

European Physical Society Conference, July 20–27, 2011

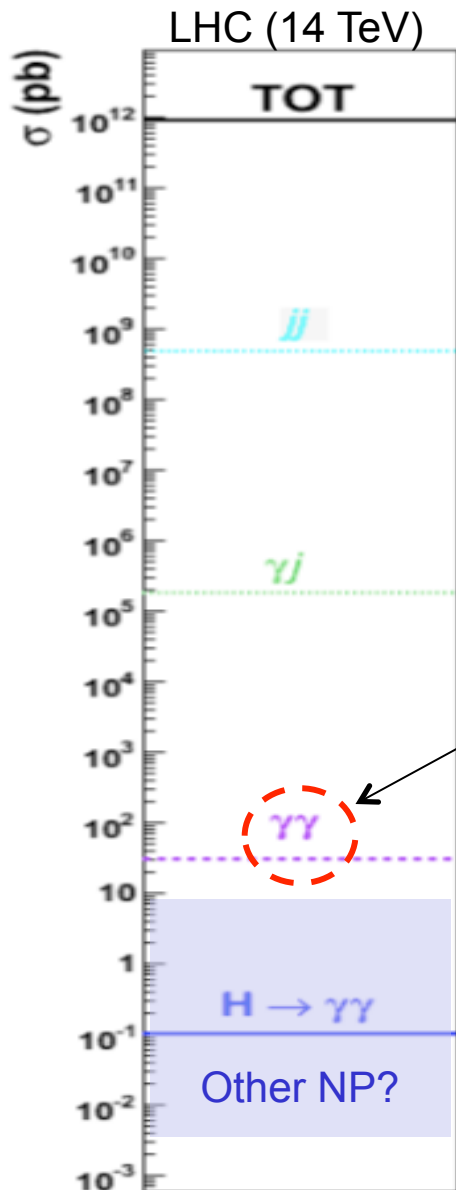
Measurement of the Cross section for Prompt Isolated Diphoton Production in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ TeV

Costas Vellidis

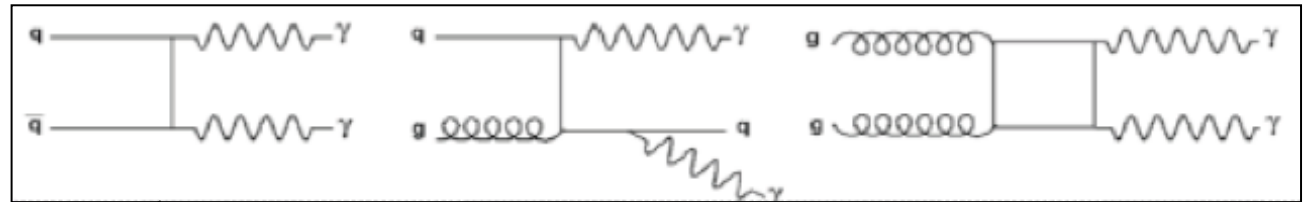
FNAL

On behalf of the CDF Collaboration

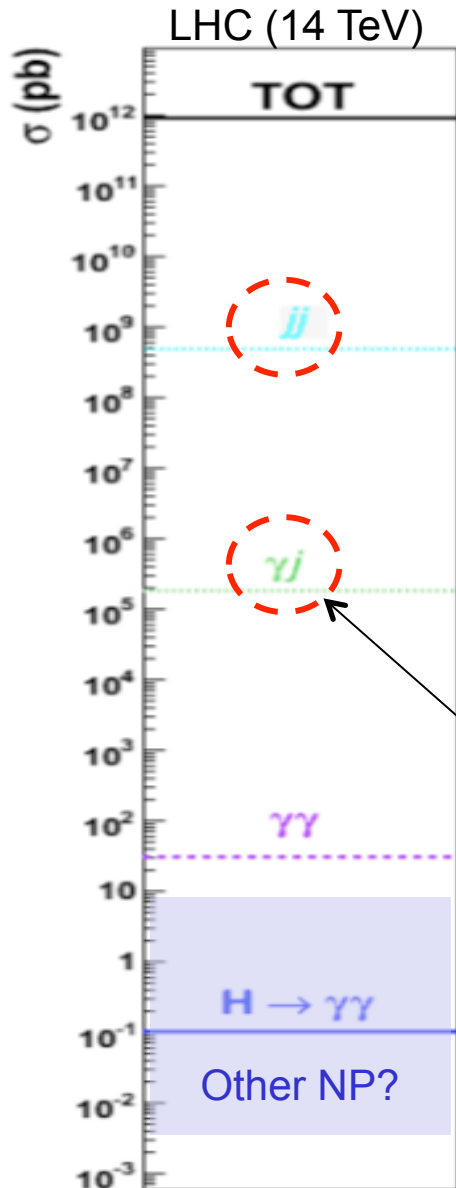
Prompt diphoton production at hadron colliders



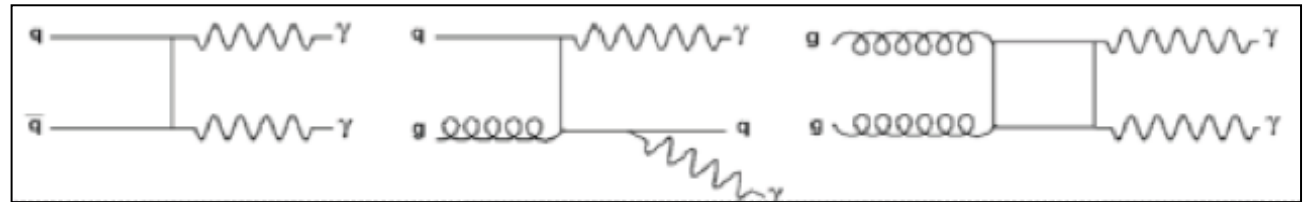
- **Prompt photons** = photons produced directly in perturbative scattering or via parton fragmentation (as opposed to non-perturbative photon production in meson decays).
- Main source of prompt diphoton production at hadron colliders via QCD interactions.



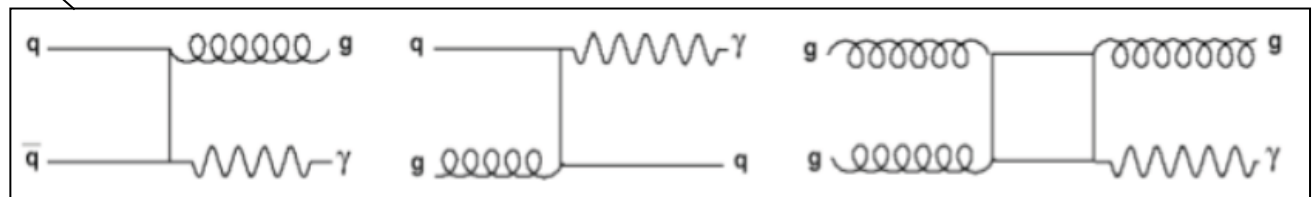
Prompt diphoton production at hadron colliders



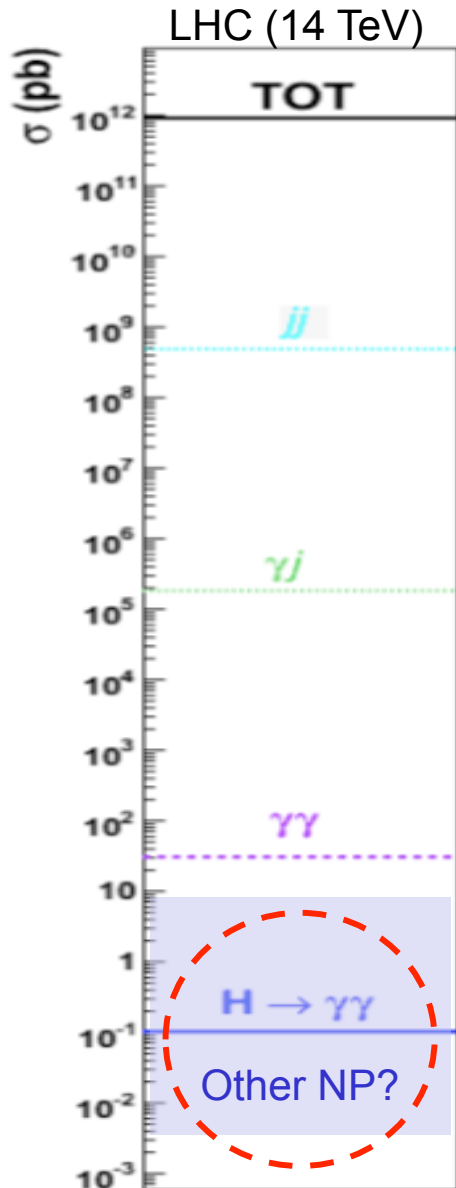
- **Prompt photons** = photons produced directly in perturbative scattering or via parton fragmentation (as opposed to non-perturbative photon production in meson decays).
- Main source of prompt diphoton production at hadron colliders via QCD interactions.



- **Main background: γ +jet and dijet**, with one or two jets misidentified as photons \rightarrow reducible background.

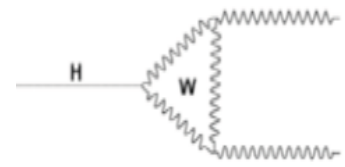


Prompt diphoton production at hadron colliders

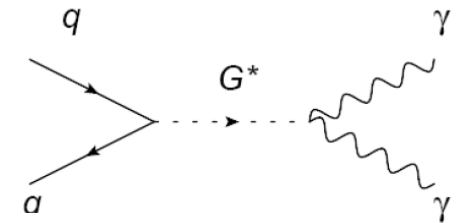


- **Prompt photons** = photons produced directly in perturbative scattering or via parton fragmentation (as opposed to non-perturbative photon production in meson decays).
- At much smaller rate, prompt diphotons may originate from more exotic (and exciting!) production mechanisms:

- Higgs decay

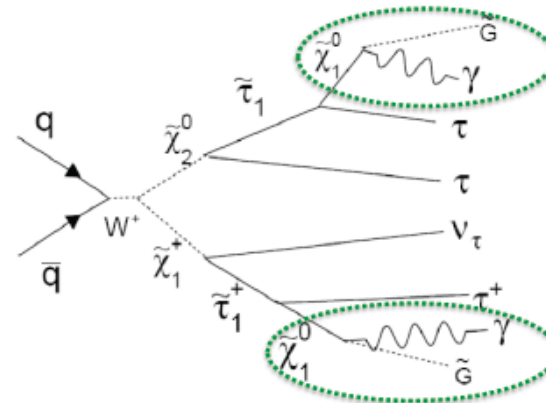


- Extra dimensions

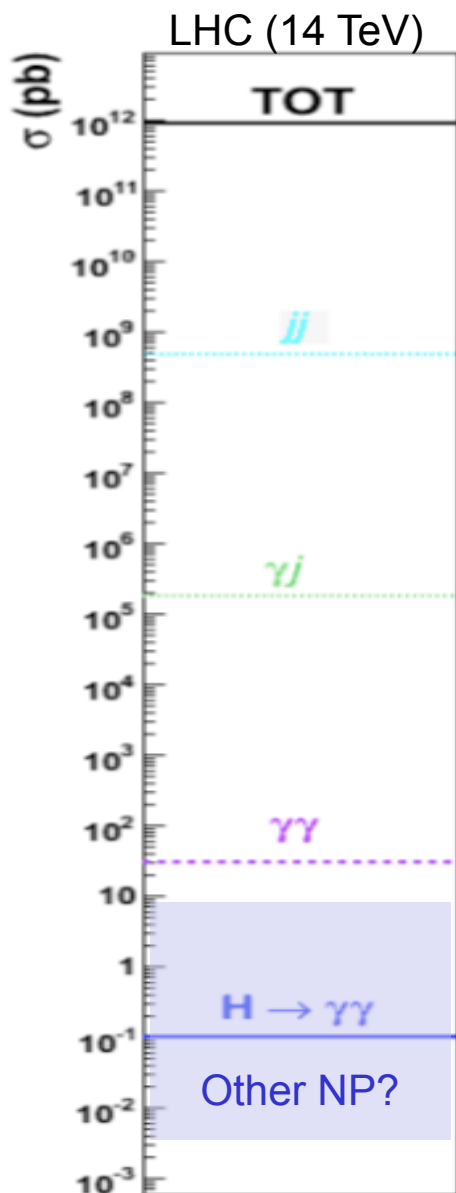


- SUSY

- ...



Prompt diphoton production at hadron colliders



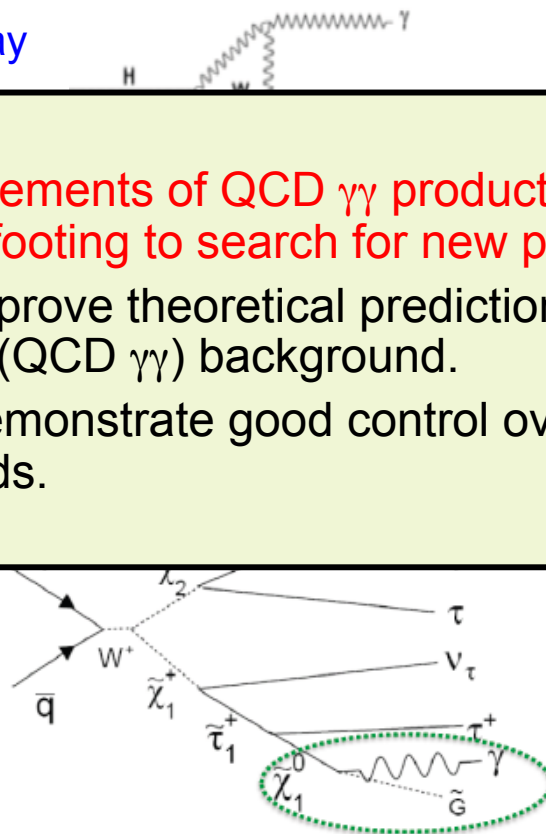
- **Prompt photons** = photons produced directly in perturbative scattering or via parton fragmentation (as opposed to non-perturbative photon production in meson decays).
- At much smaller rate, prompt diphotons may originate from more exotic (and exciting!) production mechanisms:
 - Higgs decay

Precise measurements of QCD $\gamma\gamma$ production should put us on solid footing to search for new physics:

- Validate/improve theoretical predictions for irreducible (QCD $\gamma\gamma$) background.
- Develop/demonstrate good control over reducible backgrounds.

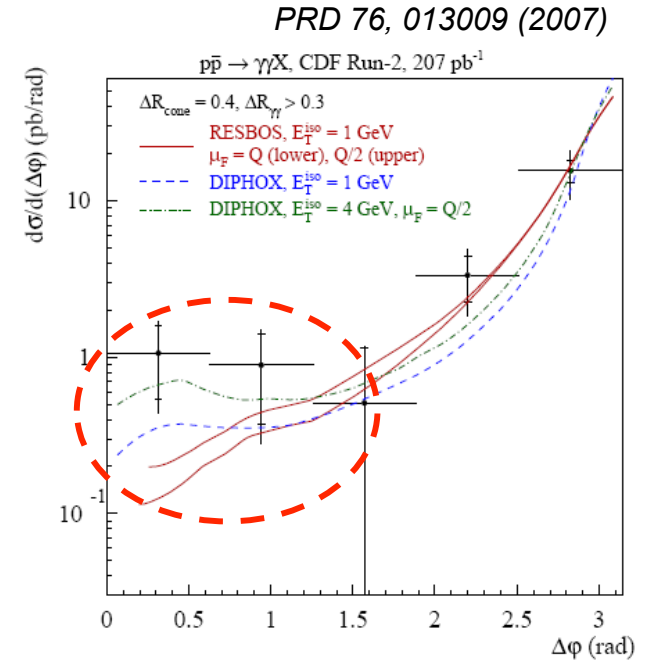
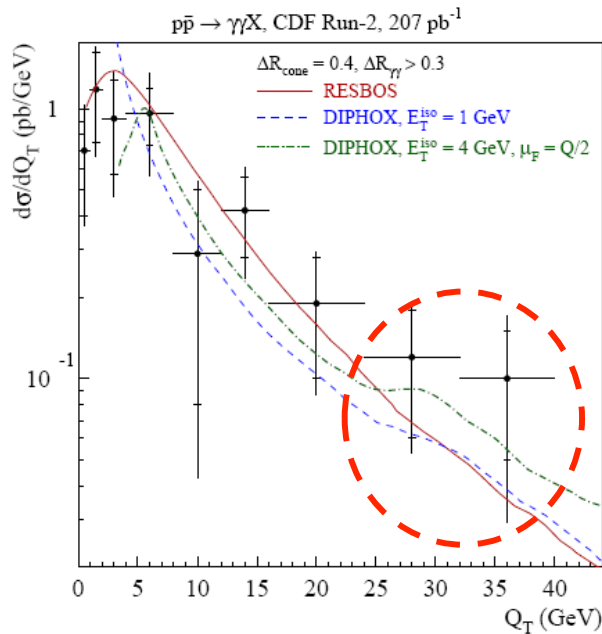
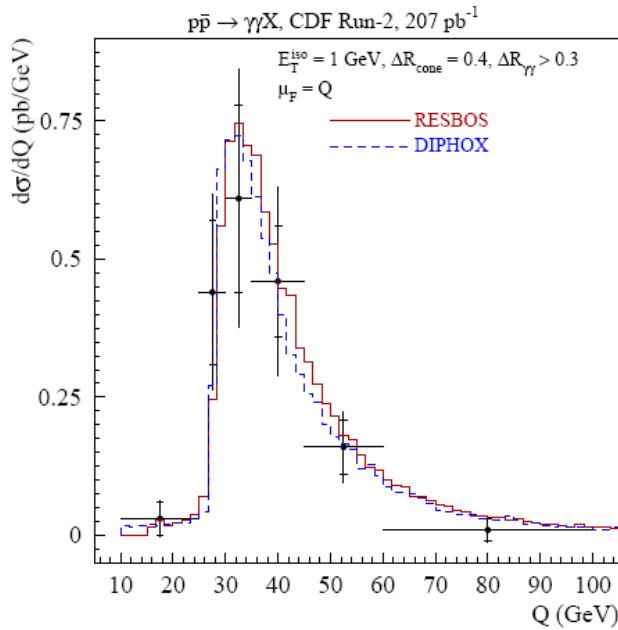
- SUSY

- ...



Previous Tevatron measurements

- CDF publication in Run II with 207 pb⁻¹. *PRL 95, 022003 (2005)*
- Event selection: $p_{T1(2)}=14(13)$ GeV, $|\eta_{1,2}|<0.9$, $\Delta R(\gamma,\gamma)<0.3$, $E_T^{\text{iso}}<1$ GeV.

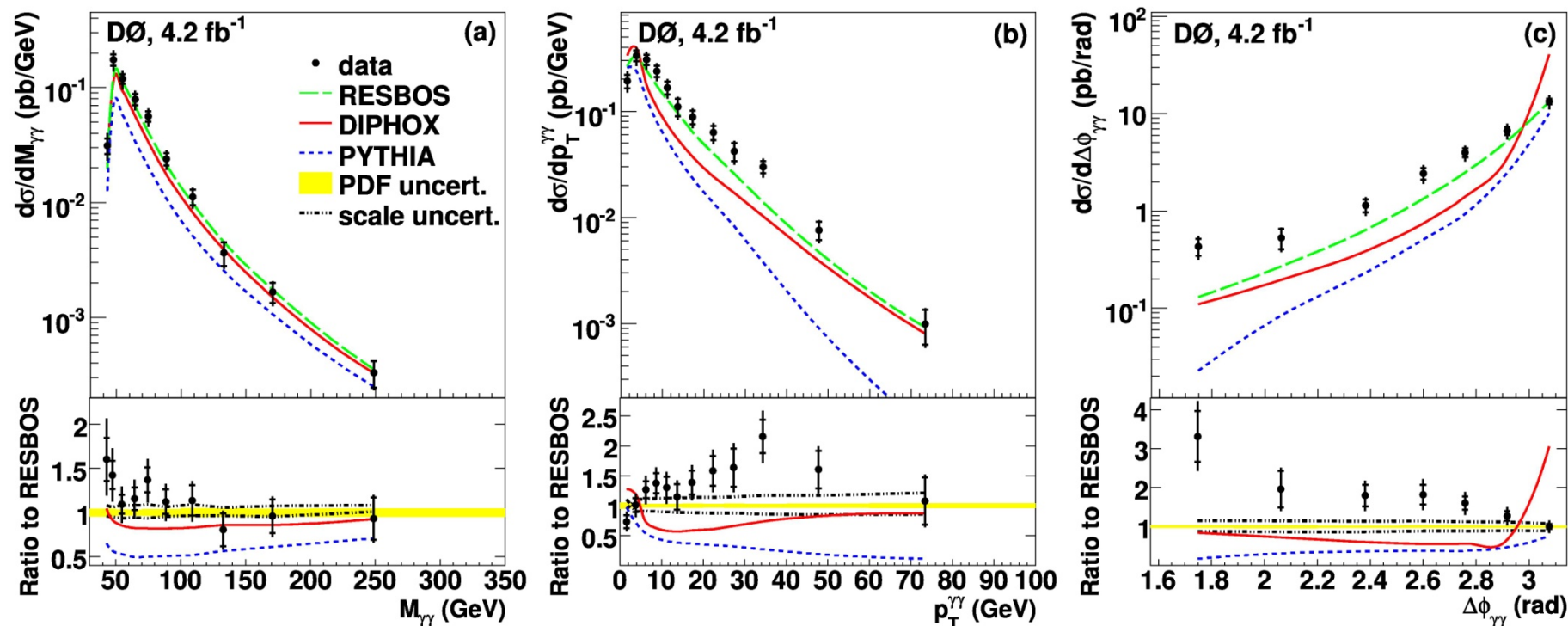


- $p_T(\gamma\gamma) > 25$ GeV region in data dominated by events with $p_T(\gamma\gamma) > M(\gamma\gamma)$ and $\Delta\phi(\gamma,\gamma) < \pi/2 \rightarrow$ potentially large fragmentation contributions.
- Large sensitivity of theoretical prediction on isolation requirement.

Here the Pythia prediction uses only matrix element based production of photons

Previous Tevatron measurements

- D0 publication in Run II with 4.2 fb^{-1} *PLB 690, 108 (2010)*
- $p_{T1(2)} = 21(20) \text{ GeV}/c$, $|\eta_{1,2}| < 1$, $\Delta R(\gamma, \gamma) > 0.4$, $(E_{\text{tot}}^{R=0.4} - E_{\text{em}}^{R=0.2}) / E_{\text{em}}^{R=0.2} < 0.1$, $p_T(\gamma\gamma) < M(\gamma\gamma)$

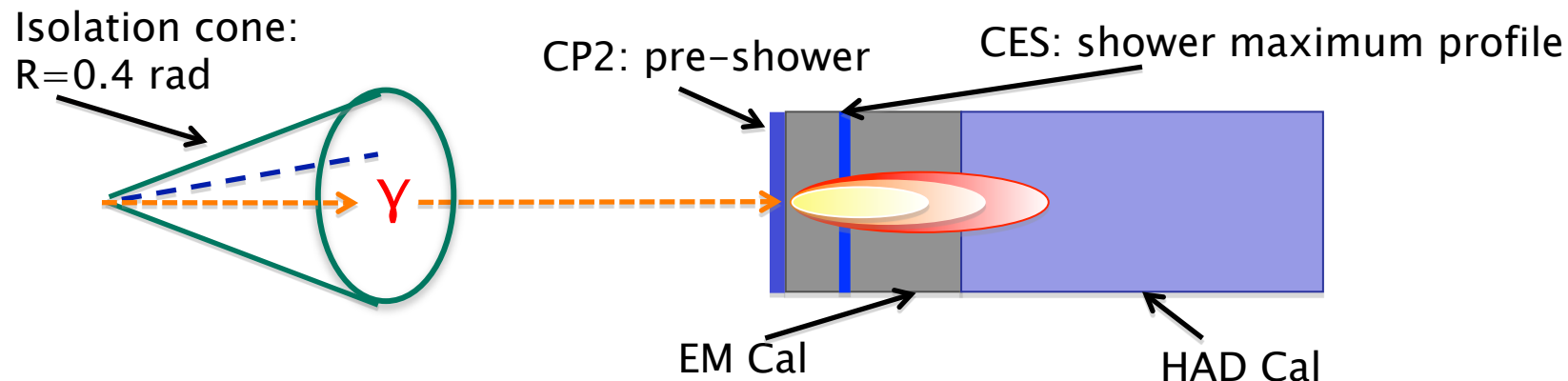


- Good agreement between data and RESBOS for $M_{\gamma\gamma} > 50 \text{ GeV}/c^2$
- Need for a resummed calculation
- Data spectrum harder than predicted
- Observable nearly insensitive to experimental effects
- Supports conclusion from $p_T(\gamma\gamma)$ measurement

(*) Overall normalization uncertainty (7.3%) not included in data error bars.

Here the Pythia prediction uses only matrix element based production of photons

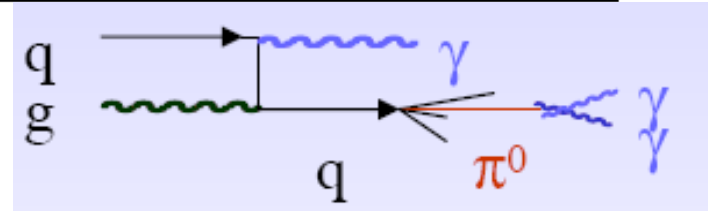
Photon identification and event selection



- Photons are selected offline from EM clusters, reconstructed within a cone of radius $R=0.4$ in the η - ϕ plane, and requiring:
 - Fiducial to the central calorimeter: $|\eta| < 1.1$ Avoids divergence in NLO calculation
 - $E_T \geq 17$ GeV (1st γ in the event), 15 GeV (2nd γ)
 - Isolated in the calorimeter: $I_{\text{cal}} = E_{\text{tot}}(R=0.4) - E_{\text{EM}}(R=0.4) \leq 2$ GeV
 - Low HAD fraction: $E_{\text{HAD}}/E_{\text{EM}} \leq 0.055 + 0.00045 \times E_{\text{tot}}/\text{GeV}$
 - At most one track in cluster with $p_T^{\text{trk}} \leq 1$ GeV/c + $0.005 \times E_T^\gamma/\text{c}$
 - Shower profile: $\chi^2_{\text{CES}} \leq 20$
 - E_T of 2nd CES cluster ≤ 2.4 GeV + $0.01 \times E_T$
- } Imply that $\Delta R(\gamma, \gamma) \geq 0.4$

Background

$$\frac{d\sigma}{dX} = \frac{N_{\gamma\gamma}}{\varepsilon \cdot A \cdot L \cdot \Delta}$$



Jets misidentified as photons: dijet and γ +jet

→ Fluctuations in jet fragmentation to leading π^0 or η^0 meson ($\pi^0, \eta^0 \rightarrow \gamma\gamma$)

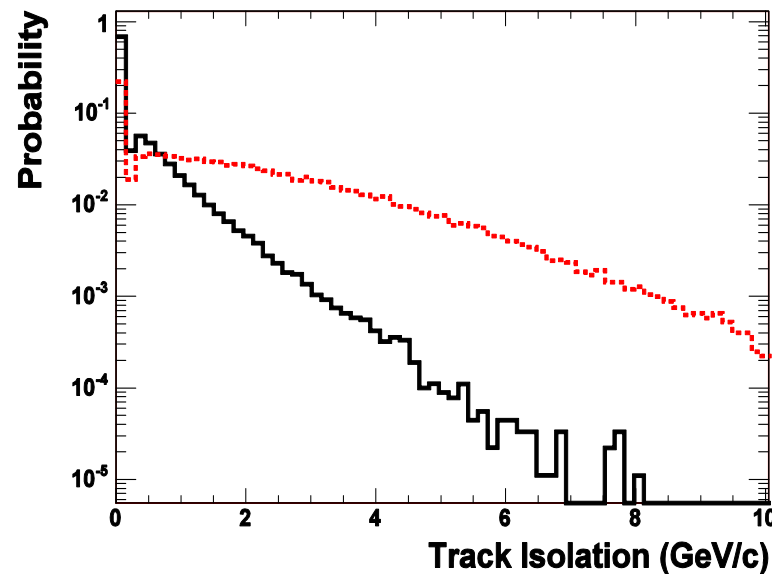
→ Normalization and shape estimated from MC using **track isolation**: $I_{trk} = \sum_{\text{tracks in } R < 0.4}^{z_{vtx} - z_{trk} < 5cm} p_T^{trk}$

→ Sensitive only to underlying event and jet fragmentation (for fake γ), immune to multiple interactions (due to z-cut) and calorimeter leakage

→ Good resolution in low- E_T region, where background is most important

→ Uses charged particles only

Substantially different shape of signal and background I_{trk} distributions can be used to characterize true and fake γ



Background estimation: 4×4 matrix method

- Use the track isolation cut for each photon to compute a per-event weight under the different hypotheses ($\gamma\gamma$, γ +jet and dijet):

$$\begin{pmatrix} w_{jj} \\ w_{j\gamma} \\ w_{\gamma j} \\ w_{\gamma\gamma} \end{pmatrix} = E^{-1} \times \begin{pmatrix} w_{ff} \\ w_{fp} \\ w_{pf} \\ w_{pp} \end{pmatrix} \begin{array}{l} \text{Both photons fail} \\ \text{Leading fail, trailing passes} \\ \text{Leading passes, trailing fails} \\ \text{Both photons pass} \end{array}$$

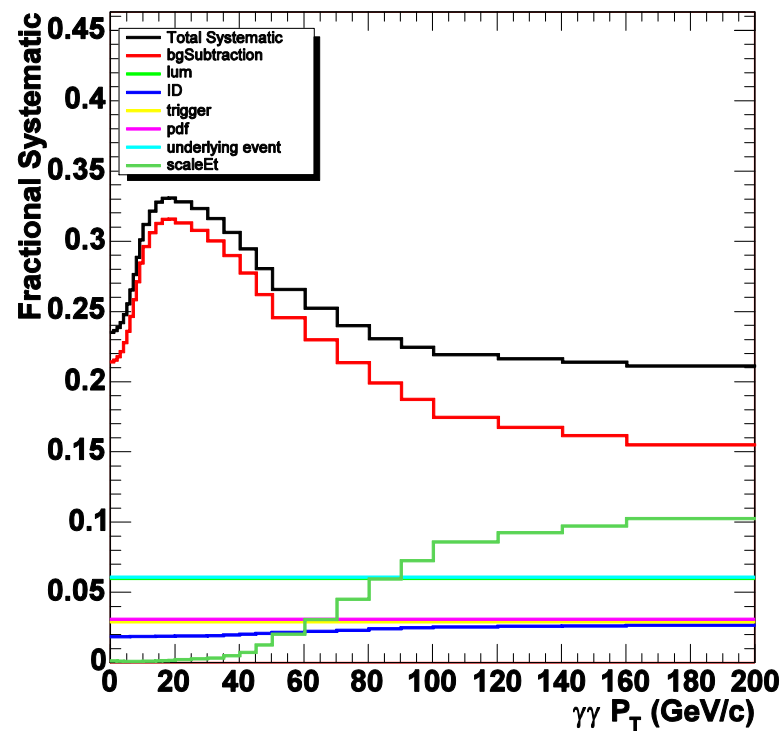
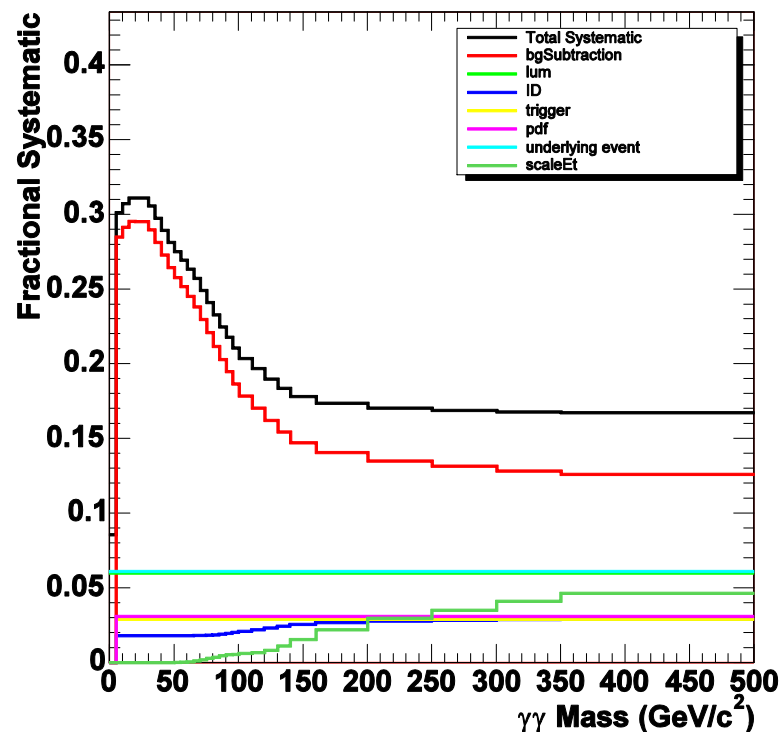
$$E = \begin{pmatrix} (1-\epsilon_{j1})(1-\epsilon_{j2}) & (1-\epsilon_{j1})(1-\epsilon_{\gamma2}) & (1-\epsilon_{\gamma1})(1-\epsilon_{j2}) & (1-\epsilon_{\gamma1})(1-\epsilon_{\gamma2}) \\ (1-\epsilon_{j1})\epsilon_{j2} & (1-\epsilon_{j1})\epsilon_{\gamma2} & (1-\epsilon_{\gamma1})\epsilon_{j2} & (1-\epsilon_{\gamma1})\epsilon_{\gamma2} \\ \epsilon_{j1}(1-\epsilon_{j2}) & \epsilon_{j1}(1-\epsilon_{\gamma2}) & \epsilon_{\gamma1}(1-\epsilon_{j2}) & \epsilon_{\gamma1}(1-\epsilon_{\gamma2}) \\ \epsilon_{j1}\epsilon_{j2} & \epsilon_{j1}\epsilon_{\gamma2} & \epsilon_{\gamma1}\epsilon_{j2} & \epsilon_{\gamma1}\epsilon_{\gamma2} \end{pmatrix}$$

- For instance, if leading passes/trailing fails, the event weight is:
- Estimated number of prompt diphoton events bin-by-bin is given by the sum of $\gamma\gamma$ weights:

$$\begin{pmatrix} w_{ff} \\ w_{fp} \\ w_{pf} \\ w_{pp} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

$$N_{\gamma\gamma} = \sum_{i=1}^{N_{data}} w_{\gamma\gamma}^i$$

Experimental systematic uncertainties



- Total systematic uncertainty ~15-30%, smoothly varying with the kinematic variables considered
- Main source is background subtraction, followed by overall normalization (efficiencies: 7%; integrated luminosity: 6%; UE correction: 6%)

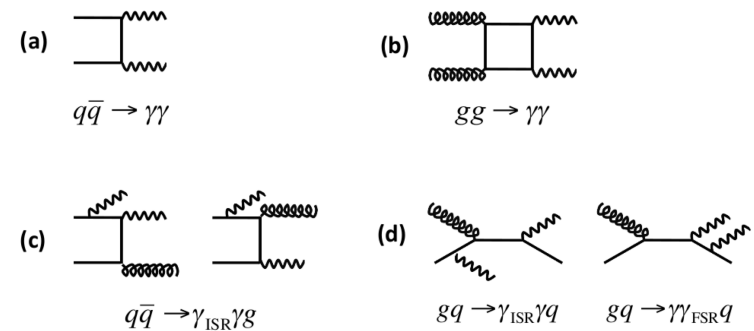
Theoretical predictions

- **DIPHOX**: Fixed-order NLO calculation including non-perturbative fragmentations (T. Binoth *et al.*, Phys. Rev. D **63**, 114016 (2001))
- **RESBOS**: Low- p_T resummed calculation smoothly matched to high- p_T NLO (T. Balazs *et al.*, Phys. Rev. D **76**, 013008 (2007))
- **PYTHIA** 6.2.16 parton-shower calculation (no k-factor applied) (T. Sjöstrand *et al.*, Comp. Phys. Comm. **135**, 238 (2001))

Theoretical predictions

- **DIPHOX**: Fixed-order NLO calculation including non-perturbative fragmentations (T. Binoth *et al.*, Phys. Rev. D **63**, 114016 (2001))
- **RESBOS**: Low- p_T resummed calculation smoothly matched to high- p_T NLO (T. Balazs *et al.*, Phys. Rev. D **76**, 013008 (2007))
- **PYTHIA** 6.2.16 parton-shower calculation (no k-factor applied) (T. Sjöstrand *et al.*, Comp. Phys. Comm. **135**, 238 (2001))

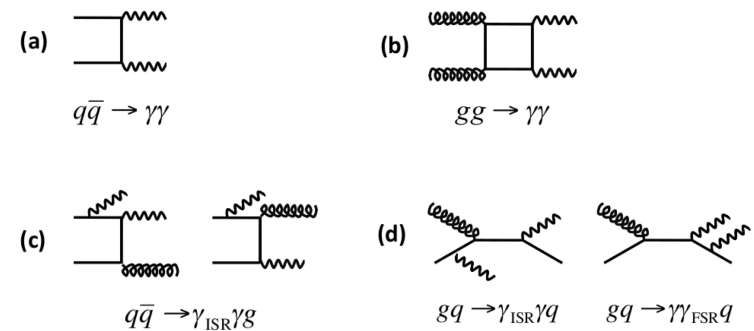
Two separate calculations, one involving (a – b) only (“PYTHIA $\gamma\gamma$ ”) and one involving (a – d) (“PYTHIA $\gamma\gamma+\gamma j$ ”), are compared with the data



Theoretical predictions

- **DIPHOX**: Fixed-order NLO calculation including non-perturbative fragmentations (T. Binoth *et al.*, Phys. Rev. D **63**, 114016 (2001))
- **RESBOS**: Low- p_T resummed calculation smoothly matched to high- p_T NLO (T. Balazs *et al.*, Phys. Rev. D **76**, 013008 (2007))
- **PYTHIA** 6.2.16 parton-shower calculation (no k-factor applied) (T. Sjöstrand *et al.*, Comp. Phys. Comm. **135**, 238 (2001))

Two separate calculations, one involving (a – b) only (“PYTHIA $\gamma\gamma$ ”) and one involving (a – d) (“PYTHIA $\gamma\gamma+\gamma j$ ”), are compared with the data



- NLO theoretical uncertainties:
 - PDFs: 3-6%; use 44 eigenvectors from CTE6.1M
 - Renormalization/factorization/fragmentation scales: ~10-20% depending on the observable; all scales simultaneously varied by $\times 2$ up and down

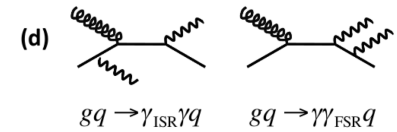
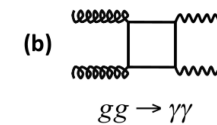
Theoretical predictions

- **DIPHOX**: Fixed-order NLO calculation including non-perturbative fragmentations (T. Binoth *et al.*, Phys. Rev. D **63**,114016 (2001))

- **RESBOS**: Low- p_T resummed calculation smoothly matched to high- p_T NLO (T. Balazs *et al.*, Phys. Rev. D **76**, 013008 (2007))

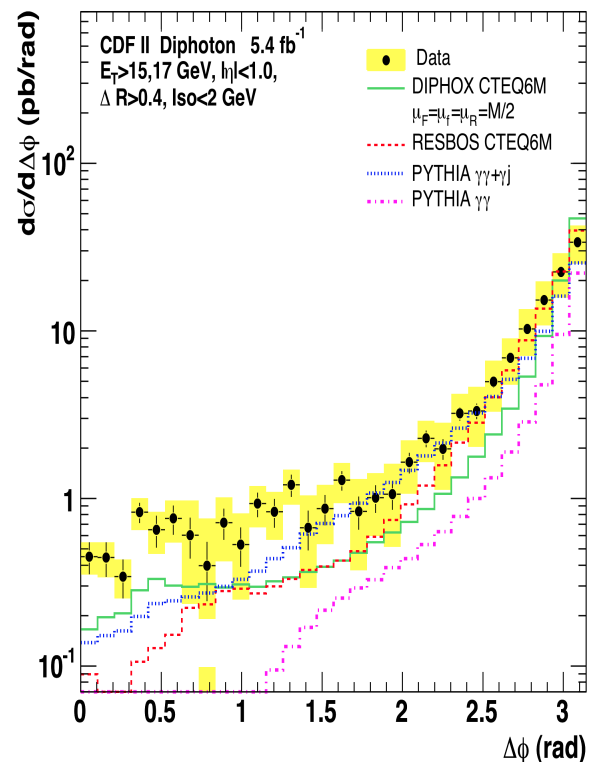
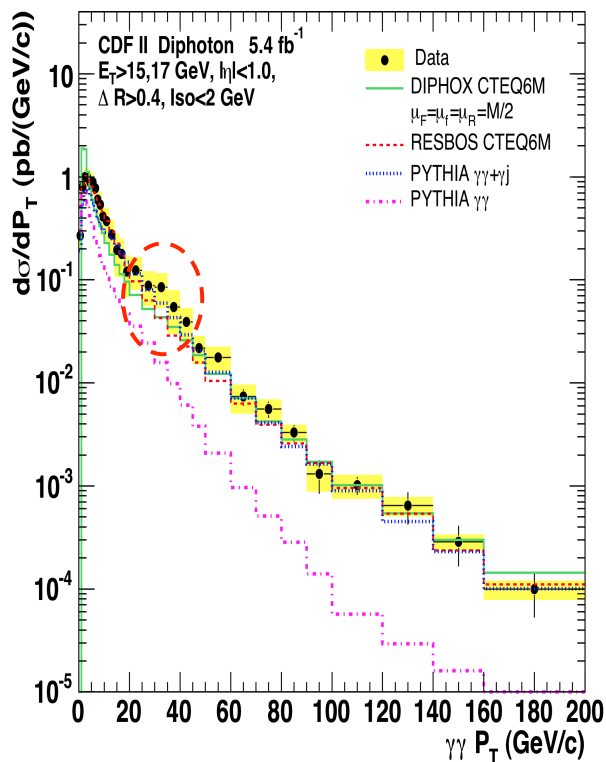
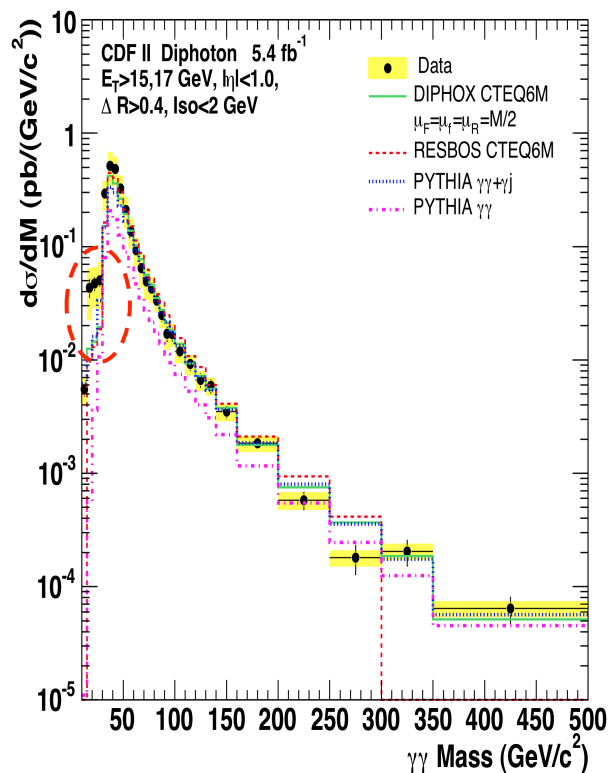
- **PYTHIA** 6.2.16 p8
(T.Sjöstrand *et al.*,
Two separate calc
(a – b) only (“PYTHIA
(a – d) (“PYTHIA
the data

	Total cross section (pb)
Data	$12.5 \pm 0.2_{\text{stat}} \pm 3.7_{\text{syst}}$
RESBOS	$11.3 \pm 2.4_{\text{syst}}$
DIPHOX	$10.6 \pm 0.6_{\text{syst}}$
PYTHIA $\gamma\gamma + \gamma j$	9.2
PYTHIA $\gamma\gamma$	5.0



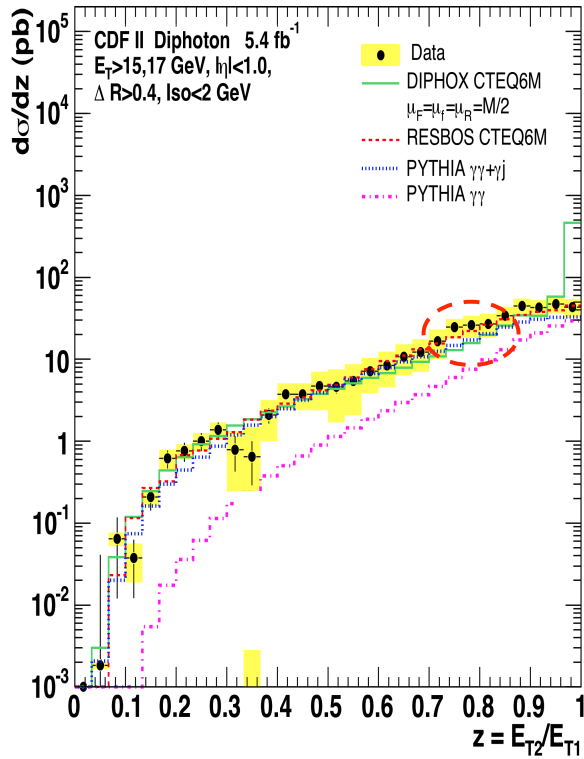
- NLO theoretical uncertainties:
 - PDFs: 3-6%; use 44 eigenvectors from CTE6.1M
 - Renormalization/factorization/fragmentation scales: ~10-20% depending on the observable; all scales simultaneously varied by $\times 2$ up and down

Differential cross sections

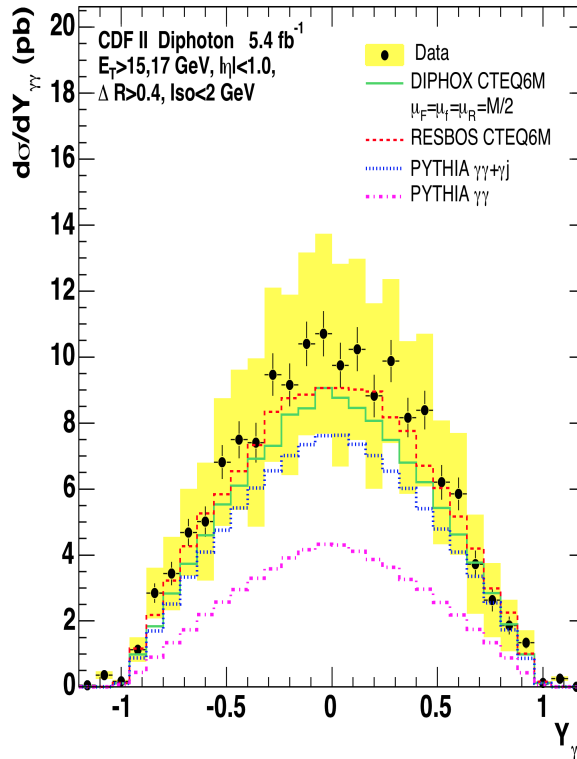


- Good agreement between data and theory for $M_{\gamma\gamma} > 30$ GeV/c²
- Resummation important for $p_T(\gamma\gamma) > 20$ GeV/c
- Fragmentations cause excess of data over theory for $p_T(\gamma\gamma) = 20 - 50$ GeV/c
- Resummation important for $\Delta\phi_{\gamma\gamma} > 2.2$ rad
- Data spectrum harder than predicted

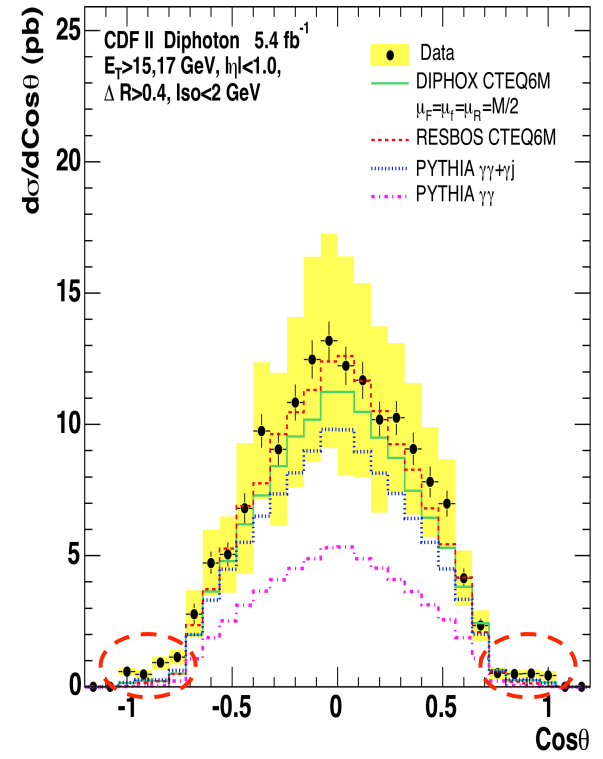
Differential cross sections



- Good agreement between data and RESBOS
- Good agreement between data and DIPHOX, except for $0.7 < z < 0.8$

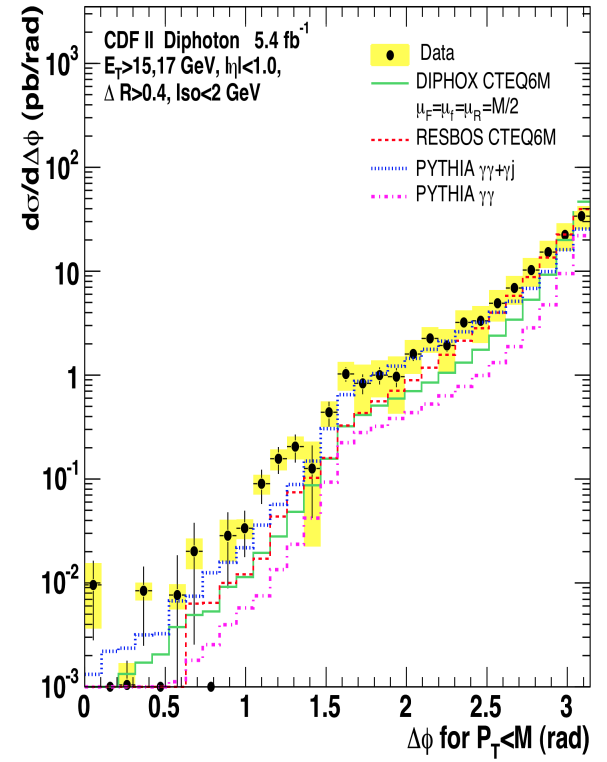
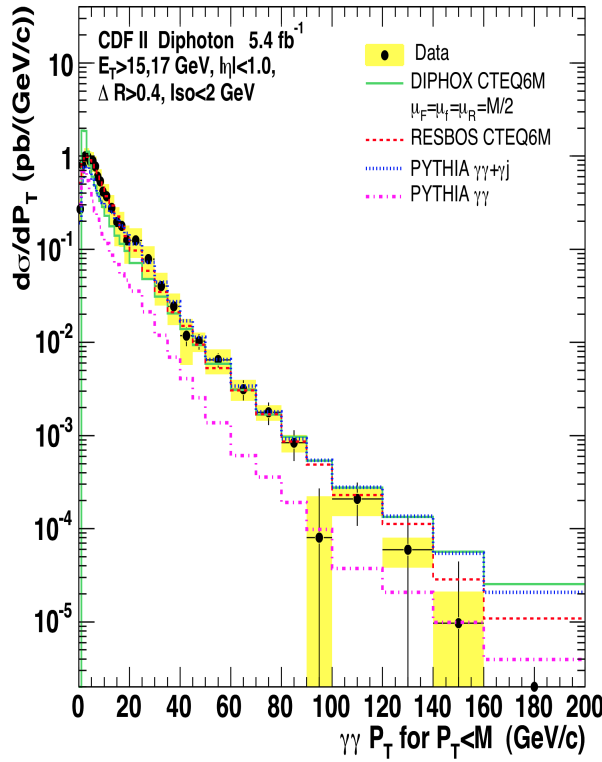
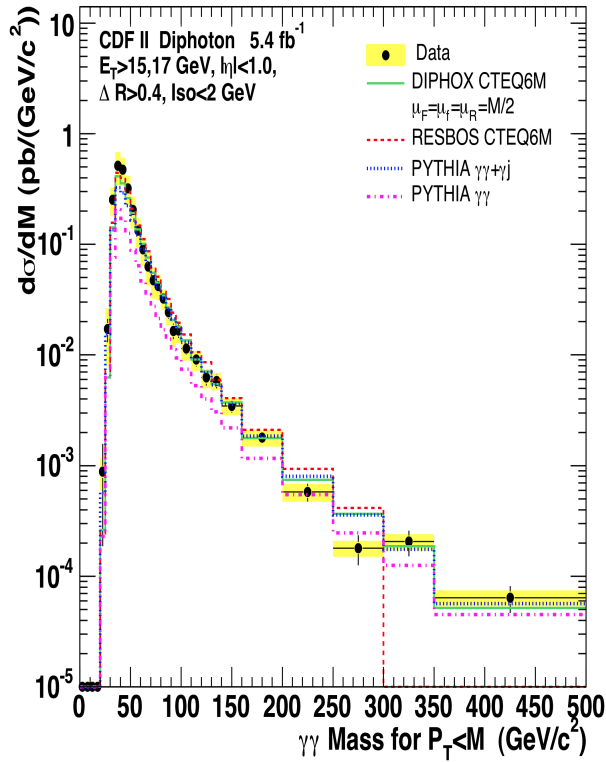


- Good agreement between data and theory



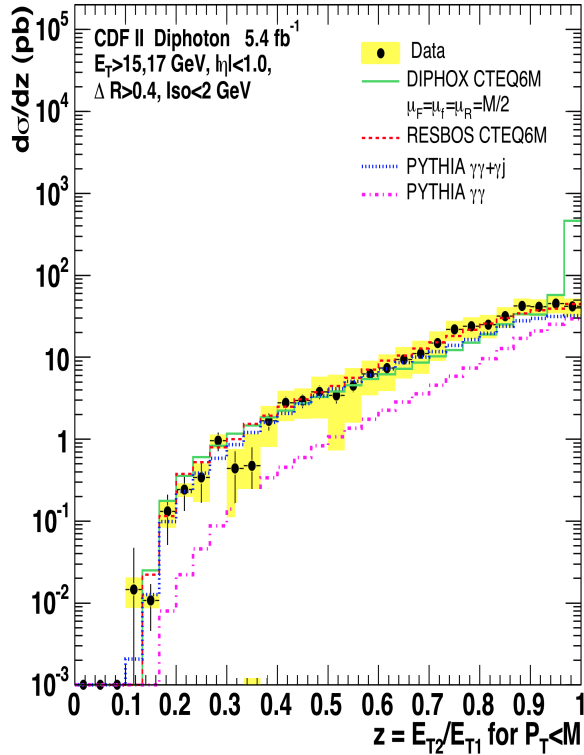
- Observable sensitive to PDFs
- Good agreement between data and theory, except for $|\cos\theta^*| \rightarrow 1$

Differential cross sections for $p_T(\gamma\gamma) < M(\gamma\gamma)$

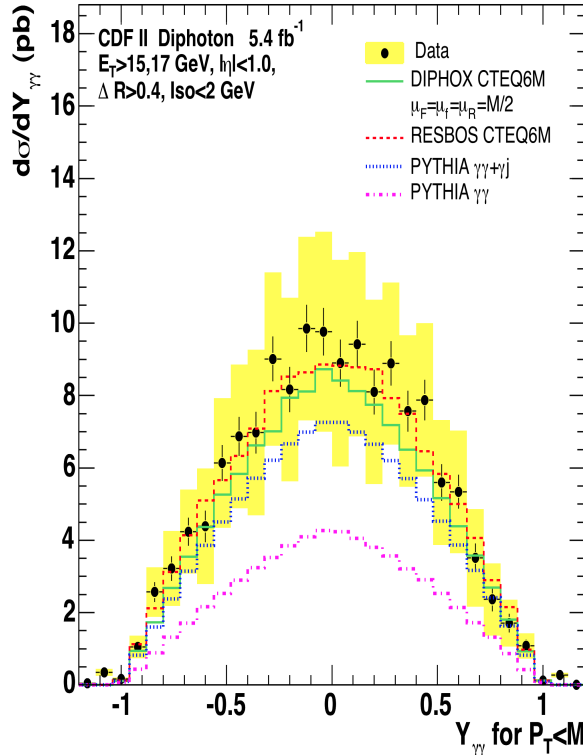


- Good agreement between data and theory
- “Shoulder” in data for $p_T(\gamma\gamma) = 20 - 50$ GeV/c significantly reduced
- Discrepancies between data and theory for $\Delta\phi_{\gamma\gamma} < 1.7$ rad reduced

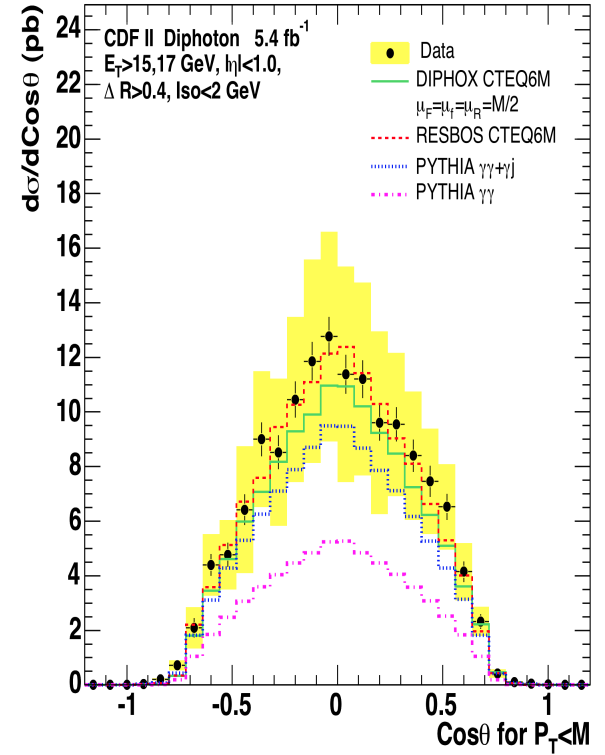
Differential cross sections for $p_T(\gamma\gamma) < M(\gamma\gamma)$



- Good agreement between data and RESBOS
- Good agreement between data and DIPHOX, except for $0.7 < z < 0.8$



- Good agreement between data and theory



- Good agreement between data and theory

Summary and conclusions

- Reported measurements of differential cross sections for direct diphoton production at $\sqrt{s}=1.96$ TeV using 5.4 fb^{-1} .

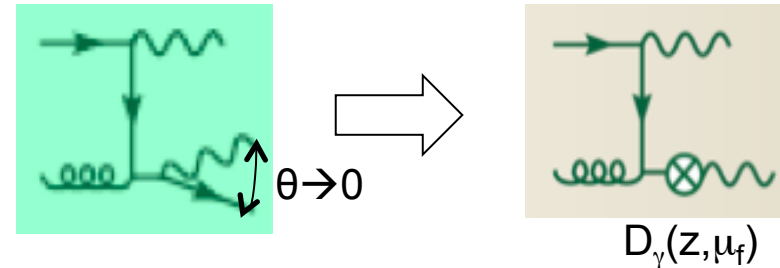
$$\frac{d\sigma}{dM_{\gamma\gamma}} \quad \frac{d\sigma}{dp_T^{\gamma\gamma}} \quad \frac{d\sigma}{d\Delta\phi_{\gamma\gamma}} \quad \frac{d\sigma}{dz} \quad \frac{d\sigma}{dy_{\gamma\gamma}} \quad \frac{d\sigma}{d\cos\theta^*}$$

- Measurements are compared to state-of-art theoretical predictions such as **DIPHOX**, **RESBOS**, and **PYTHIA**. Overall agreement between data and theory, within known limitations, is observed.
- Resummation** matched with NLO pQCD calculations works well at low $p_T(\gamma\gamma)$ ($<20 \text{ GeV}/c$) and large $\Delta\phi_{\gamma\gamma}$ ($>2.2 \text{ rad}$).
- Fragmentations** appear to be not under good control in sensitive kinematic regions [$M(\gamma\gamma) < 60 \text{ GeV}/c^2$, $20 \text{ GeV}/c < p_T(\gamma\gamma) < 50 \text{ GeV}/c$, $\Delta\phi_{\gamma\gamma} < 1 \text{ rad}$].
- Data-to-theory comparisons show best agreement for $p_T(\gamma\gamma) < M(\gamma\gamma)$, where theoretical uncertainties are smaller and predictions are less sensitive to the isolation requirement.
- Parton-shower PYTHIA** Monte Carlo, which in previous analyses limited to matrix-element-based simulations was found to fail reproducing the data, now provides a description of the data competitive with full NLO calculations by including **ISR and FSR photons**
- A PRL ([arXiv:1106.5123](https://arxiv.org/abs/1106.5123)) and a PRD ([arXiv:1106.5131](https://arxiv.org/abs/1106.5131)) have been submitted

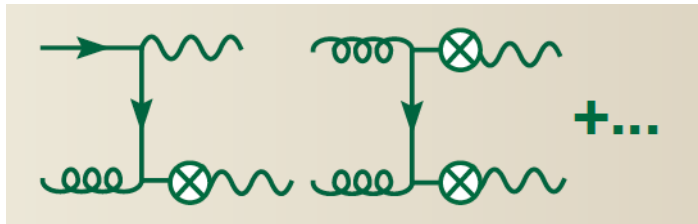
Backup slides

Fragmentation contributions

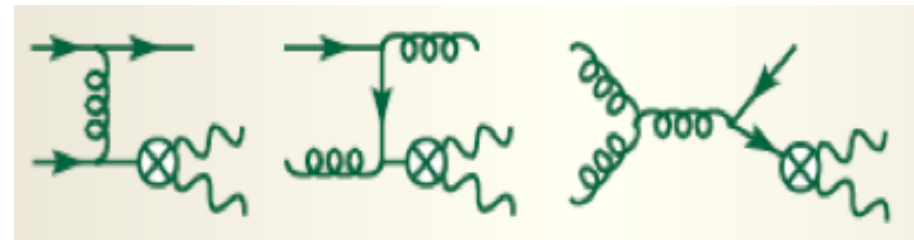
- Collinear singularity in final state photon radiation off a parton can be handled e.g. via fragmentation functions.



Single-photon fragmentation



Double-photon fragmentation



Low-mass/small-angle diphoton pairs



Not included in any theoretical prediction!

- Fragmentation contributions can be suppressed via:
 - experimental photon isolation requirements (can only be approximated in theory)
 - $p_T(\gamma\gamma) < M(\gamma\gamma)$

$$E_T^{iso} = \sum_{\substack{\text{partons or hadrons} \\ \text{within } R < 0.4}} E_T - E_{T\gamma}$$

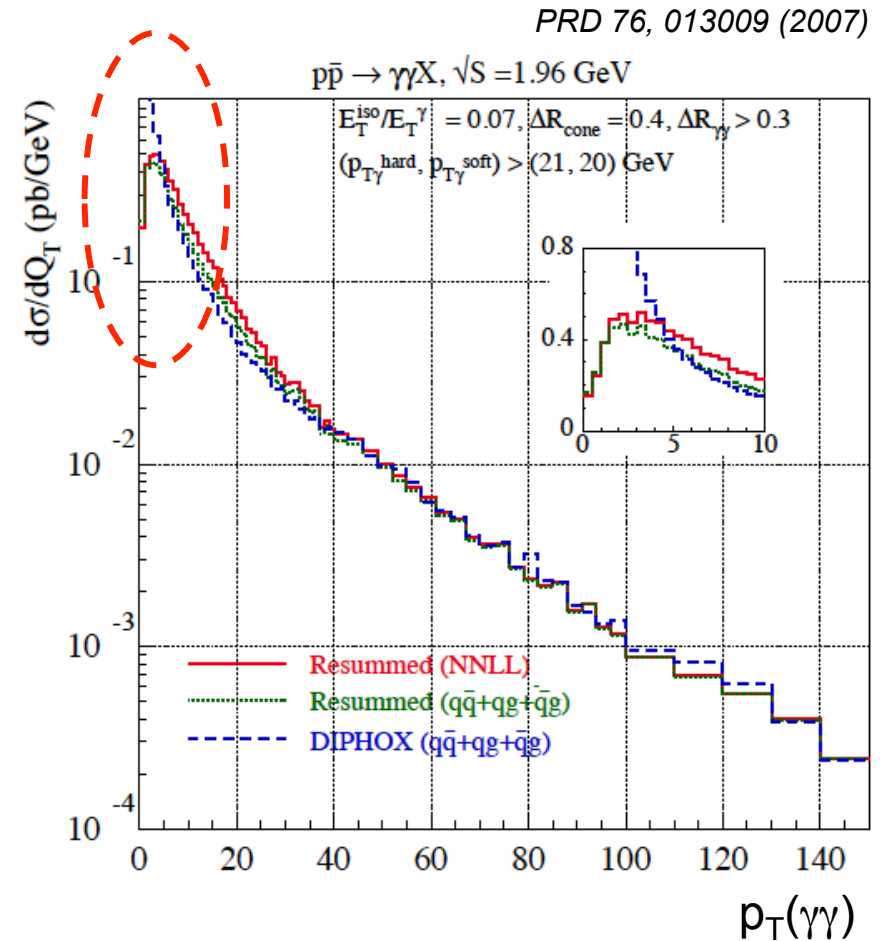
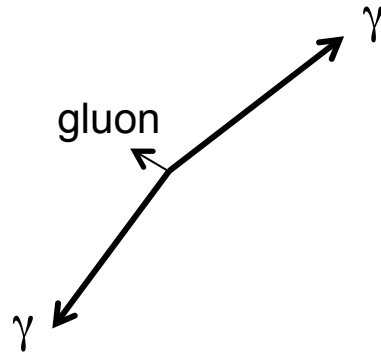
Resummation of initial state gluons

- At fixed $M(\gamma\gamma)$, the differential cross section as a function of $p_T(\gamma\gamma)$ at $O(\alpha_s)$ given by:

$$\frac{d\sigma}{dp_{T\gamma\gamma}^2} = \sigma_0 \frac{\alpha_s}{\pi} \frac{1}{p_{T\gamma\gamma}^2} \left[a \ln \left(\frac{M_{\gamma\gamma}^2}{p_{T\gamma\gamma}^2} \right) + a_0 \right]$$

Fixed-order calculation less reliable for $p_T(\gamma\gamma) \ll M(\gamma\gamma)$ and diverges as $p_T(\gamma\gamma) \rightarrow 0$.

[Also when $\Delta\phi(\gamma,\gamma) \rightarrow \pi$.]

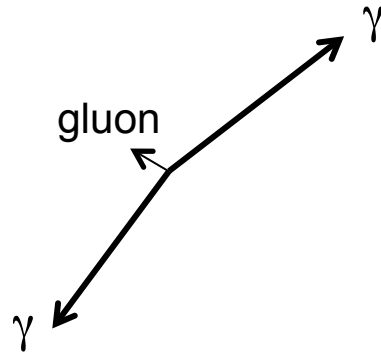


Resummation of initial state gluons

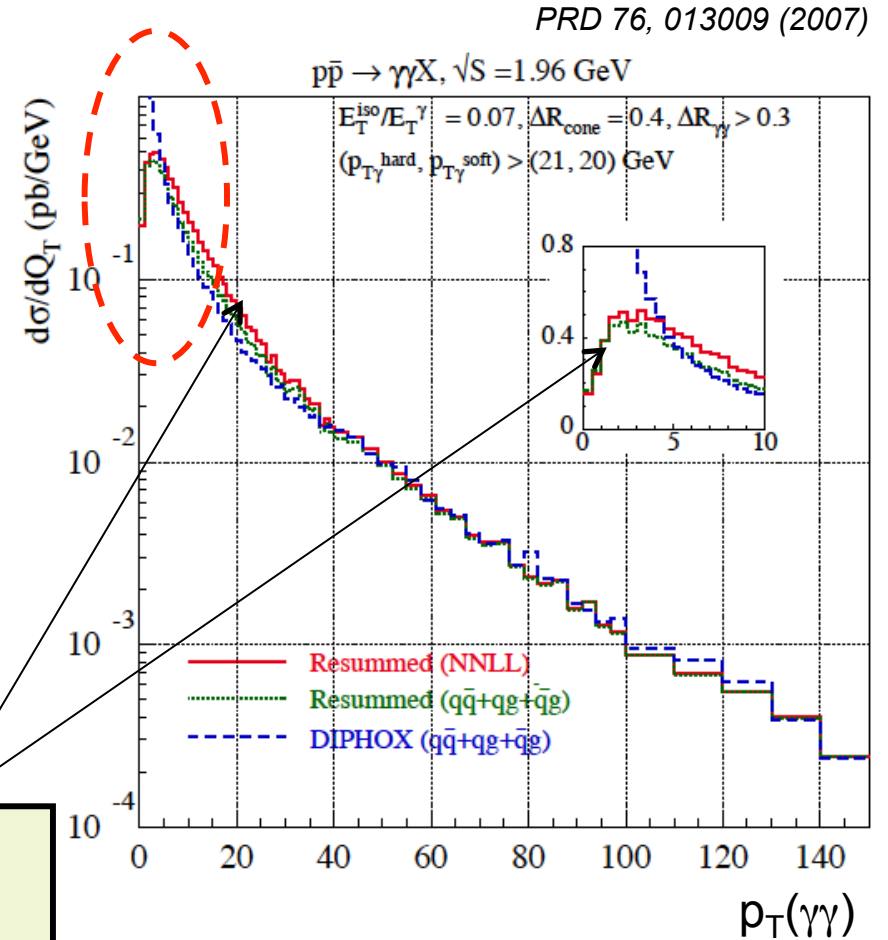
- At fixed $M(\gamma\gamma)$, the differential cross section as a function of $p_T(\gamma\gamma)$ is given by:

$$\frac{d\sigma}{dp_{T\gamma\gamma}^2} = \sigma_0 \frac{\alpha_s}{\pi} \frac{1}{p_{T\gamma\gamma}^2} \left[a \ln \left(\frac{M_{\gamma\gamma}^2}{p_{T\gamma\gamma}^2} \right) + a_0 \right]$$

Fixed-order calculation less reliable for $p_T(\gamma\gamma) \ll M(\gamma\gamma)$ and diverges as $p_T(\gamma\gamma) \rightarrow 0$.
[Also when $\Delta\phi(\gamma,\gamma) \rightarrow \pi$.]



Physical description of the $p_T(\gamma\gamma)$ and $\Delta\phi(\gamma,\gamma)$ distributions requires all-order resummation of soft and collinear logarithms.



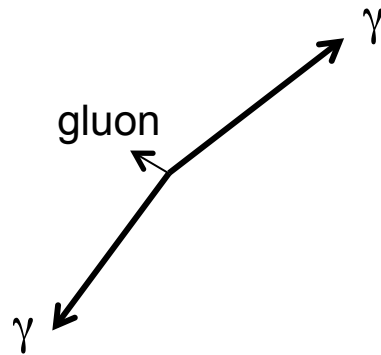
Resummation of initial state gluons

- At fixed $M(\gamma\gamma)$, the differential cross section as a function of $p_T(\gamma\gamma)$ is given by:

$$\frac{d\sigma}{dp_{T\gamma\gamma}^2} = \sigma_0 \frac{\alpha_s}{\pi} \frac{1}{p_{T\gamma\gamma}^2} \left[a \ln \left(\frac{M_{\gamma\gamma}^2}{p_{T\gamma\gamma}^2} \right) + a_0 \right]$$

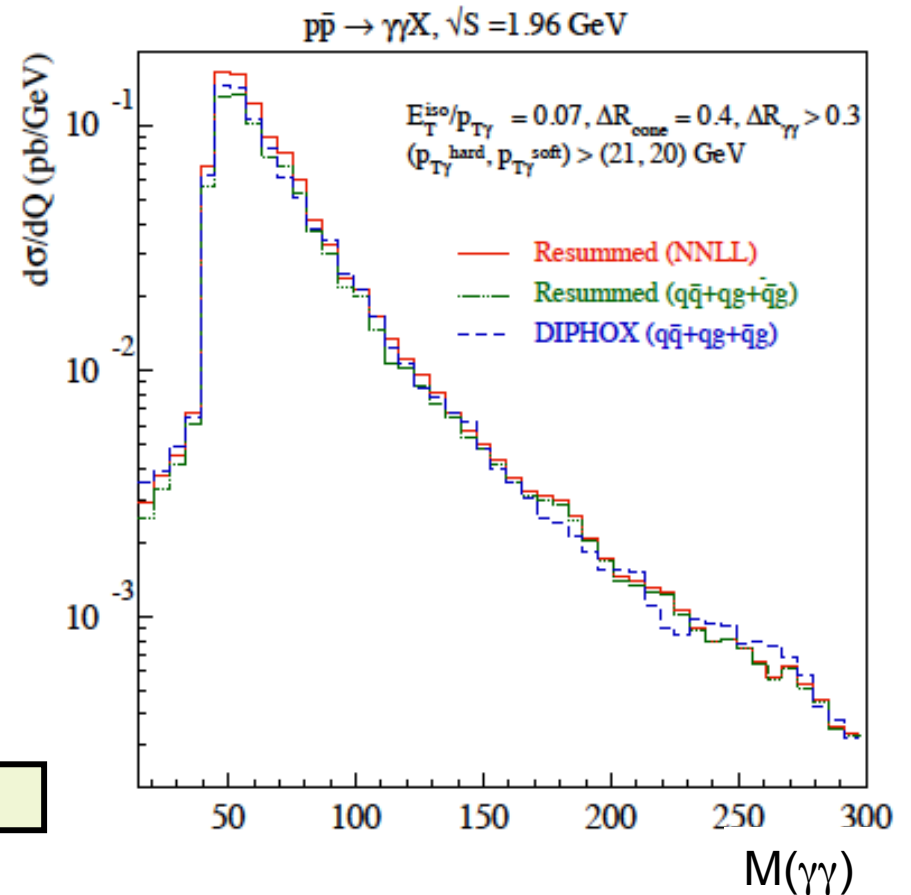
Fixed-order calculation less reliable for $p_T(\gamma\gamma) \ll M(\gamma\gamma)$ and diverges as $p_T(\gamma\gamma) \rightarrow 0$.

[Also when $\Delta\phi(\gamma, \gamma) \rightarrow \pi$.]

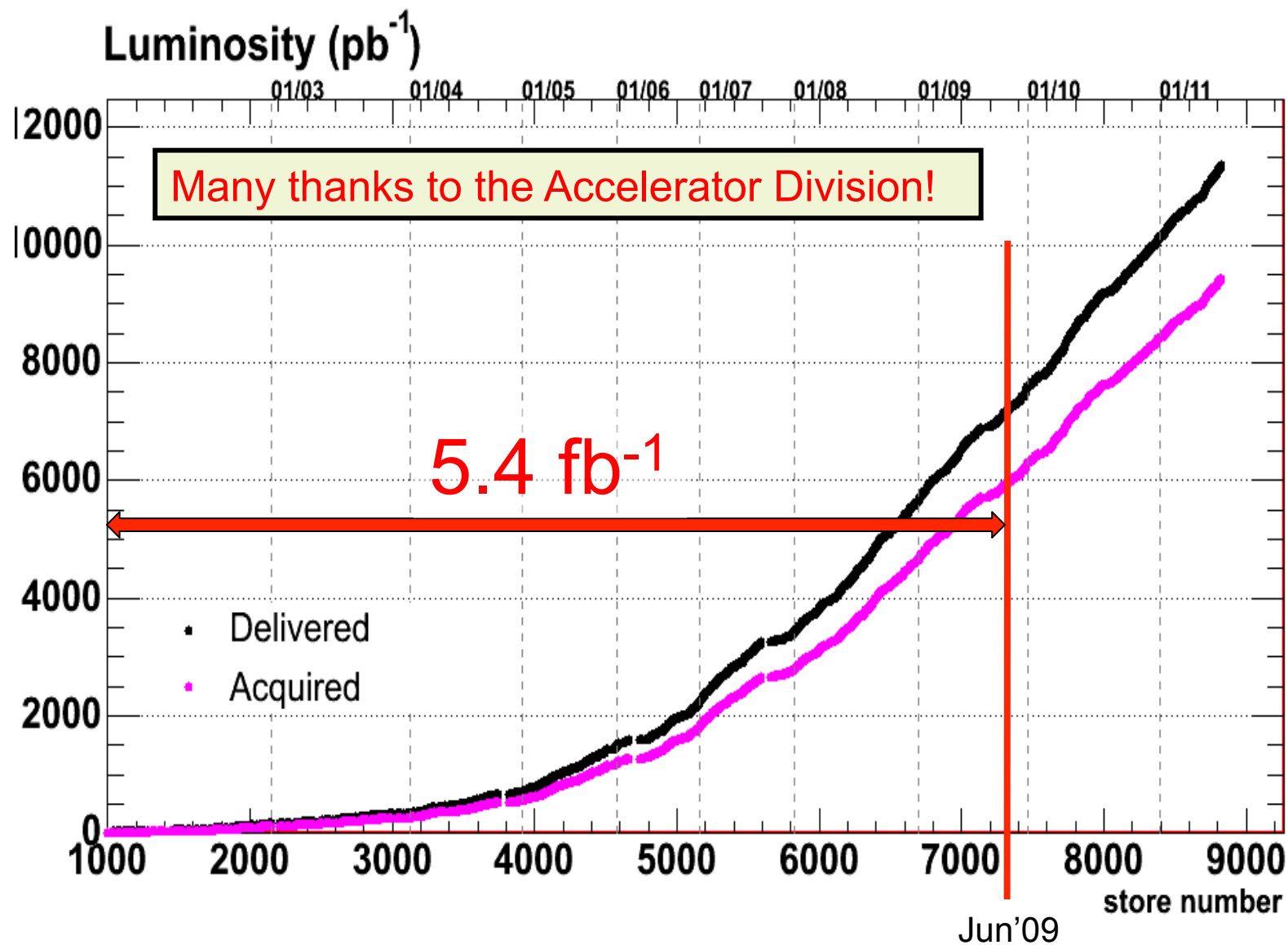


Only small effect on $M(\gamma\gamma)$ from resummation

PRD 76, 013009 (2007)



Data set



x27 more luminosity than previous CDF publication!

Triggers

Diphoton-12

- **L1:**
 - EM $E_T \geq 8$ GeV
 - $E_{\text{HAD}}/E_{\text{EM}} \leq 0.125$
 - $N_{\text{cluster}} = 2$
- **L2:**
 - EM $E_T \geq 10$ GeV
 - $E_{\text{HAD}}/E_{\text{EM}} \leq 0.125$
 - $N_{\text{cluster}} = 2$
 - Isolation ≤ 3 GeV
or IsoFraction ≤ 0.15
- **L3:**
 - EM $E_T \geq 12$ GeV
 - $E_{\text{HAD}}/E_{\text{EM}} \leq 0.055 + 0.00045 \times E_{\text{tot}}/\text{GeV}$
 - $N_{\text{cluster}} = 2$
 - Isolation ≤ 2 GeV
or IsoFraction ≤ 0.1
 - Shower profile: $\chi^2_{\text{CES}} \leq 20$

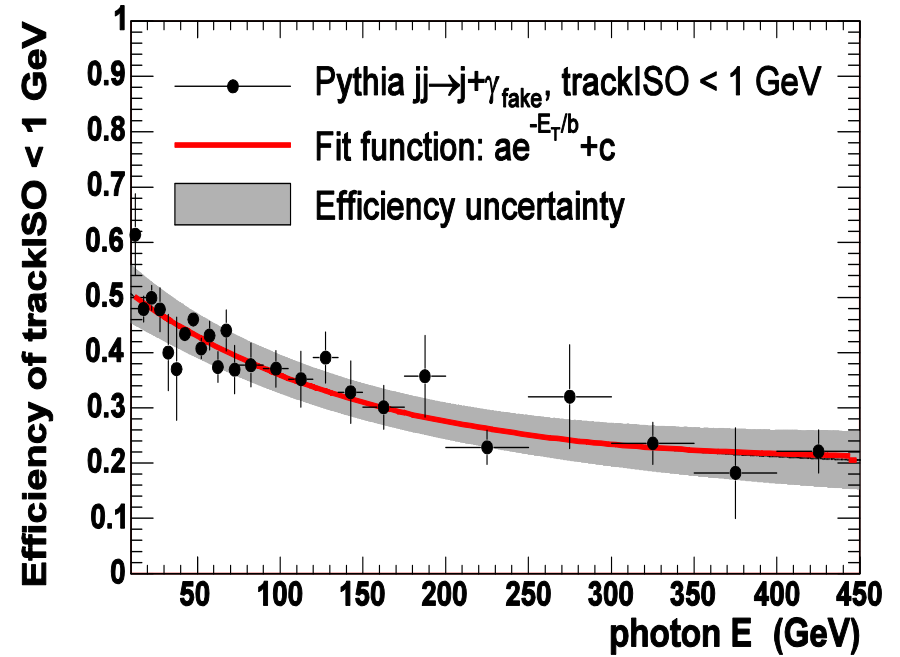
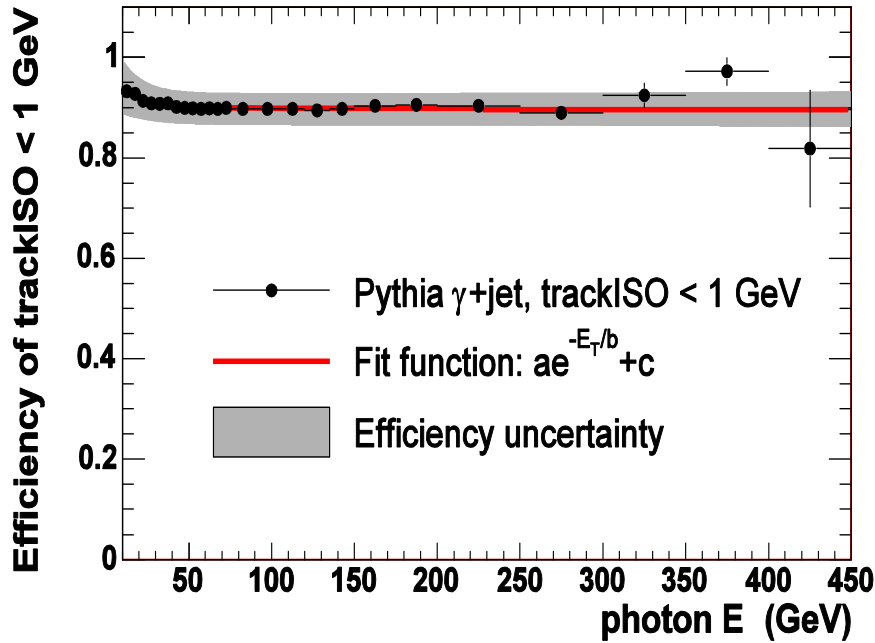
“OR”

Diphoton-18 Same as diphoton-12 except:

- **L2:**
 - EM $E_T \geq 16$ GeV
 - No isolation
- **L3:**
 - EM $E_T \geq 18$ GeV
 - No isolation

Trigger efficiency after offline event selection: 100% for $E_T \geq 15$ GeV

Photon characterization using track isolation



For a single γ , a weight can be defined to characterize it as signal or background:

$$\mathcal{W} = \frac{\mathcal{E} - \mathcal{E}_b}{\mathcal{E}_s - \mathcal{E}_b}$$

→ $\mathcal{E} = 1$ (0) if $I_{\text{trk}} < (\geq) 1 \text{ GeV}/c$

→ \mathcal{E}_s = signal efficiency for $I_{\text{trk}} < 1 \text{ GeV}/c$

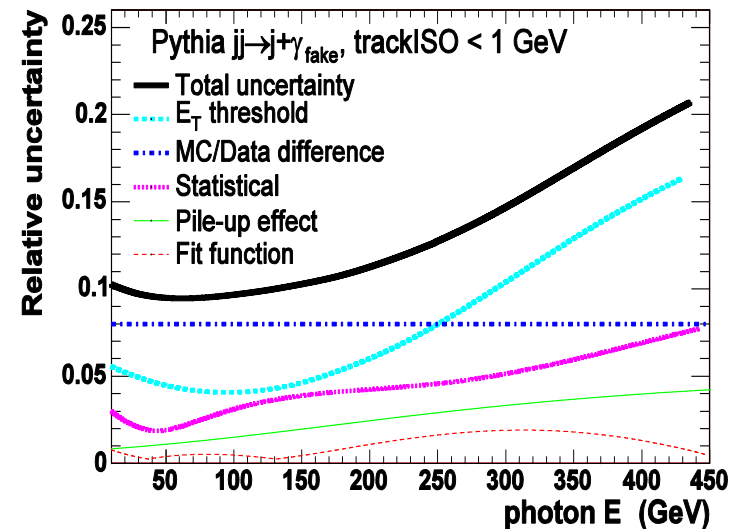
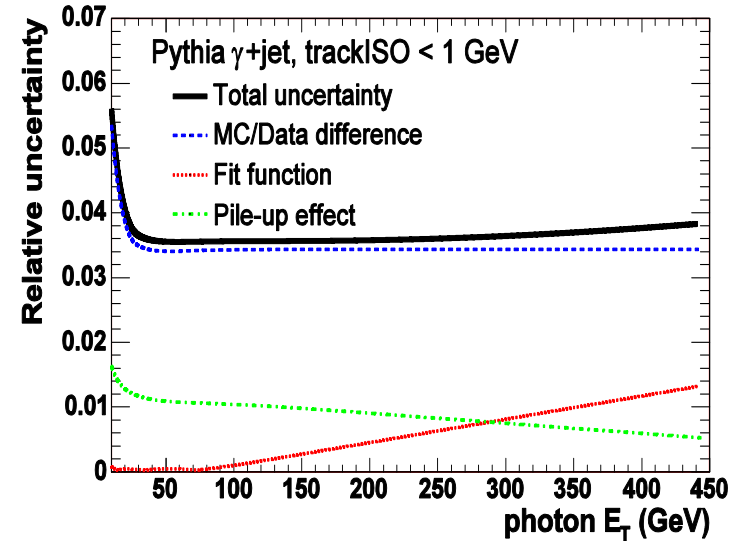
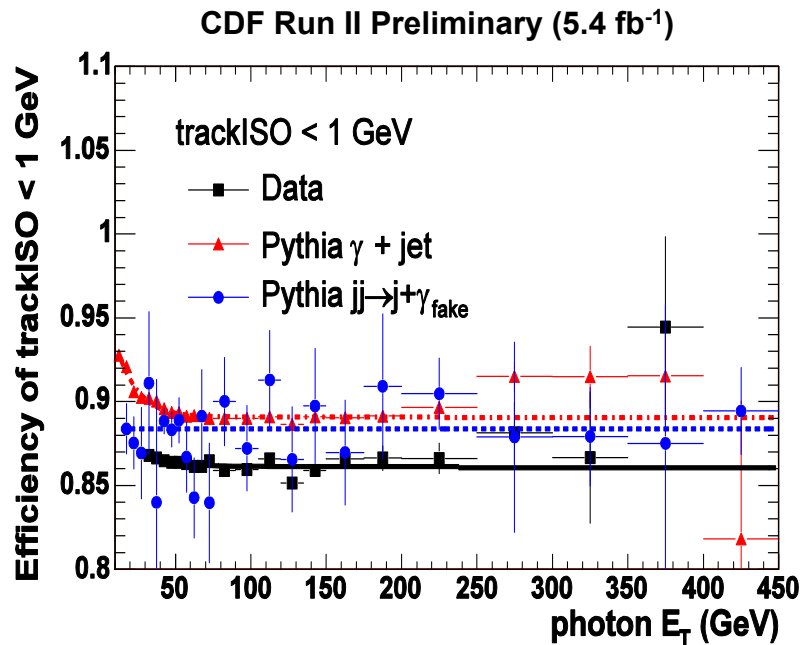
→ \mathcal{E}_b = background efficiency for $I_{\text{trk}} < 1 \text{ GeV}/c$

Both modeled by $ae^{-E_T/b} + c$

Cut chosen at $I_{\text{trk}} = 1 \text{ GeV}/c$, where $\mathcal{E}_s - \mathcal{E}_b = \text{max}$, to optimize resolution

Background estimation: 4×4 matrix method

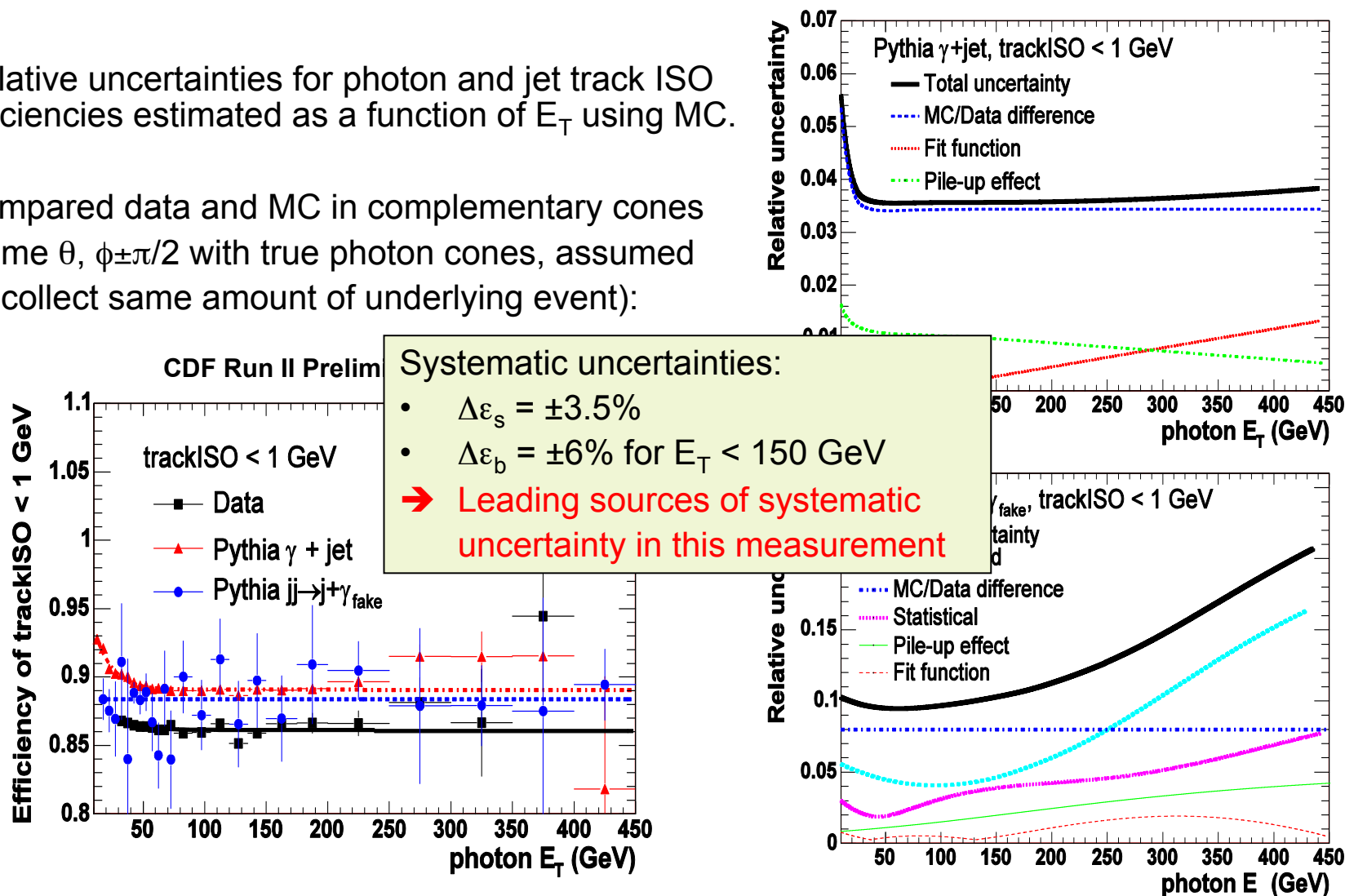
- Relative uncertainties for photon and jet track ISO efficiencies estimated as a function of E_T using MC.
- Compared data and MC in complementary cones (same θ , $\phi \pm \pi/2$ with true photon cones, assumed to collect same amount of underlying event):



→ Data and MC consistent to within 3%.

Background estimation: 4x4 matrix method

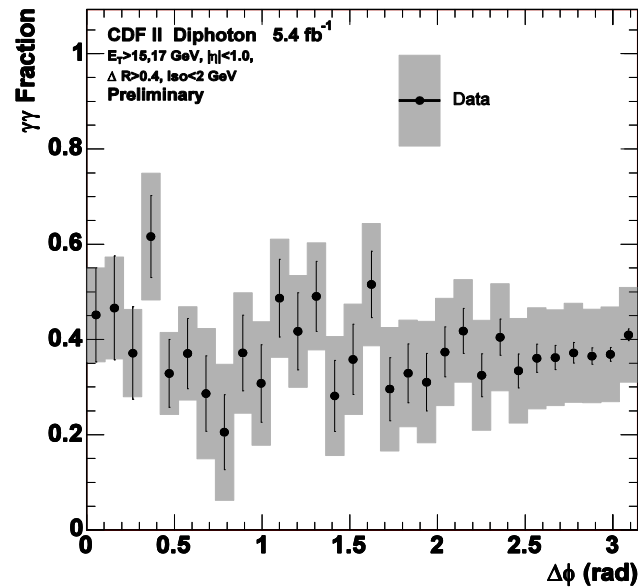
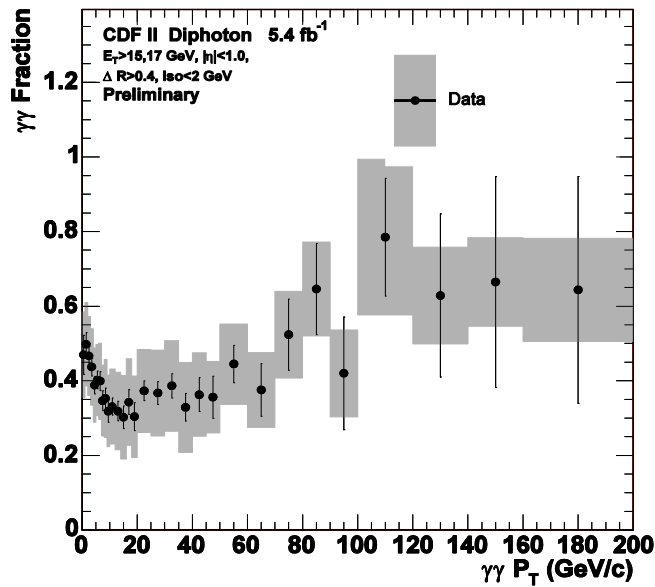
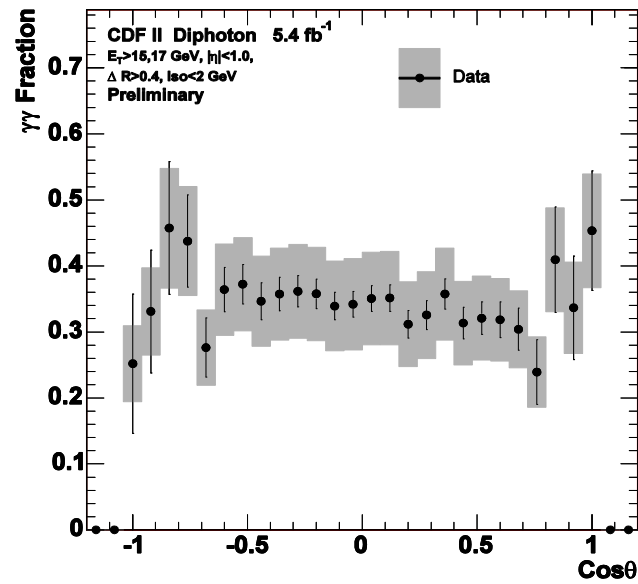
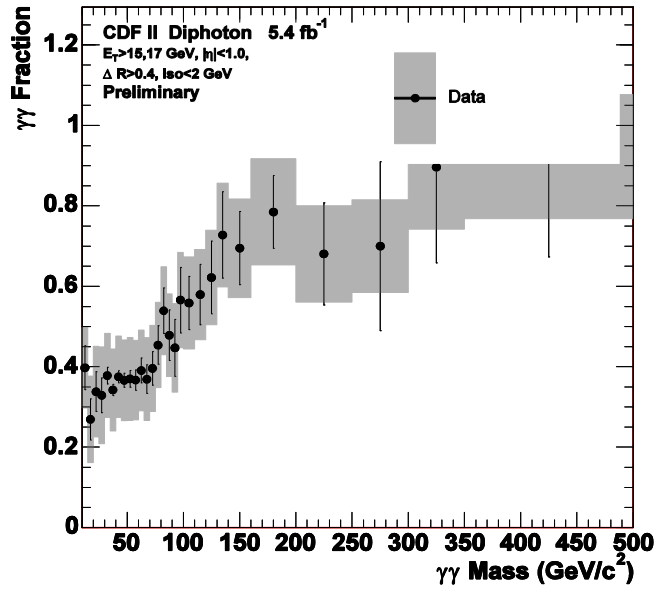
- Relative uncertainties for photon and jet track ISO efficiencies estimated as a function of E_T using MC.
- Compared data and MC in complementary cones (same θ , $\phi \pm \pi/2$ with true photon cones, assumed to collect same amount of underlying event):



→ Data and MC consistent to within 3%.

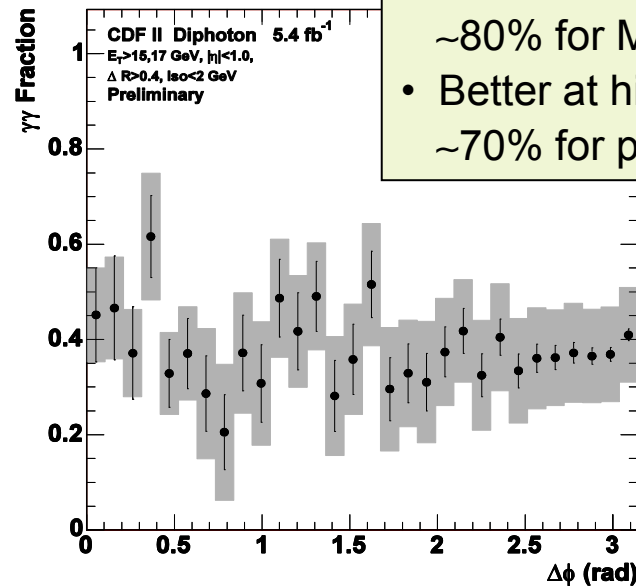
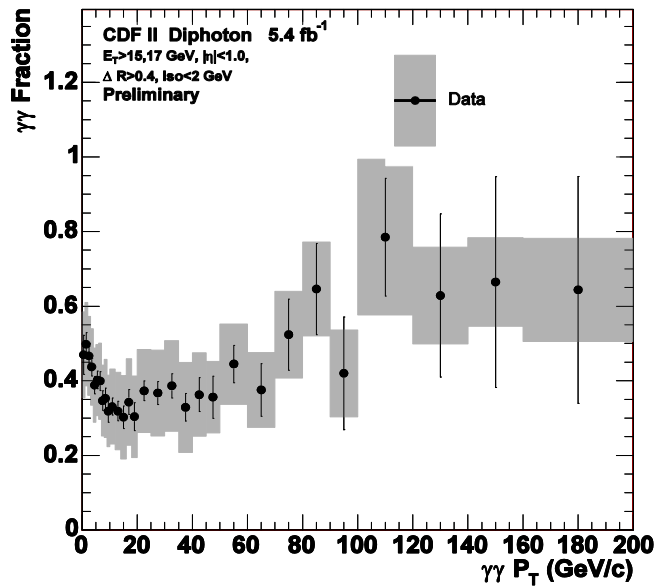
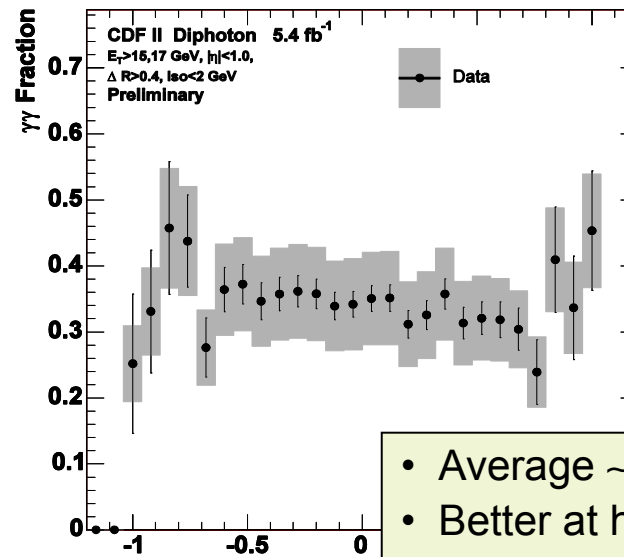
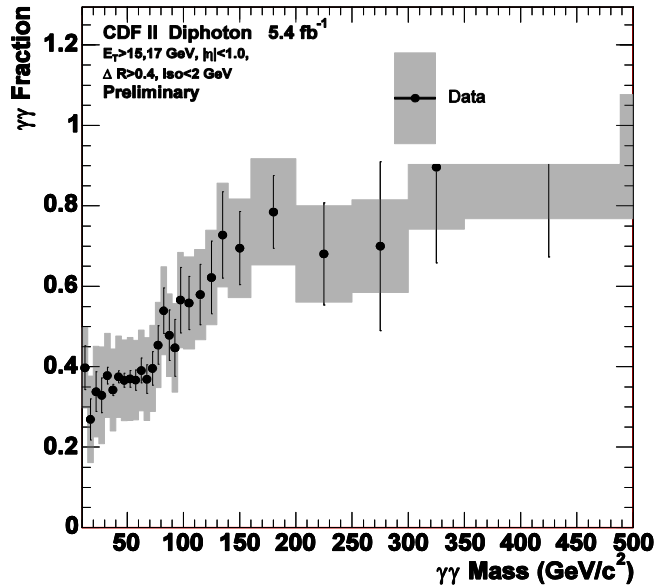
Signal fraction

$$\text{Signal fraction} = \frac{N_{\gamma\gamma}}{N_{data}}$$



Signal fraction

$$\text{Signal fraction} = \frac{N_{\gamma\gamma}}{N_{data}}$$



- Average ~40%
- Better at high mass:
 60-80% for $M(\gamma\gamma) \sim 80-150 \text{ GeV}/c^2$
 ~80% for $M(\gamma\gamma) > 150 \text{ GeV}/c^2$
- Better at high p_T(γγ):
 ~70% for p_T(γγ) > 100 GeV/c

Acceptance × efficiency

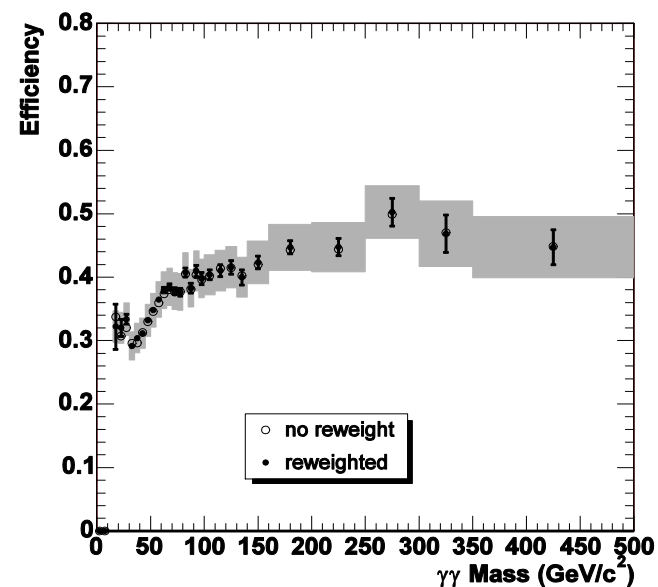
$$\frac{d\sigma}{dX} = \frac{N_{\gamma\gamma}}{\varepsilon \cdot A \cdot L \cdot \Delta}$$

- Defined as:

Number of events with two reconstructed EM clusters passing all cuts

Number of events with two generator-level photons passing kinematic and isolation cuts

- Estimated using detector- and trigger-simulated and reconstructed PYTHIA events reweighted to match the data



Acceptance × efficiency

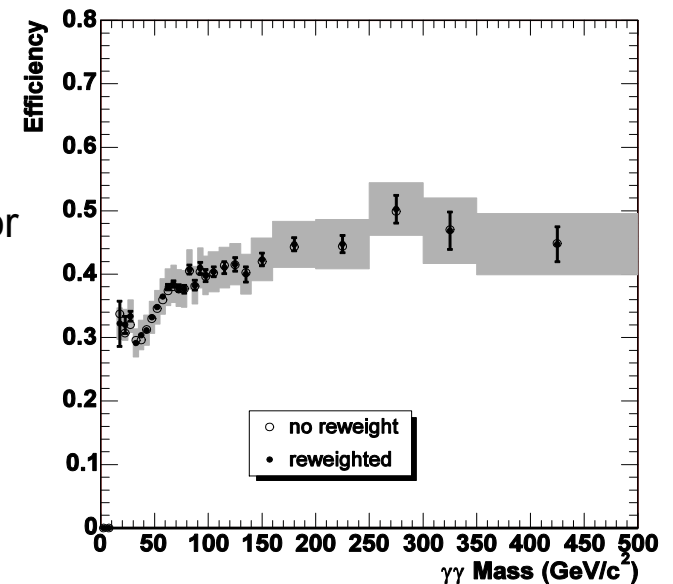
$$\frac{d\sigma}{dX} = \frac{N_{\gamma\gamma}}{\varepsilon \cdot A \cdot L \cdot \Delta}$$

- Defined as:

Number of events with two reconstructed EM clusters passing all cuts

Number of events with two generator-level photons passing kinematic and isolation cuts

- Estimated using detector- and trigger-simulated and reconstructed PYTHIA events reweighted to match the data
- RESBOS and DIPHOX do not include non-perturbative effects: underlying event and hadronization
 → lower efficiency of the isolation cut relative to PYTHIA
 (PYTHIA events are removed from the isolated denominator of the efficiency due to the underlying event)
- Correction estimated by convoluting PYTHIA UE isolation energy with DIPHOX energy in the isolation cone
 → constant per event factor of 0.88 applied to the data



Acceptance × efficiency

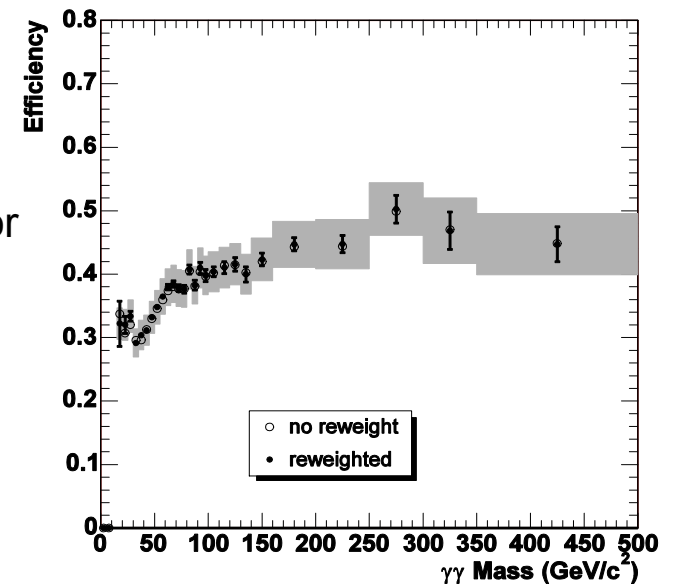
$$\frac{d\sigma}{dX} = \frac{N_{\gamma\gamma}}{\varepsilon \cdot A \cdot L \cdot \Delta}$$

- Defined as:

Number of events with two reconstructed EM clusters passing all cuts

Number of events with two generator-level photons passing kinematic and isolation cuts

- Estimated using detector- and trigger-simulated and reconstructed PYTHIA events reweighted to match the data
- RESBOS and DIPHOX do not include non-perturbative effects: underlying event and hadronization
 → lower efficiency of the isolation cut relative to PYTHIA
 (PYTHIA events are removed from the isolated denominator of the efficiency due to the underlying event)
- Correction estimated by convoluting PYTHIA UE isolation energy with DIPHOX energy in the isolation cone
 → constant per event factor of 0.88 applied to the data



Uncertainties in the efficiency estimation:

- 3% from material uncertainty
- 1.5% from the EM energy scale
- 3% from trigger efficiency uncertainty
- 6% (3% per photon) from UE correction

Average efficiency ~40%
 Total systematic uncertainty: ~7-15%
 Comparable statistical uncertainty

Corrections and tests

- EM energy scale set by tuning the reconstructed $Z^0 \rightarrow e^+e^-$ mass to the world average by Gaussian fitting in the window $M_{ee} = 86-96 \text{ GeV}/c^2$
 - ➔ correction applied as a function of time before event selection to account for a few events below the energy threshold which the correction pushes above threshold

Corrections and tests

- EM energy scale set by tuning the reconstructed $Z^0 \rightarrow e^+e^-$ mass to the world average by Gaussian fitting in the window $M_{ee} = 86-96 \text{ GeV}/c^2$
 - ➔ correction applied as a function of time before event selection to account for a few events below the energy threshold which the correction pushes above threshold
- Measurement of the $Z^0 \rightarrow e^+e^-$ cross section tests the cross section measurement procedures:
 - ➔ Trigger efficiency
 - ➔ Ability of MC to predict event selection efficiency
 - ➔ Efficiency corrections
 - ➔ Luminosity

“Photon-like” e^+e^- selection applied with special requirements:

 - ➔ Two tracks allowed in cluster
 - ➔ Leading p_T^{trk} cut applied on the 2nd track in cluster
 - ➔ Track isolation corrected subtracting leading p_T^{trk}
 - ➔ $0.8 \leq E/p \leq 1.2$ cut applied to eliminate hard radiation

Measured/published ratio in the window $M_{ee} = 65-115 \text{ GeV}/c^2$: 1.007 ± 0.01 with 5% RMS over time

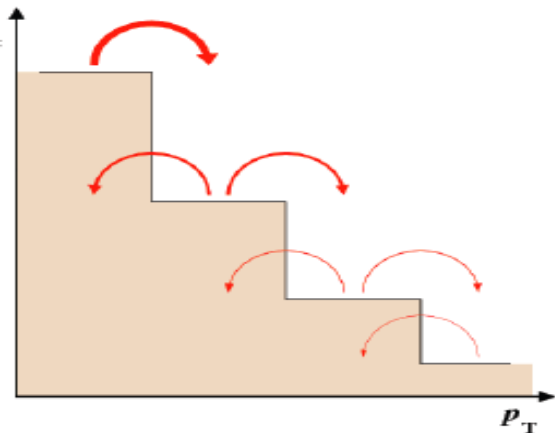
Corrections and tests

- EM energy scale set by tuning the reconstructed $Z^0 \rightarrow e^+e^-$ mass to the world average by Gaussian fitting in the window $M_{ee} = 86-96 \text{ GeV}/c^2$
 - ➔ correction applied as a function of time before event selection to account for a few events below the energy threshold which the correction pushes above threshold
- Measurement of the $Z^0 \rightarrow e^+e^-$ cross section tests the cross section measurement procedures:
 - ➔ Trigger efficiency
 - ➔ Ability of MC to predict event selection efficiency
 - ➔ Efficiency corrections
 - ➔ Luminosity

“Photon-like” e^+e^- selection applied with special requirements:

- ➔ Two tracks allowed in cluster
- ➔ Leading p_T^{trk} cut applied on the 2nd track in cluster
- ➔ Track isolation corrected subtracting leading p_T^{trk}
- ➔ $0.8 \leq E/p \leq 1.2$ cut applied to eliminate hard radiation

Measured/published ratio in $M_{ee} = 65-115 \text{ GeV}/c^2$ window: 1.007 ± 0.01 with 5% RMS over time

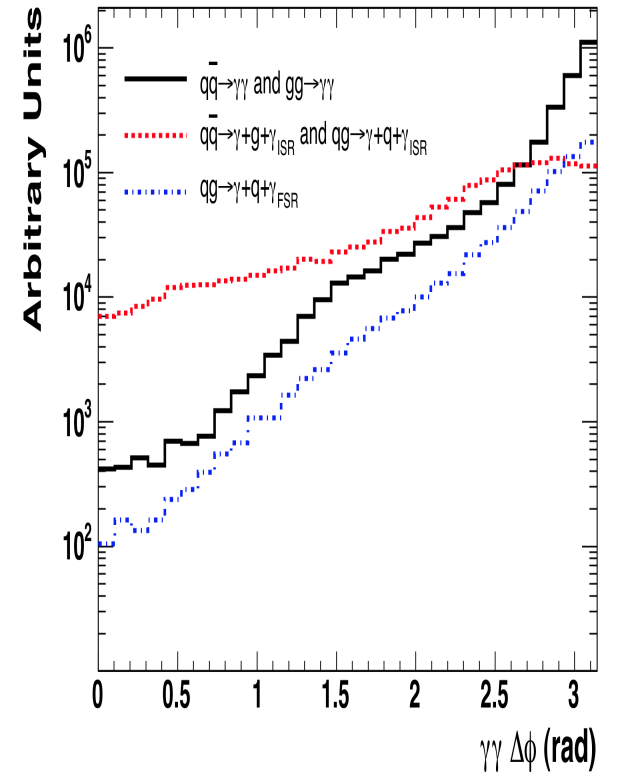
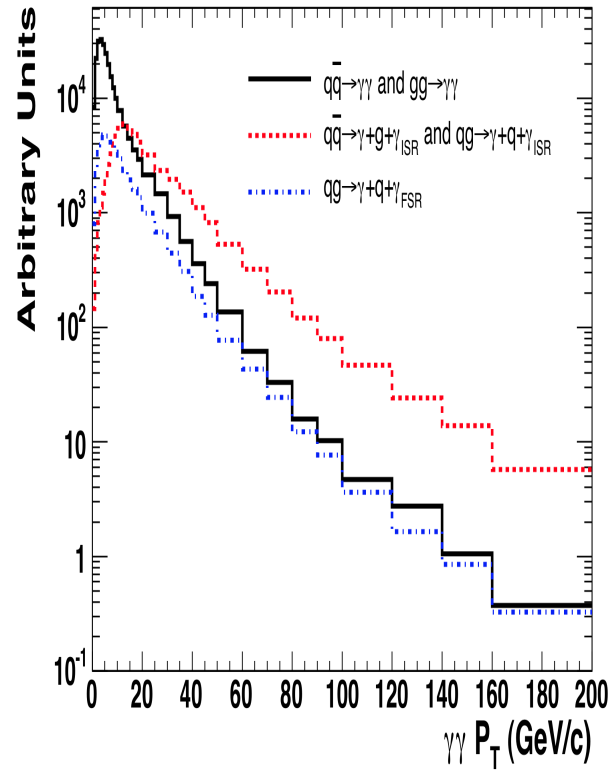
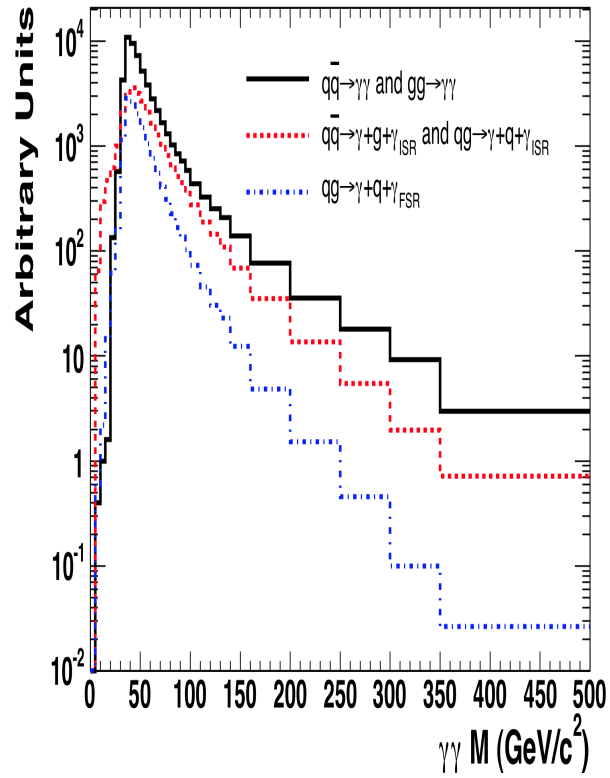


- Experimental effects (photon energy resolution, misvertexing) lead to event migration

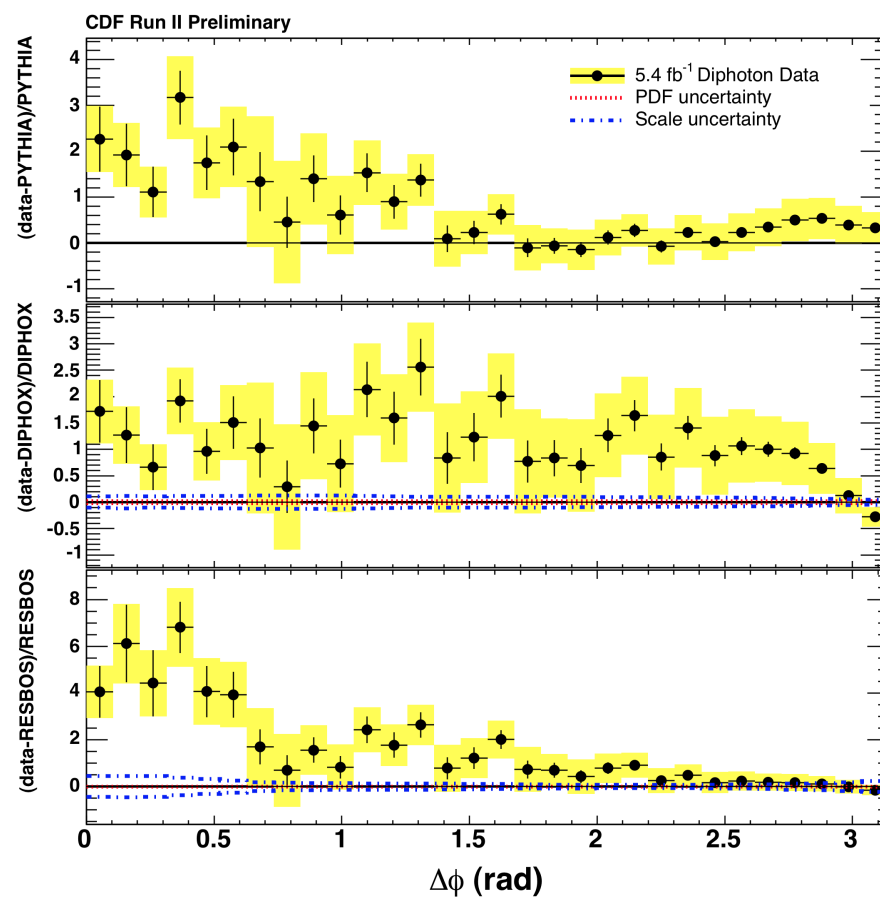
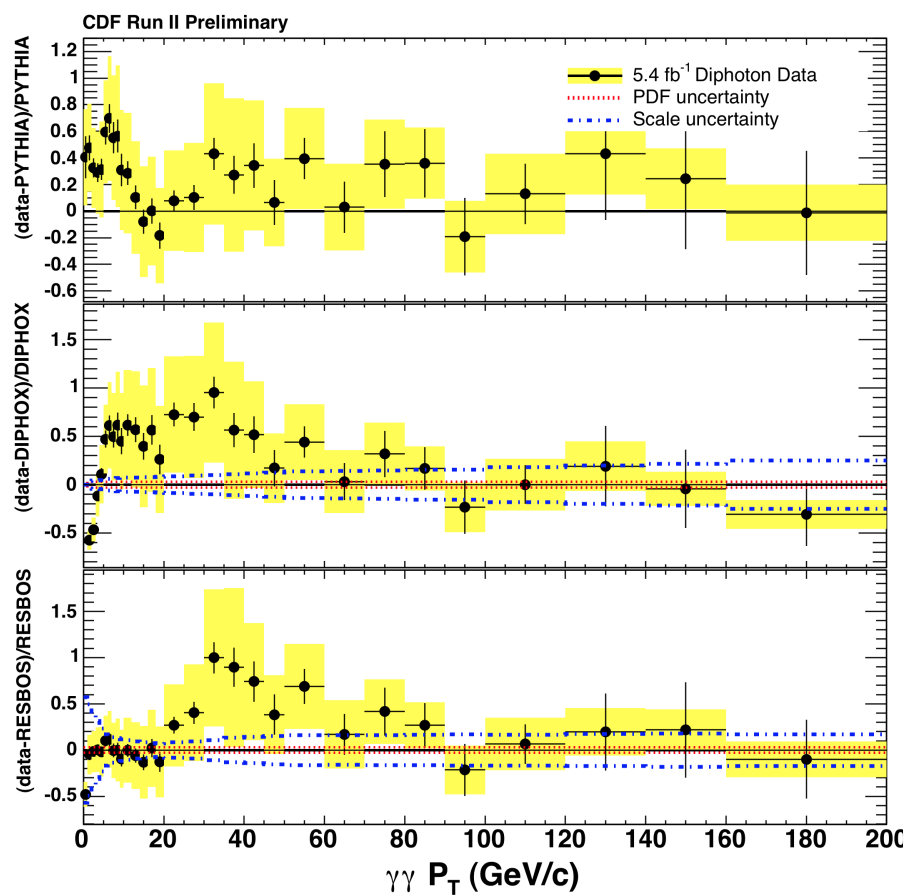
Purity (bin i) = $N(\text{gen bin } i \text{ AND reco bin } i) / N(\text{reco bin } i)$

➔ The acceptance correction also accounts for this

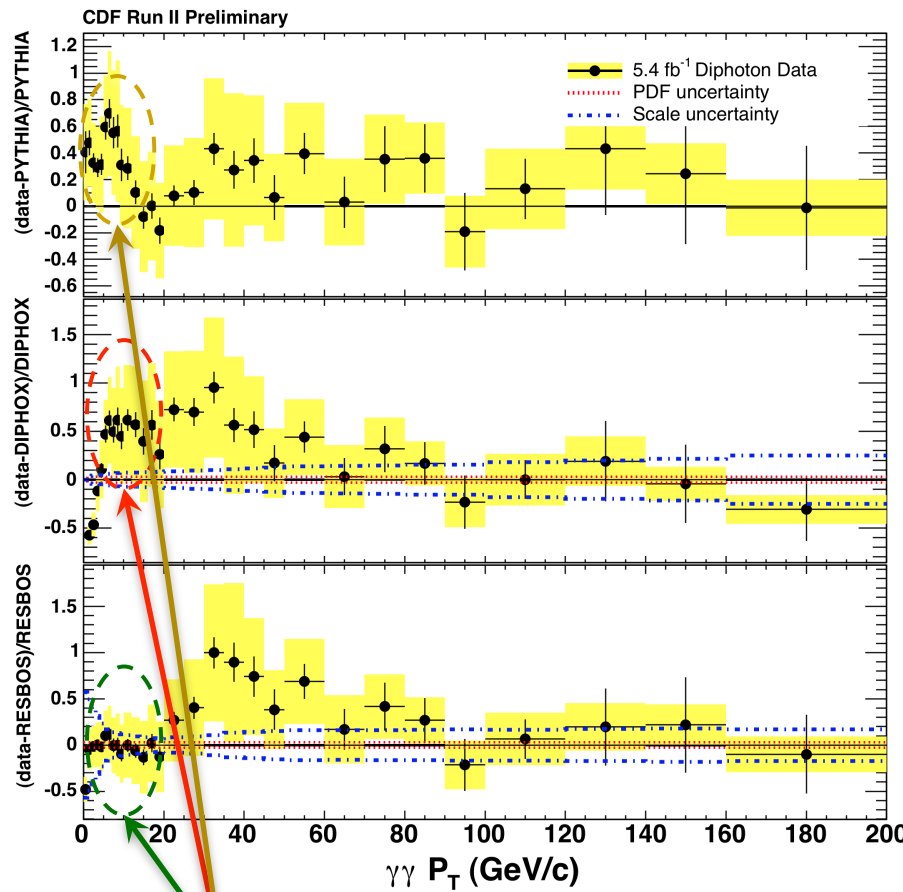
Matrix element and radiation contributions in Pythia



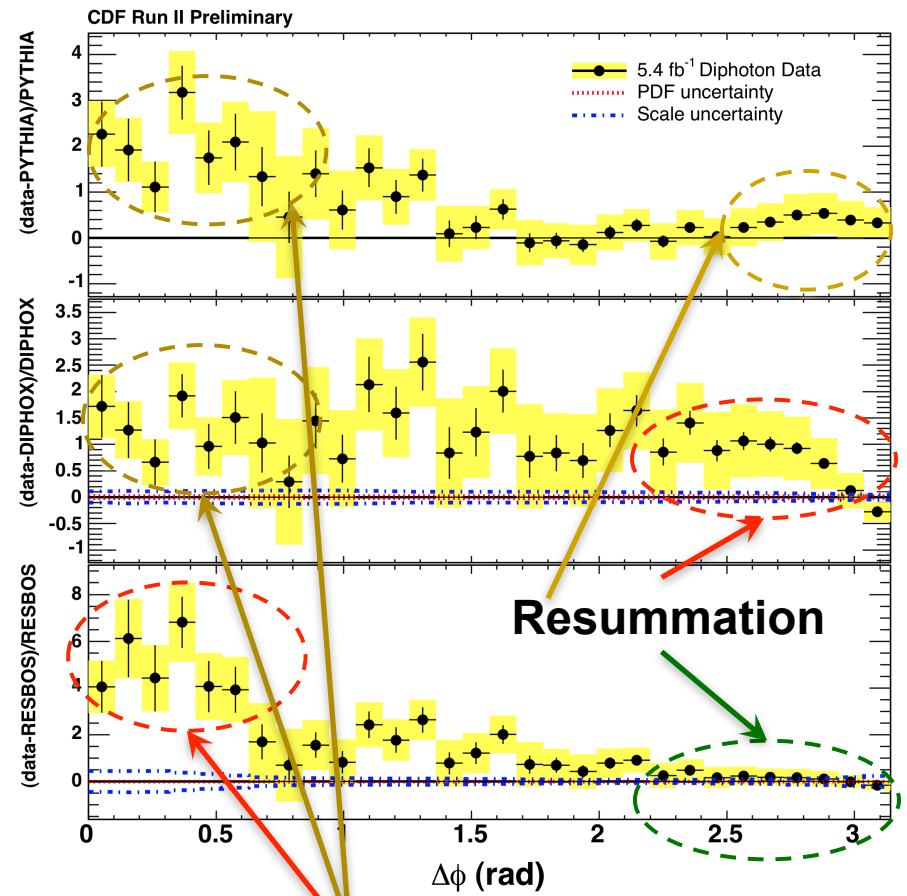
Data-to-theory cross section ratios



Data-to-theory cross section ratios



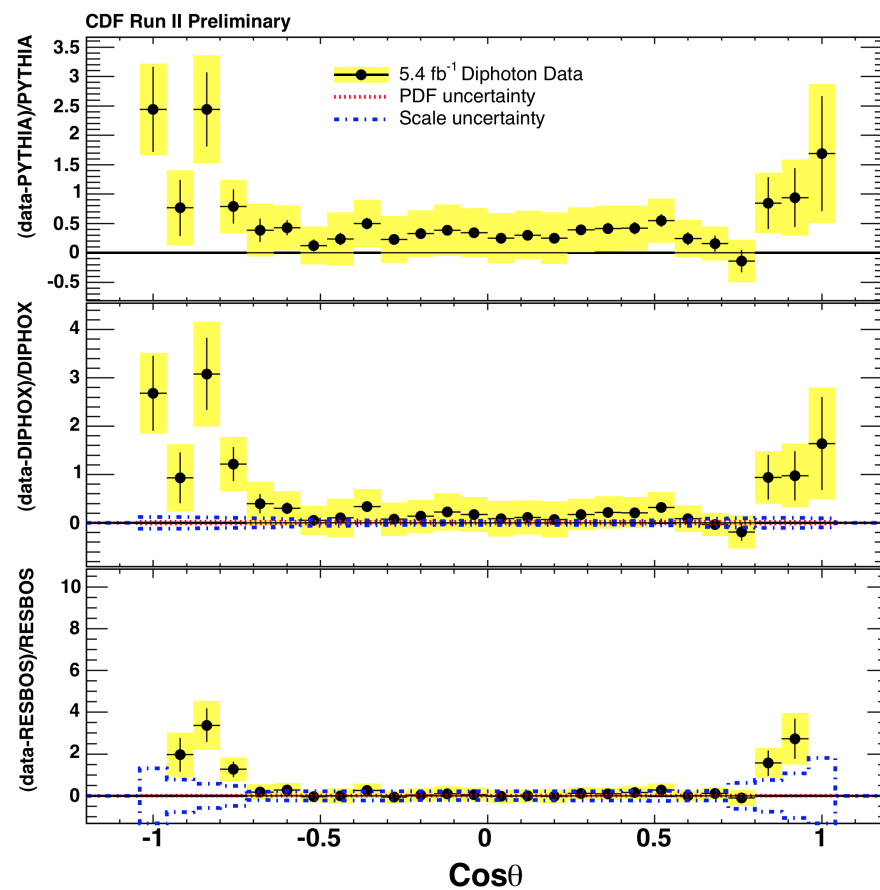
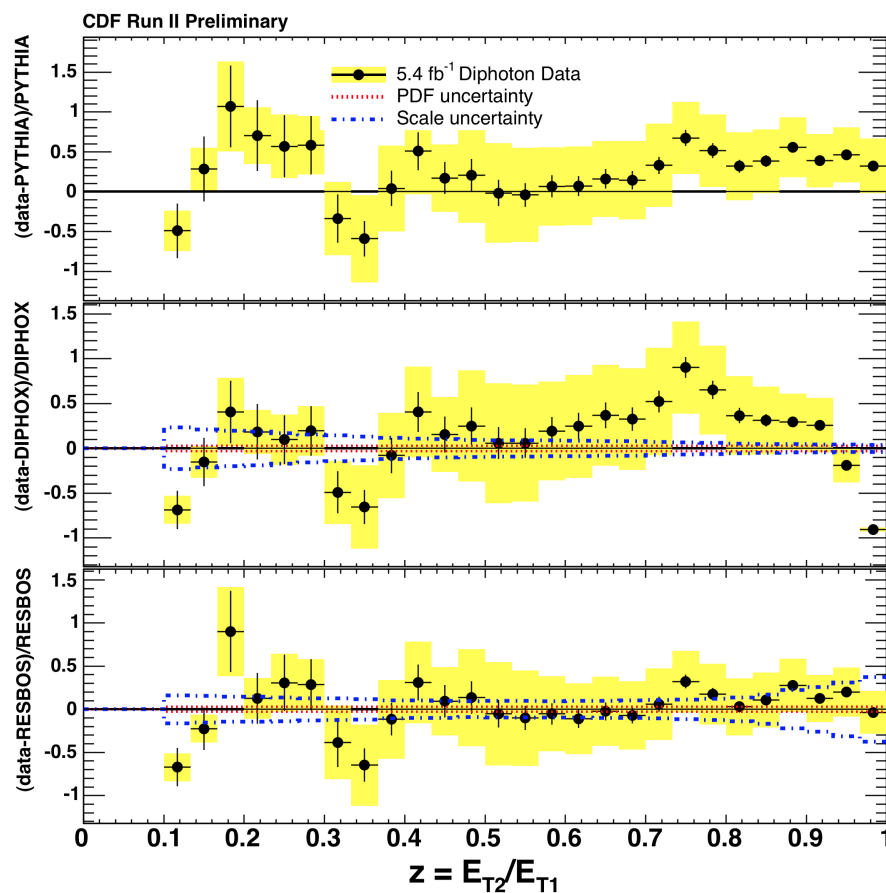
Resummation



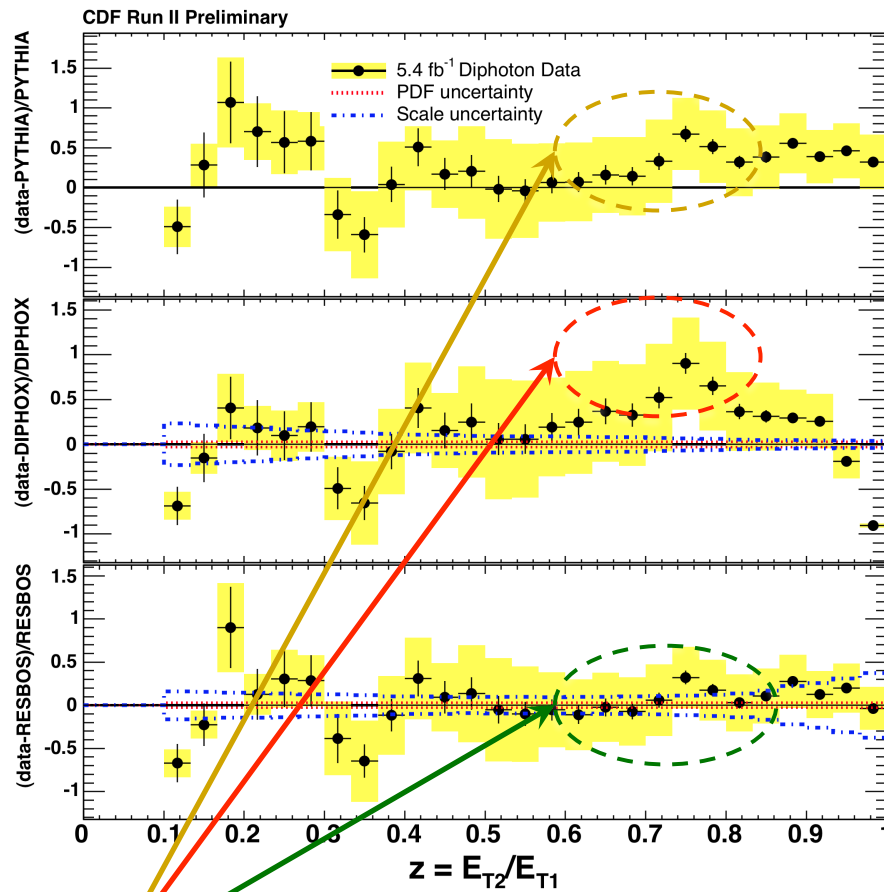
Resummation

Fragmentations

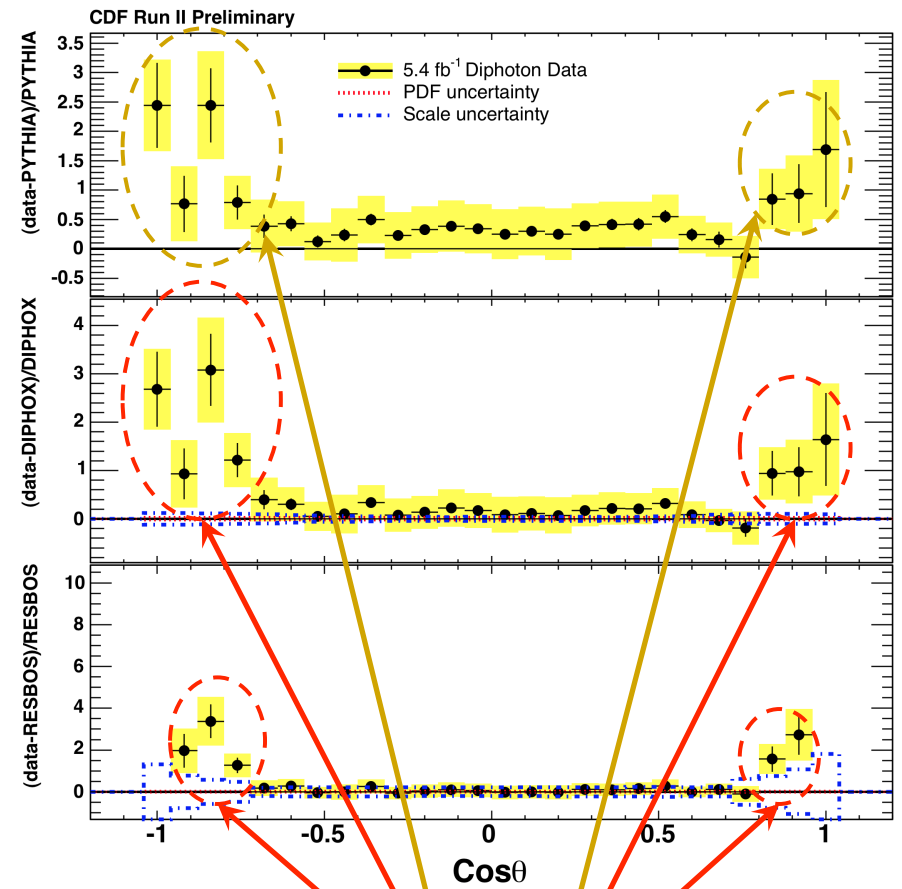
Data-to-theory cross section ratios



Data-to-theory cross section ratios

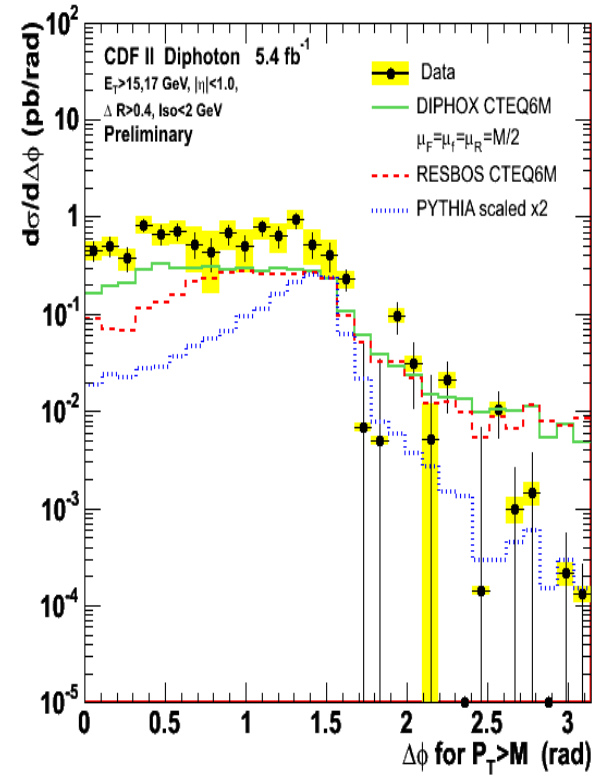
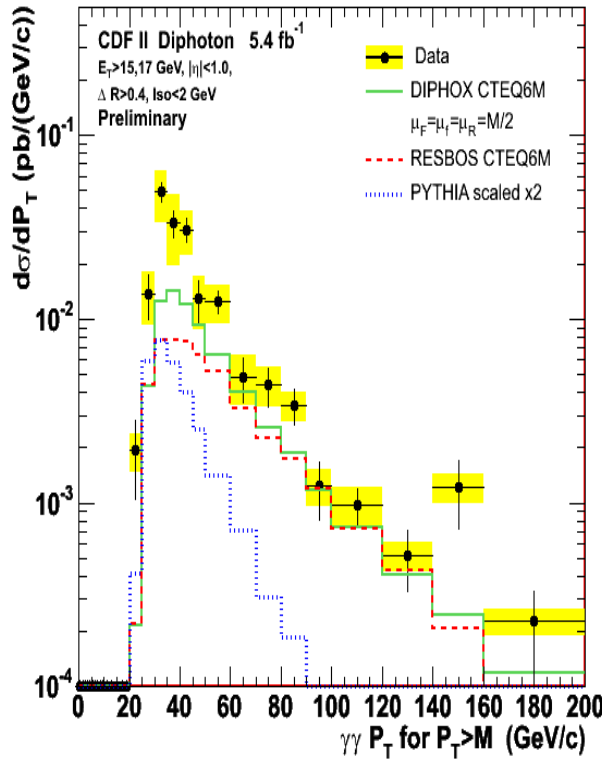
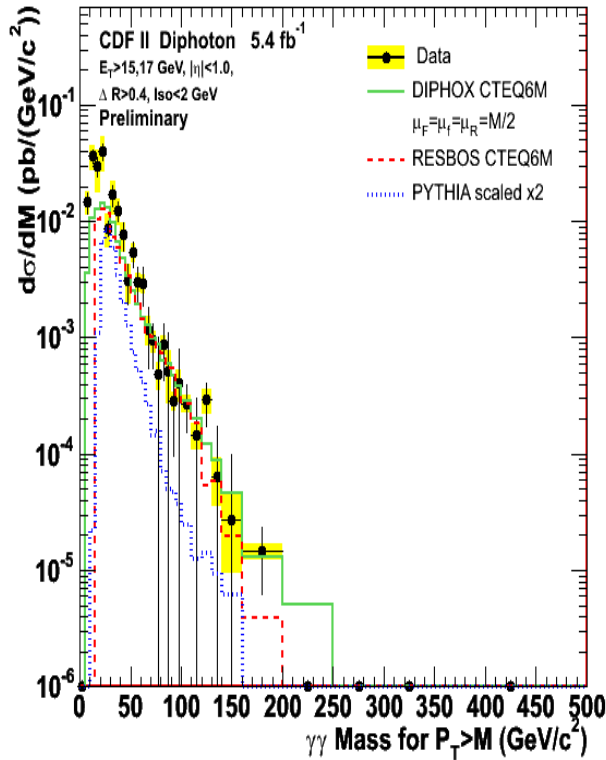


Resummation



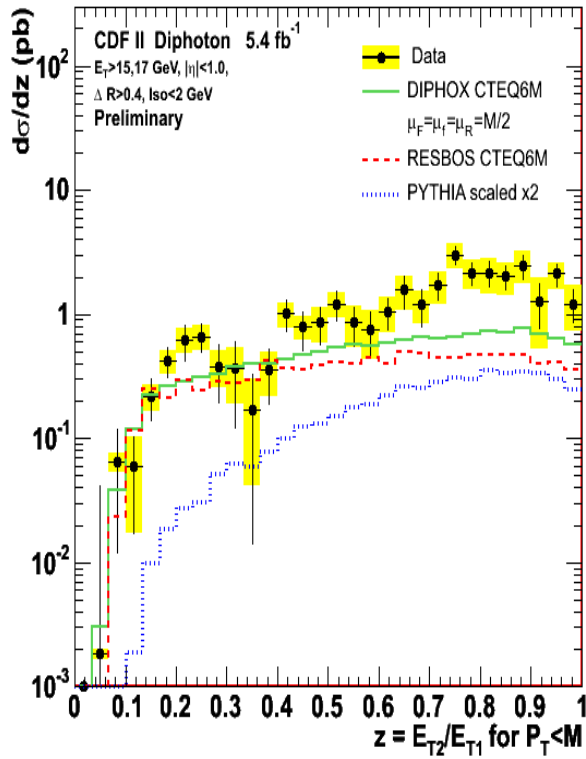
Fragmentations

Differential cross sections for $p_T(\gamma\gamma) > M(\gamma\gamma)$

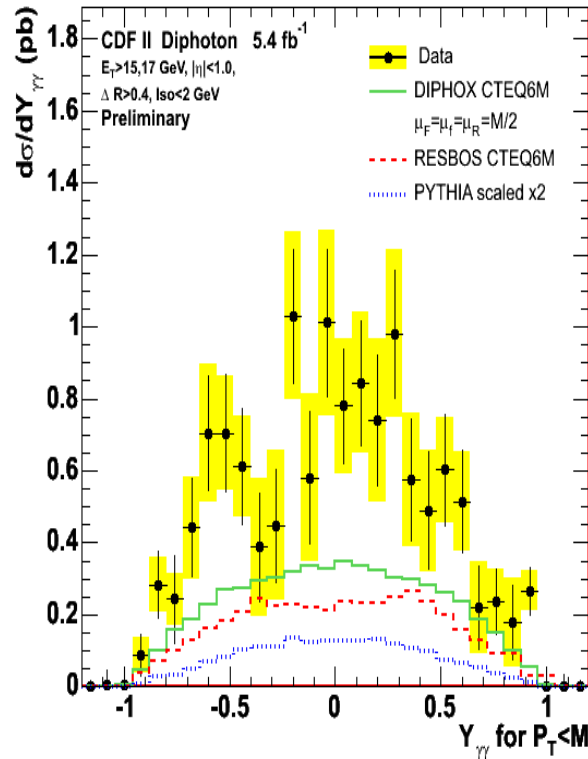


- Theory underestimates the data at the peak $M_{\gamma\gamma} \sim 30$ GeV/c²
- Theory underestimates the data for $p_T(\gamma\gamma) < 90$ GeV/c
- Theory underestimates the data for $\Delta\phi_{\gamma\gamma} < 1.7$ rad

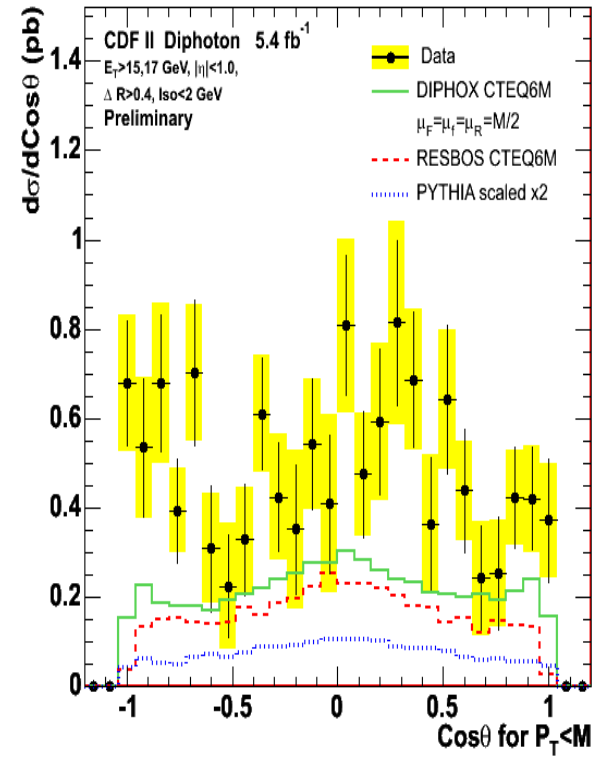
Differential cross sections for $p_T(\gamma\gamma) > M(\gamma\gamma)$



- Theory underestimates the data



- Theory underestimates the data



- Theory underestimates the data