

# The Asymptotic Safety Program for Quantum Gravity

**Frank Saueressig**

*Group for Theoretical High Energy Physics (THEP)  
Institute of Physics*



International Europhysics Conference for High Energy Physics

Grenoble, Rhône-Alpes France, July 21-27 2011

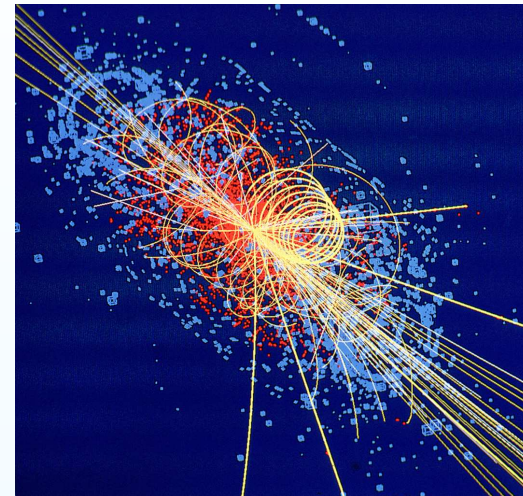
# Introduction

standard model of particle physics:

THE STANDARD MODEL					
	Fermions			Bosons	
Quarks	$u$ up	$c$ charm	$t$ top	$\gamma$ photon	Force carriers
	$d$ down	$s$ strange	$b$ bottom	$Z$ Z boson	
Leptons	$\nu_e$ electron neutrino	$\nu_\mu$ muon neutrino	$\nu_\tau$ tau neutrino	$W$ W boson	
	$e$ electron	$\mu$ muon	$\tau$ tau	$g$ gluon	
	Higgs boson*				

\*Yet to be confirmed

Source: AAAS



- describes: electromagnetic/strong/weak force + interactions with matter
- extremely well tested

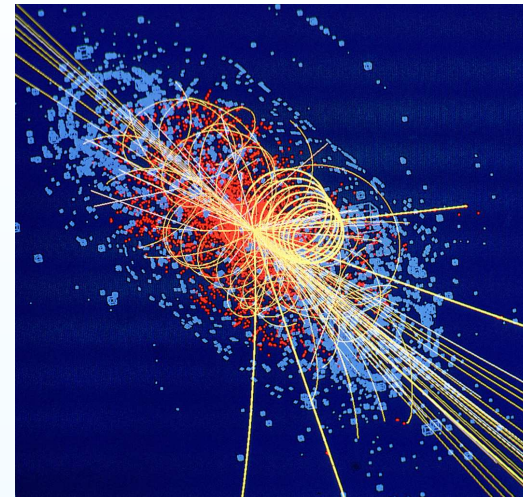
# Introduction

standard model of particle physics:

THE STANDARD MODEL					
Fermions			Bosons		
Quarks	$u$ up	$c$ charm	$t$ top	$\gamma$ photon	Force carriers
	$d$ down	$s$ strange	$b$ bottom	$Z$ Z boson	
Leptons	$\nu_e$ electron neutrino	$\nu_\mu$ muon neutrino	$\nu_\tau$ tau neutrino	$W$ W boson	
	$e$ electron	$\mu$ muon	$\tau$ tau	$g$ gluon	
	Higgs boson*				

\*Yet to be confirmed

Source: AAAS



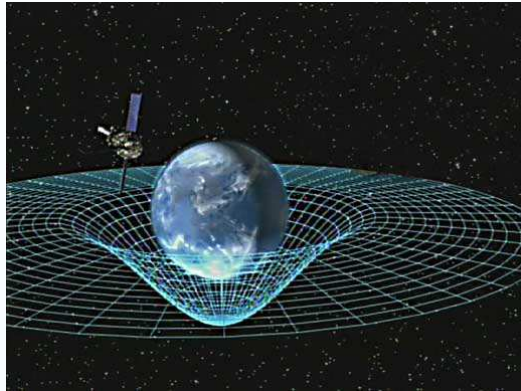
- describes: electromagnetic/strong/weak force + interactions with matter
- extremely well tested

theoretical basis: quantum field theory

- includes only relevant and marginal couplings  
⇒ renormalizable quantum field theory

# Introduction

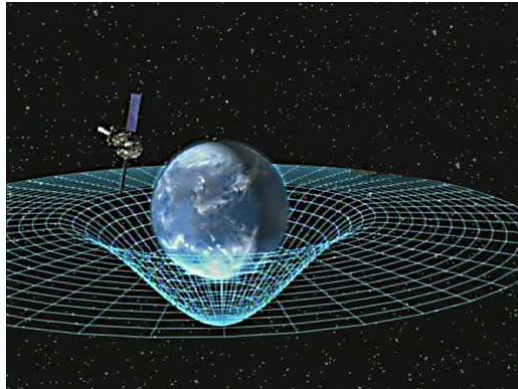
General Relativity:



- describes: gravity + interactions with matter
- extremely well tested

# Introduction

General Relativity:



- describes: gravity + interactions with matter
- extremely well tested

theoretical basis: classical theory

$$\underbrace{R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R}_{\text{space-time curvature}} = \underbrace{-\Lambda g_{\mu\nu} + 8\pi G_N T_{\mu\nu}}_{\text{matter content}}$$

- Newton constant  $G_N$  has **negative** mass-dimension  
 $\implies$  perturbatively non-renormalizable quantum field theory

# General Relativity: perturbatively non-renormalizable

perturbative quantization of General Relativity:

- $G_N$  has negative mass-dimension:
  - infinite number of counterterms
  - General Relativity is perturbatively non-renormalizable

Possible conclusions:

- a) Treat General Relativity as **effective field theory**:
  - compute corrections in  $E^2/M_{\text{Pl}}^2 \ll 1$  (independent of UV-completion)
  - breaks down at  $E^2 \approx M_{\text{Pl}}^2$

# General Relativity: perturbatively non-renormalizable

perturbative quantization of General Relativity:

- $G_N$  has negative mass-dimension:
  - infinite number of counterterms
  - General Relativity is perturbatively non-renormalizable

Possible conclusions:

- a) Treat General Relativity as **effective field theory**:
  - compute corrections in  $E^2/M_{\text{Pl}}^2 \ll 1$  (independent of UV-completion)
  - breaks down at  $E^2 \approx M_{\text{Pl}}^2$
- b) UV-completion requires new physics:
  - supersymmetry, extra dimensions, ...
  - possibly: extension of QFT-framework

# General Relativity: perturbatively non-renormalizable

perturbative quantization of General Relativity:

- $G_N$  has negative mass-dimension:
  - infinite number of counterterms
  - General Relativity is perturbatively non-renormalizable

Possible conclusions:

- a) Treat General Relativity as **effective field theory**:
  - compute corrections in  $E^2/M_{\text{Pl}}^2 \ll 1$  (independent of UV-completion)
  - breaks down at  $E^2 \approx M_{\text{Pl}}^2$
- b) UV-completion requires new physics:
  - supersymmetry, extra dimensions, ...
  - possibly: extension of QFT-framework
- c) **Gravity makes sense as Quantum Field Theory**:
  - UV-completion requires going beyond perturbation theory



# General Relativity: perturbatively non-renormalizable

perturbative quantization of General Relativity:

- $G_N$  has negative mass-dimension:
  - infinite number of counterterms
  - General Relativity is perturbatively non-renormalizable

Possible conclusions:

a) Treat General Relativity as **effective field theory**:

- compute corrections in  $E^2/M_{\text{Pl}}^2 \ll 1$  (independent of UV-completion)
- breaks down at  $E^2 \approx M_{\text{Pl}}^2$

b) UV-completion requires new physics:

- supersymmetry, extra dimensions, ...
- possibly: extension of QFT-framework

c) **Gravity makes sense as Quantum Field Theory**:

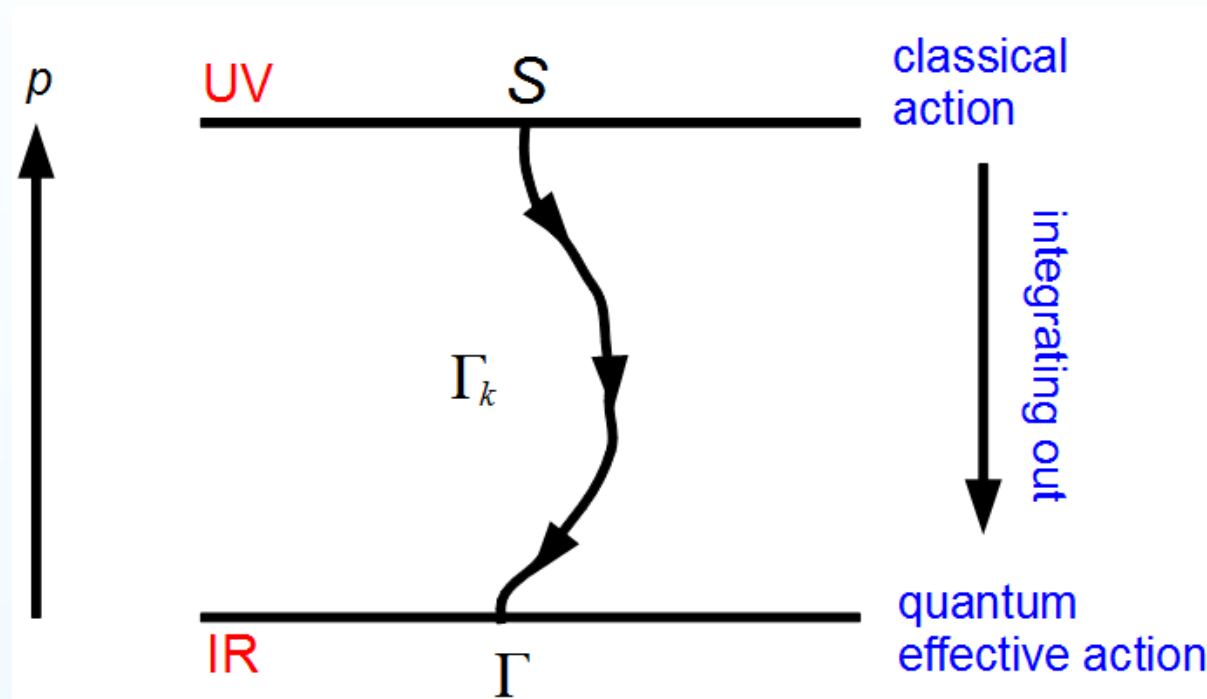
- UV-completion requires going beyond perturbation theory

# Wilsonian renormalization and asymptotic safety

## basic concepts

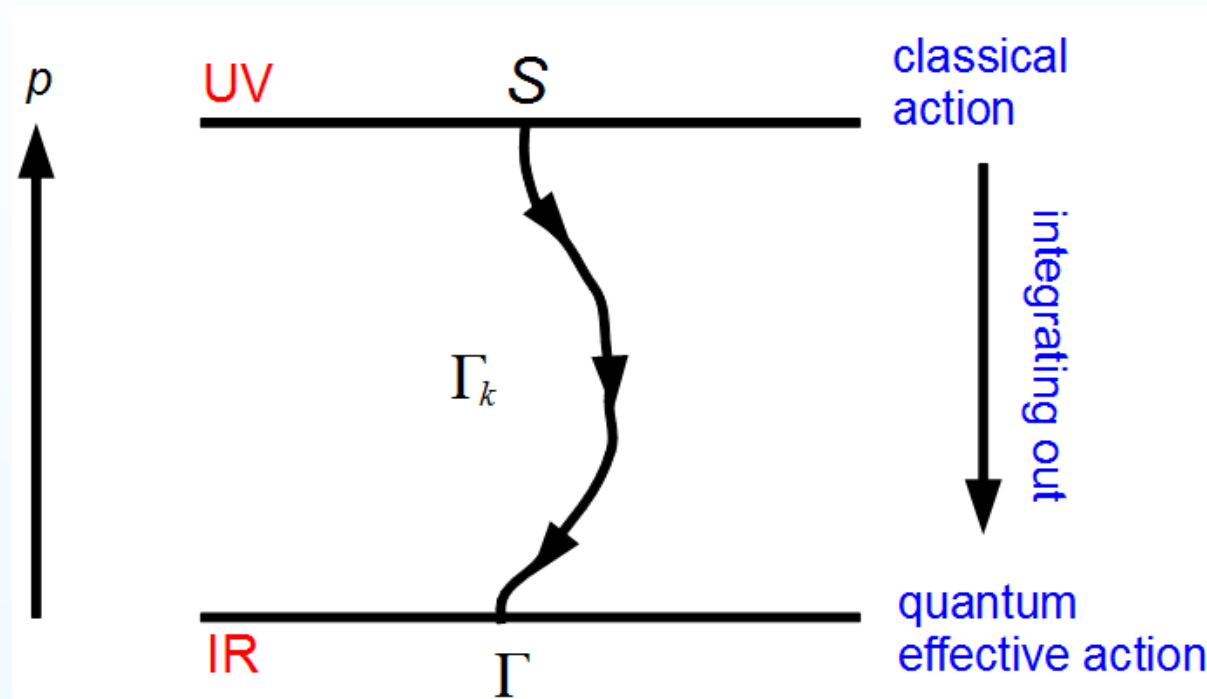
# Wilson's modern picture of renormalization

central idea: integrate out quantum fluctuations shell-by-shell in momentum-space



# Wilson's modern picture of renormalization

central idea: integrate out quantum fluctuations shell-by-shell in momentum-space



implementation:

- action with scale-dependent couplings ( $G_N, \Lambda, \dots$ ):  $g_i(k)$
- scale-dependence governed by  $\beta$ -functions:  $k\partial_k g_i = \beta_{g_i}(\{g_i\})$

# Ensuring good UV-behavior: fixed points of the RG-flow

amplitudes depend on dimensionless couplings only

- RG-flow for dimensionless running couplings:  $g_i(k)$

Fixed points  $g_i^*$ :

- $\beta$ -functions vanish:

$$\beta_{g_i}(\{g_i^*\}) \stackrel{!}{=} 0$$

$g_i^*$  remain finite

- RG-trajectory captured by fixed point in UV:

$\implies$  physical quantities remain free of unphysical divergences

# Ensuring good UV-behavior: fixed points of the RG-flow

amplitudes depend on dimensionless couplings only

- RG-flow for dimensionless running couplings:

$$g_i(k)$$

Fixed points  $g_i^*$ :

- $\beta$ -functions vanish:

$$\beta_{g_i}(\{g_i^*\}) \stackrel{!}{=} 0$$

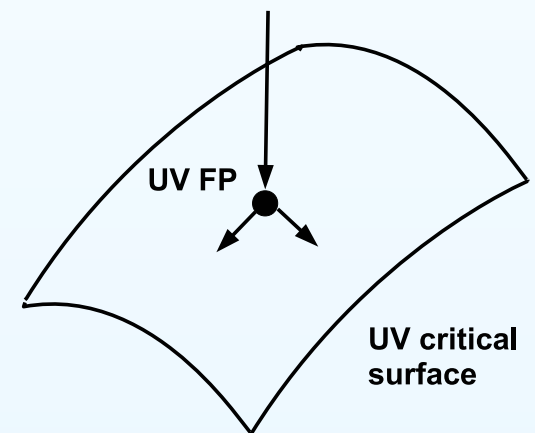
$g_i^*$  remain finite

- RG-trajectory captured by fixed point in UV:

$\implies$  physical quantities remain free of unphysical divergences

Concepts associated with UV-fixed points:

- trajectories emanating from fixed point in UV  
 $\equiv$  span UV critical surface
- predictivity:  
 $\equiv$  UV critical surface has finite dimension



# Renormalization: asymptotic freedom and asymptotic safety

Wilsonian formulation:

- UV fixed points allow two classes of renormalizable Quantum Field Theories
- **Gaussian Fixed Point (GFP):**
  - **perturbatively renormalizable field theories**
  - UV-limit: free theory
  - asymptotic freedom (example: QCD)

# Renormalization: asymptotic freedom and asymptotic safety

Wilsonian formulation:

- UV fixed points allow two classes of renormalizable Quantum Field Theories
- **Gaussian Fixed Point (GFP):**
  - **perturbatively renormalizable field theories**
  - UV-limit: free theory
  - asymptotic freedom (example: QCD)
- **non-Gaussian Fixed Point (NGFP):**
  - **non-perturbatively renormalizable field theories**
  - UV-limit: interacting theory
  - asymptotic safety



# Renormalization: asymptotic freedom and asymptotic safety

Wilsonian formulation:

- UV fixed points allow two classes of renormalizable Quantum Field Theories
- **Gaussian Fixed Point (GFP):**
  - **perturbatively renormalizable field theories**
  - UV-limit: free theory
  - asymptotic freedom (example: QCD)
- **non-Gaussian Fixed Point (NGFP):**
  - **non-perturbatively renormalizable field theories**
  - UV-limit: interacting theory
  - asymptotic safety

Wilsonian picture: generalization of perturbative renormalization

**asymptotic safety** as predictive as **asymptotic freedom**

# Examples: Asymptotically Safe Theories

Theories with non-Gaussian UV fixed point

- $O(N)$ -sigma model ( $d = 2 + \epsilon$ )

[Brézin, Zinn-Justin '76]

- critical exponents of Heisenberg ferromagnets

- Gross-Neveu model ( $d = 2 + \epsilon$ )

[Gawedzki, Kupiainen '85]

- Grosse-Wulkenhaar model (non-commutative  $\phi^4$ -theory)

[Grosse, Wulkenhaar '05; Disertori, Gurau, Magnen, Rivasseau '07]

- Gravity in  $2 + \epsilon$  dimensions

[Christensen, Duff; Gastmans, Kallosh, Truffin '78]

# Examples: Asymptotically Safe Theories

Theories with non-Gaussian UV fixed point

- $O(N)$ -sigma model ( $d = 2 + \epsilon$ )

[Brézin, Zinn-Justin '76]

- critical exponents of Heisenberg ferromagnets

- Gross-Neveu model ( $d = 2 + \epsilon$ )

[Gawedzki, Kupiainen '85]

- Grosse-Wulkenhaar model (non-commutative  $\phi^4$ -theory)

[Grosse, Wulkenhaar '05; Disertori, Gurau, Magnen, Rivasseau '07]

- Gravity in  $2 + \epsilon$  dimensions

[Christensen, Duff; Gastmans, Kallosh, Truffin '78]

Weinberg's asymptotic safety conjecture (1979):

gravity in  $d = 4$  has non-Gaussian UV fixed point

# Renormalizing gravity

Wilsonian formulation:

- UV fixed points allow two classes of renormalizable Quantum Field Theories
- **Gaussian Fixed Point (GFP):**
  - perturbatively renormalizable field theories
  - UV-limit: free theory
  - asymptotic freedom
- **non-Gaussian Fixed Point (NGFP):**
  - non-perturbatively renormalizable field theories
  - UV-limit: interacting theory
  - asymptotic safety



Gravity

Weinberg's asymptotic safety conjecture (1979):

gravity in  $d = 4$  has non-Gaussian UV fixed point

Testing asymptotic safety:

Functional Renormalization Group Equations (FRG)

Causal Dynamical Triangulations (CDT)

# Functional Renormalization Group Equation for gravity

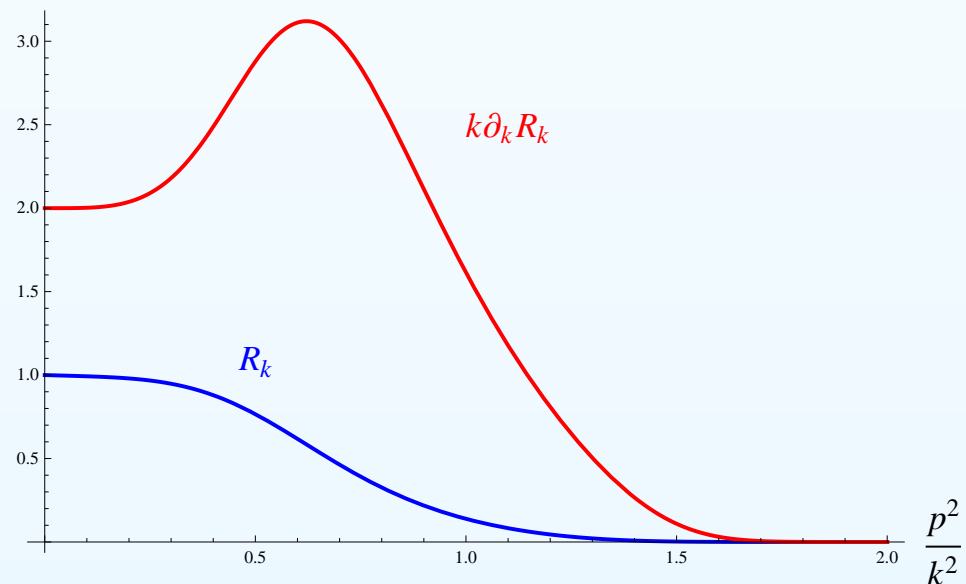
[C. Wetterich, Phys. Lett. B301 (1993) 90]

[M. Reuter, Phys. Rev. D 57 (1998) 971, hep-th/9605030]

scale-dependence of  $\Gamma_k$  governed by exact RG equation

$$k\partial_k\Gamma_k[\phi, \bar{\phi}] = \frac{1}{2}\text{Tr} \left[ \left( \frac{\delta^2\Gamma_k}{\delta\phi\delta\phi} + \mathcal{R}_k \right)^{-1} k\partial_k\mathcal{R}_k \right]$$

- $\mathcal{R}_k(p^2) =$  IR momentum-cutoff at scale  $k$



# Functional Renormalization Group Equation for gravity

[C. Wetterich, Phys. Lett. B301 (1993) 90]

[M. Reuter, Phys. Rev. D 57 (1998) 971, hep-th/9605030]

scale-dependence of  $\Gamma_k$  governed by exact RG equation

$$k\partial_k\Gamma_k[\phi, \bar{\phi}] = \frac{1}{2}\text{Tr} \left[ \left( \frac{\delta^2\Gamma_k}{\delta\phi\delta\phi} + \mathcal{R}_k \right)^{-1} k\partial_k\mathcal{R}_k \right]$$

limits of the RG-flow:

- $k = \Lambda$ : initial (boundary) condition  $\Gamma_{k=\Lambda} = \Gamma_\Lambda$
- $k = 0$ : all quantum fluctuations integrated out  $\Gamma_{k=0} = \Gamma$

$$\Gamma = \Gamma_\Lambda + \lim_{k \rightarrow 0} \int_\Lambda^k d\hat{k} \partial_{\hat{k}} \Gamma_{\hat{k}} \left[ \Gamma_{\hat{k}}^{(2)}, \mathcal{R}_{\hat{k}} \right]$$

in between:

- regulator ensures finiteness of flow

# Functional Renormalization Group Equation for gravity

[C. Wetterich, Phys. Lett. B301 (1993) 90]

[M. Reuter, Phys. Rev. D 57 (1998) 971, hep-th/9605030]

scale-dependence of  $\Gamma_k$  governed by exact RG equation

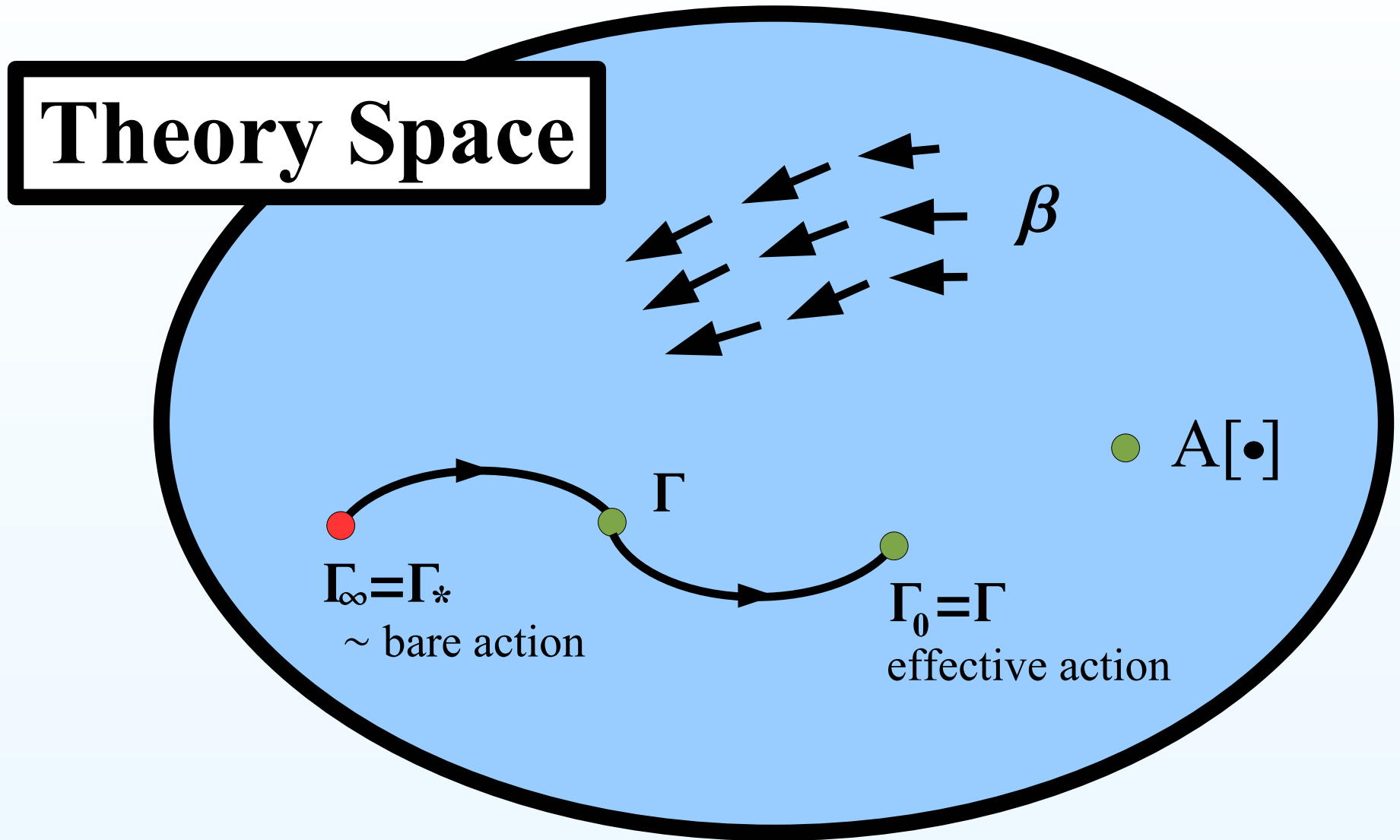
$$k\partial_k\Gamma_k[\phi, \bar{\phi}] = \frac{1}{2}\text{Tr} \left[ \left( \frac{\delta^2\Gamma_k}{\delta\phi\delta\phi} + \mathcal{R}_k \right)^{-1} k\partial_k\mathcal{R}_k \right]$$

renormalizability:

- if  $\Gamma_{\Lambda \rightarrow \infty} = \Gamma_*$  exists,  $\Gamma_*$  qualifies as fundamental theory
  - perturbatively renormalizable theory:  $\Gamma_*$  is free theory (e.g. QCD)
  - non-perturbatively renormalizable theory:  $\Gamma_*$  is interacting
  - non-renormalizable:  $\Gamma_*$  does not exist
- predictivity: provided by fixed point



# Theory space underlying the Functional Renormalization Group



## Non-perturbative approximation: derivative expansion of $\Gamma_k$

- caveat: FRGE cannot be solved exactly

$\Leftrightarrow$  gravity: need non-perturbative approximation scheme

## Non-perturbative approximation: derivative expansion of $\Gamma_k$

- caveat: FRGE cannot be solved exactly  
 $\iff$  gravity: need non-perturbative approximation scheme
- expand  $\Gamma_k$  in derivatives and truncate series:

$$\Gamma_k[\Phi] = \sum_{i=1}^N \bar{u}_i(k) \mathcal{O}_i[\Phi]$$

$\implies$  substitute into FRGE

$\implies$  projection of flow gives  $\beta$ -functions for running couplings

$$k\partial_k \bar{u}_i(k) = \beta_i(\bar{u}_i; k)$$

# Non-perturbative approximation: derivative expansion of $\Gamma_k$

- caveat: FRGE cannot be solved exactly  
 $\iff$  gravity: need non-perturbative approximation scheme
- expand  $\Gamma_k$  in derivatives and truncate series:

$$\Gamma_k[\Phi] = \sum_{i=1}^N \bar{u}_i(k) \mathcal{O}_i[\Phi]$$

$\implies$  substitute into FRGE

$\implies$  projection of flow gives  $\beta$ -functions for running couplings

$$k \partial_k \bar{u}_i(k) = \beta_i(\bar{u}_i; k)$$

- testing the reliability:
  - within a given truncation:  
cutoff-scheme dependence of physical quantities (= vary  $\mathcal{R}_k$ )
  - stability of results within extended truncations

Letting things flow

The Einstein-Hilbert truncation

# The Einstein-Hilbert truncation: setup

Einstein-Hilbert truncation: two running couplings:  $G(k), \Lambda(k)$

$$\Gamma_k = \frac{1}{16\pi G(k)} \int d^4x \sqrt{g} [-R + 2\Lambda(k)] + S^{\text{gf}} + S^{\text{gh}}$$

- project flow onto  $G$ - $\Lambda$ -plane

# The Einstein-Hilbert truncation: setup

Einstein-Hilbert truncation: two running couplings:  $G(k), \Lambda(k)$

$$\Gamma_k = \frac{1}{16\pi G(k)} \int d^4x \sqrt{g} [-R + 2\Lambda(k)] + S^{\text{gf}} + S^{\text{gh}}$$

- project flow onto  $G$ - $\Lambda$ -plane

explicit  $\beta$ -functions for dimensionless couplings  $g_k := k^2 G(k)$ ,  $\lambda_k := \Lambda(k) k^{-2}$

- Particular choice of  $\mathcal{R}_k$  (optimized cutoff)

$$k \partial_k g_k = (\eta_N + 2) g_k,$$

$$k \partial_k \lambda_k = -(2 - \eta_N) \lambda_k - \frac{g_k}{2\pi} \left[ 5 \frac{1}{1-2\lambda_k} - 4 - \frac{5}{6} \frac{1}{1-2\lambda_k} \eta_N \right]$$

- anomalous dimension of Newton's constant:

$$\eta_N = \frac{g B_1}{1 - g B_2}$$

$$B_1 = \frac{1}{3\pi} \left[ 5 \frac{1}{1-2\lambda} - 9 \frac{1}{(1-2\lambda)^2} - 7 \right], \quad B_2 = -\frac{1}{12\pi} \left[ 5 \frac{1}{1-2\lambda} + 6 \frac{1}{(1-2\lambda)^2} \right]$$

## Einstein-Hilbert truncation: Fixed Point structure

$\beta$ -functions for  $g_k := k^2 G(k)$ ,  $\lambda_k := \Lambda(k)k^{-2}$

$$k\partial_k g_k = (\eta_N + 2)g_k,$$

$$k\partial_k \lambda_k = -(2 - \eta_N) \lambda_k - \frac{g_k}{2\pi} \left[ 5 \frac{1}{1-2\lambda_k} - 4 - \frac{5}{6} \frac{1}{1-2\lambda_k} \eta_N \right]$$

microscopic theory  $\iff$  fixed points of the  $\beta$ -functions

$$\beta_g(g^*, \lambda^*) = 0, \quad \beta_\lambda(g^*, \lambda^*) = 0$$



# Einstein-Hilbert truncation: Fixed Point structure

$\beta$ -functions for  $g_k := k^2 G(k)$ ,  $\lambda_k := \Lambda(k)k^{-2}$

$$k\partial_k g_k = (\eta_N + 2)g_k,$$

$$k\partial_k \lambda_k = -(2 - \eta_N)\lambda_k - \frac{g_k}{2\pi} \left[ 5 \frac{1}{1-2\lambda_k} - 4 - \frac{5}{6} \frac{1}{1-2\lambda_k} \eta_N \right]$$

microscopic theory  $\iff$  fixed points of the  $\beta$ -functions

$$\beta_g(g^*, \lambda^*) = 0, \quad \beta_\lambda(g^*, \lambda^*) = 0$$

- Gaussian Fixed Point:
  - at  $g^* = 0, \lambda^* = 0 \iff$  free theory
  - saddle point in the  $g$ - $\lambda$ -plane

# Einstein-Hilbert truncation: Fixed Point structure

$\beta$ -functions for  $g_k := k^2 G(k)$ ,  $\lambda_k := \Lambda(k)k^{-2}$

$$k\partial_k g_k = (\eta_N + 2)g_k,$$

$$k\partial_k \lambda_k = -(2 - \eta_N) \lambda_k - \frac{g_k}{2\pi} \left[ 5 \frac{1}{1-2\lambda_k} - 4 - \frac{5}{6} \frac{1}{1-2\lambda_k} \eta_N \right]$$

microscopic theory  $\iff$  fixed points of the  $\beta$ -functions

$$\beta_g(g^*, \lambda^*) = 0, \quad \beta_\lambda(g^*, \lambda^*) = 0$$

- Gaussian Fixed Point:
  - at  $g^* = 0, \lambda^* = 0 \iff$  free theory
  - saddle point in the  $g$ - $\lambda$ -plane
- non-Gaussian Fixed Point ( $\eta_N^* = -2$ ):
  - at  $g^* > 0, \lambda^* > 0 \iff$  “interacting” theory
  - UV attractive in  $g_k, \lambda_k$

# Einstein-Hilbert truncation: Fixed Point structure

$\beta$ -functions for  $g_k := k^2 G(k)$ ,  $\lambda_k := \Lambda(k)k^{-2}$

$$k\partial_k g_k = (\eta_N + 2)g_k,$$

$$k\partial_k \lambda_k = -(2 - \eta_N)\lambda_k - \frac{g_k}{2\pi} \left[ 5 \frac{1}{1-2\lambda_k} - 4 - \frac{5}{6} \frac{1}{1-2\lambda_k} \eta_N \right]$$

microscopic theory  $\iff$  fixed points of the  $\beta$ -functions

$$\beta_g(g^*, \lambda^*) = 0, \quad \beta_\lambda(g^*, \lambda^*) = 0$$

- Gaussian Fixed Point:
  - at  $g^* = 0, \lambda^* = 0 \iff$  free theory
  - saddle point in the  $g$ - $\lambda$ -plane
- non-Gaussian Fixed Point ( $\eta_N^* = -2$ ):
  - at  $g^* > 0, \lambda^* > 0 \iff$  “interacting” theory
  - UV attractive in  $g_k, \lambda_k$

Asymptotic safety: non-Gaussian Fixed Point is UV completion for gravity

## Einstein-Hilbert truncation: Stability properties

Ref.	$g^*$	$\lambda^*$	$g^*\lambda^*$	$\theta' \pm i\theta''$	gauge	$\mathcal{R}_k$
BMS	0.902	0.109	0.099	$2.52 \pm 1.78i$	geometric	II, opt
RS	0.403	0.330	0.133	$1.94 \pm 3.15i$	harmonic	I, sharp
LR	0.272	0.348	0.095	$1.55 \pm 3.84i$	harmonic	I, exp
	0.344	0.339	0.117	$1.86 \pm 4.08i$	Landau	I, exp
L	1.17	0.25	0.295	$1.67 \pm 4.31i$	Landau	I, opt
CPR	0.707	0.193	0.137	$1.48 \pm 3.04i$	harmonic	I, opt
	0.556	0.092	0.051	$2.43 \pm 1.27i$	harmonic	II, opt
	0.332	0.274	0.091	$1.75 \pm 2.07i$	harmonic	III, opt

BMS: Benedetti, Machado, Saueressig, 2009.

RS: Reuter, Saueressig, 2002.

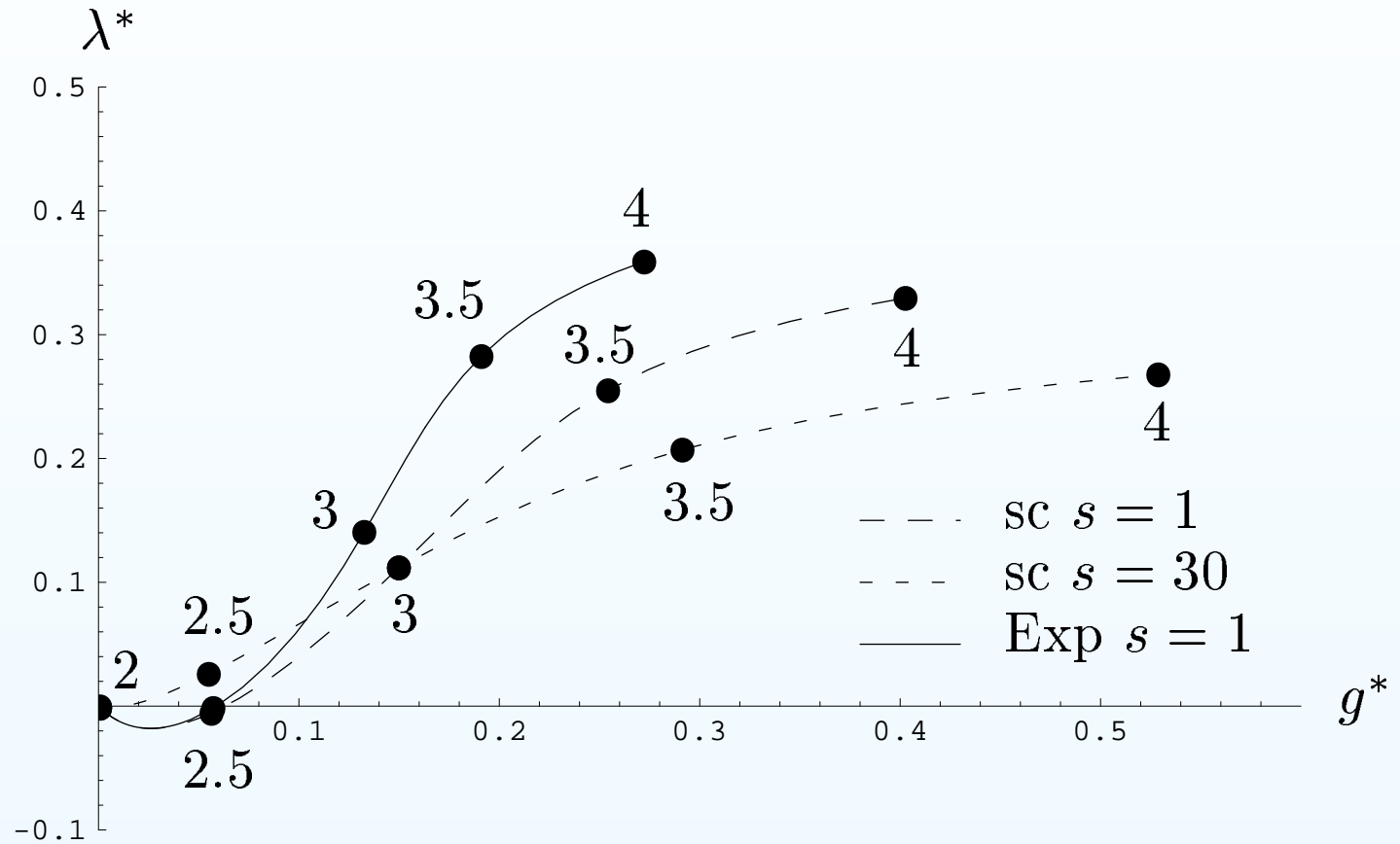
LR: Lauscher, Reuter, 2002.

L: Litim, 2004.

CPR: Codello, Percacci, Rahmede, 2009.

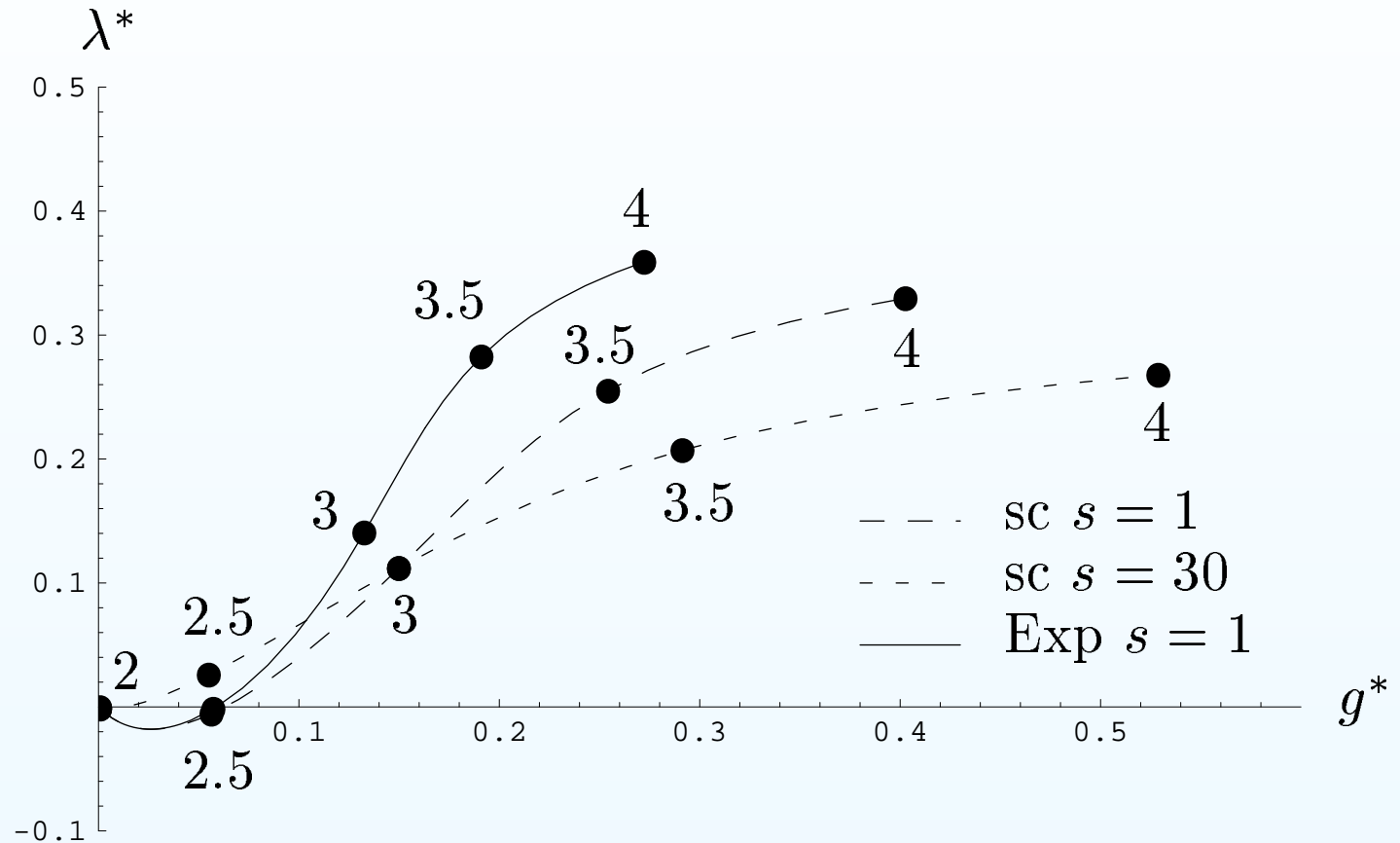
# Einstein-Hilbert truncation: NGFP in $d = 2 + \epsilon$

$\beta$ -functions continuous in  $d \iff$  reproduce perturbative fix point in  $d = 2 + \epsilon$



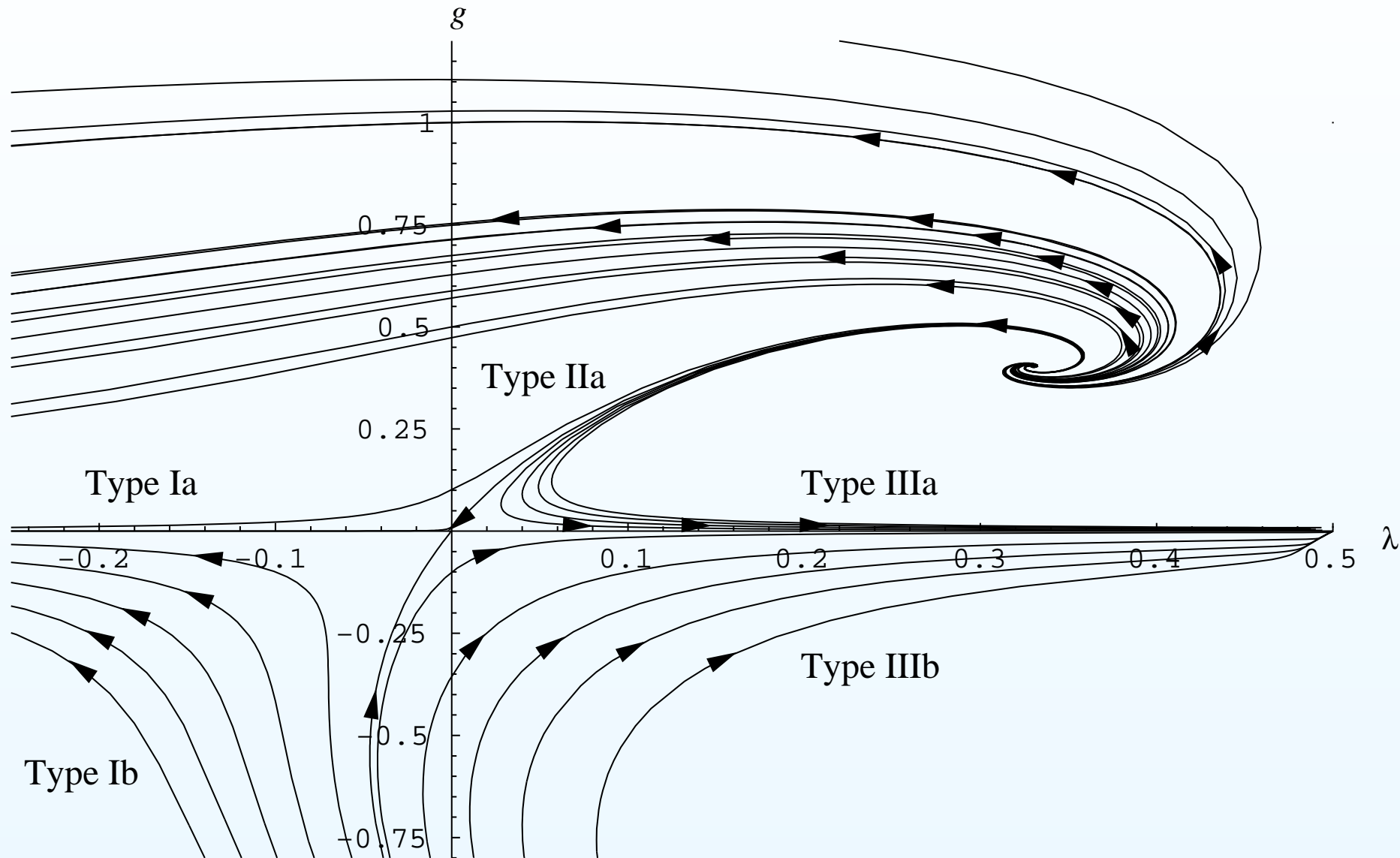
# Einstein-Hilbert truncation: NGFP in $d = 2 + \epsilon$

$\beta$ -functions continuous in  $d \iff$  reproduce perturbative fix point in  $d = 2 + \epsilon$



NGFP in  $d = 4 \iff$  analytic continuation of NGFP in  $d = 2 + \epsilon$

# Einstein-Hilbert-truncation: the phase diagram



where FRG and CDT meet  
spectral dimension of space-time



# Spectral dimension for classical manifolds

- Heat-equation: diffusion of scalar test particle on manifold with metric  $g$

$$\partial_T K_g(x, x'; T) = \Delta_g K_g(x, x'; T)$$

- define average return probability

$$\begin{aligned} P_g(T) &\equiv \frac{1}{V} \int d^d x \sqrt{g(x)} K_g(x, x; T) \\ &= \frac{1}{V} \text{Tr} [\exp(T \Delta_g)] \\ &= \left( \frac{1}{4\pi T} \right)^{d/2} \sum_{n=0}^{\infty} A_n T^n \end{aligned}$$

- asymptotic expansion: space-time dimension seen by diffusion process

$$d = -2 \left. \frac{d \ln P_g(T)}{d \ln T} \right|_{T=0}$$

# Spectral dimension for classical manifolds

- Heat-equation: diffusion of scalar test particle on manifold with metric  $g$

$$\partial_T K_g(x, x'; T) = \Delta_g K_g(x, x'; T)$$

- define average return probability

$$\begin{aligned} P_g(T) &\equiv \frac{1}{V} \int d^d x \sqrt{g(x)} K_g(x, x; T) \\ &= \frac{1}{V} \text{Tr} [\exp(T \Delta_g)] \\ &= \left( \frac{1}{4\pi T} \right)^{d/2} \sum_{n=0}^{\infty} A_n T^n \end{aligned}$$

- asymptotic expansion: space-time dimension seen by diffusion process

$$d = -2 \left. \frac{d \ln P_g(T)}{d \ln T} \right|_{T=0}$$

**$P(T)$ : accessible in CDT and FRG!**

# Spectral dimension of QEG space-times

- in QEG: metric of manifold is  $k$ -dependent
  - $\implies$  diffusion process “with momentum  $k$ ” sees metric  $\langle g_{\mu\nu} \rangle_k$
  - $\implies$  diffusion equation and return probability will become  $k$ -dependent
- Computation of the spectral dimension:
  1. determine  $k$ -dependence of  $\Delta(k)$
  2. solve the  $k$ -dependent heat equation
  3. evaluate “quantum return probability”  $P(T)$
  4. obtain spectral dimension

$$\mathcal{D}_s = -2 \left. \frac{d \ln P(T)}{d \ln T} \right|_{T=0}$$

# Spectral dimension $\mathcal{D}_s$ of QEG space-times

- spectral dimension

$$\mathcal{D}_s = -2 \left. \frac{d \ln P(T)}{d \ln T} \right|_{T=0}$$

- Quantum return probability:

$$P(T) = \int \frac{d^d p}{(2\pi)^d} \exp[-p^2 F(p^2) T], \quad F(p^2) = \Lambda(p)/\Lambda(k_0)$$

- classical regime: no running,  $F(p^2) = 1$ :

$$P(T)|_{T=0} \propto T^{-d/2} \implies \mathcal{D}_s = d$$

- fixed point regime:  $\Lambda(p) \propto p^2 \rightarrow F(p^2) \propto p^2$ :

$$P(T)|_{T=0} \propto T^{-d/4} \implies \mathcal{D}_s = d/2$$

# Spectral dimension $\mathcal{D}_s$ of QEG space-times

- spectral dimension

$$\mathcal{D}_s = -2 \left. \frac{d \ln P(T)}{d \ln T} \right|_{T=0}$$

- Quantum return probability:

$$P(T) = \int \frac{d^d p}{(2\pi)^d} \exp[-p^2 F(p^2) T], \quad F(p^2) = \Lambda(p)/\Lambda(k_0)$$

- classical regime: no running,  $F(p^2) = 1$ :

$$P(T)|_{T=0} \propto T^{-d/2} \implies \mathcal{D}_s = d$$

- fixed point regime:  $\Lambda(p) \propto p^2 \rightarrow F(p^2) \propto p^2$ :

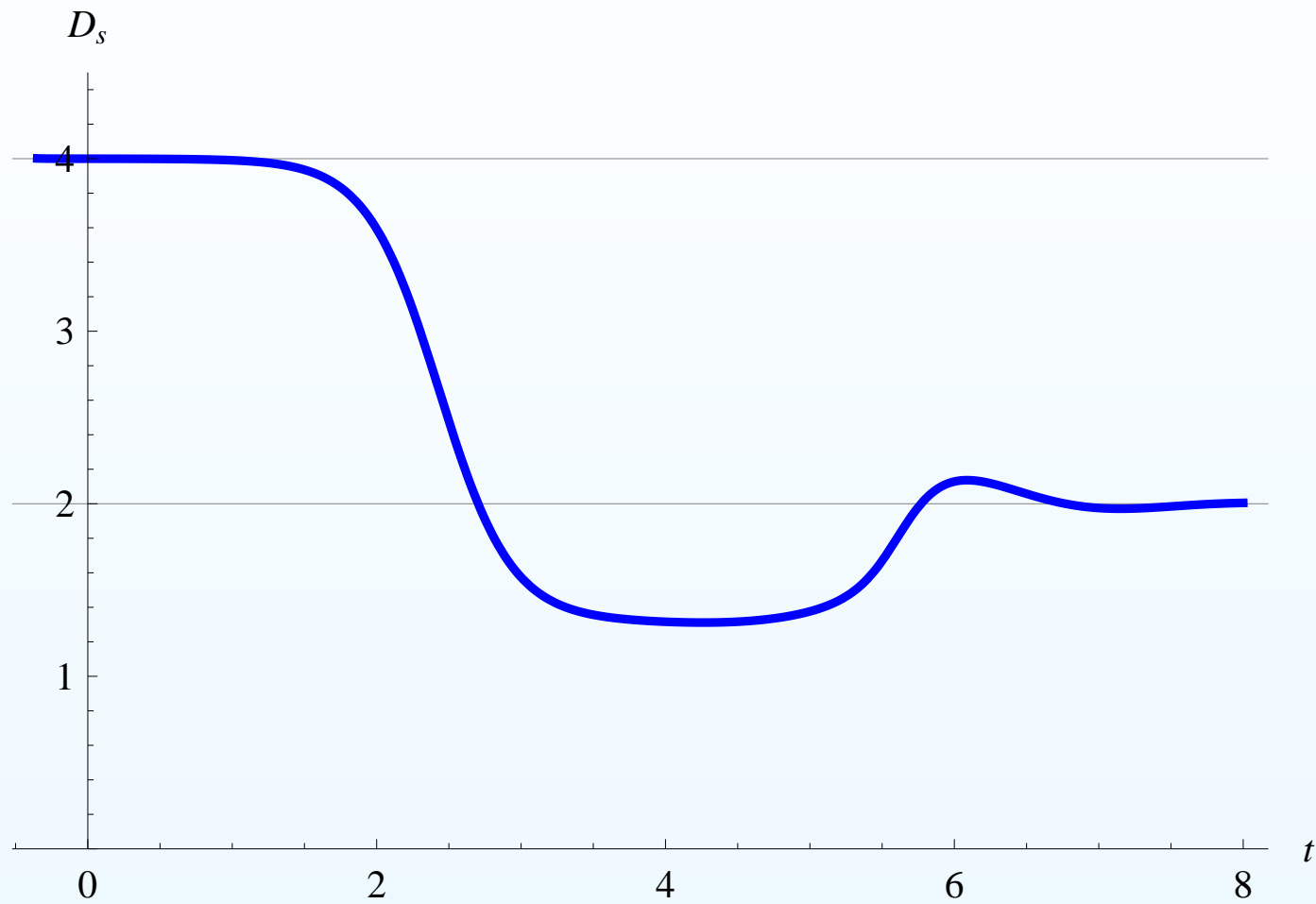
$$P(T)|_{T=0} \propto T^{-d/4} \implies \mathcal{D}_s = d/2$$

$d = 4$ : QEG predicts continuous change of fractal dimension

$$\mathcal{D}_s = 4 \text{ macroscopically} \implies \mathcal{D}_s = 2 \text{ microscopically}$$

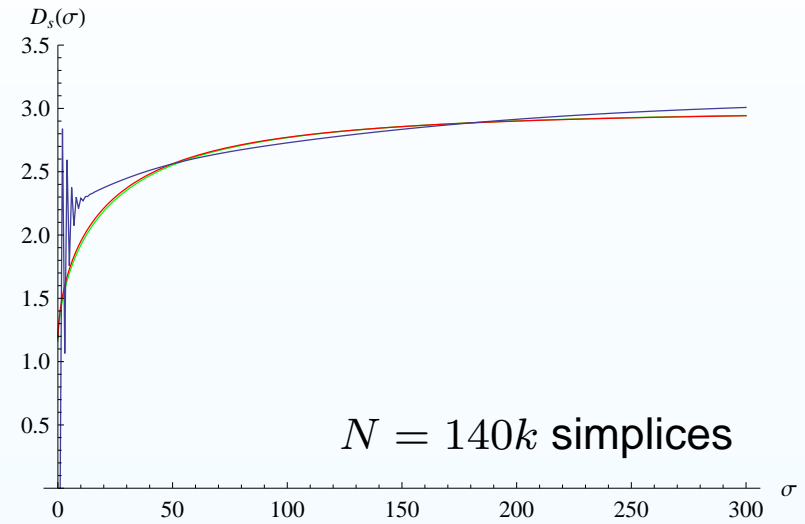
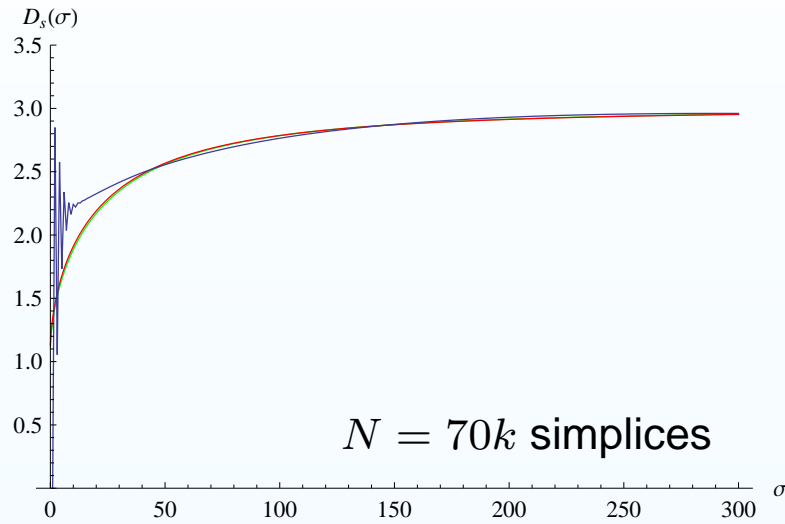
# Spectral dimension $\mathcal{D}_s$ of QEG space-times

Flow of spectral dimension along a typical RG-trajectory



# Matching FRG and 3d-lattice-simulations

[3d-CDT-Data: D. Benedetti and J. Henson, Phys. Rev. D80 (2009) 124036]



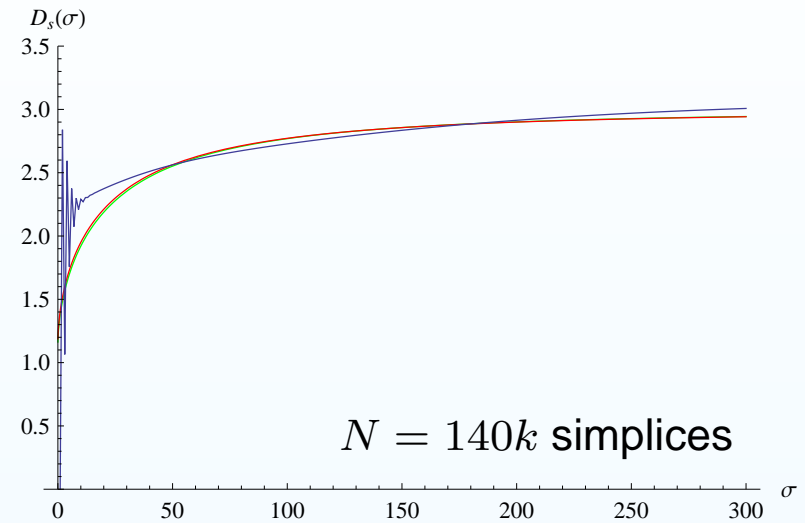
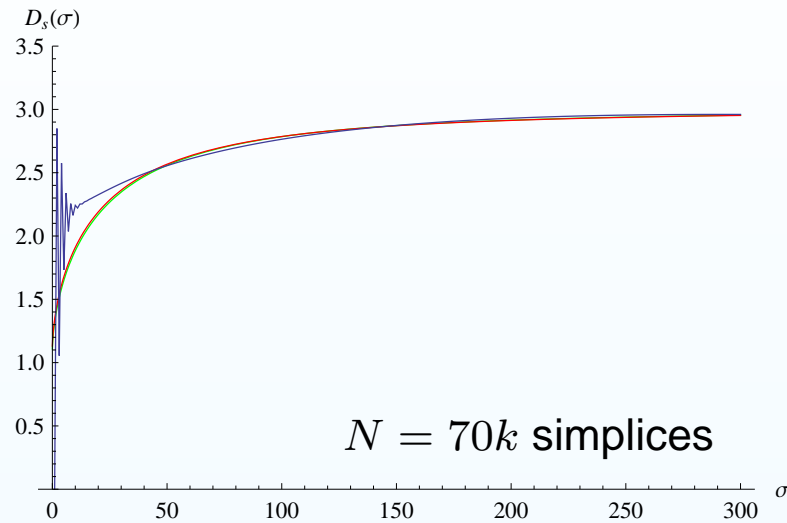
Semiclassical regime: parameterize leading quantum corrections:

$$\Lambda(p) = \Lambda_0 (1 + a p^\delta)$$

FRG-prediction:  $\delta = 3$   $\approx$  Lattice fit:  $\delta \approx 3.3$

# Matching FRG and 3d-lattice-simulations

[3d-CDT-Data: D. Benedetti and J. Henson, Phys. Rev. D80 (2009) 124036]



Semiclassical regime: parameterize leading quantum corrections:

$$\Lambda(p) = \Lambda_0 (1 + a p^\delta)$$

FRG-prediction:  $\delta = 3$   $\approx$  Lattice fit:  $\delta \approx 3.3$

can compare and fit CDT Data and RG-results



# Summary

gravitational asymptotic safety program

UV completion of gravity provided by non-trivial RG fixed point

Functional Renormalization Group equations:

- all computations support the existence of this fixed point
- UV-critical surface has finite dimension  $\iff$  predictivity

Causal Dynamical Triangulations

- spectral dimension allows to compare lattice and continuum results!

# Summary

gravitational asymptotic safety program

UV completion of gravity provided by non-trivial RG fixed point

Functional Renormalization Group equations:

- all computations support the existence of this fixed point
- UV-critical surface has finite dimension  $\iff$  predictivity

Causal Dynamical Triangulations

- spectral dimension allows to compare lattice and continuum results!

Acknowledgments:

- D. Benedetti, J. Henson providing CDT-data
- K. Groh, S. Rechenberger and O. Zanusso for collaboration

# Summary

gravitational asymptotic safety program

UV completion of gravity provided by non-trivial RG fixed point

Functional Renormalization Group equations:

- all computations support the existence of this fixed point
- UV-critical surface has finite dimension  $\iff$  predictivity

Causal Dynamical Triangulations

- spectral dimension allows to compare lattice and continuum results!

Acknowledgments:

- D. Benedetti, J. Henson providing CDT-data
- K. Groh, S. Rechenberger and O. Zanusso for collaboration

Thank you!