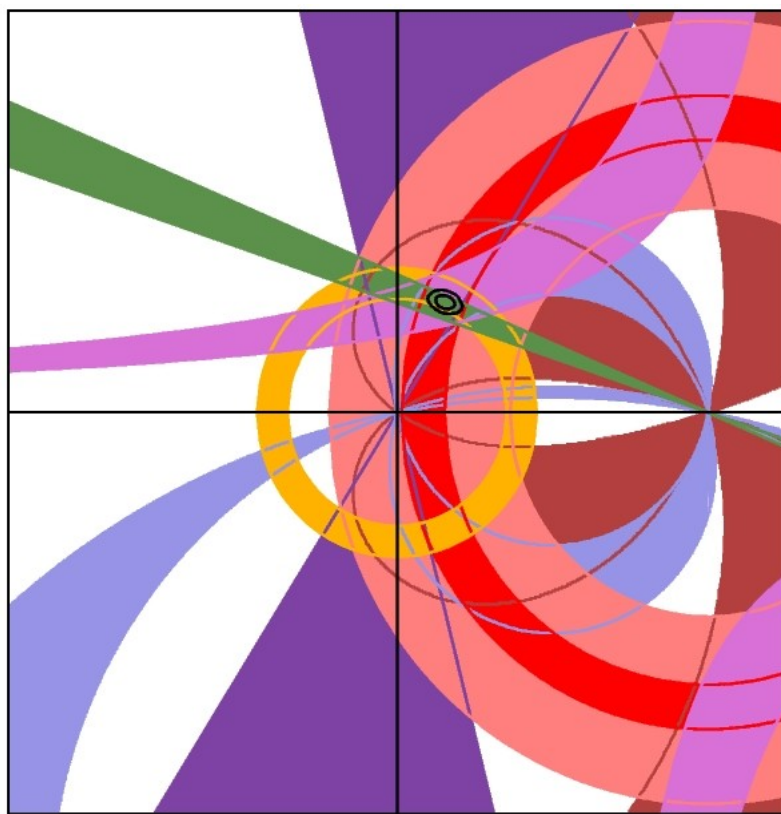


Standard Model updates and new physics analysis with the Unitarity Triangle fit



Marcella Bona



Queen Mary
University of London

**EPS 2011,
Grenoble, France
July 22nd, 2011**



M.Bona *et al.*, UTfit
JHEP0507:028, 2005

www.utfit.org

**A. Bevan, M. Bona, M. Ciuchini,
D. Derkach, E. Franco, V. Lubicz,
G. Martinelli, F. Parodi, M. Pierini,
C. Schiavi, L. Silvestrini, A. Stocchi,
V. Sordini, C. Tarantino and V. Vagnoni**

unitarity Triangle analysis in the SM

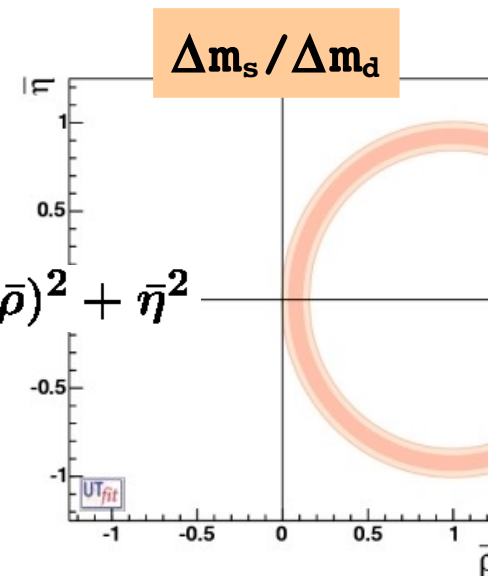
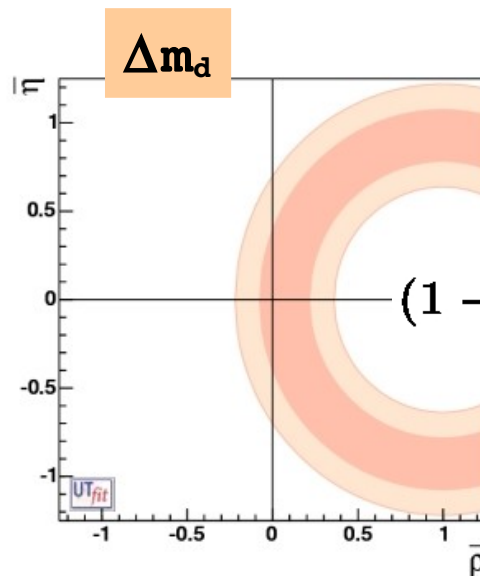
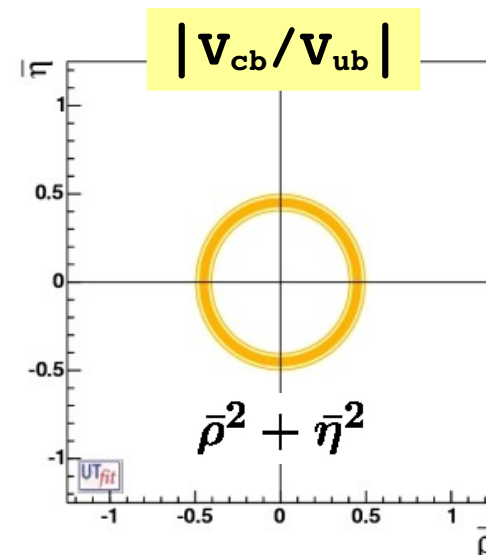
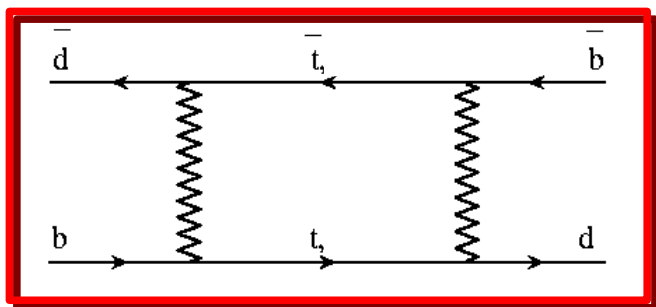
→ SM UT analysis:

- provide the best determination of CKM parameters
- test the consistency of the SM (“direct” vs “indirect” determinations)
- provide predictions for SM observables (ex. $\sin 2\beta$, Δm_s , ...)

CP-conserving inputs

$|V_{ub}|/|V_{cb}| \sim R_b$ (tree-level)

B_d - B_d and B_s - B_s mixing



V_{cb} and V_{ub}

should update to
 39.5 ± 1.0

Laiho *et al*

$V_{cb} (excl) = (38.6 \pm 1.2) 10^{-3}$

HFAG

$V_{cb} (incl) = (41.7 \pm 0.7) 10^{-3}$

$\sim 2.2\sigma$ discrepancy

new preliminary value from HFAG

3.26 ± 0.30

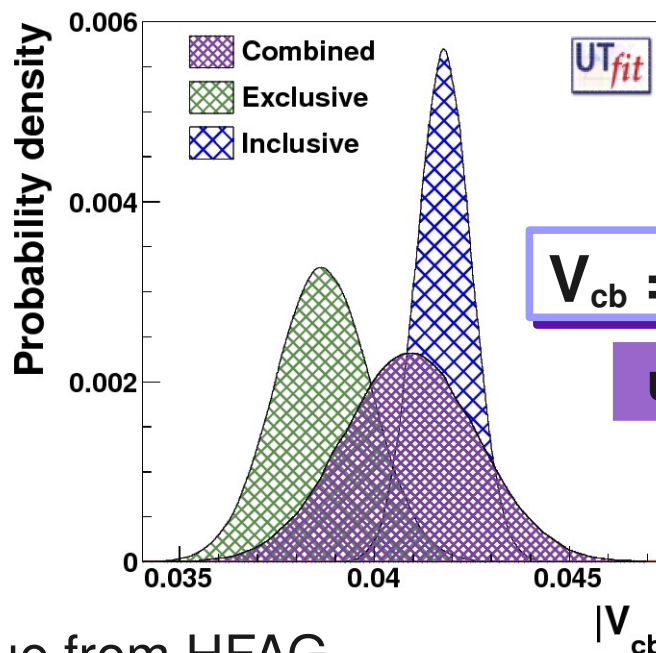
Laiho *et al*

$V_{ub} (excl) = (3.42 \pm 0.37) 10^{-3}$

UTfit from HFAG

$V_{ub} (incl) = (4.32 \pm 0.38) 10^{-3}$

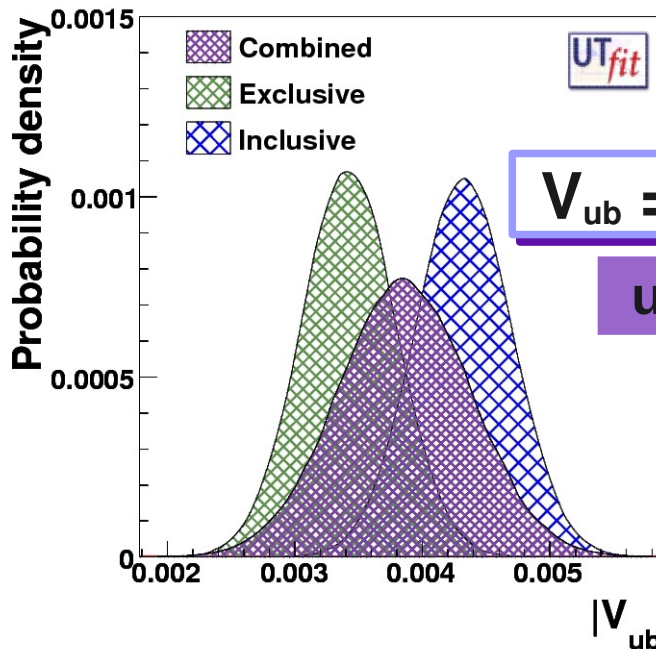
$\sim 1.7\sigma$ discrepancy



UTfit input value

$V_{cb} = (40.9 \pm 1.7) 10^{-3}$

uncertainty $\sim 4\%$



UTfit input value

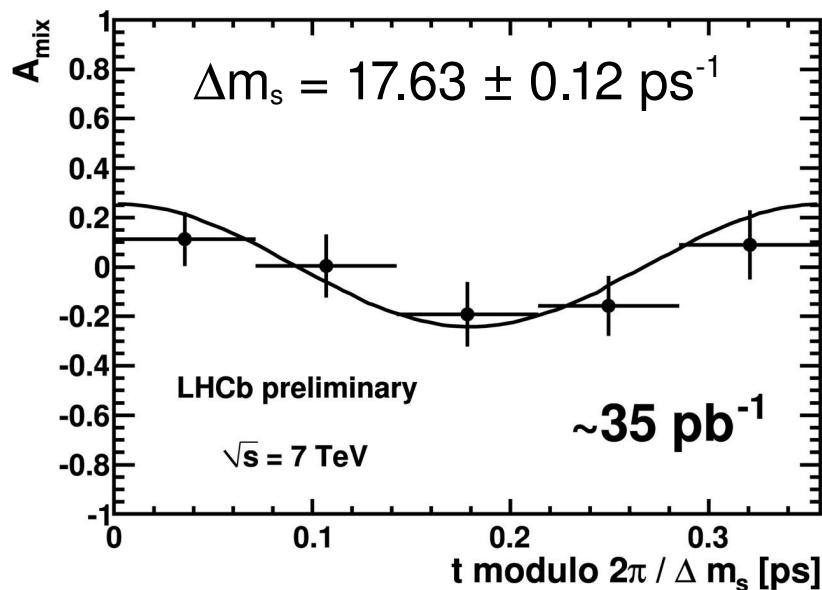
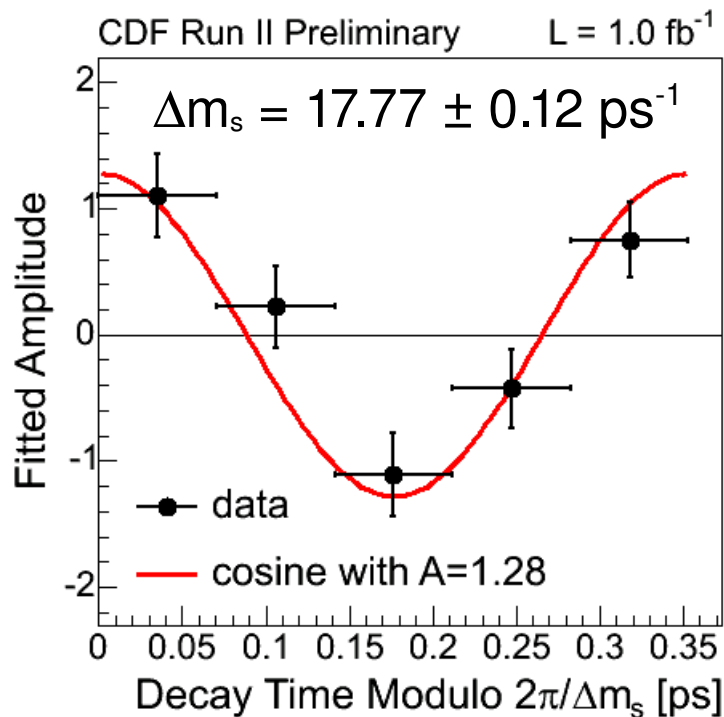
$V_{ub} = (3.86 \pm 0.52) 10^{-3}$

uncertainty $\sim 13\%$

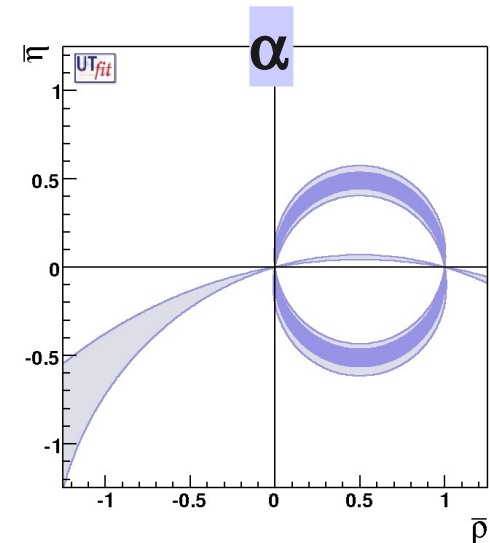
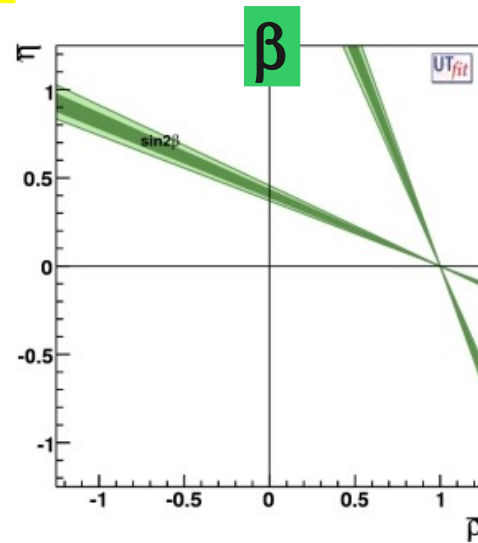
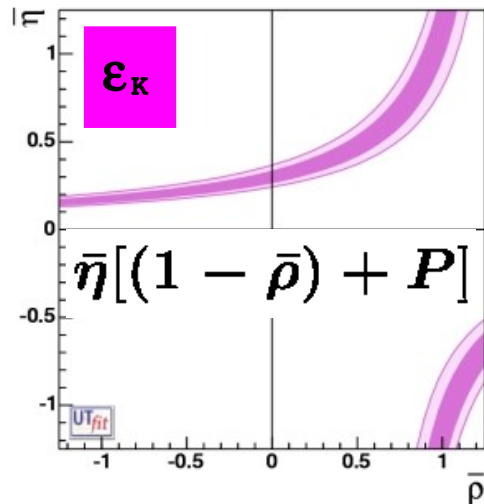
B_q - B_q mixing

- $\Delta m_d = (0.507 \pm 0.005) \text{ ps}^{-1}$
- $\Delta m_s = (17.70 \pm 0.08) \text{ ps}^{-1}$

New world average from CDF and LHCb



CP-violating inputs



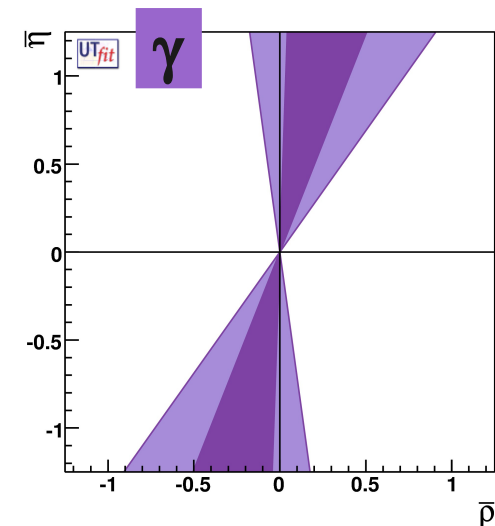
ε_K from K-K mixing

$$- B_K = 0.731 \pm 0.036$$

$\sin 2\beta$ from $B \rightarrow J/\psi K^0$ + theory

α from $\pi\pi, \rho\rho, \pi\rho$ decays:
combined: $(91 \pm 6)^\circ$

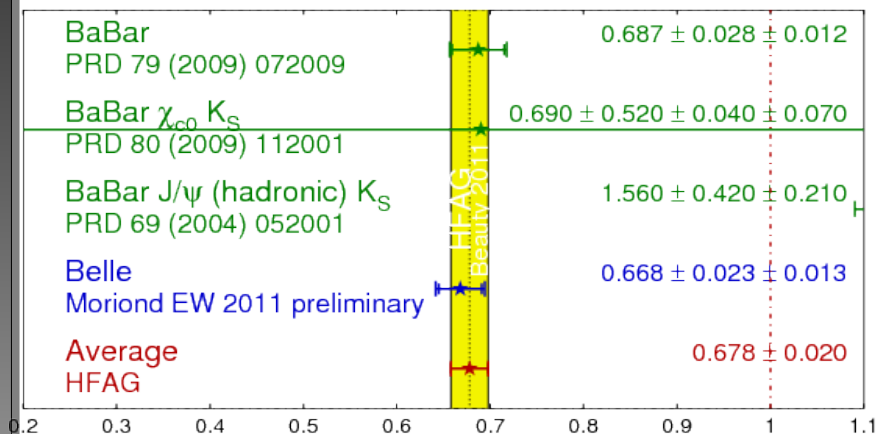
γ from $B \rightarrow DK$ decays (tree level)



Latest $\sin 2\beta$ results:

$$\sin(2\beta) \equiv \sin(2\phi_1) \quad \text{HFAG}$$

Beauty 2011
PRELIMINARY



BABAR Collaboration
Physical Review D 79:072009, 2009

BaBar with $465 \cdot 10^6$ $\bar{B}B$ pairs

$$\sin 2\beta = 0.666 \pm 0.031 \pm 0.013$$

Belle with $772 \cdot 10^6$ $\bar{B}B$ pairs

$$\sin 2\beta = 0.663 \pm 0.025 \pm 0.013$$

Belle Collaboration
Moriond EW 2011, preliminary

UTfit input value

$$\sin 2\beta(J/\psi K^0) = 0.664 \pm 0.022$$

EPS2011 update from Belle:

$$\sin 2\beta(J/\psi K^0) = 0.668 \pm 0.026$$

so new average should be

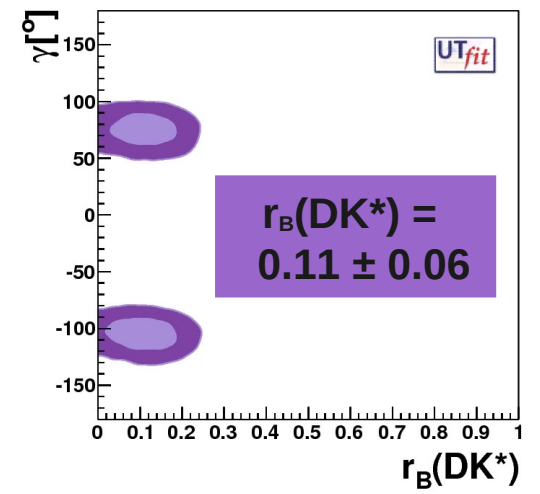
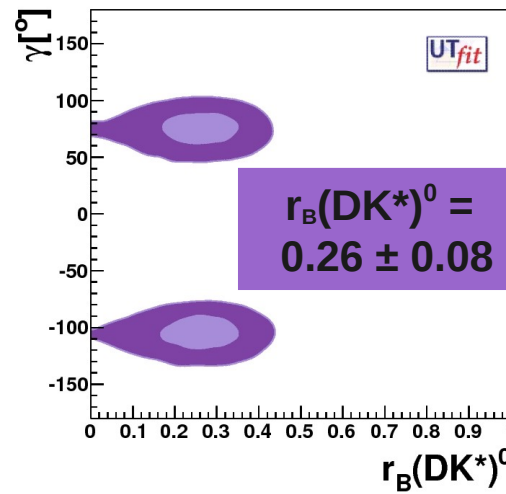
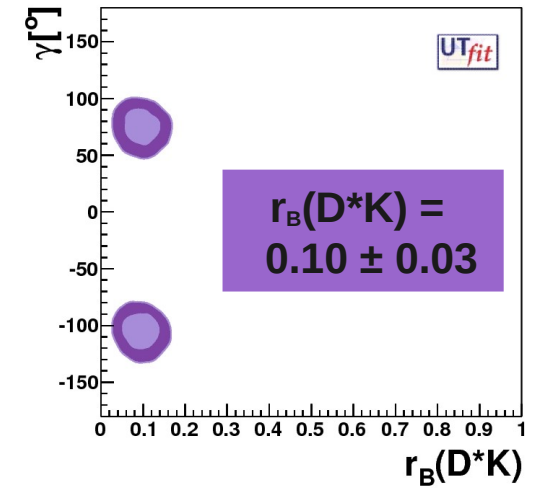
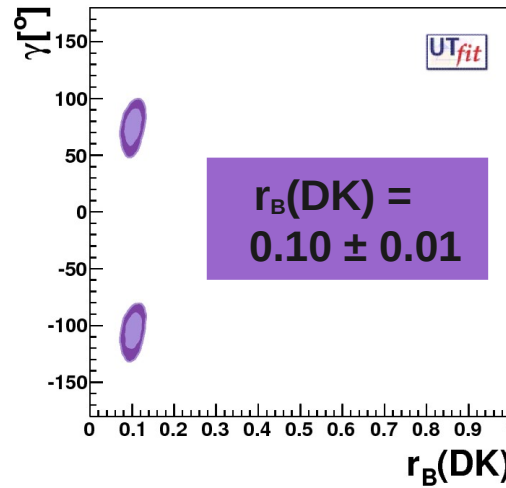
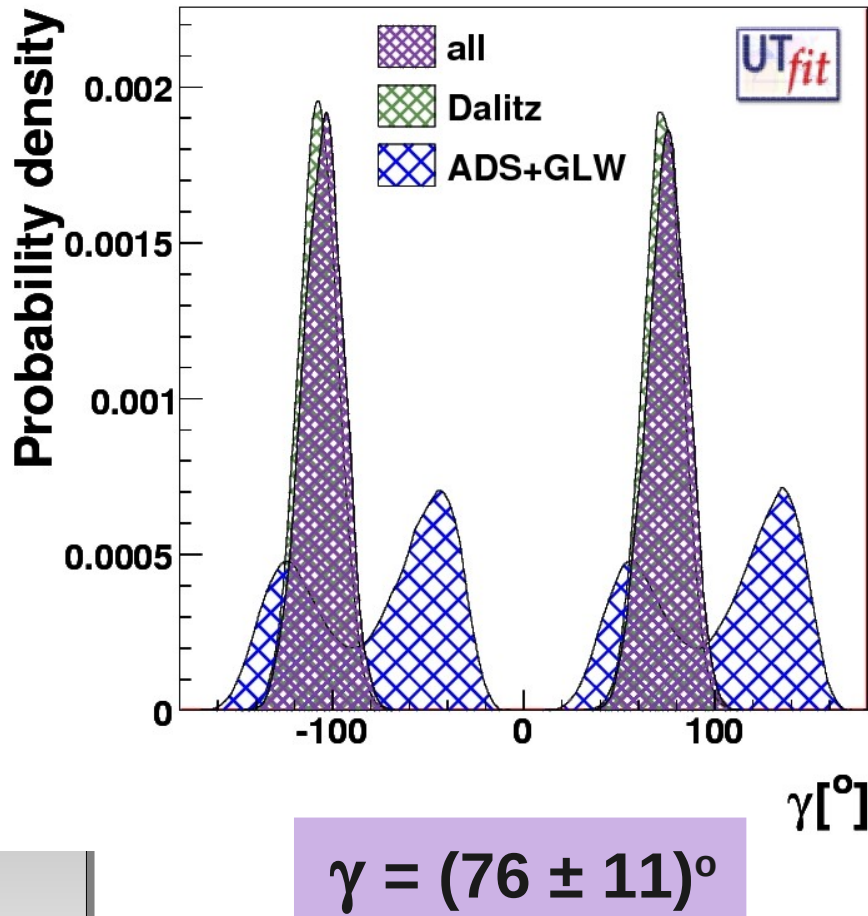
$$\sin 2\beta(J/\psi K^0) = 0.667 \pm 0.021$$

data-driven theoretical uncertainty

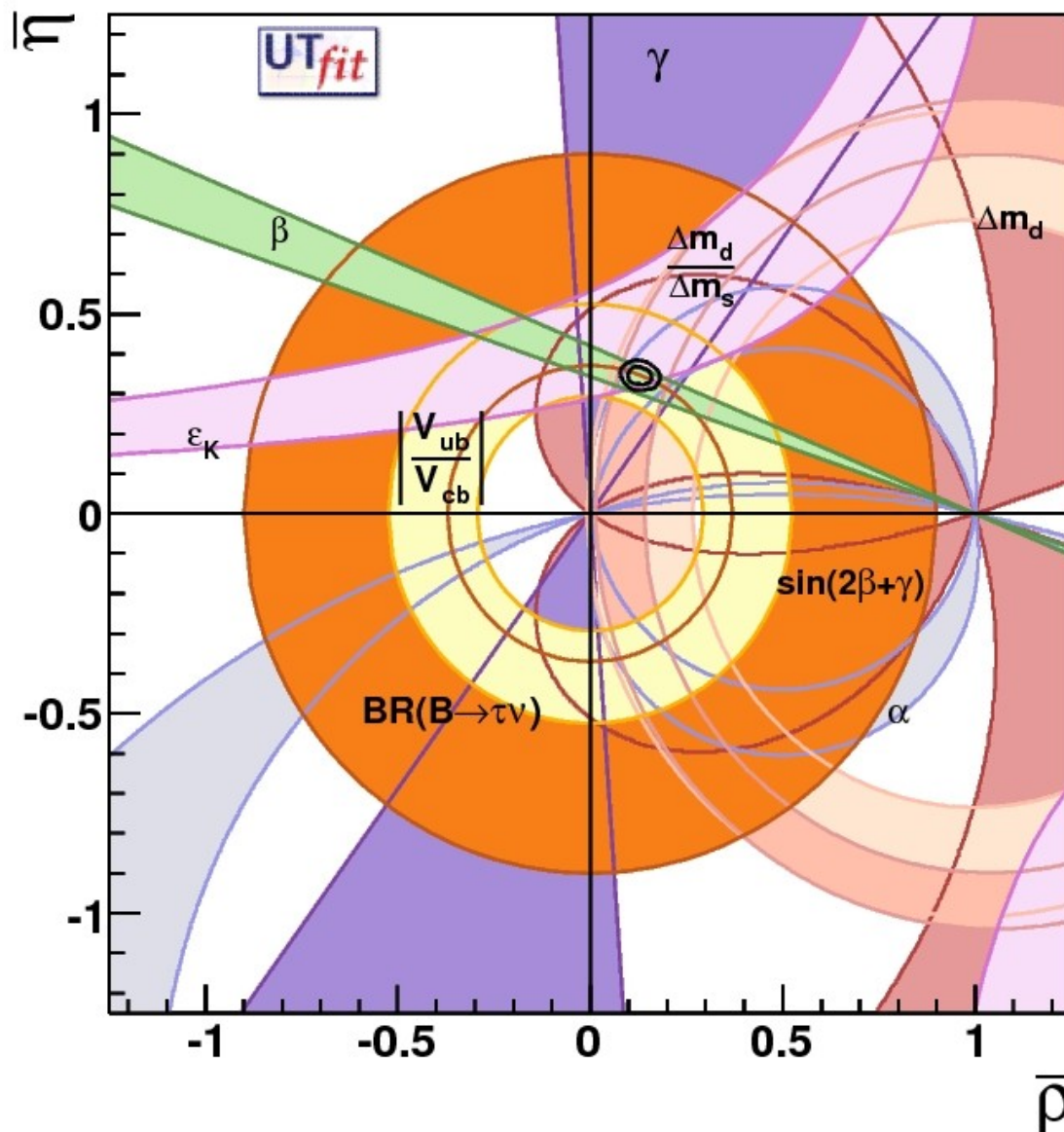
$$\Delta S = 0.000 \pm 0.012$$

M.Ciuchini, M.Pierini, L.Silvestrini
Phys. Rev. Lett. 95, 221804 (2005)

γ and DK trees



unitarity Triangle analysis in the SM



levels @
95% Prob

$$\bar{\rho} = 0.129 \pm 0.022$$

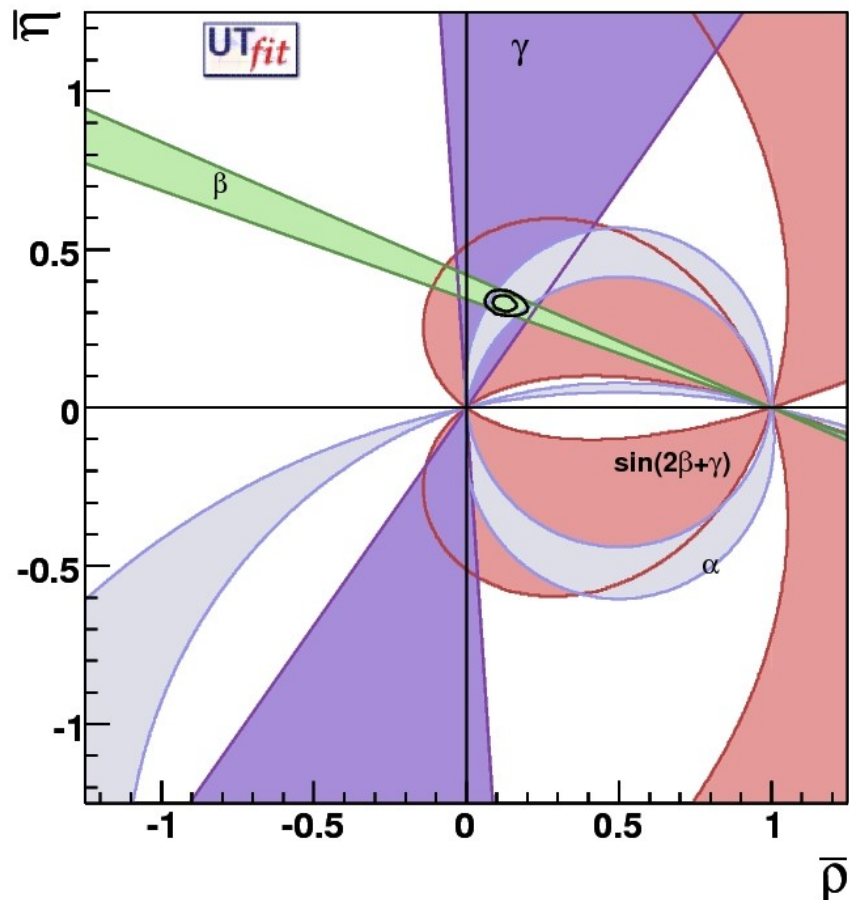
$$\bar{\eta} = 0.346 \pm 0.015$$

$$\beta = (22 \pm 1)^\circ$$

$$\gamma = (69 \pm 3)^\circ$$

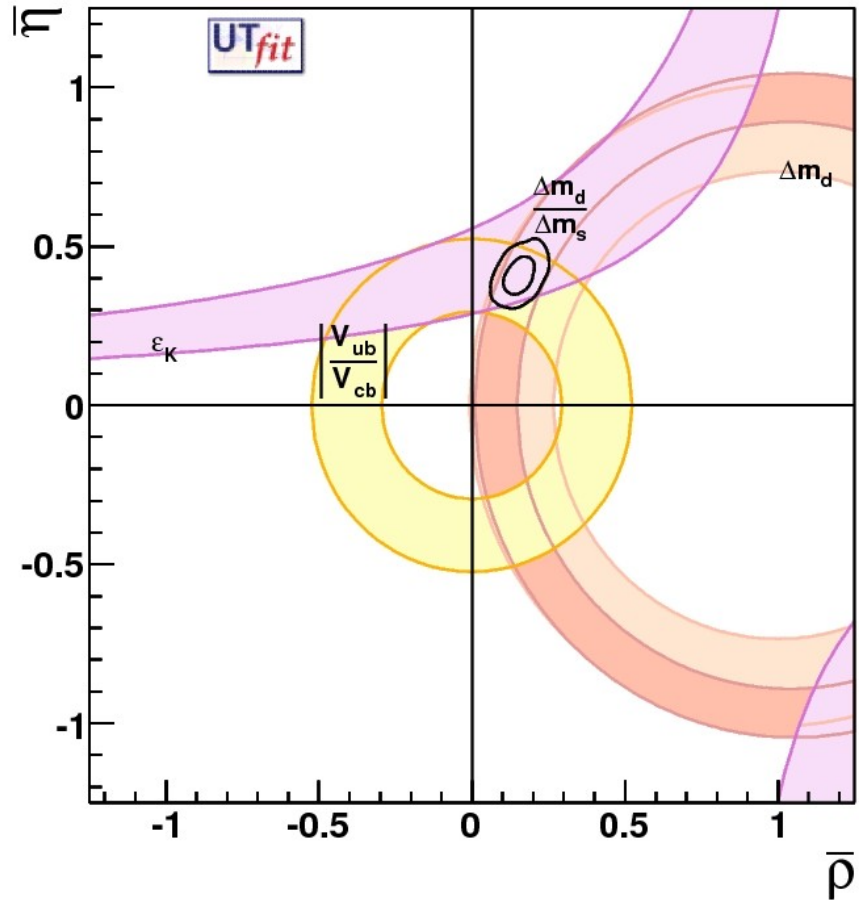
$$\alpha = (89 \pm 3)^\circ$$

angles vs the others



$$\bar{\rho} = 0.127 \pm 0.027$$

$$\bar{\eta} = 0.329 \pm 0.016$$



$$\bar{\rho} = 0.157 \pm 0.037$$

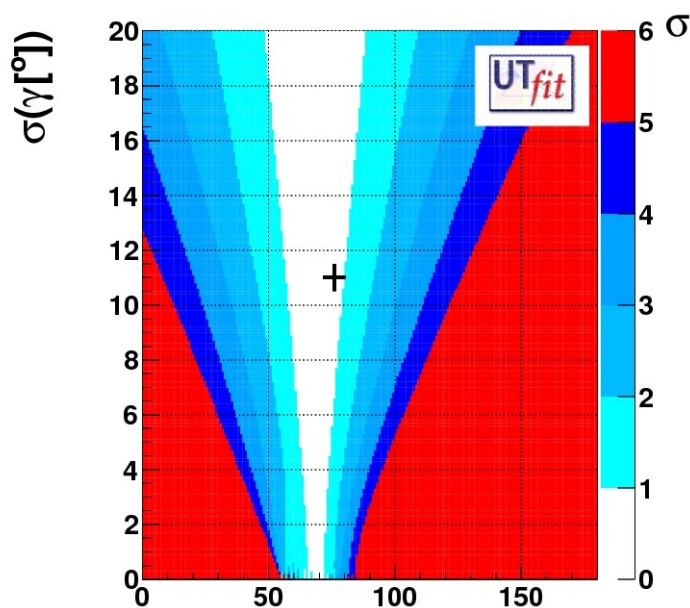
$$\bar{\eta} = 0.409 \pm 0.043$$

compatibility plots

A way to “measure” the agreement of a single measurement with the indirect determination from the fit using all the other inputs: test for the SM description of the flavor physics

Color code: agreement between the predicted values and the measurements at better than 1, 2, ... $n\sigma$

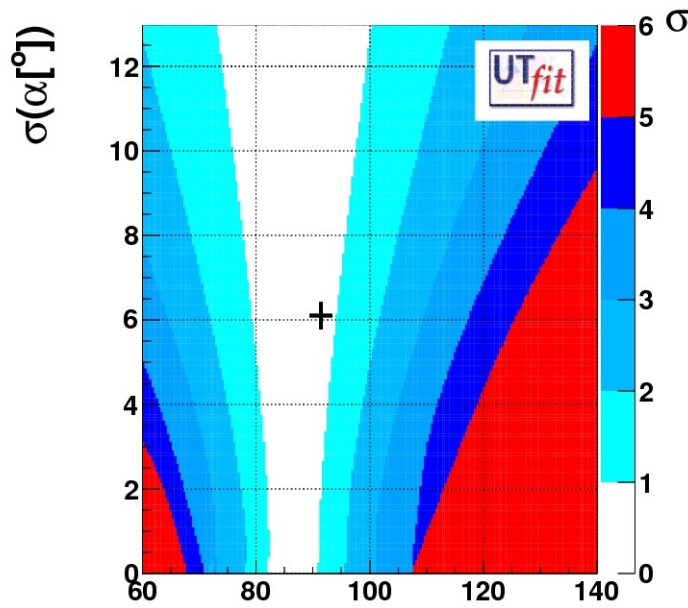
The cross has the coordinates (x,y)=(central value, error) of the direct measurement



$$\gamma_{\text{exp}} = (76 \pm 11)^\circ$$

$$\gamma_{\text{UTfit}} = (69 \pm 3)^\circ$$

$<1\sigma$

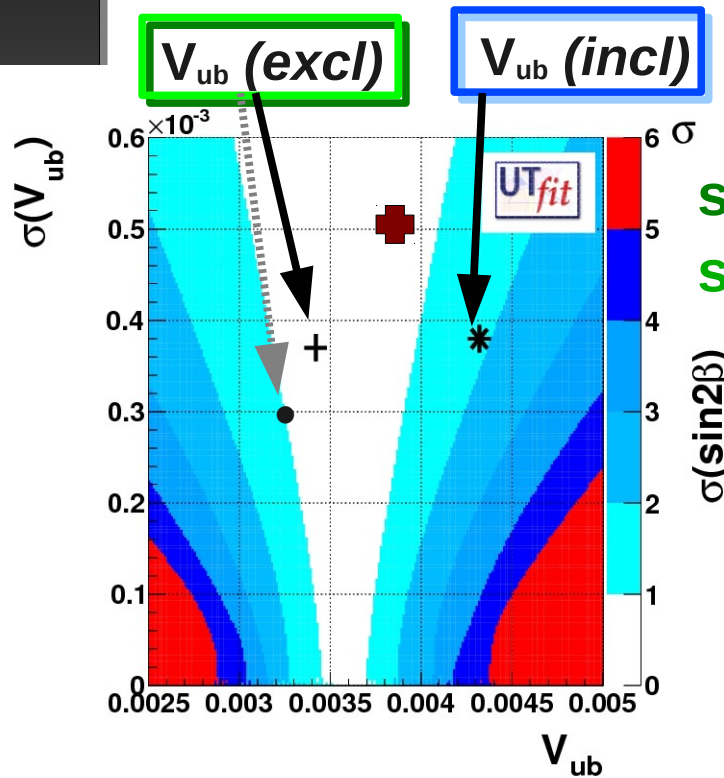


$$\alpha_{\text{exp}} = (91 \pm 6)^\circ$$

$$\alpha_{\text{UTfit}} = (86 \pm 4)^\circ$$

$<1\sigma$

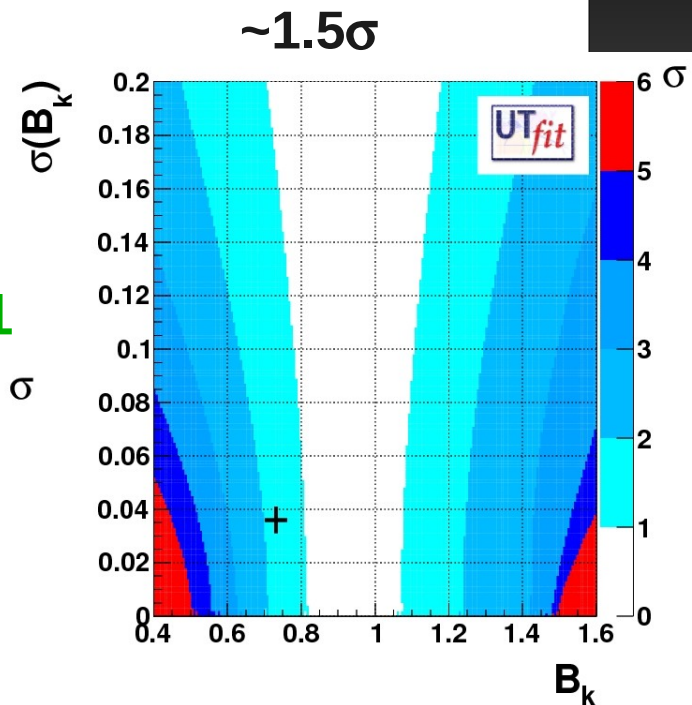
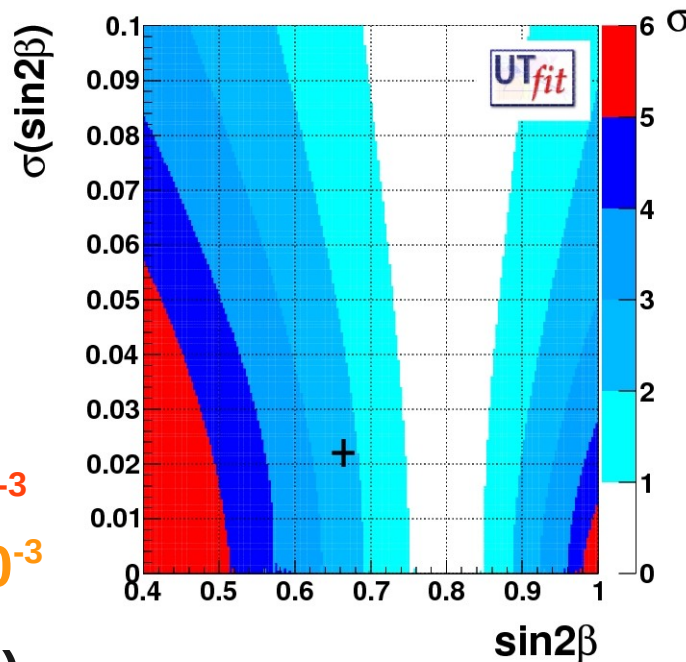
tensions



$V_{ub_{exp}} = (3.86 \pm 0.52) \cdot 10^{-3}$
 $V_{ub_{UTfit}} = (3.57 \pm 0.14) \cdot 10^{-3}$

$<1\sigma$ (incl $\sim 1.8\sigma$)

$\sim 2.3\sigma$
 $\sin 2\beta_{exp} = 0.664 \pm 0.022$
 $\sin 2\beta_{UTfit} = 0.803 \pm 0.051$

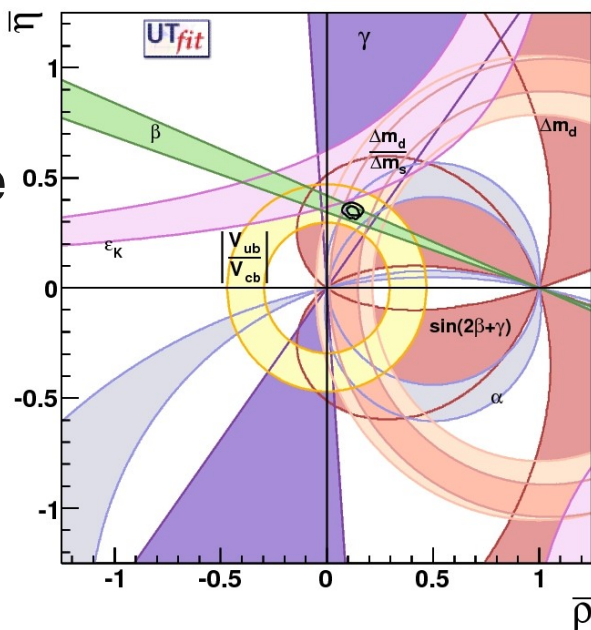


$BK_{exp} = 0.731 \pm 0.036$

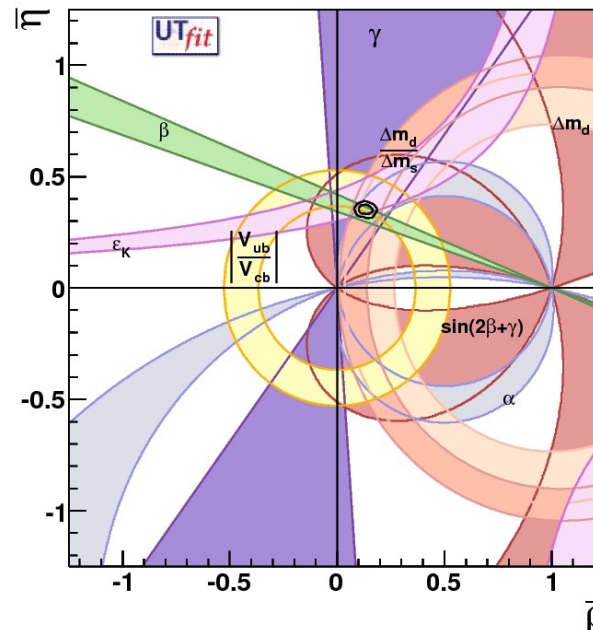
$BK_{UTfit} = 0.93 \pm 0.13$

$BK_{nolattice} = 0.85 \pm 0.14$

only
exclusive
values

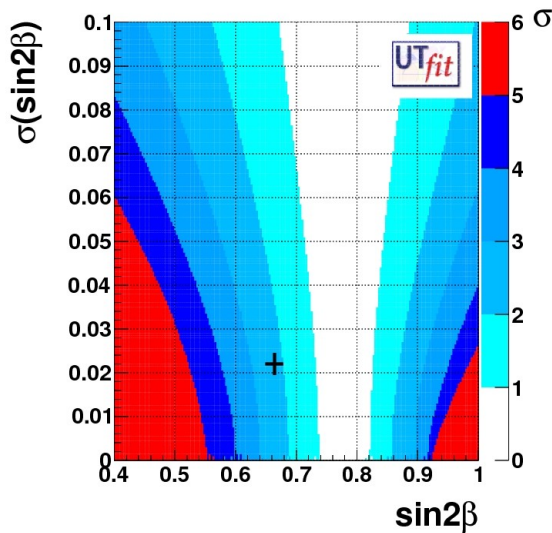


only
inclusive
values



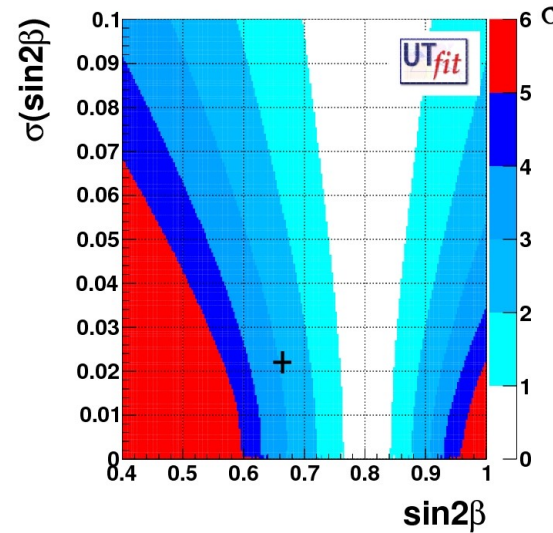
$\sin 2\beta_{UTfit} = 0.783 \pm 0.043$

$\sim 2.3\sigma$



$\sin 2\beta_{UTfit} = 0.804 \pm 0.040$

$\sim 2.9\sigma$



$\sin 2\beta_{UTfit} = 0.846 \pm 0.062 \rightarrow$ no semileptonic

$\sim 2.4\sigma$

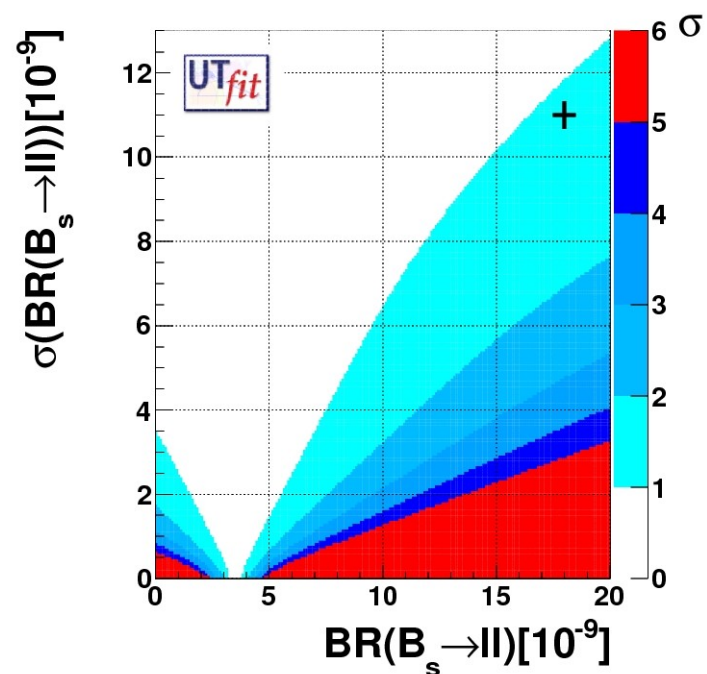
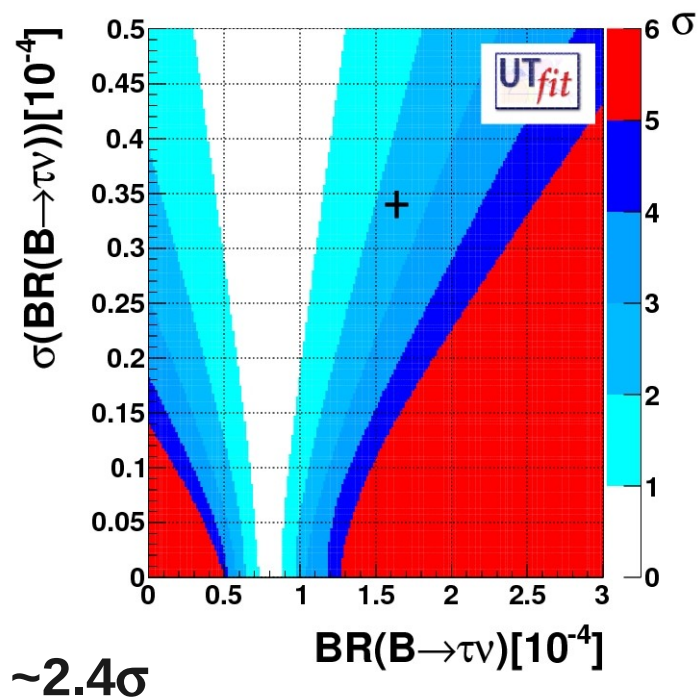
more standard model predictions:

current HFAG world average

$$\text{BR}(B \rightarrow \tau \nu) = (1.64 \pm 0.34) 10^{-4}$$

latest CDF result

$$\text{BR}(B_s \rightarrow \mu \mu) = (18^{+11}_{-9}) 10^{-9}$$



indirect determinations from UT

$$\text{BR}(B \rightarrow \tau \nu) = (0.79 \pm 0.08) 10^{-4}$$

$$\text{BR}(B_s \rightarrow \mu \mu) = (3.54 \pm 0.29) 10^{-9}$$

UTfit beyond the MFV:

fit simultaneously for the CKM and the NP parameters (generalized UT fit)

- add most general loop NP to all sectors
- use all available experimental info
- find out NP contributions to $\Delta F=2$ transitions

B_d and B_s mixing amplitudes (2+2 real parameters):

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

$$\Delta m_{q/K} = C_{B_q/\Delta m_K} (\Delta m_{q/K})^{SM}$$

$$A_{CP}^{B_d \rightarrow J/\psi K_s} = \sin 2(\beta + \phi_{B_d})$$

$$A_{SL}^q = \text{Im}(\Gamma_{12}^q / A_q)$$

$$\varepsilon_K = C_\varepsilon \varepsilon_K^{SM}$$

$$A_{CP}^{B_s \rightarrow J/\psi \phi} \sim \sin 2(-\beta_s + \phi_{B_s})$$

$$\Delta \Gamma^q / \Delta m_q = \text{Re}(\Gamma_{12}^q / A_q)$$

new-physics-specific constraints

semileptonic asymmetry:

sensitive to NP effects in both size and phase

$$A_{SL}^s \times 10^2 = -0.17 \pm 0.91$$

D0

Phys.Rev.D82:012003,2010

same-side dilepton charge asymmetry:

admixture of B_s and B_d so sensitive to NP effects in both systems

$$A_{SL}^{\mu\mu} \times 10^3 = -7.9 \pm 2.0$$

D0 arXiv:1106.6308

lifetime τ^{FS} in flavour-specific final states:

average lifetime is a function to the width and the width difference
(independent data sample)

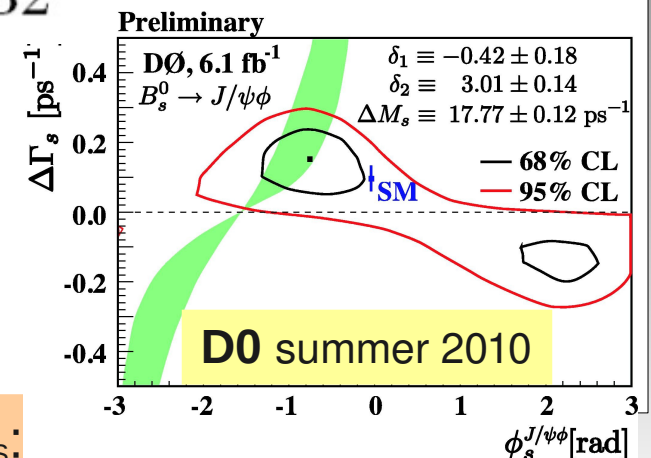
HFAG

$$\tau_{B_s}^{FS} [\text{ps}] = 1.461 \pm 0.032$$

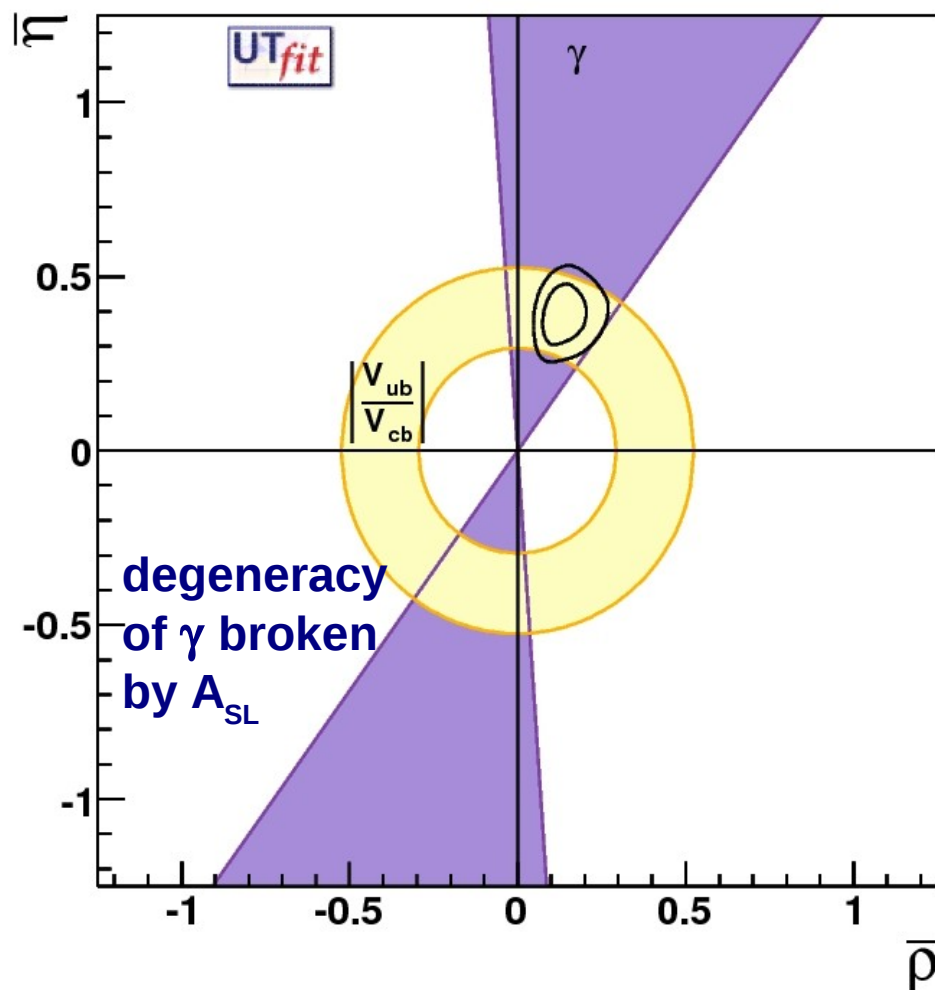
$\phi_s = 2\beta_s$ vs $\Delta\Gamma_s$ from $B_s \rightarrow J/\psi\phi$

angular analysis as a function of proper time
and b-tagging

additional sensitivity from the $\Delta\Gamma_s$ terms



NP analysis results



$$\bar{\rho} = 0.145 \pm 0.045$$

$$\bar{\eta} = 0.389 \pm 0.054$$

SM is

$$\bar{\rho} = 0.129 \pm 0.022$$

$$\bar{\eta} = 0.346 \pm 0.015$$

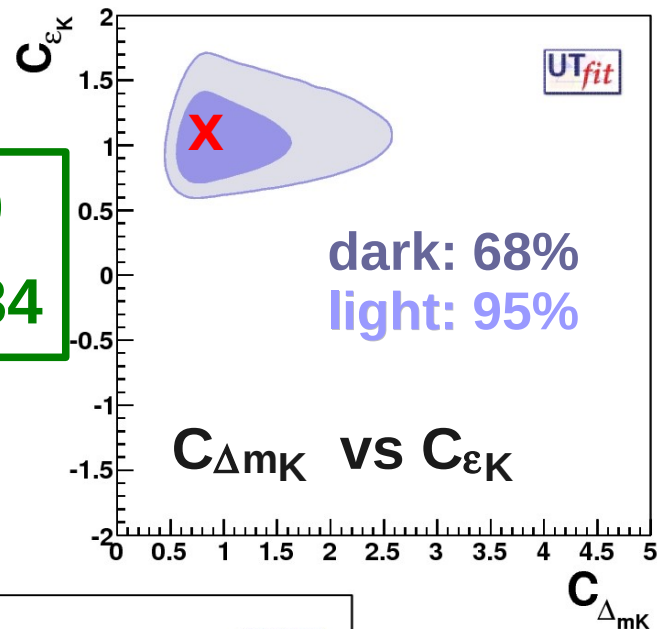
NP parameter results

$$C_{B_d} = 0.82 \pm 0.14$$

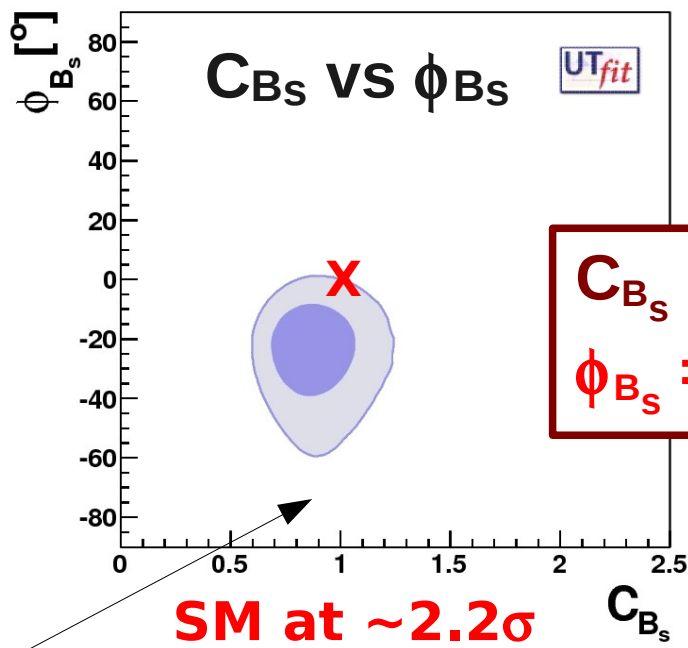
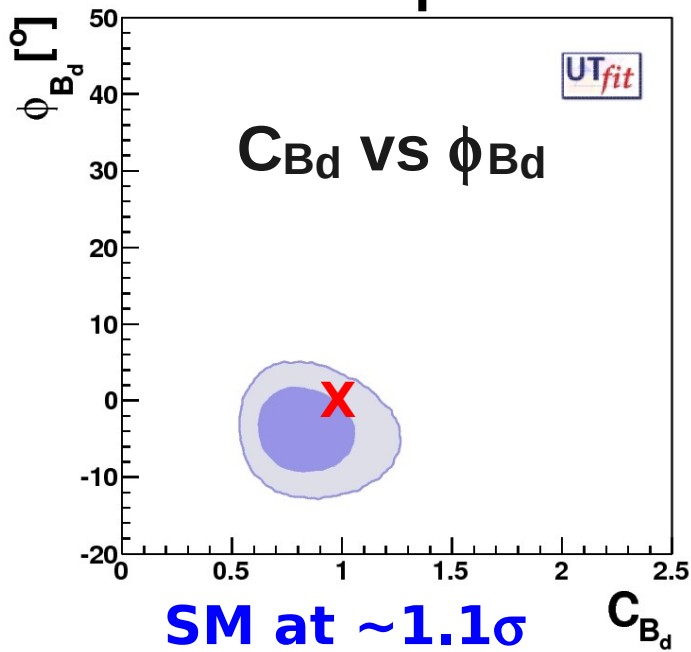
$$\phi_{B_d} = (-3.9 \pm 3.6)^\circ$$

$$C_{\varepsilon_K} = 1.02 \pm 0.19$$

$$C_{\Delta m_K} = 0.96 \pm 0.34$$



X SM expectation



$$C_{B_s} = 0.87 \pm 0.12$$

$$\phi_{B_s} = (-23 \pm 10)^\circ$$

one ambiguity missing as D0 uses the sign of $B_d \rightarrow J/\psi K^*$ strong phase to eliminate it

conclusions

- SM analysis displays good overall consistency but some tension in $\sin 2\beta$, B_k and $B \rightarrow \tau \nu$
- Extraction of SM predictions with different scenarios: semileptonic inclusive vs exclusive
- The tensions pull $|V_{ub}|$ in opposite directions
- General UTA provides a precise determination of CKM parameters and NP contributions to $\Delta F=2$ amplitudes

we would love to use the results from all the collaborations: please release the likelihood!

NEVER EVER
EVER
GIVE UP!



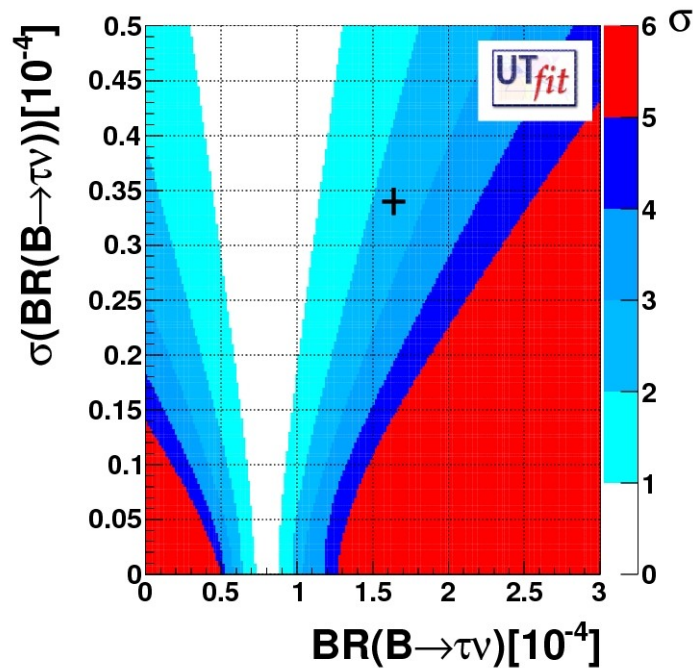


backup

more standard model determinations: $B_d \rightarrow \tau \nu$

current HFAG world average
 $BR(B \rightarrow \tau \nu) = (1.64 \pm 0.34) 10^{-4}$

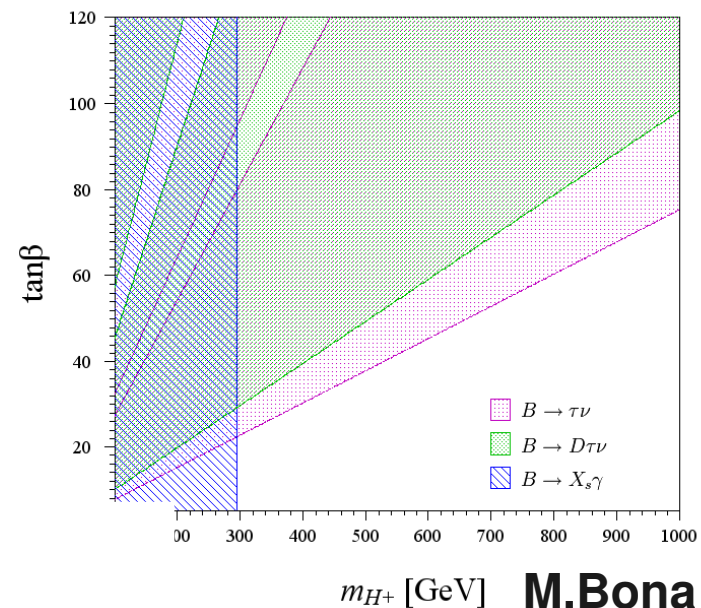
$$\mathcal{B}(B \rightarrow l \nu) = \frac{G_F^2 m_B m_l^2}{8\pi} \left(1 - \frac{m_l^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B$$



indirect determination from UT_b
 $BR(B \rightarrow \tau \nu) = (0.79 \pm 0.08) 10^{-4}$

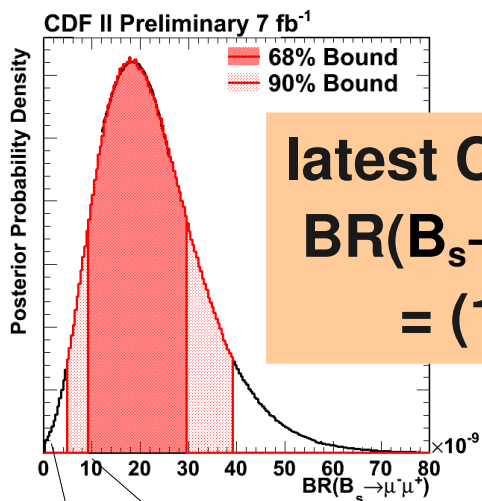
SM prediction enhanced or reduced by factor r_H :

$$R_{2HDM} = \left(1 - \tan^2 \beta \frac{m_B^2}{m_{H^+}^2}\right)^2$$

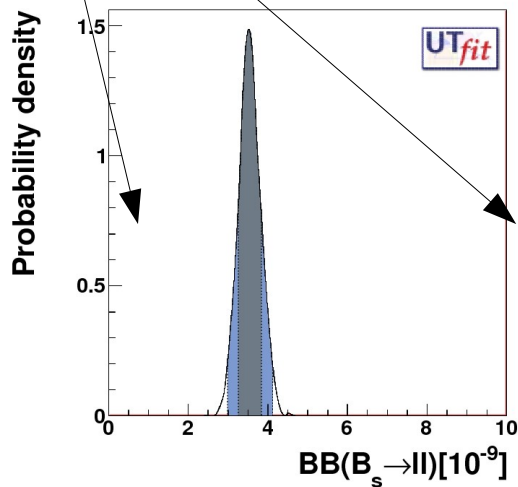


m_{H^+} [GeV] M. Bona et al

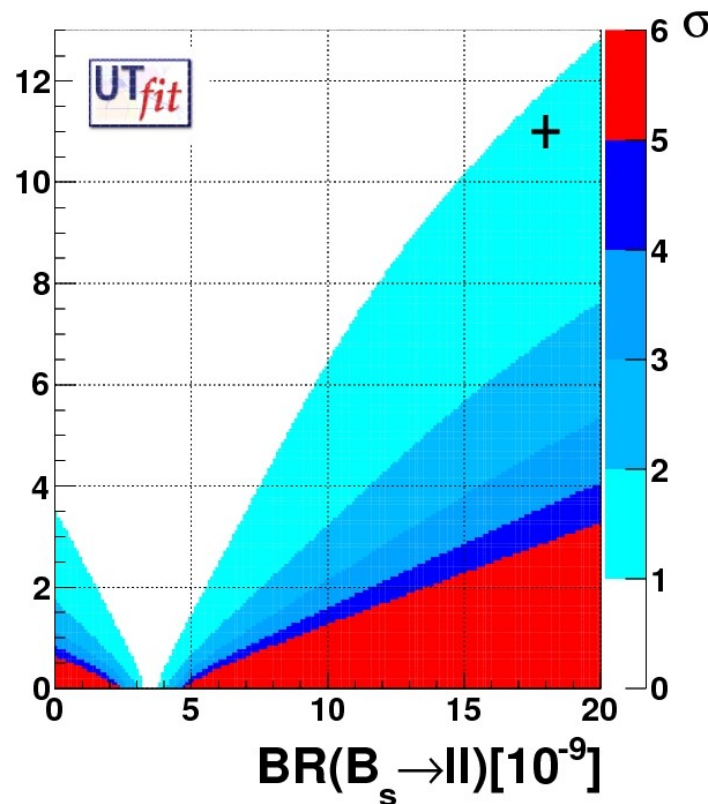
more standard model determinations: $B_s \rightarrow \mu\mu$



latest CDF result:
 $BR(B_s \rightarrow \mu\mu)$
 $= (18^{+11}_{-9}) 10^{-9}$

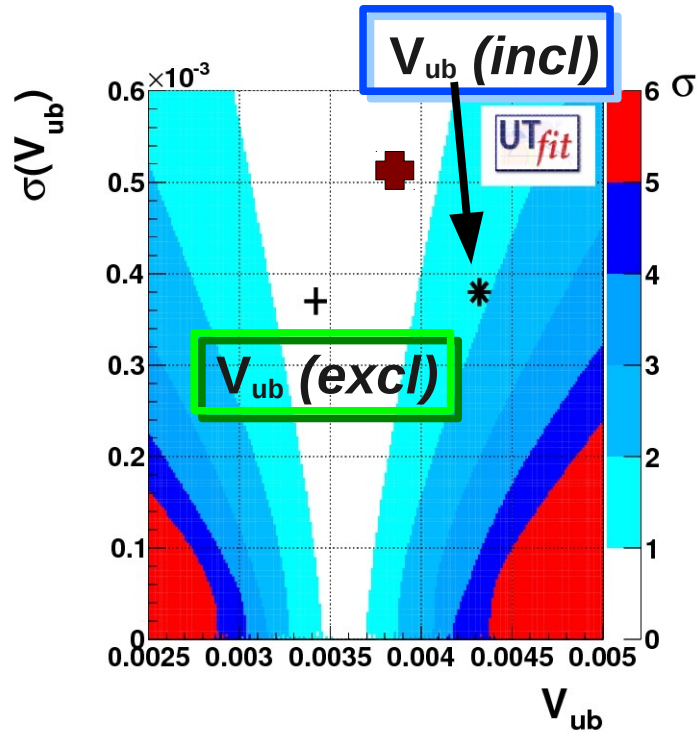


$\sigma(BR(B_s \rightarrow ll)) [10^{-9}]$



indirect determination from UT
 $BR(B_s \rightarrow ll) = (3.54 \pm 0.29) 10^{-9}$

tensions



$$V_{ub_{\text{exp}}} = (3.86 \pm 0.52) \cdot 10^{-3}$$

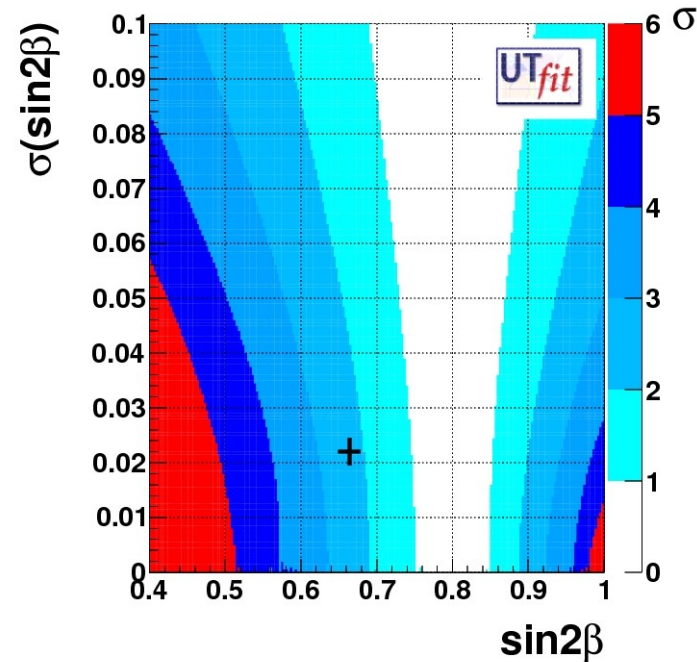
$$V_{ub_{\text{UTfit}}} = (3.57 \pm 0.14) \cdot 10^{-3}$$

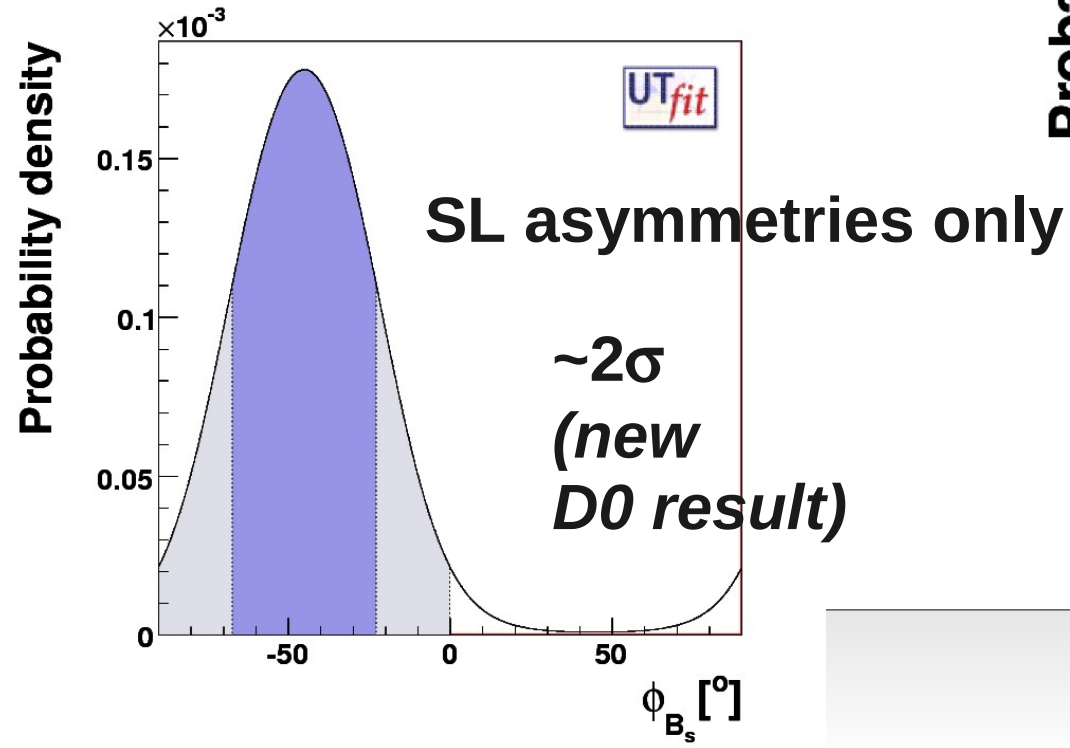
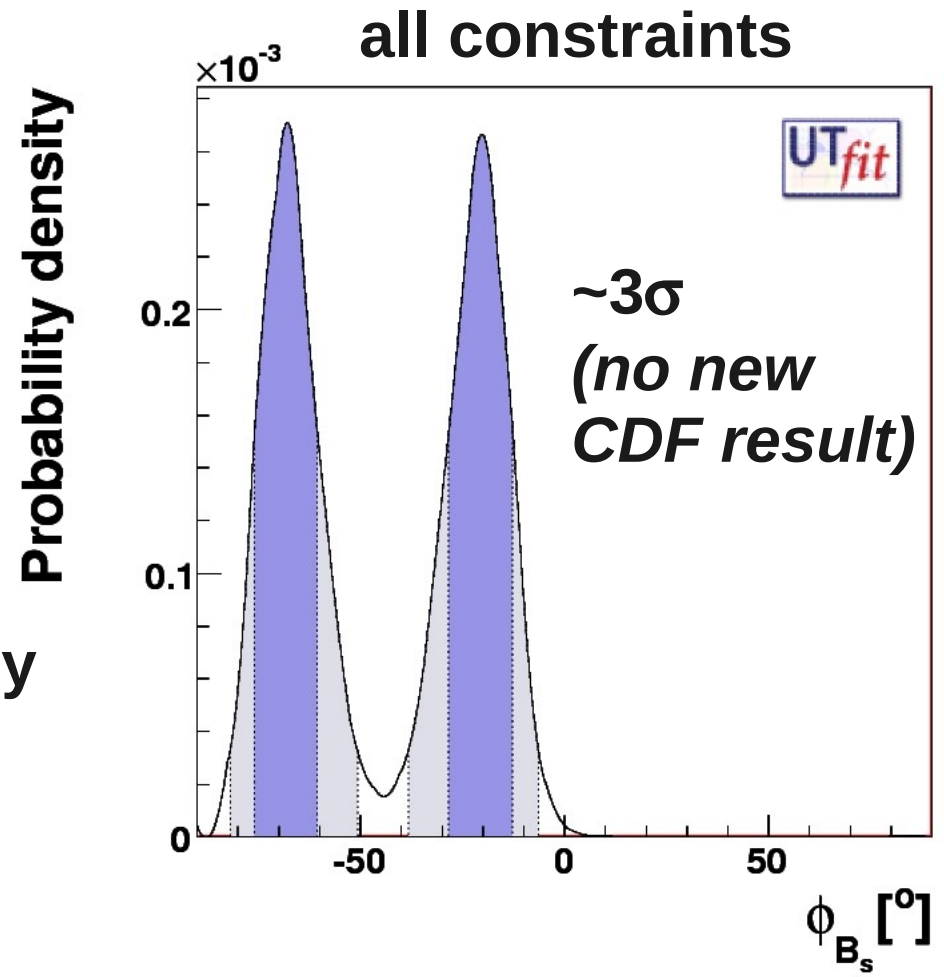
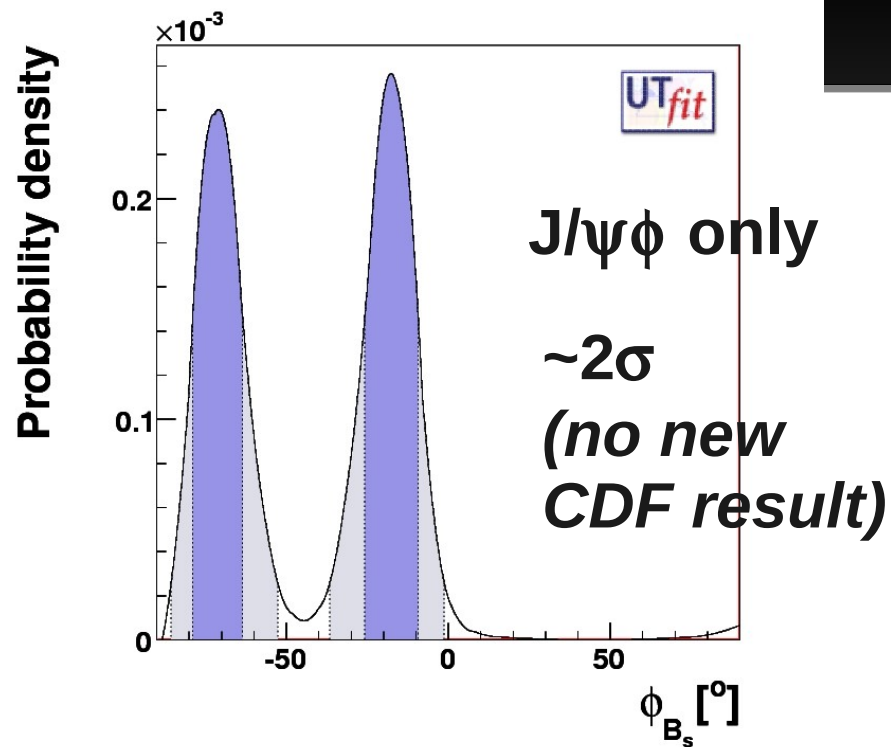
$<1\sigma$ (incl $\sim 1.8\sigma$)

$\sim 2.3\sigma$

$$\sin 2\beta_{\text{exp}} = 0.664 \pm 0.022$$

$$\sin 2\beta_{\text{UTfit}} = 0.803 \pm 0.051$$



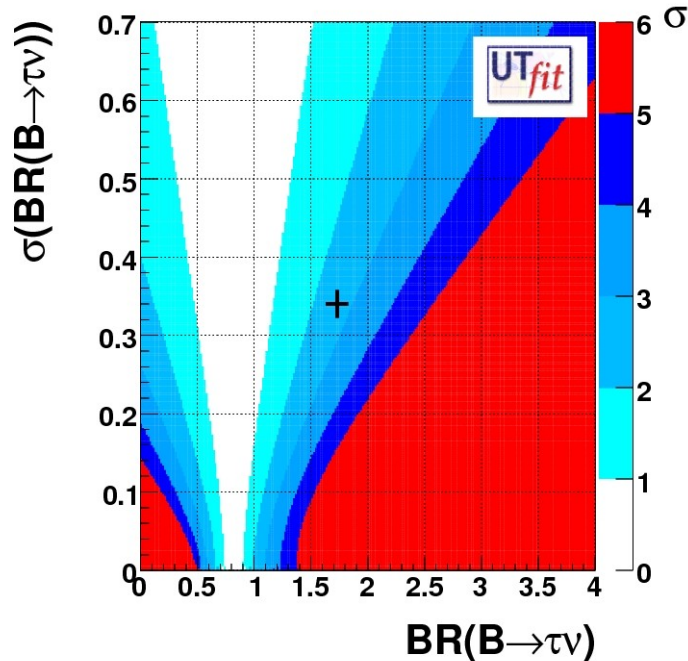


Consider MFV models

Define a Universal Unitarity Triangle using only observables unaffected by MFV-NP:

R_b & angles

Define \bar{BR} as the prediction obtained assuming NO NP effect in the decay amplitude



$$BR(B \rightarrow \tau\nu)_{\text{exp}} = (1.74 \pm 0.34) \cdot 10^{-4}$$

$$BR(B \rightarrow \tau\nu)_{\text{UTfit}} = (0.79 \pm 0.07) \cdot 10^{-4}$$

$\sim 2.7\sigma$

$$R_{\text{UUT}}^{\text{exp}} = 2.1 \pm 0.5$$

where

$$R_{\text{UUT}}^{\text{exp}} = BR_{\text{exp}} / \bar{BR}_{\text{UUT}}$$

to be compared with the $|V_{ub}|$ - and f_B -independent theory calculation of R_{UUT} in specific MFV models

Consider Two Higgs Doublet model II

$$R_{2\text{HDM}} = \left(1 - \tan^2 \beta \frac{m_B^2}{m_{H^+}^2} \right)^2$$

→ bounds on $\tan\beta/m_{H^+}$

Two regions selected:

1. small $\tan\beta/m_{H^+}$: $R < 1$ disfavoured at $\sim 2\sigma$
2. “fine-tuned” region for $\tan\beta/m_{H^+} \sim 0.3$:
positive correction, $R \sim R_{\text{exp}}$ can be obtained

incompatible with semileptonic decays

$$\text{BR}(B \rightarrow D\tau\nu) / \text{BR}(B \rightarrow D\ell\nu) = (49 \pm 10)\%$$

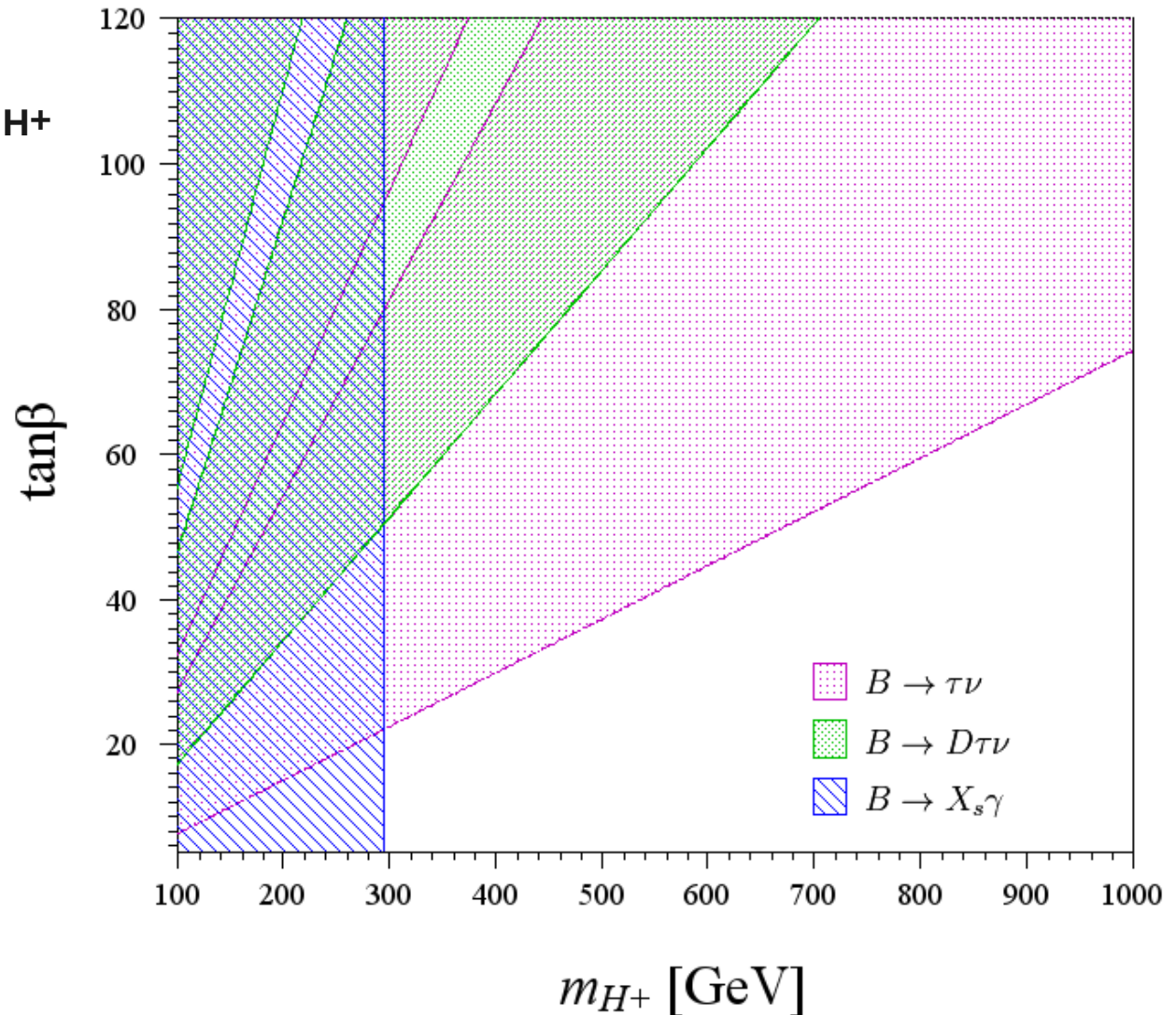
$B \rightarrow X_s \gamma$ gives a lower bound on m_{H^+} :

Consider Two Higgs Doublet model II

$$R_{2\text{HDM}} = \left(1 - \tan^2 \beta \frac{m_B^2}{m_{H^+}^2} \right)^2$$

→ bounds on $\tan\beta/m_{H^+}$

$$\tan \beta < 7.4 \frac{m_{H^+}}{100 \text{ GeV}}$$



At the high scale

new physics enters according to its specific features

At the low scale

use OPE to write the most general effective Hamiltonian. the operators have different chiralities than the SM

NP effects are in the Wilson Coefficients C

NP effects are enhanced

- up to a factor 10 by the values of the matrix elements especially for transitions among quarks of different chiralities
- up to a factor 8 by RGE

$$\mathcal{H}_{\text{eff}}^{\Delta B=2} = \sum_{i=1}^5 C_i Q_i^{bq} + \sum_{i=1}^3 \tilde{C}_i \tilde{Q}_i^{bq}$$

$$Q_1^{q_i q_j} = \bar{q}_{jL}^{\alpha} \gamma_{\mu} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} \gamma^{\mu} q_{iL}^{\beta} ,$$

$$Q_2^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jR}^{\beta} q_{iL}^{\beta} ,$$

$$Q_3^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jR}^{\beta} q_{iL}^{\alpha} ,$$

$$Q_4^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} q_{iR}^{\beta} ,$$

$$Q_5^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jL}^{\beta} q_{iR}^{\alpha} .$$

M. Bona *et al.* (UTfit)

JHEP 0803:049,2008

arXiv:0707.0636

Effective BSM Hamiltonian for $\Delta F=2$ transitions

Most general form of the effective Hamiltonian for $\Delta F=2$ processes

$$\mathcal{H}_{\text{eff}}^{K-\bar{K}} = \sum_{i=1}^5 C_i Q_i^{sd} + \sum_{i=1}^3 \tilde{C}_i \tilde{Q}_i^{sd}$$

$$\mathcal{H}_{\text{eff}}^{B_q-\bar{B}_q} = \sum_{i=1}^5 C_i Q_i^{bq} + \sum_{i=1}^3 \tilde{C}_i \tilde{Q}_i^{bq}$$

The Wilson coefficients C_i have in general the form

$$C_i(\Lambda) = F_i \frac{L_i}{\Lambda^2}$$

Putting bounds on the Wilson coefficients give insights into the NP scale in different NP scenarios that enter through F_i and L_i

F_i : function of the NP flavour couplings

L_i : loop factor (in NP models with no tree-level FCNC)

Λ : NP scale (typical mass of new particles mediating $\Delta F=2$ transitions)

Contribution to the mixing amplitudes

analytic expression for the contribution to the mixing amplitudes

$$\langle \bar{B}_q | \mathcal{H}_{\text{eff}}^{\Delta B=2} | B_q \rangle_i = \sum_{j=1}^5 \sum_{r=1}^5 \left(b_j^{(r,i)} + \eta c_j^{(r,i)} \right) \eta^{a_j} C_i(\Lambda) \langle \bar{B}_q | Q_r^{bq} | B_q \rangle$$

Lattice QCD

arXiv:0707.0636: for "magic numbers" a, b and c , $\eta = \alpha_s(\Lambda)/\alpha_s(m_t)$

analogously for the K system

$$\langle \bar{K}^0 | \mathcal{H}_{\text{eff}}^{\Delta S=2} | K^0 \rangle_i = \sum_{j=1}^5 \sum_{r=1}^5 \left(b_j^{(r,i)} + \eta c_j^{(r,i)} \right) \eta^{a_j} C_i(\Lambda) R_r \langle \bar{K}^0 | Q_r^{sd} | K^0 \rangle$$

to obtain the p.d.f. for the Wilson coefficients $C_i(\Lambda)$ at the new-physics scale, we switch on **one coefficient at a time** in each sector and calculate its value from the result of the NP analysis.

Testing the TeV scale

$$C_i(\Lambda) = \frac{F_i L_i}{\Lambda^2}$$

The dependence of C on Λ changes on flavor structure. we can consider different flavour scenarios:

- **Generic:** $C(\Lambda) = \alpha/\Lambda^2$ $F_i \sim 1$, arbitrary phase
- **NMFV:** $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ $F_i \sim |F_{SM}|$, arbitrary phase
- **MFV:** $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ $F_1 \sim |F_{SM}|$, $F_{i \neq 1} \sim 0$, SM phase

$\alpha (L_i)$ is the coupling among NP and SM

- ⊙ $\alpha \sim 1$ for strongly coupled NP
- ⊙ $\alpha \sim \alpha_w (\alpha_s)$ in case of loop coupling through **weak (strong)** interactions

F_{SM} is the combination of CKM factors for the considered process

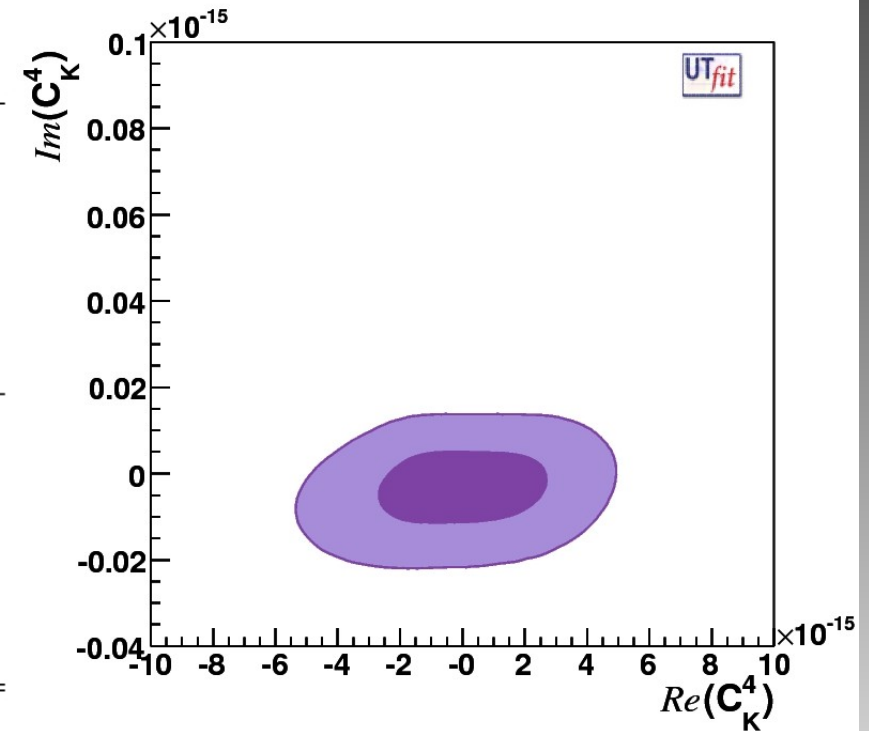
If no NP effect is seen
lower bound on NP scale Λ
if NP is seen
upper bound on NP scale Λ

Results from the Wilson coefficients

the results obtained for the flavour scenarios:

In deriving the lower bounds on the NP scale, we assume $Li = 1$, corresponding to strongly-interacting and/or tree-level NP.

Parameter	95% allowed range (GeV^{-2})	Lower limit on Λ (TeV)	
		for arbitrary NP	for NMFV
$\text{Re}C_K^1$	$[-9.6, 9.6] \cdot 10^{-13}$	$1.0 \cdot 10^3$	0.35
$\text{Re}C_K^2$	$[-1.8, 1.9] \cdot 10^{-14}$	$7.3 \cdot 10^3$	2.0
$\text{Re}C_K^3$	$[-6.0, 5.6] \cdot 10^{-14}$	$4.1 \cdot 10^3$	1.1
$\text{Re}C_K^4$	$[-3.6, 3.6] \cdot 10^{-15}$	$17 \cdot 10^3$	4.0
$\text{Re}C_K^5$	$[-1.0, 1.0] \cdot 10^{-14}$	$10 \cdot 10^3$	2.4
$\text{Im}C_K^1$	$[-4.4, 2.8] \cdot 10^{-15}$	$1.5 \cdot 10^4$	5.6
$\text{Im}C_K^2$	$[-5.1, 9.3] \cdot 10^{-17}$	$10 \cdot 10^4$	28
$\text{Im}C_K^3$	$[-3.1, 1.7] \cdot 10^{-16}$	$5.7 \cdot 10^4$	19
$\text{Im}C_K^4$	$[-1.8, 0.9] \cdot 10^{-17}$	$24 \cdot 10^4$	62
$\text{Im}C_K^5$	$[-5.2, 2.8] \cdot 10^{-17}$	$14 \cdot 10^4$	37

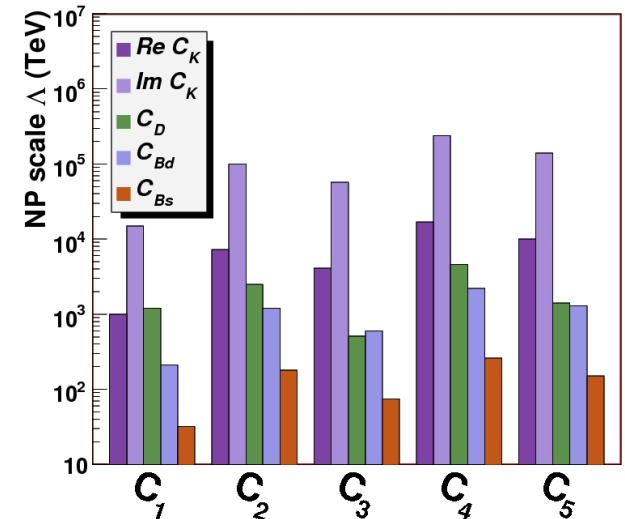


To obtain the lower bound for loop-mediated contributions, one simply multiplies the bounds by $\alpha_s \sim 0.1$ or by $\alpha_w \sim 0.03$.

Upper and lower bound on the scale

Lower bounds on NP scale from K and B_d physics (in TeV at 95% prob.)

Scenario	strong/tree	α_s loop	α_W loop
MFV	5.5	0.5	0.2
NMFV	62	6.2	2
General	24000	2400	800



Upper bounds on NP scale from B_s :

Scenario	strong/tree	α_s loop	α_W loop
NMFV	35	4	2
General	800	80	30

- the **general** case was already problematic (well known flavour puzzle)
- NMFV** has problems with the size of the B_s effect vs the (insufficient) suppression in B_d and (in particular) K mixing
- MFV** is OK for the size of the effects, but the B_s phase cannot be generated