NUCLEAR PARTON DISTRIBUTION FUNCTIONS

Karol Kovařík

Institut für Theoretische Physik, KIT, Karlsruhe

in collaboration with I.Schienbein, F.Olness, J.Owens, J.Morfin, C.Keppel, J.Y.Yu, T. Stavreva

Phys.Rev.D80 094004, 2009 Phys.Rev.Lett.106 122301 (2011)

OUTLINE

- 1. Nuclear effects in PDFs
- 2. Overview of global nPDF analysis
- 3. Generalized CTEQ framework
- 4. Neutrino DIS data
- 5. Outlook

NUCLEAR EFFECTS

What are nuclear parton density functions (nPDF) ?

- parton densities for partons in bound proton & neutron

Where are nuclear parton density functions useful ?

1. Strange quark content of the proton

strange PDF from neutrino DIS with heavy nuclei - nuclear effects very important

crucial for: W-boson production at the LHC (standard candle process)

NUCLEAR EFFECTS

What are nuclear parton density functions (nPDF) ?

- parton densities for partons in bound proton & neutron

Where are nuclear parton density functions useful ?

1. Strange quark content of the proton

strange PDF from neutrino DIS with heavy nuclei - nuclear effects very important

crucial for: determining weak mixing angle from NuTeV experiment

NUCLEAR EFFECTS

What are nuclear parton density functions (nPDF) ?

- parton densities for partons in bound proton & neutron

Where are nuclear parton density functions useful ?

1. Strange quark content of the proton

2. Collisions of protons and nuclei at RHIC, ALICE & CMS

- 3. Neutrino scattering experiments e.g. MINERvA
- 4. Neutrino oscillations experiments e.g. MINOS

Factorization & PDFs

Parton distribution functions (PDFs) $f_{A\to a}(x,\mu_F)$

- universal, non-perturbative objects - describe the structure of hadrons (in terms of partons - quarks & gluons) - obey DGLAP evolution equations

> The hard cross-section $\hat{\sigma}$ $ab \rightarrow c$

- free of long distance effects
- calculable in pQCD
- process dependant

 $\sigma = f_{A\rightarrow a}\otimes f_{B\rightarrow b}\otimes \hat{\sigma}$ $ab \rightarrow c$ from experiment from pQCD →→→

Collins, Soper, Sterman hep-ph/0409313

Universality of PDFs - same parton distribution functions for all processes

- Deep Inelastic Scattering (DIS)

$$
F_2^A(x,\mu^2) = \sum_i \left[f_i^A \otimes C_{2,i}\right](x,\mu^2)
$$

- Drell-Yan processes (DY)

$$
\sigma_{A+B\to l^++l^-+X} = \sum_{i,j} f_i^A \otimes f_j^B \otimes \hat{\sigma}^{i+j \to l^++l^-+X}
$$

- hadron production

$$
\sigma_{A+B\to H+X} = \sum_{i,j,k} f_i^A \otimes f_j^B \otimes \hat{\sigma}^{i+j\to k+X} \otimes D_k^H
$$

 PDFs give predictions for unexplored kinematic regions and for new physics at the LHC

CTEQ framework to fit PDFs from experimental data

CTEQ6M hep-ph/0201195

- the input scale set to $\mu_0 = Q_0 = 1.3 \, \text{GeV}$

- parameterization of the PDFs in \times

 $x f_k(x,Q_0) = c_0 x^{c_1} (1-x)^{c_2} e^{c_3 x} (1+e^{c_4} x)^{c_5}$ $k = u_v, d_v, g, \bar{u} + \bar{d}, s, \bar{s}$ $d(x, Q_0)/\bar{u}(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} + (1+c_3x)(1-x)^{c_4}$

- make sure # of free parameters not too high - CTEQ approx. 20 free params

- carefully choose data sets & kinematic cuts to constrain free parameters

7

- perform χ^2 fit to data

Which data sets are included ?

- Deep Inelastic Scattering $(l^{\pm}p, l^{-}d, \nu N, \bar{\nu}N)$ $p \longrightarrow \bigcup_{\mathcal{F}} Y$
- Neutrino DIS di-muon production
- Drell-Yan & vector boson production $(W^\pm, Z^0, \gamma$)
- hadronic jet data

Review of existing global analyses of nuclear PDF

- first differentiating factor how to relate nuclear PDF to proton PDF
	- 1. Multiplicative nuclear correction factor

 $f_i^A(x_N, Q_0^2) = R_i(x_N, Q_0, A, Z) f_i(x_N, Q_0^2)$ λ bound parton density free parton density

Hirai, Kumano, Nagai [PRC76(2007)065207] arXiv: 0709.0338 Eskola, Paukkunen, Salgado [JHEP0904(2009)065] arXiv: 0902.4154

2. Convolution relation

$$
f_i^A(x_N, Q_0^2) = \int_{x_N}^A \frac{dy}{y} W_i(y, A, Z) f_i(x_N/y, Q_0^2)
$$

nucleon density in nucleus with y/A mom. fraction

de Florian, Sassot [PRD69(2004)074028] hep-ph/0311227

- second differentiating factor - data sets included in the analysis

Review of existing global analyses of nuclear PDF

DE FLORIAN, SASSOT'04 [PRD69(2004)074028] LO, NLO

$$
\chi^2/\text{dof} = 0.76
$$

- first NLO analysis of nuclear data
- the only group using convolution relation

$$
f_i^A(x_N, Q_0^2) = \int_{x_N}^A \frac{dy}{y} W_i(y, A, Z) f_i(x_N/y, Q_0^2)
$$

- typical nucleon density for valence quarks

$$
W_v(y, A, Z) = A[a_v \delta(1 - \epsilon_v - y) + (1 - a_v)\delta(1 - \epsilon_{v'} - y)]
$$

$$
+ n_v \left(\frac{y}{A}\right)^{\alpha_v} \left(1 - \frac{y}{A}\right)^{\beta_v} + n_s \left(\frac{y}{A}\right)^{\alpha_s} \left(1 - \frac{y}{A}\right)^{\beta_s}
$$

- the only framework using evolution in Mellin space & have PDFs also for *x^N >* 1

- only standard DIS data sets (semi-global)
- no error analysis

FIG. 3: Data on nuclear Drell Yan cross sections rates to

Noticeably, while some parameters show a clear depen-

dence on the size of the nucleus, such as the shifts in the \blacksquare

at moderate and large x^N , those related to the shape of the nucleus effective densities at small x^N , such as αv, $\overline{\mathbb{R}}$. well known A dependence of shadowing effects at small and dependence of shadowing effects at small and dependence of small at small and dependence of small and dependence of small and dependence of shadowing effects at sma

deuterium and those computed with NLO nPDF.

OVERVIEW NPDF 0.95 1 10 100 0.8 1 10 100) Q2 (GeV2

F2Ca / F2

 Review of existing global analyses of nuclear PDF compared in the LO and NLO and NLO and NLO and \mathcal{C} The dashed analyses of nuclear σ and σ is an and σ and σ are shown by the data of σ $\frac{1}{2}$ examples of medium and large nuclei, we take the γ^2 dof-12 x^2 dof = 1.2

HIRAI, KUMANO, NAGAI'07 [PRC76(2007)065207] LO, NLO, ERROR PDFS to decree with increasing \mathcal{L}

- uses multiplicative factor are different. This kind of different with in-

$$
f_i^A(x_N, Q_0^2) = R_i(x_N, Q_0, A, Z) f_i(x_N, Q_0^2)
$$

 \overline{C} and \overline{C} and \overline{C} approximate in large in where proton PDF in MRST 1998 and factor

$$
R_i(x, A, Z) = 1 + \left(1 - \frac{1}{A^{\alpha}}\right) \frac{a_i + b_i x + c_i x^2 + d_i x^3}{(1 - x)^{\beta_i}}
$$

- neglects region x>1 is still slightly underestimated at x = 0.0125; however,

 $t \in \mathbb{R}^n$ as \mathbb{R}^n are not as \mathbb{R}^n - includes all current DIS & DY data set (same as our T and T is a region of T are compared with the LO ones are compared with the LO ones are compared with the LO ones are T analysis - discussed later)

in Fig. 7 for the ratio F Ca ² /F ^D ² . The LO and NLO ratios - use Hessian method to produce error PDFs

q¯, and g) are shown in the NLO for all the analyzed nuclei

(6). The extreme values (x⁺

⁰i, ^x[−]

dependent of A in our current and a in our current in our current in \mathcal{A} Sec. II A, so that they are the same in Fig. 8. Although 1982.

culating the NPDFs and their uncertainties at given x

0i) are assumed to be in-

-0.2

[R(data) — R(theory)] / R(theory)

-0.2

 $\frac{1}{\alpha}$

Ξ.

F ^A¹

or σpA¹

DY /σpD

2 /F A2

FIG. 3: (Color online) Comparison with Drell-Yan data of

 \blacksquare in the NLO. The theoretical ratios are calculated at the Q² points of the experimental data. The uncertainties are esti-

used data in Figs. 1, 2, and 3 for the ratios F ^A

differences between experimental and theoretical values

NLO results are used and they are calculated at the experimental Q² points. The uncertainty bands are also shown in the NLO, and they are calculated at Q2=10 $\frac{1}{2}$ 50 GeV² for the Drell-Yan processes. The scale Q2=10 GeV² is taken because the average of all the F² data is of the order of this value. The scale is Q2=50 GeV² for

mated at $\overline{}$

(Rexp−Rtheo)/Rtheo, where R is R = F ^A

scale 20 GeV² is taken for the ratio of the σpA

the Drell-Yan ratios of the σpA

Because the NLO contributions are obvious only in

by the shaded areas.

Review of existing global analyses of nuclear PDF

ESKOLA, PAUKKUNEN, SALGADO'09 [JHEP0904(2009)065] LO, NLO, ERROR PDFS

- uses multiplicative factor

 $f_i^A(x_N, Q_0^2) = R_i(x_N, Q_0, A, Z) f_i(x_N, Q_0^2)$

where proton PDF in CTEQ6.1M and factor is a complicated piecewise defined function

 $R_i(x, A, Z) =$ $\sqrt{ }$ $\frac{1}{2}$ \overline{a} $a_0 + (a_1 + a_2x)(e^{-x} - e^{-x_a}) \quad x \leq x_a$ $b_0 + b_1x + b_2x^2 + b_3x^3$ $x_a \leq x \leq x_e$ $c_0 + (c_1 - c_2x)(1 - x)^{-\beta}$ $x_e \leq x \leq 1$

with A-dependent parameters

- neglects region x>1

- includes all current DIS & DY data set & π^0 RHIC data to constrain gluon

- use Hessian method to produce error PDFs

Figure 3: The nuclear modifications R^V , RS, R^G for Carbon (upper group of panels) and

Why another set of NPDFs ?

- nuclear PDFs different from proton PDFs PDF parameters contain more information but the fit has less data to constrain them
- big source of uncertainty systematic uncertainty connected to assumptions made about PDF parameters & the parametric form of the PDF at Q_0

Why another set of NPDFs ?

- nuclear PDFs different from proton PDFs - PDF parameters contain more information but the fit has less data to constrain them

- big source of uncertainty - systematic uncertainty connected to assumptions made about PDF parameters & the parametric form of the PDF at Q_0

1. nCTEQ framework entirely different from previous nPDF frameworks useful to study parameterization dependence & estimate systematic uncertainty

2. nCTEQ analysis is in a close relation to the existing CTEQ proton analysis allows to calculate nuclear correction factors in a flexible way (Q dependent & based on global analysis)

3. Our analysis aims at including also neutrino DIS data

nCTEQ framework for nuclear PDF - based on CTEQ6M proton fit

- functional form for bound protons same as for free proton PDF (restrict x to 0<x<1)

 $x f_k(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} e^{c_3 x} (1+e^{c_4 x})^{c_5}$ $k = u_v, d_v, g, \bar{u} + \bar{d}, s, \bar{s}$ $d(x, Q_0)/\bar{u}(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} + (1+c_3x)(1-x)^{c_4}$

- coefficients with A-dependance (reduces to proton for A=1)

$$
c_k \to c_k(A) \equiv c_{k,0} + c_{k,1} (1 - A^{-c_{k,2}}), \quad k = \{1, \ldots, 5\}
$$

- PDF for a nucleus with A-nucleons out of which Z-protons

$$
f_i^{(A,Z)}(x,Q) = \frac{Z}{A} f_i^{p/A}(x,Q) + \frac{A-Z}{A} f_i^{n/A}(x,Q)
$$

Note: PDF of neutron are related to the proton by isospin symmetry

- Input scale and other input parameters as in CTEQ6M proton analysis

$$
Q_0 = m_c = 1.3
$$
 GeV $m_b = 4.5$ GeV $\alpha_s(m_Z) = 0.118$

Experiments included in the analysis:

CERN BCDMS & EMC & NMC DESY Hermes SLAC E-139 & E-049 FNAL E-665 *N* = (D*,* Al*,* Be*,* C*,* Ca*,* Cu*,* Fe*,* Li*,*Pb*,* Sn*,*W) $N = (D, C, Ca, Pb, Xe)$ $N = (D, He, N, Kr)$ $N = (D, Ag, Al, Au, Be, C, Ca, Fe, He)$

FNAL E-772 & E-886 $N = (D, C, Ca, Fe, W)$

NPDF fit properties:

- we fit nuclear data with NLO QCD predictions & include heavy quark effects (ACOT)
- added nuclear observables to CTEQ fitting routines (need to treat 2 nuclei at once)
	- *DIS:* F_2^A/F_2^A' Drell-Yan: $\sigma_{DY}^{pA}/\sigma_{DY}^{pA'}$
- applied standard CTEQ kinematical cuts Q>2GeV & W>3.5GeV

NPDF fit results:

- 708 (1233) data points after (before) cuts
- 32 free parameters 675 degrees of freedom
- overall χ^2 /dof = 0.95
- individually for different data sets
	- for F_2^A/F_2^D χ^2 /pt=0.92
	- for $F_2^A/F_2^{A'}$ χ^2 /pt=0.69
	- for $\sigma_{DY}^{pA}/\sigma_{DY}^{pA'}$ χ^2 /pt=1.08

16

 \odot CTEQ parameters dependent on atomic number A - $c_k(A)/c_{k,0}$

for parton distributions

 $\int d_v u_v$ $g \quad \bar{u} + \bar{d}$ "

 \odot Parton density functions for bound partons as a function of \times

19

Example of different assumptions in nuclear gluon PDF

nCTEQ estimate of gluon nPDF uncertainty

- vary gluon nPDF assumptions & parametrizations
- large uncertainty for low x<0.1 in nCTEQ framework
- need further data to constrain gluon nPDF

20

 \odot Comparison of iron F_2 from neutrino and charged lepton DIS

$$
R[F_2^{Fe}] = F_2^{Fe}/F_2^D
$$

Phys.Rev.D80 094004, 2009 Phys.Rev.D77 054013, 2008

Re-analyze neutrino data within the same framework as for charged lepton

Neutrino DIS cross-section data

NuTeV & di-muon **CHORUS** $N = Fe$ $N = Pb$ \rightarrow 2310 data points \rightarrow 824 data points

All charged lepton DIS & Drell-Yan data \rightarrow 708 data points

Challenges in combining the neutrino & charged lepton data

- deal with the disparity of number of data points assigning weights to neutrino data
- neutrino DIS data only with 2 heavy nuclei insufficient to get a reliable A-dependance
- do all neutrino data show the different behavior or only NuTeV ?

Properties of neutrino fits

- CHORUS data are in good agreement with the charged lepton data combined: χ^2 /pt=1.03

- NuTeV data (with correlated errors) difficult to fit alone or with the charged lepton data alone: χ^2 /pt=1.35 combined: χ^2 /pt=1.33
- Neutrino data dominate the combined fit without re-weighting final result depends from the weight chosen

with only NuTeV Consistency check

Analysis of fits with different weights of neutrino DIS (corr. errors)

- Nuclear correction factors - $R = F_2^{Fe}/F_2^{Fe,0}$

w	$l^{\pm}A$	χ^2 (/pt)	νA	χ^2 (/pt)	total χ^2 (/pt)
$\vert 0 \vert$	708	638(0.90)	THE REAL		638 (0.90)
1/7		708 645 (0.91)		3134 4710 (1.50)	5355(1.39)
1/2	708	680 (0.96)		$3134 \mid 4405 \; (1.40)$	5085(1.32)
$\left \begin{array}{c} 1 \\ \end{array} \right $		708 736 (1.04)		3134 4277 (1.36)	5014(1.30)
∞	\sim - \sim 100 \pm	$\mathcal{L} = \mathcal{L} \times \mathcal{L}$		$3134 \mid 4192 \; (1.33)$	4192(1.33)

Analysis of fits with different weights of neutrino DIS (corr. errors)

- Nuclear correction factors - $R = F_2^{Fe}/F_2^{Fe,0}$

w	$l^{\pm}A$	χ^2 (/pt)	νA	χ^2 (/pt)	total χ^2 (/pt)
$\vert 0 \vert$	708	638 (0.90)			638(0.90)
1/7	708	645 (0.91)	3134	4710 (1.50)	5355(1.39)
1/2	708	680 (0.96)	3134	4405(1.40)	5085(1.32)
1	708	736(1.04)		3134 4277 (1.36)	5014(1.30)
∞	$\mathbf{r} = \mathbf{r}$			$3134 \mid 4192 \; (1.33)$	4192(1.33)

24

Analysis of fits with different weights of neutrino DIS (corr. errors)

- χ^2 - distribution criterion

 $P(\chi^2, N) = \frac{(\chi^2)^{N/2 - 1} e^{-\chi^2/2}}{2^{N/2} \Gamma(\Lambda^2)}$ $2^{N/2}\Gamma(N/2)$

CTEQ hep-ph/0101051 MSTW arXiv:0901.0002 [hep-ph]

 \int ξ 90 *P PDFs defined as 90% C.L.* \longrightarrow $\int^{320} P(\chi^2, N) d\chi^2 = 0.90$

Analysis of fits with different weights of neutrino DIS (corr. errors)

- χ^2 - distribution criterion

 $P(\chi^2, N) = \frac{(\chi^2)^{N/2 - 1} e^{-\chi^2/2}}{2^{N/2} \Gamma(\Lambda^2)}$ $2^{N/2}\Gamma(N/2)$

CTEQ hep-ph/0101051 MSTW arXiv:0901.0002 [hep-ph]

 \int ξ 90 *P*(χ^2 , N) $d\chi^2 = 0.90$ error PDFs defined as 90% C.L. \longrightarrow $\int^{300} P(\chi^2, N) d\chi^2 = 0.90$

Analysis of fits with neutrino DIS (uncorrelated errors)

- Nuclear correction factors - $R = F_2^{Fe}/F_2^{Fe,0}$

26

uncorrelated errors

correlated errors

OUTLOOK

Global nuclear CTEQ fit is able to describe the charged lepton data well

- some challenges on the way to a comprehensive public nPDF release
- relax kinematical cuts and fit Fermi motion peak in a natural way
- error PDFs & realistic estimate of uncertainty

 \odot Incompatibility of neutrino DIS with charged lepton DIS

- incompatibility a "precision" effect the result changes when using uncorrelated errors
- tension in NuTeV data \rightarrow high χ^2 of the fit to NuTeV alone \rightarrow problem of NuTeV data ?
- NOMAD data can help decide

The impact of nuclear PDF from neutrino DIS on proton PDF

- how does the incompatibility of neutrino DIS impact the uncertainty of strange quark PDF ?

27

