

The proton charge radius: a persistent mystery

Paul Indelicato



✓ CREMA: Muonic Hydrogen Collaboration





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_**√_LKB** 8 juillet 2010



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From collisions to atomic physics

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Proton size

vendredi 18 novembre 2011

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Proton form factor

2 quarks up (2/3 e) + 1 quark down (-1/3 e) + strong interaction (gluons)

Vertex EM interaction: Dirac and Pauli Form factors (S, P: spin and 4-momentum of nucleon, f: quark flavor)

$$\begin{split} \langle P', S' | V_{(f)}^{\mu} | P, S \rangle &= \bar{U}(P', S') \bigg[\gamma^{\mu} F_{1}^{(f)}(Q^{2}) \\ &+ i \sigma^{\mu\nu} \frac{q_{\nu}}{2M_{N}} F_{2}^{(f)}(Q^{2}) \bigg] U(P, S), \\ V_{(f)}^{\mu} &= \bar{\psi}_{(f)} \gamma^{\mu} \psi_{(f)}, \end{split}$$

Physical charge density are derived from the Sachs Form factors

$$G_E(Q^2) = F_1(Q^2) - \frac{Q^2}{(2M_N)^2} F_2(Q^2),$$

$$G_M(Q^2) = F_1(Q^2) + F_2(Q^2).$$

Measure the moments of the charge distribution:



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 $G_N(q^2) = \int dr e^{-i\boldsymbol{q}\cdot\boldsymbol{r}} \frac{\rho_N(\boldsymbol{r})}{4\pi},$

 $< r^n > = \int_0^\infty r^{2+n} \rho(r) dr,$

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Metrology in hydrogen

From collisions to muonic atoms

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Proton size



Hydrogen



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Hydrogen



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L'hydrogène



Lamb 1948

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L'hydrogène





Zoom sur l'hydrogène





L'hydrogène



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The Rydberg constant

Relative uncertainty in $R_{\scriptscriptstyle \infty}$

The Rydberg constant is deduced from the measurements of the 1S-2S and 2S-nD transitions and from the scaling law 1/n³ of the Lamb shift. The limitation is due to the uncertainty in the 2S-nD measurement.



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Proton size effect and VP potential





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Lamb shift = self-energy + vacuum polarization + proton radius

2S-2P	self-energy	vacuum pol.	rp	total
е-р	1084.1 MHz	-26.9 MHz	0.146 MHz	1057.8 MHz
µ-p	0.17 THz	-50.94 THz	0.96 THz	-49,81THz

electronic hydrogen : e-p

muonic hydrogen : µ-p





Exotic atom



Experiment



Challenges

- production of muonic hydrogen in 2S
- powerful triggerable 6µm laser
- small signal analysis

Aim: better determination of proton radius r_p

Principle of the experiment





The muonic hydrogen experiment

Getting up close and personal with the proton!

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The PSI accelerator facility

Proton beam, 570 MeV, 2 mA, 1.14 MW



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The cyclotron trap





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muonic hydrogen source







Slowing down pions

Pion lifetime is 80 ns



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Slowing down pions





Precision Determination of the d $\pi {\leftrightarrow}$ NN Transition

Strength at Threshold, T. Strauch, F.D. Amaro, D.F.
Anagnostopoulos, P. B hler, D.S. Covita, H. Gorke, D.
Gotta, A. Gruber, A. Hirtl, P. Indelicato, E.O. Le Bigot,
M. Nekipelov, J.M.F. dos Santos, S. Schlesser, P.
Schmid, L.M. Simons, M. Trassinelli, J.F.C.A. Veloso and
J. Zmeskal. Phys. Rev. Lett. 104, 142503 (2010).



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Pionic and muonic Hydrogens X-ray spectroscopy



Line Shape of the mu H(3p-1s) Hyperfine Transitions, D.S. Covita, D.F. Anagnostopoulos, H. Gorke, D. Gotta, A. Gruber, A. Hirtl, T. Ishiwatari, P. Indelicato, E.-O.L. Bigot, M. Nekipelov, J.M.F.d. Santos, P. Schmid, L.M. Simons, M. Trassinelli, J.F.C.A. Veloso and J. Zmeskal. Phys. Rev. Lett. 102, 023401 (2009).



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The laser trigger signal





A muon's Odyssey

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A muon's Odyssey



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Laser chain : thin-disk Yb: YAG laser


Laser chain : Ti:Sa





Rubidiu

Cesium

lode

- home made Ti:Sa lasers
- cw-Ti:Sa frequency controlled with FP, atom, molecule (abs. freq. <50MHz)
- short length pulsed oscillator seeded with Cw-Ti:Sa (1.2 mJ, 6 ns, delay 50 ns, $\Delta v = 200$ MHz)
- multipass amplifier ´ 6 (12 mJ, 5 ns)

Antognini et al, Opt. Comm. 253 (2005) p.362

Laser chain : Raman cell



Laser chain : multipass cavity

 \rightarrow illuminate at 6 µm all the muon stopping volume (5'15'190 mm³)





X-rays analysis \rightarrow event gate sorting \rightarrow noise rejection

Example : FP 900 - 11 hrs meas.

1.56 million detector events

expected 2-3 laser induced events/hour !



time signature in LAAPD

- photon < 10keV \rightarrow 1 shot in the LAAPD
- e⁻ in B = 5T \rightarrow many counts in detectors



energy signature in LAAPD

- E > 8keV ⇔ electron
- 1keV < E < 8keV ⇔ X ray
- E<1keV ⇔ neutron



X-rays analysis \rightarrow noise rejection

Example : FP 900 - 11 hrs meas.

- ≽ 400 µ⁻/s
- > 240 laser shot/s
- 860 000 laser shot/hour
- 1.56 million detector clicks
- > 19600 clicks in the laser region
- expected 2-3 laser induced events/hour !



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Time spectra



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Resonance searches (2002, 2003, 2007)

2002 : Run with all the parts at the same time (search for resonance 10 hrs).



 $\mu p: 2S_{1/2}(F=1)- 2P_{3/2}(F=2)$ movie resonance search



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muonic hydrogen : $2S_{1/2}(F=1) - 2P_{3/2}(F=2)$



R. Pohl, A. Antognini, F. Nez, et al., Nature 466, 213 (2010).

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Proton charge radius (~0.1%)

49.95

muonic hydrogen: 2S_{1/2}(F=0)- 2P_{3/2}(F=1)



~ at the position deduced with new $r_p \rightarrow hfs$: Zeemach radius (few %)

muonic deuterium : $2S_{1/2}(F=3/2) - 2P_{3/2}(F=5/2)$



Not at the position estimated with new r_p and isotopic shift \rightarrow deuteron polarizability

muonic deuterium : $2S_{1/2}(F=1/2)-2P_{3/2}(F=3/2)$ and $2S_{1/2}(F=1/2)-2P_{3/2}(F=1/2)$



Laser chain : frequency calibration

FSR measured/controlled in cw with I₂ (1 ph abs), Cs (2 ph fluo), Rb (2 ph fluo), lines



 $\mu p (2S_{1/2}(F=1)-2P_{3/2}(F=2))$: uncertainty budget

Statistics

uncertainty on position (fit)

541 MHz (~ 4 % of Γ_{nat})

 $\Delta v_{\text{experimental}}$ = 20 (1) GHz (Γ_{nat} = 18.6 GHz)

 Sources : Laser frequency (H₂0 calibration) AC and DC stark shift Zeeman shift (5 Telsa) Doppler shift Collisional shift 	300 MHz < 1 MHz < 30 MHz < 1 MHz 2 MHz
TOTAL UNCERTAINTY ON FREQUENCY	618 MHz
 Broadening : 6 µm laser line width Doppler Broadening Collisional broadening 	~ 2 GHz < 1 GHz 2.4 MHz

Updated: v (µp : $2S_{1/2}(F=1)-2P_{3/2}(F=2)$) = 49 881.16 (62) GHz (12.5 ppm) Nature: v (µp : $2S_{1/2}(F=1)-2P_{3/2}(F=2)$) = 49 881.88 (76) GHz (16 ppm) $\mu p (2S_{1/2}(F=0)-2P_{3/2}(F=1))$: uncertainty budget

Statisticsuncertainty on position (fit)	960 MHz
 Sources : Laser frequency (H₂0 calibration) AC and DC stark shift Zeeman shift (5 Telsa) Doppler shift Collisional shift 	300 MHz < 1 MHz < 30 MHz < 1 MHz 2 MHz
TOTAL UNCERTAINTY ON FREQUENCY	1006 MHz
 Broadening : 6 µm laser line width Doppler Broadening Collisional broadening 	~ 2 GHz < 1 GHz 2.4 MHz

$v (\mu p : 2S_{1/2}(F=0) - 2P_{3/2}(F=1)) = 54611.87 (1.0) \text{ GHz}$ (18.5ppm)





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Proton size



















• The two lines obey to:

$$E_{{}^{5}P_{3/2}} - E_{{}^{3}S_{1/2}} = \Delta E_{\rm LS} + \Delta E_{\rm FS} + \frac{3}{8}\Delta E_{\rm HFS}(2p_{3/2}) - \frac{1}{4}\Delta E_{\rm HFS}(2s)$$
$$E_{{}^{3}P_{3/2}} - E_{{}^{1}S_{1/2}} = \Delta E_{\rm LS} + \Delta E_{\rm FS} - \frac{5}{8}\Delta E_{\rm HFS}(2p_{3/2}) + \frac{3}{4}\Delta E_{\rm HFS}(2s)$$

 $\Delta E_{\rm HFS}(2s)^{\rm Exp.} = 22.8082 \pm 0.0078 \,\mathrm{meV} = 5514.99 \,\mathrm{GHz}$

 $\Delta E_{\rm HFS}(2s)^{\rm Mart.} = 22.8148 \pm 0.0078 \,\mathrm{meV}$ $\Delta E_{\rm HFS}(2s)^{\rm Carlson} = 22.8146 \pm 0.0049 \,\mathrm{meV}$



$$\Delta \tilde{E} = 209.9779(49) - 5.2262 r_p^2 + 0.0347 r_p^3 \text{ meV}$$

R. Pohl, A. Antognini, F. Nez, et al., Nature 466, 213 (2010).

$E_{{}^{5}P_{3/2}} - E_{{}^{3}S_{1/2}} = 209.9759 - 5.22888R^{2} + 0.0357376R^{3} - 0.000045R^{4}$

P. Indelicato, P. Mohr, to be published

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Proton size



$$\Delta \tilde{E} = 209.9779(49) - 5.2262 r_p^2 + 0.0347 r_p^3 \text{ meV}$$

0.84184(67) fm

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$$\Delta \tilde{E} = 209.9779(49) - 5.2262 \ r_{\rm p}^2 + 0.0347 \ r_{\rm p}^3 \ {\rm meV}$$

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$$E_{\rm th} = \left(209.9974(48) - 5.2262 \frac{r_p^2}{\rm fm^2}\right) \,\rm{meV}$$

U. D. Jentschura, Annals of Physics 326, 500 (2011).

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P. Indelicato, P. Mohr. to be published 0.84145(66) fm

0.84184(67) fm



$$\Delta \tilde{E} = 209.9779(49) - 5.2262 \ r_{\rm p}^2 + 0.0347 \ r_{\rm p}^3 \ {\rm meV}$$

R. Pohl, A. Antognini, F. Nez, et al., Nature 466, 213 (2010).

Model dependent

Martynenko: 2s Hyperfine structure: 24.1819-0.16018 Rz meV

 $E_{\,{}^5P_{3/2}}-E_{\,{}^3S_{1/2}}=209.6341-5.22888 Rp^2+0.0357376 Rp^3-0.000045 Rp^4+0.32036 Rz$

 $E_{\,{}^3P_{3/2}}-E_{\,{}^1S_{1/2}}=230.5679-5.22888 Rp^2+0.0357376 Rp^3-0.000045 Rp^4-0.96108 Rz$

$$R_{Z,e,e}^{\text{Exp.},1} = \frac{3R^4 + 9R^3R_{\text{M}} + 11R^2R_{\text{M}}^2 + 9RR_{\text{M}}^3 + 3R_{\text{M}}^4}{2\sqrt{3}(R + R_{\text{M}})^3}$$

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Proton size

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$$\Delta \tilde{E} = 209.9779(49) - 5.2262 r_p^2 + 0.0347 r_p^3 \text{ meV} \qquad [1.84]84(67) \text{ fm}$$

R. Pohl, A. Antognini, F. Nez, et al., Nature **466**, 213 (2010).
Model dependent

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$$R_{Z,e,e}^{\text{Exp.},1} = \frac{3R^4 + 9R^3R_{\text{M}} + 11R^2R_{\text{M}}^2 + 9RR_{\text{M}}^3 + 3R_{\text{M}}^4}{2\sqrt{3}(R + R_{\text{M}})^3}$$

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Proton size



Zemach radius

 Perturbation theory on HFS leads to a combined finite charge correction and magnetic moment distribution distribution (Zemach, 1958):

$$\rho_{em}(\boldsymbol{r}) = \int \rho(\boldsymbol{r} - \boldsymbol{u}) \mu(\boldsymbol{u}) d\boldsymbol{u}$$

$$R_{\mathbf{Z}} = \langle r_{\mathbf{Z}} \rangle = \int r \rho_{em}(\mathbf{r}) d\mathbf{r}.$$



$$\langle R_{Z,a,b}^{n} \rangle = \int d^{3}r_{1}d^{3}r_{2}\rho_{a}(r_{1})\rho_{b}(r_{2})|\boldsymbol{r}_{1}-\boldsymbol{r}_{2}|^{n}$$

$$\Delta E_{\rm HFS} = E_{\rm F} \left(1 - 2\alpha m_{\mu} \left\langle r_{Z,e,m}^1 \right\rangle \right)$$

Using the Dipole model for the charge distribution

$$\langle R_{Z,e,e}^3 \rangle^2 = \frac{3675}{256} \left[\langle r^2 \rangle \right]^3$$

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Proton size



• From the muonic hydrogen LS:

- Zemach radius: Charge radius: 0.84185 fm
- Zemach radius: 1.079 (41) fm
- Magnetic radius (assuming a dipole model) distribution: 0.867 (66) fm

• Mainz experiment:

- Magnetic moment distribution: 0:777(13)_{stat}(9)_{syst}(5)_{model}(2)_{group} fm
- Zemach radius: 1.085 (3) fm cited as 1.045 in [1]
- Proton-structure corrections to hyperfine splitting in muonic hydrogen, C.E.
 Carlson, V. Nazaryan et K. Griffioen. Phys. Rev. A 83, 042509 (2011).



TABLEI. Hyperfine splitting for the 2S state of muonic hydrogen, using different modern analytic fits in the terms that involve elastic form factors.

Form factor fit	$E_{\rm HFS}^{2S}$ (meV)	r_Z (fm)
AMT [18]	22.8123	1.080
Kelly [19]	22.8141	1.069
AS [20]	22.8105	1.091
Mainz 2010 [21-23]	22.8187	1.045
Muonic Hydrogen	1.079 (41)	
Hydrogen (Volotka et al. 2005)	1.045 (16)	
Protons	size LKB 04/07/2011	



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$$\langle R_{Z,a,b}^{n} \rangle = \int d^{3}r_{1}d^{3}r_{2}\rho_{a}(r_{1})\rho_{b}(r_{2})|\boldsymbol{r}_{1}-\boldsymbol{r}_{2}|^{n}$$

$$\Delta E_{\rm HFS} = E_{\rm F} \left(1 - 2\alpha m_{\mu} \left\langle r_{Z,e,m}^{1} \right\rangle \right)$$

$$\rho_{em}(\boldsymbol{r}) = \int \rho(\boldsymbol{r} - \boldsymbol{u}) \mu(\boldsymbol{u}) d\boldsymbol{u}$$

$$R_{\mathbf{Z}} = \left\langle r_{\mathbf{Z}} \right\rangle = \int r \rho_{em}(\boldsymbol{r}) d\boldsymbol{r}.$$

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$$\langle R_{Z,e,e}^3 \rangle^2 = \frac{3675}{256} \left[\langle r^2 \rangle \right]^3$$

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2s Hyperfine structure



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How to get the radius from hydrogen

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Proton size

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Hydrogen



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Hydrogen+Deuterium



Analysis by F. Biraben



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- Muonic hydrogen: 0.84184(67) fm
- Hydrogen mean value: 0.8752(71) fm: 4.7 σ
- Combined Hydrogen+Deuterium mean value: 0.861(16) fm: 6.6 σ
- Mainz electron-proton scattering: 0.879(8) fm: 4.7 σ
- CODATA 2006: 0.8768(69) fm: 5.0 σ
- CODATA 2011: 0.8775(51) fm (uses latest scattering data): 6.9 σ



What could it be?

Douter de tout ou tout croire sont deux solutions également commodes, qui l'une et l'autre nous dispensent de réfléchir. [H. Poincaré]

Doubting everything or believing everything are two easy solutions, which both allow not to think

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• QED

- Electron-proton elastic scattering data analysis
- Under-estimated systematic errors in some hydrogen measurements
 - possible, but many different kind of experiments (microwave, 1s-3s, 2s-ns and 2s-nd)
- New physics

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New, very accurate experiments, but...

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Using experimental proton form factors



Dispersion analysis of the nucleon form factors including meson continua, M.A. Belushkin, H.W. Hammer et U.-G. Meißner. Phys. Rev. C 75, 035202 (2007).

p-QCD approach 0.830 (0.822 . . . 0.835) fm

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Using experimental proton form factors



Dispersion analysis of the nucleon form factors including meson continua, M.A. Belushkin, H.W. Hammer et U.-G. Meißner. Phys. Rev. C 75, 035202 (2007).

Super-convergence approach 0.844 (0.840 . . . 0.852) fm

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Using experimental shapes



W. Melnitchouk et J.A. Tjon. Phys. Rev. Extracted radius: 0.850 fm

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- Dipole model lead to exponential charge density. The high moment of the distribution contributes a lot at high q.
- There is no region of q where the <r²> dominates the finite size effect to >98%, and the finite size effect is large enough to get an accuracy <2%
- People have been letting float the F(0) value in fits (should be exactly 1), which spoils exact behavior where it is important for getting <r²>.

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Model-independent extraction of the proton charge radius from electron scattering, R.J. Hill and G. Paz. Physical Review D 82, 113005 (2010).

Third Zemach moment of the proton, I.C. Cloët and G.A. Miller. Phys. Rev. C 83, 012201 (2011).

The RMS charge radius of the proton and Zemach moments, M.O. Distler, J.C. Bernauer and T. Walcher. Physics Letters B 696, 343-347 (2011).

QED is not endangered by the proton's size, A. De Rújula. Physics Letters B **693**, 555-558 (2010). *QED confronts the radius of the proton*, A. De Rújula. Physics Letters B **697**, 26-31 (2011).

Precision Measurement of $\mu G_{E}^{p}/G_{M}^{p}$ at Low Q2, S. Gilad. Few-Body Systems 1-3 (2011).

High-Precision Determination of the Electric and Magnetic Form Factors of the Proton, A1 Collaboration, J.C. Bernauer, P. Achenbach, C. Ayerbe Gayoso, R. Böhm, D. Bosnar, L. Debenjak, M.O. Distler, L. Doria, A. Esser, H. Fonvieille, J.M. Friedrich, J. Friedrich, M. Gómez Rodríguez de la Paz, M. Makek, H. Merkel, D.G. Middleton, U. Müller, L. Nungesser, J. Pochodzalla, M. Potokar, S. Sánchez Majos, B.S. Schlimme, S. Širca, T. Walcher and M. Weinriefer. Phys. Rev. Lett. **105**, 242001 (2010).

Troubles with the Proton rms-Radius, I. Sick. Few-Body Systems 1-3 (2011).

High Precision Measurement of the Proton Elastic Form Factor Ratio pGE=GM at Low Q2 (arXiv:1102.0318v1 [nucl-ex] 1 Feb 2011), X. Zhan et al (Jefferson Lab)





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	SC approach	Explicit pQCD app.	Ref. [23]	Recent determ.
r_E^p (fm)	0.844 (0.840 0.852)	0.830 (0.8220.835)	0.848	0.886(15) [72-74]
r_M^p (fm)	0.854 (0.8490.859)	0.850 (0.8430.852)	0.857	0.855(35) [73,75]
$(r_E^n)^2$ (fm ²)	-0.117 (-0.110.128)	-0.119 (-0.1080.13)	-0.12	-0.115(4) [52]
r_M^n (fm)	0.862 (0.8540.871)	0.863 (0.8590.871)	0.879	0.873(11) [76]

Resonance	Mass (GeV)	a_1 (GeV ²)	a_2 (GeV ²)	Γ (GeV)
ω	0.782	0.755960	0.370592	_
ϕ	1.019	-0.776537	-2.913229	_
<i>s</i> ₁	1.124860	0.902379	2.484859	_
<i>s</i> ₂	2.019536	0.022798	-0.130622	5.158635
v_1	1.062128	-0.127290	-2.162533	_
v_2	1.300946	-1.243412	3.704233	_
v_3	1.493630	4.191380	-7.091021	_
v_4	1.668522	-3.176013	3.723858	_
v_5	2.915451	0.048987	0.075965	19.088297

Dispersion analysis of the nucleon form factors including meson continua, M.A. Belushkin, H.W. Hammer and U.-G. Meißner. Phys. Rev. C **75**, 035202 (2007).

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[1] QED is not endangered by the proton's size, A. De Rújula. Physics Letters B 693, 555-558 (2010)

[2] QED confronts the radius of the proton, A. De Rújula. Physics Letters B 697, 26-31 (2011)

[3] The RMS charge radius of the proton and Zemach moments, M.O. Distler, J.C. Bernauer et T. Walcher. Physics Letters B 696, 343-347 (2011).

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Comparison of Zemach's 3rd moments



• $\langle r^3 \rangle_{(2)}=3.789 \langle r^2 \rangle^{3/2}$ Dipole • $\langle r^3 \rangle_{(2)}=1.960 \langle r^2 \rangle^{3/2}$ Gauss • $\langle r^3 \rangle_{(2)}=3.91 \langle r^2 \rangle^{3/2}$ Arrington (2007) • $\langle r^3 \rangle_{(2)}=3.78(13) \langle r^2 \rangle^{3/2}$ Friar & Sick (2005) • $\langle r^3 \rangle_{(2)}=4.18(13) \langle r^2 \rangle^{3/2}$ Distler, Bernauer, Walcher (2011) • $\langle r^3 \rangle_{(2)}=36.6\pm7.3=51 \langle r^2 \rangle^{3/2}$ De Rujula (2010), retracted (?) in 2011...

[1] QED is not endangered by the proton's size, A. De Rújula. Physics Letters B 693, 555-558 (2010)

[2] QED confronts the radius of the proton, A. De Rújula. Physics Letters B 697, 26-31 (2011)

[3] The RMS charge radius of the proton and Zemach moments, M.O. Distler, J.C. Bernauer

et T. Walcher. Physics Letters B 696, 343-347 (2011).

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_∕∕_lkb

$\Delta \tilde{E} = 209.9779(49) - 5.2262 r_p^2 + 0.0347 r_p^3 \text{ meV}$



Third Zemach moment of the proton, I.C. Cloët and G.A. Miller. Phys. Rev. C 83, 012201 (2011).

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MAMI A1 experiment





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MAMI A1 experiment

11/04/2011

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Are scattering data analysis reliable?



105, 242001 (2010).

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Are scattering data analysis reliable?



Are scattering data analysis reliable?





Are scattering data analysis reliable?



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Proton size results summary







Predicted radius much smaller, but not accurate

11/04/2011



QCD vacuum and the proton...



http://www.physics.adelaide.edu.au/~dleinweb/VisualQCD/QCDvacuum/welcome.html

11/04/2011

Should we go beyond charge distributions?



vendredi 18 novembre 2011

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Proton polarization

- Several calculations
 - Rosenfelder (1999)



 $b E^{W_1(0,Q^2)} + 0016 - 0.0127(5) \text{ meV},$

$$\Delta E_{2s}^{\text{p.pol}} = -\frac{136 \pm 30}{n^3} \,\mu\text{eV} = -0.017 \pm 0.004 \,\text{meV}.$$

- Pachucki (1999) $\Delta E_{2s}^{\text{p.pol}} = -0.012 \pm 0.002 \text{ meV},$

- Martynenko (2006) $\Delta E_{2s}^{\text{p.pol}} = \Delta E^{\text{subt.}} + \Delta E^{\text{inel.}}$ = 0.0023 - 0.01613 meV = -0.0138(29) meV,
- Carlson and Vanderhaeghen (2011) $\Delta E_{2s}^{\text{p.pol}} = \Delta E^{\text{subt.}} + \Delta E^{\text{inel.}} + \Delta E^{\text{el.}}$
- = 0.0053(19) 0.0127(5) 0.0295(13) meV= -0.0074(20) 0.0295(13) meV. $\Delta E_{2s}^{\text{p.pol}} = \Delta E^{\text{subt.}} + \Delta E^{\text{inel.}}$ $= \left[\delta E^{W_1(0,Q^2)} + \delta E^{\text{proton pole}} \right] + \delta E^{\text{continuum}}$

Could be wrong by 0.04 meV

DPF 2011



G. A. Miller, A. W. Thomas, J. D. Carroll, and J. Rafelski: Natural Resolution of the Proton Size Puzzle (arXiv:1101.4073v2 [physics.atom-ph], March 29th, 2011)



FIG. 1: Direct two-photon exchange graph corresponding to the hitherto neglected term. The dashed line denotes the lepton; the solid line, the nucleon; the wavy lines photons; and the ellipse the off-shell nucleon.

0.31 meV for μH and 9Hz for hydrogen (with model dependent parametrization)



New physics

Predictions are always difficult, in particular about the Future [N. Bohr]

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Proton size

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Where could it come from?

- A muon edm? If $d_{\mu} = 2 \times 10^{-19}$ e·cm would shifts the energy level < 200 MHz
- Charge equality between e^- and μ^- generation? Checked to $u_r = 10^{-8}$ (from $\mu^+ e^-$)
- Deviation from Coulomb's law: probe of hidden sector
 - Test of Coulomb law via spectroscopy is very clean and model independent probe of new particles It is independent on stability and decay channel.

$$V(r) = -\frac{Z\alpha}{r} \left(1 + \alpha' e^{-mr}\right) \quad \text{or} \quad V(r) = -\frac{Z\alpha}{r} \left(1 + \alpha'' (\mathbf{s_1} \cdot \mathbf{s_2}) e^{-mr}\right)$$

- From simple atoms there are constraints on light bosons with ultra-weak coupling:

 $m \in [1eV, MeV]$ and $\alpha' < 10^{-13}, \alpha'' < 10^{-17}$ [PRL 104,220406 (2010), arXiv:1008.3536v2]

• Minicharge particles? [Jaeckel and Roy (2010), Jentschura(2010)]

Vacuum polarization with pair production of light fermions with $q = \varepsilon e$ and masses $m_{\varepsilon} < m_e$

No parameter found explaining $r_{
m p}$ puzzle without contradicting, simple atoms spectroscopy, $g_{e/\mu} - 2$, α ...

- New bound-state QED theory?
 - Non-local in time interaction...(R. K. Gainutdinov)?

11/04/2011



Proton Size Anomaly, V. Barger, C.-W. Chiang, W.-Y. Keung et al. Phys. Rev. Lett. 106, 153001 (2011):

We explore the possibility that new scalar, pseudoscalar, vector, and tensor flavorconserving nonuniversal interactions may be responsible for the discrepancy.We consider exotic particles that, among leptons, couple preferentially to muons and mediate an attractive nucleon-muon interaction. We find that the many constraints from low energy data disfavor new spin-0, spin-1, and spin-2 particles as an explanation.

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Muonic hydrogen and MeV forces, D. Tucker-Smith et I. Yavin. Physical Review D 83, 101702 (2011).

We explore the possibility that a new interaction between muons and protons is responsible for the discrepancy between the CODATA value of the proton-radius and the value deduced from the measurement of the Lamb shift in muonic hydrogen. We show that a new force carrier with roughly MeV-mass can account for the observed energyshift as well as the discrepancy in the muon anomalous magnetic moment. However, measurements in other systems constrain the couplings to electrons and neutrons to be suppressed relative to the couplings to muons and protons, which seems challenging from a theoretical point of view. One can nevertheless make predictions for energy shifts in muonic deuterium, muonic helium, and true muonium under the assumption that the new particle couples dominantly to muons and protons.

DPF 2011
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Muonic hydrogen and MeV forces, D. Tucker-Smith et I. Yavin. Physical Review D 83, 101702 (2011).





New Parity-Violating Muonic Forces and the Proton Charge Radius, B. Batell, D. McKeen et M. Pospelov. Phys. Rev. Lett. 107, 011803 (2011).

The recent discrepancy between proton charge radius measurements extracted from electron-proton versus muon-proton systems is suggestive of a new force that differentiates between lepton species. We identify a class of models with gauged right-handed muon number, which contains new vector and scalar force carriers at the 100 MeV scale or lighter, that is consistent with observations. Such forces would lead to an enhancement by several orders-of-magnitude of the parity-violating asymmetries in the scattering of low-energy muons on nuclei. The relatively large size of such asymmetries, $O(10^{-4})$, opens up the possibility for new tests of parity violation in neutral currents with existing low-energy muon beams.

DPF 2011

B. Batell, D. McKeen, and M. Pospelov, Phys. Rev. Lett. 107, 011803 (2011).

$$\Delta r_p^2|_{e-\mathrm{H}} = -\frac{6\kappa^2}{m_V^2}; \quad \Delta r_p^2|_{\mu-\mathrm{H}} = -\frac{6(\kappa^2 + \eta)}{m_V^2}f(am_V) \quad (11)$$

where $a = (\alpha m_{\mu} m_{p})^{-1} (m_{\mu} + m_{p})$ is the μ -H Bohr radius, $f(\hat{x}) \equiv \hat{x}^{4} (1 + \hat{x})^{-4}$, and $\eta \equiv \kappa g_{R}/(2e)$. The difference $\Delta r_{p}^{2}|_{e-H} - \Delta r_{p}^{2}|_{\mu-H}$ must be consistent with the observed pattern (5) and requires η to be positive. In the scaling regime of $am_{V} \gg 1$ one has

$$\frac{\eta}{m_V^2} \simeq \frac{\Delta r^2}{6} \simeq 0.01 \text{ fm}^2 \simeq \frac{2.5 \times 10^{-5}}{(10 \text{ MeV})^2}.$$
 (12)

In the same regime, the model predicts that future experiments with μ -He would detect the effective charge radius of the helium nucleus shifted down by $\Delta r_{\text{He}}^2 = -0.06 \text{ fm}^2$.

DPF 2011



What's next

Deuterium: deuton polarization too large Helium?

11/04/2011

Prospect : muonic helium spectroscopy at PSI



improve He spectroscopy



CREMA 2011





- We have performed a 15 ppm measurement of the Lamb-shift in muonic hydrogen
- The deduced proton radius using a Dipole model is 5 standard deviations away from the hydrogen and electron-proton elastic scattering data
- Better modeling of the proton form-factor required to confirm or reduce the disagreement
- Experiment checked with 2nd µH line and 3 µD lines
- Muonic He in 2013 (check of theory, different laser wavelength-in the red)

11/04/2011



One day on the life of a Proton Size Investigator







One day on the life of a Proton Size Investigator









Proton Size Investigators thank you for your attention



Hydrogen spectroscopy

Transition energy :
$$(E_n - E_p) = (E_n^{Dirac} - E_p^{Dirac}) + (E_n^{Re\,coil} - E_p^{recoil}) + (L_n - L_p)$$

in which : E_{Dirac} and E_{Recoil} have an exact formulation prop. to R_{∞}
 $L_n = F(QED, R_p)$ Lamb shift

✓ Q.E.D. = series expansion of α , Z α , ln(Z α) which depends on n as 1/n³

$$\Delta E = \left(\frac{\alpha}{\pi}\right)^2 \frac{(Z\alpha)^4}{n^3} m_e c^2 \left[B_{40} + (Z\alpha)B_{50} + (Z\alpha)^2 \left(L^3 B_{63} + L^2 B_{62} + L B_{61} + B_{60}\right) + \dots\right]$$

$$\approx \dots \left[0,54 - 0,16 + 5 \times 10^{-5} \left(-280 - 50 + 500 - 100\right) + \dots\right]$$

avec
$$L = \ln \left[(Z\alpha)^{-2} \right]$$

✓ Distribution charge radius contributions vary as 1/n³

$$E_{NS} = \frac{2}{3} \left(\frac{\mu}{m_e}\right)^3 \frac{(Z\alpha)^2}{n^3} m_e c^2 \left(\frac{Z\alpha R_p}{\lambda_c}\right)^2 \delta_{l0}$$

Q.E.D. calculations are right ! : 2 experimental data R, R_p
 are Q.E.D. calculations right ? : 3 experimental data R, R_p, Q.E.D.



The end!

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Proton size

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Laser chain : principle



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Lamb shift - form factor - polarizability



A. Antognini thesis, P.Indelicato P.Mohr in preparation

Muonic hydrogen Lamb shift

Name	Reference	formula	value		Pachucki [1]	Borie [3]
NR One loop electron VP		$\frac{8\pi(Z_{2})^{2}}{5\pi\pi^{2}}Q_{n1}^{(3)}m_{r}$	205.00740		205.0074	
Relativistic one loop VP	[3]					205.0282
NR two-loop electron VP		$\frac{m(2-n)}{2} Q_{n1}^{(2)} m_{r}$	1.50790		1.5079	1.5081
NR three-loop electron VP	[12]	$0.08353(1) \frac{\alpha^2 (Z \alpha)^2}{\kappa^3} m_{\tau}$	0.00530		check	
Polarization insertion in two coulomb lines	[4,1]	$0.01244 \frac{a^{2}(2a)^{2}}{2a^{2}} m_{\pi}$	0.15090		0.1509	0.1510
Polarisation insertion in two and three coulomb lines	[12]	$0.036506(4) \frac{n^2(Zn)^2}{n^2} m_{\pi}$	0.00230			
Relativistic correction	[2,4,1,3]	$\alpha(Z\alpha)^4$	0.01690		0.0109	
Wichmann-Kroll	[13, 14, 3]	$\frac{m(\pi_{\pi})}{m}Q_{\mu}^{WK}m_{\pi}$	-0.00103			-0.00303
Radiative photon and electron polarization in the Coulomb line $\alpha^2 (Z \alpha)^4$	[4,1]		-0.00500	0.0010	-0.005	
Electron loop in the radiative photon		$\left[1.09426\left(-\frac{m}{5m_{r}}-1\right)-4 \times 2.21656\right]$	-0.00150			
of order $\alpha^2 (Z\alpha)^4$		$\times \frac{n^{4}(2n)^{4}}{2^{2n^{2}}} (\frac{m_{r}}{n})^{3} m$				
Mixed electron and muon loops	[15]	$\frac{44\alpha^2(2\alpha)^2}{45\pi^2\pi^2} \left(\frac{m_e}{m}\right)^2 Q_{n1}^{(3)} m_e$	0.00007			0.00007
Hadronic polarization $\alpha(Z\alpha)^4m$	[16]	$-0.638(22)\frac{4\alpha(Z\alpha)^{*}}{15\pi^{2}n^{2}}(\frac{m_{\pi}}{m})^{2}m\delta_{10}$	0.01050	0.0004	0.0113	0.0110
Hadronic polarization $\alpha(Z\alpha)^*m$			0.00005			
Hadronic polarization in the radiative shoton $\alpha^2 (Z\alpha)^4 m$			-0.00002			
Becoil contribution	[17]	$\frac{(Z_{n})^{4}m^{2}}{2}(-\frac{1}{2}, -\frac{1}{2})(1-\delta_{n})$	0.05750		0.0575	0.0575
Badiative corrections of order $\alpha^{n}(Z\alpha)^{k}m$	0.11	Tables 32 33 37 38 39	-0.06770		-0.0077	-0.66788
Recail corrections of order $\alpha(Z\alpha)^*$ $\Im m$	[11]	Table 4.1	-0.04470		-0.045	-0.04497
Radiative Recoil corrections of	[11]	Table 5.1	-0.00968		-0.0099	-0.0096
order $\alpha(Z\alpha)^n \frac{m}{M}m$						
Nuclear structure correction of order $(Z\alpha)^5$	[1]	$-\frac{0.001(18)}{n^2}\delta_{20} \text{ meV}$	0.01200		0.012	0.015
Light by light electron loop contribution		-0.00112 /	0.00135	0.00135		
Polarization operator induced correction		2 meV	0.00019			
Radiative photon induced correction		$\frac{4.4004992}{2}\delta_{10} \text{ meV}$	-0.00001	d·4 G·	·1	
to nuclear polarizability $\alpha(Z\alpha)^5m$ Form	n factor	(dipole Gaussian)				
Sum			206.04310	0.00275	206.0353	206.04739
terms deperdent	on the nucle	ar radius from Ref. [11] (contributions i	in meV			
Leading nuclear size contribution		$\frac{2(2\pi\pi)}{5\pi^3}m_{\pi}^2 < r^2 > \delta_{10}$	-5.1975	$< r_{p}^{2} >$		
Nuclear size correction of order $(Z\alpha)^{6}$	[4,1]	$-\frac{(\pi - m)}{2m^2} - m_{\pi} < r^3 > (2) \delta_{10}$	0.034	$< r_{\xi}^{2} > -6$	$0.0363 < \tau_{p}^{2} >$	$0.0347 < r_{g}^{2} >$
Ramsche corrections to nuclear innite size effect	[1]	$\alpha(Z\alpha)$ m $< r_p >$	-0.0275	$< r_{F}^{*} >$		$-0.0273 < r_{p}^{*} >$
Nuclear size correction of order $(Z\alpha)^{\alpha} < \tau_{p}^{z} >$	[1]	$-\langle \ln \frac{n}{2} \rangle \frac{r_{e}^{2}}{2n^{2}} m_{\nu}^{2} < r_{p}^{2} >$		$< r_{p}^{2} >$		
Nuclear size correction of order $(Z\alpha)^{\circ} < r_{p}^{\circ} >$	[1]	$\frac{1}{3} \frac{m_{\mu}}{3m^{3}} - m_{\mu}^{*}(< r_{\mu}^{2} >)^{2}$		$< r_{p}^{2} >^{2}$		
Total < P" > contribution			-5.225			
	Table 7.	All known contributions from Eides et a	d m b			

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Most important contributions 1

Diagram	Value (in meV)	Name & order	References
O*	205.0074(1)	leading order VP $[\alpha(Z\alpha)^2]$	Galanin and Pomeranchuk [1952], Pachucki [1999, Table I], Eides et al. [2001, Eq. (208)]
	-3.862(108)	leading nuclear size contribution $[(Z\alpha)^4 m_{\rm T}^3 \langle r^2 \rangle]$	Eides et al. [2001, § 9.6], Pachucki [1999, Table I]
-00-*	1.5079(1)	two-loop EVP $[\alpha^2 (Z\alpha)^2 m]$	Di Giacomo [1969], Pachucki [1999, Table I], Eides et al. [2001, Eq. (213)]
	-0.6677(1)	muon self-energy + muon VP	Pachucki [1999, Table I], Eides et al. [2001, § 9.5]
	0.1509(1)	double EVP $[\alpha^2 (Z\alpha)^2]$	Pachucki [1999, Tabl0J3268 Eides et al. [2001, Eq. (215)], Pachucki [1996, Eq. (31)]
~~~~	0.0594(1)	rel. corr. to EVP $[\alpha(Z\alpha)^4]$	Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (223)]
	0.0575(1)	recoil of order $\alpha^4$ [ $\alpha^4$ ]	Barker and Glover [1955], Pachucki [1999, Table I], Eides et al. [2001, Table 11]
	-0.0440(1)	recoil corrections of order $(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5], Pachucki [1999, Table I]

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### Most important contributions 1

Diagram	Value (in meV)	Name & order	References
-0-*	205.0074(1)	leading order VP $[\alpha(Z\alpha)^2]$	Galanin and Pomeranchuk [1952], Pachucki [1999, Table I], Eides et al. [2001, Eq. (208)]
	-3.862(108)	leading nuclear size contribution $[(Z\alpha)^4 m_{\rm r}^3 \langle r^2 \rangle]$	Eides et al. [2001, § 9.6], Pachucki [1990, Table I]
-0-0-+	1.5079(1)	two-loop EVP $[\alpha^2 (Z\alpha)^2 m]$	Di Giacomo [1969], Pachucki [1999, Table I], Eides et al. [2001, Eq. (213)]
	-0.6677(1)	muon self-energy + muon VP	Pachucki [1999, Table I], Eides et al. [2001, § 9.5]
	0.1509(1)	double EVP $[\alpha^2(Z\alpha)^2]$	Pachucki [1999, Tabl0J3268 Eides et al. [2001, Eq. (215)], Pachucki [1996, Eq. (31)]
0	0.0594(1)	rel. corr. to EVP $[\alpha(Z\alpha)^4]$	Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (223)]
	0.0575(1)	recoil of order $\alpha^4$ [ $\alpha^4$ ]	Barker and Glover [1955], Pachucki [1999, Table I], Eides et al. [2001, Table 11]
	-0.0440(1)	recoil corrections of order $(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5], Pachucki [1999, Table I]

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	-3.862(108)	leading nuclear size contribution $[(Z\alpha)^4 m_r^3 \langle r^2 \rangle]$	Eides et al. [2001, § 9.6], Pachucki [1999, Table I]
-0-0-*	1.5079(1)	two-loop EVP $[\alpha^2 (Z\alpha)^2 m]$	Di Giacomo [1969], Pachucki [1999, Table I], Eides et al. [2001, Eq. (213)]
	-0.6677(1)	muon self-energy + muon VP	Pachucki [1999, Table I], Eides et al. [2001, § 9.5]
	0.1509(1)	double EVP $[\alpha^2(Z\alpha)^2]$	Pachucki [1999, Tabk013268 Eides et al. [2001, Eq. (215)], Pachucki [1996, Eq. (31)]
9	0.0594(1)	rel. corr. to EVP $[\alpha(Z\alpha)^4]$	Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (223)]
	0.0575(1)	recoil of order $\alpha^4$ [ $\alpha^4$ ]	Barker and Glover [1955], Pachucki [1999, Table I], Eides et al. [2001, Table 11]
	-0.0440(1)	recoil corrections of order $(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5], Pachucki [1999, Table I]

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### Most important contributions 1

Diagram	Value (in meV)	Name & order	References
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	-3.862(108)	leading nuclear size contribution $[(Z\alpha)^4 m_r^3 \langle r^2 \rangle]$	Eides et al. [2001, § 9.6], Pachucki [1999, Table I]
-0-0-*	1.5079(1)	two-loop EVP $[\alpha^2 (Z\alpha)^2 m]$	Di Giacomo [1969], Pachucki [1999, Table I], Eides et al. [2001, Eq. (213)]
	-0.6677(1)	muon sen-energy + muon VP	Pachucki [1999, Table I], Eides et al. [2001, § 9.5]
-0-*	0.1509(1)	double EVP $[\alpha^2(Z\alpha)^2]$	Pachucki [1999, Tabk0J3268 Eides et al. [2001, Eq. (215)], Pachucki [1996, Eq. (31)]
0	0.0594(1)	rel. corr. to EVP $[\alpha(Z\alpha)^4]$	Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (223)]
	0.0575(1)	recoil of order $\alpha^4$ [ $\alpha^4$ ]	Barker and Glover [1955], Pachucki [1999, Table I], Eides et al. [2001, Table 11]
	-0.0440(1)	recoil corrections of order $(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5], Pachucki [1999, Table I]

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### Most important contributions 1

Diagram	Value (in meV)	Name & order	References
O-*	205.0074(1)	leading order VP $[\alpha(Z\alpha)^2]$	Galanin and Pomeranchuk [1952], Pachucki [1999, Table I], Eides et al. [2001, Eq. (208)]
	-3.862(108)	leading nuclear size contribution $[(Z\alpha)^4 m_{\rm r}^3 \langle r^2 \rangle]$	Eides et al. [2001, § 9.6], Pachucki [1999, Table I]
-0-0	1.5079(1)	two-loop EVP $[\alpha^2 (Z\alpha)^2 m]$	Di Giacomo [1969], Pachucki [1999, Table I], Eides et al. [2001, Eq. (213)]
	-0.6677(1)	muon seit-energy + muon VP	Pachucki [1999, Table I], Eddes et al. [2001, § 9.5]
	0.1509(1)	double EVP $[\alpha^2(Z\alpha)^2]$	Pachucki [1999, Tabl013268 Eides et al. [2001, Eq. (215)], Pachucki [1996, Eq. (31)]
0	0.0594(1)	rel. corr. to EVP $[\alpha(Z\alpha)^4]$	Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (223)]
	0.0575(1)	recoil of order $\alpha^4$ [ $\alpha^4$ ]	Barker and Glover [1955], Pachucki [1999, Table I], Eides et al. [2001, Table 11]
	-0.0440(1)	recoil corrections of order $(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5], Pachucki [1999, Table I]

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### Most important contributions 2

Π	0.03(3)	light by light electron-loop contribution of order $\alpha^2(Z\alpha)^3m$ $[\alpha^2(Z\alpha)^3m]$	Pachucki [1996, p. 2095], Eides et al. [2001, § 9.3.2]	
×	0.0232(15)	nuclear size correction of order $(Z\alpha)^5$ $[(Z\alpha)^5m_t^3\langle r^3\rangle_{(2)}m]$	Pachucki [1996], Faustov and Martynenko [2000], Eides et al. [2001, Eq. (256)], Pachucki [1999]	
2	-0.0126(1)	(part of the) EVP with finite size $[\alpha(Z\alpha)^4m_t^3\langle r^2\rangle]$	Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (268)]	
	0.012(2)	proton polarizability $[(Z\alpha)^5m]$	Startsev et al. [1976], Rosenfelder [2000], Faustov and Martynenko [2000], Pachucki [1999, Table I], Eides et al. [2001, Eq. (261)]	
- *	0.0108(4)	$\begin{array}{l} \text{hadronic} \\ \text{polarization, order} \\ \alpha(Z\alpha)^4m \\ [\alpha(Z\alpha)^4m] \end{array}$	Folomeshkin [1974], Friar et al. [1999], Faustov and Martynenko [1999], Eides et al. [2001, Eq. (252)], Pachucki [1999, Table I]	
	-0.0099(1)	proton self-energy	Pachucki [1999, Table I]	200
	-0.0095(1)	radiative-recoil corrections of order $\alpha(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5]	
•••	-0.0083(1)	(part of the) EVP with finite size $[\alpha(Z\alpha)^4m_T^3\langle r^2\rangle]$	Friar [1979a], Friar [1981], Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (266)]	
-0-	-0.006(1)	muon self-energy with electron VP $[\alpha^2(Z\alpha)^4]$	Pachucki [1996, Eqs. (40) and (45)], Pachucki [1999, Table I], Eides et al. [2001, Eq. (237)]	1
•000-	0.0053(1)	three-loop electron polarization contribution, order $\alpha^3(Z\alpha)^2$ $[\alpha^3(Z\alpha)^2m]$	Kinoshita and Nio [1999], Eides et al. [2001, Eq. (214)], Pachucki [1999, Table I]	

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### Most important contributions 2

	0.03(3) 0.0232(15)	light by light electron-loop contribution of order $\alpha^2(Z\alpha)^3m$ $[\alpha^2(Z\alpha)^3m]$ nuclear size correction of order $(Z\alpha)^5$	Pachucki [1996, p. 2095], Eides et al. [2001, § 9.3.2] Pachucki [1996], Faustov and Martynenko [2000], Eides et al. [2001 Eq. (256)]	
2 0	-0.0126(1)	$[(Z\alpha)^5 m_t^3 \langle r^3 \rangle_{(2)} m]$ (part of the) EVP with finite size $[\alpha (Z\alpha)^4 m_t^3 \langle r^2 \rangle]$	Pachucki [1999] Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (268)]	
<u>}</u>	0.012(2)	proton polarizability $[(Z\alpha)^5m]$	Startsev et al. [1976], Rosenfelder [2000], Faustov and Martynenko [2000], Pachucki [1999, Table I], Eides et al. [2001, Eq. (261)]	
- *	0.0108(4)	hadronic polarization, order $\alpha(Z\alpha)^4m$ $[\alpha(Z\alpha)^4m]$	Folomeshkin [1974], Friar et al. [1999], Faustov and Martynenko [1999], Eides et al. [2001, Eq. (252)], Pachucki [1999, Table I]	
	-0.0099(1)	proton self-energy	Pachucki [1999, Table I]	200
	-0.0095(1)	radiative-recoil corrections of order $\alpha(Z\alpha)^n \frac{m}{M}m$	Eides et al. [2001, § 9.5]	
• O •	-0.0083(1)	(part of the) EVP with finite size $[\alpha(Z\alpha)^4m_r^3\langle r^2\rangle]$	Friar [1979a], Friar [1981], Pachucki [1996], Pachucki [1999, Table I], Eides et al. [2001, Eq. (266)]	
-0	-0.006(1)	muon self-energy with electron VP $[\alpha^2(Z\alpha)^4]$	Pachucki [1996, Eqs. (40) and (45)], Pachucki [1999, Table I], Eides et al. [2001, Eq. (237)]	
•000•	0.0053(1)	three-loop electron polarization contribution, order $\alpha^3(Z\alpha)^2$ $[\alpha^3(Z\alpha)^2m]$	Kinoshita and Nio [1999], Eides et al. [2001, Eq. (214)], Pachucki [1999, Table I]	

New calculation: S. G. Karshenboim, V. G. Ivanov, E. Y. Korzinin, et al., Phys. Rev. A 81, 060501 (2010) and 1 S. Karshenboim, E. Korzinin, V. Ivanov, et al., JETP Letters 92, 8 (2010). 0.00115(1) meV

1



# Check of QED calculations

- Non-perturbative effects due to proton size? Several papers mention «distortion of the wavefunction at the origin»...
  - Dirac + Vacuum polarization
  - Källén and Sabry
  - Self-energy
- Non-perturbative effects in all-order re-summation

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- One question is, of course, whether the QED corrections are complete or whether some higher order contributions are not yet sufficiently well studied. The most recent, comprehensive discussion of these possible deficiencies can be found in the paper by Edith Borie [22]. One might add that it could be necessary that the distorted relativistic wave functions had to be treated in a perturbative sum over multiphoton exchanges. The convergence of such sums is not proven in QED. Such a calculation may imply the solution of a Bethe-Salpeter equation and appears hardly possible [24]. However, further theoretical studies may finally show that the approximations made so far are adequate
- In «The RMS charge radius of the proton and Zemach moments», M.O. Distler, J.C. Bernauer et T. Walcher. Physics Letters B 696, 343-347 (2011).

_√Lkb

Direct numerical solution of Dirac equation, numerical grid, 10000 points, ~4000 inside the proton

$$\left(c\boldsymbol{\alpha}\cdot\boldsymbol{p}+\beta\mu_{\mathrm{r}}c^{2}+V_{\mathrm{Nuc}}(\boldsymbol{r})\right)\Phi_{n\kappa\mu}(\boldsymbol{r})=\mathcal{E}_{n\kappa\mu}\Phi_{n\kappa\mu}(\boldsymbol{r}),$$

$$V_{11}^{pn}(r) = -\frac{\alpha(Z\alpha)}{3\pi} \int_{1}^{\infty} dz \sqrt{z^2 - 1} \left(\frac{2}{z^2} + \frac{1}{z^4}\right) \frac{e^{-2m_e rz}}{r}$$
$$= -\frac{2\alpha(Z\alpha)}{3\pi} \frac{1}{r} \chi_1\left(\frac{2}{\lambda_e}r\right)$$



$$\chi_n(x) = \int_1^\infty dz e^{-xz} \frac{1}{z^n} \left( \frac{1}{z} + \frac{1}{2z^3} \right) \sqrt{z^2 - 1}.$$

$$V_{11}(r) = -\frac{2\alpha(Z\alpha)}{12\pi} \frac{1}{r} \int_0^\infty dr' \, r' \rho(r') \\ \times \left[ \chi_2 \left( \frac{2}{\lambda_e} \mid r - r' \mid \right) - \chi_2 \left( \frac{2}{\lambda_e} \mid r + r' \mid \right) \right].$$

Analytical expression for the evaluation of vacuum-polarization potentials in muonic atoms, S. Klarsfeld. Physics Letters 66B, 86-88 (1977).

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### Check point nucleus

#### Non relativistic

Pachucki 1996 calculation	205.006
Pachucki 1999 calculation	205.0074
Numerical calc with Pachucki data	205.007385
Numerical calc with CODATA 2006 data	205.007359
Relativistic	
Borie calc with ? data	205.0282
PI calc with CODATA 2006 data	205.02820
PM calc with CODATA 2006 data	205.028 201

Energies in meV

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## Charge radius dependence Coul.+VP



**Fig. 3.** Dependence of  $\frac{\Delta E_{V11FN}}{R^2}$  as a function of *R* in meV/fm² for different charge distribution models.

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Model	$a (\mathrm{meV/fm^2})$	$b ({\rm meV/fm^3})$
Uniform	-5.2284	0.0313
Dipole	-5.2271	0.0353
Fermi	-5.2271	0.0324
Gauss	-5.2265	0.0328
Ref.[34], Dip.	-5.2248	0.0347
Ref.[32]	-5.225	0.0347
Ref.[34], Gauss	-5.2248	0.0317



#### Källen and Sabry



#### $\Delta E = aR^2 + bR^3 + cR^4$

a         -0.0002145         -0.0002145         -0.0002           b         0.0000078         0.0000086         0.0000	mi Gaussian	Fermi	Exponential	Uniform	Model
<i>b</i> 0.0000078 0.000086 0.0000	46 -0.0002145	-0.0002146	-0.0002145	-0.0002145	а
0.000000 0.00000 0.0000	82 0.0000083	0.0000082	0.0000086	0.0000078	b
<i>c</i> -0.000008 -0.000009 -0.0000	08 -0.0000009	-0.000008	-0.0000009	-0.000008	С

_**∕**_lkb

#### Finite size correction on muon self-energy

#### All-orders calculations



z

$$E_{SE-NS} = \left(4\ln 2 - \frac{23}{4}\right)\alpha(Z\alpha)\mathcal{E}_{NS}$$

$$\mathcal{E}_{NS} = \frac{2}{3} \left(\frac{\mu_r}{m_{\mu}}\right)^3 \frac{(Z\alpha)^2}{n^3} m_{\mu} \left(\frac{Z\alpha < r >}{\mathcal{X}_{\rm C}}\right)^2$$

$$E_{SE-NS} = -0.000824 < r^2 >$$

(All-orders calculations- $E_{SE-NS} < r^2 >$ )/Z⁵ < $r^3 >$  1.8±1 x 10⁻⁵







Vacuum Polarization

H-like "One Photon" order ( $\alpha/\pi$ )

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H-like "Two Photon" order  $(\alpha/\pi)^2$ 



Two-loop self-energy (1s)

V. A. Yerokhin, P. Indelicato, and V. M. Shabaev, Phys. rev. A 71, 040101(R) (2005).

$$\Delta E_{\text{SESE}} = m \left(\frac{\alpha}{\pi}\right)^2 (Z\alpha)^4 \{B_{40} + (Z\alpha)B_{50} + (Z\alpha)^2 \\ \times [L^3 B_{63} + L^2 B_{62} + L B_{61} + G_{\text{SESE}}^{\text{h.o.}}(Z)]\}$$



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Proton size

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**Evolution at low-Z** 

All order numerical calculations

V. A. Yerokhin, Phys. Rev. A 80, 040501 (2009)

$$G_{\text{SESE}}^{\text{h.o.}}(Z=0) \equiv B_{60} = -84(15),$$

V. A. Yerokhin, P. Indelicato, and V. M. Shabaev, Phys. Rev. A 71, 040101(R) (2005).

Analytic calculations

$$B_{40} = 1.409244, B_{50} = -24.2668(31),$$
  
 $B_{63} = -8/27, B_{62} = 16/27 - (16/9) \ln 2,$ 

K. Pachucki and U. D. Jentschura, Phys. Rev. Lett. **91**, 113005 (4) (2003).

 $B_{60} = -61.6(9.2).$ 

U. D. Jentschura, A. Czarnecki, and K. Pachucki, Physical Review A 72, 062102 (2005).

$$B_{61} = 48.388913$$
  
 $\delta B_{61} = -1.4494...$ 

-127(42)

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Dyson: the expansion in  $\alpha$  of QED has zero convergence radius... e+-e, plus d'état lié!

Divergence of Perturbation Theory in Quantum Electrodynamics, F.J. Dyson. Physical Review 85, 631–632 (1952).

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## Example of unsolved problems

$$V_{\rm VP\infty}(k^2) = \frac{1}{k^2} \left[ 1 + \Pi(k^2) + \Pi(k^2) \frac{1}{k^2} \Pi(k^2) + \Pi(k^2) \frac{1}{k^2} \Pi(k^2) \frac{1}{k^2} \Pi(k^2) + \cdots \right]$$
$$= \frac{\Pi(k^2)}{1 - \Pi(k^2)}$$

$$\Pi\left(k^{2}\right) = \frac{2\alpha k^{2}}{3\pi} \int_{1}^{\infty} dz \frac{1}{z} \left(\frac{1}{z} + \frac{1}{2z^{3}}\right) \frac{\sqrt{z^{2} - 1}}{4m_{e}^{2}z^{2} + k^{2}}.$$

Singularity at  $k_0 = e^{\frac{3\pi}{2\alpha} + \frac{5}{6}} \approx 6.53 \times 10^{280}$  huge momenta = very short distances S. Brodsky, P. J. Mohr, P. Indelicato

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Proton size