



DIRECT CP ASYMMETRIES IN HADRONIC D DECAYS

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5th FCPPL workshop

arXiv:1203.3120[hep-ph]

Outline

- ◆ Motivation
- ◆ Branching Ratios
- ◆ Penguin parameterization
- ◆ Predict direct CP asymmetries
- ◆ Summary

Evidence of CPV

- ◆ First evidence of CP violation in charmed meson decays by LHCb, with 3.5σ [PRL 108 (2012) 111602]

$$\begin{aligned}\Delta A_{CP}^{exp} &\equiv A_{CP}(K^+K^-) - A_{CP}(\pi^+\pi^-) \\ &= [-0.82 \pm 0.21(\text{stat}) \pm 0.11(\text{syst})]\%\end{aligned}$$

- ◆ Confirmed by CDF, with 2.7σ [CDF note 10784]

$$\Delta A_{CP} = [-0.62 \pm 0.21(\text{stat}) \pm 0.10(\text{syst})]\%$$

- ◆ Naively expected much smaller in the SM

$$A_{CP}^{dir} \sim \frac{|V_{cb}^* V_{ub}|}{|V_{cs}^* V_{us}|} \frac{\alpha_s}{\pi} \sim 10^{-4}$$

- ◆ Necessary to predict more precisely in the SM.

Dynamics of D decays

- ◆ To predict CPV, we have to well understand the important decay mechanism.
 - ◆ At first, branching ratios should be well explained
- *A long-standing puzzle:*

$$R = \frac{Br(D^0 \rightarrow K^+ K^-)}{Br(D^0 \rightarrow \pi^+ \pi^-)} \approx 2.8$$

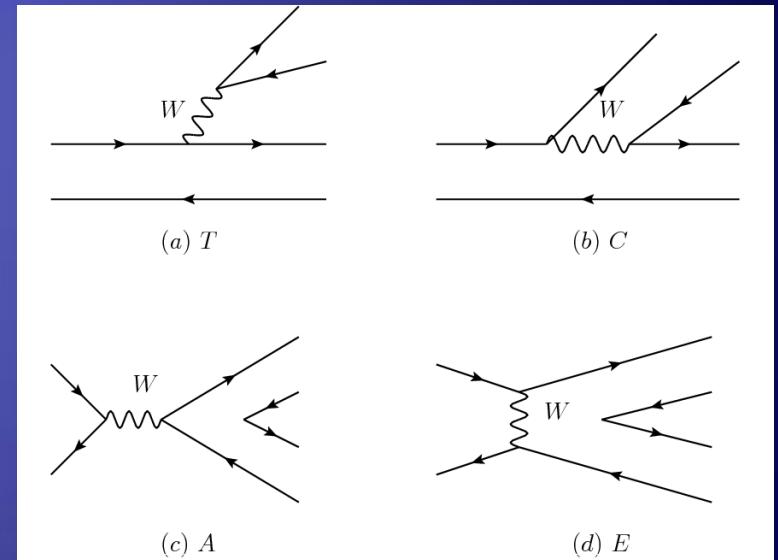
- ◆ $R=1$ in the SU(3) flavor symmetry limit
- ◆ Imply large **SU(3) breaking** effects

Guidelines

- ◆ Framework: topological amplitudes in **factorization** hypothesis:
 - ◆ Short-distance dynamics: Wilson coefficients
 - ◆ Long-distance dynamics: hadronic matrix elements
- ◆ Introduce **non-perturbative** parameters in hadronic matrix elements
 - ◆ Charm mass $m_c \approx 1.3 \text{ GeV}$, not reliable for perturbation in coupling constant and $1/m_c$
- ◆ Include large **SU(3) breaking effects.**

Topological diagrams for BRs

- ◆ According to weak interactions and flavor flows
- ◆ Include all strong interaction effects, involving FSI
- ◆ This is a complete set
- ◆ Topological diagrammatic approach was studied in the flavor SU(3) symmetry limit



[Cheng & Chiang, 2010]

➤ How SU(3) breaking ??

Mode-dependent dynamics

SU(3) breaking effects

- ◆ **Evolution of Wilson coefficients** depending on masses of decay products

- Difference of mass ratios $\frac{m_\pi}{m_D}$ vs $\frac{m_K}{m_D}$
mK~0.5 GeV
mD~1.8 GeV
- Especially for the modes with η' involved M η' ~1 GeV

- ◆ **Glauber strong phase** associated with **pion** in nonfactorizable amplitudes [H.n Li, S. Mishima, 2009]

- Pion: as a massless Goldstone boson, and as a $q\bar{q}$ bound state ?

Nambu-Goldstone bosons

- Pion : massless Goldstone boson, and $q\bar{q}$ bound state?
 - **Massless** boson => huge spacetime
=> large separation between $q\bar{q}$
=> **high mass** due to confinement => contradiction!
 - Reconciliation : Tight bound $q\bar{q}$, but multi-parton
=> soft cloud (Lepage, Brodsky 79; Nussinov, Shrock 08; Duraisamy, Kagan 08)
 - **Glauber phase** corresponds to soft cloud [H.n Li, S. Mishima, 09]
 - Pion is unique
 - **SU(3) breaking** effects: distinguish **pions** from other final states

Branching ratios

Include important dynamics

For SU(3) breaking effects

Cabibbo-favored BRs

well consistent with data

Modes	Br(FSI)	Br(diagram)	Br(pole)	Br(exp)	Br(this work)
$D^0 \rightarrow \pi^0 \bar{K}^0$	1.35	2.36 ± 0.08	2.4 ± 0.7	2.38 ± 0.09	2.37
$D^0 \rightarrow \pi^+ K^-$	4.03	3.91 ± 0.17	3.9 ± 1.0	3.891 ± 0.077	3.71
$D^0 \rightarrow \bar{K}^0 \eta$	0.80	0.98 ± 0.05	0.8 ± 0.2	0.96 ± 0.06	1.00
$D^0 \rightarrow \bar{K}^0 \eta'$	1.51	1.91 ± 0.09	1.9 ± 0.3	1.90 ± 0.11	1.68
$D^+ \rightarrow \pi^+ \bar{K}^0$	2.51	3.08 ± 0.36	3.1 ± 2.0	3.074 ± 0.096	3.17
$D_S^+ \rightarrow K^+ \bar{K}^0$	4.79	2.97 ± 0.32	3.0 ± 0.9	2.98 ± 0.08	2.99
$D_S^+ \rightarrow \pi^+ \eta$	1.33	1.82 ± 0.32	1.9 ± 0.5	1.84 ± 0.15	1.65
$D_S^+ \rightarrow \pi^+ \eta'$	5.89	3.82 ± 0.36	4.6 ± 0.6	3.95 ± 0.34	3.44
$D_S^+ \rightarrow \pi^+ \pi^0$		0	0	< 0.06	0

Cabibbo-suppressed BRs

our advantage

Modes	Br(FSI)	Br(diagram)	Br(pole)	Br(exp)	Br(this work)
$D^0 \rightarrow \pi^+ \pi^-$	1.59	2.24 ± 0.10	2.2 ± 0.5	1.45 ± 0.05	1.44 ↙
$D^0 \rightarrow K^+ K^-$	4.56	1.92 ± 0.08	3.0 ± 0.8	4.07 ± 0.10	4.19 ↙
$D^0 \rightarrow K^0 \bar{K}^0$	0.93	0	0.3 ± 0.1	0.320 ± 0.038	0.35
$D^0 \rightarrow \pi^0 \pi^0$	1.16	1.35 ± 0.05	0.8 ± 0.2	0.81 ± 0.05	0.55
$D^0 \rightarrow \pi^0 \eta$	0.58	0.75 ± 0.02	1.1 ± 0.3	0.68 ± 0.07	0.94
$D^0 \rightarrow \pi^0 \eta'$	1.7	0.74 ± 0.02	0.6 ± 0.2	0.91 ± 0.13	0.64
$D^0 \rightarrow \eta \eta$	1.0	1.44 ± 0.08	1.3 ± 0.4	1.67 ± 0.18	1.48
$D^0 \rightarrow \eta \eta'$	2.2	1.19 ± 0.07	1.1 ± 0.1	1.05 ± 0.26	1.52
$D^+ \rightarrow \pi^+ \pi^0$	1.7	0.88 ± 0.10	1.0 ± 0.5	1.18 ± 0.07	0.88
$D^+ \rightarrow K^+ \bar{K}^0$	8.6	5.46 ± 0.53	8.4 ± 1.6	6.12 ± 0.22	5.97
$D^+ \rightarrow \pi^+ \eta$	3.6	1.48 ± 0.26	1.6 ± 1.0	3.54 ± 0.21	3.37 ↙
$D^+ \rightarrow \pi^+ \eta'$	7.9	3.70 ± 0.37	5.5 ± 0.8	4.68 ± 0.29	4.54 ↙
$D_S^+ \rightarrow \pi^0 K^+$	1.6	0.86 ± 0.09	0.5 ± 0.2	0.62 ± 0.23	0.65
$D_S^+ \rightarrow \pi^+ K^0$	4.3	2.73 ± 0.26	2.8 ± 0.6	2.52 ± 0.27	2.21
$D_S^+ \rightarrow K^+ \eta$	2.7	0.78 ± 0.09	0.8 ± 0.5	1.76 ± 0.36	1.00
$D_S^+ \rightarrow K^+ \eta'$	5.2	1.07 ± 0.17	1.4 ± 0.4	1.8 ± 0.5	1.92

Puzzle $D^0 \rightarrow \pi^+ \pi^-$ vs $D^0 \rightarrow K^+ K^-$

- Revisited: $R_{\text{exp}} = 2.8$, $R = 1$ in SU(3) flavor symmetry

Modes	Br(FSI)	Br(diagram)	Br(pole)	Br(exp)	Br(this work)
$D^0 \rightarrow \pi^+ \pi^-$	1.59	2.24 ± 0.10	2.2 ± 0.5	1.45 ± 0.05	1.44
$D^0 \rightarrow K^+ K^-$	4.56	1.92 ± 0.08	3.0 ± 0.8	4.07 ± 0.10	4.19

- Large SU(3) breaking effects
- Glauber phase associated with pions
 - dominate source of the difference between these two modes

Must Explain BRs well
otherwise, some important dynamics
may be missed



penguin parameterization and
predict direct CP asymmetries

Direct CP asymmetry

- ◆ Definition:

$$A_{CP}(f) = \frac{\Gamma(D \rightarrow f) - \Gamma(\bar{D} \rightarrow \bar{f})}{\Gamma(D \rightarrow f) + \Gamma(\bar{D} \rightarrow \bar{f})}$$

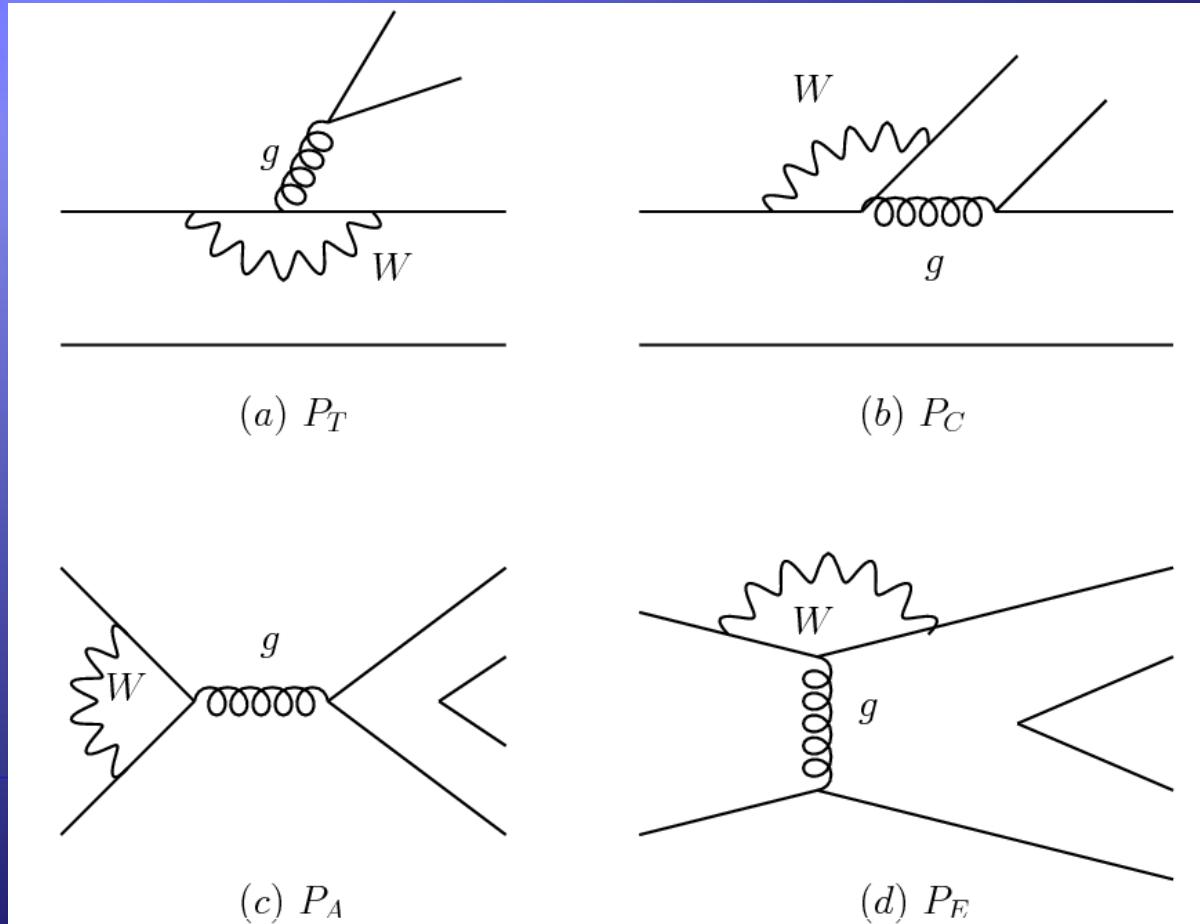
- ◆ Occurs only in singly Cabibbo-suppressed decays
- ◆ Interference of Tree and Penguin contributions

$$\mathcal{A}(D \rightarrow f) = V_{cd}^* V_{ud} T - V_{cb}^* V_{ub} P$$

$$V_{CKM} = \begin{pmatrix} 1 - \lambda^2/2 - \lambda^4/8 & \lambda & A\lambda^3(\bar{\rho} - i\bar{\eta}) + A\lambda^5(\bar{\rho} - i\bar{\eta})/2 \\ -\lambda + A^2\lambda^5[1 - 2(\bar{\rho} + i\bar{\eta})]/2 & 1 - \lambda^2/2 - \lambda^4(1 + 4A^2)/8 & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 + A\lambda^4[1 - 2(\bar{\rho} + i\bar{\eta})]/2 & 1 - A^2\lambda^4/2 \end{pmatrix}$$

Penguin topologies

- ◆ All topological penguin diagrams for $D \rightarrow PP$



Penguin parameterization

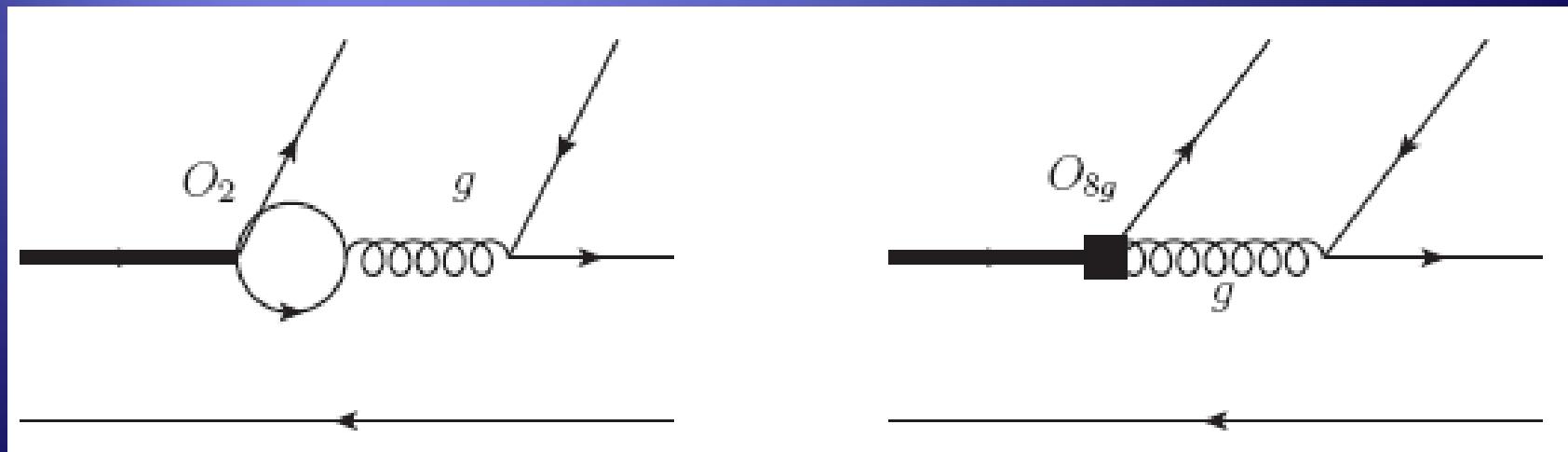
- ◆ Long-distance hadronic parameters, related to tree level, fixed by the data of branching ratios
- ◆ Combine the short-distance dynamics associated with penguin operators
- ◆ Then predict direct CP asymmetries

Strategy: relate penguin to tree

- ◆ At tree level, operators are all $(V-A)(V-A)$
- ◆ For penguins, the hadronic matrix elements with $(V-A)(V-A)$ operators are the same as tree level
- ◆ $(V-A)(V+A)$ can be related to tree by a sign, since either V or A contribute to $D \rightarrow PP$
- ◆ $(S-P)(S+P)$ are different and subtle
- ◆ Related to tree parameters by chiral enhancement, or neglected by power suppression

Quark loops & Magnetic penguin

- ◆ quark loops and magnetic penguin are absorbed into short-distance Wilson coefficients



Predictions

- ◆ Penguin matrix elements are either related to tree matrix elements (fixed by BRs), or factorizable and calculable, or power suppressed, if not able to be related
- ◆ formulate penguin contribution without introducing additional free parameters
- ◆ **Unambiguous** predictions of direct CP asymmetries

Predictions of Direct CP asymmetries

Modes	$a_{\text{CP}}(\text{FSI})$	$a_{\text{CP}}(\text{diagram})$	$a_{\text{CP}}^{\text{tree}}$	$a_{\text{CP}}^{\text{tot}} (\times 10^{-3})$
$D^0 \rightarrow \pi^+ \pi^-$	0.02 ± 0.01	0.86	0	0.74 ↙
$D^0 \rightarrow K^+ K^-$	0.13 ± 0.8	-0.48	0	-0.54 ↙
$D^0 \rightarrow \pi^0 \pi^0$	-0.54 ± 0.31	0.85	0	0.26
$D^0 \rightarrow K^0 \bar{K}^0$	-0.28 ± 0.16	0	0.69	0.90
$D^0 \rightarrow \pi^0 \eta$	1.43 ± 0.83	-0.16	-0.29	-0.61
$D^0 \rightarrow \pi^0 \eta'$	-0.98 ± 0.47	-0.01	0.43	1.67
$D^0 \rightarrow \eta \eta$	0.50 ± 0.29	-0.71	0.29	0.18
$D^0 \rightarrow \eta \eta'$	0.28 ± 0.16	0.25	-0.30	0.97
$D^+ \rightarrow \pi^+ \pi^0$		0	0	-0.23
$D^+ \rightarrow K^+ \bar{K}^0$	-0.51 ± 0.30	-0.38	-0.08	-0.93
$D^+ \rightarrow \pi^+ \eta$		-0.65	-0.46	0.63
$D^+ \rightarrow \pi^+ \eta'$		0.41	0.30	1.28
$D_S^+ \rightarrow \pi^0 K^+$		0.88	0.17	0.76
$D_S^+ \rightarrow \pi^+ K^0$		0.52	-0.01	0.87
$D_S^+ \rightarrow K^+ \eta$		-0.19	0.75	0.76
$D_S^+ \rightarrow K^+ \eta'$		-0.41	-0.48	1.83

Difference of CP asymmetries

- ◆ The prediction in the SM

$$\Delta A_{CP}^{SM} \equiv A_{CP}(K^+K^-) - A_{CP}(\pi^+\pi^-) = -0.13\%$$

- ◆ Enhanced from naively expectation in SM 10^{-4}
- ◆ LHCb $\Delta A_{CP} = [-0.82 \pm 0.24]\%$
- ◆ CDF $\Delta A_{CP} = [-0.62 \pm 0.23]\%$
- ◆ The prediction is still smaller than experimental measurements
- ◆ If CPV remains the current central value ($\sim 1\%$), may be a signal of new physics

Summary

- ◆ Explain branching ratios at tree level
- ◆ Parameterize penguin contributions, related to tree amplitudes under factorization hypothesis
 - ◆ Fix hadronic parameters using data of branching ratios
 - ◆ Combine short-distance dynamics associated with penguin operators
- ◆ Unambiguous predictions of direct CP asymmetries in $D \rightarrow PP$ in the SM
- ◆ Especially $\Delta A_{CP}^{SM} = -0.13\%$
- ◆ Much smaller than LHCb and CDF measurements

THANK YOU!

Improvement of BRs involving η'

- ◆ Benefited from the scale-dependent Wilson coefficients, predictions with all the η' involved modes are improved, compared to the pole model and diagrammatic approach

Modes	Br(FSI)	Br(diagrammatic)	Br(pole)	Br(exp)	Br(this work)
$D^0 \rightarrow \bar{K}^0 \eta'$	1.51	1.91 ± 0.09	1.9 ± 0.3	1.90 ± 0.11	1.84
$D_S^+ \rightarrow \pi^+ \eta'$	5.89	3.82 ± 0.36	4.6 ± 0.6	3.95 ± 0.34	3.25
$D^0 \rightarrow \pi^0 \eta'$	1.7	0.74 ± 0.02	0.6 ± 0.2	0.91 ± 0.13	0.70
$D^0 \rightarrow \eta \eta'$	2.2	1.19 ± 0.07	1.1 ± 0.1	1.05 ± 0.26	1.58
$D^+ \rightarrow \pi^+ \eta'$	7.9	3.70 ± 0.37	5.5 ± 0.8	4.68 ± 0.29	4.78 
$D_S^+ \rightarrow K^+ \eta'$	5.2	1.07 ± 0.17	1.4 ± 0.4	1.8 ± 0.5	1.67
$D^+ \rightarrow K^+ \eta'$		0.91 ± 0.17	1.0 ± 0.1	1.76 ± 0.22	1.07

Penguin operators

$$O_3 = \sum_{q'=u,d,s} (\bar{u}_\alpha c_\alpha)_{V-A} (\bar{q}'_\beta q'_\beta)_{V-A}, \quad (\textbf{V-A})(\textbf{V-A})$$

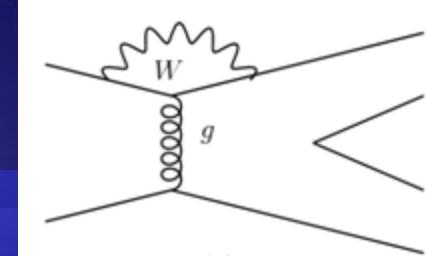
$$O_4 = \sum_{q'=u,d,s} (\bar{u}_\alpha c_\beta)_{V-A} (\bar{q}'_\beta q'_\alpha)_{V-A},$$

$$O_5 = \sum_{q'=u,d,s} (\bar{u}_\alpha c_\alpha)_{V-A} (\bar{q}'_\beta q'_\beta)_{V+A}, \quad (\textbf{V-A})(\textbf{V+A})$$

$$O_6 = \sum_{q'=u,d,s} (\bar{u}_\alpha c_\beta)_{V-A} (\bar{q}'_\beta q'_\alpha)_{V+A}$$

(S+P)(S-P)

Penguin exchange annihilation amplitude



- ◆ The hadronic matrix elements with the operators of $(S-P)(S+P)$ are absent in tree level
- ◆ No helicity suppression for $(S-P)(S+P)$, so their contributions can not be neglected
- ◆ Can not use the hadronic parameters at tree level
- ◆ Assumed to be dominated by scalar resonances