#### Particle Cosmology Center for the University of Pennsylvania

# Classical and Quantum Stability of Chameleon Theories Justin Khoury (UPenn)

#### w. Amol Upadhye, Junpu Wang, Wayne Hu, Lam Hui



## A Richer Dark Sector

Dark energy candidates:  $\Lambda$ , quintessence... Ratra & Peebles (1988); Wetterich (1988); Caldwell, Dave & Steinhardt (1998)



A Richer Dark Sector Dark energy candidates: Λ , quintessence... Ratra & Peebles (1988); Wetterich (1988); Caldwell, Dave & Steinhardt (1998)



Tantalizing prospect: dark sector includes new light fields (e.g. quintessence) that couple to both dark and baryonic matter.

Scalar fields can "hide" themselves from local exp'ts through screening mechanisms

$$
\rho_{\rm here} \sim 10^{30} \rho_{\rm cosmos}
$$

Large in a cosmological sense, but small in a particle physics sense



ο To ne degree To neutralize  $\Lambda$  to accuracy  $H_0^z$  , new degrees of freedom must be light:  $\Lambda$  to accuracy  $H_0^2$ 



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 $m_\phi \lesssim H_0$ 

These degrees of freedom must couple to Standard Model fields:

 $\phi$ 

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To ne degree<br>degree<br>The to Sto<br>The to Sto<br>The to Sto<br>The to Sto<br>Incel tests of GR Must rely on screening mechanism for consistency with local tests of GR

=⇒

Experimental Program  $U(r) = -g$  $\overline{M}$  $8\pi M_\text{P}^2$ Pl  $e^{-r/\lambda}$ r





Screening mechanisms have rich phenomenology for tests of GR: Forced us to rethink implications of existing data Inspired design of novel experimental tests

#### Chameleon/symmetron/dilaton:

 $\mathcal{L}=-\frac{1}{2}$ 2  $(\partial \phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$ 

Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:



#### ● Galileon/Vainshtein:

$$
\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 + \frac{1}{\Lambda^3} (\partial \phi)^2 \Box \phi + g \frac{\phi}{M_{\text{Pl}}} T^\mu_{\ \mu}
$$

Deffayet, Dvali, Gabadadze & Vainshtein (2001); Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

- Screening condition:

$$
|\partial^2 \phi| \gg \Lambda^3
$$

Chameleon/symmetron/dilaton:

 $\mathcal{L}=-\frac{1}{2}$ 2  $(\partial \phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$ 

 $\frac{1}{4\Lambda^4}(\partial\phi)^4+g$ 

 $\phi$ 

 $T^{\mu}_{\,\,\mu}$ 

 $M_{\rm Pl}$ 

Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:  $\mid \phi \ll \phi_c \mid$ 



Kinetic/k-mouflage: Babichev, Deffayet & Ziour (2009)

- Screening condition:  $\left|\,\left|\overline{\partial \phi}\right|\gg\Lambda^2\,\right|$ 

 $\mathcal{L}=-\frac{1}{2}$ 

2

 $\left(\partial\phi\right)$ 

 $\frac{2}{4} - \frac{1}{4 \Lambda}$ 

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- Kinetic/k-mouflage:  $\mathcal{L}=-\frac{1}{2}$ Babichev, Deffayet & Ziour (2009)
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- Shift symmetry:  $\phi \rightarrow \phi + c$
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 $-$  Cutoff:  $\Lambda \sim {\rm mm}^{-1}$ 

 $\phi$ 

 $(\partial \phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$ 

 $T^{\mu}_{\,\,\mu}$ 

 $\Lambda \sim (1000\;{\rm km})^{-1}$ 

 $M_{\rm Pl}$ 

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- Screening condition:

- Galileon symmetry:  $\phi \rightarrow \phi + c + b_\mu x^\mu$  - Superluminality

 $|\partial^2\phi|\gg\Lambda^3$ 

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 $T^{\mu}_{\,\,\mu}$ 

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 $\frac{1}{4\Lambda^4}(\partial\phi)^4+g$ 

1

$$
\frac{1}{\sqrt{1-\frac{1}{2}}}
$$

 $|\partial \phi| \gg \Lambda^2$ 

2

 $\left(\partial\phi\right)$ 

 $\mathcal{L}=-\frac{1}{2}$ 

 $\bigodot^1_4$ 

2

 $(\partial \phi)^2 +$ 

#### Screening Mechanisms Chameleon/symmetron/dilaton: **Galileon/Vainshtein:** Kinetic/k-mouflage:  $\mathcal{L}=-\frac{1}{2}$  $\mathcal{L}=-\frac{1}{2}% \sum_{i=1}^{n^{\prime}}% \frac{1}{i}(\mathbf{r}_{i}+\$ 2  $(\partial \phi)^2 +$ 1  $\frac{1}{\Lambda^3}(\partial\phi)^2\Box\phi+g$  $\phi$  $M_{\rm Pl}$  $T^{\mu}_{\,\,\mu}$ 2  $\left(\partial\phi\right)$  $\bigodot^1_4$  $\frac{1}{4\Lambda^4}(\partial\phi)^4+g$  $\phi$  $M_{\rm Pl}$  $T^{\mu}_{\,\,\mu}$ - Screening condition:  $\mid \phi \ll \phi_c \mid$ - Quantum stability? - Screening condition: - Shift symmetry:  $\partial^3\phi|\gg1$ Superluminality - Cutoff: - Cutoff:  $\Lambda \sim {\rm mm}^{-1}$  $|\partial^3 \phi| \gg \Lambda^4$  ? – Cutoff:  $\Lambda \sim {\rm mm}^{-1}$ Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010) Babichev, Deffayet & Ziour (2000) Deffayet, Dvali, Gabadadze & Vainshtein (2001) Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004);  $\mathcal{L}=-\frac{1}{2}$ 2  $(\partial \phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$

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 $\Lambda \sim (1000\;{\rm km})^{-1}$ 

# Should we be worried?

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## Should we be worried?



## "The absence of evidence is not evidence of absence."

#### Chameleon Mechanism

Khoury & Weltman (2003); Gubser & Khoury (2004); Brax, van de Bruck, Davis, Khoury and Weltman (2004); Mota and Shaw (2006).



Consider scalar field  $\phi$  with potential  $V(\phi)$  and coupled to matter:

$$
\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-V(\phi)+g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\;\mu}
$$

where  $T^{\mu}_{~\mu}$  is stress tensor of all matter (Baryonic and Dark)

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 $V(\phi)$ 

 $~\sim \rho \phi$ 

 $\phi$ 

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$$
\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-V(\phi)+g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\;\mu}
$$

 $V_{\text{eff}}(\phi)$ 

where  $T^{\mu}_{~\mu}$  is stress tensor of all matter (Baryonic and Dark) For non-relativistic matter,  $T^{\mu}_{~\mu} \approx -\rho$  , hence

$$
\nabla^2\phi=V_{,\phi}+\frac{g}{M_{\rm Pl}}\rho
$$

$$
\longrightarrow \left| V_{\text{eff}}(\phi) = V(\phi) + g \frac{\phi}{M_{\text{Pl}}} \rho \right|
$$

## Density-dependent mass *V*<sub>eff</sub>(φ)

$$
V_{\text{eff}}(\phi) = V(\phi) + g \frac{\phi}{M_{\text{Pl}}} \rho
$$

$$
\textbf{e.g.}\quad V(\phi)=\frac{M^{4+n}}{\phi^n}
$$



Thus  $\ m = m(\rho)$  increases with increasing density Laboratory tests => set  $m^{-1}(\rho_{\text{local}}) \lesssim$  $≤ \text{mm}$ 

Generally implies:  $m^{-1}(\rho_{\rm cosmos}) \lesssim$  $\lesssim {\rm Mpc}$ 

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Generally implies:  $m^{-1}(\rho_{\rm cosmos}) \lesssim$  $\lesssim {\rm Mpc}$ 

Meanwhile,  $m^{-1}(\rho_{\rm solar\ system}) \lesssim 10-10^4\ \rm{AU}$ =⇒ ruled out by post-Newtonian tests?

*R*  $\rho = \rho_{\rm out}$   $R \wedge \rho = \rho_{\rm in}$ 







$$
\longrightarrow \left[\phi(r > R) \sim \frac{\Delta R}{R} \times \frac{g G_{\rm N} M}{r}\right]
$$

where  $\frac{\Delta R}{\Delta}$ 

R =  $\phi_{\rm out}-\phi_{\rm in}$  $\frac{\displaystyle \phi_{\rm out}-\phi_{\rm in}}{\displaystyle 6g M_{\rm Pl}\Phi_{\rm N}}\ll 1 \;\implies$  thin-shell screening



But small objects  $\implies$  no thin-shell

#### Thin-shell condition depends on environment!



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#### Thin-shell condition depends on environment!





# $G_{\rm N}^{\rm eff} = G_{\rm N} (1 + 2 g^2)$ between small objects in space !

#### Smoking Guns

 Satellite Energy Exchange (SEE) Mission  $\Delta G_{\rm N}$  $G_N$  $< 10^{-6}$ 



#### $\Delta G_{\rm N}$  $G_N$  $\sim \mathcal{O}(1)$

#### Smoking Guns

#### Satellite Energy Exchange (SEE) Mission  $\Delta G_{\rm N}$  $G_N$  $< 10^{-6}$

#### MICROSCOPE (2015)





 Satellite Test of the Equivalence Principle (STEP)







 $\Delta G_{\rm N}$  $G_N$  $\sim \mathcal{O}(1)$ ∆*a a >* 10−<sup>13</sup>

Symmetron Mechanism K. Hinterbichler and J. Khoury, Phys. Rev. Lett. (2010); Olive and Pospelov (2008); Brax, van de Bruck, Davis and Shaw (2010).

Instead of  $m(\rho)$ , here coupling to matter depends on density.

 $\mathcal{L}=-\frac{1}{2}$ 2  $\left(\partial\phi\right)$  $\frac{2}{2} - V(\phi) + \frac{\phi^2}{2M}$  $\frac{\varphi}{2M^2} T^\mu_{\,\,\,\mu}$ *µ*

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 $V\left(\phi\right)$ 

 $\phi$ 

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$$
\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-V(\phi)+\frac{\phi^2}{2M^2}T^{\mu}_{\ \mu}
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Potential is of the spontaneous-symmetrybreaking form:

$$
V(\phi) = -\frac{1}{2}\mu^2\phi^2 + \frac{1}{4}\lambda\phi^4
$$

Most general renormalizable potential with  $\phi \rightarrow -\phi$  symmetry.

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$$
\Longrightarrow \Bigg| V_{\text{eff}}(\phi) = \frac{1}{2}\left( \frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{1}{4}\lambda \phi^4
$$

∴ Whether symmetry is broken or not depends on local density

Perturbations  $\delta \phi$  around local background value couple as: *L*coupling ∼  $\bar{\phi}$  $\frac{\varphi}{M^2}\delta\phi\, \rho$ Density-dependent coupling

**O** In voids, where symmetry is broken,  $\overline{\bullet}$  Symmetron decoupled in high-density regions (where  $\overline{\phi} \simeq 0$  )

 $\phi$ 

λ



# Inspiration...

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# Symmetron Couch (\$9500.00)

"NASA-style gravity reduction."

"Offers a unique multi-phase wave experience."



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∼

#### Chameleon Searches

Eot-Wash

Adelberger et al., Phys. Rev. Lett. (2008)





#### CHameleon Afterglow SEarch (CHASE), Fermilab

Chou et al., Phys. Rev. Lett. (2008,2010)





#### Chameleon Searches (cont'd)

Axion Dark Matter eXperiment (ADMX)

P. Sikivie & co., Phys. Rev. Lett. (2010)



#### Astrophysical signatures

Macroscopic violations of the Equivalence Principle Hui, Nicolis & Stubbs (2009); Jain & Vanderplas (2011)





Modified stellar evolution  $\circledcirc$ 



Chang & Hui (2010); Davis, Lim, Sakstein & Shaw (2011); Jain, Vikram & Sakstein (2012)

cf. Eugime Lim's talk

#### Astrophysical signatures (cont'd)

Saturn's rings

Enhanced self-gravity due to chameleon should impact propagation of density waves.



$$
v_g = \frac{\pi G_{\rm N}\sigma_0}{\kappa}
$$

 $\sigma_0 \equiv$  surface density  $\kappa \equiv$  epicyclic frequency



Self-Acceleration?

Wang, Hui & J. Khoury, to appear

 $\widetilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}$ 

"Jordan frame" "Einstein frame"

Can cosmic acceleration result from  $\Delta A \sim \mathcal{O}(1)$  even though Einstein-frame metric is NOT accelerating?

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 $\tilde{g}_{\mu\nu} = A^2(\phi) g_{\mu\nu}$ 

Take chameleon with  $\,\,A(\phi)=1+\frac{g\phi}{\pi\,\pi}$  $M_{\rm Pl}$  $+ \ldots$ 

Force condition:  $\frac{1}{\sqrt{D}}=2g^2$  in unscreened regions  $F_{\phi}$  $F_{\rm N}$  $=2g^2$  in unscreened regions  $\implies$   $g\lesssim \mathcal{O}(1)$ 

Screening condition:

 $\Delta R$ R  $\approx \frac{\phi_0}{6 a M_{\rm D}}$  $6gM_\mathrm{Pl}\Phi_\mathrm{N}$ 

 $\lt 1$ 

Milky Way potential  $\Phi_{\rm N} \sim 10^{-6} \implies \boxed{\phi_{\rm 0} \lesssim 10^{-6} M_{\rm Pl}}$ 

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∴<br>∴

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Can cosmic acceleration result from  $\Delta A \sim \mathcal{O}(1)$  even though Einstein-frame metric is NOT accelerating?

Take chameleon with  $\,\,A(\phi)=1+\frac{g\phi}{\pi\,\pi}$  $M_{\rm Pl}$  $+ \ldots$ 

 $\Delta A \lesssim g$ 

Force condition:  $\frac{1}{\sqrt{D}}=2g^2$  in unscreened regions Screening condition:  $F_{\phi}$  $F_{\rm N}$  $=2g^2$  in unscreened regions  $\implies$   $g\lesssim \mathcal{O}(1)$  $\Delta R$ R  $\approx \frac{\phi_0}{6 a M_{\rm D}}$  $6gM_\mathrm{Pl}\Phi_\mathrm{N}$  $\lt 1$ Milky Way potential  $\Phi_{\rm N} \sim 10^{-6} \implies \boxed{\phi_{\rm 0} \lesssim 10^{-6} M_{\rm Pl}}$  $\phi_0$ 

 $M_{\rm Pl}$ 

 $\ll 1$  (NO self-acceleration)

The Mpc Barrier  $\,$  Can scalar field have  $\,m_\phi \sim H_0 \,$  for  $\Delta t \gtrsim H_0^{-1} \,$  ? Wang, Hui & J. Khoury, to appear The Mpc Barrier  $\,$  Can scalar field have  $\,m_\phi \sim H_0 \,$  for  $\Delta t \gtrsim H_0^{-1} \,$  ? Cosmological evolution:  $\ddot{\phi}+3H\dot{\phi}=-\dfrac{\text{d} V}{\text{d}\phi}-\dfrac{g}{M_{\text{Pl}}}$  $\rho$ Wang, Hui & J. Khoury, to appear The Mpc Barrier  $\,$  Can scalar field have  $\,m_\phi \sim H_0 \,$  for  $\Delta t \gtrsim H_0^{-1} \,$  ? Cosmological evolution:  $\ddot{\phi}+3H\dot{\phi}=-\dfrac{\text{d} V}{\text{d}\phi}-\begin{pmatrix} g \ M_{\text{Pl}} \end{pmatrix}$  $\rho$  $\sim H^2M_{\rm Pl}$ ∴ Pretty large force! Wang, Hui & J. Khoury, to appear

Hence  ${\rm d}V/{\rm d}\phi$  &  $g\rho/M_{\rm Pl}$  must cancel to good accuracy

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Hence  $dV/d\phi$  &  $g\rho/M_{\rm Pl}$  must cancel to good accuracy

Cancellation must hold over last Hubble time:

$$
\Delta \left( \frac{\mathrm{d}V}{\mathrm{d}\phi} + \frac{g}{M_{\mathrm{Pl}}} \rho \right) \sim H_0^2 M_{\mathrm{Pl}} \sim m_0^2 \Delta \phi
$$

$$
\quad \longrightarrow \quad
$$

$$
\implies \qquad m_0 \sim H_0 \sqrt{\frac{M_{\rm Pl}}{\Delta \phi}} > 10^3 \; H_0
$$

Argument generalizes to very wide class of chameleon/symmetron theories, including many fields!

Quantum Stability<br>to appear in Phys. Rev. Lett. to appear

Focus on scalar loops

$$
\therefore \quad \Delta V_{\rm 1-loop} = \frac{m_\phi^4(\phi)}{64\pi^2}
$$

$$
\frac{n_{\phi}^{4}(\phi)}{64\pi^{2}}\ln\left(\frac{m_{\phi}^{2}(\phi)}{\mu_{0}^{2}}\right)
$$

 $\implies$  Tension: Need small  $m_\phi$  to keep loop corrections under control, BUT chameleon mechanism relies on  $m_\phi\uparrow$  with  $\rho$ Tension: Need small  $m_{\phi}$  to keep loop corrections under control,

Quantum Stability<br>Reappears in Phys. Rev. Lett to appear in Phys. Rev. Lett.  $\frac{m_\phi^4(\phi)}{64\pi^2}\ln\Bigg(\frac{m_\phi^2}{\mu}\Bigg)$ "

 $\frac{2}{\phi}(\phi)$ 

 $\mu_0^2$ 

Focus on scalar loops:  $\Delta V_{\rm 1-loop}=$ 

**E** Tension: Need small  $m_{\phi}$  to keep loop corrections under control, BUT chameleon mechanism relies on  $m_{\phi} \uparrow$  with  $\rho$ BUT chameleon mechanism relies on  $m_{\phi} \uparrow$  with  $\rho$ 

 $m_\phi^4(\phi)$ 

0.01 0.1 1 10 100 1000 10000 matter coupling  $g$  0.0001 0.001 0.01 0.1 mass [eV] excluded by Eot-Wash large quantum corrections (ez1) Using  $V_{,\phi}(\phi_{\rm min})=-\frac{\partial V}{\partial I}$  to relate  $\phi_{\rm min}$  and  $\rho$  :  $J_{,\phi}(\phi_\text{min})=-\frac{g\rho}{M_\text{D}}$  $M_{\rm Pl}$  $\phi_{\rm min}$  and  $\rho$  $m_\phi^{-1}$  $\phi$   $\sim$  $\gtrsim 27 \left( \frac{g \rho_{\rm lab}}{10} \right)$ 10 g cm−<sup>3</sup>  $\sqrt{-1/3}$  $\rm \mu m$ (Model-independent)

Quantum Stability<br>Reappears in Phys. Rev. Lett to appear in Phys. Rev. Lett.  $\frac{m_\phi^4(\phi)}{64\pi^2}\ln\Bigg(\frac{m_\phi^2}{\mu}\Bigg)$ "

 $\frac{2}{\phi}(\phi)$ 

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**E** Tension: Need small  $m_{\phi}$  to keep loop corrections under control, BUT chameleon mechanism relies on  $m_{\phi} \uparrow$  with  $\rho$ BUT chameleon mechanism relies on  $m_{\phi} \uparrow$  with  $\rho$ 

 $m_\phi^4(\phi)$ 

0.01 0.1 1 10 100 1000 10000 matter coupling  $g$  0.0001 0.001 0.01 0.1 mass [eV] excluded by Eot-Wash large quantum corrections (ez1) Using  $V_{,\phi}(\phi_{\rm min})=-\frac{\partial V}{\partial I}$  to relate  $\phi_{\rm min}$  and  $\rho$  :  $J_{,\phi}(\phi_\text{min})=-\frac{g\rho}{M_\text{D}}$  $M_{\rm Pl}$  $\phi_{\rm min}$  and  $\rho$  $m_\phi^{-1}$  $\phi$   $\sim$  $\gtrsim 27 \left( \frac{g \rho_{\rm lab}}{10} \right)$ 10 g cm−<sup>3</sup>  $\sqrt{-1/3}$  $\rm \mu m$ (Model-independent)

∴ Factor of  $\approx 2$  improvement in lab bounds on  $m_\phi^{-1}$  would close<br>∵ the window around  $\,g \sim 1$ the window around  $g \sim 1$  $m_\phi^{-1}$  $\approx 2$  improvement in lab bounds on  $m_\phi^-$ 

 $V_{\text{eff}}(\phi)$  $\phi$  $V(\phi)$  $~\sim \rho \phi$  $V(\phi) = \frac{M^5}{4} \qquad M = 10^{-3} \; \text{eV}$ Expand around minimum:  $V = \overline{V} + \ldots + \frac{\delta \phi^n}{\Delta n}$  $\frac{1}{\Lambda^{n-4}} + \ldots$ where  $\Lambda$  $\overline{M}$ =  $\int \bar{\phi}$  $\overline{M}$ "  $n+1$ *n*−4 =  $\int M^2$  $m^2$ " *n*+1 3(*n*−4)  $>$  $\bigwedge^{1/2}$  $m^2$ " 1 3 Cosmologically: Lab:  $m \sim \text{Mpc}^{-1}$  $\implies \Lambda \sim 10^5~{\rm GeV}$  $m \sim 10^{-3} \text{ eV} \implies \Lambda \sim 10^{-3} \text{ eV}$  $\phi$ Upadhye, Hu & Khoury, 1204.3906 [hep-ph] Quantum Stability (cont'd)

#### Conclusions

 If new forces are associated with dark sector, then some screening mechanism is required by local tests of gravity

Chameleon and Symmetron mechanisms rely on densitydependent mass and coupling, respectively.

 Rich phenomenology for laboratory, solar-system and cosmological tests of gravity

Open questions

 UV completion? cf. Hinterbichler, J. Khoury, & Nastase, JHEP (2011)

Symmetron topological defects?



