The Persistence of Attraction: the Dipolar Efimov Effect

Jose P. D'Incao

JILA, University of Colorado at Boulder and National Institute of Standards and Technology

(in colaboration with Y. Wang and C. H. Greene)



University of Colorado



Air Force Office of Scientific Research MURI

National Science Foundation

(Figure credit: Brad Baxley)





Universal Few-Body Physics



From the theoretical side:

- ✓ Signatures of Efimov Physics,
- Four- and More-bodies universal states,
- ▶ New families of universal states,

Ultracold Quantum Gases

(<peV)



clean and accurate experiments
 CONTROL of interactions (*B*-field)
 can explore the universal regime (low T and strong interactions)



Universal Few-Body Physics



From the theoretical side:

- ✓ Signatures of Efimov Physics,
- Four- and More-bodies universal states,
- ▶ New families of universal states,

Ultracold Quantum Gases

(<peV)

From the experimental side:

clean and accurate experiments
 CONTROL of interactions (*B*-field)
 can explore the universal regime (low T and strong interactions)

Allows for the use of simple models !!!



Universal Few-Body Physics



From the theoretical side:

- ✓ Signatures of Efimov Physics,
- Four- and More-bodies universal states,
- ▶ New families of universal states,

Ultracold Dipolar Quantum Gases

- CONTROL of interactions (*E*-field)
 new phases, quantum computing, etc ...
 long range anisotropic interactions
 - The Efimov effect persist !!!
 Dipolar interaction is extremely beneficial !!!
 Possible *new* few-body physics !!!

Few-body physics in Ultracold Gases (Why do Experiments care about it?)



Atomic/Molecular Losses

(three-body recombination, ...)





Few-body Physics in Ultracold Gases





Few-body Physics in Ultracold Gases











Efimov Physics

appearance of an *attractive* or *repulsive* three-body effective interaction ... in the strongly interacting regime $(|a| \gg r_0)$



Control of the few-body interactions



Scattering length dependence on 3-body collision rates





 \bigcirc

6

(a > 0)

 $F(10^4 r_0/a)$



appearance of an *attractive* or *repulsive* three-body effective interaction ... in the strongly interacting regime $(|a| \gg r_0)$

$$E_n = \frac{E_0}{[e^{\pi/s_0}]^{2n}}, \quad n = 1, 2, 3, \dots$$
$$a_n = [e^{\pi/s_0}]^n a_0 \qquad e^{\pi/s_0} \approx 22.7 \quad \text{(geometric factor)}$$

parameter,
of the full
n
$$3, ...$$

 $60 \approx 22.7$ (geometric factor)
 10^{-7}
 10^{-6}
 10^{-6}
 10^{-6}
 10^{-6}
 10^{-6}
 10^{-6}
 10^{-6}
 10^{-6}
 10^{-7}
 10^{-6}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}
 10^{-7}





- Universal properties of the properties of the four-body system with large scattering lengths, Hammer & Platter, *EPJA* 32, 113 (2007)
- Signatures of universal four-body phenomena and its relation to the Efimov effect von Stecher, D'Incao, and Greene, *Nat. Phys.* (2009)



Two four-boson states for each
 Efimov trimer [... see also Deltuva, FBS 50, 391 (2011)]

$$E_{4b}^{(n,m)} = c_m E_{3b}^{(n)} \qquad m = 1, 2$$

$$n = 1, 2, ..., \infty$$

$$(c_1 \approx 4.58, \ c_2 \approx 1.01)$$



- Universal properties of the properties of the four-body system with large scattering lengths, Hammer & Platter, *EPJA* 32, 113 (2007)
- Signatures of universal four-body phenomena and its relation to the Efimov effect von Stecher, D'Incao, and Greene, *Nat. Phys.* (2009)





- Universal properties of the properties of the four-body system with large scattering lengths, Hammer & Platter, *EPJA* 32, 113 (2007)
- Signatures of universal four-body phenomena and its relation to the Efimov effect von Stecher, D'Incao, and Greene, *Nat. Phys.* (2009)





- Universal properties of the properties of the four-body system with large scattering lengths, Hammer & Platter, *EPJA* 32, 113 (2007)
- Signatures of universal four-body phenomena and its relation to the Efimov effect von Stecher, D'Incao, and Greene, *Nat. Phys.* (2009)





- Universal properties of the properties of the four-body system with large scattering lengths, Hammer & Platter, EPJA 32, 113 (2007)
- Signatures of universal four-body phenomena and its relation to the Efimov effect von Stecher, D'Incao, and Greene, *Nat. Phys.* (2009)



The arrival of dipolar gases







Ultracold Dipolar gases

Atomic (magnetic) dipoles: weak interactions !!! Molecular (electric) dipoles: strong interactions !!! (challenging experiments)





... also : Innsbruck, Yale, Heidelberg, Hanover, ...

	d_ℓ^{\max}	kd_{ℓ}^{\max} (T = 100nK)
RbK	$\approx 6 \times 10^3 a_0$	2.4
m RbCs	$\approx 5 \times 10^4 a_0$	23
LiCs	$\approx 6 \times 10^5 a_0$	238
SrO	$\approx 1.1 \times 10^6 a_0$	467
IK	$\approx 2.8 \times 10^6 a_0$	1100







... also : Innsbruck, Yale, Heidelberg, Hanover, ...

"Universal" dipolar molecules: RbK, LiNa, LiK,LiRb, LiCs (unit probability of reactive collisions)

... reaction dynamics (ultracold controllable chemistry), ...

[see: Idziaszek and Julienne PRL 104, 113202 (2010); Zuchwski and Hutson, arXiv:1003.1418]

"Non-universal" dipolar molecules:

RbCs, NaK, NaRb, NaCs, KCs (no reactive collisions)

..., dipolar quantum phases, ...

Dipolar Few-body Physics (... what changes ?)













Dipole-Dipole Model Interaction













How about Few-body Physics ?


Ultracold Dipolar Few-body Physics



Three-body Dipolar Physics





Efimov Physics for dipoles ?

non-dipolar	dipolar	
large <i>s</i> -wave scattering length	<i>l</i> is not conserved (multiple partial waves contributes)	
ZRM offers an analytical solution	ZRM for dipoles (?) (no effective-range)	
Identical Bosons: J=0	J is also not conserved !!!	



Strong Losses: Tichnor & Rittenhouse PRL (2010) [pertubative treatment]





For N particles ...

$$\hat{H} = -\frac{1}{2\mu}\nabla_T^2 + \sum_{i < j} V(r_{ij})$$

... angles + set of non-compact coordinates $r_{ij} \rightarrow [0, \infty]$



For N particles ...

$$\hat{H} = -\frac{1}{2\mu} \nabla_T^2 + \sum_{i < j} V(r_{ij})$$

... the hyperspherical way !!!

$$\hat{H} = -\frac{1}{2\mu} \frac{d^2}{d^2 R} + \frac{\Lambda^2(\Omega)}{2\mu R^2} + V(R,\Omega)$$

... angles + set of non-compact coordinates $r_{ij} \rightarrow [0, \infty]$



hyperradius R: overall size (collective motion)

 $R \to [0,\infty]$

hyperangles $\{\Omega\}$: internal motion $\{\Omega\} \rightarrow [0, \propto \pi]$





Thursday, October 13, 2011



Bound and Scattering Properties

$$\left[-\frac{1}{2\mu}\frac{d^2}{dR^2} + U_{\nu}(R) - E\right]F_{\nu}(R) + \sum_{\nu'}W_{\nu\nu'}(R)F_{\nu'}(R) = 0$$

(Hyperradial Schrodinger Equation)



Thursday, October 13, 2011



Bound and Scattering Properties

$$\left[-\frac{1}{2\mu}\frac{d^2}{dR^2} + U_{\nu}(R) - E\right]F_{\nu}(R) + \sum_{\nu'}W_{\nu\nu'}(R)F_{\nu'}(R) = 0$$

(Hyperradial Schrodinger Equation)





Bound and Scattering Properties

$$\left[-\frac{1}{2\mu}\frac{d^2}{dR^2} + U_{\nu}(R) - E\right]F_{\nu}(R) + \sum_{\nu'}W_{\nu\nu'}(R)F_{\nu'}(R) = 0$$

(Hyperradial Schrodinger Equation)









$$\Psi_{M}(R,\Omega) = \sum_{\nu} F_{\nu}^{M}(R) \Phi_{\nu}^{M}(R;\Omega) \qquad (J \text{ is not conserved})$$
$$\Phi_{\nu}^{M}(R;\Omega) = \sum_{J} \sum_{K} \phi_{\nu K}^{J}(R;\theta,\varphi) D_{KM}^{J}(\alpha,\beta,\gamma)$$

$$\sum_{J'K'} \left[\left\langle JMK \mid \frac{\hat{\Lambda}(\Omega)}{2\mu R^2} \mid J'MK' \right\rangle \delta_{JJ'} + \left\langle JMK \mid \sum_{i>j} v_{dip}(\vec{r}_{ij}) \mid J'MK' \right\rangle - U_{\nu}(R) \right] \phi_{\nu K'}^{J'}(R;\theta,\varphi) = 0$$





$$\Psi_{M}(R,\Omega) = \sum_{\nu} F_{\nu}^{M}(R) \Phi_{\nu}^{M}(R;\Omega) \qquad (J \text{ is not conserved})$$
$$\Phi_{\nu}^{M}(R;\Omega) = \sum_{J} \sum_{K} \phi_{\nu K}^{J}(R;\theta,\varphi) D_{KM}^{J}(\alpha,\beta,\gamma)$$

$$\sum_{J'K'} \left[\left\langle JMK \mid \frac{\hat{\Lambda}(\Omega)}{2\mu R^2} \mid J'MK' \right\rangle \delta_{JJ'} + \left\langle JMK \mid \sum_{i>j} v_{dip}(\vec{r}_{ij}) \mid J'MK' \right\rangle - U_{\nu}(R) \right] \phi_{\nu K'}^{J'}(R;\theta,\varphi) = 0$$

$$\begin{split} \langle J'K'M'|v_{dd}(\vec{r}_{ij})|JKM\rangle &= \frac{d_{\ell}}{\mu_{2b}r_{ij}^{3}}(-1)^{K+M}\delta_{MM'}\sqrt{(2J+1)(2J'+1)} \times \\ & \left[\delta_{KK'} \begin{pmatrix} J & 2 & J' \\ K & 0 & -K' \end{pmatrix} \begin{pmatrix} J & 2 & J' \\ M & 2 & -M' \end{pmatrix} \right. \\ & \left. -\delta_{K-2,K'}\frac{3}{\sqrt{6}} \begin{pmatrix} J & 2 & J' \\ -K & 2 & K' \end{pmatrix} \begin{pmatrix} J & 2 & J' \\ -M & 0 & M' \end{pmatrix} \left(\frac{r_{ij}^{x} - i r_{ij}^{y}}{2}\right)^{2} \\ & \left. -\delta_{K+2,K'}\frac{3}{\sqrt{6}} \begin{pmatrix} J & 2 & J' \\ K & 2 & -K' \end{pmatrix} \begin{pmatrix} J & 2 & J' \\ M & 0 & -M' \end{pmatrix} \left(\frac{r_{ij}^{x} + i r_{ij}^{y}}{2}\right)^{2} \right] \end{split}$$





$$\Psi_{M}(R,\Omega) = \sum_{\nu} F_{\nu}^{M}(R) \Phi_{\nu}^{M}(R;\Omega) \qquad (J \text{ is not conserved})$$
$$\Phi_{\nu}^{M}(R;\Omega) = \sum_{J} \sum_{K} \phi_{\nu K}^{J}(R;\theta,\varphi) D_{KM}^{J}(\alpha,\beta,\gamma)$$

$$\sum_{J'K'} \left[\left\langle JMK \mid \frac{\hat{\Lambda}(\Omega)}{2\mu R^2} \mid J'MK' \right\rangle \delta_{JJ'} + \left\langle JMK \mid \sum_{i>j} v_{dip}(\vec{r}_{ij}) \mid J'MK' \right\rangle - U_{\nu}(R) \right] \phi_{\nu K'}^{J'}(R;\theta,\varphi) = 0$$

 $U_{\nu}(R)$: Adiabatic Potentials $P_{\nu}(R)$ and $Q_{\nu\nu'}(R)$: nonadiabatic couplings





$$\Psi_{M}(R,\Omega) = \sum_{\nu} F_{\nu}^{M}(R) \Phi_{\nu}^{M}(R;\Omega) \qquad (J \text{ is not conserved})$$
$$\Phi_{\nu}^{M}(R;\Omega) = \sum_{J} \sum_{K} \phi_{\nu K}^{J}(R;\theta,\varphi) D_{KM}^{J}(\alpha,\beta,\gamma)$$

$$\sum_{J'K'} \left[\left\langle JMK \mid \frac{\hat{\Lambda}(\Omega)}{2\mu R^2} \mid J'MK' \right\rangle \delta_{JJ'} + \left\langle JMK \mid \sum_{i>j} v_{dip}(\vec{r}_{ij}) \mid J'MK' \right\rangle - U_{\nu}(R) \right] \phi_{\nu K'}^{J'}(R;\theta,\varphi) = 0$$

 $U_{\nu}(R)$: Adiabatic Potentials $P_{\nu}(R)$ and $Q_{\nu\nu'}(R)$: nonadiabatic couplings

 $J_{\rm max} = 14$

50x50 system of 2D-PDE !!!

(very challenging !!!)





Wang, D'Incao, Greene, PRL 106, 233201 (2011)





Wang, D'Incao, Greene, PRL 106, 233201 (2011)





Wang, D'Incao, Greene, PRL 106, 233201 (2011)





Wang, D'Incao, Greene, PRL 106, 233201 (2011)



So what ?





(schematic representation)





(schematic representation)





(schematic representation)





(schematic representation)









 a_{3b}^{*-} : resonances in three-dipoles recombination a_{3b}^{*+} : interference in three-dipoles recombination a_{Dd}^{*} : resonances in dipole-dipolar dimer collisions





(3-body)	$a_{3\mathrm{b}}^{*-}/d_\ell \approx -8.1$	$a_{3\mathrm{b}}^{*+}/d_\ell \approx 1.8$	$a_{Dd}^*/d_\ell \approx 8.6$
(4-body)	$a_{4{\rm b},1}^{*-}/d_{\ell} \approx -3.5$	$a_{4{ m b},2}^{*-}/d_{\ell} \approx -7.3$	
	$a_{dd,1}^*/d_\ell \approx 20.$	$a_{dd,2}^*/d_\ell \approx 57.$	$a_{Dd}^c/d_\ell \approx 58.$









Three-fermionic dipoles effective potentials







Three-fermionic dipoles effective potentials

 * "Universal three-body physics for fermionic dipoles", Wang, D'Incao, Greene, arXiv:1106.6133 (2011)







Three-fermionic dipoles effective potentials

 * "Universal three-body physics for fermionic dipoles", Wang, D'Incao, Greene, arXiv:1106.6133 (2011)



Single Universal State !!! d_{ℓ}

Probability Density:

$$\rho(\vec{r}) = \frac{1}{3} \langle \Psi | \sum_{i} \delta(\vec{r} - \vec{r_i}) | \Psi \rangle$$



50



Probability Density:



Long-Lived, universal states !!!

$E_{2d} \rightarrow 0$	$d_\ell \; (r_0)$	$md_\ell^2 E_{3d}$	$m d_\ell^2 \Gamma$
	39.7	171	42
	58.2	135	43
	100	139	17
$\Theta \approx 16^{\circ}$	$b_l \approx 0.26 d_\ell$	$b_s \approx 0.14 d_\ell$	$\Delta \approx 15^{\circ}$












• Both the theoretical and experimental advances in ultracold quantum gases make these systems ideal candidates to explore universal few-body physics





• Both the theoretical and experimental advances in ultracold quantum gases make these systems ideal candidates to explore universal few-body physics

•The persistence of the Efimov effect for dipolar systems represents an important landmark for the characterization of universal properties as well as the expected stability of ultracold dipolar gases.





• Both the theoretical and experimental advances in ultracold quantum gases make these systems ideal candidates to explore universal few-body physics

•The persistence of the Efimov effect for dipolar systems represents an important landmark for the characterization of universal properties as well as the expected stability of ultracold dipolar gases.

• The effect of dipolar interaction might be responsible for the origin of a NEW class of universal few-body states





• Both the theoretical and experimental advances in ultracold quantum gases make these systems ideal candidates to explore universal few-body physics

•The persistence of the Efimov effect for dipolar systems represents an important landmark for the characterization of universal properties as well as the expected stability of ultracold dipolar gases.

• The effect of dipolar interaction might be responsible for the origin of a NEW class of universal few-body states

• Future experiments in ultracold dipolar gases are expected to offer a much better scenario to probe few-body universal states.