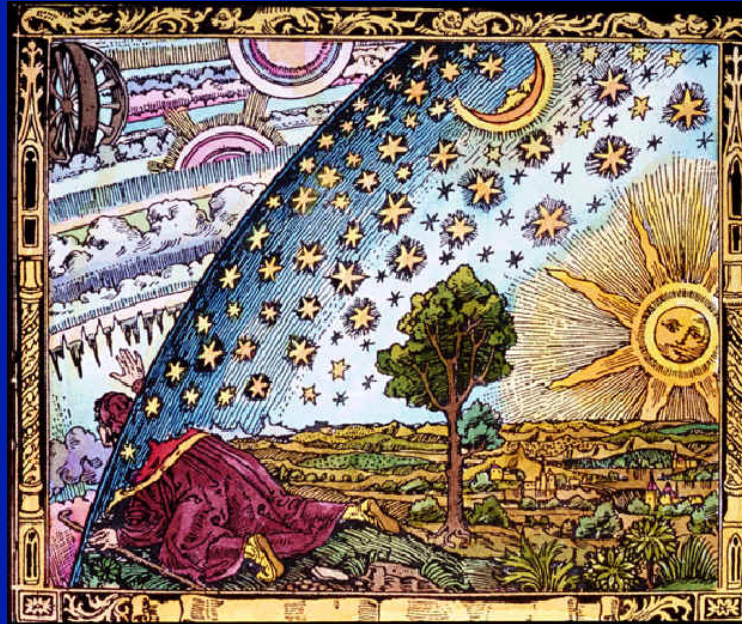


# Some phenomenological aspects of Loop Quantum Cosmology



**Aurélien Barrau**

Laboratoire de Physique Subatomique et de Cosmology  
CNRS/IN2P3 – University Joseph Fourier – Grenoble, France

# Why going beyond GR ?

## Dark energy (and matter) / quantum gravity

- **Observations : the acceleration of the Universe**
- **Theory : singularity theorems**

**Successful techniques of QED do not apply to gravity. Something new has to be invented.**

**Which gedankenexperiment ? (as is QM, SR and GR) Which paradoxes ?**

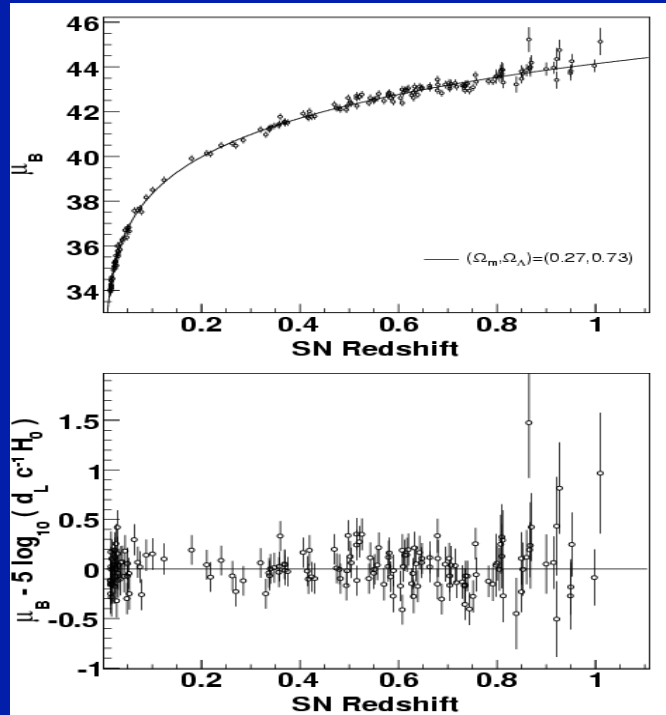
**Quantum black holes and the early universe are privileged places to investigate such effects !**

- \* **Entropy of black holes**
- \* **End of the evaporation process, IR/UV connection**
- \* **the Big-Bang**

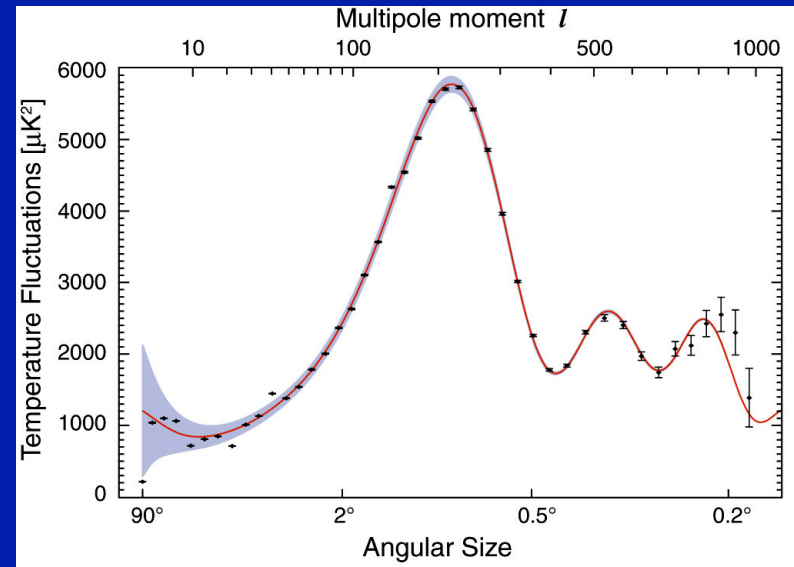
→ **Many possible approaches : strings, covariant approaches (effective theories, the renormalization group, path integrals), canonical approaches (quantum geometrodynamics, loop quantum gravity), etc. See reviews par C. Kiefer**

→ **I will focus on LQG**

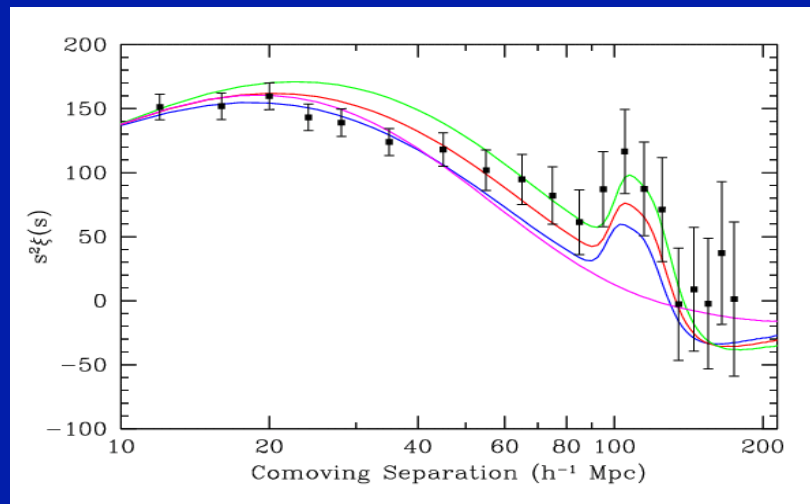
# The observed acceleration



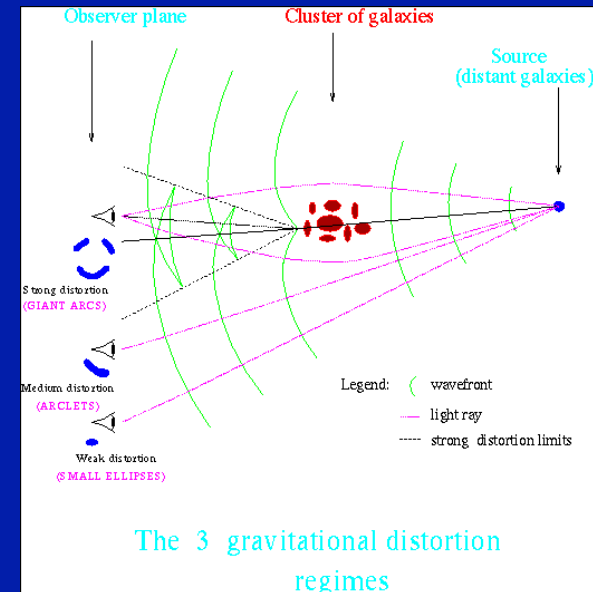
SNLS, Astier et al.



WMAP, 5 ans



SDSS, Eisenstein et al. 2005



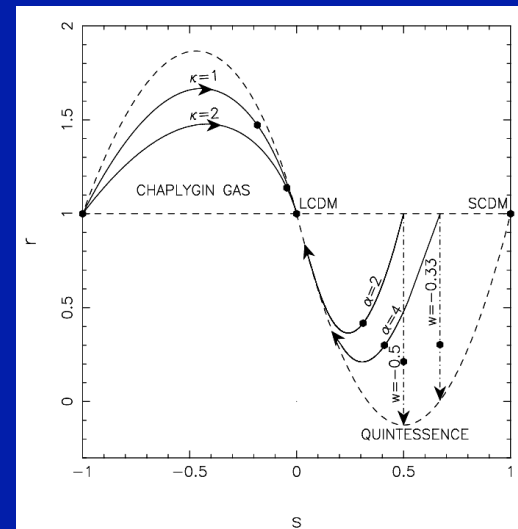
$$\Lambda / 8\pi G \sim 10^{-47} \text{ GeV}^4$$

$$H^2 = \frac{8\pi G}{3} \left( \sum_a \rho_a + \rho_{DE} \right) - \frac{k}{a^2},$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} \left( \sum_a (\rho_a + 3p_a) + \rho_{DE} + 3p_{DE} \right)$$

$$a(t) = a(t_0) + \dot{a}|_0(t-t_0) + \frac{\ddot{a}|_0}{2}(t-t_0)^2 + \frac{\dddot{a}|_0}{6}(t-t_0)^3 + \dots$$

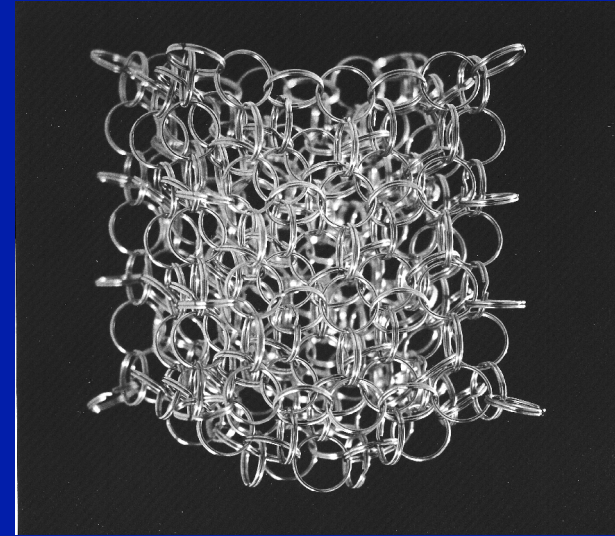
Level	Geometrical Parameter	Physical Parameter
1	$H(z) \equiv \frac{\dot{a}}{a}$	$\rho_m(z) = \rho_{0m}(1+z)^3,$ $\rho_{DE} = \frac{3H^2}{8\pi G} - \rho_m$
2	$q(z) \equiv -\frac{\ddot{a}a}{\dot{a}^2} = -1 + \frac{d \log H}{d \log(1+z)}$ $q(z) _{\Lambda\text{CDM}} = -1 + \frac{3}{2}\Omega_m(z)$	$V(z), T(z) \equiv \frac{\dot{\phi}^2}{2}, w(z) = \frac{T-V}{T+V},$ $\Omega_V = \frac{8\pi G V}{3H^2}, \Omega_T = \frac{8\pi G T}{3H^2}$
3	$r(z) \equiv \frac{\ddot{a}a^2}{\dot{a}^3}, s \equiv \frac{r-1}{3(q-1/2)}$ $\{r, s\} _{\Lambda\text{CDM}} = \{1, 0\}$	$\Pi(z) \equiv \dot{V} = \dot{\phi}V', \Omega_\Pi = \frac{8\pi G \dot{V}}{3H^3}$



Alam et al., MNRAS  
344 (2003) 1057

# Not that easy in the UV limit...

« Can we construct a quantum theory of spacetime based only on the experimentally well confirmed principles of general relativity and quantum mechanics ? » L. Smolin, hep-th/0408048



## DIFFEOMORPHISM INVARIANCE

**Loops (solutions to the WDW) = space**

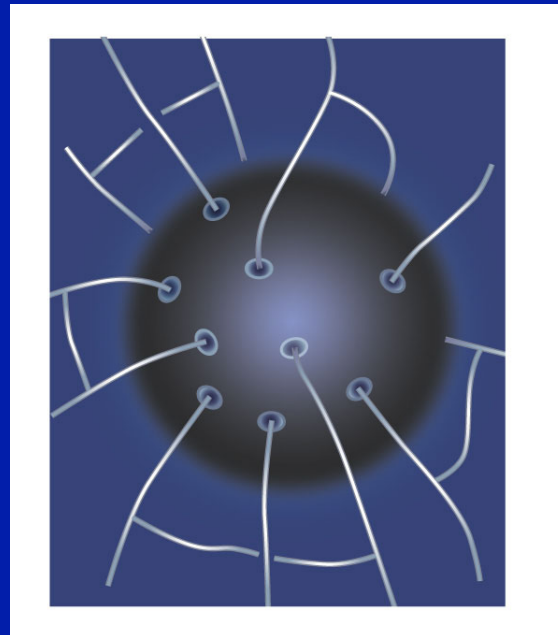
- Mathematically well defined
- Singularities
- Black holes

**In QFT, one need  $U(O)$  where  $O$  is a region of  $(M,g)$  to define (anti)commutation relations between spacelike separated regions !!!**

# How to build Loop Quantum Gravity ?

## 1) If you are a relativist...

- Foliation  $\rightarrow$  space metric and conjugate momentum
- Constraints (difféomorphism, hamiltonian +  $SO(3)$ )
- Quantization « à la Dirac »  $\rightarrow$  WDW  $\rightarrow$  Ashtekar variables
- « smearing »  $\rightarrow$  holonomies and fluxes



See e.g. the book « Quantum Gravity » by C. Rovelli

# How to build Loop Quantum Gravity ?

## 2) If you are a particle physicist...

- **Think of lattice QCD**
- **Define a graph and the Hilbert space :  $L^2(G^L/G^N)$ . The Fock space is obtained by taking the appropriate limit.**
- **In gravity you do the same :  $H\Gamma = L^2[SU(2)^L/SU(2)^N]$ . Then  $\tilde{H}\Gamma = H\Gamma / \sim$  (automorphism group)**
- **Define « natural » operators on  $L^2[SU(2)]$**
- **Gauge invariance + Penrose theorem lead to a simple geometrical interpretation in the classical limit.**
- **Define the spin-network basis (diagonalizes the area and volume operators)**

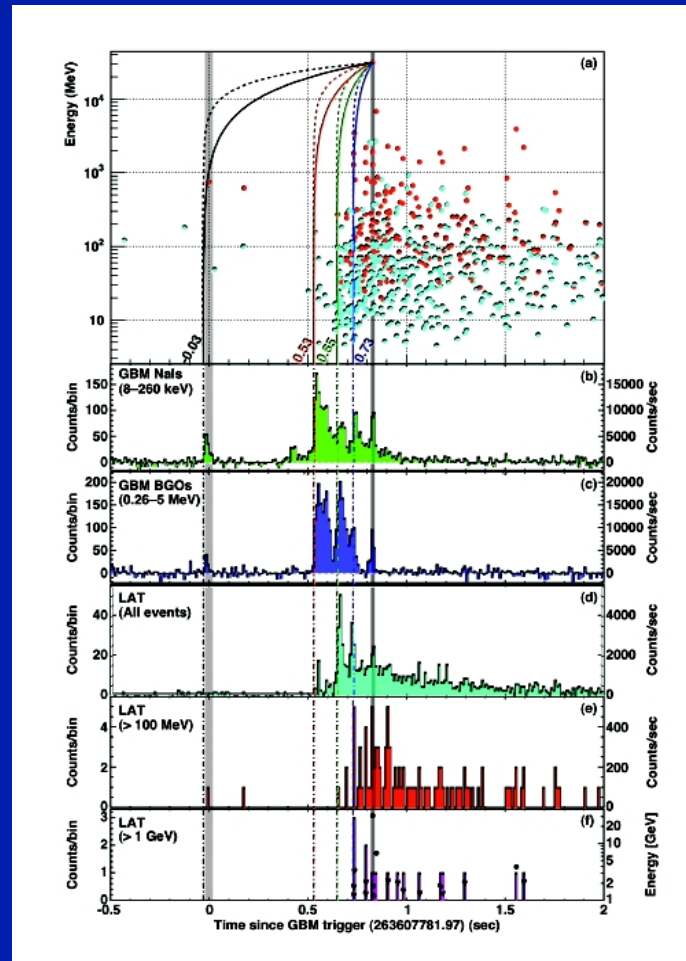
See e.g. the book « Quantum Gravity » by C. Rovelli

- Mathematically well defined**
- Singularities**
- Black holes**



# Experimental tests

- **High energy gamma-ray** (Amélineo-Camelia et al.)



Not very conclusive however



# Experimental tests

- **Discrete values for areas and volumes** (Rovelli et al.)
- **Observationnal cosmology** (... , et al.)

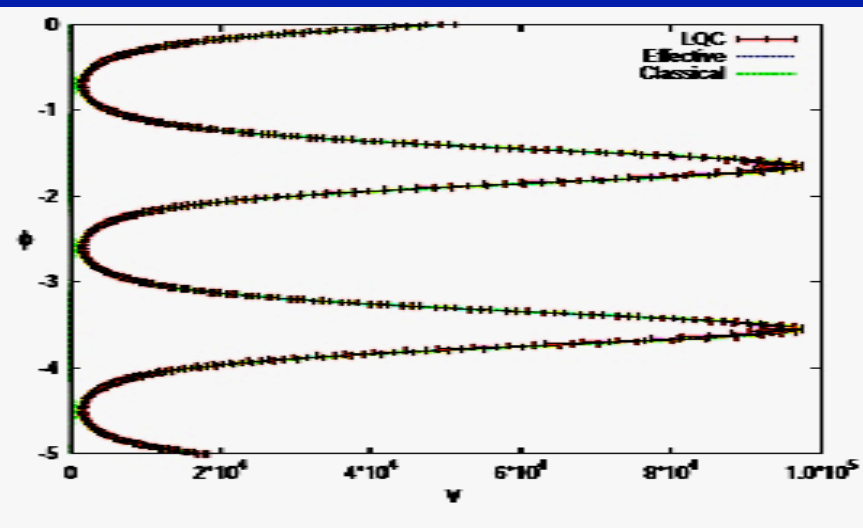
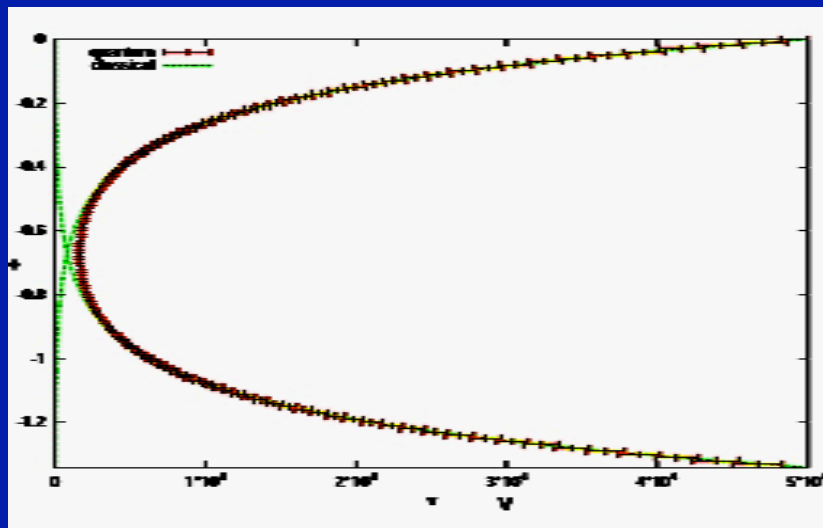
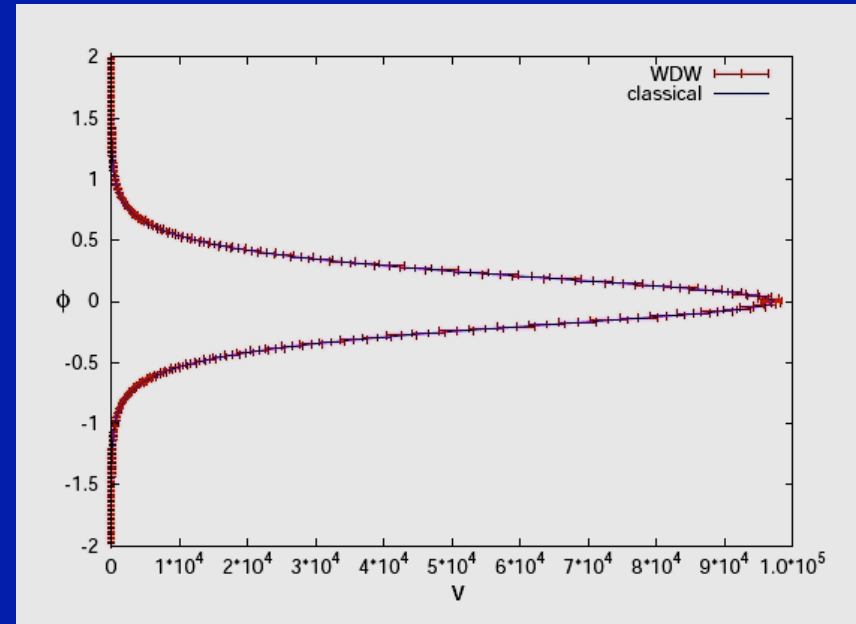
**LQC :**

- **IR limit**
- **UV limit (bounce)**
- **inflation**

# WDW vs LQC

**WDW : The IR test is passed with flying colors.  
But the singularity is not resolved. →**

**In LQC, no Big Bang, no new principle, no  
new principle required, « other side »  
opened, huge « quantum geometrical »  
effects @  $10^{94}$  g/cm<sup>3</sup>**



Plots from Ashtekar

## LQC: a few results

- The volume of the Universe takes its minimum value at the bounce and scales as  $p(\Phi)$
- The recollapse happens at  $V_{\max}$  which scales as  $p(\Phi)^{3/2}$ . GR is OK.
- The states remain sharply peaked for a very large number of cycles. Determinism is kept even for an infinite number of cycles.
- The dynamics can be derived from effective Friedmann equations

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- The LQC correction naturally comes with the correct sign. This is non-trivial.
- Furthermore, one can show that the upper bound of the spectrum of the density operator coincides with  $\rho_{\text{crit}}$

→ Role of the high symmetry assumed ? (string entropy ?)

# LQC & inflation

## -Inflation

- success (paradoxes solved, perturbations, etc.)
- difficulties (no fundamental theory, initial conditions, etc.)

## -LQC

- success (background-independant quantization of GR, BB Singularity resolution, good IR limit)
- difficulties (very hard to test !)

*Could it be that considering both LQC and inflation within the same framework allows to cure simultaneously all the problems ?*

*Bojowald, Hossain, Copeland, Mulryne, Nunes, Shaeri, Tsujikawa, Singh, Maartens, Vandersloot, Lidsey, Tavakol, Mielczarek .....*

# First approach: LQC corrections to the modes in a classical background

« standard » inflation

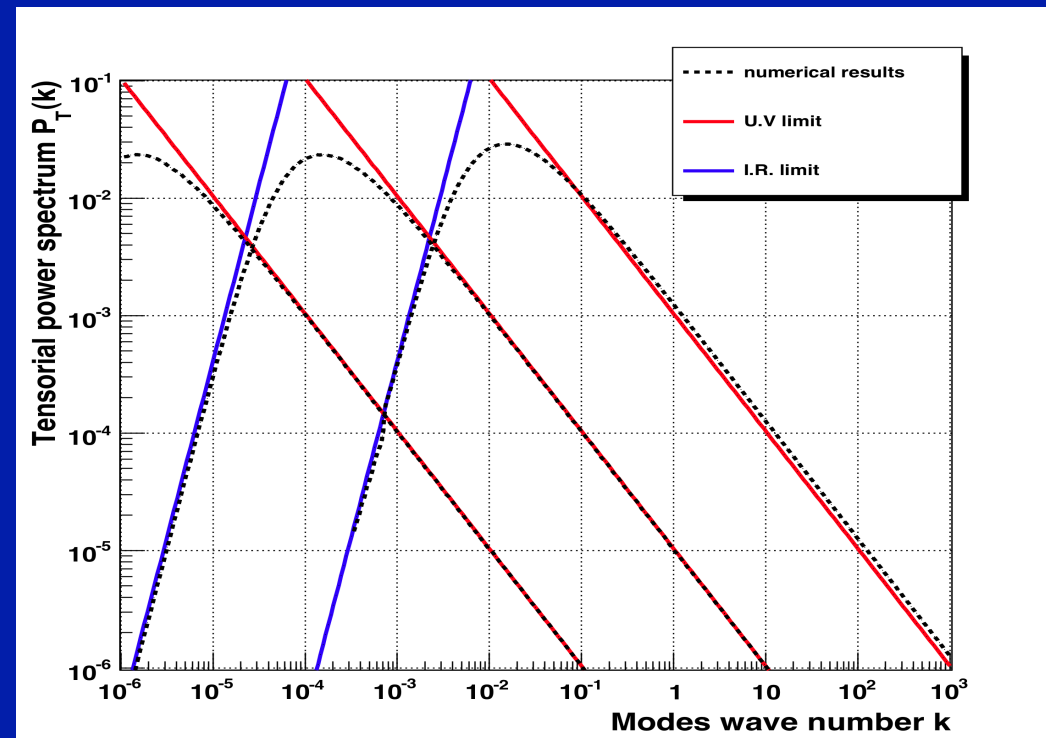
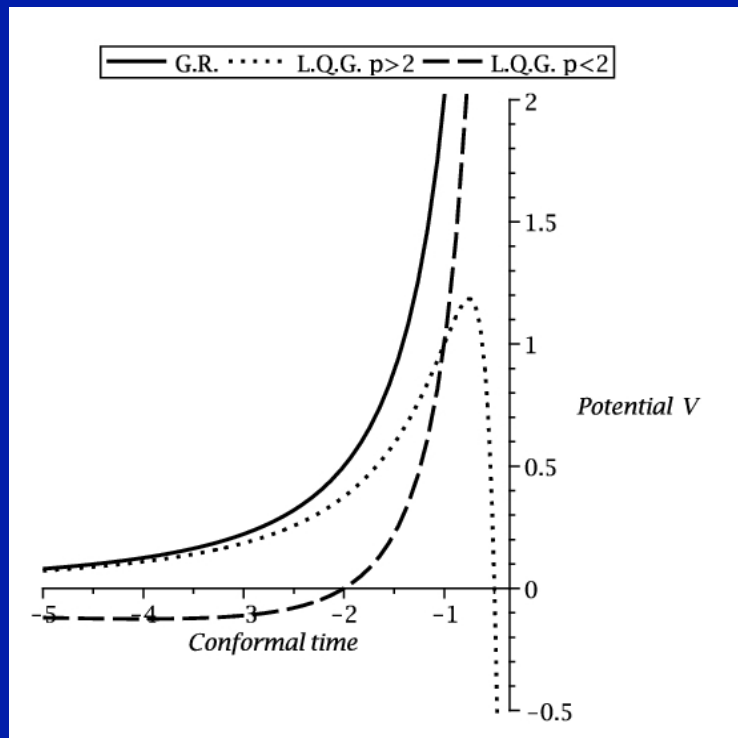
-decouples the effects  
-happens after superinflation

Bojowald & Hossain, Phys. Rev. D 77, 023508 (2008)

A.B. & Grain, Phys. Rev. Lett. , 102, 081321 (2009)

J. Grain, A.B., A. Gorecki, Phys. Rev. D , 79, 084015 (2009)

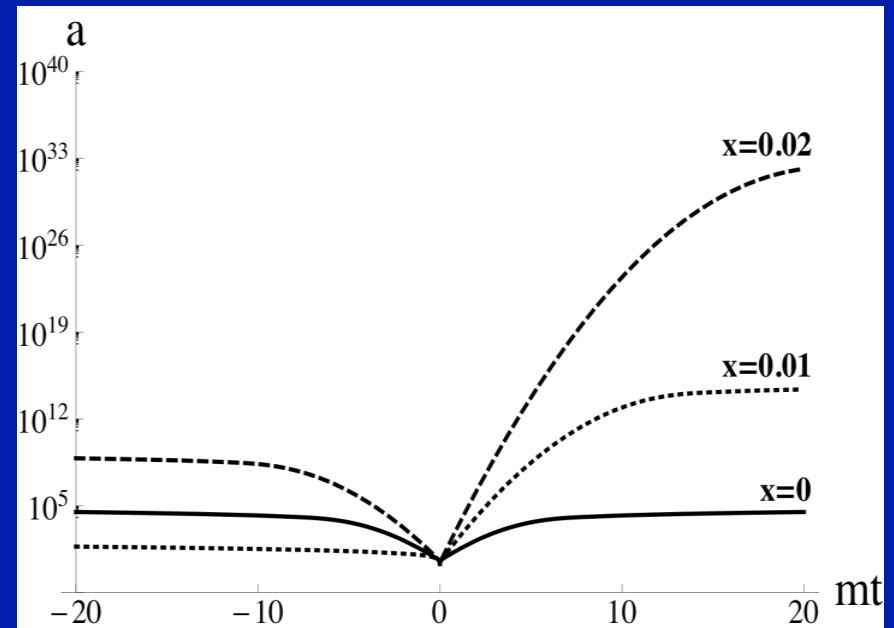
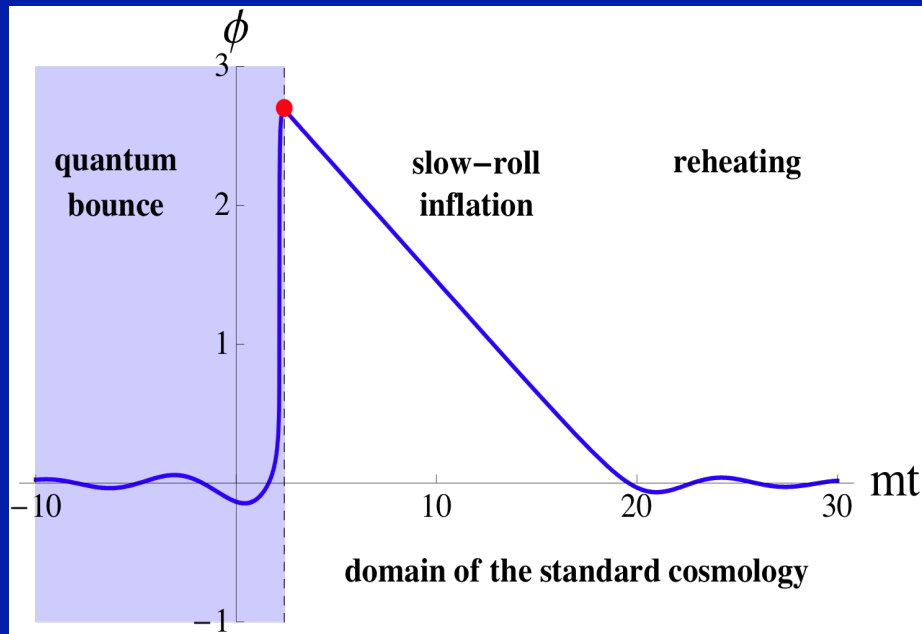
J. Grain, T. Cailleteau, A.B., Phys. Rev. D, 81, 024040 (2010)



# Second approach : Taking into account the background modifications

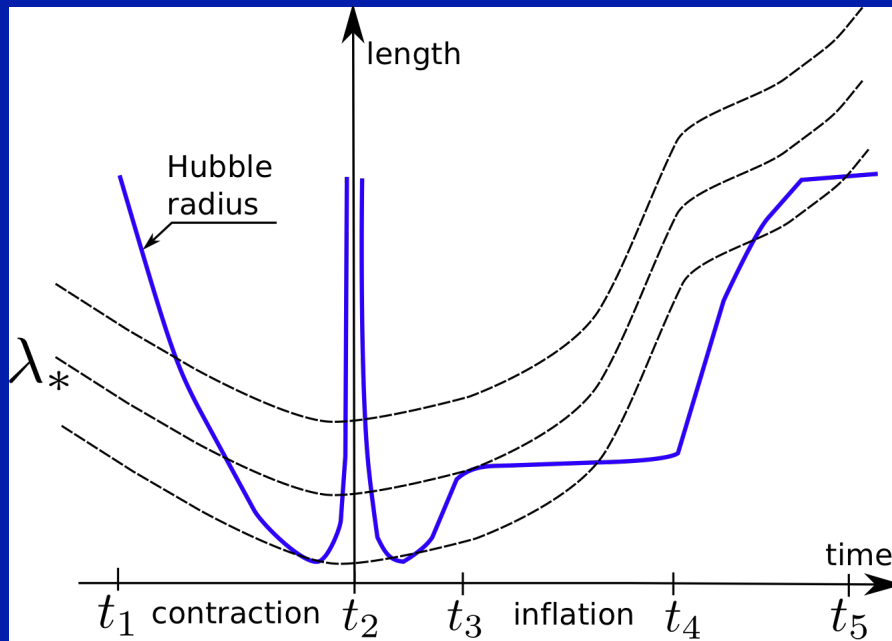
H changes sign in the KG equation  $\phi'' + 3H\phi' + m^2\phi = 0$

→ Inflation inevitably occurs !

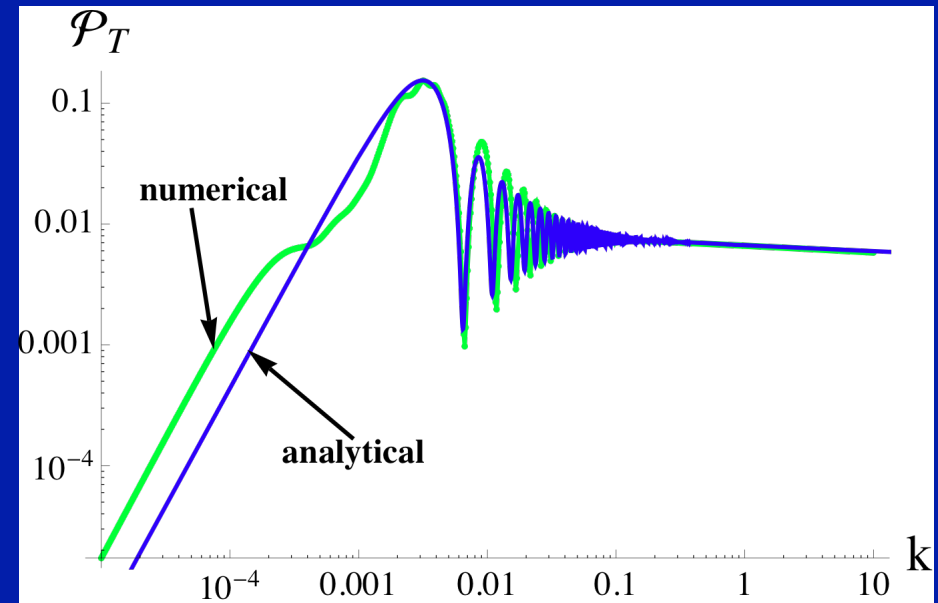


Mielczarek, Cailleteau, Grain, A.B., Phys. Rev. D, 81, 104049, 2010

# A tricky horizon history...



**Physical modes may cross the horizon several times...**

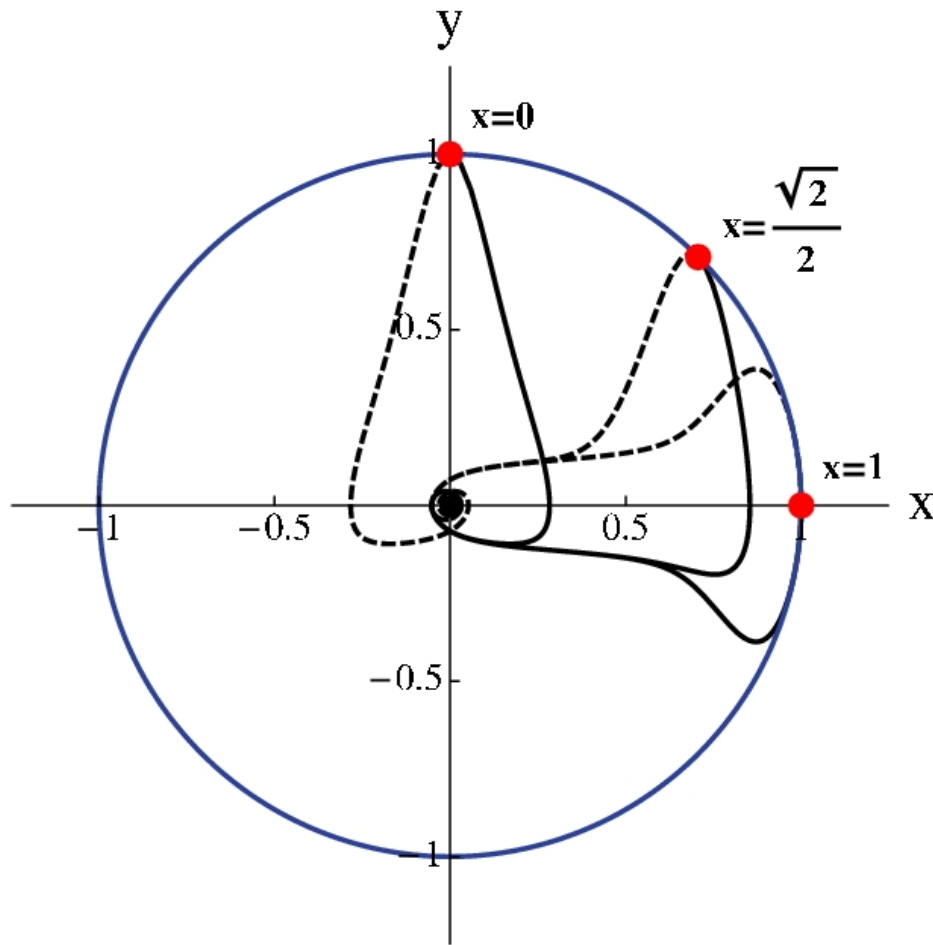


**Computation of the primordial power spectrum:**

- Bogolubov transformations**
- Full numerical resolution**

- The power is suppressed in the infra-red (IR) regime. This is a characteristic feature associated with the bounce
- The UV behavior agrees with the standard general relativistic picture.
- Damped oscillations are superimposed with the spectrum around the "transition" momentum  $k^*$  between the suppressed regime and the standard regime.
- The first oscillation behaves like a "bump" that can substantially exceed the UV asymptotic value.





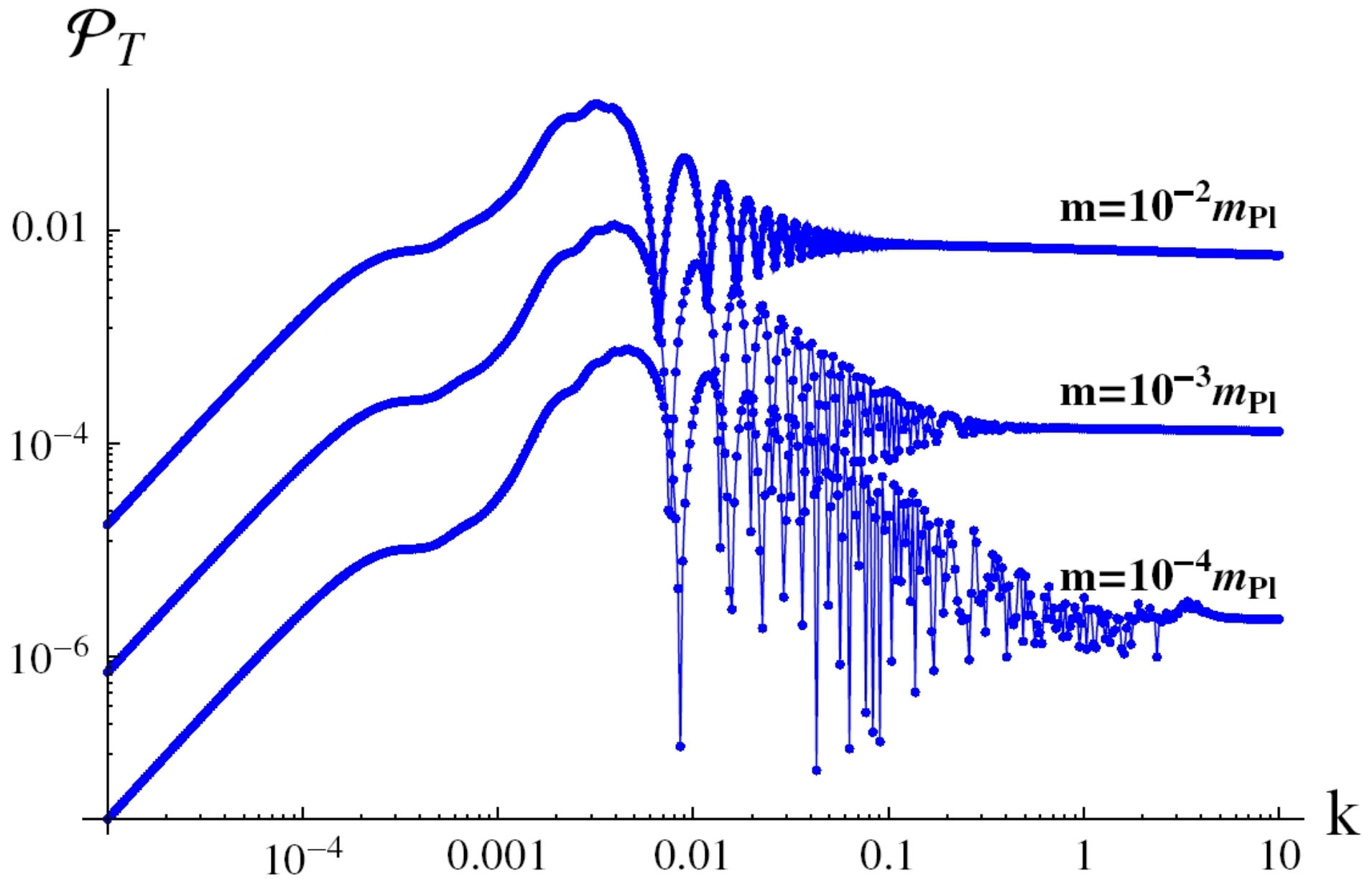
Effective description with a Bogoliubov transformation :

- Frequency of the oscillations controlled by  $\Delta(\eta)$ , the width of the bounce
- Amplitude of the oscillations controlled by  $k_0$ , the effective mass at the bounce

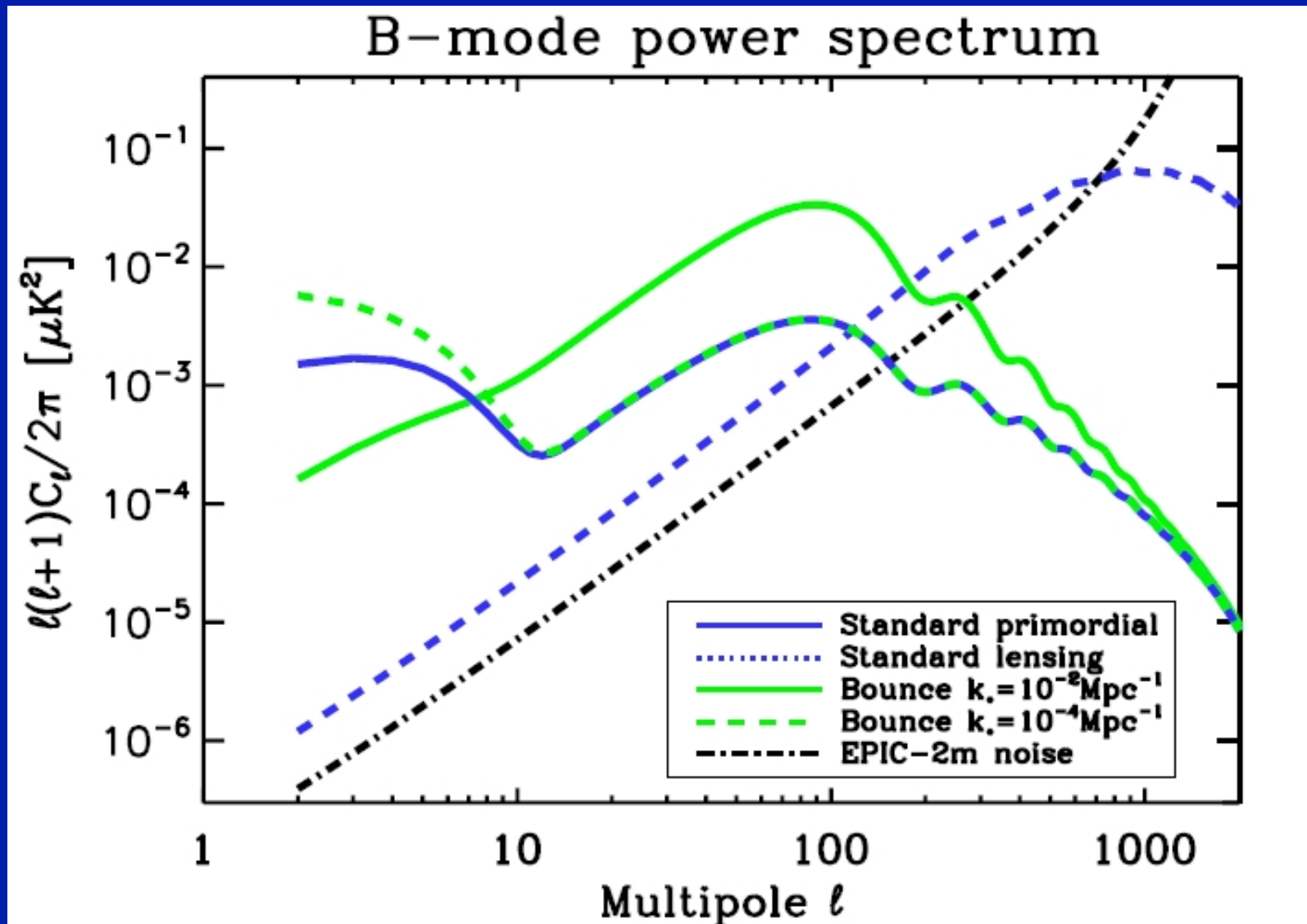
Fundamental description :

- $R$  driven by the field mass
- $k^*$  driven by initial conditions

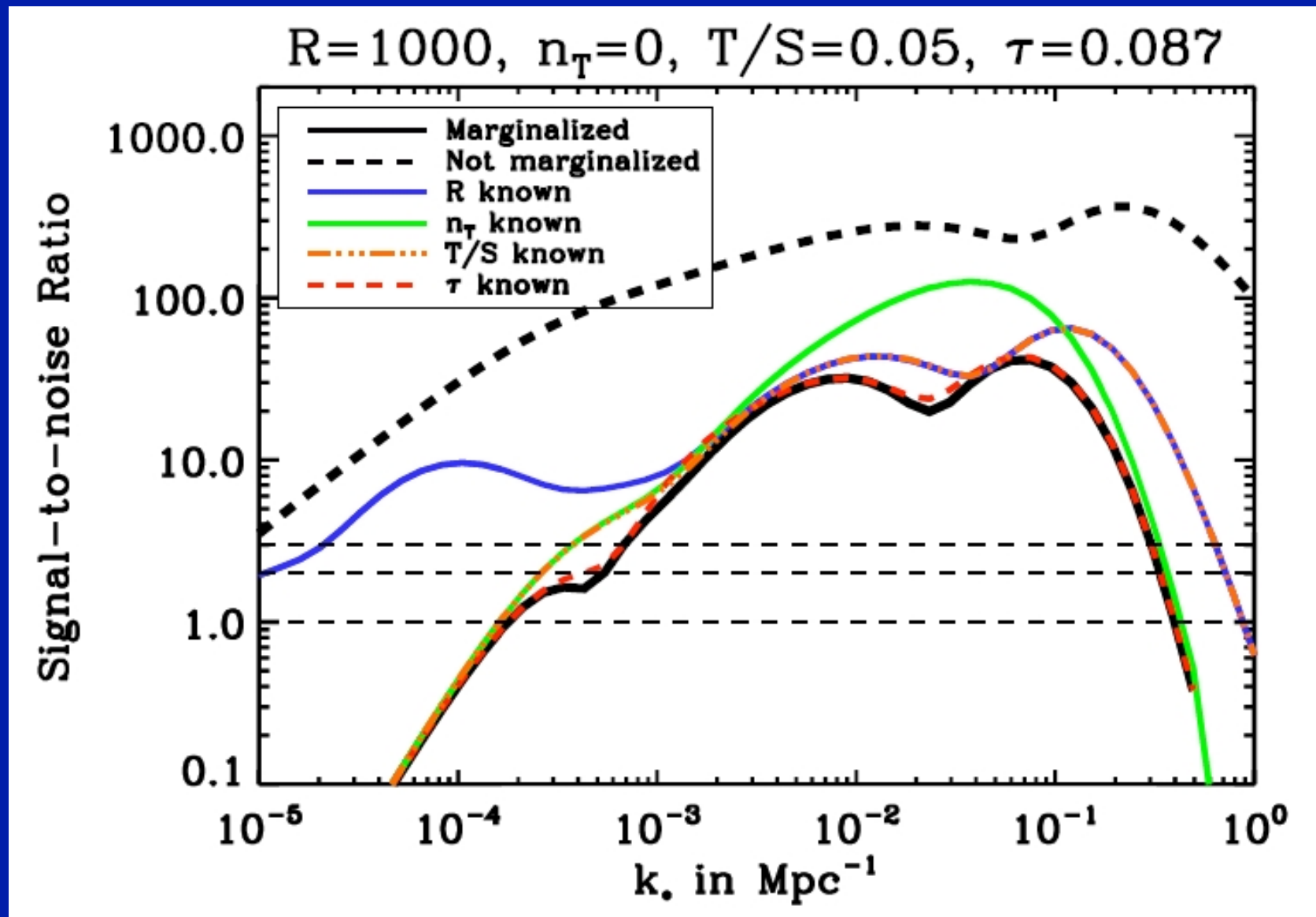
**Initial conditions are critical**



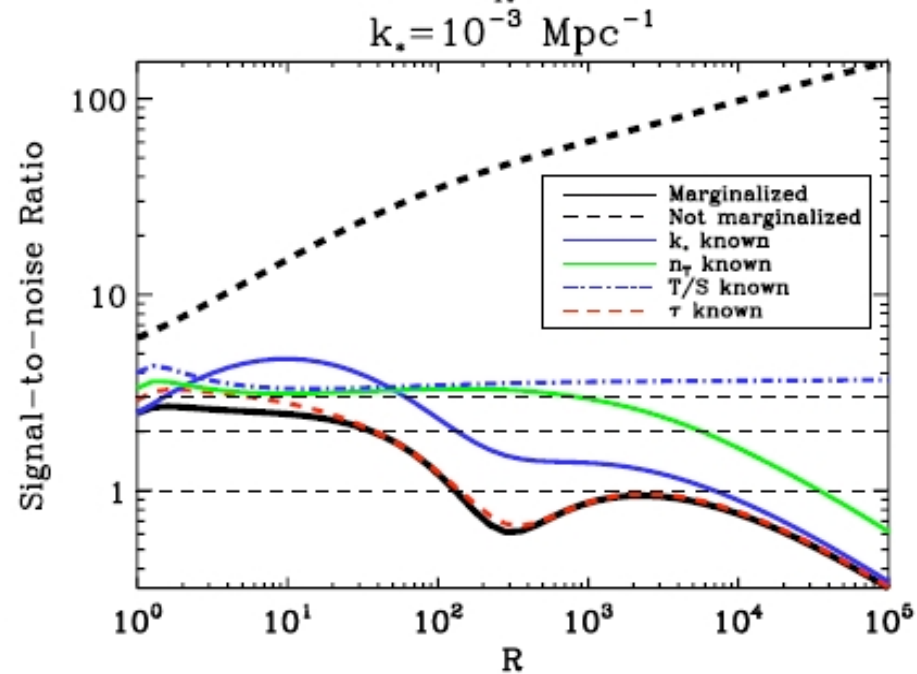
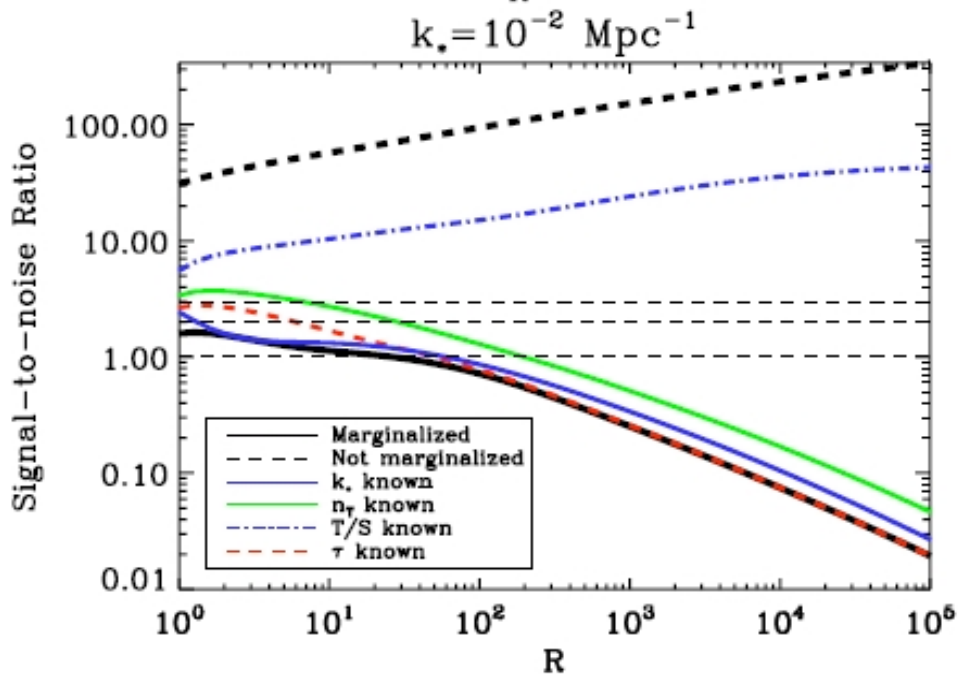
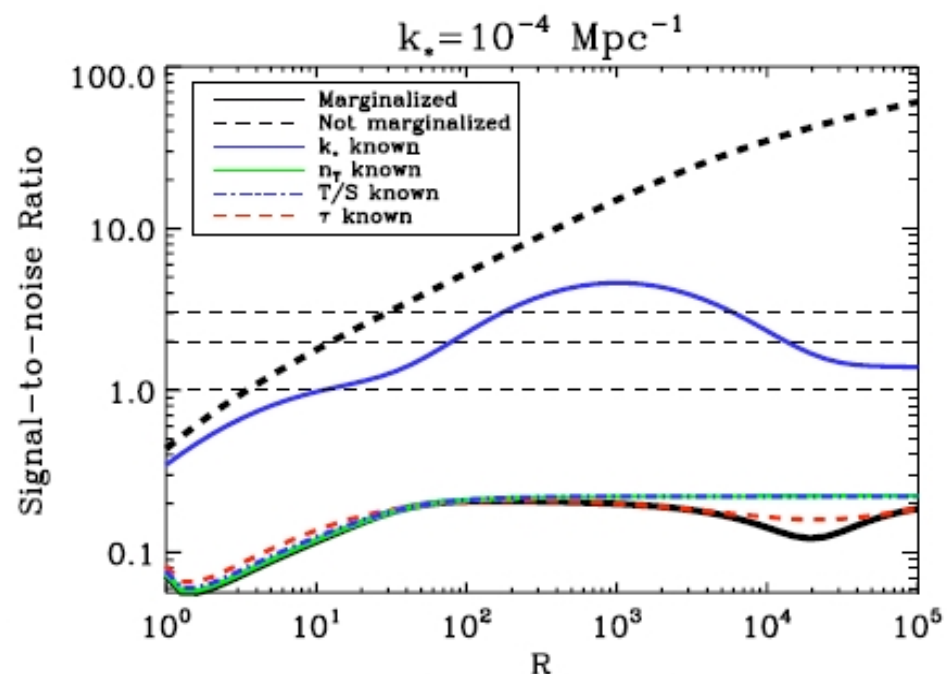
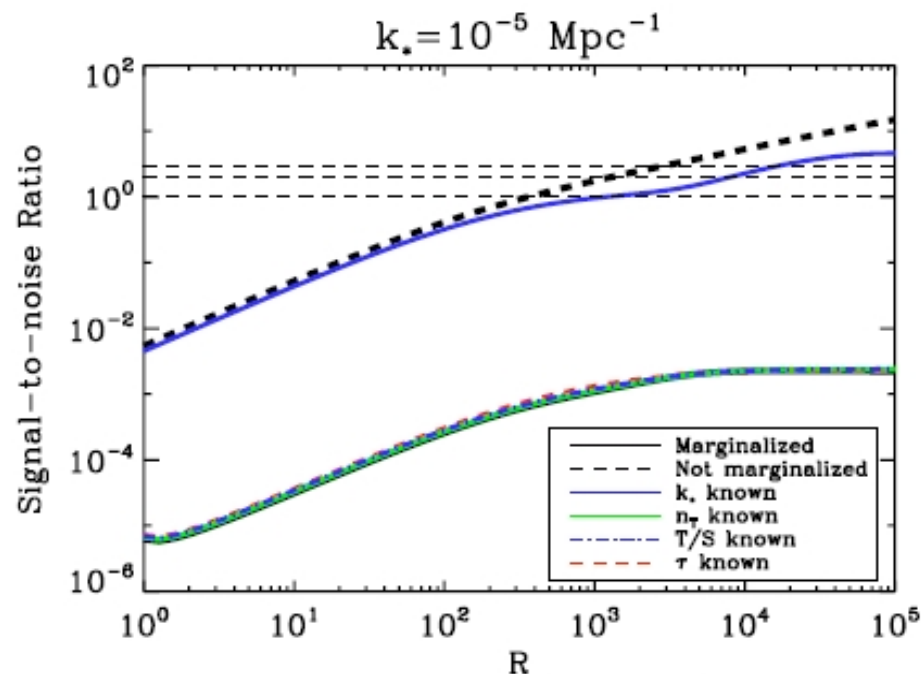
# CMB consequences...



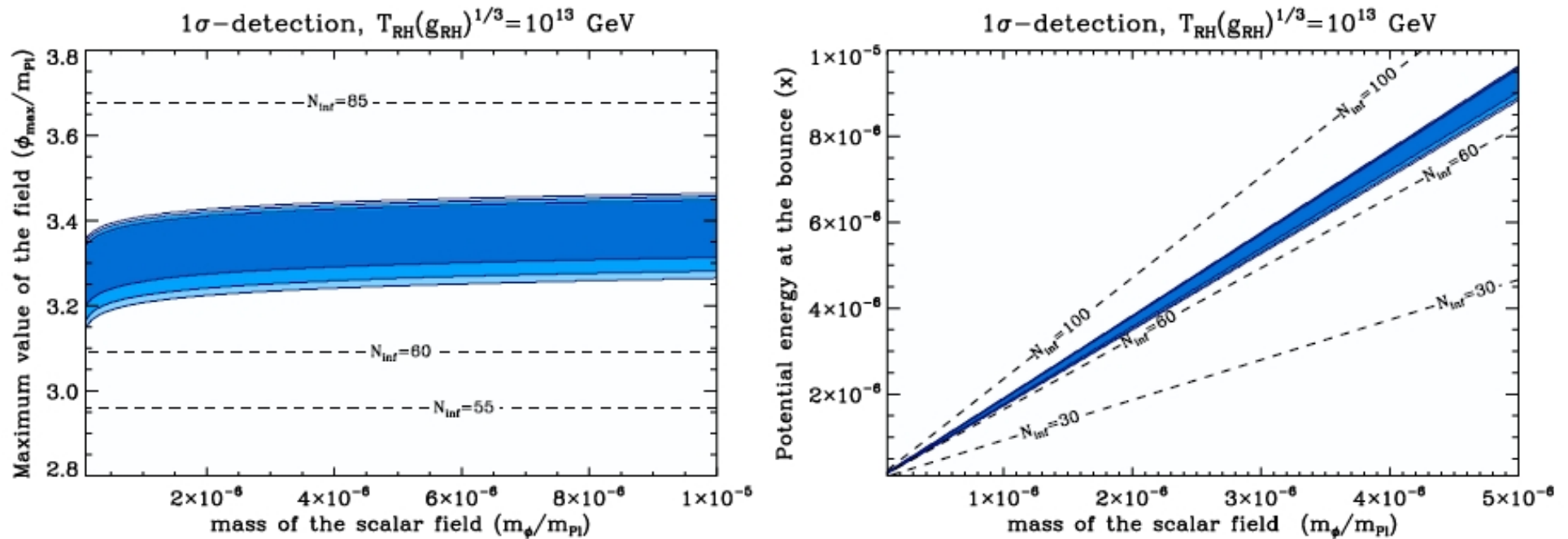
Grain, A.B., Cailleteau, Mielczarek, Phys. Rev. D, 82, 123520 (2010)



A.B., Grain et al.



# CMB consequences



**If the scalar spectrum is assumed not to be affected :  
one needs  $x < 2E-6$  to probe the model**

Grain, A.B., Cailleteau, Mielczarek, Phys. Rev. D, 82, 123520 (2010)

Is a  $N > 78$  inflation probable ?  
What is the probability to be  
compatible with WMAP data ?

**YES. Good news from Ashtekar and Sloan !**

Can B-mode be used to distinguished  
with string inflation ?

**I think yes.**



# Anomaly-free vector algebra for holonomy corrections

$$\begin{aligned}
 \{S_{\text{tot}}[N_1], S_{\text{tot}}[N_1]\} &= 0, \\
 \{D_{\text{tot}}[N_1^a], D_{\text{tot}}[N_2^a]\} &= 0, \\
 \{S_{\text{tot}}[N], D_{\text{tot}}[N^a]\} &= \frac{\bar{N}}{\sqrt{\bar{p}}} \mathcal{B} D^Q[N^a] \\
 &+ \frac{\bar{N}}{\kappa \sqrt{\bar{p}}} \int_{\Sigma} d^3x \delta N^c \delta_c^k (\partial_d \delta E_k^d) \delta E_k^d \mathcal{A} \\
 + [\cos(v_2 \bar{\mu} \gamma \bar{k}) - 1] &\frac{\sqrt{\bar{p}}}{2} \left( \frac{\bar{\pi}^2}{2\bar{p}^3} - V(\bar{\varphi}) \right) \times \\
 &\times \int_{\Sigma} d^3x \bar{N} \partial_c (\delta N^a) \delta_a^j \delta E_j^c \\
 &+ \frac{\bar{\pi}}{\bar{p}^{3/2}} \int_{\Sigma} d^3x \bar{N} (\partial_a \delta N^a) \delta \pi \\
 &- \bar{p}^{3/2} V_{\varphi}(\bar{\varphi}) \int_{\Sigma} d^3x \bar{N} (\partial_a \delta N^a) \delta \varphi
 \end{aligned}$$

## Including matter

The counterterms can be computed together with the integers. As expected  $v_2=0$  (no diffo correction).

→ The algebra is determined (on need  $\mathbf{B}=0$  – and of course  $\mathbf{A}=0$ )

# Perspective I : closing the algebra for scalar modes (with holonomy corrections)

$$\begin{aligned} \{H_G^Q[N], D_G^Q[N^a]\} &= \frac{\bar{N}}{\sqrt{\bar{p}}} \left[ \frac{\sin v_2 \bar{\mu} \gamma \bar{k}}{v_2 \bar{\mu} \gamma} \right. \\ &+ \left. \frac{\sin v_1 \bar{\mu} \gamma \bar{k}}{v_1 \bar{\mu} \gamma} - 2 \frac{\sin 2\bar{\mu} \gamma \bar{k}}{2\bar{\mu} \gamma} \right] D_G^Q[N^a] \\ &+ \frac{1}{\kappa} \int_{\Sigma} d^3x \frac{\bar{N}}{\sqrt{\bar{p}}} \mathcal{A}^{HD} (\partial_d \delta N^c) \delta E_k^d \delta_c^k \end{aligned} \quad (12)$$

where

$$\begin{aligned} \mathcal{A}^{HD} &= \bar{k}^2 \left[ \frac{1}{2} \left( \frac{\sin x}{x} \right)^2 (1 + \cos v_2 x) \right. \\ &+ 2 \left( \frac{\sin v_2 x}{v_2 x} \right) \left( \frac{\sin 2x}{2x} - \frac{\sin v_1 x}{v_1 x} \right) \\ &+ 2 \left( \cos x - \frac{\sin x}{x} \right) \left( \frac{\sin x}{x} \right) \cos(v_2 x) \left( \frac{\bar{p}}{\bar{\mu}} \frac{\partial \bar{\mu}}{\partial \bar{p}} \right) \\ &- \left( \frac{\sin v_2 x}{v_2 x} \right)^2 \\ &\left. - 2 \left( \cos v_2 x - \frac{\sin v_2 x}{v_2 x} \right) \left( \frac{\sin 2x}{2x} \right) \left( \frac{\bar{p}}{\bar{\mu}} \frac{\partial \bar{\mu}}{\partial \bar{p}} \right) \right]. \end{aligned} \quad (13)$$

**The solution seems to be uniquely determined.**

Complementary to Bojowald, Hossain, Kagan and Shankaranarayanan

Cailleteau, Mielczarek, A.B., Grain, preliminary

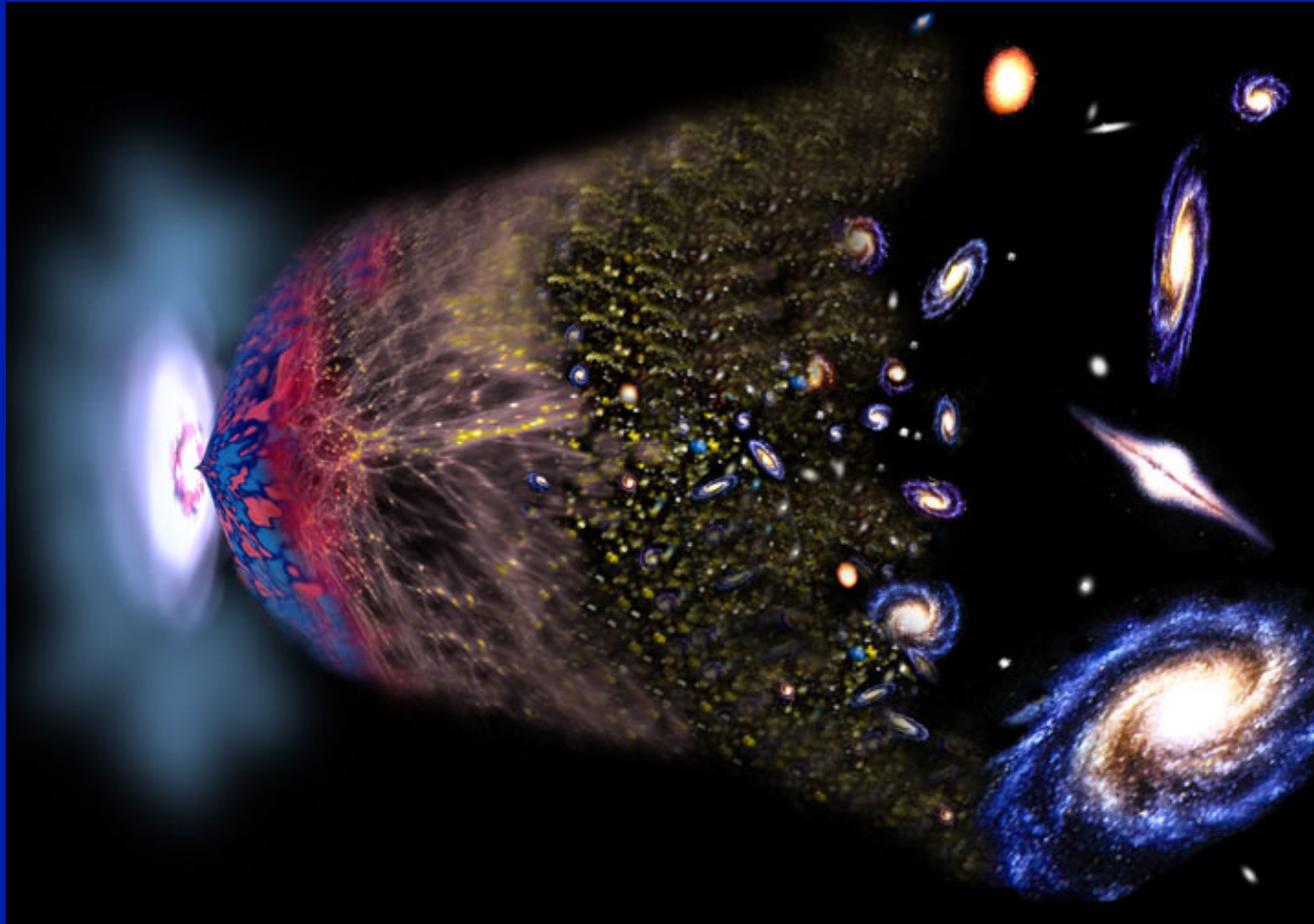
# Perpsectices II (tensor modes) : IV + holonomy for background + modes

$$\begin{aligned}\dot{\Phi} &= \frac{d\Phi}{dt} = \{\Phi, p_{\Phi}\} \frac{\partial H_m}{\partial p_{\Phi}} = D \frac{p_{\Phi}}{\bar{p}^{\frac{3}{2}}} \\ \dot{p}_{\Phi} &= \frac{dp_{\Phi}}{dt} = -\{\Phi, p_{\Phi}\} \frac{\partial H_m}{\partial \Phi} = -\bar{p}^{\frac{3}{2}} V_{,\Phi}.\end{aligned}$$

$$\begin{aligned}\ddot{\Phi} &= \dot{D} \frac{p_{\Phi}}{\bar{p}^{\frac{3}{2}}} + D \frac{\dot{p}_{\Phi}}{\bar{p}^{\frac{3}{2}}} - \frac{3}{2} \frac{\dot{\bar{p}}}{\bar{p}} D \frac{p_{\Phi}}{\bar{p}^{\frac{3}{2}}} \\ &= \frac{\dot{D}}{D} \frac{D p_{\Phi}}{\bar{p}^{\frac{3}{2}}} - D \partial_{\Phi} V(\Phi) - 3 H D \frac{p_{\Phi}}{\bar{p}^{\frac{3}{2}}} \\ &= \frac{\dot{D}}{D} \dot{\Phi} - D \partial_{\Phi} V(\Phi) - 3 H \dot{\Phi}\end{aligned}$$

**Simulation in progress**

A.B., Cailleteau, Grain, in progress



*Toward a loop – inflation paradigm ?*