

Dark energy : Myth or Reality ?

Alain Blanchard



Annecy, May 17, 2011



Standard Cosmological model:

Homogenous (on large scale) $\rightarrow$ RW metric:

$$ds^2 = -c^2 dt^2 + R(t)^2 \left[r^2 (d\theta^2 + \sin^2 \theta d\phi^2) + \frac{dr^2}{1 - kr^2} \right]$$

General Relativity gives the dynamics $\rightarrow R(t)$

Observable quantity $\frac{\dot{R}}{R} = H(t)$

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Testing the models

Geometry:

From the light geodesic equation:

$$c^2 dt^2 - R^2(t) \frac{dr^2}{1 - kr^2} = 0$$

General Mattig relation:

$$\int_{t_S}^{t_0} \frac{cdt}{R(t)} = \int_0^{r_S} \frac{dr}{(1 - kr^2)^{1/2}} = S_k^{-1}(r_S)$$

with:

$$S_k(u) = \begin{cases} \sin(u) & \text{if } k = +1 \\ u & \text{if } k = 0 \\ \sinh(u) & \text{if } k = -1 \end{cases}$$

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Geometrical tests

Most cosmological tests rely on a combination of $r(z)$ and $R(t)$

Angular distance:

$$\theta = \frac{d}{D_{\text{ang}}}$$

with

$$D_{\text{ang}} = R(t_S) r$$

Luminosity distance $l = \frac{L}{4\pi D^2}$

$$\begin{aligned} D_{\text{lum}} &= R(t_0) r (1+z) \\ &= R(t_S) r (1+z)^2 \\ &= D_{\text{ang}} (1+z)^2 \end{aligned}$$

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Application to SNIa

SNIa are bright objects that can be detected at large distances.

Up to $z \sim 2$ i.e. $t(z) \sim 3$ Gyr

SNIa are “standardizable”.

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Distant SNIa

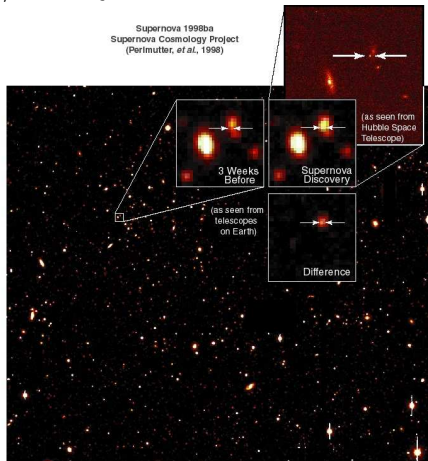
Just look for distant supernovae...

Distant SNIa

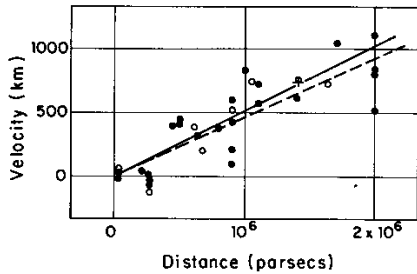
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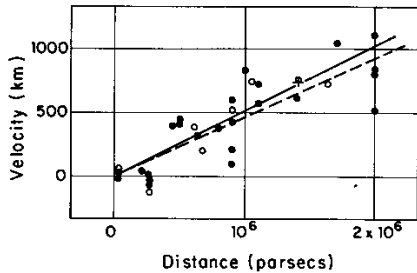
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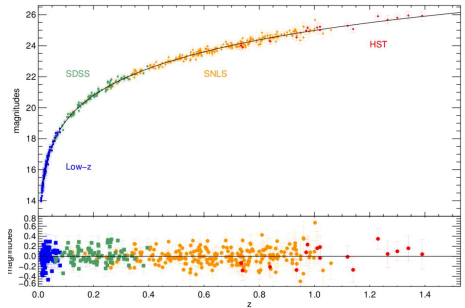
SN Ia Hubble diagram



SN Ia Hubble diagram



Acceleration!!



What does it mean?

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COSMOLOGY MARCHES ON



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In GR, the source of gravity is ρ and P :

$$\ddot{R} \propto -(\rho + 3P)R$$

Observations need $P \approx -\rho$

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So that the gravity strength is repulsive and proportional to R

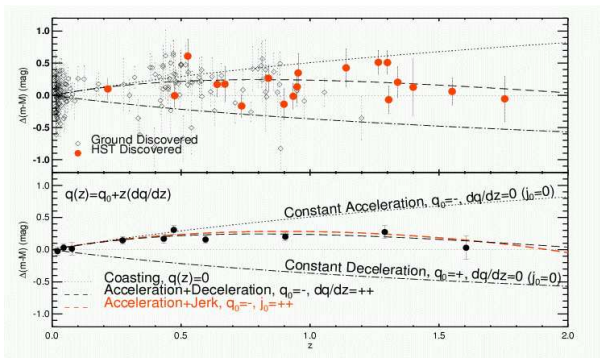
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$$\ddot{R} \propto -(\rho_m(1+z)^3 - 2\rho_v)R$$

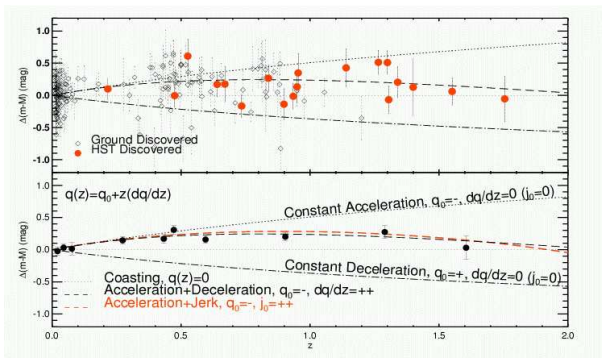
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→ Acceleration+decceleration!!

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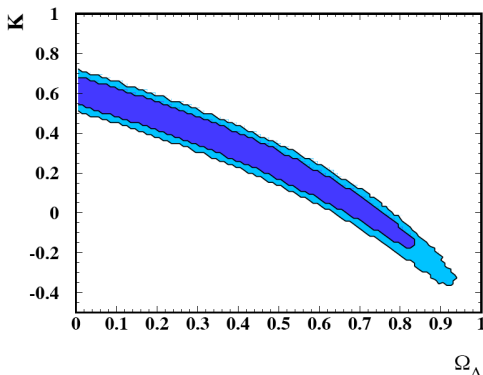
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Fit the Hubble diagram with K and Λ

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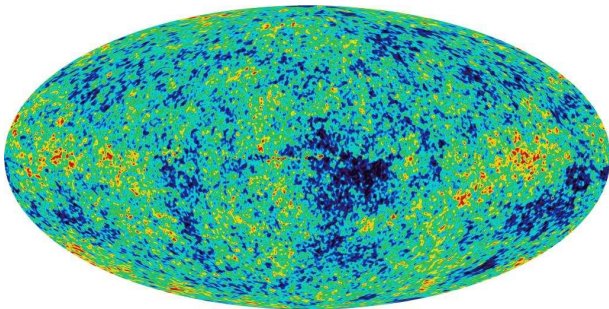
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L. Ferramacho, A. Blanchard, Y. Zolnierowski (2009)

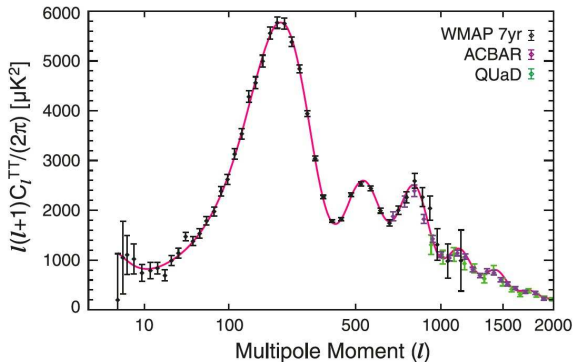
Cosmic microwave radiation fluctuations

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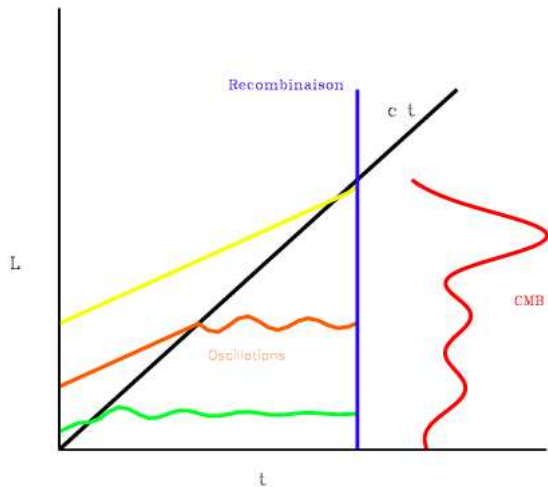


WMAP 1, 3, 5, 7,...

Cosmic microwave radiation fluctuations



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Cosmic microwave radiation fluctuations

Essentially geometric:

$$z_{lss} \approx 1090$$

Angular distance to the CMB is the key parameter, combined with the acoustic scale r_s corresponding to the sound horizon at l_s :

$$l_A = \frac{D_{ang}(z_{lss})}{r_s}$$

The shift parameter:

$$R = \sqrt{\Omega_m H_0^2 D_{ang}(z_{lss})} = 1.710 \pm 0.019$$

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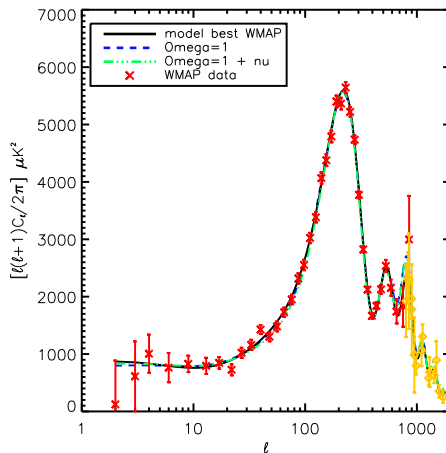
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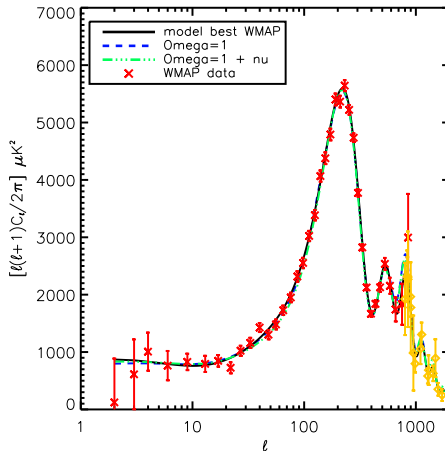
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Degeneracy in parameters allows to reproduce the C_ℓ
Blanchard et al., 2003

Conlusion (at this point)

Conclusion (at this point)

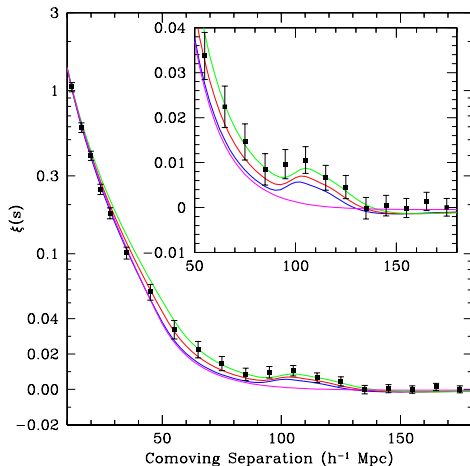
Neither SNIa nor CMB strongly require acceleration!

But...

The sound horizon is also imprinted in the matter distribution:

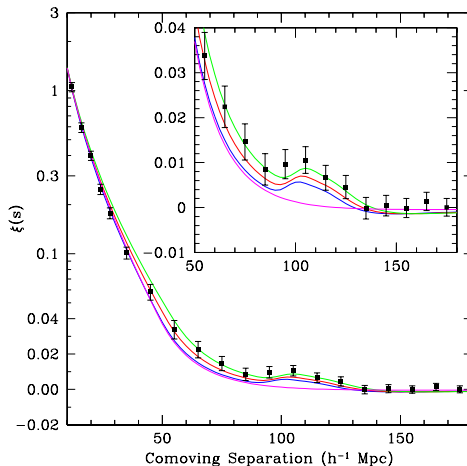
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This was a prediction of Λ CDM

Trensverse dilatation

The scale appears along the line of sight and trensverse to the line of sight:

$$D(z) = \left[(1+z)^2 D_{ang}^2 \frac{cz}{H(z)} \right]^{1/3}$$

From the SDSS LRG:

$$D(0.35) \sim 0.469 \pm 0.019$$

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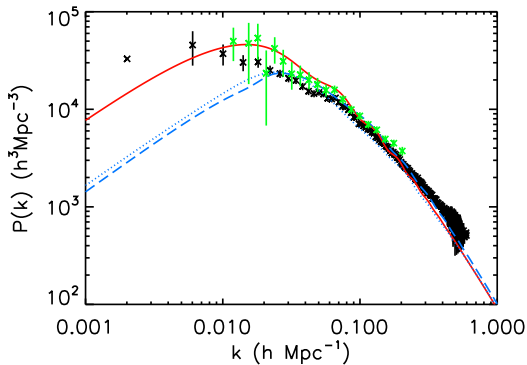
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Parameters in Λ CDM

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Parameters in Λ CDM

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SNIa, CMB, P(k)

Parameter	Vanilla	Vanilla + Ω_k	Vanilla + w	Vanilla + $\Omega_k + w$
$\Omega_b h^2$	0.0227 ± 0.0005	0.0227 ± 0.0006	0.0228 ± 0.0006	0.0227 ± 0.0005
$\Omega_c h^2$	0.112 ± 0.003	0.109 ± 0.005	0.109 ± 0.005	0.109 ± 0.005
θ	1.042 ± 0.003	1.042 ± 0.003	1.042 ± 0.003	1.042 ± 0.003
τ	0.085 ± 0.017	0.088 ± 0.017	0.087 ± 0.017	0.088 ± 0.017
n_s	0.963 ± 0.012	0.964 ± 0.013	0.967 ± 0.014	0.964 ± 0.014
$\log(10^{10} A_s)$	3.07 ± 0.04	3.06 ± 0.04	3.06 ± 0.04	3.06 ± 0.04
Ω_k	0	-0.005 ± 0.007	0	-0.005 ± 0.0121
w	-1	-1	-0.965 ± 0.056	-1.003 ± 0.102
Ω_Λ	0.738 ± 0.015	0.735 ± 0.016	0.739 ± 0.014	0.733 ± 0.020
Age	13.7 ± 0.1	13.9 ± 0.4	13.7 ± 0.1	13.9 ± 0.6
Ω_m	0.262 ± 0.015	0.270 ± 0.019	0.261 ± 0.020	0.272 ± 0.029
σ_8	0.806 ± 0.023	0.791 ± 0.030	0.816 ± 0.014	0.788 ± 0.042
z_{re}	10.9 ± 1.4	11.0 ± 1.5	11.0 ± 1.5	11.0 ± 1.4
h	0.716 ± 0.014	0.699 ± 0.028	0.713 ± 0.015	0.698 ± 0.037

L. Ferramacho, A. Blanchard, Y. Zolnierowski (2009)

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$\Omega_c h^2$	0.110 ± 0.004	0.109 ± 0.005	0.113 ± 0.005	0.111 ± 0.005
θ	1.042 ± 0.003	1.042 ± 0.003	1.042 ± 0.003	1.042 ± 0.003
τ	0.088 ± 0.017	0.087 ± 0.017	0.085 ± 0.017	0.085 ± 0.016
n_s	0.968 ± 0.013	0.965 ± 0.013	0.963 ± 0.014	0.960 ± 0.014
$\log(10^{10} A_s)$	3.07 ± 0.04	3.06 ± 0.04	3.07 ± 0.04	3.06 ± 0.04
Ω_k	0	-0.002 ± 0.007	0	-0.017 ± 0.013
w	-1	-1	-1.112 ± 0.148	-1.33 ± 0.242
K	-0.042 ± 0.042	-0.035 ± 0.042	-0.105 ± 0.091	-0.133 ± 0.077
Ω_Λ	0.747 ± 0.017	0.745 ± 0.020	0.756 ± 0.022	0.744 ± 0.022
Age	13.6 ± 0.1	13.7 ± 0.4	13.6 ± 0.1	14.5 ± 0.7
Ω_m	0.253 ± 0.017	0.257 ± 0.025	0.244 ± 0.022	0.272 ± 0.029
σ_8	0.801 ± 0.026	0.794 ± 0.029	0.846 ± 0.068	0.867 ± 0.060
z_{re}	11.1 ± 1.5	11.0 ± 1.4	10.9 ± 1.5	10.8 ± 1.4
h	0.725 ± 0.017	0.720 ± 0.036	0.748 ± 0.038	0.703 ± 0.042

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Possible origins:

The Einstein Cosmological constant...

Vaccum energy density:

Nerst-Weinberg argument:

$$\rho_v \sim \int^{k_c} k d^3\mathbf{k} \propto k_c^4 \approx 10^{121} \rho_{obs}$$

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Inhomogeneities

Non linear effect of GR

Standard model: background + linear perturbations

Non linear effects could contribute to the background

However $\delta h^2 \sim 10^{-10}$

Lemaître-Tolman-Bondi models.

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More fashionable solutions

A new component with like a scalar field such that :

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quintessence, ...

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Conclusion

Λ CDM has been predictive and successful!

The situation will improve with the next generation of experiments (BigBOSS, LSST, Euclid, Wfirst, ...)

A fundamentally new step is to test GR on large scale.

Measuring growth rate of structures: testing GR

Evidence for the Fifth Element, Astrophysical status of Dark Energy, 2010, Blanchard, A., *A&A Rev.*, **18**, 595.

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Evidence for the Fifth Element, Astrophysical status of Dark Energy, 2010, Blanchard, A., A&A Rev., 18, 595.

Λ CDM has been predictive and successful!

The situation will improve with the next generation of experiments (BigBOSS, LSST, Euclid, Wfirst, ...)

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