

# ***Some cosmological implications of “MiniBoone”***

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***RPP Clermont-Ferrand 14/01/2011***

# Status of neutrino mixing, circa 2010

mismatch between flavour and mass basis expressed via mixing matrix  $U$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \quad s_{lk} \equiv \sin \theta_{lk}, \quad c_{lk} \equiv \cos \theta_{lk}$$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}.$$

$$(\Delta m_{21}^2)_{\text{BF}} = 7.65 \times 10^{-5} \text{ eV}^2,$$

$$7.05 \times 10^{-5} \text{ eV}^2 \leq \Delta m_{21}^2 \leq 8.34 \times 10^{-5} \text{ eV}^2,$$

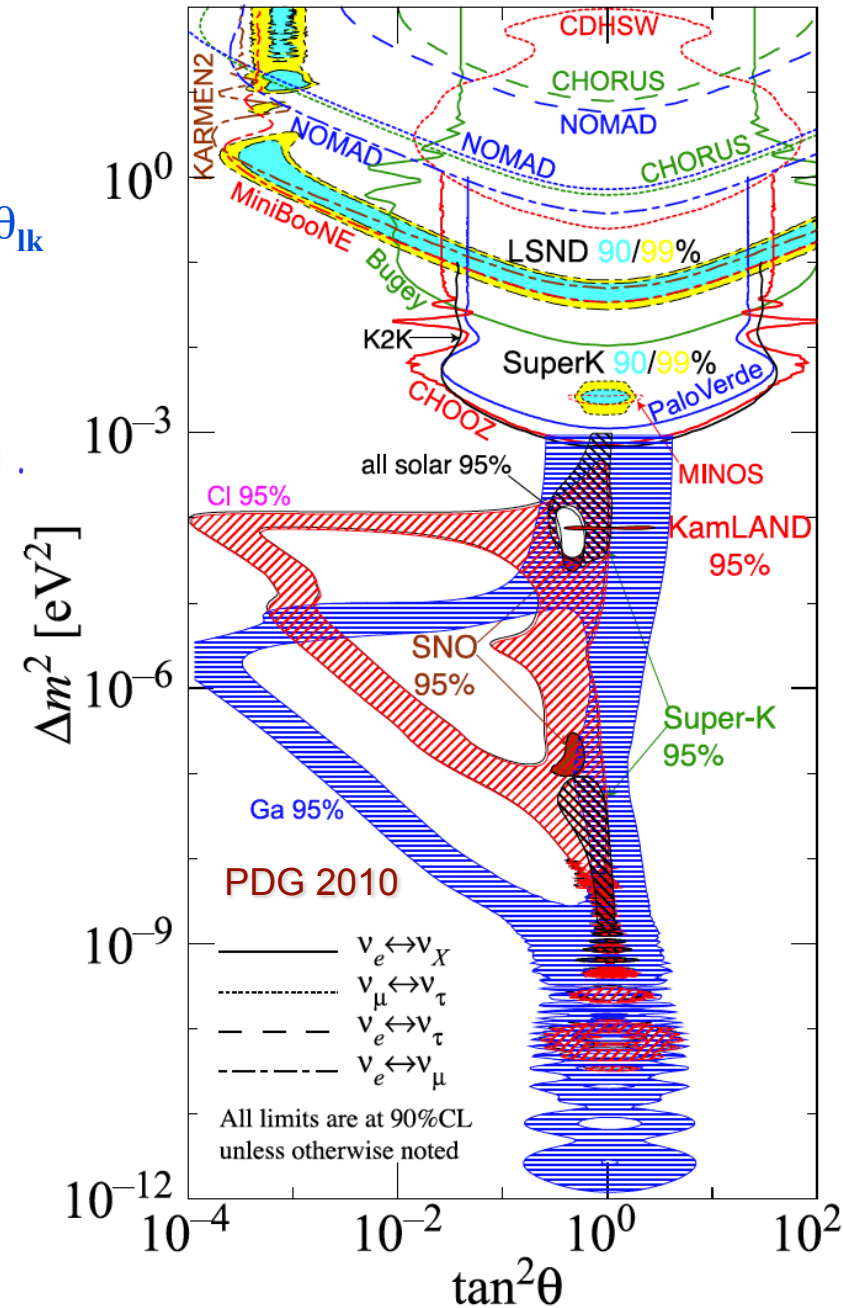
$$(\sin^2 \theta_{12})_{\text{BF}} = 0.304, \quad 0.25 \leq \sin^2 \theta_{12} \leq 0.37,$$

$$(|\Delta m_{31}^2|)_{\text{BF}} = 2.40 \times 10^{-3} \text{ eV}^2,$$

$$2.07 \times 10^{-3} \text{ eV}^2 \leq |\Delta m_{31}^2| \leq 2.75 \times 10^{-3} \text{ eV}^2,$$

$$(\sin^2 \theta_{23})_{\text{BF}} = 0.5, \quad 0.36 \leq \sin^2 \theta_{23} \leq 0.67.$$

$$\sin^2 \theta_{13} < 0.035 \text{ (0.056)} \quad \text{at } 90\% \text{ (99.73\%)} \text{ C.L.}$$



# Preamble: What does MiniBoone(\*) claim?

In a  $\nu_\mu$  beam above 475 MeV, we see no evidence for an excess of  $\nu_e$ -like events. (This is the region of maximal sensitivity if the LSND signal is L/E and CPT invariant.)

In a  $\nu_\mu$  beam below 475 MeV, we see a  $3\sigma$  excess ( $128 \pm 43$ ) of  $\nu_e$  signal candidates that don't fit well to a  $2\nu$  mixing hypothesis.

In a anti- $\nu_\mu$  beam below 475 MeV, we see a small excess ( $18 \pm 14$ ). It rules out some explanations of the  $\nu_\mu$  beam low-E excess. In a anti- $\nu_\mu$  beam above 475 MeV, we see an excess of events. The null hypothesis in the 475-1250 MeV region is only 0.5% probable. A  $2\nu$  fit prefers an LSND-like signal at 99.4% CL.

(\*) a dedicated experiment to test the LSND claim of  $\text{anti-}\nu_\mu \rightarrow \text{anti-}\nu_e$  appearance

from a Talk by Chris Polly (FNAL)  
@ NOW, 09/2010

# It does not fit in SM+3 massive $\nu$ : Interpretations?

- “Background” issue? E.g. something wrong with fluxes and  $\sigma$ 's estimated?
- Of course, one may solve the discrepancy by a radical departure from known physics (e.g. CPT violation, getting rid of QFT as we know it...)
- Less radical alternatives invoked include either (3+2) with CP, or 3+1 with NSI, but no real good fit/explanation of all the data is on the market!

With the addition of the new MiniBooNE data sets, **there are clear incompatibilities between neutrino and antineutrino experiments under a (3+1) oscillation hypothesis. A better description** of all short-baseline data over a (3+1) is **provided by a (3+2) oscillation hypothesis with CP violation. However, we still find large incompatibilities among appearance and disappearance experiments**, consistent with previous analyses, as well as incompatibilities between neutrino and antineutrino experiments.

*Karagiorgi et al. 0906.1997v3*

We have compared the quality of the (3+1) NSI fits to an updated fit in the (3+2) oscillation scheme[...] **Similarly to (3+1) NSI, in (3+2) the appearance experiments can be described very well.** However, we confirm previous results that **for (3+2) oscillations significant tension remains in the global fit between appearance and disappearance experiments.** [...] Let us mention also that in **none of the scenarios considered here we can explain the MiniBooNE low energy excess of events** when disappearance data are taken into account.

*E. Akhmedov and T. Schwetz. JHEP 1010, 115 (2010)*

# Purpose of the present talk

- ❖ To argue that even the most “conservative” BSM interpretations of these data have significant cosmological consequences.
- ❖ My main goal is to illustrate the basic physics through which new  $\nu$  states affect observables in BBN, CMB, LSS mostly via “extra radiation” and “extra mass”
- ❖ I will first review how the standard scenario goes... and also briefly address why adding NSI does not have a major impact.
- ❖ I will quickly present the present constraints on sterile  $\nu$ 's and rather stress the power of forthcoming data (PLANCK & co.)

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# The Birth of cosmological $\nu$ 's

**$T \gg 1 \text{ MeV}$**   
**Neutrinos in equilibrium**

$$f_\nu(p, T) = f_{FD}(p, T) = \frac{1}{e^{p/T} + 1} \quad T_\nu = T_e = T_\gamma$$

Above  $\sim \text{MeV}$ -scale temperatures,  $e^\pm$  pairs  
can be created “Boltzmann unsuppressed”.  
 $\nu$ 's are populated (& reach a thermal distribution)  
via reactions of the kind

$$\nu_a \nu_b \leftrightarrow \bar{\nu}_a \bar{\nu}_b$$

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**They decouple from the plasma at  $T \sim \mathcal{O}(1) \text{ MeV}$**

*Rate of weak processes*

$$\Gamma_w \approx n \sigma c \approx g a^{-3} G_F^2 E^2 \approx g G_F^2 T^5$$

*Hubble expansion rate*

$$H \approx \sqrt{G_N \rho} \approx \sqrt{g G_N} T^2$$

$$\frac{\Gamma_w}{H} \approx \left( \frac{T}{\text{MeV}} \right)^3$$

After this epoch ( $\sim \mathcal{O}(1) \text{ s}$  after Big Bang)  $\nu$ 's evolve only due to gravity



# “Detection” of the CνB

- Pseudo-thermal distribution:  $T_\nu = 1.95 \text{ K}$
- Number density ( $\nu + \bar{\nu}$ ):  $112 \text{ cm}^{-3}/\text{flavour}$
- Mean kinetic energy:  $\ll \text{meV}$

*lower than 2.7 K of  
CMB due to later  
 $e^+ e^- \rightarrow \gamma \gamma$   
(heating of photons)*

Direct searches hopeless?



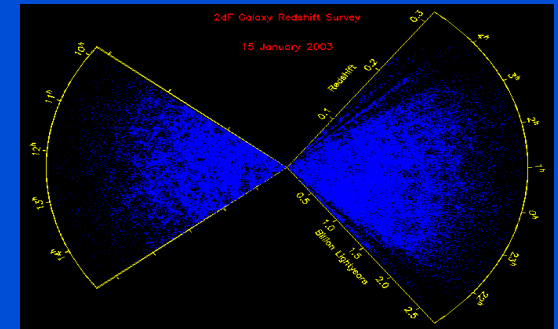
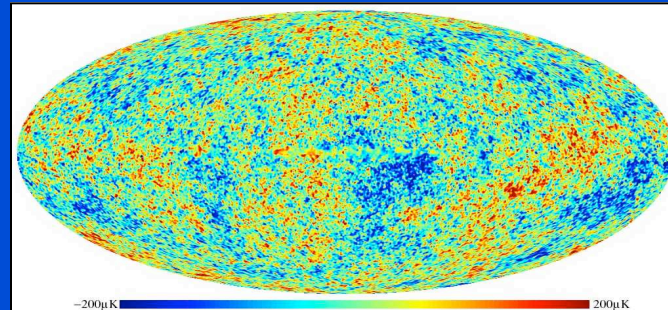
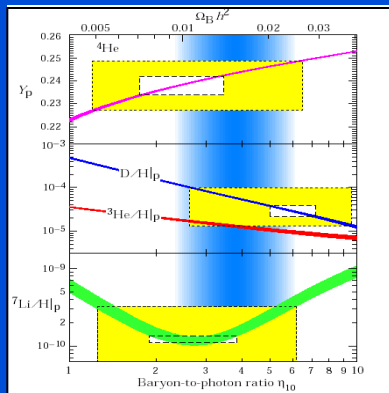
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Indirect searches: Cosmological observables



**BBN**

$T \sim \text{MeV}$

$\nu_e$  vs.  $\nu_{\mu,\tau}$        $N_{\text{eff}}$

**CMB**

$T \sim \text{eV}$

Gravity only (no flavor discr.)

**LSS**

$N_{\text{eff}}$  &  $m_\nu$

# Neutrinos & BBN: How do $\nu$ 's enter the game?

## Hubble Expansion Law

$$H = \frac{\dot{a}}{a} = \left( \frac{8\pi G_N}{3} \right)^{1/2} (\rho_\gamma + \rho_e + \rho_b + \rho_\nu + \rho_X)^{1/2}$$

$$\rho_\nu + \rho_X \rightarrow \frac{7}{8} \frac{4^{1/3}}{11^{1/3}} N_{\text{eff}} \rho_\gamma$$

$$N_{\text{eff}} = 3$$

(SM only & instantaneous decoupling)

*Gravity only, mostly integral quantity, extra relativistic species*

For a review, see e.g. F. Iocco et al.

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## Weak Rates: $p \leftrightarrow n$ equilibrium



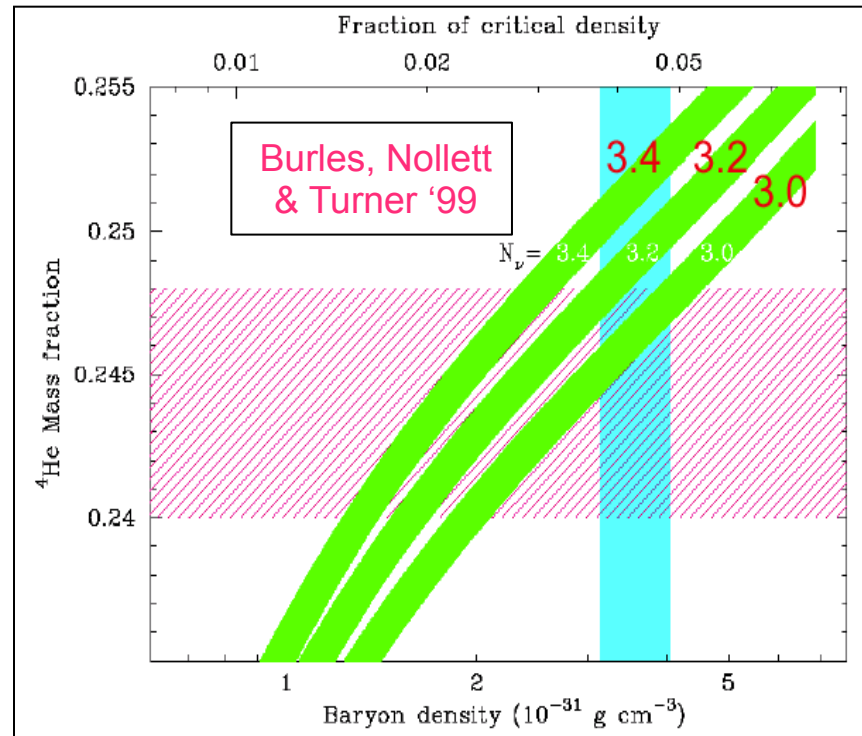
*Very sensitive to weak interactions (only e-flavour matters), energy spectrum.*

Final  $n/p$  (& hence  ${}^4\text{He}$ , where most neutrons are ultimately locked) depends on “when”  $\Gamma_w = H$

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# Estimating $^4\text{He}$ response to parameter changes



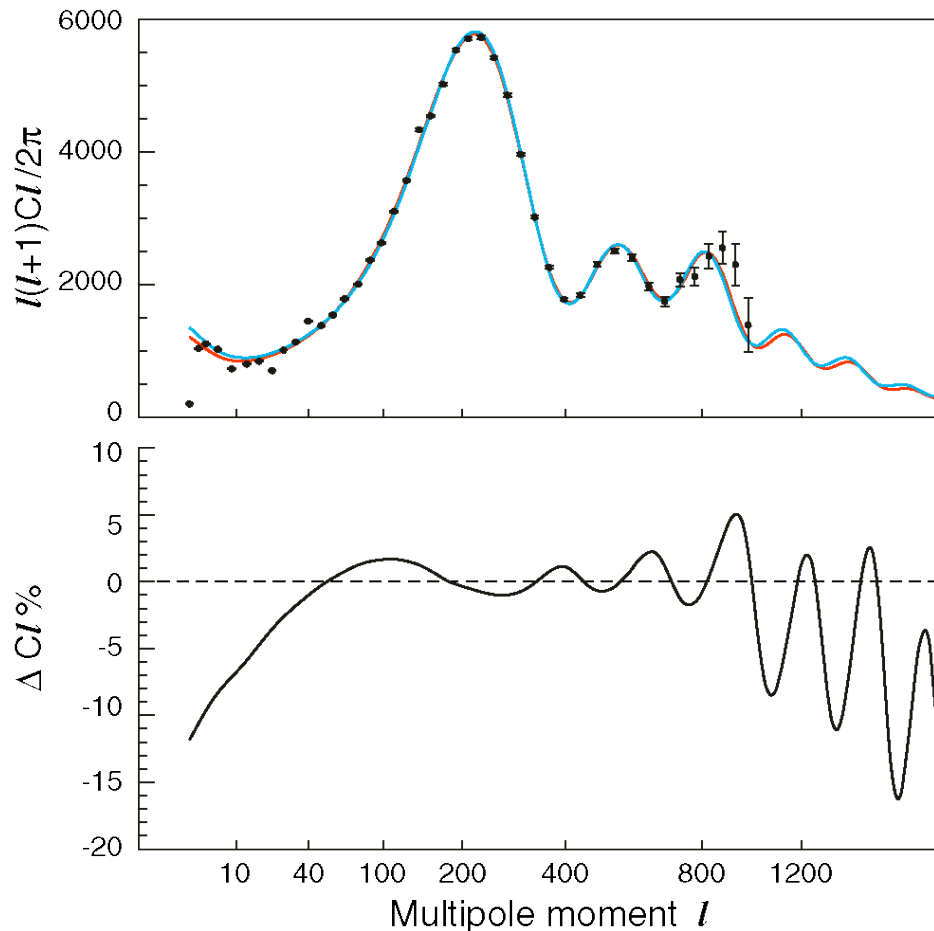
- High  $N_{eff} \rightarrow$  High  $H \rightarrow$  early freeze out ( $\Gamma_{pn} \sim H$  at high  $T$ )  $\rightarrow$  high  $n/p \rightarrow$  high  $Y_p$
- $\nu_e > \bar{\nu}_e \rightarrow \nu_e n \rightarrow e^- p$  favored over  $\bar{\nu}_e p \rightarrow e^+ n \rightarrow$  low  $n/p$  at fr.out  $\rightarrow$  low  $Y_p$   
(chemical potential  $\mu_{\nu_e} > 0$ )

• ...

# Neutrinos & CMB

We know that neutrinos are light (eV-scale at most); for this range, both  $m_\nu$  and  $N_{\text{eff}}$  mostly affect the time of matter-radiation equality. All the rest fixed:

- Raising  $N_{\text{eff}}$  means more radiation, hence delayed equality.
- Raising  $m_\nu$  means that part of the total that we call now (dark) matter was behaving as  $\sim$ radiation at CMB formation, hence delayed equality.



$$N_{\text{eff}} = 3.046$$

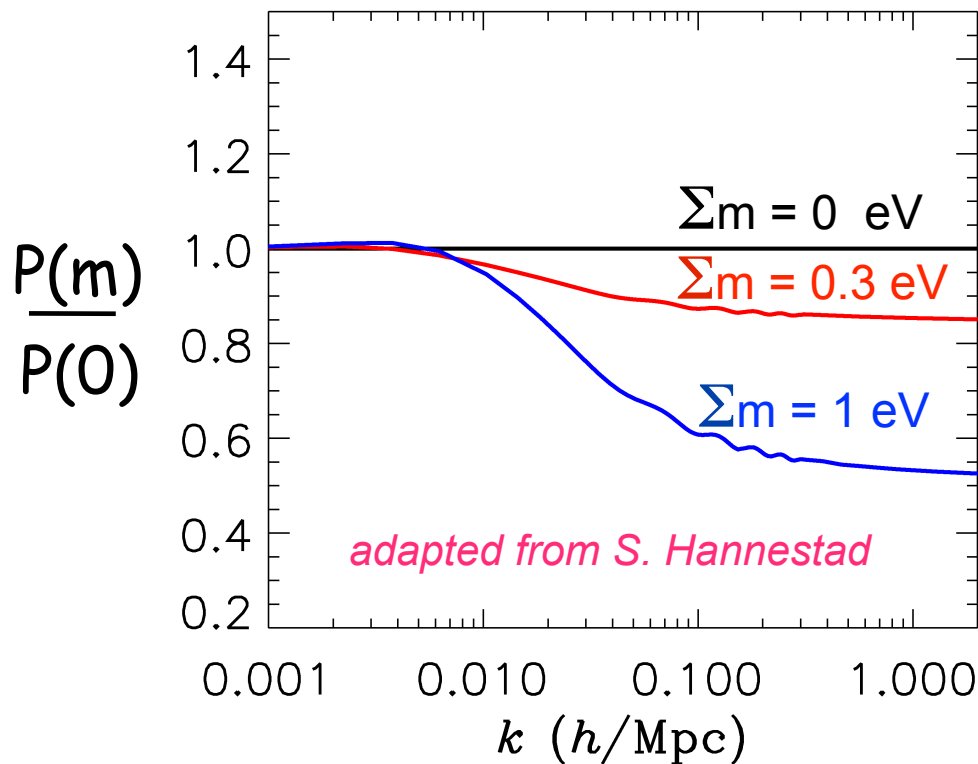
$$N_{\text{eff}} = 0$$

WMAP-7  
Dunkley, et al., 2009,  
*ApJS*, 180, 306-329

# Suppression of power-spectrum due to Neutrinos

$\nu$ 's do not contribute to gravitational clustering below the free-streaming scale, but they do contribute to the homogeneous expansion. This “unbalance” introduces a peculiar spectral suppression. In linear theory one finds

$$\frac{\Delta P}{P} \approx -8 \frac{\Omega_\nu}{\Omega_m} \approx -0.8 \frac{\Sigma m_i}{1 \text{ eV}} \frac{0.1}{\Omega_m h^2} \quad @ k > k_{NR} \approx 0.015 (\Sigma m_{eV} \times \Omega_m h^2)^{1/2} \text{ Mpc}^{-1}$$



*This is the key effect used to derive bounds on massive neutrinos from LSS*

Alteration in presence of NSI

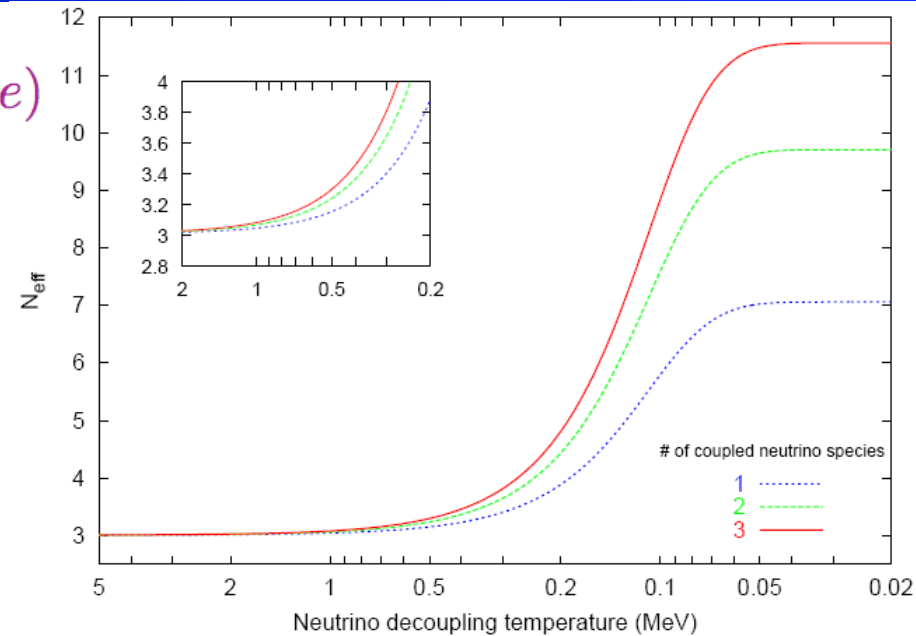


# “Altered Decoupling”: Non-Standard Interactions

$$\mathcal{L}_{\text{NSI}}^{\alpha\beta} = -2\sqrt{2}G_F \sum_P \epsilon_{\alpha\beta}^P (\bar{\nu}_\alpha \gamma^\mu L \nu_\beta) (\bar{e} \gamma_\mu P e)$$

❖ If FCNC or non-universal 4f Fermi-like operators present,  $\nu$  decoupling might change (typically delayed) :  $\nu$ 's get more energy from  $e^+e^-$  annihilations;  
Both **collisional** and **refractive** terms affected!

*G. Mangano et al.  
Nucl.Phys. B 756, 100 (2006)*

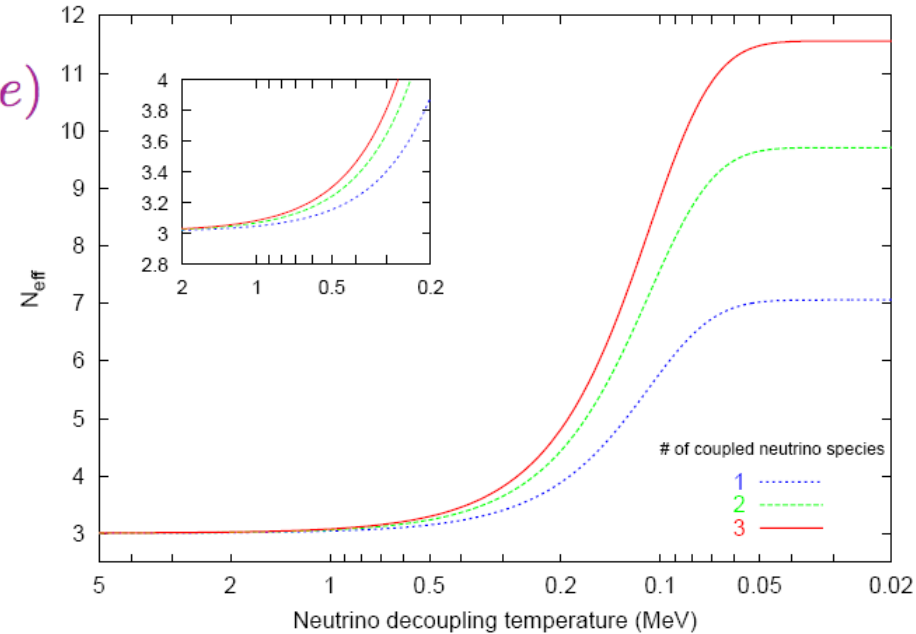


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- ❖ We found that no more than  $\Delta N_{\text{eff}} = \mathcal{O}(0.1)$  (but anyway  $\Delta Y_p = 0.001$  due to cancellations) might be obtained for current upper bounds, for  $\epsilon$ 's  $\sim \mathcal{O}(1)$ . Likely, improvements in the  $\nu_\tau$  sector bounds (e.g. from OPERA) will exclude these cases.

NSI of  $\mathcal{O}(0.01)$  required to fit Miniboone, see e.g. *E. Akhmedov and T. Schwetz, JHEP 1010, 115 (2010)* have at most a sub-leading effect in the early universe

Adding sterile states...

# The Quantum Zeno effect (for production via osc.)

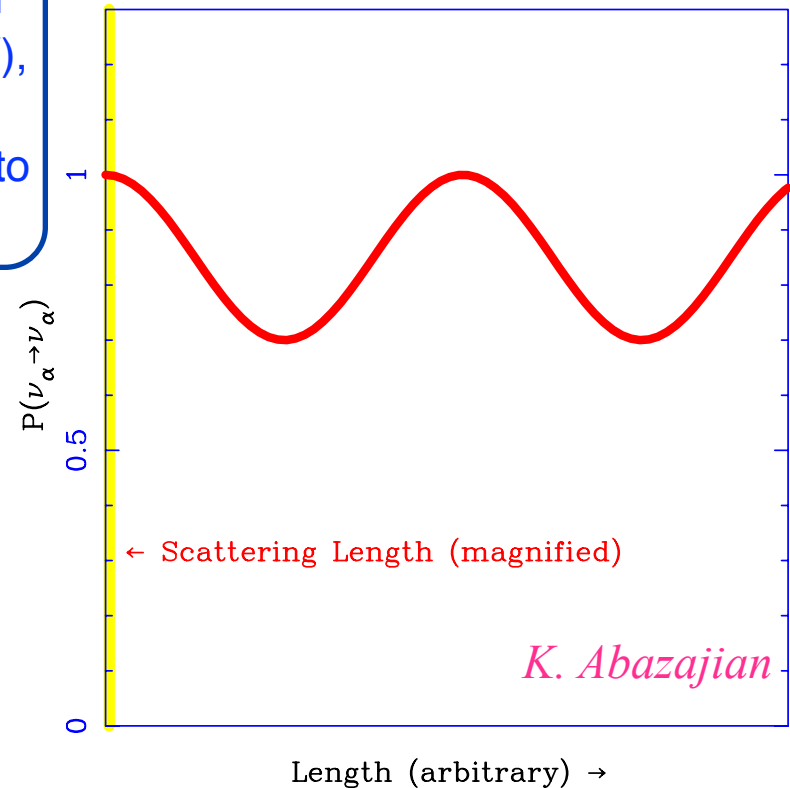
Each CC scattering of a  $\nu$  acts as a “measurement” of its flavor state. At high temperatures (say,  $T \geq 100$  MeV),  $\lambda_{\text{scatt}}$  is extremely short compared to  $\lambda_{\text{osc}}$ . Therefore, a population of active  $\nu$ 's won't have time to evolve into sterile  $\nu$ 's, but in small amounts.

$$\lambda_{\text{scatt}} = [\sigma n]^{-1} \sim E^{-2} T^{-3} \propto T^{-5}$$

$$\lambda_{\text{osc}} = \frac{4\pi E}{\Delta m^2} (\text{in vacuo}) \propto T$$

$$P_{\alpha\alpha}(\lambda_{\text{scatt}}) = 1 - \sin^2(2\theta) \sin^2 \left( \pi \frac{\lambda_{\text{scatt}}}{\lambda_{\text{osc}}} \right)$$

Suppression of  $\nu_s$  Production at Early Times



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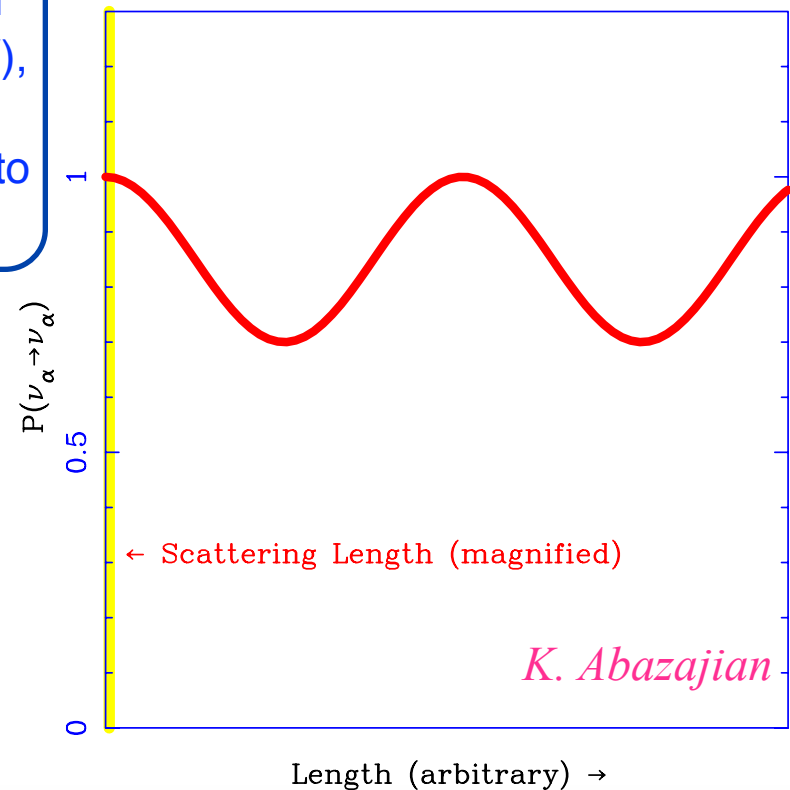
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Suppression of  $\nu_s$  Production at Early Times



As the universe expands, cools & becomes less dense,  $\lambda_{\text{scatt}} \nearrow$ . Then,  $P_{\text{as}} = (1 - P_{\alpha\alpha}) \nearrow$

☞ The larger  $\Delta m^2$ , the faster  $\nu$ 's oscillate, the higher the conversion  $P_{\text{as}}$

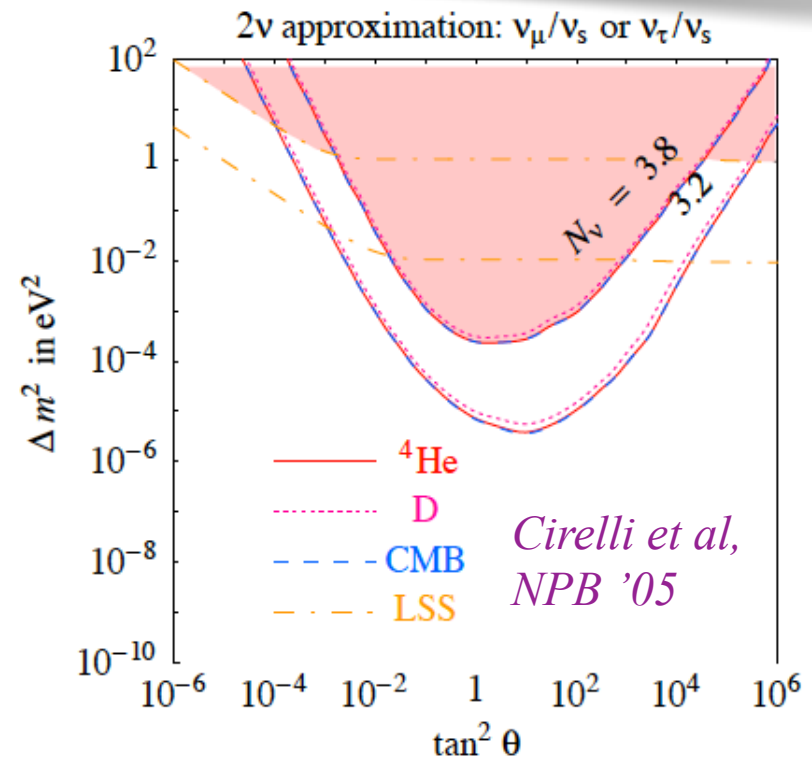
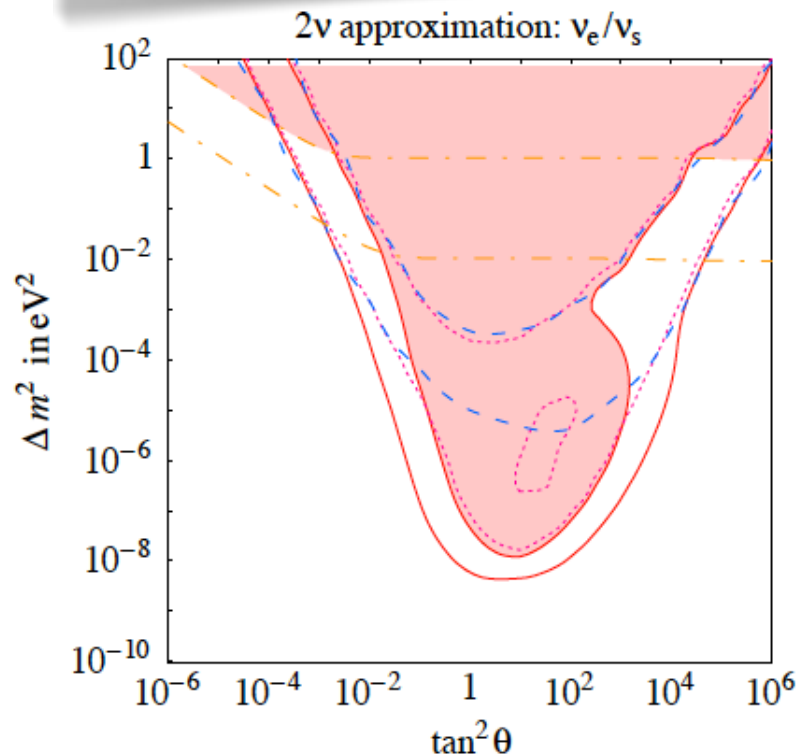
☞ Also, the larger  $\theta^2$ , the larger  $P_{\text{as}}$

⇒ Oscillation constraints from Early Universe exclude upper-right portion in a  $\theta^2$ - $\Delta m^2$  plot

# Sterile neutrinos are born

- \* If oscillations are effective before decoupling: the additional species can be brought into equilibrium:  $N_{\text{eff}}=4$
- \* If oscillations are effective after decoupling:  $N_{\text{eff}}=3$  but the spectrum of active neutrinos is distorted (direct effect on n/p equilibrium!)

Matter effects are responsible for the hierarchy dependence (resonant vs. non-resonant case) See e.g. Kirilova '03, Dolgov & Villante, NPB 679 (2004)...



# “Realistic” 3+2 $\nu$ 's scenarios for LSND/MB

If multiple sterile states are present (as in 3+2 schemes) in general there is: i) partial thermalization ( $4 < N_{\text{eff}} < 5$ ) ii) spectral *distortions* at BBN times  
see e.g. Melchiorri et al. JCAP 01 (2009) 036

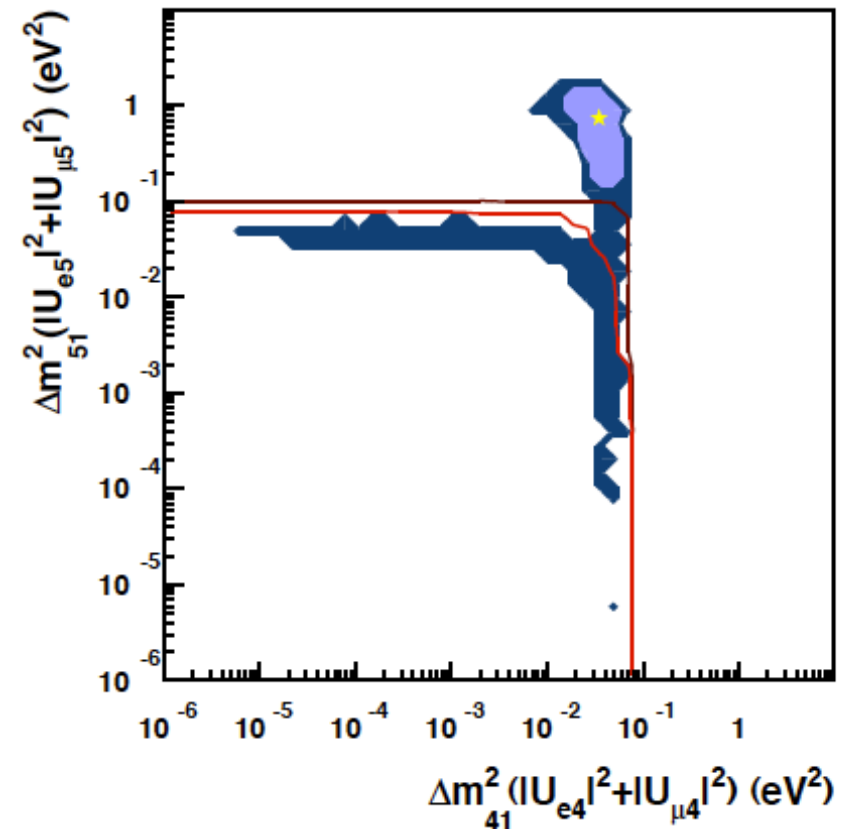


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👉 Simultaneous fits are still possible (goodness of fit dominated by Lab exp.)



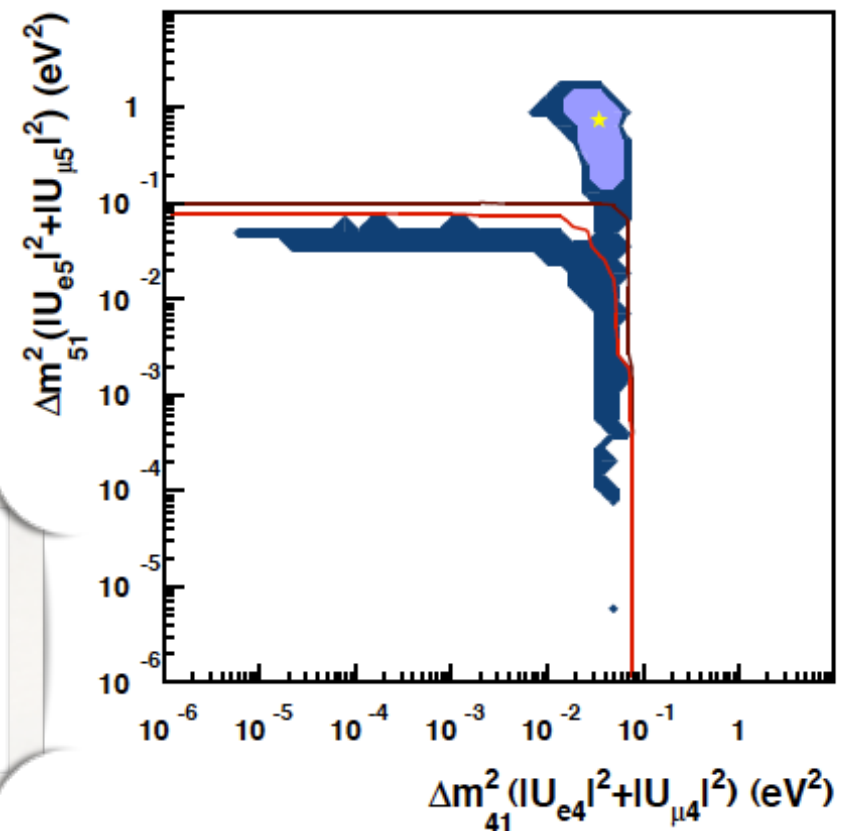
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**Note:** the required CP-violation might further affect weak rates at BBN time (yet to be studied). A general-scan approach to the problem is numerically challenging (many scales involved...) and perhaps unjustified at present.

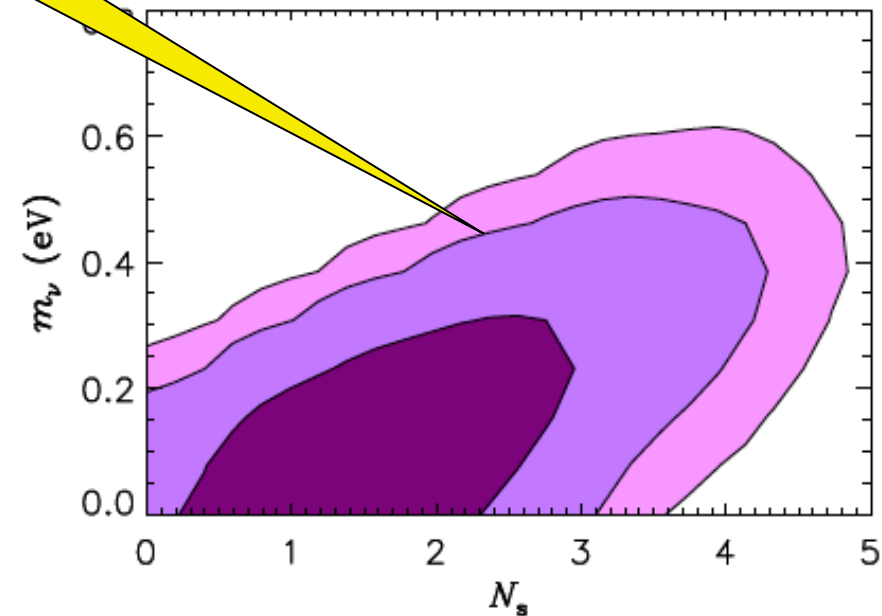
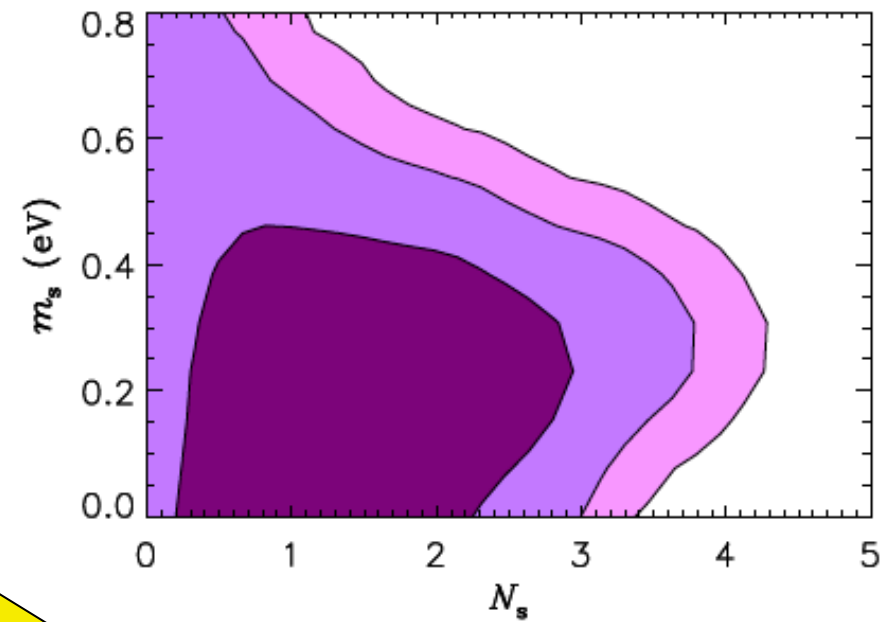


# Present status on $N_{\text{eff}}$ ... and $\Sigma m_\nu$

From cosmology alone, all data are consistent with expectations within (less than)  $2\sigma$

Not surprisingly, a  $m_\nu - N_{\text{eff}}$  degeneracy is due to CMB sensitivity to the quantities

*J. Hamann et al., PRL 105, 181301 (2010)*



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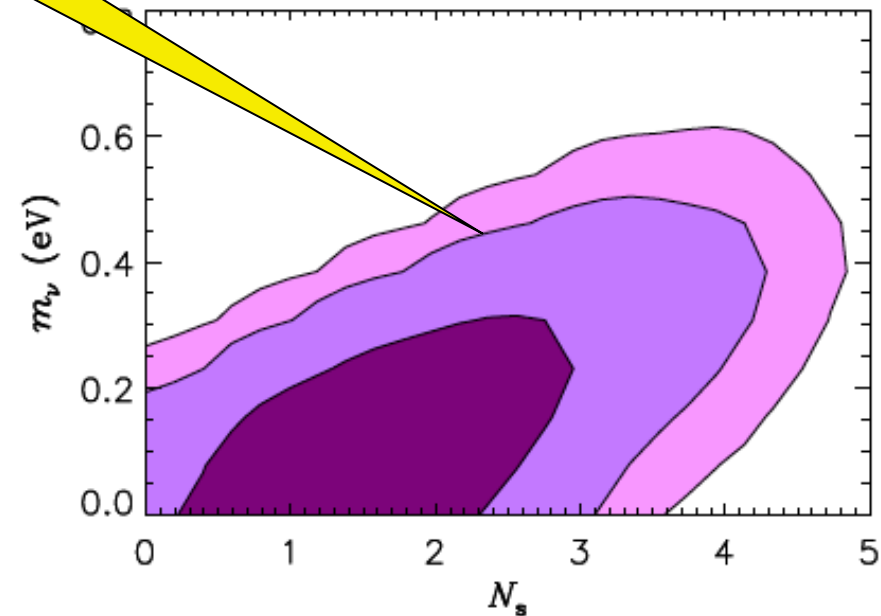
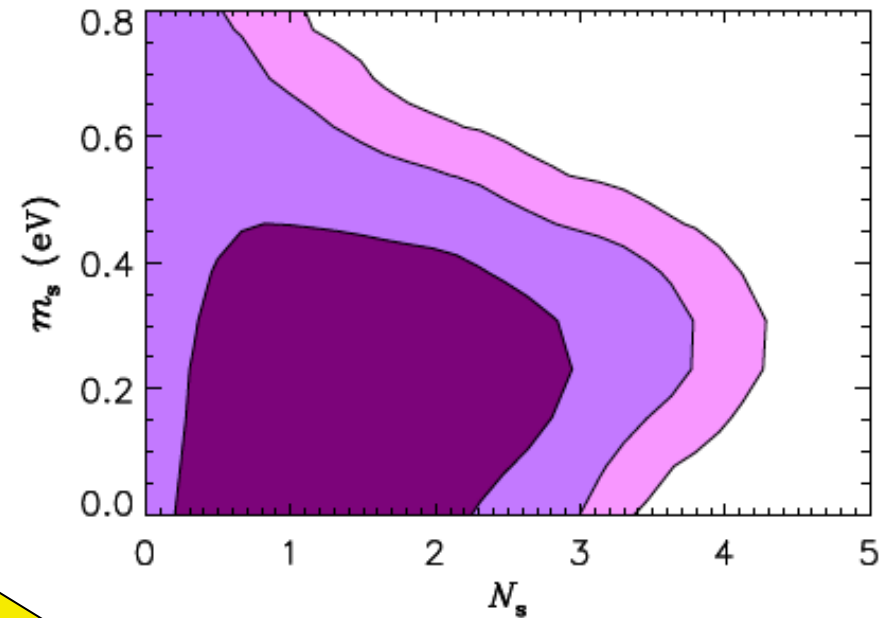
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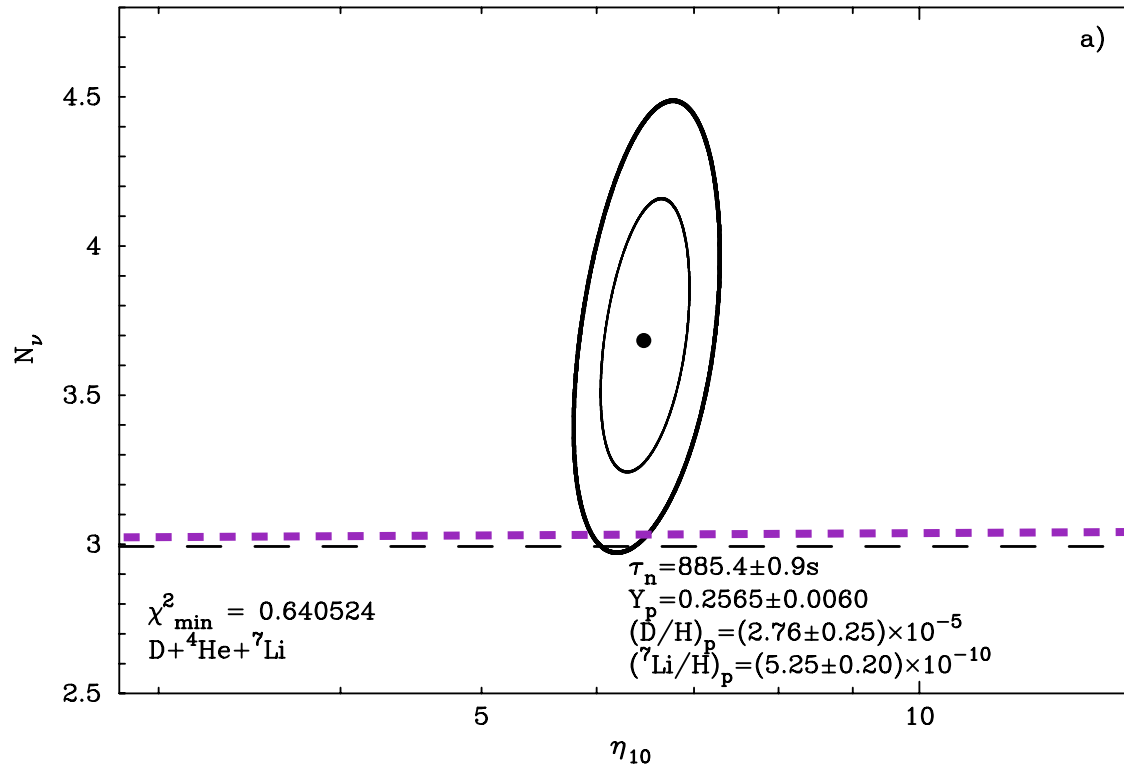
*J. Hamann et al., PRL 105, 181301 (2010)*

In the literature one finds also several (sometimes strong) claims about “hints for non-standard values”.

I would call cosmological “discrepancies at 1 to 2  $\sigma$ ’s” simply as “agreement”



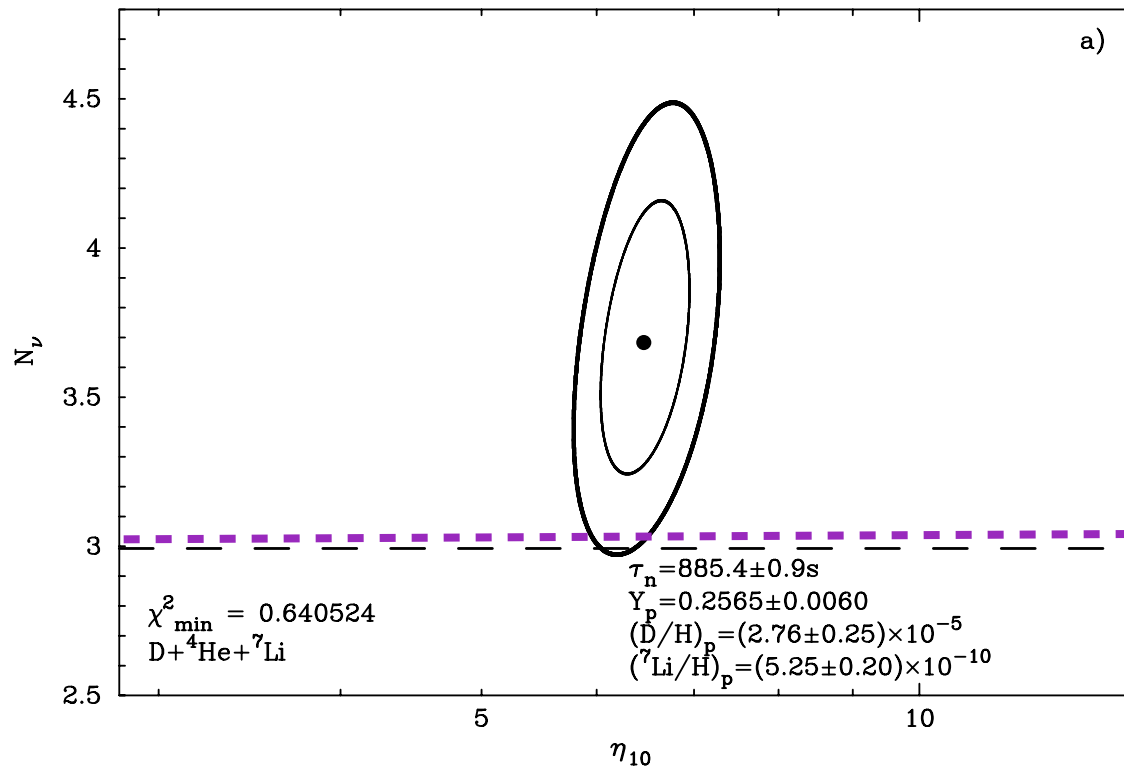
# A hint of a large $N_{\text{eff}}$ from BBN?



*Izotov and Thuan,  
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for non-standard big bang nucleosynthesis,”  
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$$Y_p = 0.2565 \pm 0.0060$$

Not surprisingly, in *Aver et al. JCAP 1004,029 (2010)*  
such an “evidence” is claimed to follow from  
underestimated systematics...

# Forecast on future reach on $N_{\text{eff}}$ ...

*It is very difficult to make accurate predictions, especially about the future  
(Niels Bohr)*



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## PLANCK CMB mission

Sensitivity  $\Delta N_{\text{eff}} \sim 0.05-0.1$

Realistically (including degeneracies)  $\Delta N_{\text{eff}} \sim 0.2$

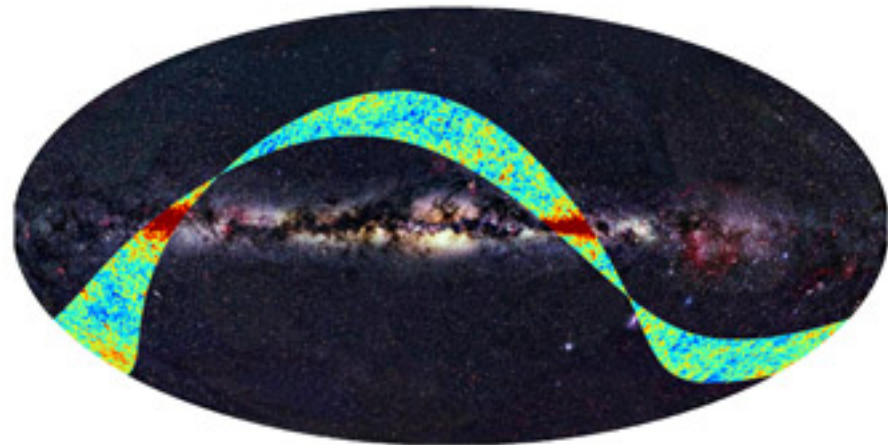
Lopez et al. PRL 82 3952 '99

Bowen et al., MNRAS 334 '02

Bashinsky & Seljak, PRD 69 '04

Hamann et al JCAP 07 (2010) 022

In a few years we should know if there is any other major ( $>10\%$ ) component of light degrees of freedom (like axions, gravitinos,  $\nu_R$ ) in the “cosmic soup”



## ...and $\Sigma m_\nu$

	Planck	P+BAO	P+HPS	P+HST	P+HST+BAO	P+HST+HPS
$\omega_{\text{dm}}$	0.22	0.24	0.20	0.21	0.21	0.19
$N_{\text{eff}}$	0.21	0.21	0.22	0.21	0.21	0.22
$\Sigma m_\nu$	0.68	0.81	0.44	0.67	0.73	0.44
$w$	2.14	1.16	0.72	0.74	0.76	0.55
$n_s$	0.46	0.48	0.49	0.46	0.48	0.48

Hamann et al JCAP 07 (2010) 022

Very likely, extra states would leave their imprints also due to their mass,  
if heavier than 0.3-0.5 eV...

In addition with future surveys, this number is going to improve even more!

# ...and $\Sigma m_\nu$

	Planck	P+BAO	P+HPS	P+HST	P+HST+BAO	P+HST+HPS
$\omega_{\text{dm}}$	0.22	0.24	0.20	0.21	0.21	0.19
$N_{\text{eff}}$	0.21	0.21	0.22	0.21	0.21	0.22
$\Sigma m_\nu$	0.68	0.81	0.44	0.67	0.73	0.44
$w$	2.14	1.16	0.72	0.74	0.76	0.55
$n_s$	0.46	0.48	0.49	0.46	0.48	0.48

Hamann et al JCAP 07 (2010) 022

Very likely, extra states would leave their imprints also due to their mass, if heavier than 0.3-0.5 eV...

In addition with future surveys, this number is going to improve even more!

Mission/Method	$\Sigma m_i$ (eV)	$\sigma(\Sigma m_i)$ (eV)	Ref.
PLANCK + Weak Lensing surveys	0.07	0.04	Hannestad et al. '06
Inflation Probe with Lensing	0.00-0.05	0.035	Lesgourgues et al. '06
Galaxy Cluster surveys+other	0.00-0.05	$\sim 0.034$	Wang '05
CMBpol, Lensing, Cosmic Shear	$\sim 0.05$	$\sim 0.013$	Song & Knox '04
CMB, SKA	$\sim 0.05$	$\sim 0.015$	Abdalla & Rawlings '07

# Summary

- ✓ Extra physics in neutrinos (NSI, light sterile states, w or w/o CP violation...) is recurrently invoked to explain some Laboratory neutrino anomalies, like the long-standing LSND one (recently made more puzzling by MiniBoone results).
- ✓ The Standard Hot Big Bang model predicts the existence of a CνB. Although a direct detection seems unlikely, CνB effects on cosmological probes as CMB, BBN and LSS are significant
- ✓ At present, it is fair to say that cosmology is consistent with standard expectations, although the data allow (or marginally favour) a  $\Delta N_{\text{eff}} = \mathcal{O}(1)$   
Alone,  $N_{\text{eff}} \neq 3$  does not necessarily indicate the presence of new  $\nu$  dof: might be due to NS neutrino interactions (likely too small!) asymmetries (excluded but in fine-tuned circumstances), or dof unrelated to  $\nu$  (like axions, ultralight gravitinos, etc.)
- ✓ If “LSND, MiniBoone ... anomalies” are due to sterile neutrinos, an interesting check will be allowed by forthcoming determination of  $N_{\text{eff}}$  (Planck) and a simultaneous imprint of “ $\nu_s$  mass” in CMB/LSS (sterile  $\nu$ ’s are not pure radiation!)

Extra slides

# What do we know about ${}^4\text{He}$ ?

## Main problem

We cannot observe *primordial* abundances:  
Stars have altered the primordial composition.  
For  ${}^4\text{He}$ , stars mostly burn H into He  $\rightarrow Y > Y_p$

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*Observe systems with little chemical processing*

HeII  $\rightarrow$  HeI recombination lines in HII regions  
(about  $\sim 80$  such regions known) of Blue  
Compact Dwarf Galaxies\*



\*small galaxies ( $\sim 1/10$  MW) containing large clusters of young, hot, massive stars. Among the least chemically evolved objects known.



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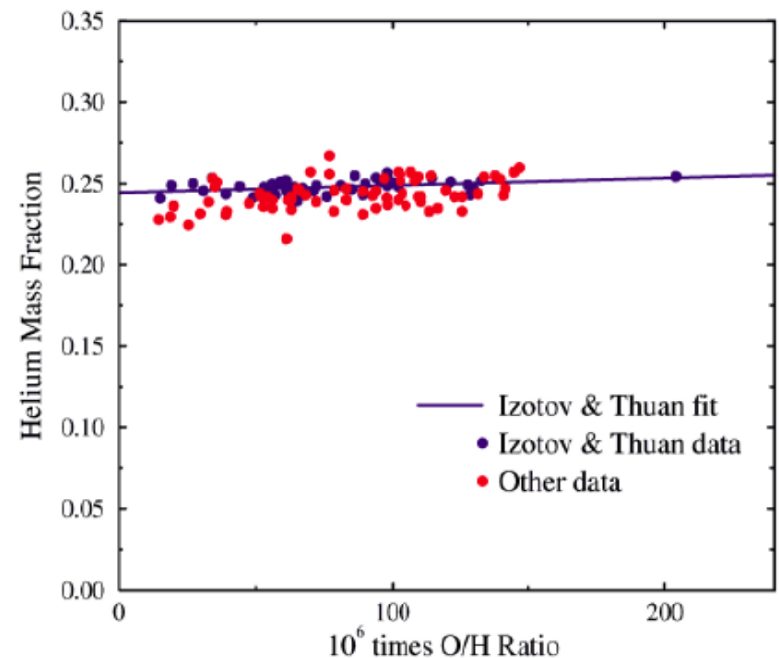
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*Correct for chemical evolution*

Extrapolate *linearly* to “zero  
metallicity” in  $Y_p$  vs O/H, N/H plots



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# Observational Status over the last years

Not simple environment, systematics dominate!

Crucial to infer (consistently, if possible) the properties of the medium: photoionization, excitations, recombinations... must be modelled and fit!

✓ Izotov et al. '04,  $Y_p = 0.2421 \pm 0.0021$

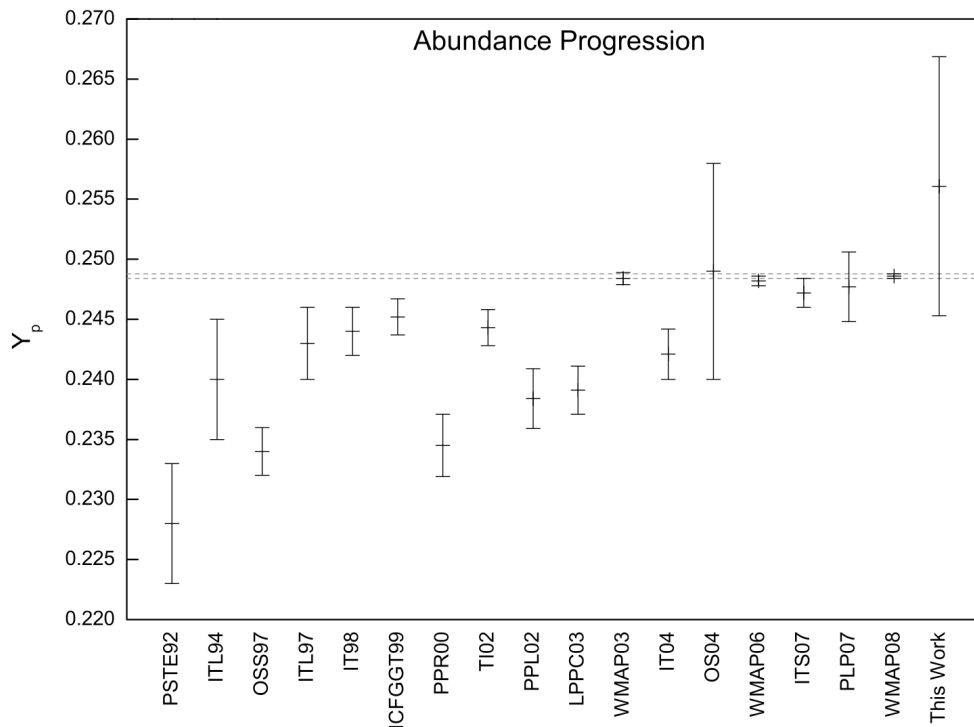
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New atomic physics, new observations  
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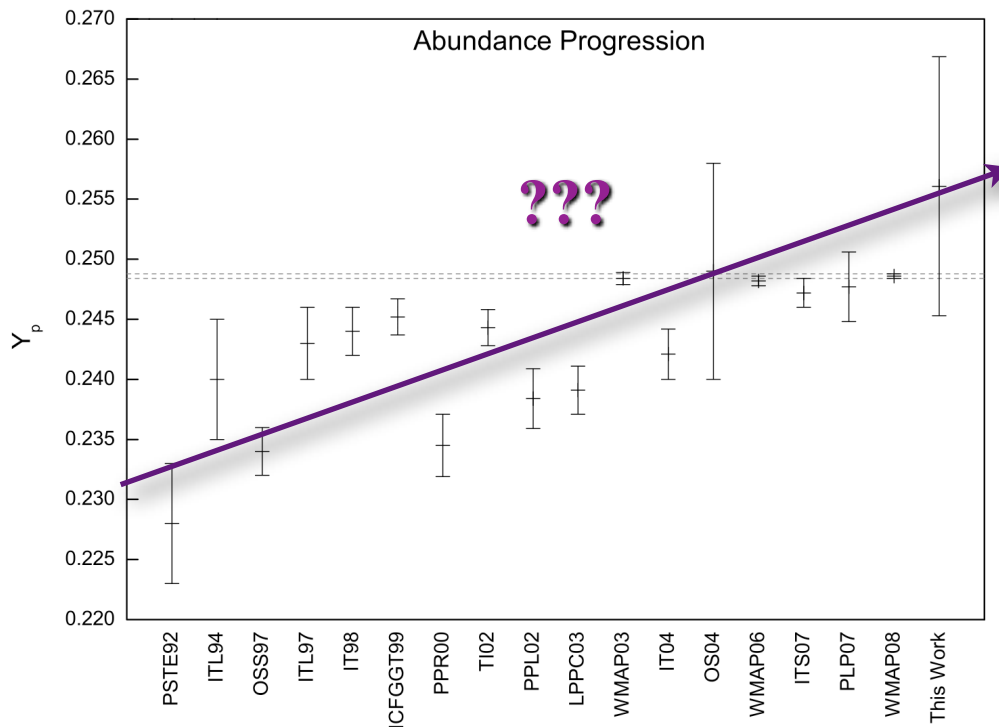
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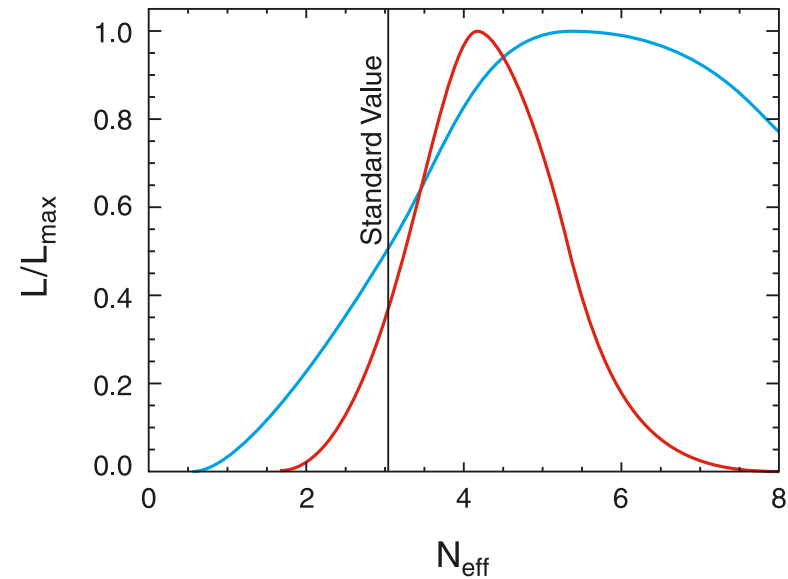
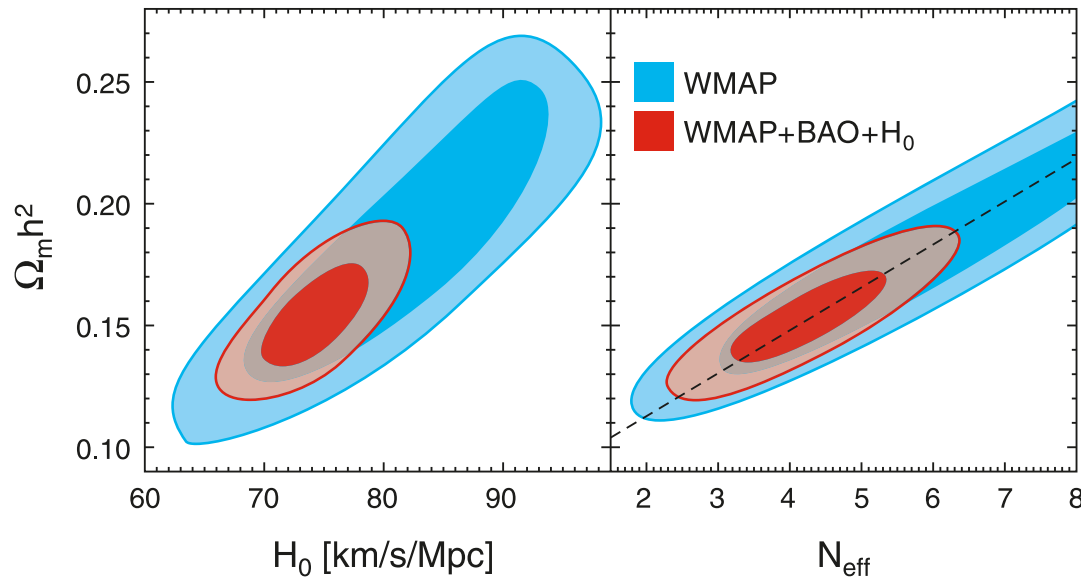
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Despite this... some bold statements!

*Izotov and Thuan, Astrophys. J. 710, L67 (2010)*  
"The primordial abundance of 4He: evidence for non-standard big bang nucleosynthesis."

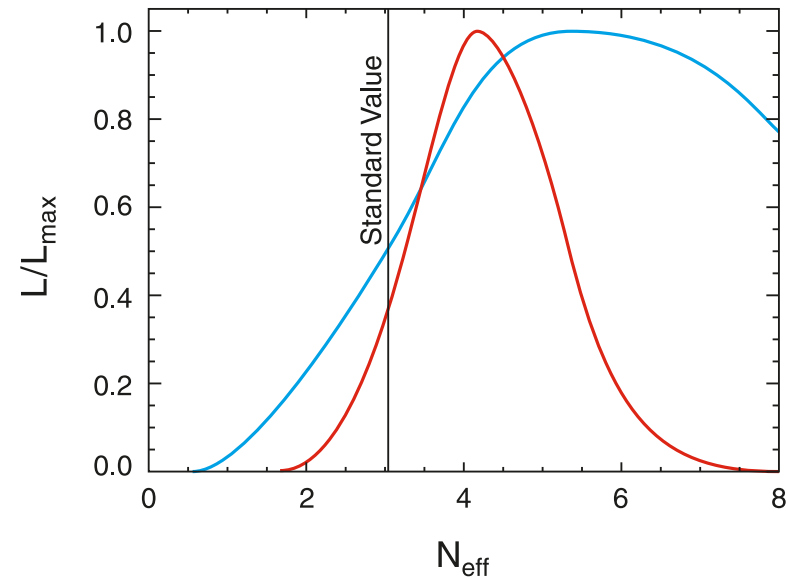
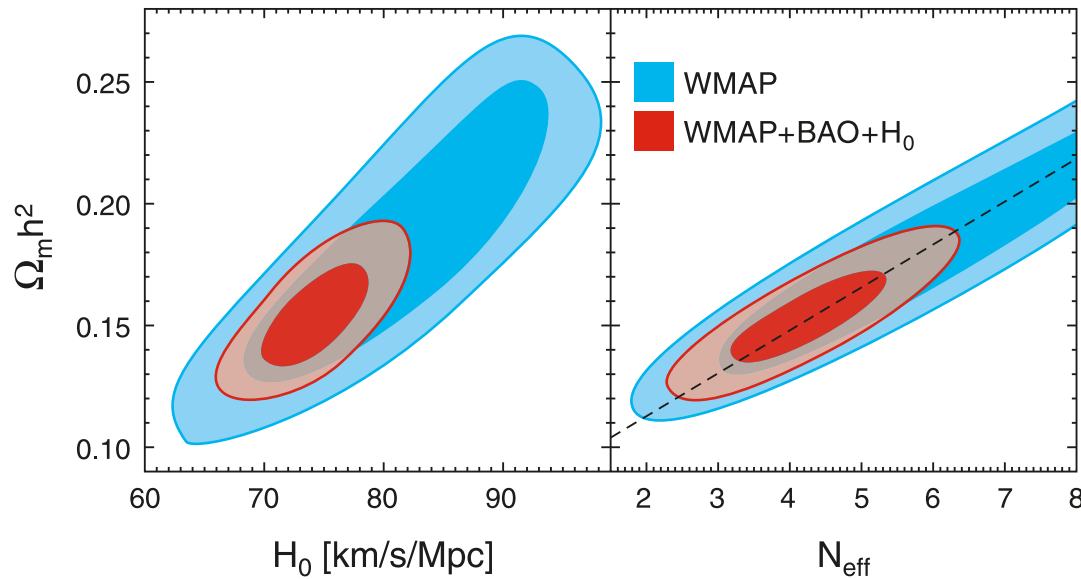
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WMAP [7-year], arXiv:1001.4538

Again consistent with expectations within  $< 2 \sigma$ 's (in cosmology!):  
you judge if one can really speak of "hints"!

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Again consistent with expectations within  $< 2 \sigma$ 's (in cosmology!):  
you judge if one can really speak of “hints”!

Note that, on the other end, there is not yet evidence for massive neutrinos power spectrum suppression in LSS

# EOM for $\nu$ evolution in the Early Universe

Effect of expanding universe

Refractive term due to  $e^-e^+$  energy density

Self-refraction

Collisions

$$i (\partial_t - H p \partial_p) \varrho_p = \left[ \left( \frac{U M_d^2 U^\dagger}{2p} - \frac{8\sqrt{2}G_F p}{3m_W^2} \mathcal{E} + \varrho - \bar{\varrho} \right), \varrho_p \right] + C[\varrho_p],$$

*additional  
refraction terms  
(e.g. FCNC)*

*barring  
cancellations, new  
collisional terms!*

mismatch between flavour and mass basis expressed via mixing matrix  $U$

Diagonal mass matrix (i.e. in mass basis)

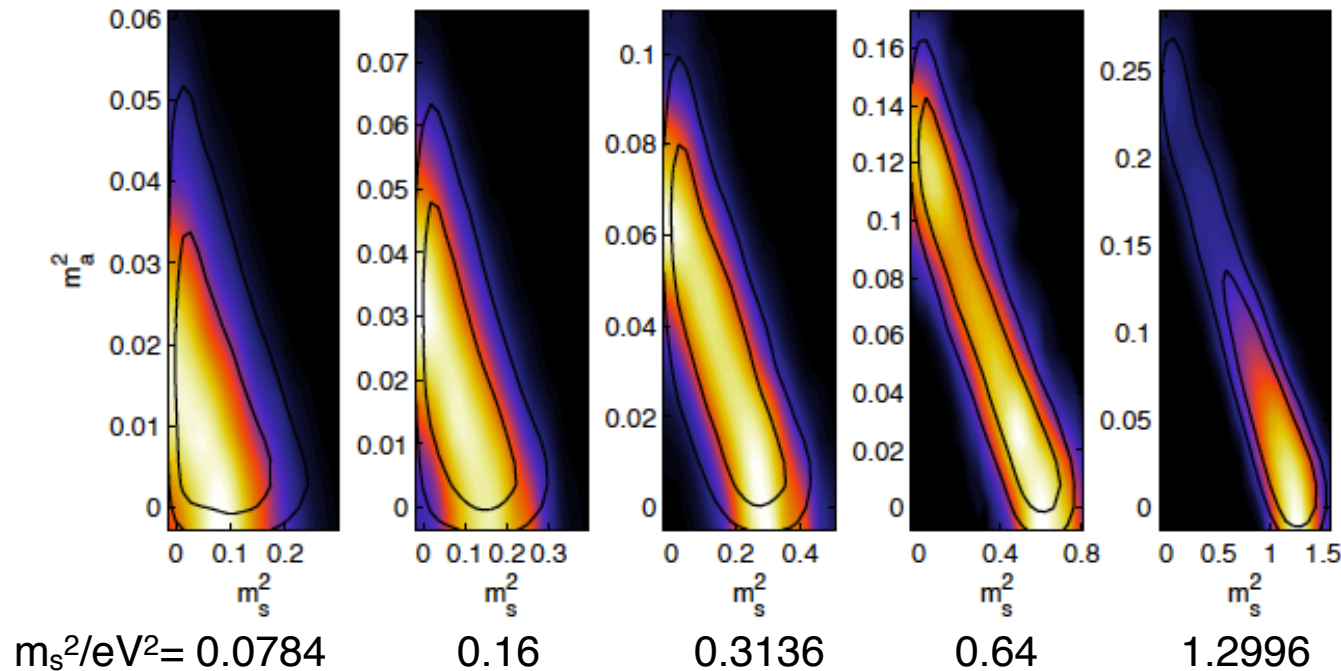
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$s_{lk} \equiv \text{Sin } \theta_{lk}, \quad c_{lk} \equiv \text{Cos } \theta_{lk}$$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}.$$

# “Direct” detection?

Sejersen Riis & Hannestad  
arXiv:1008:1495



$$|U_{es}|^2 = 0.18$$

An  $\sim eV$  scale “sterile” neutrino with a sufficiently large ( $O(0.1)$ ) mixing with the “effective”  $\nu_e$  may be detectable by the current experiment KATRIN as a distortion in the  $^3H$  endpoint spectrum

