

Detectors at LHC part.I

1. Snapshots on Physics Requirements at LHC

2. Experimental Challenges Machine Parameters Requirements for Detectors

3. Overall ATLAS and CMS Detector Design General Layout Tracking Calorimetry (Tomorrow) Muons





Bibliography Experimental Challenges in High-Luminosity Collider Physics N. Ellis and T. Virdee (Ann. Rev. Nucl. Part. Sci. 44 (1994) 609

D. Fournier and L. Serin, Experimental Techniques, European School of Particle Physics, CERN 96-04

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CERN Academic Training Lectures

ATLAS and CMS outreach pages

Important Lecture Note In this lecture I use many exemples from CMS, only because of my better knowledge of this experience. This must not be taken as a ranking between ATLAS and CMS.





Snapshots on Physics Requirements at LHC



SM Hiaas Boson

Fully hadronic final states dominate but cannot be used due to large QCD bkg. alook for final states with isolated leptons and photons despite smaller BR

At the LHC the SM Higgs provides a good benchmark to test the performance of a detector





Most promising channel for mH < 150 GeV

H
$$t^{*,W^{*}}$$
 γ
 $t^{*,W^{*}}$ $t^{*,W^{*}}$ γ

(σ .B ~ 50.10-3 pb @ mH ~ 150 GeV) a Signal: ~ 1000's of events/yr Backgrounds are large (2pb/GeV), H natural width is small (~MeV) a **excellent mass resolution** required

 $\sigma m/m = 0.5 [\sigma E1/E1 \otimes \sigma E2/E2^{X} O \cot(\theta/2)\Delta\theta]$ a energy resolution and precise vertex localisation

Typical Cuts

2 isolated photons: pT > 25, 40 GeV with $|\eta| < 2.5$

No track or em cluster with pT > 2.5 GeV in a cone size $\Delta R = 0.3$ around

 γ S Good energy resolution, measurement of photon direction, π 0 rejection, efficient photon isolation



SM HaZZ* or ZZ a 4I

130 < MH < 800 GeV

ΖZ

ΓH (MH=150 GeV) ~ 15 MeVObserved width is dominatedby instrumental mass resolution

ZZ*

 Di-muon or di-electron mass resolution should be better than
 □ Z

• Good momentum resolution for low momenta leptons

Large geometric acceptance
 Efficient lepton isolation at hi-L

 Γ H (MH=500 GeV) ~ 65 GeV For MH>350 GeV observed width is dominated by natural width





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H, h \gamma \gamma, b \overline{D} (H b \overline{D} in WH, t \overline{T} H)
h \gamma\gamma in WH, t Th 1\gamma\gamma
h, H ZZ*, ZZ 41
h, H, A \tau+\tau- (e/\mu)+ + h– + ETmiss
                        e+ + \mu- + ETmiss
                                                             inclusively and in bbHSUSY
                                   h++h-+FTmiss
H+ \tau+ \nu from t f
H+ \tau+ \nu and H+ t \overline{D} for MH>Mtop
A Zh with h bb; A \gamma\gamma
H, A \tilde{\chi}02\tilde{\chi}02, \tilde{\chi}0i\tilde{\chi}0j, \tilde{\chi}+i\tilde{\chi}-j
                                                                  Abundance of b, \tau
                                                                  and significant
H+ \tilde{\chi}+2\tilde{\chi}02
                                                                  Etmiss
qq qqH with H \tau+\tau-
                                                       Physics with Jets, (Vertex (pixel
H \tau\tau, in WH, t T
                                                       detector), Detector Hermiticity
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Experimental Challenges Machine Parameters



Enerav Frontier

New Energy Domain

Search for the unexpected Cover domain ~ 1 TeV in which SM without the Higgs (or equivalent) gives nonsense

Exploratory machine required

a hadron-hadron collider with: Largest possible primary energy Size of the tunnel x Max B field Largest possible luminosity: interesting and easily detectable final states involve leptons and photons with low σ .BR

Hadron Colliders can provide these BUT At the expense of 'clean' experimental conditions



Year of First Physics

LHC Lavout and Parameters

$$L = \frac{\gamma f k_b N_p^2}{4\pi \epsilon \beta^*} F$$

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f revolution frequency
kb no. of bunches
Np no. of protons/bunch
εn norm transverse emittance
β* betatron function
F reduction factor xing angle

Magnetic Field p (TeV) = 0.3 B(T) R(km) For p= 7 TeV, R= 4.3 km **a B = 8.4 T**

Energy at collision	Е	7	TeV
Dipole field at 7 TeV	В	8.33	т
Luminosity ² s- ¹	L	1034	cm-
Beam beam parameter	ξ	3.6	10-3
DC beam current A	lbea	m	0.56
Bunch separation		24.95	ns
No. of bunches	kb	2835	
No. particles per bunch	Np	1.1	1011
Normalized transverse	εn	3.75	μm
emittance (r.m.s.)			
Collisions			
β-value at IP	β*	0.5	m
r.m.s. beam radius at IP	σ*	16	um

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Experimental Challenges Requirements for Detectors

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Hadron colliders are broad-band exploratory machines May need to study WL-WL scattering at a cm energy of ~ 1 TeV



a LHC: pp collisions at 7 + 7 TeV

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Event Rate = L.\sigma.BR
e.g. H(1 TeV) g ZZ g 2e+2µ or 4e or 4µ
For L ~1034 , Evts/yr = 1034 10-37.10-3.107 ~ 10 /yr !!
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2. Detector Challenges I

Multipurpose :

Good Performances for leptons, jets, muons,..... Fast decaying particles (vertex) Hermitic ET_miss High Interaction Rate pp interaction rate 109 interactions/s Level-1 trigger decision will take ~2-3 μs a electronics need to store data locally (pipelining) data for only ~100 out of the 40 million crossings can be recorded per sec a fast and accurate high-level trigger and data acquisition (DAQ)

Large Particle Multiplicity

- \sim <20> superposed events in each crossing
- \sim 1000 tracks stream into the detector every 25 ns

need highly granular detectors with good time resolution for low occupancy

a large number of channels

High Radiation Lovels



Very good muon identification and momentum measurement trigger efficiently and measure sign of a few TeV muons

High energy resolution electromagnetic calorimertry With good isolation performance

Powerful inner tracking systems factor 10 better momentum resolution than at LEP With a factor 100 in sensitive surface High granularity

Hermetic calorimetry good missing ET resolution

Production Cross-sections



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At sqrt(s)=14 TeV σ tot~ 105 mb σ elastic~ 28 mb σ single diffractive~ 12 mb σ double diffractive~ few mb σ inel~ 65 mb

Evt rate = $L.\sigma = 1034 \times 65 \ 10-27 \ /s$ = $6.5 \times 108 \ /s$

Not all bunches are full (2835/3564) a events/crossing ~ 20

Operating Conditions

For every 'good' event containing a Higgs decay there are ~ 20 extra 'unwanted' minimum bias interactions superposed



Radiation Levels: Dose

Dose (Gy) in CMS for an integrated luminosity of 5.105 pb-1 (~ 10 years)



Les HOUCHES AMARAdiation Levels: Neutron Fluence

uence (E>100 keV) in CMS for an integrated luminosity of 5.105 pb-1 (~ 10 y



3.0E+17 1.0E+17 1.0E+16 1.0E+15 1.0E+14 1.0E+13 1.0E+12 1.0E+11 1.0E+10 1.0E+09 1.0E+08 2.9E+07









ATLAS and CMS Detector Design General Layout



'Cylindrical Onion-like ' Structure

The ATLAS Detector



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The CMS Detector





A Slice through CMS







ATLAS and CMS Detector Design Tracker



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Event Pileup



Answer: High Granularity

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Measurement of Momentum II

$$\frac{dp_T}{p_T} = \frac{\sigma_s}{s} = \frac{\sqrt{(3/2)} \sigma_x}{s}$$
$$\frac{dp_T}{p_T} = \frac{\sqrt{3}}{2} \sigma_x \frac{8p_T}{0.3BL^2}$$
(2)

Momentum resolution degrades linearly with increasing momentum, improves for higher field and the larger radial size of tracking cavity (quadratic in L)

Arrangement of measuring points

Uniform spacing

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$$\frac{dp_T}{p_T} = \frac{\sigma_x p_T}{0.3BL^2} \sqrt{\frac{720}{N+4}}$$

e.g. dp_T/p_T \approx 0.5% for p_T=1 GeV/c, L=1m, B=1T, σ_x = 200 µm and N=10

$$\frac{dp}{p}\Big|_{ms} \approx 0.05 \frac{1}{B\sqrt{LX_0}}$$

BUT in a real tracker errors due to multiple scattering has to be included.

Measurement of Momentum III



Figure 23.5: Quantities used to describe multiple Coulomb scattering. The particle is incident in the plane of the figure.

Apparent sagitta due to multiple scattering

$$s_p = \frac{L\theta_0}{4\sqrt{3}}$$

Relative momentum resolution due to multiple scattering is

$$\therefore \frac{s_p}{s} = \frac{dp}{p}\Big|_{ms} \approx 0.05 \frac{1}{B\sqrt{LX_0}} \quad \text{since } s = \frac{0.3BL^2}{8p} \qquad \text{B in T, L and X_0 in m}$$

i.e. Resolution is independent of p and α 1/B

If extrapolation error from one plane to next is larger than the point resolution then momentum resolution is degraded i.e. if



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ATLAS Tracker Lavout



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Hybrid Pixels: ~ 2.3 m2 of silicon sensors, 140 M pixels, 50x300 μ m2, r = 5, 10, 13 cm Si μ -strips : 60 m2 of silicon sensors, 6 M strips, 4 pts, r = 30 - 50 cm Straws TRT: 36 straws/track, Xe-CO2-CF4 ϕ =4mm, r = 56 - 107 cm

Electron Identification using Tracker in ATLAS (I)

Transition radiation is emitted when a charged particle moves from a medium of refractive index n_1 to a medium of a different index n_2 (= an apparent acceleration)

$$p = \gamma m v \implies m = \frac{1}{\beta c \gamma} p$$

$$friction = \frac{\Delta m}{m} \int_{-\infty}^{\infty} \frac{1}{\beta^2 c^2} \left(\frac{\Delta \gamma}{\gamma}\right)^2 + \left(\frac{\Delta p}{p}\right)^2 \quad \text{If } \Delta p/p \text{ is small, mass resolution}$$

$$friction = \frac{\Delta m}{m} \int_{-\infty}^{\infty} \frac{1}{\beta^2 c^2} \left(\frac{\Delta \gamma}{\gamma}\right)^2 + \left(\frac{\Delta p}{p}\right)^2 \quad \text{at high momenta is } \alpha \gamma$$

Radiated energy /boundary to vacuum

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$$W = \frac{1}{3} \alpha \ \hbar \omega_p \ \gamma \quad i.e. \ W \propto \gamma \qquad E_{ph} = \frac{\gamma}{3} \ \hbar \omega_p$$

where $\ \hbar \omega_p$ is the plasma frequency ($\approx 20 \ eV$ for polyethylene)
Hence can use it for particle identification

- X-rays are emitted at small angle ($\theta \approx 1/\gamma$). TR stays close to the charge particle track
- Number of emitted photon/boundary is small $N_{ph} \approx \frac{W}{\hbar\omega_p} \propto \alpha \approx \frac{1}{137}$!
- Need many transitions → stack of many thin foils with high Z gas to absorb the X-rays

Electron Identification using Tracker in ATLAS (II)

4mm diameter ttraw tube proportional chambers embedded in polyethylene fibres Gas 70% Xe + 20% CF₄ + 10%CO₂

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ATLAS Tracker Performance



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 $\sigma(pT)/pT \sim 0.6+18pT$ (pT in TeV)

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SCT barrel system test



Two of the SCT barrel support structures



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CSS99

Silicon Tracker



Schematic cross-section through a silicon microstrip detector. Diffusion distributes charge over multiple strips and capacitive charge division between readout amplifiers allows position interpolation.

This implies analog buffering and readout

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Challenges: Front-end Electronics



Each microstrip is read out by a charge sensitive amplifier with τ = 50 ns. The o/p voltage is sampled at the beam crossing rate of 40 MHz. Samples are stored in analog pipeline for up to Level-1 latency (3.2 µs) Following a trigger form a weighted sum of 3 samples in analog circuit This confines signal to single bx and gives pulse height Buffered pulse height data multiplexed out on optical fibres Output of laser modulated by the pulse height for each strip



CMS Tracker Performance



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To minimize the leakage currents induced by radiation (neutron) The whole tracker must be kept at low temperature \sim -18 deg.Celsius This requires not only a cooling but also a environment screen. Both bringing material....



Material in Trackers



Also very important figures for Electromagnetic Calorimeter which is behin Through electron bremstrallung and photon conversion





ATLAS and CMS Detector Design Muon System

Identification of Muons

Muons identified by their penetration through about 10 λ of calorimeter material The material of calorimeters absorbs the e's, γ 's and h[±]. Hadron Punchthough

Energy Loss in Absorber

- for $E_{_{\!\!\!\!\!\!\!\!\!\!\!}} \leq 20\text{--}30~\text{GeV}$ energy loss fluctuations dominate
- high energy muons generate their own background.

Hard bremstrahlung (catastrophic energy loss can spoil μ -tracking. The critical energy for μ in Fe is $E_c \approx 350$ GeV.

Hadron Punch-through

Debris from hadronic showers can accompany muons leading to:

- mis-identification of hadron as μ
- confusion and difficulty in matching $\mu\text{-tracks}$ (in jets)
- increase in µ trigger rate



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Each detector has 3 stations. Each station consists of 2-4 layers.



Precision chambers

Monitored Drift Tubes ($|\eta| < 2$) with a single wire resolution of 80 µm 1194 chambers, 5500m2 Cathode Strip Chambers (2 < $|\eta| < 2.7$) at higher particle fluxes 32 chambers, 27 m2

Trigger chambers

Resistive Plate Chambers ($|\eta| < 1.05$) with a good time resolution of 1 ns 1136 chambers, 3650 m2 Thin Gap Chambers (1.05 < $|\eta| < 2.4$) at higher particle fluxes 1584 chambers, 2900 m2

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End-cap MDT chamber



Figure 5-2 Drift tube operation in a magnetic field with curved drift path.

 ϕ wire = 50µm (W-Re) 3 bar, 3270V, td = 500ns Gas gain = 2.104

Measured Spatial Resolution



Figure 5-4 MDT resolution as a function of the drift distance, to an Ar/N2/CH4 (91/4/5 mixture). The curves correspond to two discriminator threshold settings. ATLAS Muon System: Performance

Multiple scattering 12 Contribution to resolution (%) Energy loss fluctuations 11 Chamber resol and align. 10 9 ∧ Total 8 $|\eta| < 1.5$ 7 6 5 4 3 2 1 0 10² 10³ 10 p_⊤ (GeV)

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Resolution limited by : m.m. and Energy Loss Fluct. @ 3% for 10 < pT < 250 GeV/cChamber Resolution and Alignment for pT > 250 GeV/c

The muon spectrometer resolution dominates for pT > 100 GeV/c





CMS Muon Detectors



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Tracking in Magnetized Iron (CMS) (in multiple scattering dominated regime)

$$\frac{\Delta p}{p} \approx \frac{40\%}{B\sqrt{L}}$$

B~1.8T, and L \approx 1.5 m BUT measurement is much better !

p resolution worsens at $|\eta| > 1.5$



Detectors at LHC part.l end

