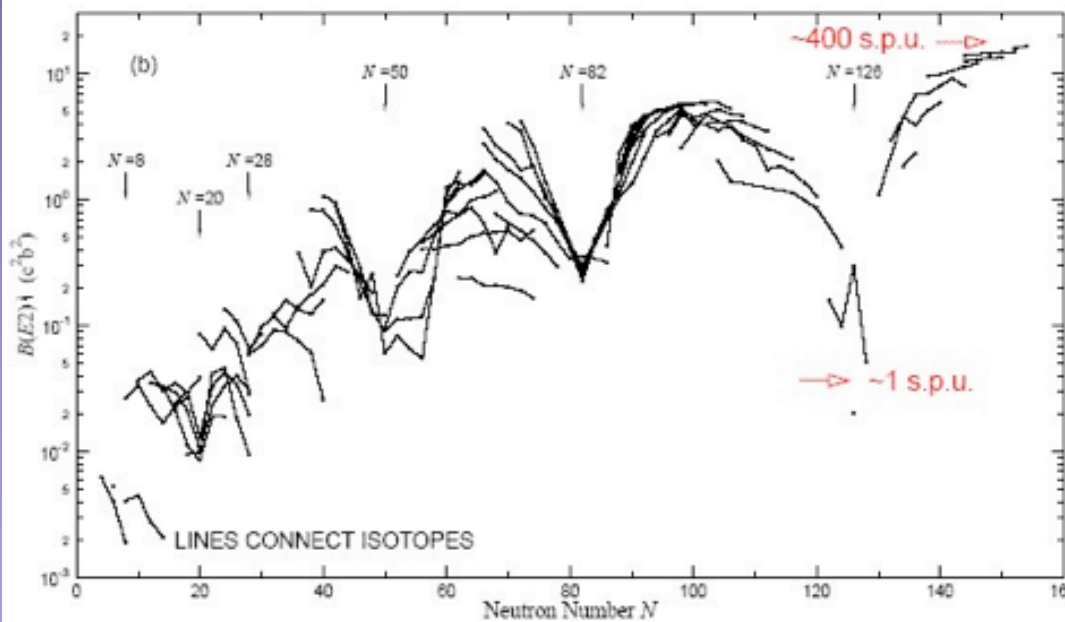
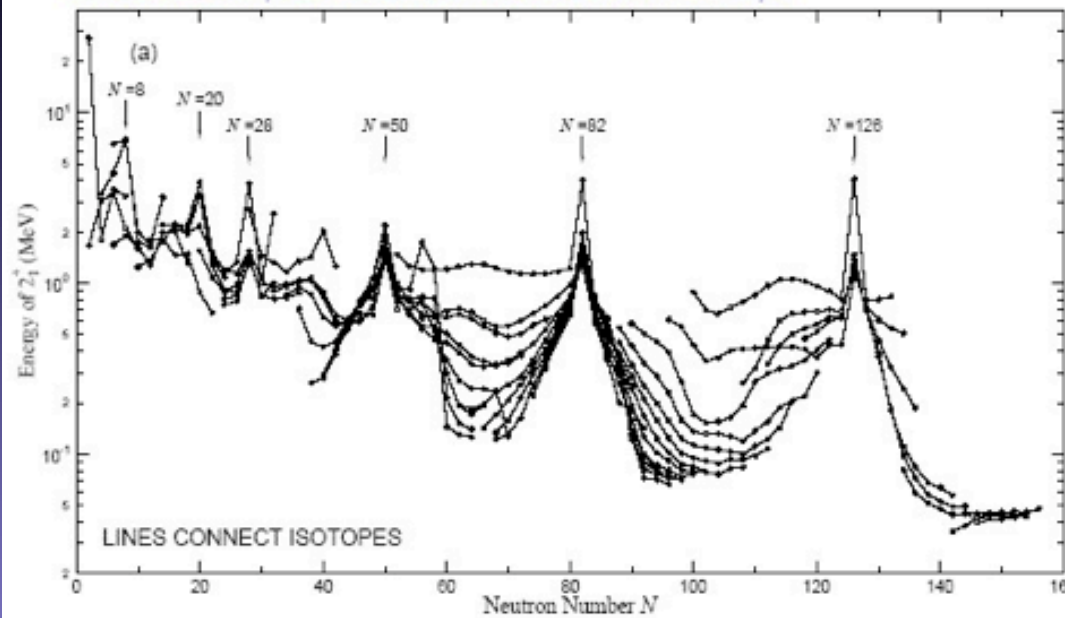


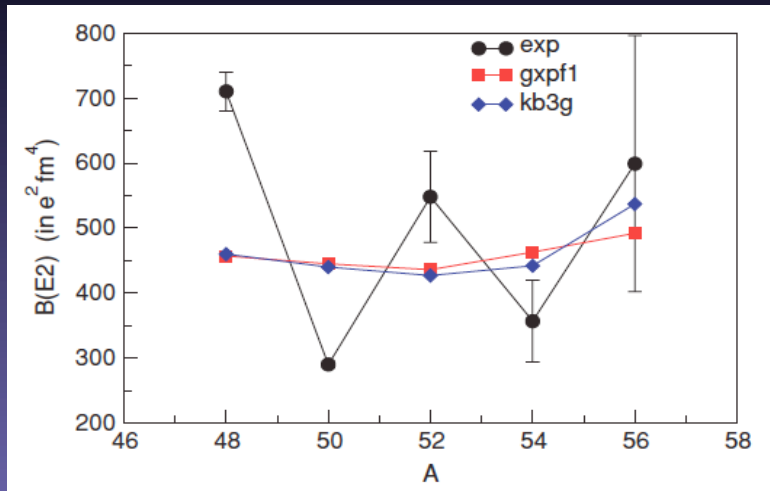
# Magnetic strength and shell evolution in light nuclei

Dennis M $\ddot{u}$ cher  
Physics Department E12  
TU M $\ddot{u}$ nchen



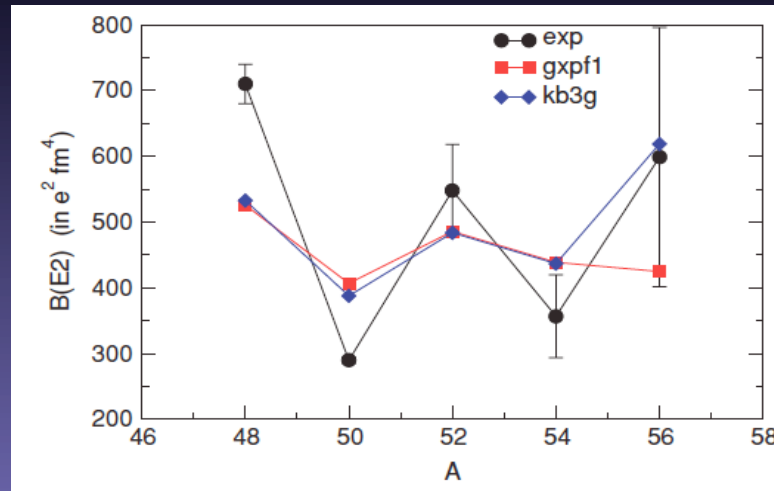


# B(E2) values off n-rich Ti-isotopes: are neutrons closed at N=34 ?



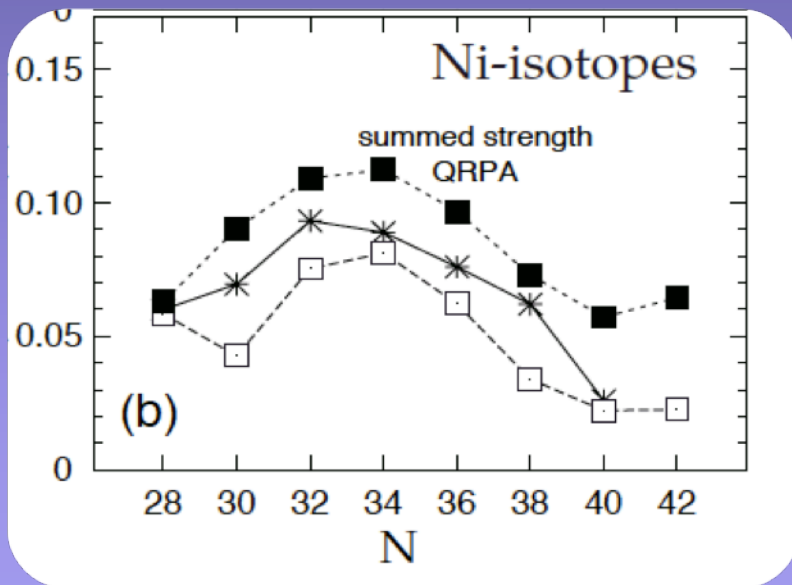
$Q_{\pi}=1.5e, Q_v=0.5e$

A. Poves, F. Nowacki, E. Caurier PRC 72, 047302 (2005)

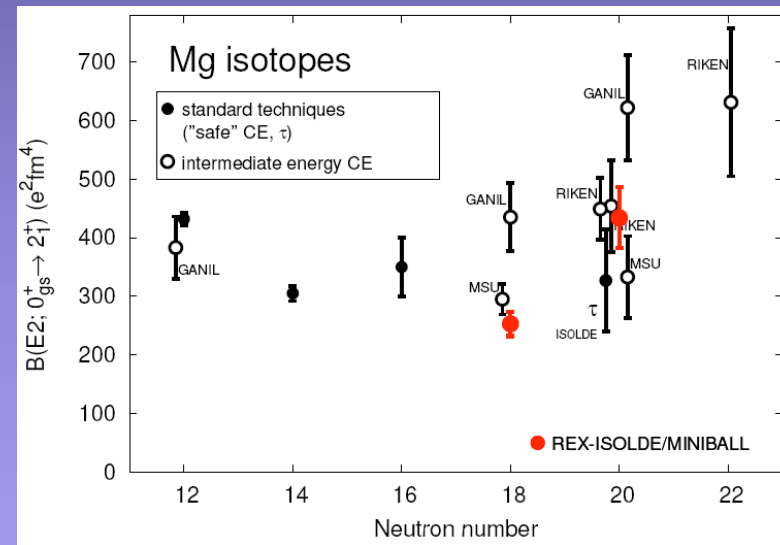


$Q_{\pi}=1.15e, Q_v=0.8e$

B(E2; 0<sup>+</sup> → 2<sup>+</sup>) (e<sup>2</sup>b<sup>2</sup>)



K. Langanke et al., Phys. Rev. C, 67, 044314 (2003)



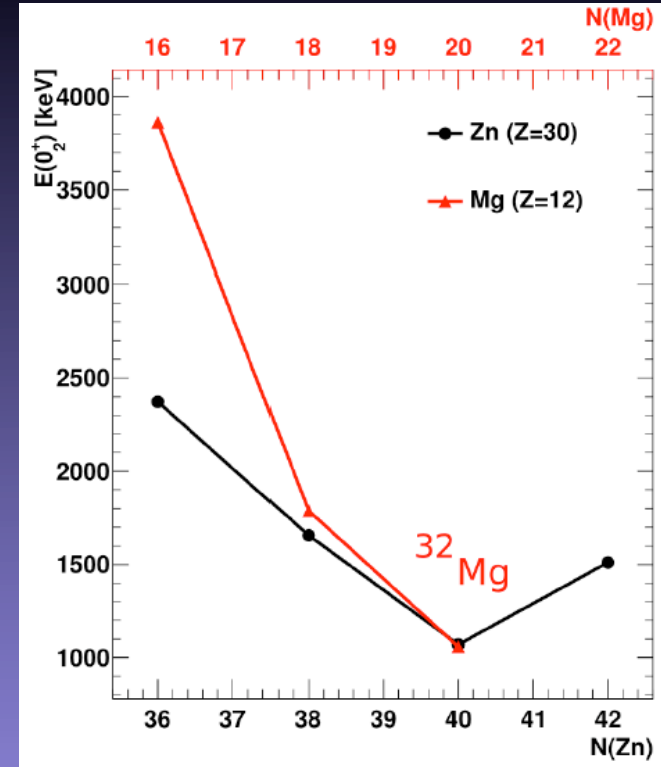
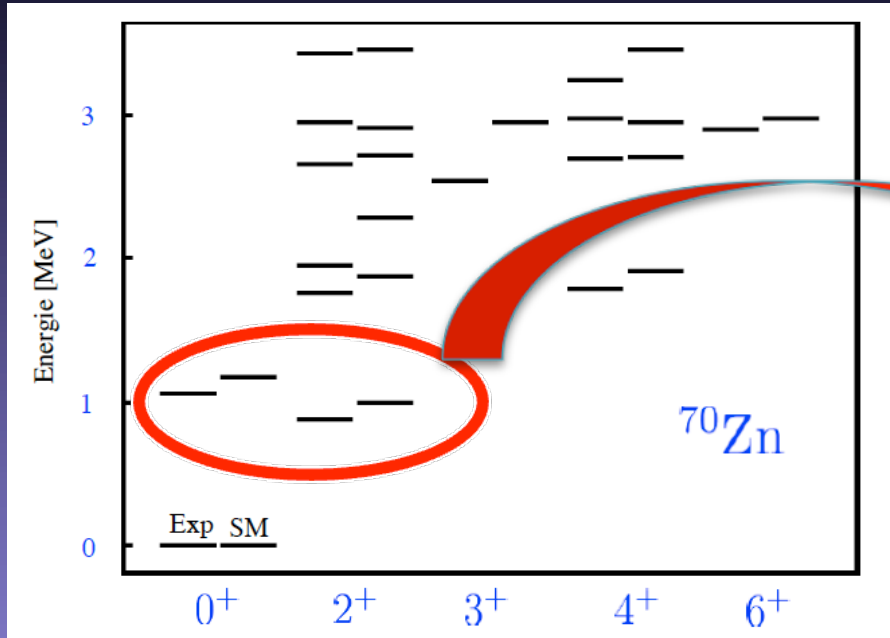
$$g_s = \alpha (fp)^2 - \beta (sd \text{ shell}), \quad 0_2^+ = \beta (fp)^2 + \alpha (sd)$$

K. Wimmer et al., Phys. Rev. Lett. 105, 2525 (2010)

H. T. Fortune, Phys. Rev. C 84, 034307 (2011)

# <sup>70</sup>Zn: yet another „Island of Inversion“

?



$f_{5/2}$   $p_{3/2}$   $p_{1/2}$   $g_{9/2}$  for  $p+n$   
 jj4c interaction (B. A. Brown)

$J^\pi$	A	$\pi(f_{5/2})$	$\pi(p_{3/2})$	$\nu(f_{5/2})$	$\nu(p_{3/2})$	$\nu(p_{1/2})$	$\nu(g_{9/2})$
$0_1^+$	0.146	0	2	4	4	2	2
	0.083	0	2	6	4	2	0
	0.065	2	0	4	4	0	4
	1.0	0.61	1.04	4.35	3.53	1.27	2.85
$0_2^+$	0.298	0	2	6	4	2	0
	0.076	2	0	4	4	0	4
	0.054	0	2	5	3	2	2
	1.0	0.54	1.17	4.95	3.54	1.44	2.07

if „closed“ configuration on top of  $0_2^+$ :  $g(2_1^+)$  large !

shell model:

$g(2_1^+) = 0.276$ , configuration „6 4 2 0“  $< 10^{-10}$

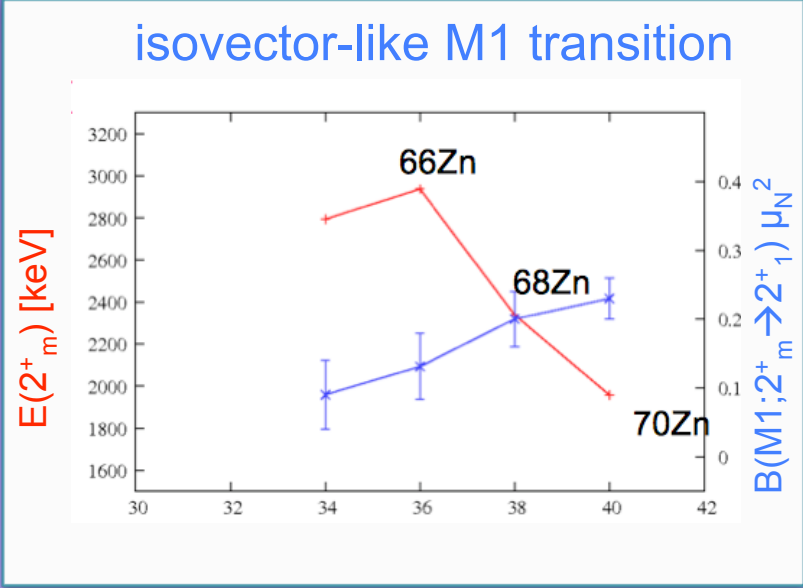
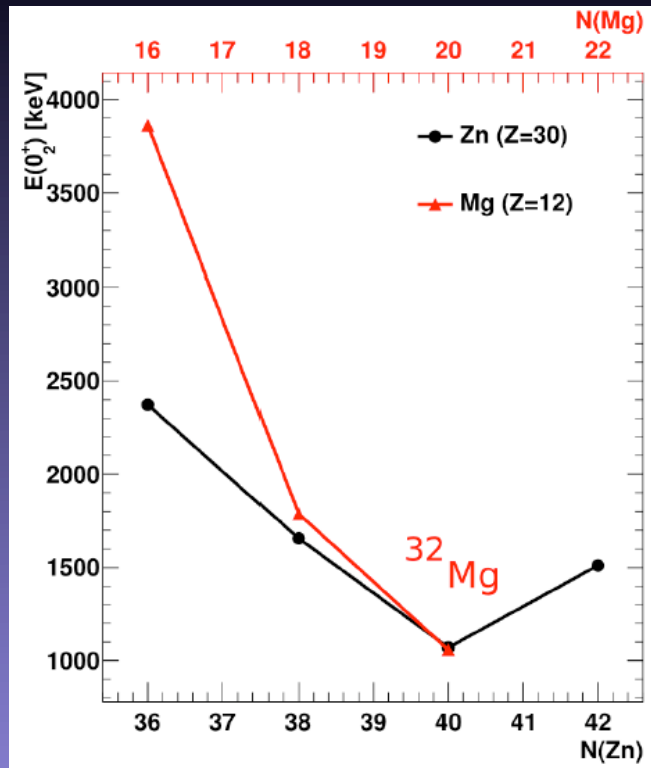
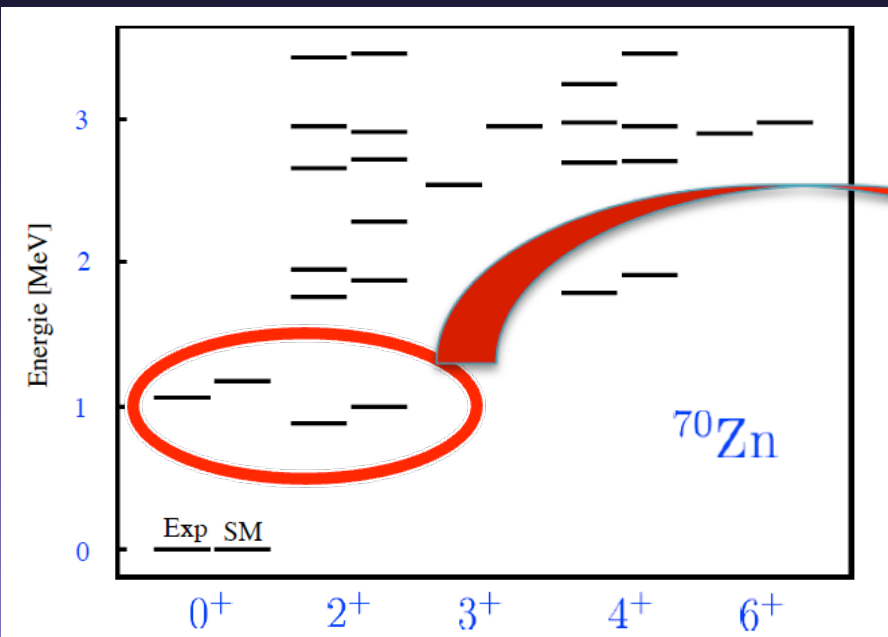
$g(2_2^+) = 0.10$ , configuration „6 4 2 0“  $< 10^{-10}$

$g(2_3^+) = 0.89$ , configuration „6 4 2 0“ largest!

(18)

$I_i^\pi$	Exp't.	FPD6	KB3	GXPFA	JJ4B
		<i>fp</i>	<i>fp</i>	<i>fp</i>	$p_{3/2}f_{5/2}p_{1/2}g_{9/2}$
$2_1^+$	+0.38(2) <sup>a</sup>	+1.52	+1.83	+1.89	+0.276

# <sup>70</sup>Zn: yet another „Island of Inversion“ ?



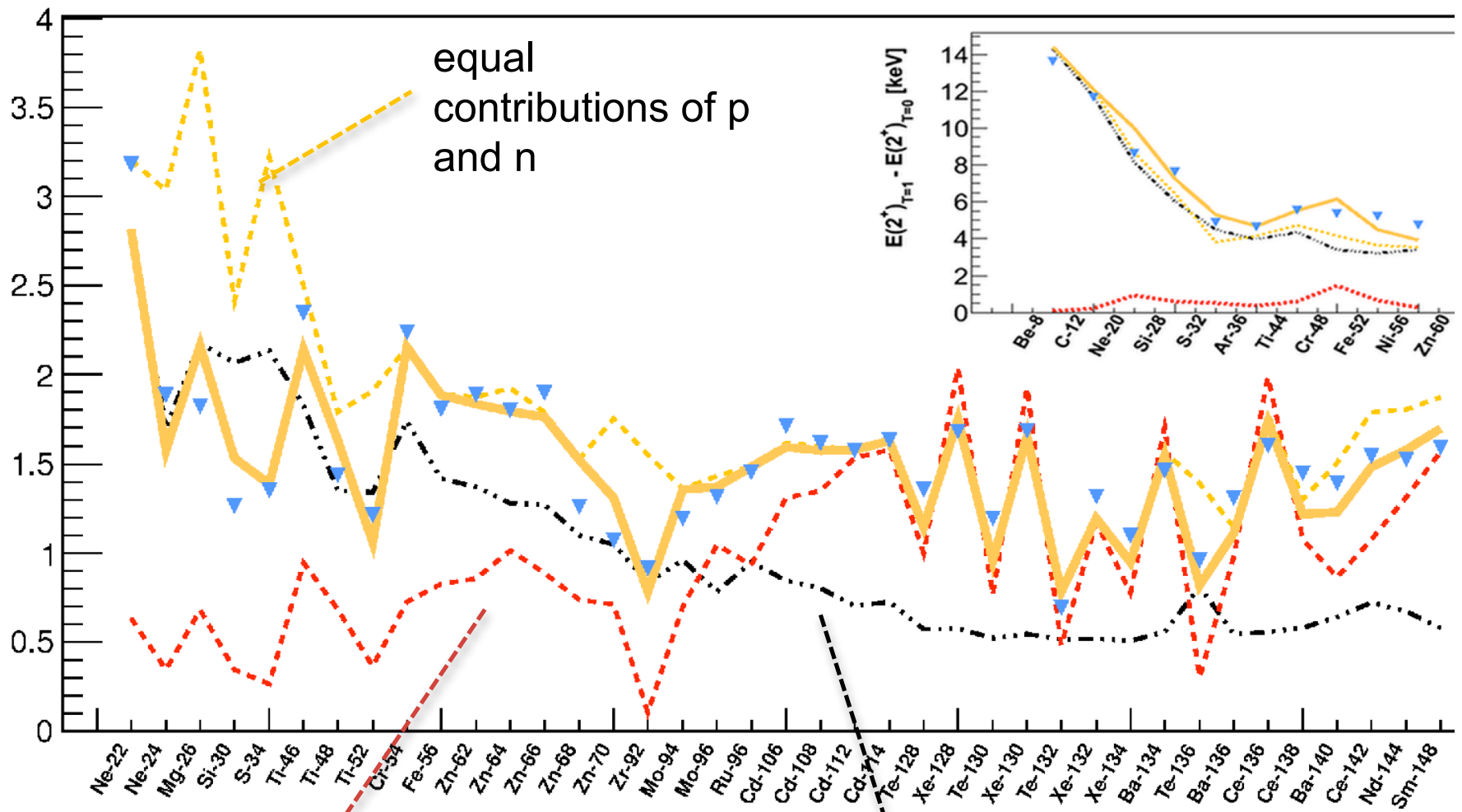
if „closed“ configuration on top of  $0^+_2$ :  $g(2^+_1)$  large !  
 shell model:  
 $g(2^+_1) = 0.276$ , configuration „6 4 2 0“  $< 10^{-10}$   
 $g(2^+_2) = 0.10$ , configuration „6 4 2 0“  $< 10^{-10}$   
 $g(2^+_3) = 0.89$ , configuration „6 4 2 0“ largest!

$I_i^\pi$	Exp't.	FPD6	KB3	GXPFA	JJ4B
$2^+_1$	+0.38(2) <sup>a</sup>	+1.52	+1.83	+1.89	+0.276

*fp*      *fp*      *fp*       $p_{3/2} f_{5/2} p_{1/2} g_{9/2}$

shell model:  $B(M1; 2^+_m \rightarrow 2^+_{-1}) = 0.18 \mu_N^2$

$E(2^+_{ms}) - E(2^+_{1})$



equal  
contributions of p  
and n

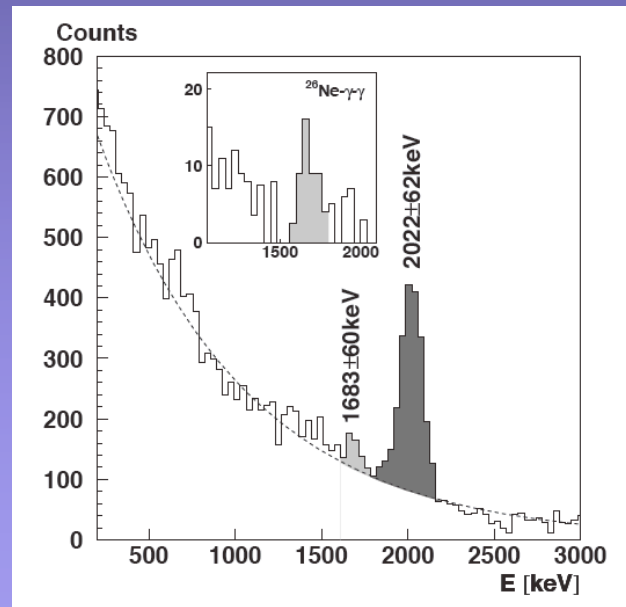
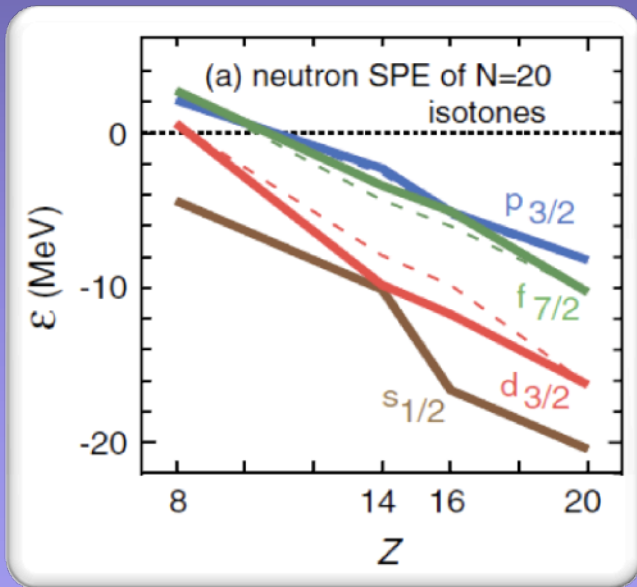
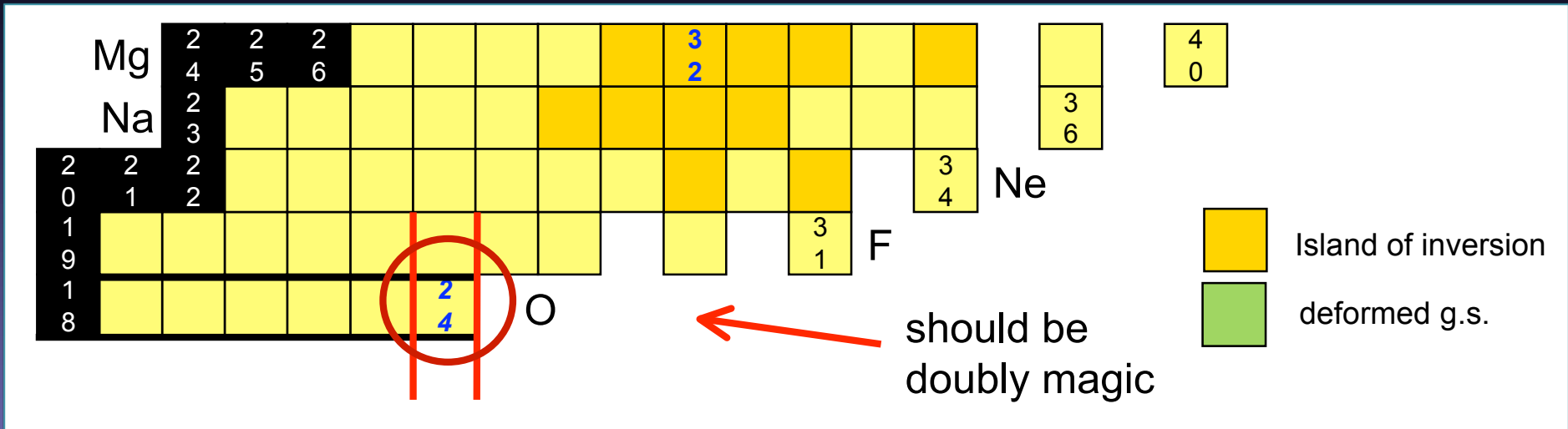
$E(2^+_{T=1}) - E(2^+_{T=0})$  [keV]

$$Q_{\pi\nu} = -\tilde{\kappa} \cdot \sqrt{\langle 0^+_{I'} | E2 | 2^+_{I'} \rangle^2 - \langle 0^+_{I'} | E2 | 2^+_{II'} \rangle^2}$$

Symmetry-energy

$$\delta V'_{pn} = \frac{4 \sum_J (2J+1) \langle j_{\pi} j_{\nu}, J | V | j_{\pi} j_{\nu}, J \rangle}{(2j_{\pi} + 1)(2j_{\nu} + 1)}$$

# Mixed Symmetry States at the dripline



$^{26}\text{Ne}$ :  
 $E(2^+_1) = 2.02 \text{ MeV}$   
 $E(2^+_2) = 3.69 \text{ MeV}$   
 $B(E2) = 6.18(8) \text{ W.u.}$

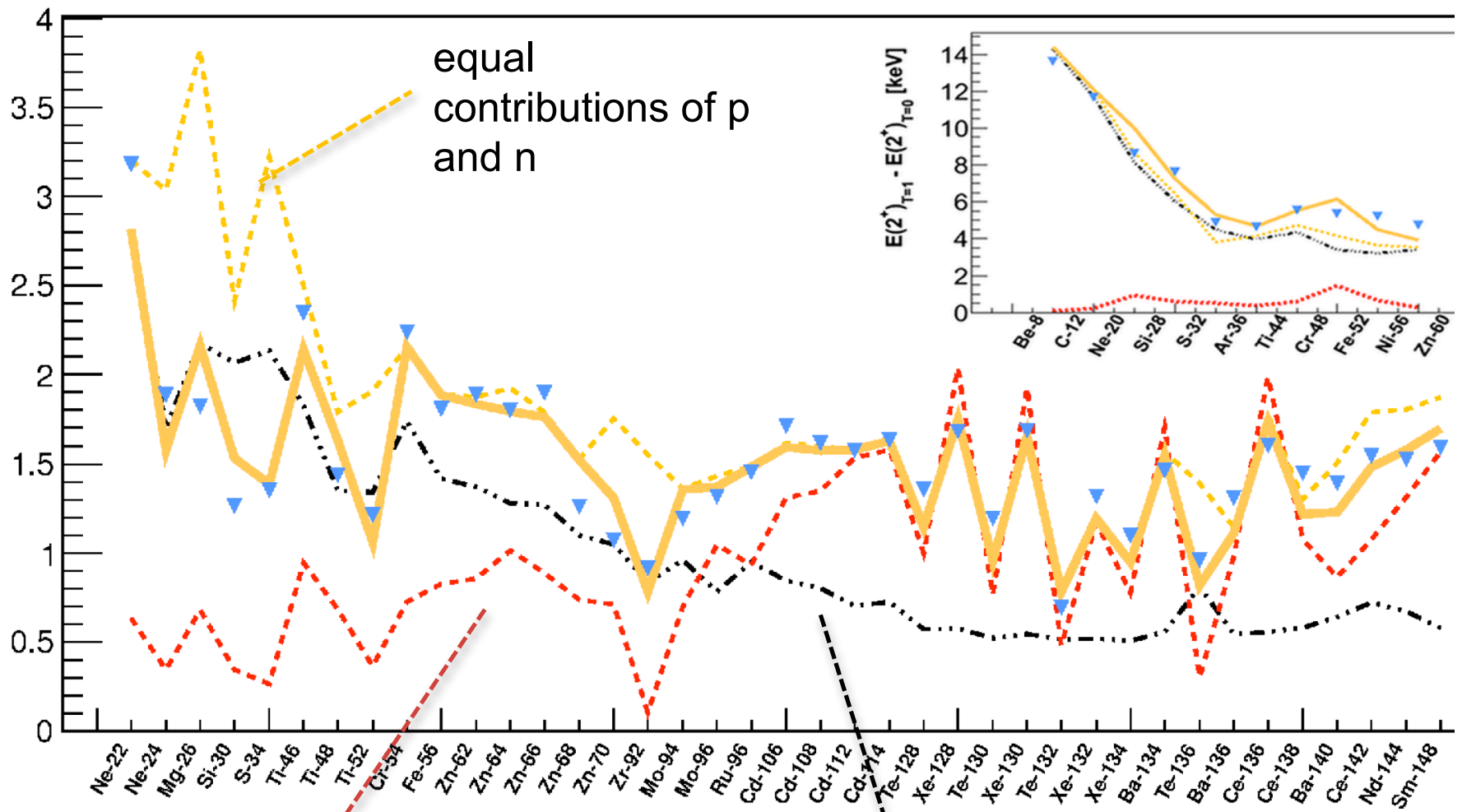
we obtain single  
 minimum with  
 $|\epsilon_{\pi} - \epsilon_{\nu}| = 2.5 \text{ MeV}$

$^{24}\text{O}$ :  $E(2^+_1) = 4.72 \text{ MeV}$   
 MeV (unbound)

N=16 ( $^{24}\text{O}$ ) is closed,  $^{28}\text{O}$  unbound!  
 → effect on semi-magic  $^{26}\text{Ne}$ ?

$^{26}\text{Ne}$ , Coulomb Excitation @ RIKEN  
 J. Gibelin et al., PRC 75, 057306 (2007)  
 shell model:  $2^+_2$  has isovector character

$E(2^+_{ms}) - E(2^+_{1})$



$$Q_{\pi\nu} = -\tilde{\kappa} \cdot \sqrt{\langle 0^+_{I'} | E2 | 2^+_{I'} \rangle^2 - \langle 0^+_{I'} | E2 | 2^+_{II'} \rangle^2}$$

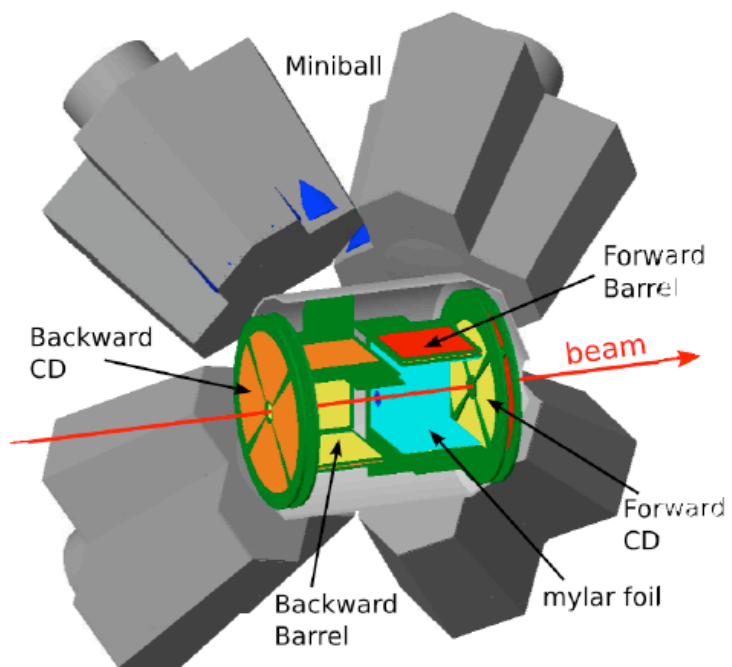
Symmetry-energy !

$$\delta V'_{pn} = \frac{4 \sum_J (2J+1) \langle j_{\pi} j_{\nu}, J | V | j_{\pi} j_{\nu}, J \rangle}{(2j_{\pi} + 1)(2j_{\nu} + 1)}$$



## Experimental setup: T-REX + MINIBALL

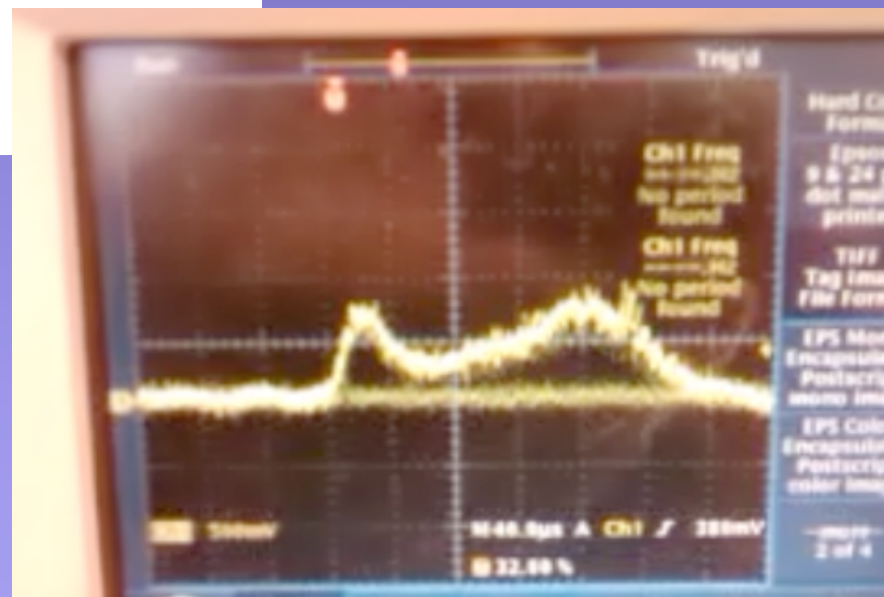
- ▶ Target:  $40 \mu\text{g}/\text{cm}^2$   $^3\text{H}$  (2n- and 1n-transfer) contained in a  $500 \mu\text{g}/\text{cm}^2$  Ti-foil (Coulex)
- ▶ Fully equipped T-REX allows to combine Coulex and transfer experiments.



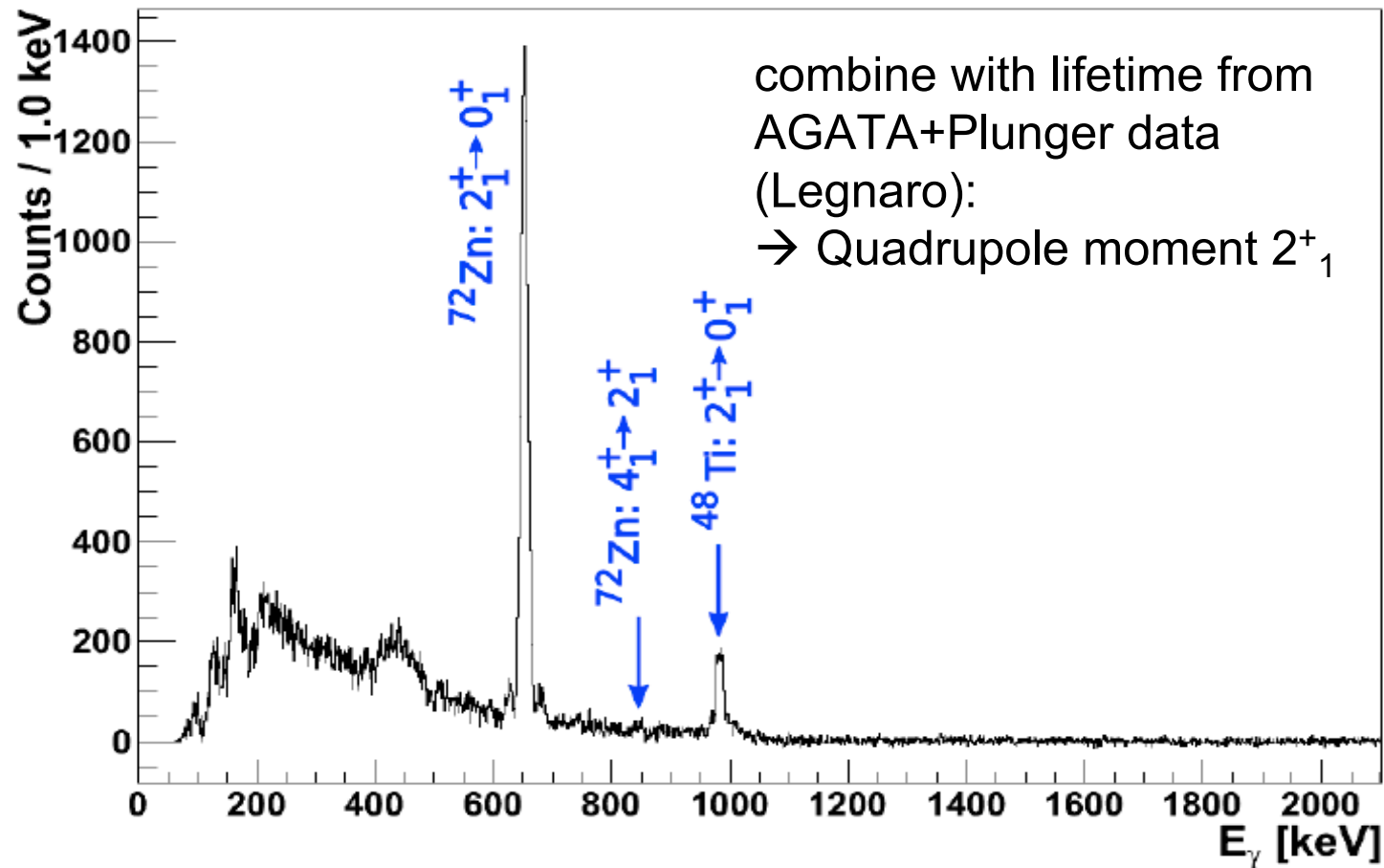
- ▶ Segmented  $\Delta E - E$  Si-telescopes for particle identification (p, d, ...)
- ▶  $12 \mu\text{m}$  mylar protection foil in front of Forward Barrel
- ▶ Segmented Forward CD for Coulex
- ▶ MINIBALL for  $\gamma$ -rays

Analysis of IS510 by S. Klupp, E12, TU Munich

high-intensity ( $10^7$  /s on target)  $^{72}\text{Zn}$  beam:  
issue with „flash“ of secondary electrons in T-REX:  
solved using high-power tritium-target-ladder (+300



## Coulex: First results

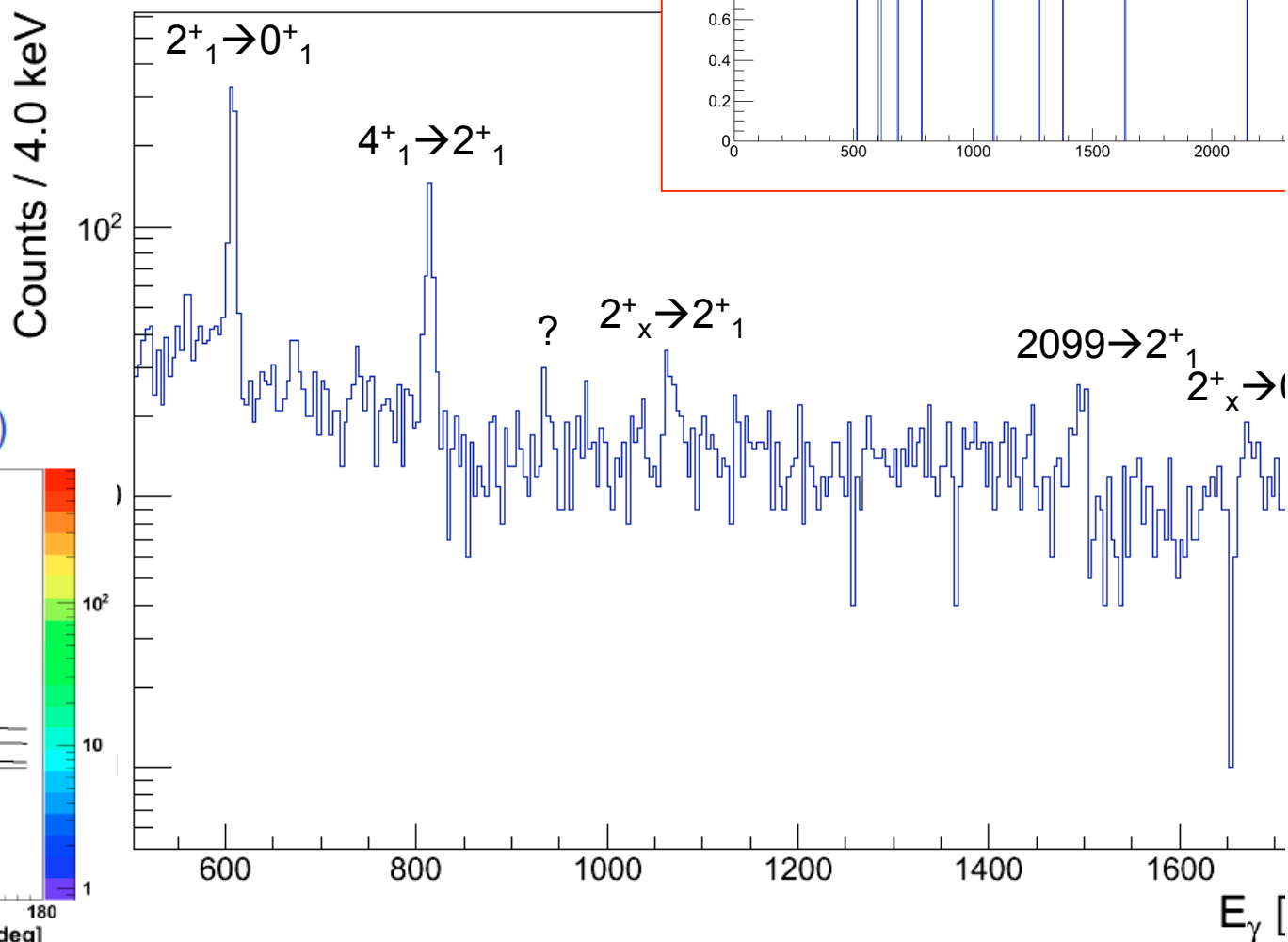
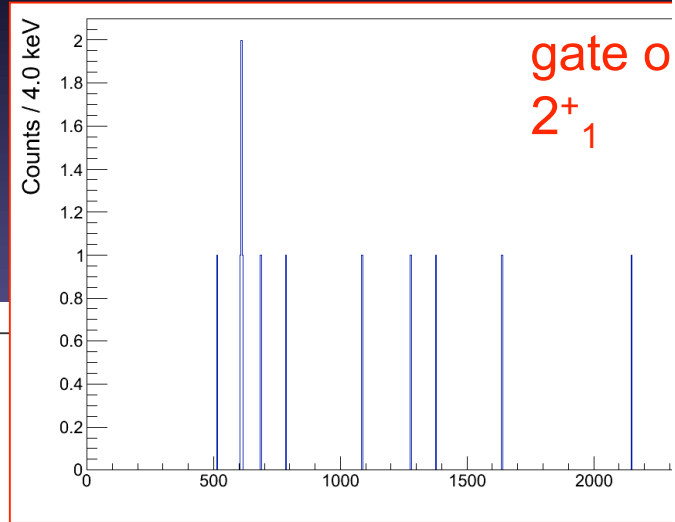


$1.95 \cdot 10^4$  counts in the  $^{72}\text{Zn}(2_1^+ \rightarrow 0_1^+)$  transition in 72h measurement time!

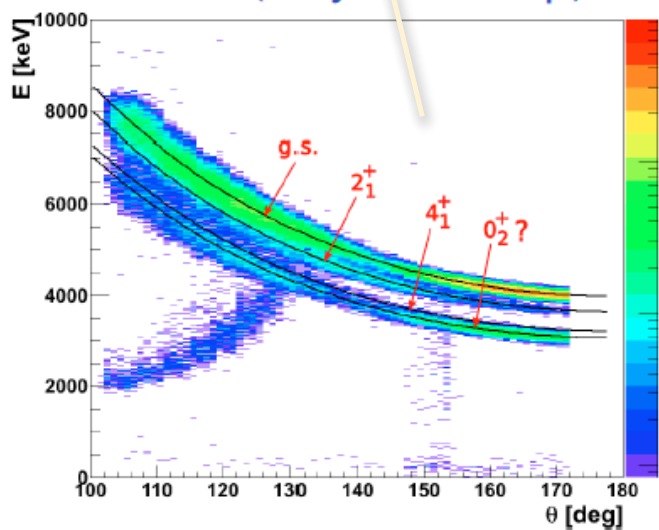
# MINIBALL-spectra of $^{74}\text{Zn}$ after 2n transfer, gated on protons

Analysis of IS510 by Stefanie Klupp, E12, TU Munich

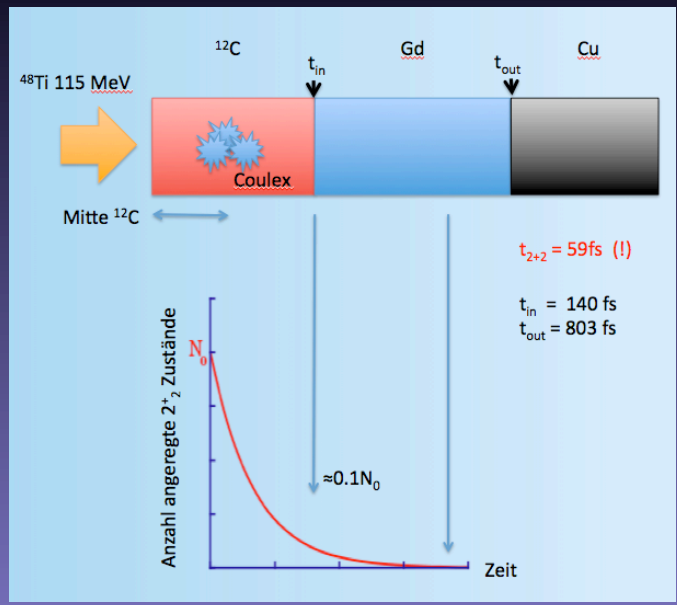
upgrade of T-REX needed for HIE ISOLDE



Simulation (only transfer p)



# Can we measure magnetic moments of short-lived $2^+_{ms}$ states?



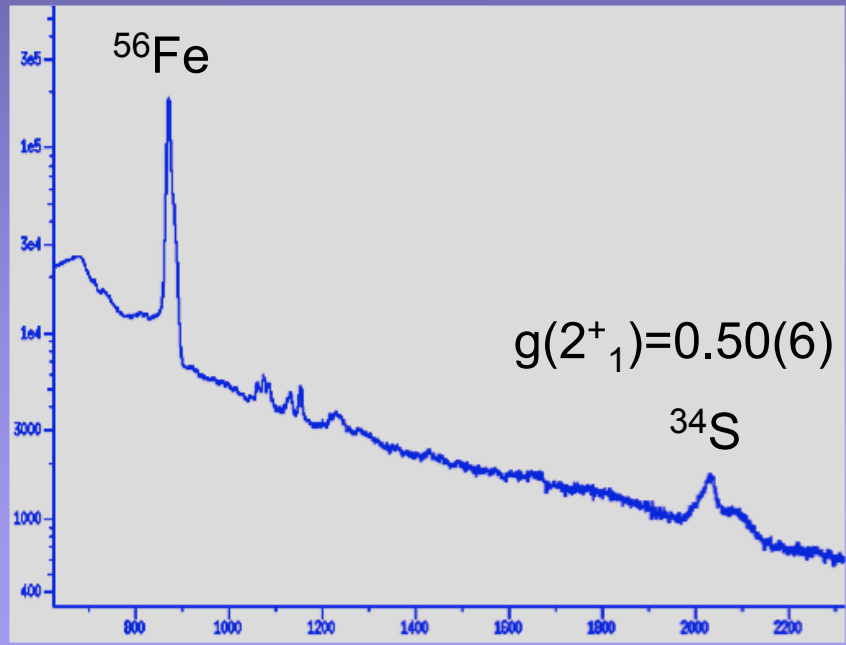
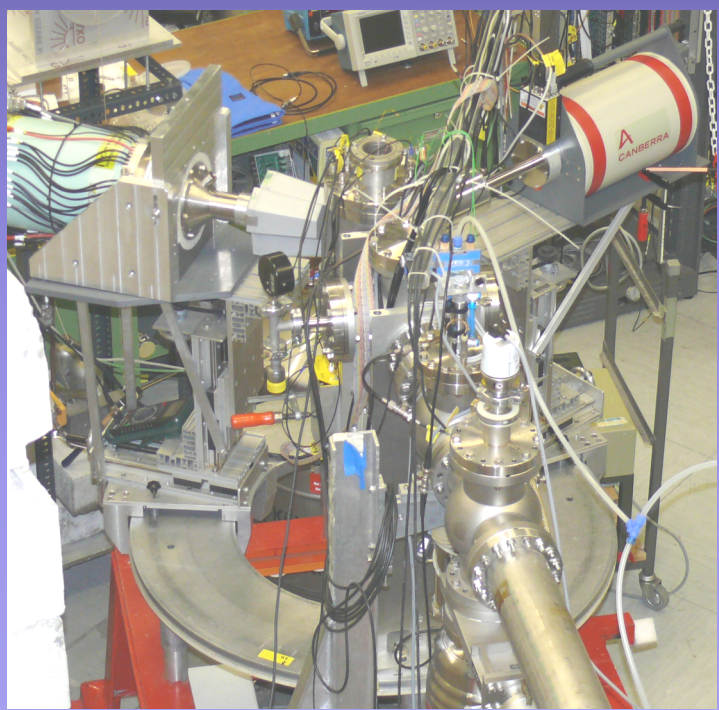
Scintillaor: LYSO (Cerium-doped Lutetium Yttrium Orthosilicate)

collab. Prof. S. Ziegler, medical physics, TU Munich

+ (i.e. glue)

Avalanche Photodiode (Hamamatsu S8664)

no radiation damage after  $10^{10}$  events, 100 kHz rate)



## Proposal for MINIBALL @ MLL Munich (15 MeV Tandem)

topics:

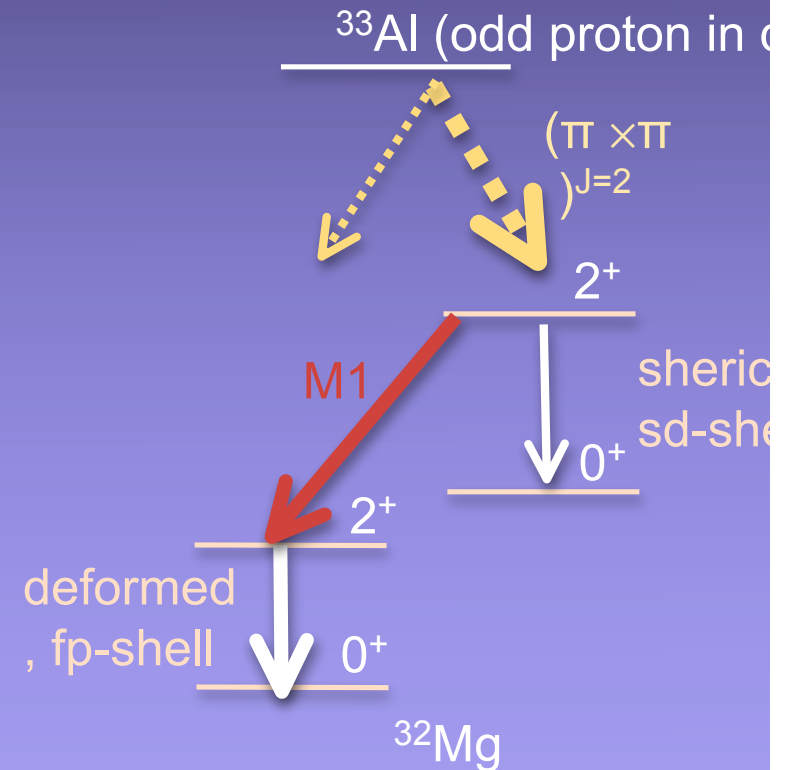
- normalisation measurements using Recoil in Vacuum after Coulex (in prep. for e.g. HIE-ISOLDE) (A. Jungclaus, Madrid)
- magnetic moments of ultra-fast states (DM)
- Lifetimes of astrophysical relevant states using cooled  $^3\text{He}$  targets (S. Bishop, E12 Munich)
- X-ray multiplicities of evaporation residues (W. Henning)

Setup:

- 4 MINIBALL triple cluster detector (72 Segments)
- trigger: all segments; Mesytec Shaper
- readout: Mesytec ADC; multievent readout

Working on a possible proposal for AGATA @ GSI:

identify  $2^+$  state build on  $0^+_2$  for nucleus in the „island of inversion)  
( $^{32}\text{Mg}$ , neutron-rich Fe)



# Thanks for your attention !

## E12 (TUM)

R. Krücken, W. Henning, R. Gernhäuser, K. Nowak, S. Klupp,  
H. Schmeiduch, S. Reichart, M. Bendel, L. Maier, C. Herlitzius,  
S. Bishop

## IS 510 (ISOLDE)

D. Mücher<sup>1</sup>, R. Krücken<sup>1</sup>, K. Wimmer<sup>1</sup>, V. Bildstein<sup>1</sup>, M. Albers<sup>2</sup>,  
L. Bettermann<sup>2</sup>, A. Blazhev<sup>2</sup>, S. Bönig<sup>3</sup>, J. Eberth<sup>2</sup>, C. Fransen<sup>2</sup>, R.  
Gernhäuser<sup>1</sup>, K. Gladnishki<sup>4</sup>, S. Das Gupta<sup>5</sup>, K. Hadynska<sup>7</sup>, M. Hass<sup>8</sup>,  
J. Iwanicki<sup>7</sup>, J. Jolie<sup>2</sup>, A. Jungclaus<sup>9</sup>, V. Kumar<sup>8</sup>, T. Kröll<sup>3</sup>, J. Leske<sup>3</sup>,  
G. Lo Bianco<sup>5</sup>, P. Napiorkowski<sup>7</sup>, B.S. Nara Singh<sup>6</sup>, K. Nowack<sup>1</sup>, R.  
Orlandi<sup>9</sup>, J. Pakarinen<sup>10</sup>, N. Pietralla<sup>3</sup>, G. Rainovski<sup>4</sup>, M. Scheck<sup>3</sup>, K.  
Singh<sup>8</sup>, J. Srebrny<sup>7</sup>, M. von Schmid<sup>3</sup>, K. Wrzosek-Lipska<sup>7</sup>, N. Warr<sup>2</sup>,  
M. Zielinska<sup>7</sup>, and the REX-ISOLDE collaboration

and

N. Pietralla, M. Scheck (TU Darmstadt)

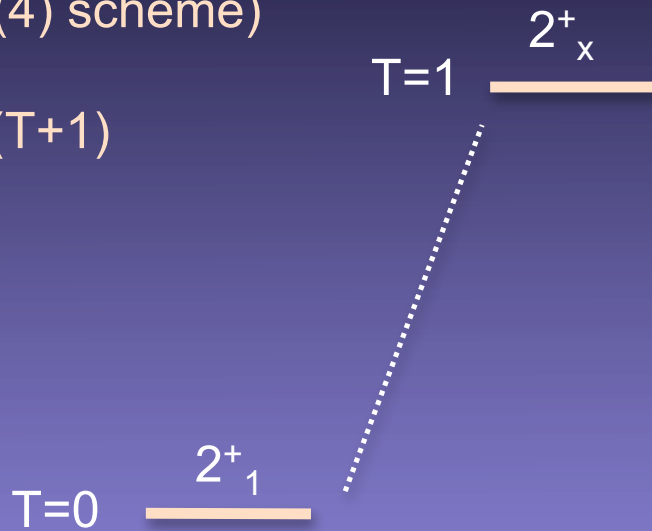
G. Rainowski (Sofia)

A. Jungclaus (Madrid)

# Maybe we can learn from the Isospin formalism ?

monopole Majorana exchange operator (Wigner, SU(4) scheme)

$$M = \sum_{i < j} P_{ij} \approx T(T+1)$$



even-even  $N=Z$  nucleus

Isospin: protons and neutrons behave the same.

algebra:  $SU_T(2)$

energy difference: (a)symmetry energy  $T(T+1)$  (+extra binding from Wigner energy)

origin: high  $T \leftrightarrow$  high permutation symmetry in charge  $\leftrightarrow$  low symmetry in space+spin

$$F \cdot F = N - N_\pi N_\nu + [(T \cdot T)^2 - T_0^4 - 2nT_0^2]$$

2 protons and 2 neutrons in same orbit

$F=0 \leftrightarrow T=1$

$F=1 \leftrightarrow T=0$  (+ 20%  $T=2$ )



even-even  $N \neq Z$  nucleus

F-Spin: proton-bosons and neutron bosons behave the same

algebra:  $SU_F(2)$

energy difference: (a)symmetry energy

$$E_s = K(Z - N)^2/A$$

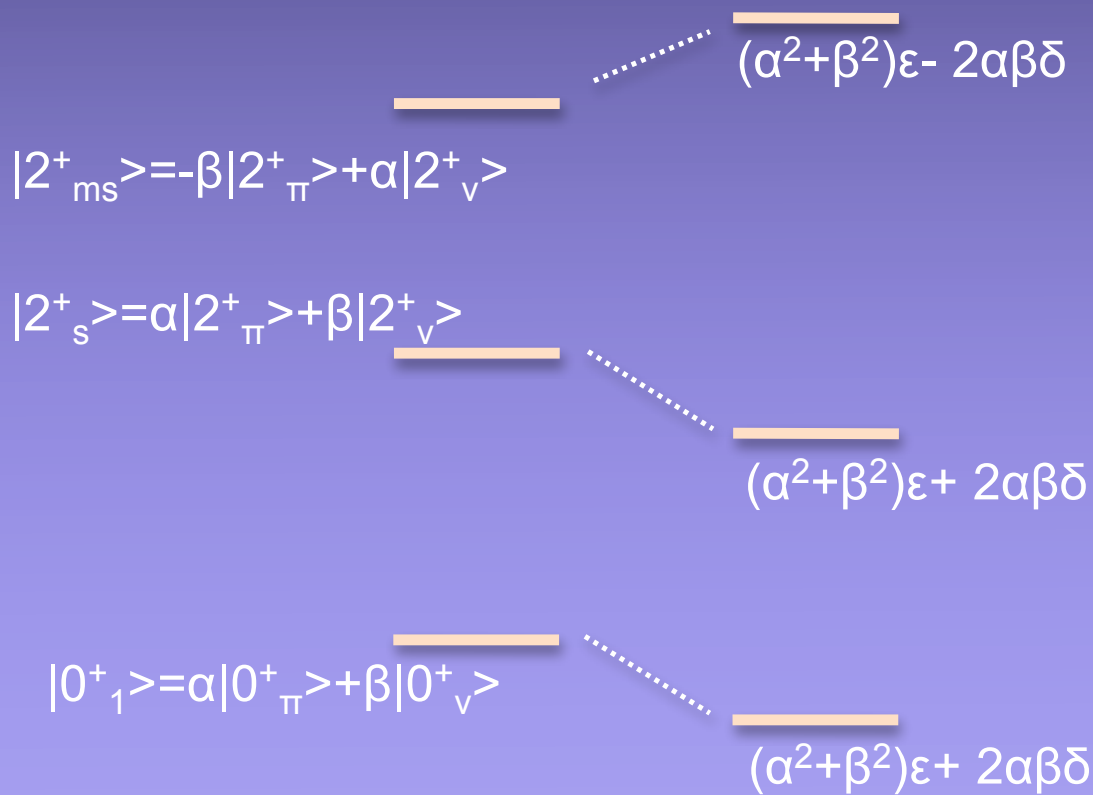
$$= \int [K(\rho_p - \rho_n)^2 / (\rho_p + \rho_n)] d\tau$$

$$V_{pn} = \sum_{j_p j_n j'_p j'_n J M} \langle j_p j_n | V_{pn} | j'_p j'_n \rangle_J A^\dagger(j_p j_n J M) A(j'_p j'_n J M).$$

$$f^{(0)}(j_p j_n, j_p j_n) = \frac{\sum_J (2J + 1) \langle j_p j_n | V_{pn} | j_p j_n \rangle_J}{\sqrt{(2j_p + 1)(2j_n + 1)}}$$

$$\langle J^+_\rho | M_{\rho\rho} | J^+_\rho \rangle = \text{const.}$$

$$\langle J^+_\rho | M_{\rho n} | J^+_{\rho'} \rangle = \text{const.}$$



shift due to monopole:  
 $E(2^+_{ms}) \rightarrow E(2^+_{ms}) + 4\alpha\beta\delta$

complete mixing:  
 $\alpha\beta = 1/2$

$$E(2^+_{ms}) \rightarrow E(2^+_{ms}) + 2\delta$$

$$\langle 0^+_{\pi} | M_{pn} | 0^+_{\nu} \rangle = \delta$$



K. Heyde, J. Sau, PRC 33, 3 (1986), p. 1050  
 seniority  $u=2$  shell-model states, single-j:

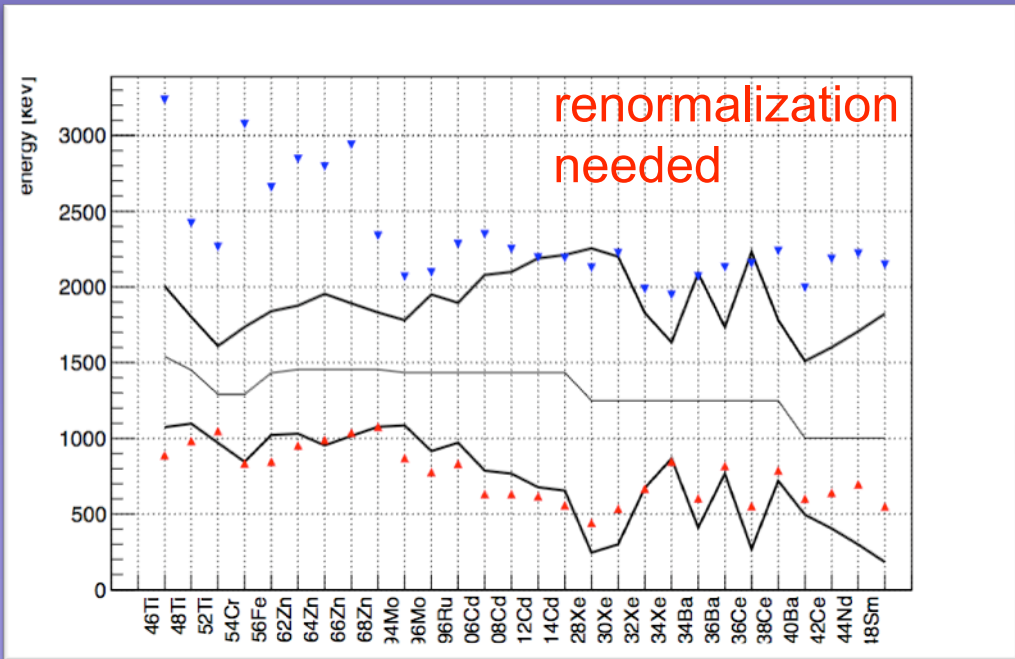
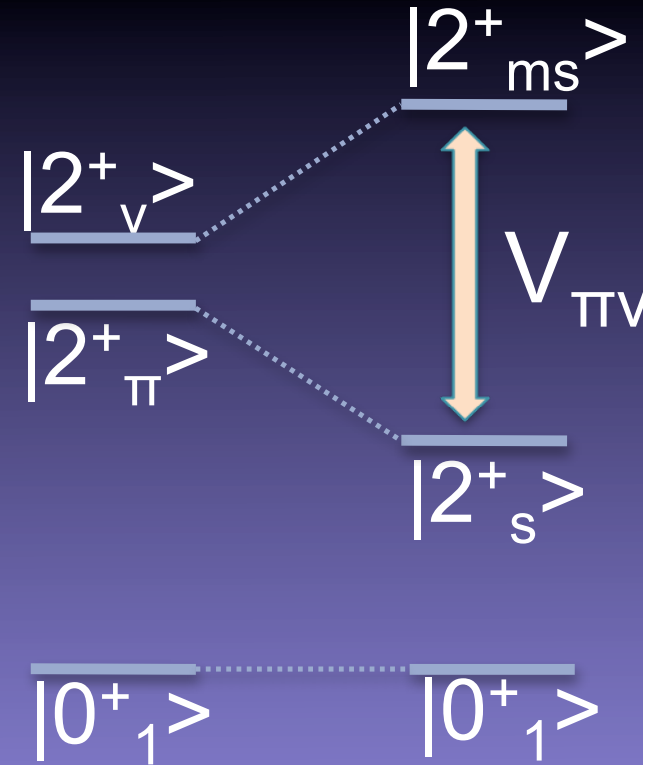
$$2_{\pi}^{+} = (j_{\pi})_{n_{\pi}=2}^{+}; 2^{+} (j)_{n=0}^{+}; 0^{+} 2^{+}$$

$$2^{+} = (j_{\pi})_{n_{\pi}=2}^{+}; 0^{+} (j)_{n=0}^{+}; 2^{+} 2^{+}$$

switch on interaction:

$$V_{\pi} = 2_{\pi}^{+} / - \frac{Q_{\pi} \cdot Q}{2^{+}}$$

$$E(2_{ms}^{+}) - E(2_1^{+}) = \frac{1}{4} (\pi - )^2 + Q_{\pi} Q$$



$$E_s = K(Z - N)^2/A$$

$$= \int [K(\rho_p - \rho_n)^2 / (\rho_p + \rho_n)] d\tau$$

A. Faessler et al, Phys. Lett 166B, 4 (1985)

K. Heyde, J. Sau, PRC 33, 3 (1986), p. 1050  
 seniority  $u=2$  shell-model states, single-j:

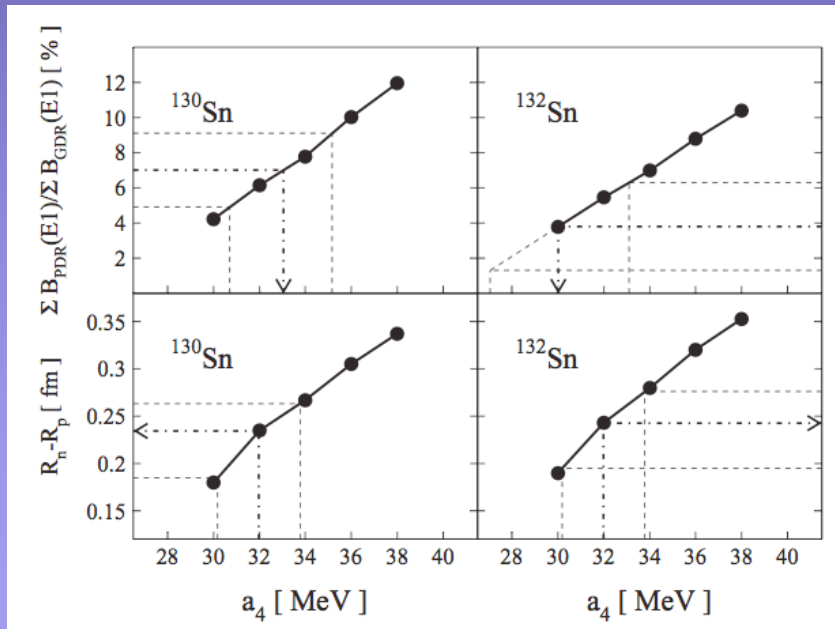
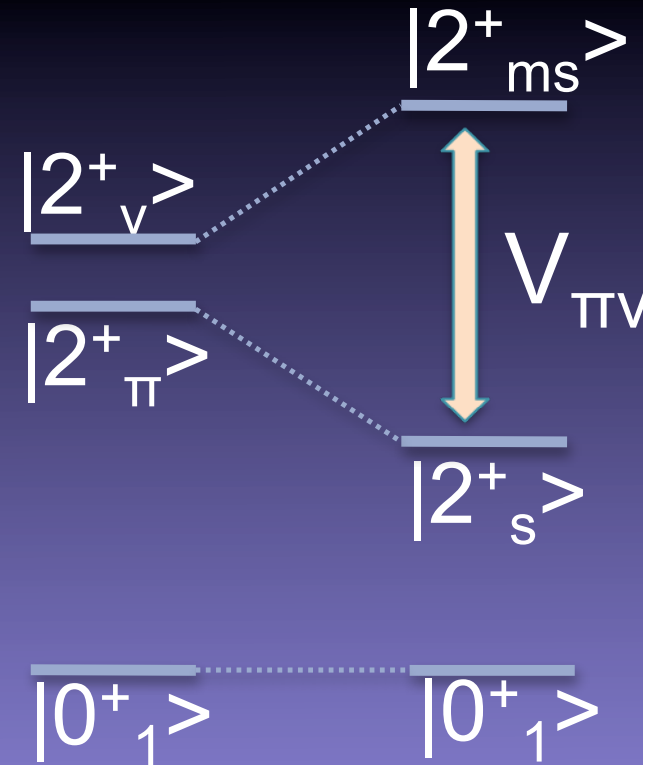
$$2_{\pi}^{+} = (j_{\pi})_{n_{\pi}=2}^{+}; 2^{+} (j)_{n=0}^{+}; 0^{+} 2^{+}$$

$$2^{+} = (j_{\pi})_{n_{\pi}=2}^{+}; 0^{+} (j)_{n=0}^{+}; 2^{+} 2^{+}$$

switch on interaction:

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$$E(2_{ms}^{+}) - E(2_1^{+}) = \frac{1}{4} (\pi - )^2 + Q_{\pi} Q$$



A. Klimkiewicz et al, PRC 76, 051603(R) 2007

$$E_s = K(Z - N)^2/A$$

$$= \int [K(\rho_p - \rho_n)^2/(\rho_p + \rho_n)] d^3r$$

A. Faessler et al, Phys. Lett 166B, 4 (1985)