Center for Particle Cosmology at the University of Pennsylvania

Classical and Quantum Stability of Chameleon Theories Justin Khoury (UPenn)

w. Amol Upadhye, Junpu Wang, Wayne Hu, Lam Hui



A Richer Dark Sector

 Dark energy candidates:
 A, quintessence...
 Ratra & Peebles (1988); Wetterich (1988); Caldwell, Dave & Steinhardt (1998)



A Richer Dark Sector Dark energy candidates: A , quintessence... Ratra & Peebles (1988); Wetterich (1988); Caldwell, Dave & Steinhardt (1998)



Tantalizing prospect: dark sector includes new light fields (e.g. quintessence) that couple to both dark <u>and</u> baryonic matter.

Scalar fields can "hide" themselves from local exp'ts through screening mechanisms

$$\rho_{\rm here} \sim 10^{30} \rho_{\rm cosmos}$$

Large in a cosmological sense, but small in a particle physics sense



 ${\ensuremath{\,^{\circ}}}$ To neutralize Λ to accuracy H_0^2 , new degrees of freedom must be light:



 ${\ensuremath{\,^{\circ}}}$ To neutralize Λ to accuracy H_0^2 , new degrees of freedom must be light:

 $m_{\phi} \lesssim H_0$

These degrees of freedom must couple to Standard Model fields:

 ${\ensuremath{\,^{\circ}}}$ To neutralize Λ to accuracy H_0^2 , new degrees of freedom must be light:

 $m_{\phi} \lesssim H_0$

These degrees of freedom must couple to Standard Model fields:

 ϕ

 ${\ensuremath{\,^{\circ}}}$ To neutralize Λ to accuracy H_0^2 , new degrees of freedom must be light:

 $m_{\phi} \lesssim H_0$

These degrees of freedom must couple to Standard Model fields:

 ϕ

Must rely on screening mechanism for consistency with local tests of GR Experimental Program $U(r) = -g \frac{M}{8\pi M_{\rm Pl}^2} \frac{e^{-r/\lambda}}{r}$





Screening mechanisms have rich phenomenology for tests of GR:
Forced us to rethink implications of existing data
Inspired design of novel experimental tests

Chameleon/symmetron/dilaton: $\mathcal{L} = -\frac{1}{2}(\partial \phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$

Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:



Galileon/Vainshtein:

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 + \frac{1}{\Lambda^3} (\partial \phi)^2 \Box \phi + g \frac{\phi}{M_{\rm Pl}} T^{\mu}_{\ \mu}$$

Deffayet, Dvali, Gabadadze & Vainshtein (2001); Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

- Screening condition:

$$\left|\partial^2\phi\right| \gg \Lambda^3$$

Chameleon/symmetron/dilaton: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$

Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:



Kinetic/k-mouflage: Babichev, Deffayet & Ziour (2009)

- Screening condition:

 $|\partial \phi| \gg \Lambda^2$

Galileon/Vainshtein:

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 + \frac{1}{\Lambda^3} (\partial \phi)^2 \Box \phi + g \frac{\phi}{M_{\rm Pl}} T^{\mu}_{\ \mu}$$

 $\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - \frac{1}{4\Lambda^4} (\partial \phi)^4 + g \frac{\phi}{M_{\rm Pl}} T^{\mu}_{\ \mu}$

Deffayet, Dvali, Gabadadze & Vainshtein (2001); Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

- Screening condition:

$$\left|\partial^2\phi\right| \gg \Lambda^3$$

Chameleon/symmetron/dilaton: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - V(\phi) + A(\phi)T^{\mu}_{\mu}$

Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition: $\phi \ll \phi_c$
- Quantum stability?
- Kinetic/k-mouflage: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 \frac{1}{4\Lambda^4}(\partial\phi)^4 + g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\mu}$ Babichev, Deffayet & Ziour (2009)
 - Screening condition: $|\partial \phi| \gg \Lambda^2$
 - Shift symmetry: $\phi
 ightarrow \phi + c$

Solution Galileon/Vainshtein: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 + \frac{1}{\Lambda^3}(\partial\phi)^2 \Box \phi + g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\mu}$ Deffayet, Dvali, Gabadadze & Vainshtein (2001);

Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

– Screening condition: $|\partial^2 \phi| \gg \Lambda^3$

– Galileon symmetry: $\phi
ightarrow \phi + c + b_{\mu} x^{\mu}$

Chameleon/symmetron/dilaton:

 $\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - V(\phi) + A(\phi) T^{\mu}_{\mu}$ Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:
- $\phi \ll \phi_c$
- Quantum stability?
- Sinetic/k-mouflage: Babichev, Deffayet & Ziour (2009)
 - Screening condition:
 - Shift symmetry: $\phi \rightarrow \phi + c$
- Galileon/Vainshtein:

 $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 + \frac{1}{\Lambda^3}(\partial\phi)^2\Box\phi + g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\mu}$ Deffayet, Dvali, Gabadadze & Vainshtein (2001); Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

- Screening condition: $|\partial^2 \phi| \gg \Lambda^3$

– Galileon symmetry: $\phi
ightarrow \phi + c + b_{\mu} x^{\mu}$

- Cutoff: $\Lambda \sim {
m mm}^{-1}$

Cutoff: $\Lambda \sim \mathrm{mm}$

- Cutoff:
$$\Lambda \sim (1000 \ {
m km})^{-2}$$

$$\mathcal{L} = -\frac{1}{2}(\partial\phi) - \frac{1}{4\Lambda^4}(\partial\phi)$$

$$(\partial \phi)^4 + g \frac{\phi}{M_{\rm Pl}} T^{\prime}$$

$$\partial \phi | \gg \Lambda^2$$

$$|\partial \phi| \gg \Lambda^2$$

$$2/1 > \sqrt{2}$$

$$-\frac{1}{2}(O\phi)^2 - \frac{1}{4\Lambda^4}$$

$$(\partial\phi)^4 + g \frac{\phi}{M_{\rm Pl}} T^{\mu}$$

Chameleon/symmetron/dilaton:

 $\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - V(\phi) + A(\phi) T^{\mu}_{\mu}$ Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:
- $\phi \ll \phi_c$
- Quantum stability?
- Kinetic/k-mouflage: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 \frac{1}{4\Lambda^4}(\partial\phi)^4 + g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\mu}$ Babichev, Deffayet & Ziour (2009)
 - Screening condition: $|\partial \phi| \gg \Lambda^2$
 - Shift symmetry: $\phi \rightarrow \phi + c$

Solution Structure
Galileon/Vainshtein: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 + \frac{1}{\Lambda^3}(\partial\phi)^2\Box\phi + g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\mu}$ Deffayet, Dvali, Gabadadze & Vainshtein (2001); Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

– Screening condition: $|\partial^2 \phi| \gg \Lambda^3$

– Galileon symmetry: $\phi
ightarrow \phi + c + b_{\mu} x^{\mu}$ – Superluminality

- Cutoff: $\Lambda \sim {
m mm}^{-1}$

- Cutoff: $\Lambda \sim {
m mm}^{-1}$

- Cutoff: $\Lambda \sim (1000 \ {
m km})^{-1}$

- Superluminality

Chameleon/symmetron/dilaton:

 $\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - V(\phi) + A(\phi) T^{\mu}_{\mu}$ Khoury & Weltman (2003); Brax et al. (2004); Hinterbichler & Khoury (2010); Brax et al. (2010)

- Screening condition:
- Quantum stability?
- Kinetic/k-mouflage: $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 \frac{1}{4\Lambda^4}(\partial\phi)^4 + g\frac{\phi}{M_{\rm Pl}}T^{\mu}_{\mu}$ Babichev, Deffayet & Ziour (2009)
 - Screening condition: $|\partial^3 \phi| \gg \Lambda^4$
 - Shift symmetry:
- Galileon/Vainshtein:

 $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 + \frac{1}{\Lambda^3}(\partial\phi)^2\Box\phi + g\frac{\phi}{M_{\rm D}}T^{\mu}_{\mu}$ Deffayet, Dvali, Gabadadze & Vainshtein (2001) Luty, Porrati & Rattazzi (2003); Nicolis & Rattazzi (2004); Nicolis, Rattazzi & Trincherini (2008)

- Screening condition: $|\partial^2 \phi| \gg \Lambda^3$

– Galileon symmetry: $\phi
ightarrow \phi + c + b_{\mu} x^{\mu}$ – Superluminality

– Cutoff:
$$\Lambda \sim {
m mm}^{-1}$$

? - Cutoff:
$$\Lambda \sim \mathrm{mm}^{-2}$$
 - Superluminality

– Cutoff:
$$\Lambda \sim (1000 \ {
m km})$$

$$\phi \ll \phi_c$$

Should we be worried?

Should we be worried?



Should we be worried?



"The absence of evidence is <u>not</u> evidence of absence."

Chameleon Mechanism

Khoury & Weltman (2003); Gubser & Khoury (2004); Brax, van de Bruck, Davis, Khoury and Weltman (2004); Mota and Shaw (2006).



Consider scalar field ϕ with potential $V(\phi)$ and coupled to matter:

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - V(\phi) + g \frac{\phi}{M_{\rm Pl}} T^{\mu}_{\ \mu}$$

where $T^{\mu}_{\ \mu}$ is stress tensor of all matter (Baryonic and Dark)

Chameleon Mechanism

Khoury & Weltman (2003); Gubser & Khoury (2004); Brax, van de Bruck, Davis, Khoury and Weltman (2004); Mota and Shaw (2006).



Consider scalar field ϕ with potential $V(\phi)$ and coupled to matter:

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - V(\phi) + g \frac{\phi}{M_{\rm Pl}} T^{\mu}_{\ \mu}$$

 $V_{
m eff}(\phi)$

where $T^{\mu}_{\ \mu}$ is stress tensor of all matter (Baryonic and Dark)

For non-relativistic matter, $~T^{\mu}_{~\mu}\approx -\rho$, hence

$$\nabla^2 \phi = V_{,\phi} + \frac{g}{M_{\rm Pl}}\rho$$

$$V_{\rm eff}(\phi) = V(\phi) + g \frac{\phi}{M_{\rm Pl}} \phi$$

Density-dependent mass

$$V_{\rm eff}(\phi) = V(\phi) + g \frac{\phi}{M_{\rm Pl}} \rho$$

e.g.
$$V(\phi) = \frac{M^{4+n}}{\phi^n}$$



Thus $m=m(\rho)$ increases with increasing density Laboratory tests => set $m^{-1}(\rho_{
m local})\lesssim {
m mm}$

Generally implies: $m^{-1}(\rho_{\rm cosmos}) \lesssim {
m Mpc}$

Density-dependent mass

$$V_{\rm eff}(\phi) = V(\phi) + g \frac{\phi}{M_{\rm Pl}} \rho$$

e.g.
$$V(\phi) = \frac{M^{4+n}}{\phi^n}$$



Thus $m = m(\rho)$ increases with increasing density Laboratory tests => set $m^{-1}(\rho_{\text{local}}) \lesssim \text{mm}$

Generally implies: $m^{-1}(\rho_{\rm cosmos}) \lesssim {
m Mpc}$

Meanwhile, $m^{-1}(\rho_{\text{solar system}}) \lesssim 10 - 10^4 \text{ AU}$ \implies ruled out by post-Newtonian tests?

 $\rho = \rho_{\rm out}$

 $= \rho_{\rm in}$ R_{i}

 $\rho = \rho_{\rm out}$



 $\rho = \rho_{\rm out}$





$$\implies \qquad \phi(r > R) \sim \frac{1}{R} \times \frac{1}{r}$$
where
$$\frac{\Delta R}{R} = \frac{\phi_{\text{out}} - \phi_{\text{in}}}{6gM_{\text{Pl}}\Phi_{\text{N}}} \ll 1 \implies \text{thin-shell screening}$$



But small objects \implies no thin-shell

Thin-shell condition depends on environment!

Thin-shell condition depends on environment!

Thin-shell condition depends on environment!

 $G_{\rm N}^{\rm eff} = G_{\rm N}(1+2g^2)$ between small objects in space !

Smoking Guns

 $^{\it \odot}$ Satellite Energy Exchange (SEE) Mission $\frac{\Delta G_{\rm N}}{G_{\rm N}} < 10^{-6}$

$\frac{\Delta G_{\rm N}}{G_{\rm N}} \sim \mathcal{O}(1)$

Smoking Guns

${}^{\it \odot}$ Satellite Energy Exchange (SEE) Mission $\frac{\Delta G_{\rm N}}{G_{\rm N}} < 10^{-6}$

MICROSCOPE (2015)

Satellite Test of theEquivalence Principle (STEP)

 $\frac{\Delta G_{\rm N}}{G_{\rm N}} \sim \mathcal{O}(1)$ $\frac{\Delta a}{a} > 10^{-13}$

Symmetron Mechanism

K. Hinterbichler and J. Khoury, Phys. Rev. Lett. (2010); Olive and Pospelov (2008); Brax, van de Bruck, Davis and Shaw (2010).

Instead of m(
ho), here coupling to matter depends on density. ${\cal L}=-rac{1}{2}(\partial\phi)^2-V(\phi)+rac{\phi^2}{2M^2}T^\mu_{\ \mu}$

Symmetron Mechanism

K. Hinterbichler and J. Khoury, Phys. Rev. Lett. (2010); Olive and Pospelov (2008); Brax, van de Bruck, Davis and Shaw (2010).

 $V(\phi)$

Instead of m(
ho), here coupling to matter depends on density.

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - V(\phi) + \frac{\phi^2}{2M^2} T^{\mu}_{\ \mu}$$

Potential is of the spontaneous-symmetrybreaking form:

$$V(\phi) = -\frac{1}{2}\mu^{2}\phi^{2} + \frac{1}{4}\lambda\phi^{4}$$

Most general renormalizable potential with $\phi \to -\phi$ symmetry.

Symmetron Mechanism

K. Hinterbichler and J. Khoury, Phys. Rev. Lett. (2010); Olive and Pospelov (2008); Brax, van de Bruck, Davis and Shaw (2010).

 $V(\phi)$

Instead of m(
ho), here coupling to matter depends on density.

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - V(\phi) + \frac{\phi^2}{2M^2}T^{\mu}_{\ \mu}$$

Potential is of the spontaneous-symmetrybreaking form:

$$V(\phi) = -\frac{1}{2}\mu^{2}\phi^{2} + \frac{1}{4}\lambda\phi^{4}$$

Most general renormalizable potential with $\phi \to -\phi$ symmetry.

$$\implies V_{\text{eff}}(\phi) = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2\right) \phi^2 + \frac{1}{4}\lambda\phi^4$$

. Whether symmetry is broken or not depends on local density

Density-dependent coupling Perturbations $\delta\phi$ around local background value couple as: $\mathcal{L}_{
m coupling}\sim rac{ar{\phi}}{M^2}\delta\phi\,
ho$

 ${\ensuremath{\, o}}$ Symmetron decoupled in high-density regions (where $\,\phi\simeq 0$) ${\ensuremath{\, o}}$ In voids, where symmetry is broken,

 $V(\phi)$

$$\mathcal{L}_{\text{coupling}} \sim \frac{\mu}{\sqrt{\lambda}M^2} \delta \phi \rho$$

$$\sim \frac{\delta \phi}{M_{\text{Pl}}} \rho$$
gravitational strength
$$\text{OTE: Tests of gravity} \implies \mu^{-1} \sim \text{Mpc}$$

Inspiration...

Inspiration...

Symmetron Couch (\$9500.00)

"NASA-style gravity reduction."

"Offers a unique multi-phase wave experience."

Inspiration...

Symmetron Couch (\$9500.00)

"NASA-style gravity reduction."

"Offers a unique multi-phase wave experience."

Chameleon Searches

Eot-Wash

Adelberger et al., Phys. Rev. Lett. (2008)

CHameleon Afterglow SEarch (CHASE), Fermilab

Chou et al., Phys. Rev. Lett. (2008,2010)

Chameleon Searches (cont'd)

Axion Dark Matter eXperiment (ADMX)

P. Sikivie & co., Phys. Rev. Lett. (2010)

Astrophysical signatures

 Macroscopic violations of the Equivalence
 Principle Hui, Nicolis & Stubbs (2009); Jain & Vanderplas (2011)

Modified stellar evolution

Chang & Hui (2010); Davis, Lim, Sakstein & Shaw (2011); Jain, Vikram & Sakstein (2012)

cf. Eugime Lim's talk

Astrophysical signatures (cont'd)

Saturn's rings

Enhanced self-gravity due to chameleon should impact propagation of density waves.

$$v_g = \frac{\pi G_{\rm N} \sigma_0}{\kappa}$$

 $\sigma_0 \equiv$ surface density $\kappa \equiv$ epicyclic frequency

Self-Acceleration?

Wang, Hui & J. Khoury, to appear

 $\tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}$

"Jordan frame"

"Einstein frame"

Can cosmic acceleration result from $\Delta A \sim \mathcal{O}(1)$ even though Einstein-frame metric is NOT accelerating?

Self-Acceleration?

Wang, Hui & J. Khoury, to appear

 $\tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}$ Frame" "Einstein frame"

"Jordan frame"

Can cosmic acceleration result from $\Delta A \sim \mathcal{O}(1)$ even though Einstein-frame metric is NOT accelerating?

Take chameleon with $A(\phi) = 1 + \frac{g\phi}{M_{\rm Pl}} + \dots$

Source condition: $\frac{F_{\phi}}{F_{\rm N}} = 2g^2 \quad \text{in unscreened regions} \implies g \lesssim \mathcal{O}(1)$ Screening condition: $\frac{\Delta R}{R} \approx \frac{\phi_0}{6 a M_{\rm Pl} \Phi_{\rm N}} < 1$ Milky Way potential $\Phi_{\rm N} \sim 10^{-6} \implies \phi_0 \lesssim 10^{-6} M_{\rm Pl}$

Self-Acceleration?

Wang, Hui & J. Khoury, to appear

 $\tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}$ "Einstein frame"

"Jordan frame"

Can cosmic acceleration result from $\Delta A \sim \mathcal{O}(1)$ even though Einstein-frame metric is NOT accelerating?

Take chameleon with $A(\phi) = 1 + \frac{g\phi}{M_{\rm Pl}} + \dots$

Force condition: $\frac{F_{\phi}}{F_{\rm N}} = 2g^2 \quad \text{in unscreened regions} \implies g \lesssim \mathcal{O}(1)$ Screening condition: $\frac{\Delta R}{R} \approx \frac{\phi_0}{6 a M_{\rm Pl} \Phi_{\rm N}} < 1$ Milky Way potential $\Phi_{\rm N} \sim 10^{-6} \implies \phi_0 \lesssim 10^{-6} M_{\rm Pl}$. $\Delta A \lesssim g rac{\phi_0}{M_{
m Pl}} \ll 1$ (NO self-acceleration)

The Mpc Barrier Wang, Hui & J. Khoury, to appear Can scalar field have $m_\phi\sim H_0$ for $\Delta t\gtrsim H_0^{-1}$?

The Mpc BarrierWang, Hui & J. Khoury, to appearCan scalar field have $m_{\phi} \sim H_0$ for $\Delta t \gtrsim H_0^{-1}$?Cosmological evolution: $\ddot{\phi} + 3H\dot{\phi} = -\frac{\mathrm{d}V}{\mathrm{d}\phi} - \frac{g}{M_{\mathrm{Pl}}}\rho$

The Mpc BarrierWang, Hui & J. Khoury, to appearCan scalar field have $m_{\phi} \sim H_0$ for $\Delta t \gtrsim H_0^{-1}$?Cosmological evolution: $\ddot{\phi} + 3H\dot{\phi} = -\frac{\mathrm{d}V}{\mathrm{d}\phi} - \begin{pmatrix}g\\M_{\mathrm{Pl}}\\\end{pmatrix}\\\sim H^2 M_{\mathrm{Pl}}\\\therefore$ Pretty large force!

Hence ${
m d}V/{
m d}\phi$ & $g
ho/M_{
m Pl}$ must cancel to good accuracy

The Mpc BarrierWang, Hui & J. Khoury, to appearCan scalar field have $m_{\phi} \sim H_0$ for $\Delta t \gtrsim H_0^{-1}$?Cosmological evolution: $\ddot{\phi} + 3H\dot{\phi} = -\frac{\mathrm{d}V}{\mathrm{d}\phi} - \underbrace{\begin{pmatrix}g\\M_{\mathrm{Pl}}\\\end{pmatrix}}_{M_{\mathrm{Pl}}} \\ \sim H^2 M_{\mathrm{Pl}} \\ \therefore$ Pretty large force!

Hence ${
m d}V/{
m d}\phi$ & $g
ho/M_{
m Pl}$ must cancel to good accuracy

Cancellation must hold over last Hubble time:

$$\Delta \left(\frac{\mathrm{d}V}{\mathrm{d}\phi} + \frac{g}{M_{\mathrm{Pl}}} \rho \right) \sim H_0^2 M_{\mathrm{Pl}} \sim m_0^2 \Delta \phi$$

$$\implies$$

$$m_0 \sim H_0 \sqrt{\frac{M_{\rm Pl}}{\Delta \phi}} > 10^3 \ H_0$$

Argument generalizes to very wide class of chameleon/symmetron theories, including many fields!

Quantum Stability

Upadhye, Hu & Khoury, 1204.3906 [hep-ph], to appear in Phys. Rev. Lett.

Focus on scalar loops:

$$\Delta V_{1-\text{loop}} = -\frac{m}{6}$$

$$\frac{(\phi)}{\pi^2} \ln\left(\frac{m_\phi^2(\phi)}{\mu_0^2}\right)$$

Tension: Need small m_ϕ to keep loop corrections under control, BUT chameleon mechanism relies on $m_\phi \uparrow$ with ρ

Quantum Stability

Upadhye, Hu & Khoury, 1204.3906 [hep-ph], to appear in Phys. Rev. Lett.

Focus on scalar loops: $\Delta V_{1-\text{loop}} = \frac{m_{\phi}^4(\phi)}{64\pi^2} \ln\left(\frac{m_{\phi}^2(\phi)}{\mu_0^2}\right)$

Tension: Need small m_{ϕ} to keep loop corrections under control, BUT chameleon mechanism relies on $m_{\phi} \uparrow$ with ho

Using $V_{,\phi}(\phi_{\min}) = -\frac{g\rho}{M_{\rm Pl}}$ to relate ϕ_{\min} and ρ : 0.1 $m_{\phi}^{-1} \gtrsim 27 \left(\frac{g\rho_{\text{lab}}}{10 \text{ g cm}^{-3}}\right)^{-1/3} \mu \text{m}$ 0.01 mass [eV] 0.001

(Model-independent)

Quantum Stability

Upadhye, Hu & Khoury, 1204.3906 [hep-ph], to appear in Phys. Rev. Lett.

Focus on scalar loops:

:
$$\Delta V_{1-\text{loop}} = \frac{m_{\phi}^4(\phi)}{64\pi^2} \ln\left(\frac{m_{\phi}^2(\phi)}{\mu_0^2}\right)$$

Tension: Need small m_ϕ to keep loop corrections under control, BUT chameleon mechanism relies on $m_\phi \uparrow$ with ho

Using $V_{,\phi}(\phi_{\min}) = -rac{g
ho}{M_{
m Pl}}$ to relate ϕ_{\min} and ho : 0.1 large quantum corrections (E>1) $m_{\phi}^{-1} \gtrsim 27 \left(\frac{g \rho_{\text{lab}}}{10 \text{ g cm}^{-3}} \right)^{-1/3} \mu \text{m}$ 0.01 mass [eV] 0.001 excluded by Eot-Wash (Model-independent) 0.0001 0.1 0.01 10 100 1000 10000 matter coupling ${\it g}$

Factor of $\approx 2\,$ improvement in lab bounds on $\,m_{\phi}^{-1}$ would close the window around $\,g\sim 1\,$

Quantum Stability (cont'd) $V_{\rm eff}(\phi)$ Upadhye, Hu & Khoury, 1204.3906 [hep-ph] $V(\phi) = \frac{M^3}{\phi} \qquad M = 10^{-3} \text{ eV}$ $V(\phi)$ Expand around minimum: $V = \bar{V} + \ldots + \frac{\delta \phi^n}{\Lambda n - 4} + \ldots$ where $\frac{\Lambda}{M} = \left(\frac{\bar{\phi}}{M}\right)^{\frac{n+1}{n-4}} = \left(\frac{M^2}{m^2}\right)^{\frac{n+1}{3(n-4)}} > \left(\frac{M^2}{m^2}\right)^{\frac{1}{3}}$ $m \sim \mathrm{Mpc}^{-1}$ $\implies \Lambda \sim 10^5 \text{ GeV}$ Cosmologically: $m \sim 10^{-3} \,\mathrm{eV} \implies \Lambda \sim 10^{-3} \,\mathrm{eV}$ @ Lab:

Conclusions

If new forces are associated with dark sector, then some screening mechanism is required by local tests of gravity

Chameleon and Symmetron mechanisms rely on densitydependent mass and coupling, respectively.

Rich phenomenology for laboratory, solar-system and cosmological tests of gravity

Open questions

UV completion? cf. Hinterbichler, J. Khoury, & Nastase, JHEP (2011)

Symmetron topological defects?

