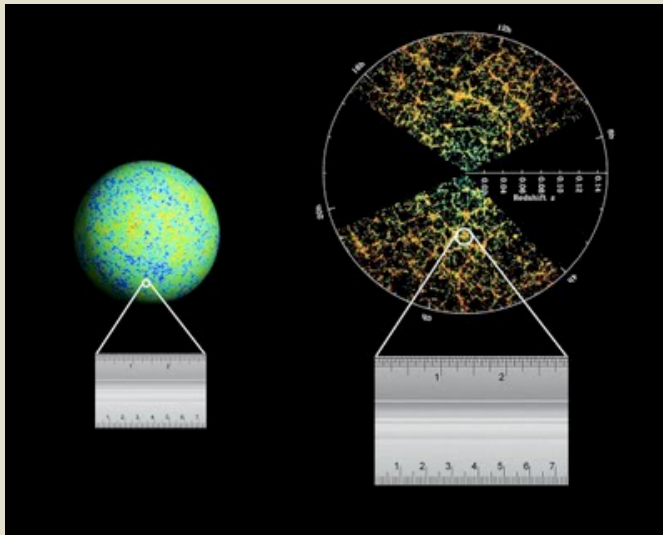


Cosmic geometry: Expansion history constraints

Standard candles



Standard rulers



Kowalski 2008

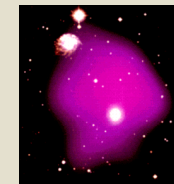
Understanding cosmic acceleration

Cosmic acceleration = a modification of Einstein's equations

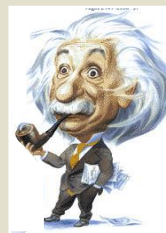


Deviations from GR?

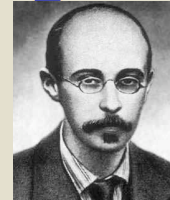
$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$



New matter?
interactions?



Λ ?



Inhomogeneous universe?

Broad aim = Phenomenology
Distinguish which sector: new gravity, new matter or Λ ?

Ambitious aim = Theoretical model
Learn something more about the underlying theory?

Distinguishing with expansion history

- Alter Friedmann and acceleration equations at late times

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho_m + 3P_m) + \textit{stuff}$$

or

$$\textit{stuff} + \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho_m + 3P_m) \quad ?$$

e.g. f(R) gravity

$$-H^2 f_R + \frac{a^2}{6} f + \frac{3}{2} H \dot{f}_R + \frac{1}{2} \ddot{f}_R + \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

e.g. DGP gravity

$$-\frac{\dot{H}}{r_c} + \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

There are benefits to asking more questions...



Three groups of extra galactic observations for testing gravity

I: Background expansion

CMB angular diameter distance

Supernovae luminosity distance

BAO angular/radial scale

II: Growth, up to some normalization

Galaxy autocorrelations

Galaxy – ISW x-corrln

Xray and SZ galaxy cluster measurements

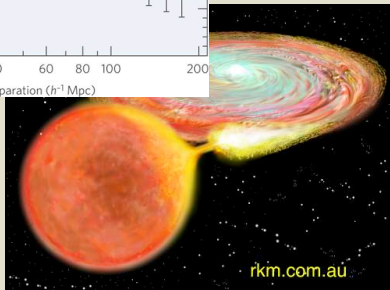
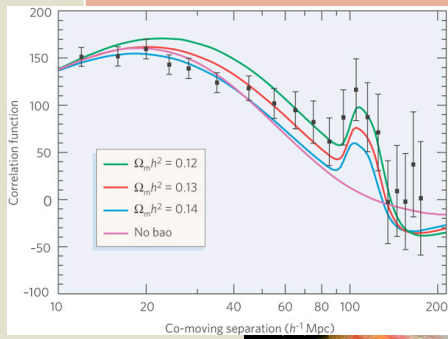
Ly-alpha measurements

III: Growth directly

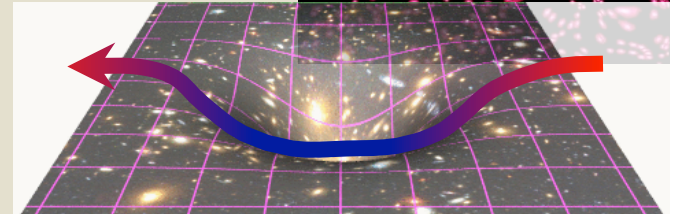
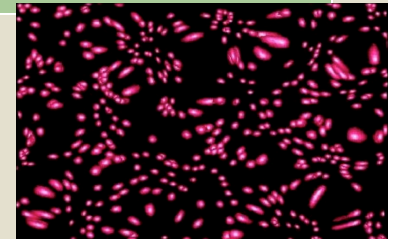
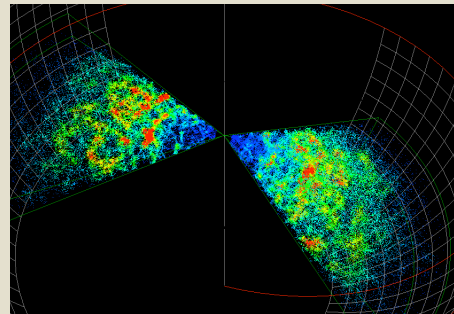
CMB ISW autocorrelation

Weak lensing autocorrelation

Peculiar velocity distribution/
bulk flows



rkm.com.au



Phenomenological model of modified gravity

- Perturbed metric $ds^2 = -(1 + 2\psi)dt^2 + a^2(1 - 2\phi)dx^2$
- Aim to describe phenomenological properties common to theories
- A modification to Poisson's equation, Q

$$k^2\phi = -4\pi G Q a^2 \rho \Delta$$

$Q \neq 1$: can be mimicked by additional (dark energy?) perturbations, or modified dark matter evolution

- An inequality between Newton's potentials, R

$$\psi = R\phi$$

$R \neq 1$: not easily mimicked.

- potential smoking gun for modified gravity?
- Significant stresses exceptionally hard to create in non-relativistic fluids e.g. DM and dark energy.

Cosmological tests of gravity

- Non-relativistic tracers: Galaxy positions and motions

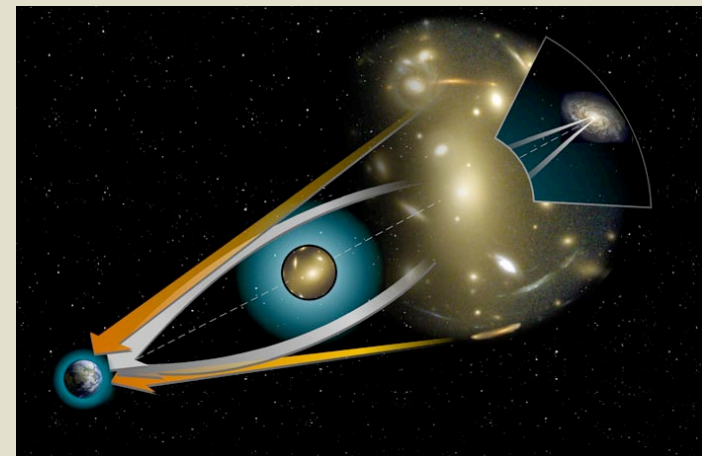
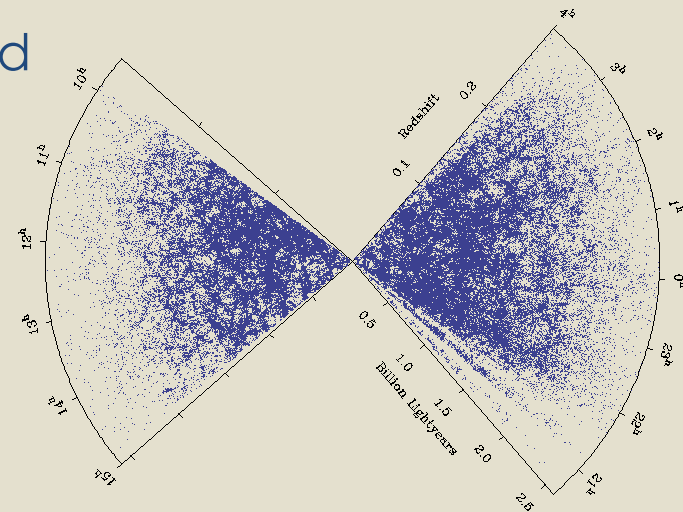
- Measure $\psi \sim G_{\text{mat}} = QRG_N$
- Biasing of tracer (galaxy) issue

$$\delta_g = b\delta_m$$

- Relativistic tracers: Weak lensing and CMB

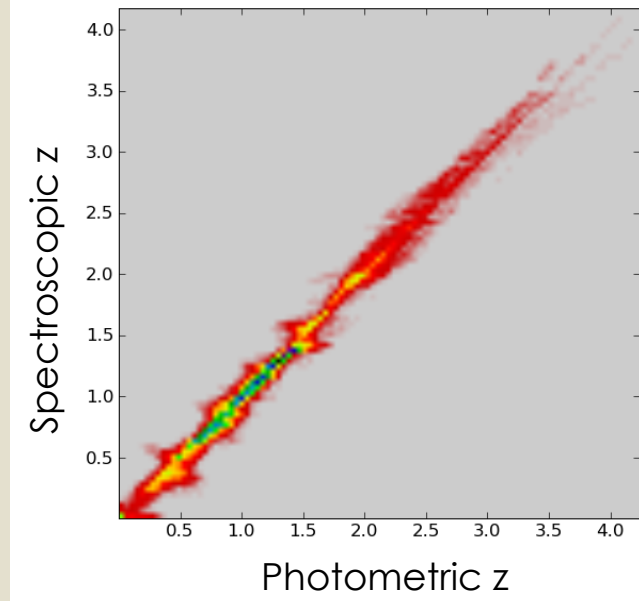
- Sensitive to $(\phi+\psi) \sim G_{\text{light}} = Q(1+R)G_N$ and time derivs
- In theory direct tracer of potential, but still uncertainties
 - stochasticity relating luminous and all mass r_g Dekel & Lahav '99
 - plenty of systematics (photoz, IAs...)

- Complementarity of tracers key to testing gravity

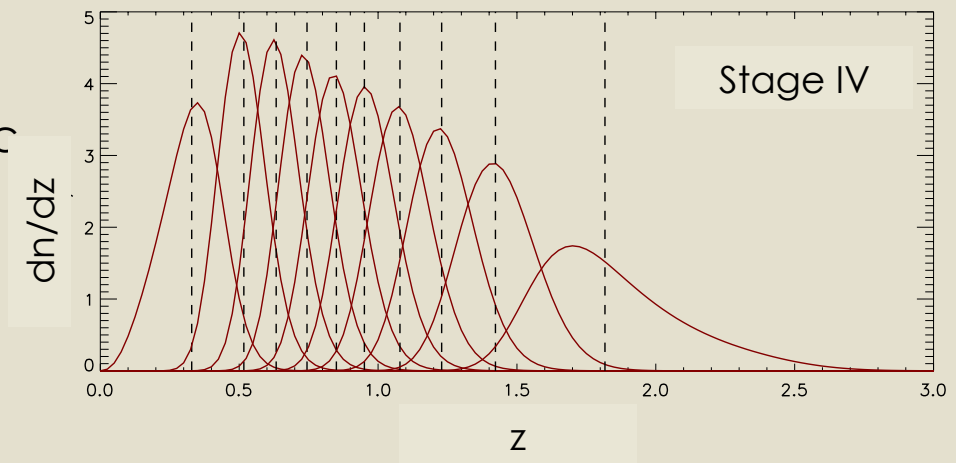


Complications : photometric redshifts

- Faster but less precise alternative to spectroscopic z
- Essential for tomography
 - Measuring evolution on dark energy
 - Cross-correlations between z bins useful for disentangling systematics and cosmology
- Sensitive to modeling
 - galaxy distribution,
 - photo- z statistical accuracy, systematic offsets and catastrophic errors



Credit: LSST Consortium

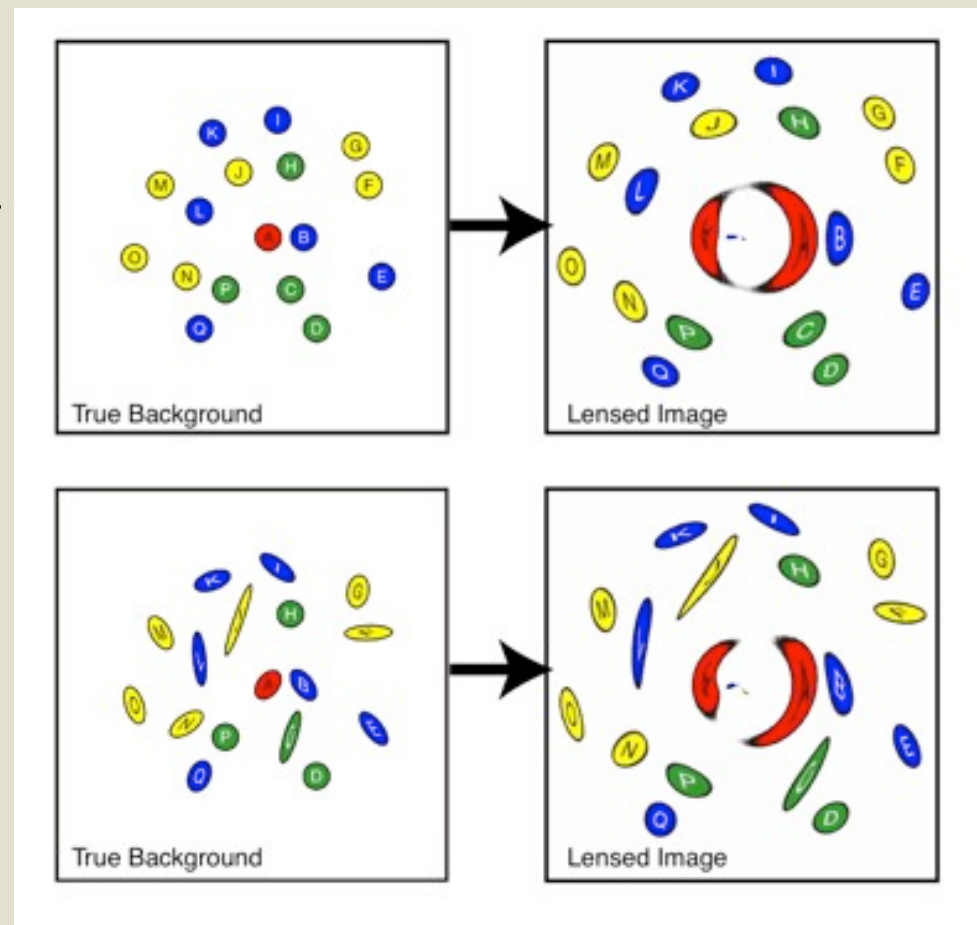


Complications : Weak lensing distortions

- 2D map on the sky of galaxy ellipticities

$$\epsilon^i(\theta) = \gamma_G^i(\theta) + \gamma_I^i(\theta) + \epsilon_{rnd}^i(\theta).$$

- Correlation in ellipticities measured statistically
 - Random ellipticities not an issue
 - Instrumental & astrophysical “contaminants” – shear calibration uncertainties
 - Correlated alignments need to be modeled and disentangled from cosmological shear



Credit: Williamson, Oluseyi, Roe 2007

Complications : Intrinsic alignments

- Significant astrophysical systematic
- Galaxies align in the potential gradient of their host halo

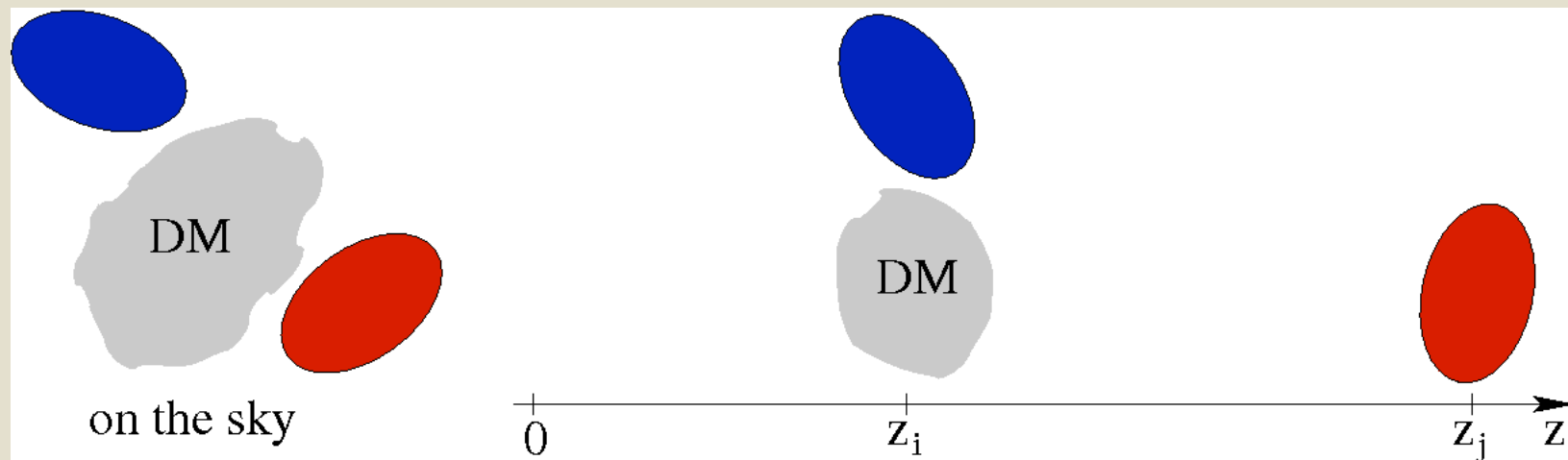
$$\langle \epsilon^i \epsilon^j \rangle = \langle \gamma_G^i \gamma_G^j \rangle + \langle \gamma_G^i \gamma_I^j \rangle + \langle \gamma_I^i \gamma_G^j \rangle + \langle \gamma_I^i \gamma_I^j \rangle$$

Correlation:

Observed
Cosmological
(GG)

Intrinsic (II)

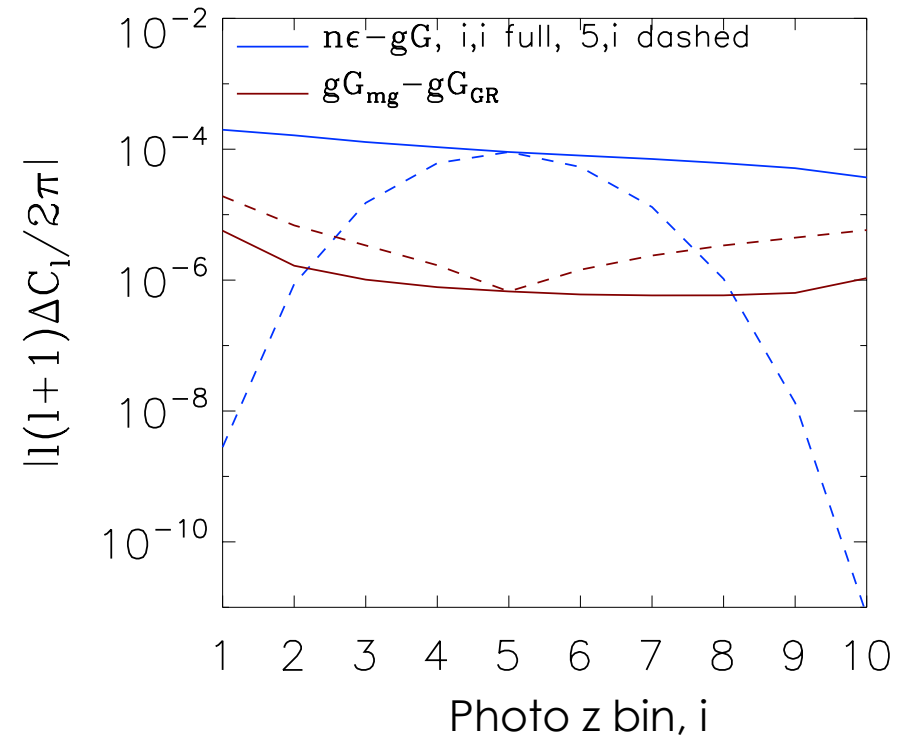
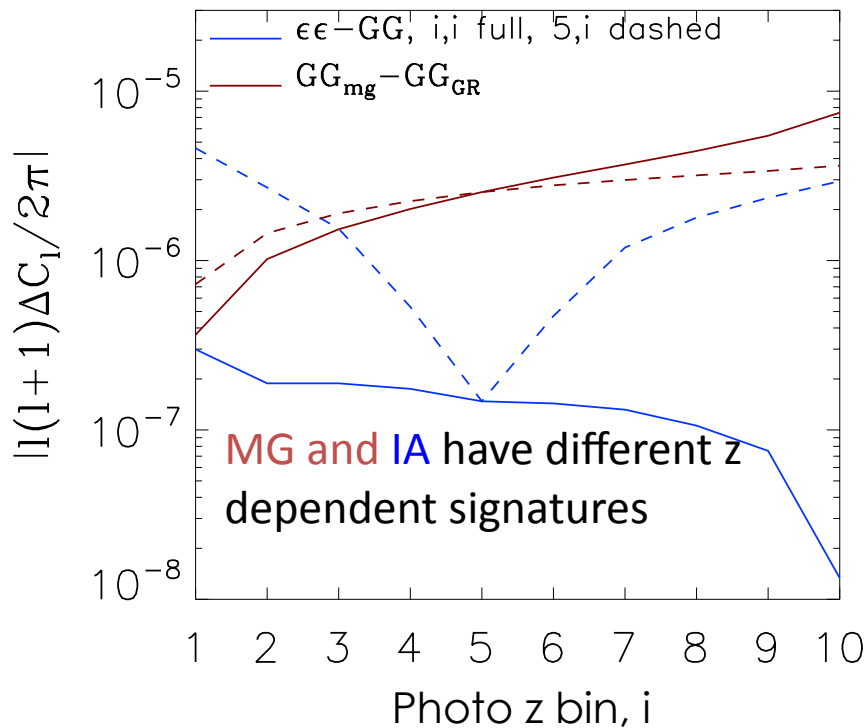
GI shear (anti) correlation



Credit: Benjamin Joachimi, iCosmo

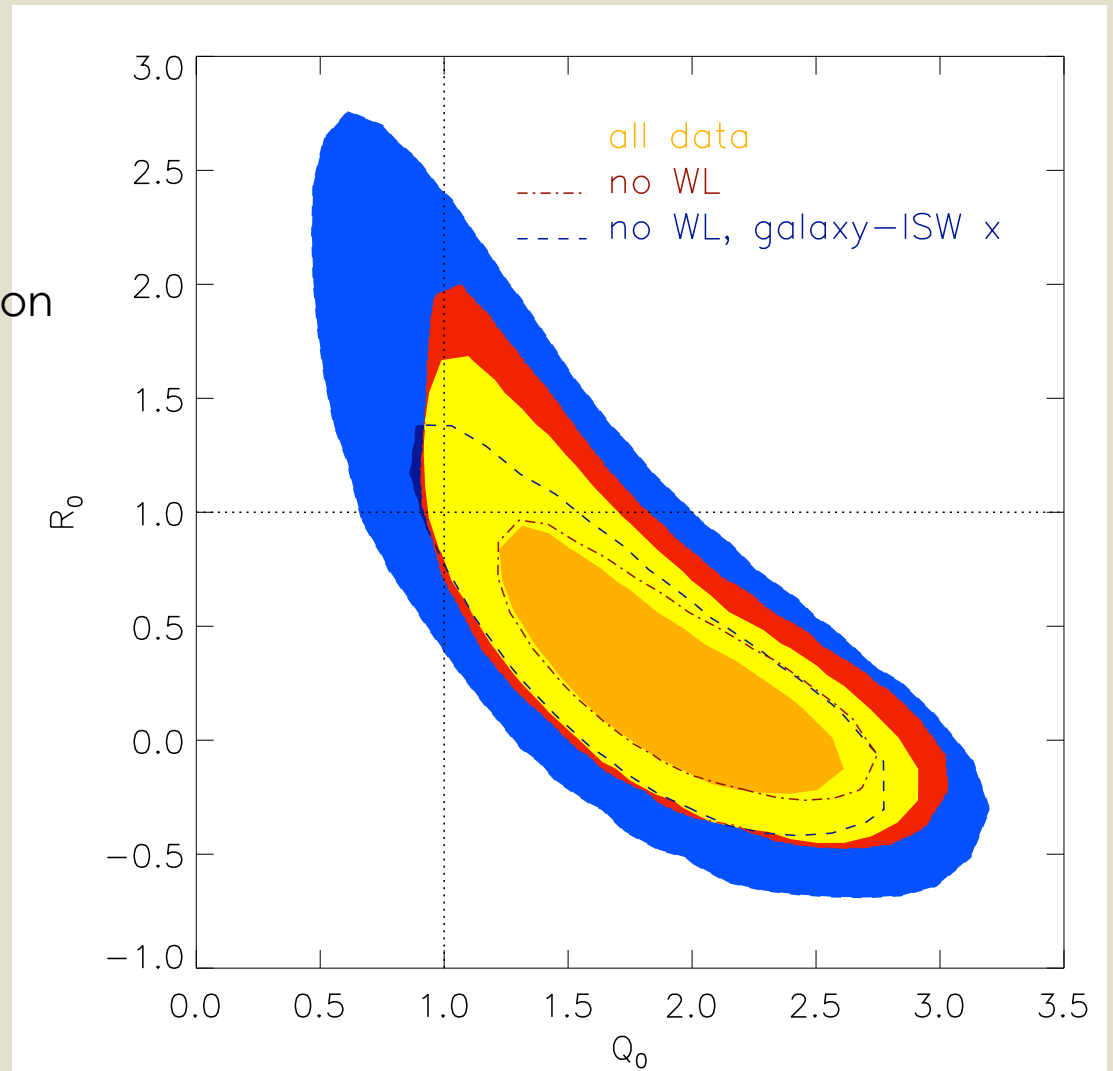
Cross- correlations and tomography: break degeneracy between systematics and theory

Plots of $C_l^{X_i Y_i}$ and $C_l^{X_5 Y_i}$



Current constraints

- Multiple data
 - WMAP CMB,
 - SDSS LRG auto
 - SDSS-WMAP ISW cross correlation
 - COSMOS weak lensing,
 - Union SN1a
- ISW + ISW-galaxy correlations drive constraints
- Principal degeneracy
 - $(\phi+\psi)$ direction $\sim Q(1+R)/2$
- “Figure of Merit”
 - $1/\text{error ellipse area}$
 - MG FoM ~ 0.03



What about future surveys?

- Fisher matrix analysis = Inverse covariance (error) matrix

$$Cov_{ij}^{-1} = F_{ij} = \frac{\partial t_a}{\partial p_i} Cov_{ab}^{-1} \frac{\partial t_b}{\partial p_j}$$

- Assumed cosmology and parameterization

$$\mathbf{p} = \{\Omega_b h^2, \Omega_m h^2, \Omega_k, \tau, w_0, w_a, Q_0, Q_0(1 + R_0)/2, n_s, \Delta_{\mathcal{R}}^2(k_0), \\ + \text{systematic nuisance parameters}\}$$

- Datasets

$$\mathbf{t} = \{C_{\ell}^{TT}, C_{\ell}^{TE}, C_{\ell}^{EE}, C_{\ell}^{Tg_1}, \dots, C_{\ell}^{Eg_1}, \dots, C_{\ell}^{g_1g_1}, C_{\ell}^{g_1g_2}, \dots, C_{\ell}^{\kappa_{N_{ph}} \kappa_{N_{ph}}}\}$$

- Survey specifications

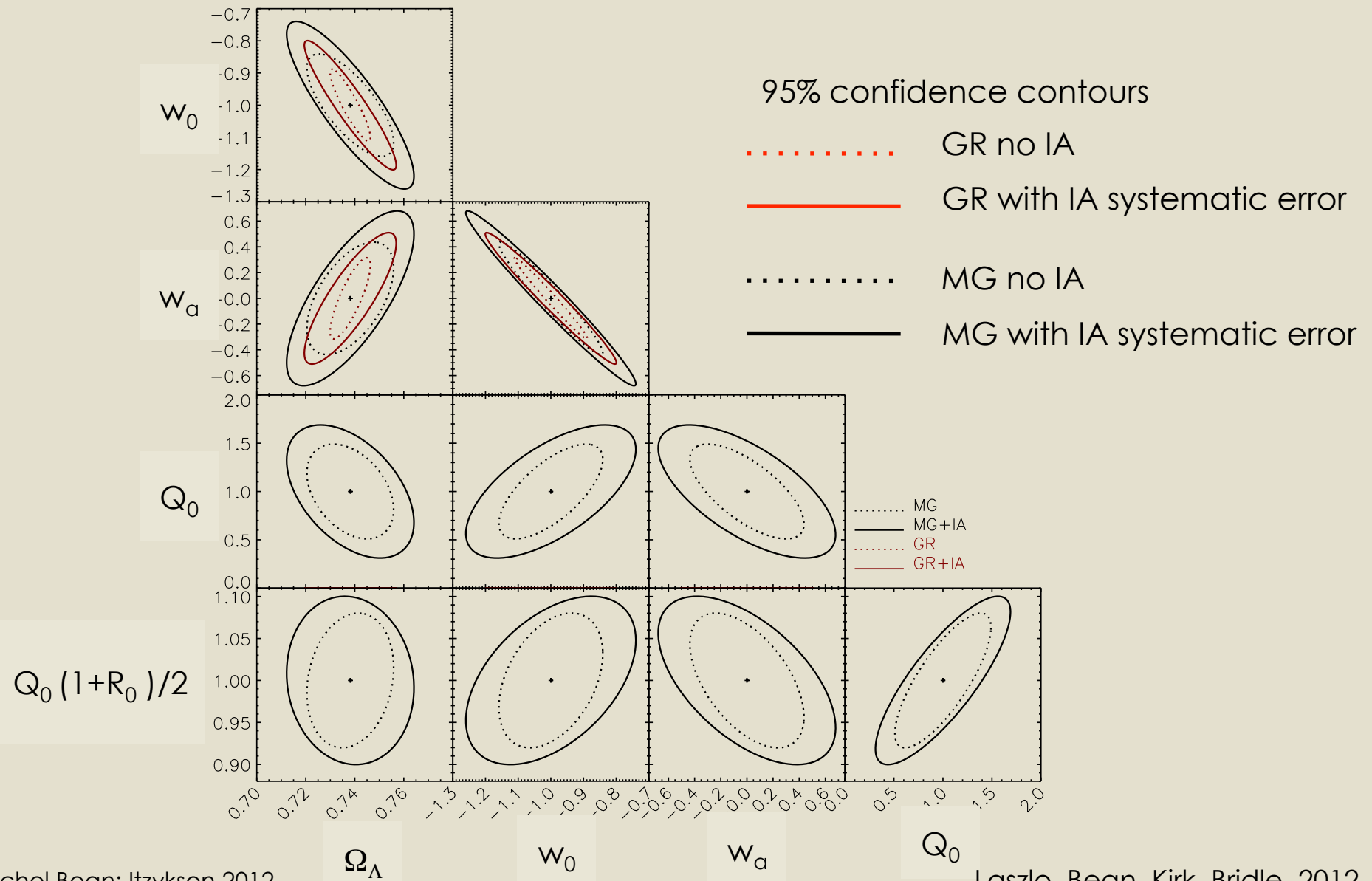
- near future (stage III) and end of decade (stage IV) surveys
- Stage III = Planck CMB + DES-like imaging + BOSS spectroscopic surveys
- Stage IV = Planck CMB + EUCLID-like imaging and spectroscopy

Forecasting: what you put in=what you get out

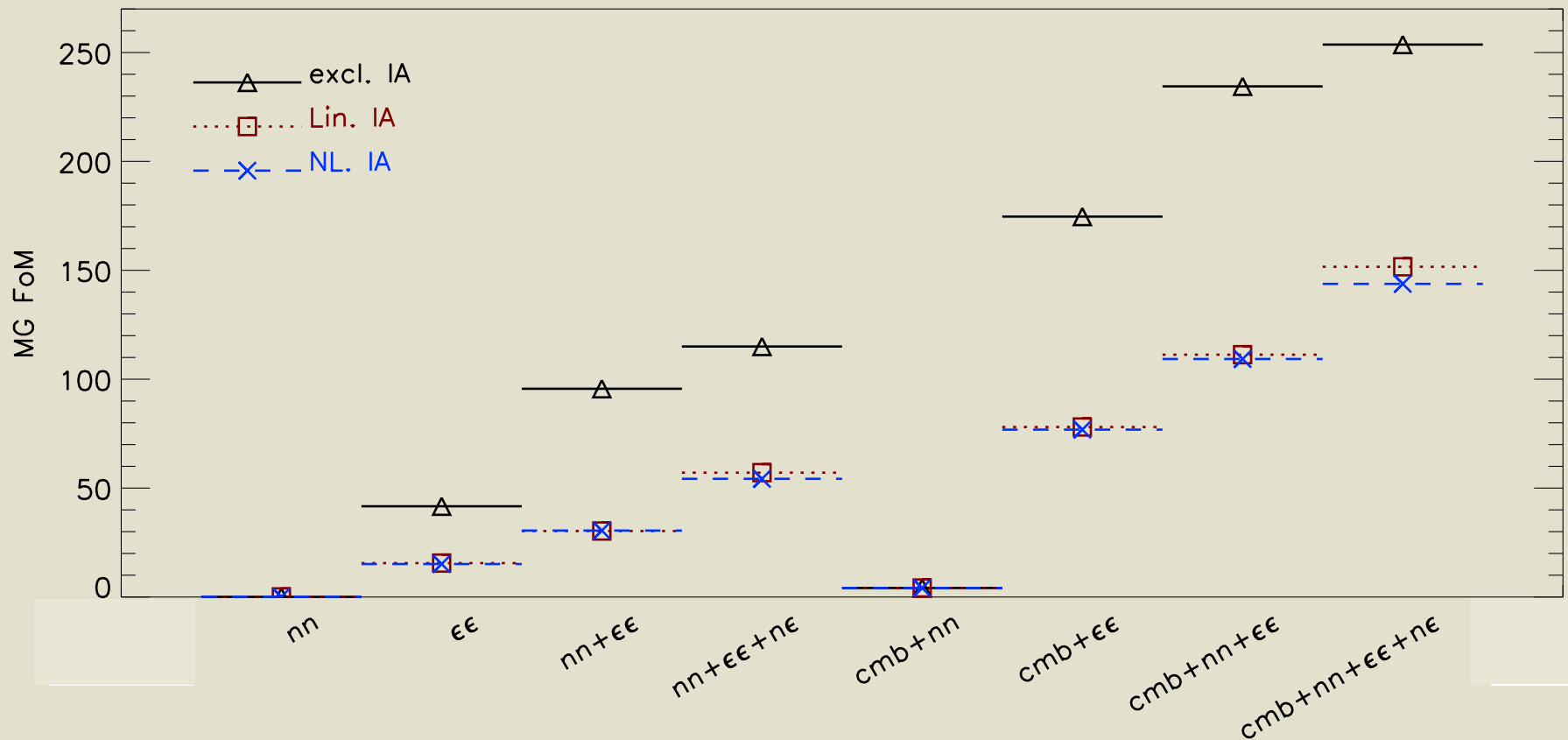
- Figures of merit /Fisher insightful but
- Model dependent – e.g. w_0/w_a or functions of z ?
- Systematic errors difficult but important!
 - Instrumental e.g. calibration uncertainties
 - Internal cross-checks: inter-filter, concurrent & repetition \neq redundancy
 - Modeling: e.g. Photo z modeling errors, nonlinearity
 - Access to ground based facilities,
 - Training sets, simulation suites
 - Astrophysical: e.g. IAs , $H\alpha$ z distribution, galaxy bias, baryonic effects
 - At what scale should one truncate the analysis?
 - Analytical modeling, gridded k & z bins, simulations?
- Buyer beware: risky to compare FoM unless apples-for-apples treatment



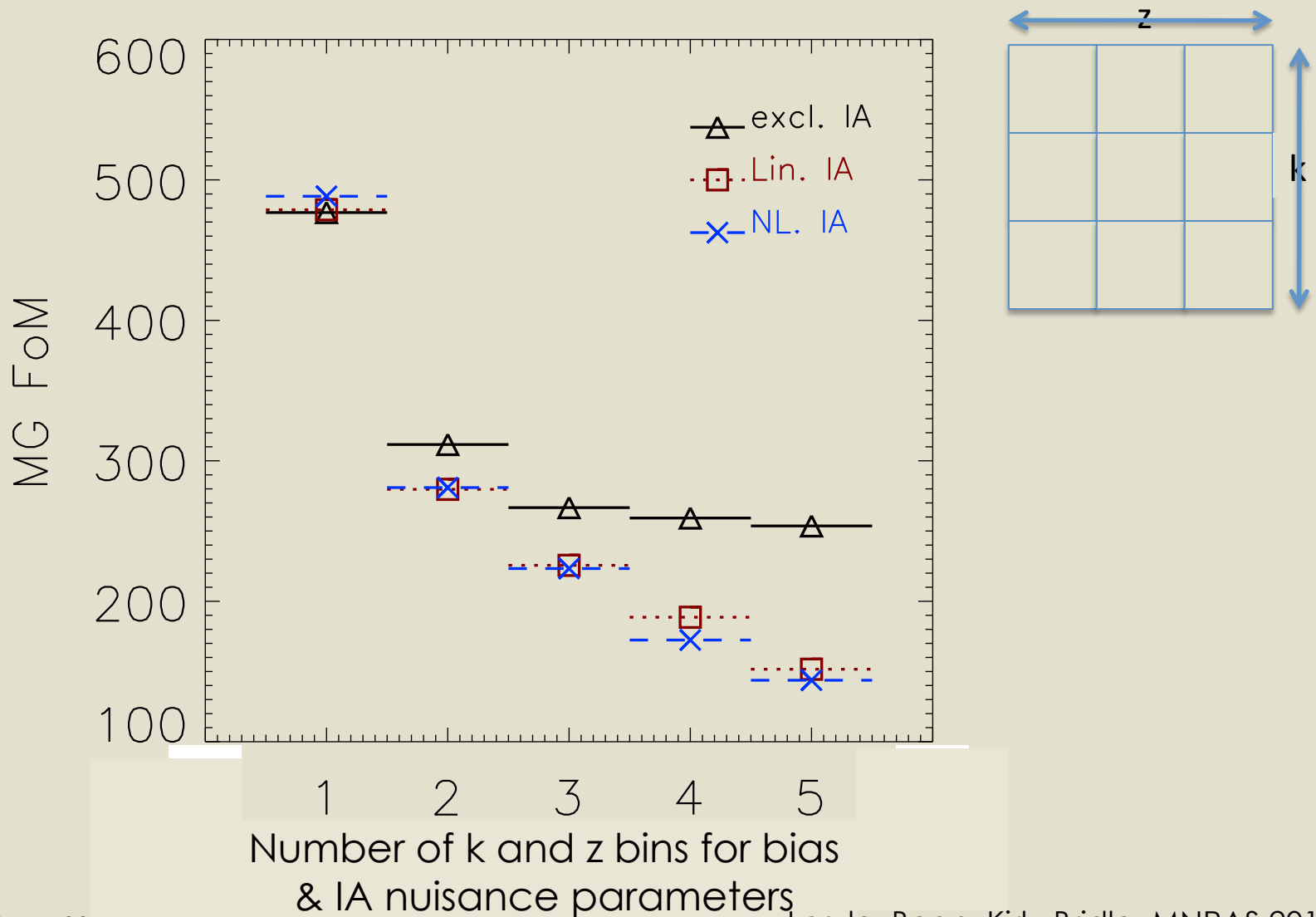
Sensitivity to theory and systematics



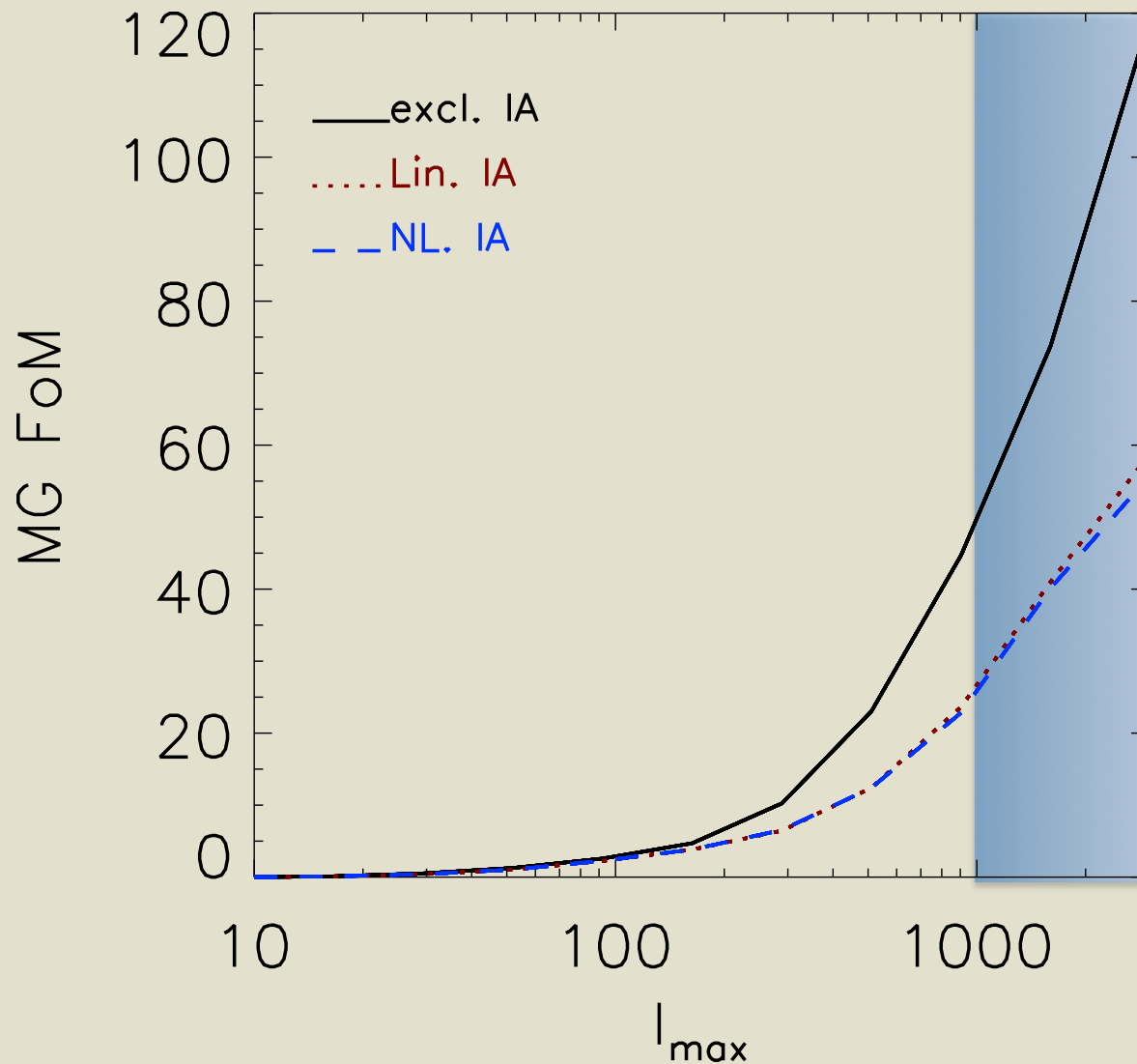
Impact of cross-correlations: reducing systematics, breaking theory degeneracies



Assumptions about bias and IA model



If you understand non-linear scales
they could make a big difference



On scales $< \sim$ a few Mpc

- Baryonic effects?
- Non-linear modeling?
- Screening effects?

Include small scale modeling uncertainties in forecasts.

Ways to modify gravity?

- Scalar tensor gravity = simple models we can model effects for

GR

$$S = \int d^4x \sqrt{-g} \frac{1}{16\pi G} R.$$

f(R) gravity

$$S = \int d^4x \sqrt{-g} \frac{1}{16\pi G} (R + f_2(R))$$

Scalar tensor gravity

$$S = \int d^4x \sqrt{-g} \frac{1}{16\pi G} f_1(\phi) R.$$

Higher dimensional gravity e.g. DGP

$$S = \int d^5x \sqrt{-g^{(5)}} \frac{1}{16\pi G^{(5)}} R^{(5)}$$

- Active area of research, many different options, no solutions, yet
- Common theme: A scalar degree of freedom

Effective field theory of acceleration

- Can we tie phenomenology/data a step closer to theory?
- Write a general action as an expansion in derivative powers

$$S = \int d^4x \sqrt{-g} \left\{ \frac{m_p^2}{2} R - \frac{1}{2} (\nabla\phi)^2 - U(\phi) \right\} + S_m[e^{\alpha(\phi)} g_{\alpha\beta}, \psi_m] \\ + \epsilon \int d^4x \sqrt{-g} \left\{ a_1 (\nabla\phi)^4 + b_2 T (\nabla\phi)^2 + c_1 G^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi \right. \\ \left. + d_3 (R^2 - 4R^{\mu\nu} R_{\mu\nu} + R_{\mu\nu\sigma\rho} R^{\mu\nu\sigma\rho}) + d_4 \epsilon^{\mu\nu\lambda\rho} C_{\mu\nu}^{\alpha\beta} C_{\lambda\rho\alpha\beta} + e_1 T^{\mu\nu} T_{\mu\nu} + e_2 T^2 \right\}.$$

- a_1, a_2 etc are all free functions of ϕ
 - non-minimally coupled to metric in Einstein frame, “modified gravity”
- What observational properties might this type of action have?

Effective field theory of acceleration

- A subset of terms particularly relevant to late time evolution
(Work with Eva-Marie Mueller and Scott Watson)

$$\begin{aligned}
 S = & \int d^4x \sqrt{-g} \left\{ \frac{m_p^2}{2} R - \frac{1}{2} (\nabla\phi)^2 - U(\phi) \right\} + S_m[e^{\alpha(\phi)} g_{\alpha\beta}, \psi_m] \\
 & + \epsilon \int d^4x \sqrt{-g} \left\{ a_1 (\nabla\phi)^4 + b_2 T(\nabla\phi)^2 + c_1 G^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi \right. \\
 & \left. + d_3 (R^2 - 4R^{\mu\nu} R_{\mu\nu} + R_{\mu\nu\sigma\rho} R^{\mu\nu\sigma\rho}) + d_4 \epsilon^{\mu\nu\lambda\rho} C_{\mu\nu}^{\alpha\beta} C_{\lambda\rho\alpha\beta} + e_1 T^{\mu\nu} T_{\mu\nu} + e_2 T^2 \right\}.
 \end{aligned}$$

- Canonical scalar field
- Non-minimally coupled matter
- A quartic term
- A Gauss-Bonnet (GB) term
- Other terms here well constrained by early time effects

Effects of extensions to minimally coupled scalar

- Study attractor behavior of this action

$$S_E = \int d^4x \sqrt{-g} \left\{ \frac{M_p^2}{2} R - \frac{1}{2} (\nabla\phi)^2 - V(\phi) + X(\dot{\phi}, H) - f(\phi) R_{GB}^2 \right\} + S_m[e^{F(\phi)} g_{\mu\nu}]$$

$$F(\phi) \equiv \exp\left(2C\sqrt{\frac{2}{3}}\frac{\phi}{M_p}\right)$$

$$X(\dot{\phi}, H) \equiv \frac{\beta}{M_p^2 H^2} \dot{\phi}^4$$

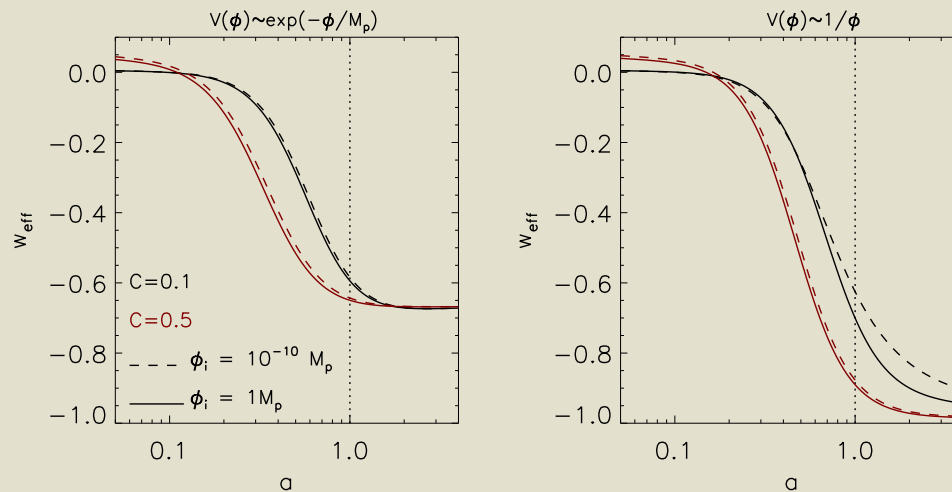
- Effect on Friedmann and equations

$$3M_p^2 H^2 = \rho_m(\phi) + \rho_\gamma + \frac{1}{2}\dot{\phi}^2 + V(\phi) + 24\dot{\phi}f'(\phi)H^3 + 3X$$

$$\ddot{\phi} + 3H\dot{\phi} + V' = \sqrt{\frac{2}{3}}\frac{C}{M_p}\rho_m - 24f'H^4\left(\frac{\dot{H}}{H^2} + 1\right) + \frac{X}{H\dot{\phi}}\left[-3\frac{\dot{X}}{XH} - 12\right]$$

To attract or not? A matter of effective potentials

- Potentials with exponential or power law forms allow attractor solutions
 - Independence to initial conditions (for better or worse e.g. $f(R)$ Amendola et al 2007)
 - Transitions from matter to accelerative era



RB, Flanagan, Laszlo, Trodden 2008

- Can postulate potentials that deny attractors, and retrofit to LCDM background.
 - Evades disadvantageous attractor behavior e.g. $f(R)$ Hu and Sawicki 2007

$$f(R) = -m^2 \frac{c_1 (R/m^2)^n}{c_2 (R/m^2)^n + 1},$$

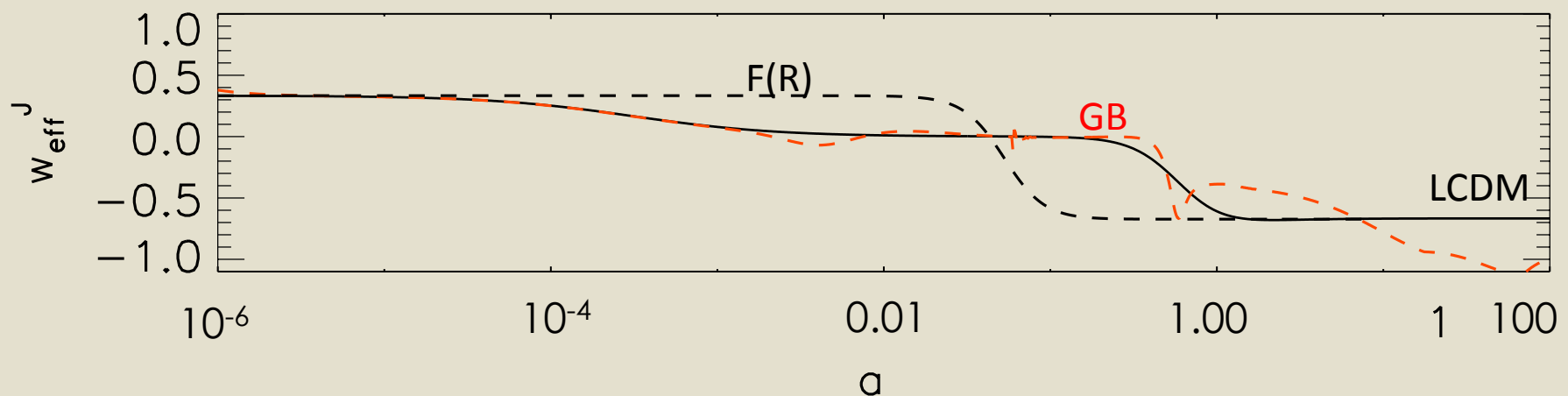
Effect on cosmological attractors

- Express dynamical equations in terms of dimensionless parameters

$$x = \frac{1}{M_p H} \frac{\dot{\phi}}{\sqrt{6}}, \quad y = \frac{1}{M_p H} \frac{\sqrt{V}}{\sqrt{3}}, \quad z = \frac{1}{M_p H} \frac{\sqrt{\rho_\gamma}}{\sqrt{3}},$$

$$\mu_X \equiv \frac{X}{3M_p^2 H^2} = 12\beta x^4 \quad \mu_{GB} \equiv \frac{f' H^2}{M_p}$$

- Find if stationary solutions exist $x'=y'=z'=\mu'=0$



Evolutionary attractor

- Matter dominated era: NMC and kinetic terms important
 - NMC cosmic grease $w > 0$
 - Quartic term adds to Hubble drag, acts to slow expansion

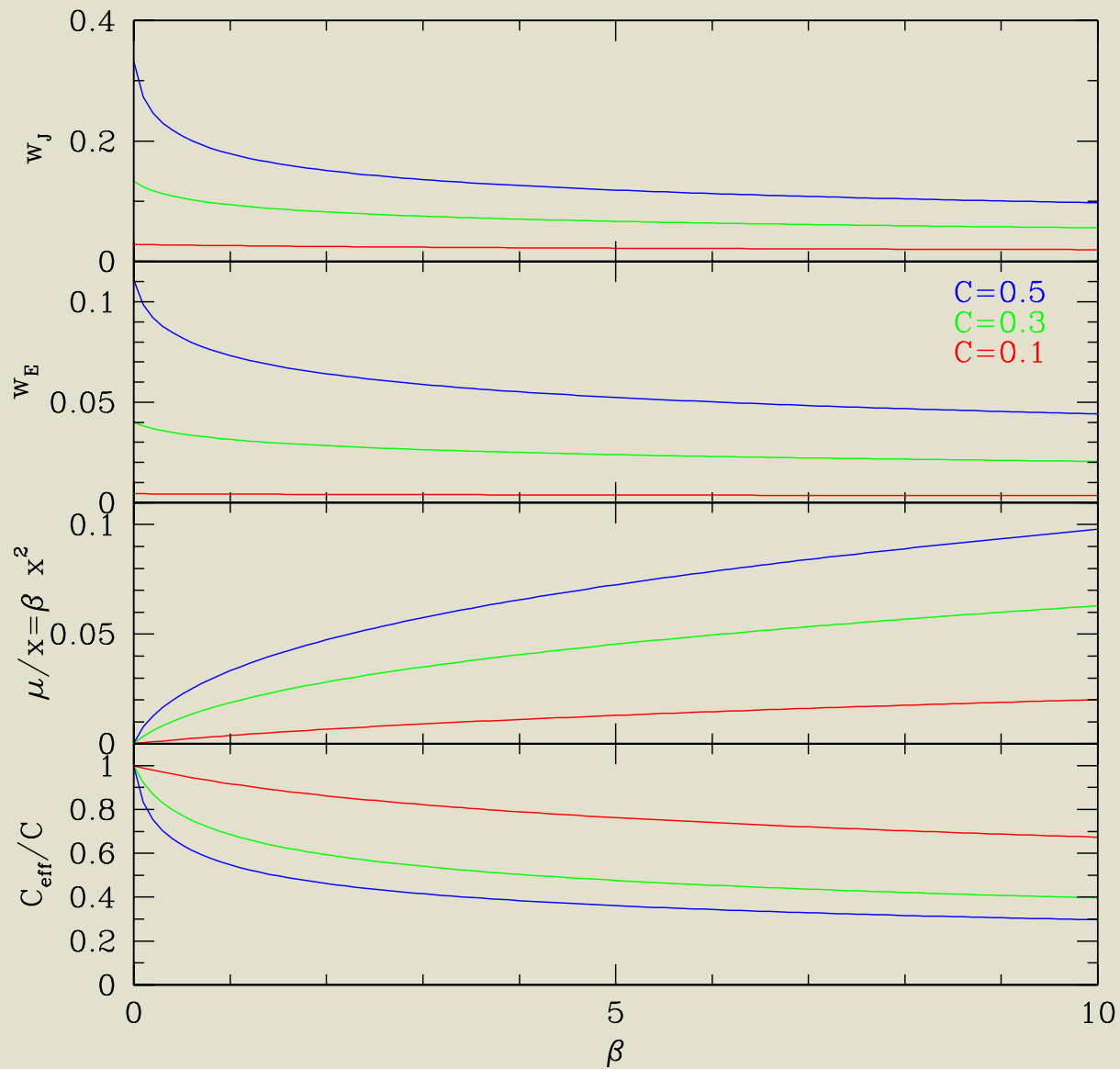
$$x = \frac{2}{3} C_{eff}(C, \beta) \quad y = 0, z = 0, \mu_{GB} = 0$$

- Accelerative era:
 - GB term gives an accelerative attractor (Koivisto and Mota 2006)

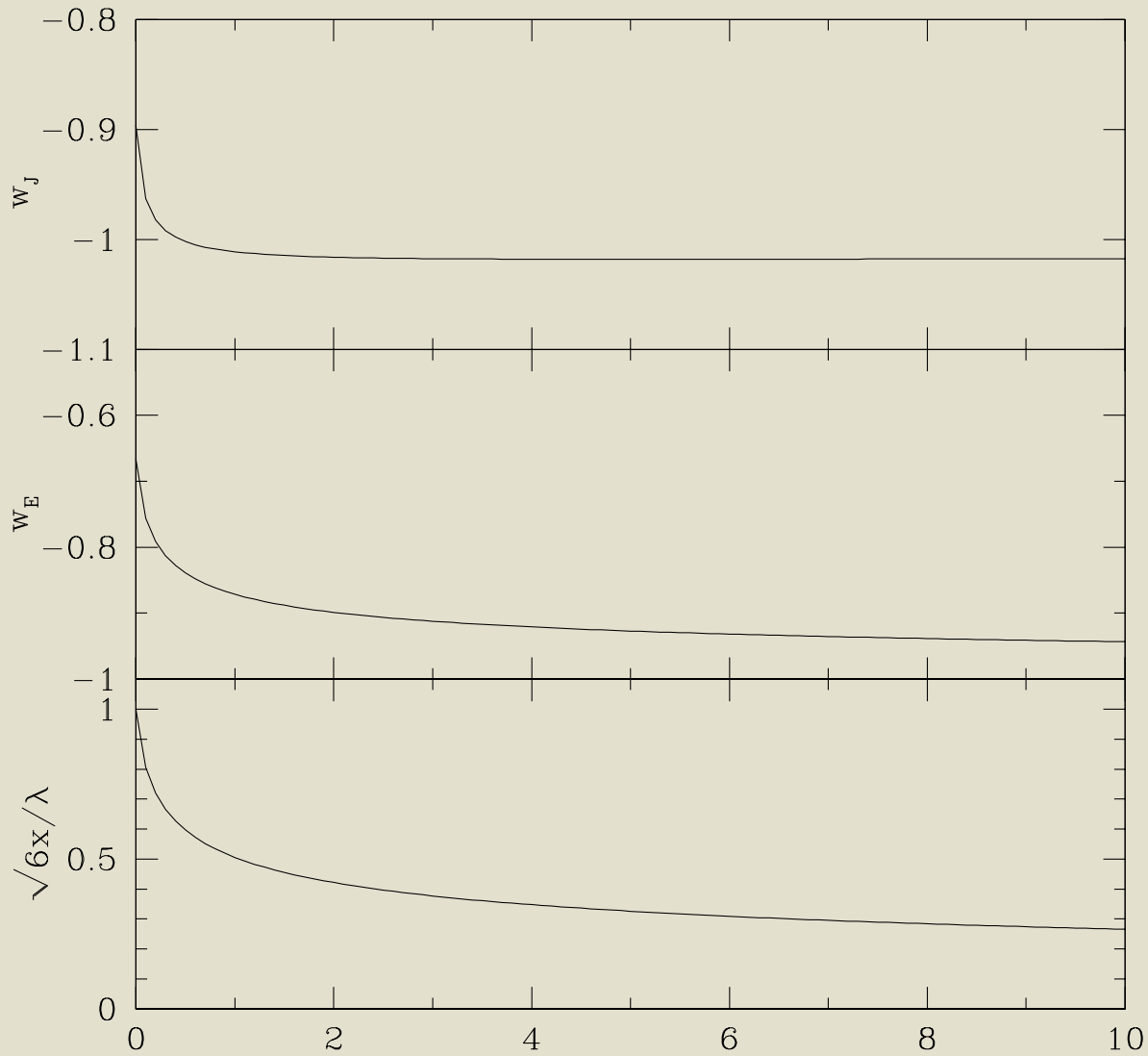
$$x = 0, y = 1, \mu_{GB} = \frac{\lambda}{8}$$

- In absence of GB term quartic allows acceleration for broader ranges of potentials with $\lambda > 2$

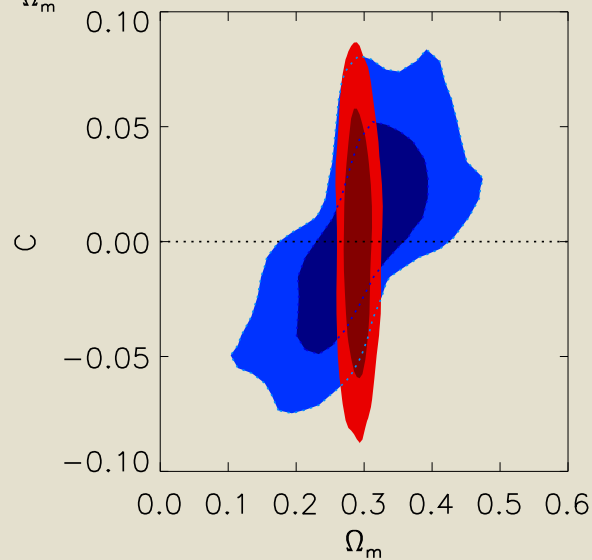
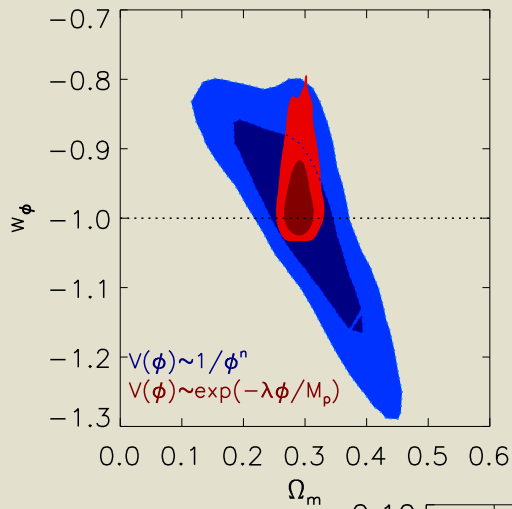
Matter era



Accelerative era

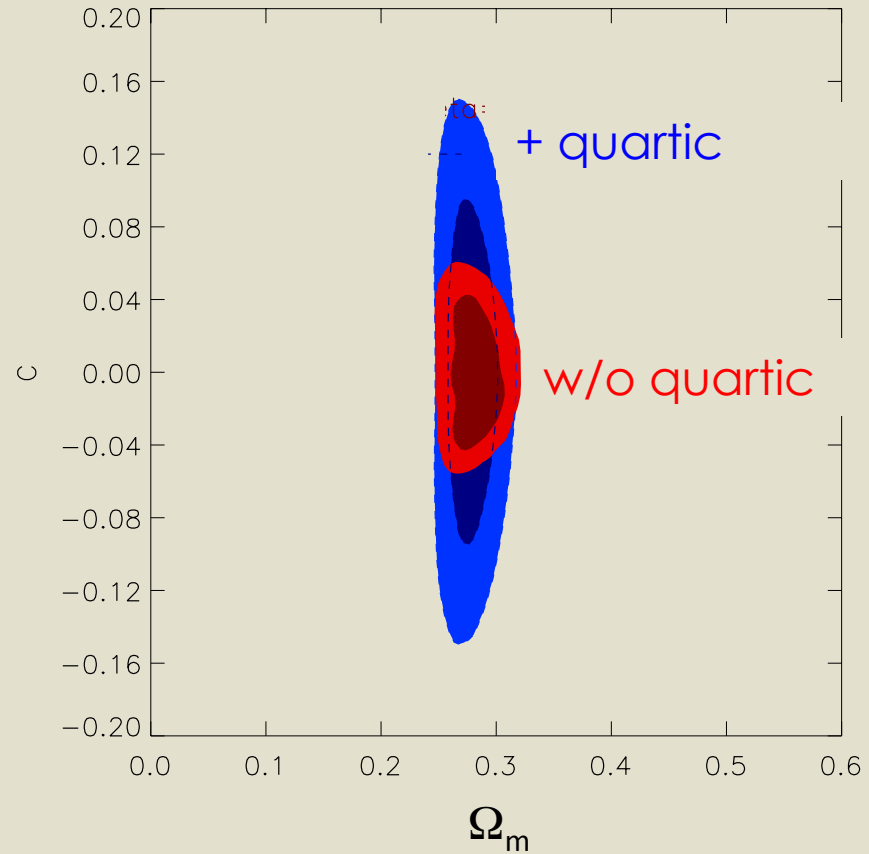


Impacts on geometric cosmological constraints



Einstein frame (just coupling to CDM)

Geometric constraints CMB + BAO + SN



Jordan frame (coupling to all matter)

Concluding thoughts

- Upcoming datasets provide invaluable opportunity to test the origins of cosmic acceleration and weak field gravity on cosmic scales
 - Complementary techniques important to break cosmological and systematic degeneracies
 - Relativistic and non-relativistic LSS tracers of gravity equivalent to complementary geometric techniques to measure expansion history/curvature
- Inclusion of systematics modeling essential to forecasting
 - Can significantly impact predictions (beware apples vs oranges)
 - Theory and systematics can be tightly coupled.
 - How important are specific systematics/ algorithms to model them to a technique's ability to constrain DE?
- Phenomenological modeling/ FoMs useful but a high pass filter on full info. Mapping to theory is the ultimate goal.
 - Surveys will give us information about z and k dependence
 - General effective field theory for DE is a first step, with interesting implications for both expansion history and growth history