



Recent advances in neutrino (astro)physics

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OUTLINE



Introduction : neutrino physics and neutrino propagation in matter

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Ongoing progress for neutrino propagation in massive stars (supernovae)

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Introduction : neutrino physics and neutrino propagation in matter

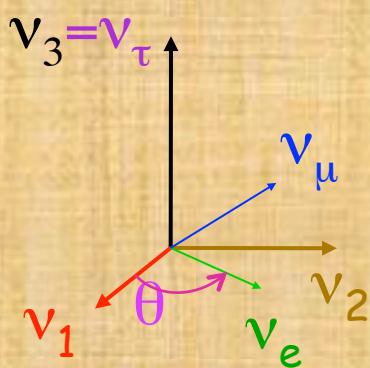


Ongoing progress for neutrino propagation in massive stars (supernovae)



CP violation effects in astrophysics and cosmology (BBN)

Neutrino Oscillations



$$\begin{pmatrix} v_e \\ v_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

flavour basis

mixing angle

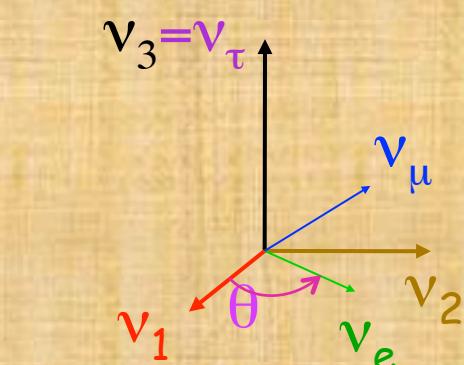
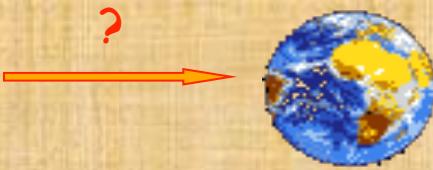
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mass basis

B. Pontecorvo, 1957

$$\Delta m^2 = m_2^2 - m_1^2$$

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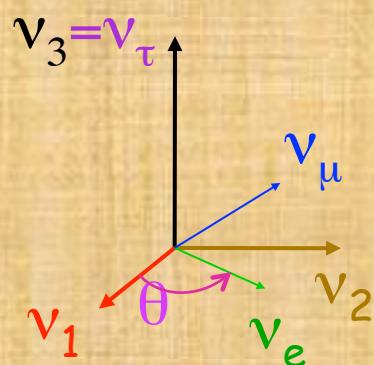
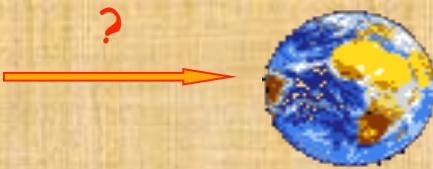


$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}_{\text{flavour basis}} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}_{\text{mixing angle}} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}_{\text{mass basis}}$$
$$\Delta m^2 = m_2^2 - m_1^2$$

- The time evolution is : $|\nu_e(t)\rangle = \cos \theta e^{-iE_1 t} |\nu_1\rangle + \sin \theta e^{-iE_2 t} |\nu_2\rangle$

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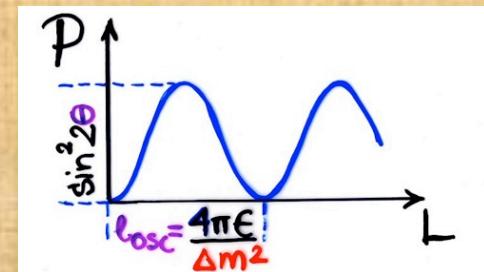
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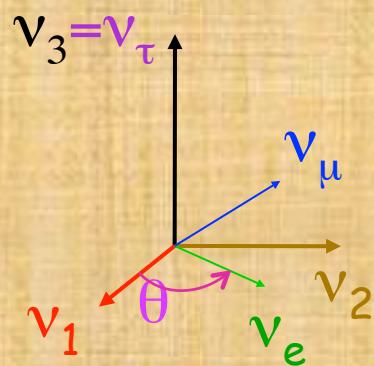
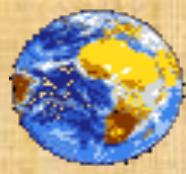
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probability for neutrino oscillations



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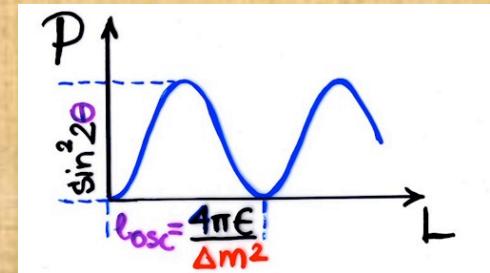
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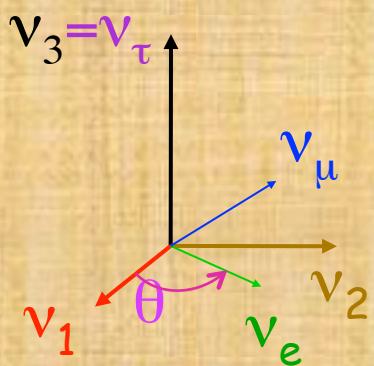
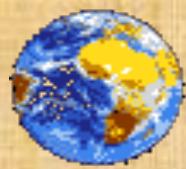
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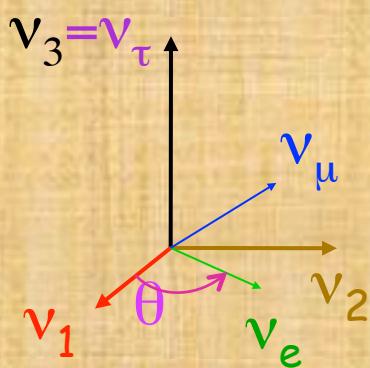
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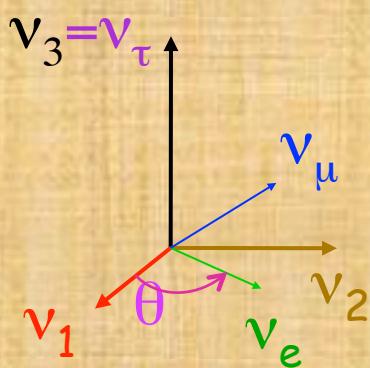
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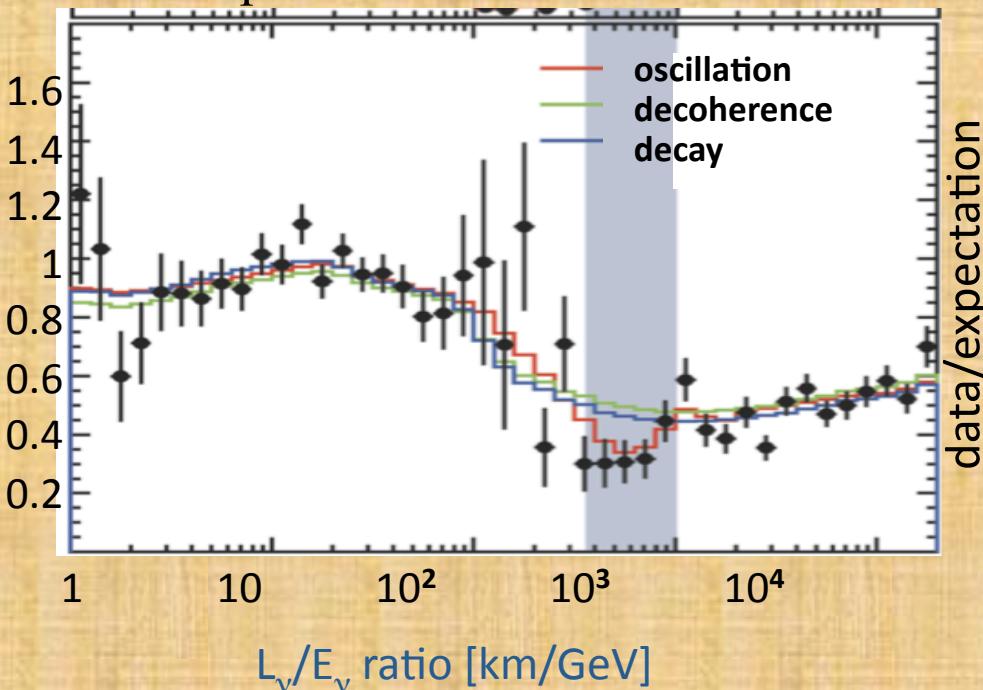
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Super-Kamiokande results .



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- The mixing matrix depends upon :
3 mixing angles, 1 CP Dirac phase, 2 Majorana phases

$$\theta_{12} \quad \sin^2 2\theta_{12} = 0.86^{+0.03}_{-0.04} \quad (\text{SNO, Kamland})$$

$$\theta_{23} \quad \sin^2 2\theta_{23} > 0.92 \quad (\text{Super-Kamiokande, K2K, MINOS, ... OPERA})$$

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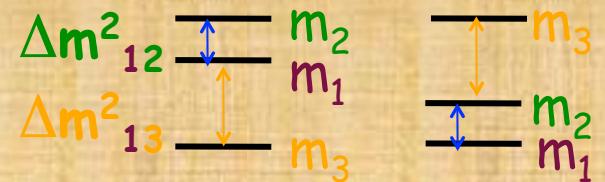
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- The mass ordering (hierarchy)?
Mass scale? \rightarrow

$$\Delta m_{32}^2 = m_3^2 - m_2^2 = 1.9 \text{ to } 3.0 \times 10^{-3} \text{ eV}^2$$

$$\Delta m_{21}^2 = m_2^2 - m_1^2 = 8.0^{+0.4}_{-0.3} \times 10^{-5} \text{ eV}^2$$



Inverted
 $(\Delta m_{13}^2 < 0)$

Normal
 $(\Delta m_{13}^2 > 0)$

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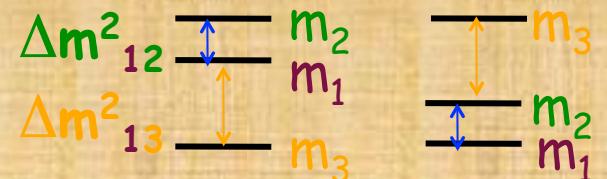
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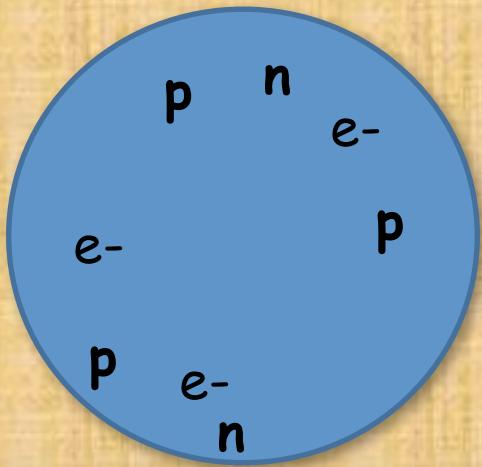
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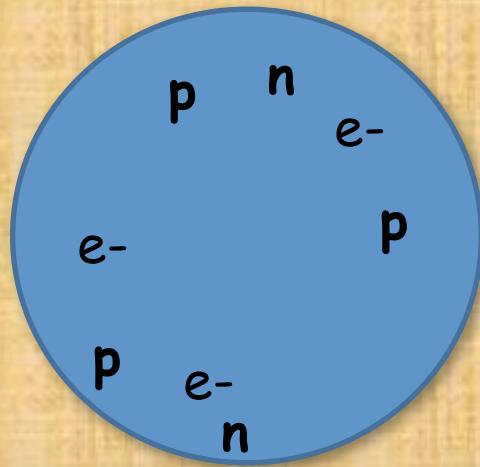
Important experimental results in the near future.



Neutrino propagation in media

THE PHYSICAL CONTEXT :

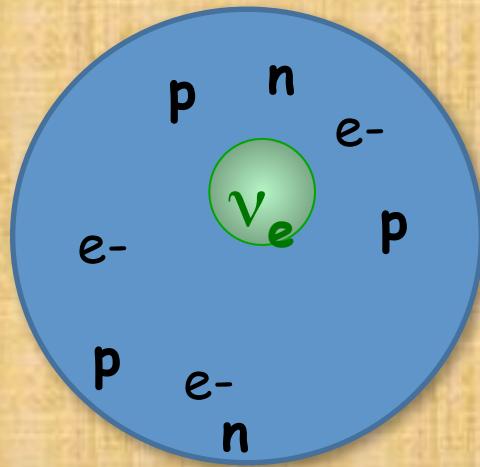
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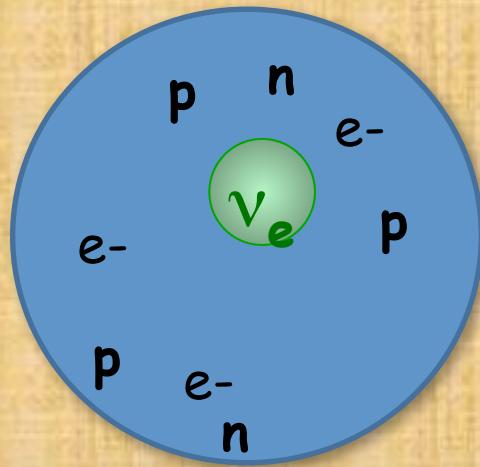
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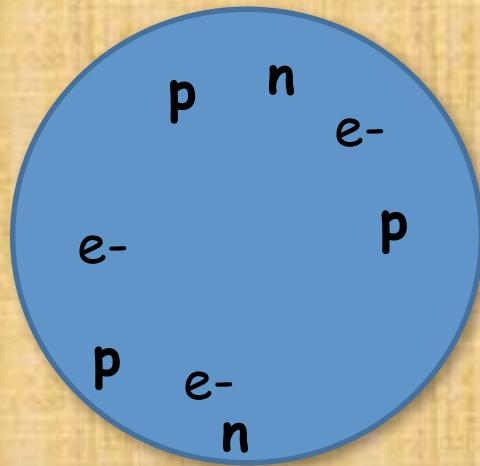
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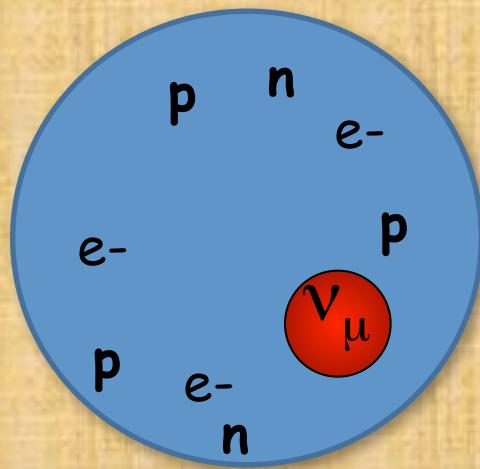
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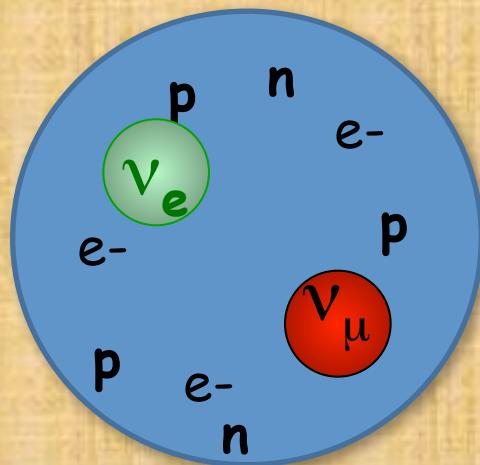
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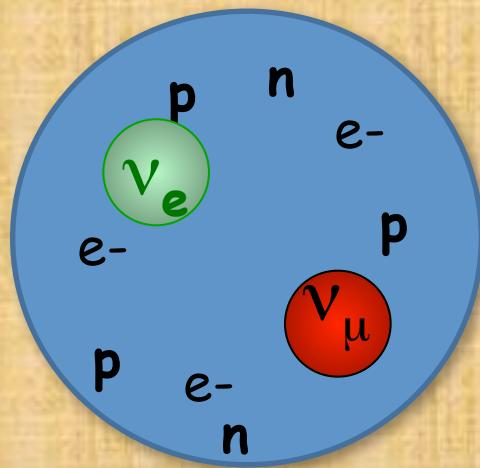
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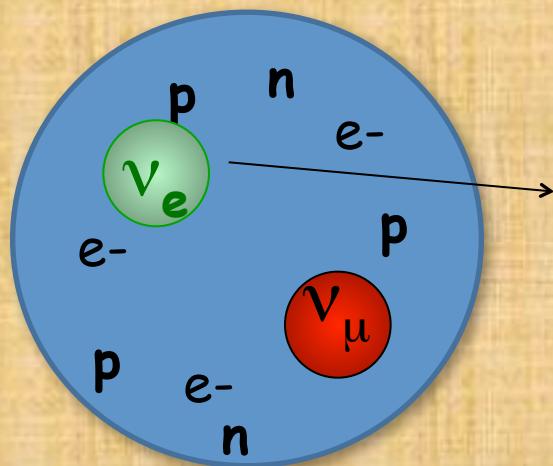
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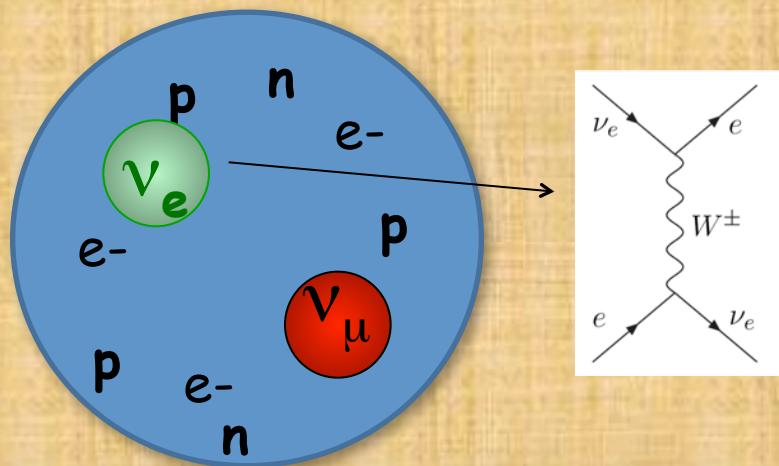
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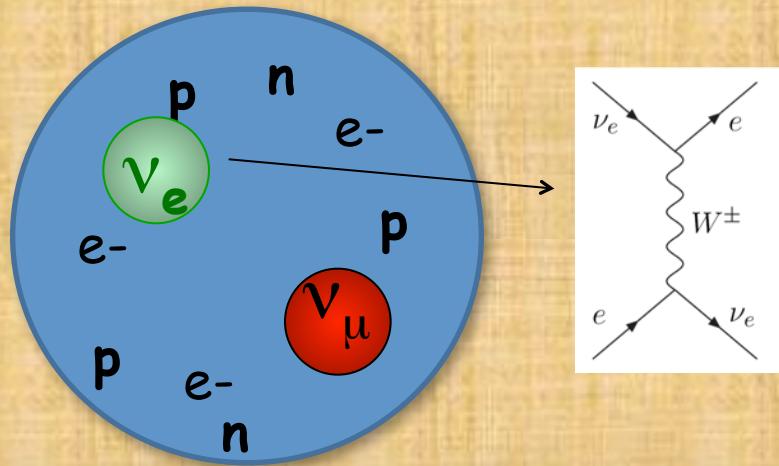
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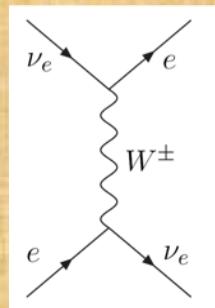
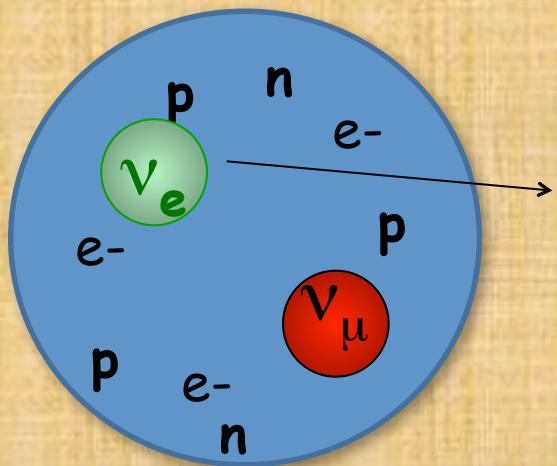
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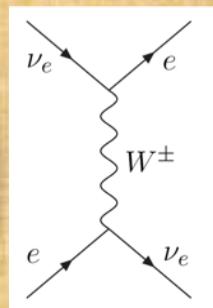
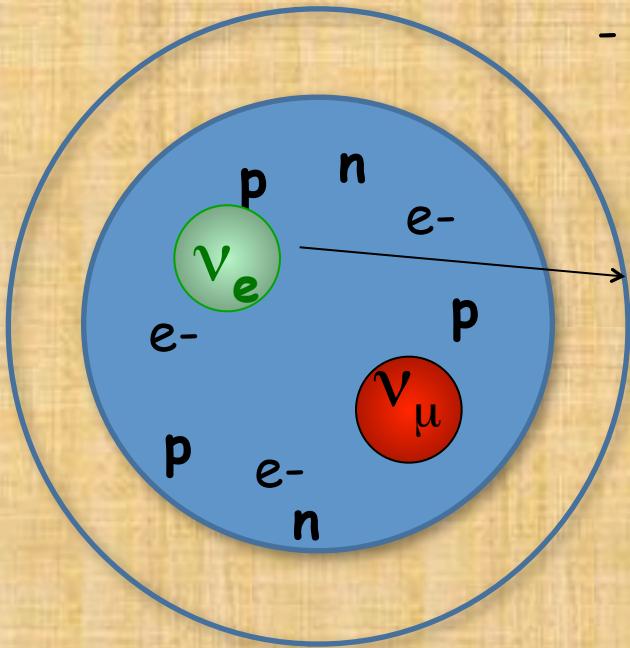
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Neutrino propagation in media

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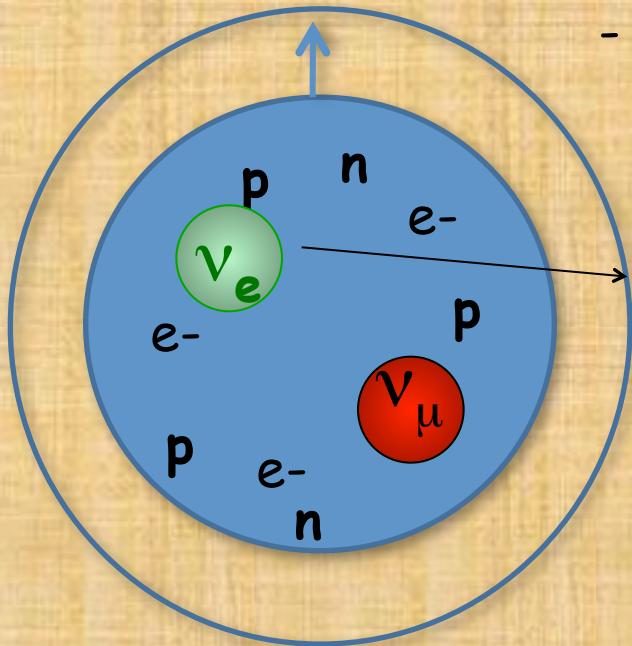
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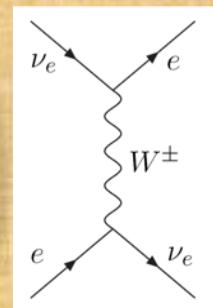
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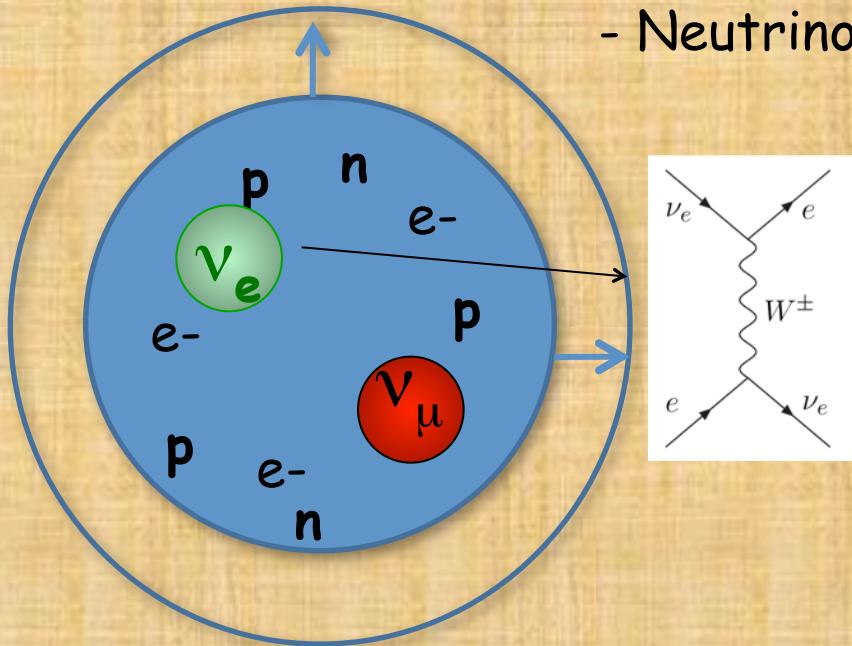
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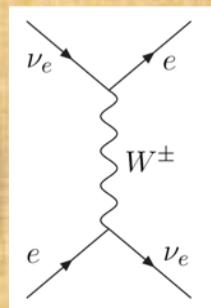
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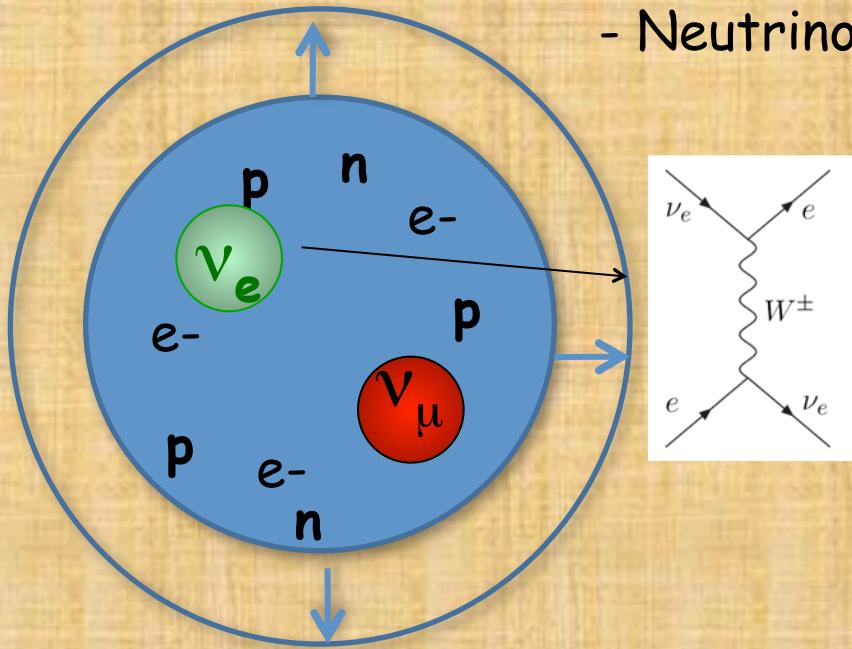
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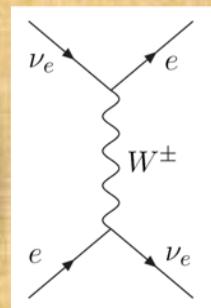
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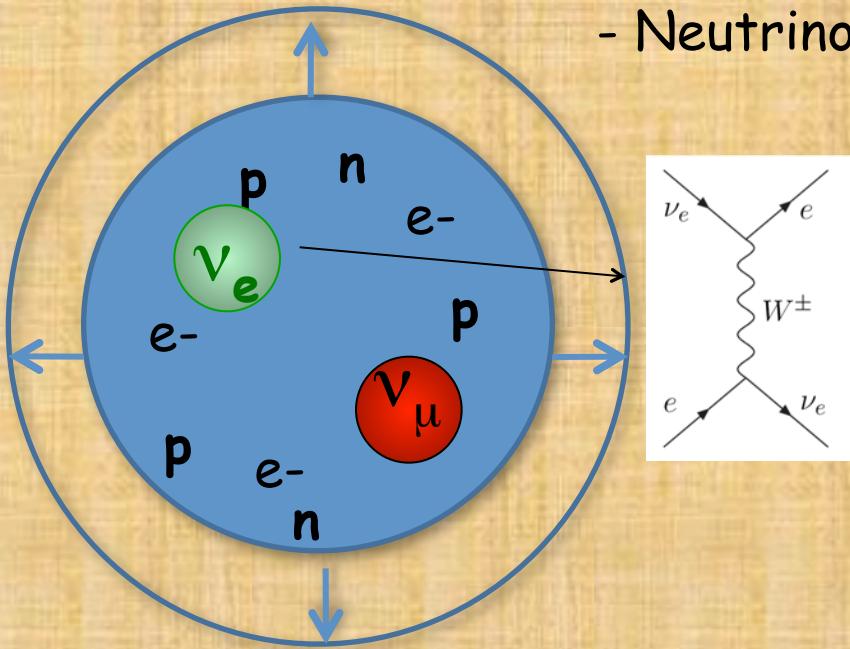
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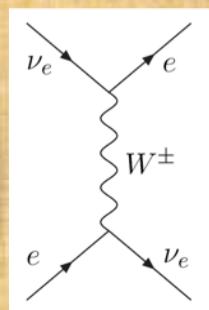
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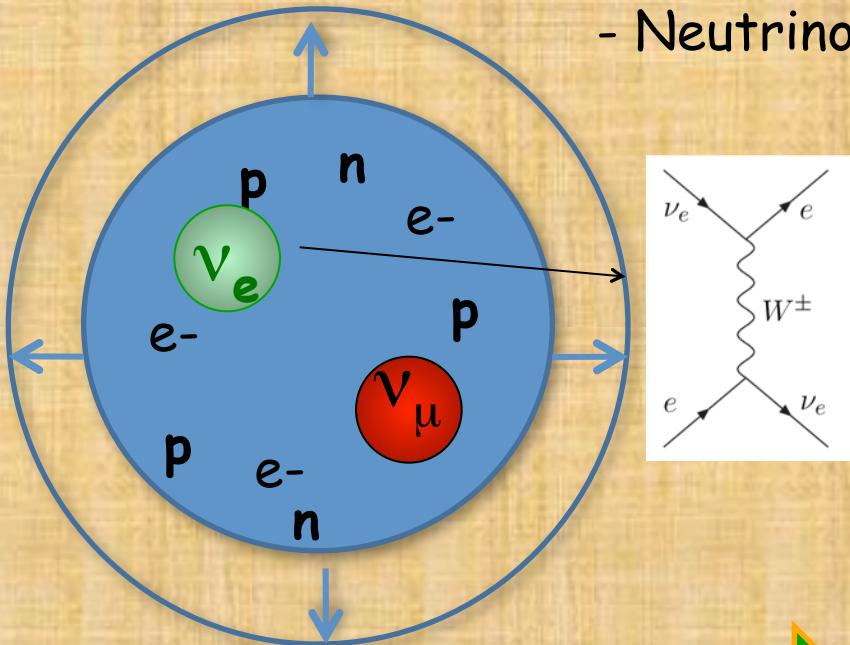
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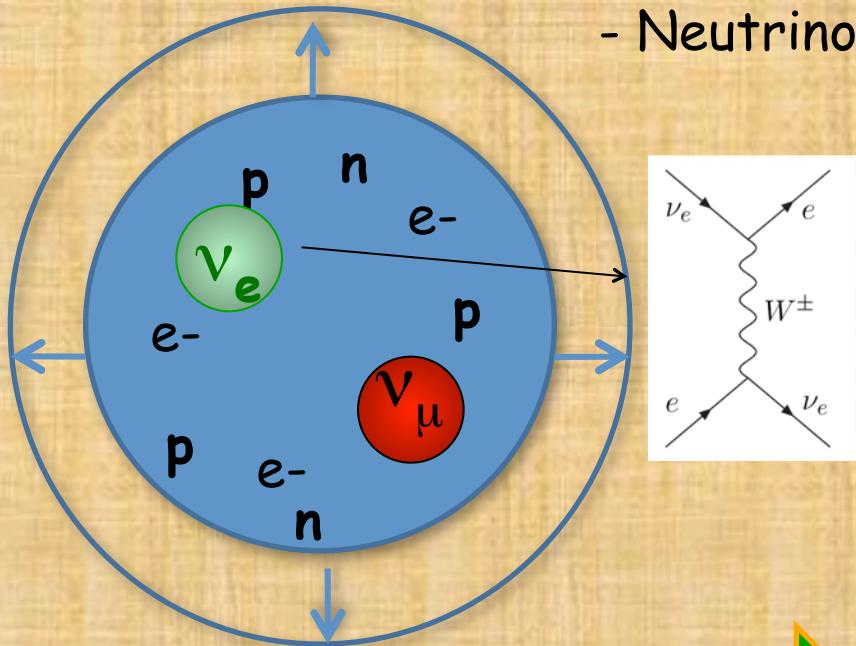
Neutrino propagation in media



Neutrinos emitted from the Sun

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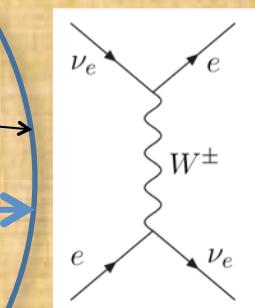
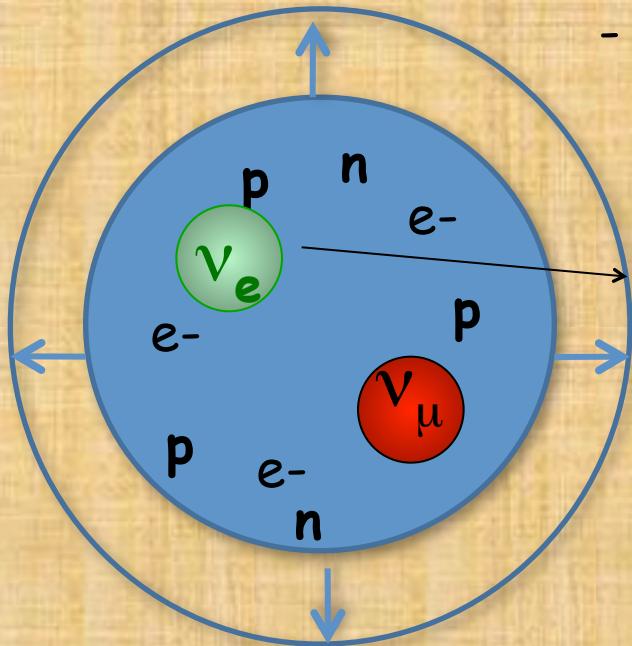
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Neutrino propagation in media

- Neutrinos emitted from the Sun
- Core-collapse supernovae and accretion disk-black holes

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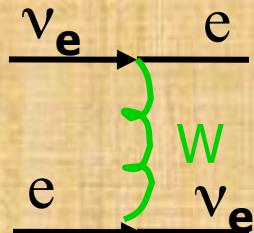


Core-collapse supernovae
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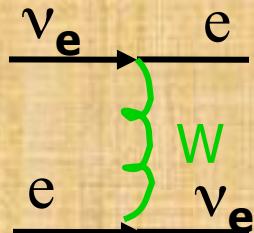
Neutrino evolution in the Early Universe
(Big-Bang Nucleosynthesis)

The Mikheev-Smirnov-Wolfenstein effect



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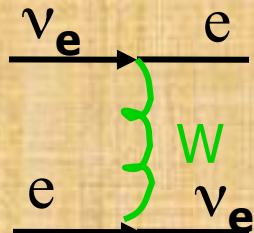


$$V(\rho) \sim G_F \rho_e.$$

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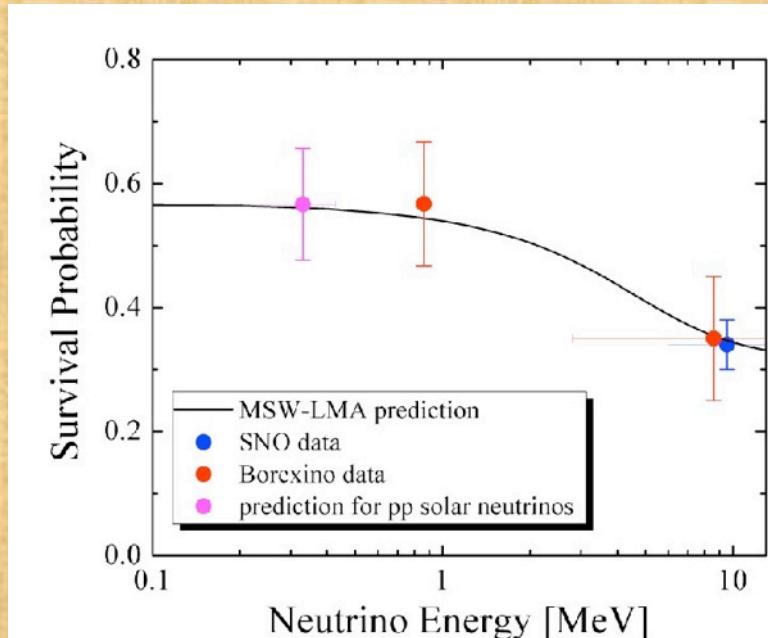
The flavour conversion depends on the adiabaticity of the propagation at the resonance (density profile and mixing parameters).

The Mikheev-Smirnov-Wolfenstein effect



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The Mikheev-Smirnov-Wolfenstein (MSW) effect ('78, '86) : neutrino coupling with matter induces a resonant flavour conversion.

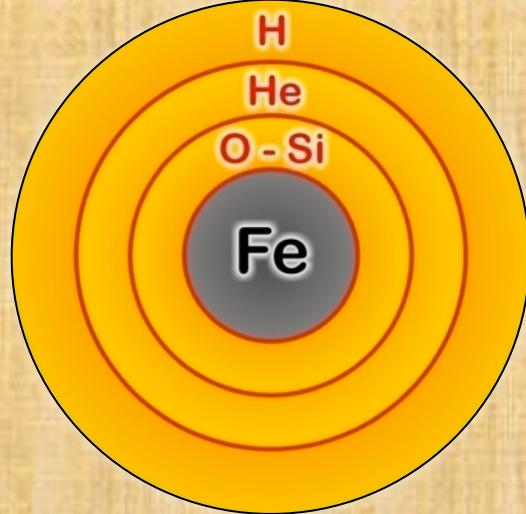


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the beautiful explanation of the « solar neutrino deficit » problem !

Borexino Collaboration, J. Conf. Ser. 202, 012028 (2010)

Core-collapse supernovae (SN)



99 % of the energy is emitted as neutrinos of all flavours in a short burst of about 10 s.

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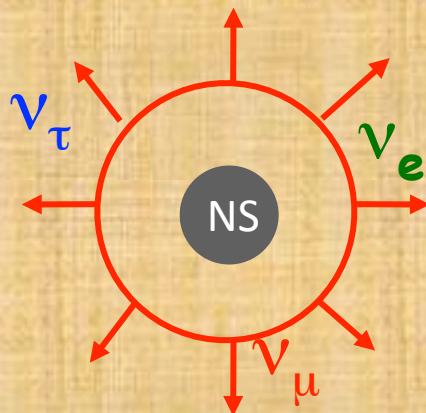
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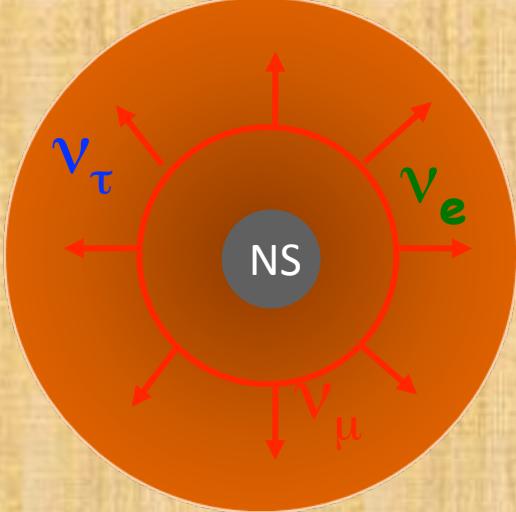
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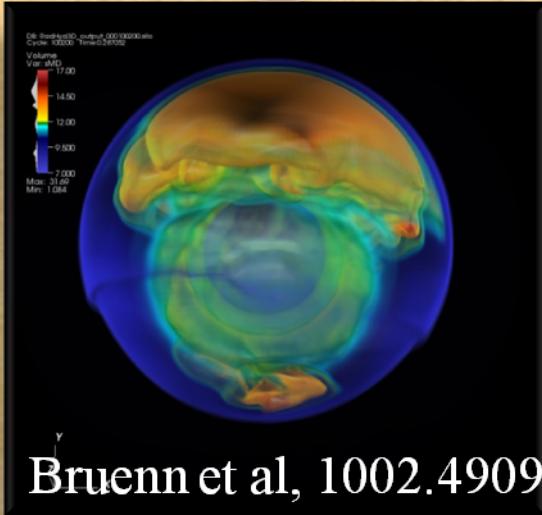
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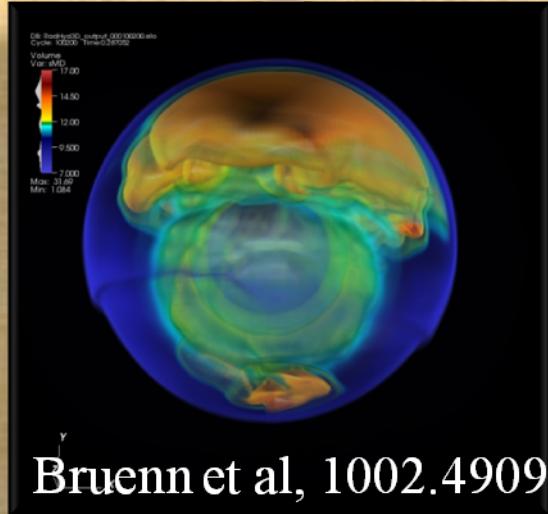
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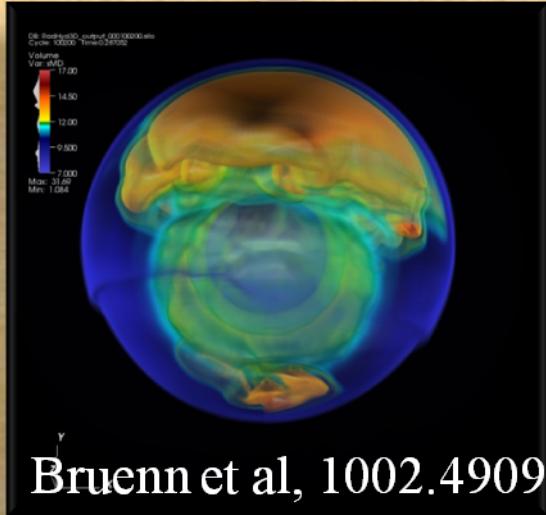
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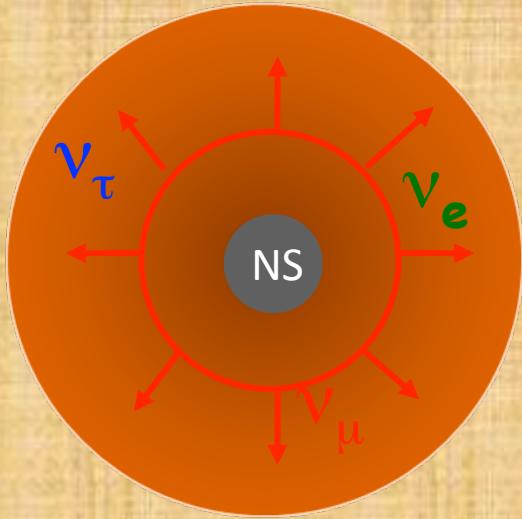
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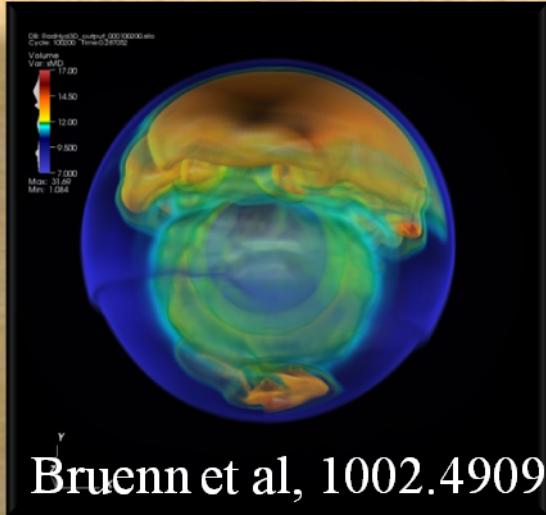


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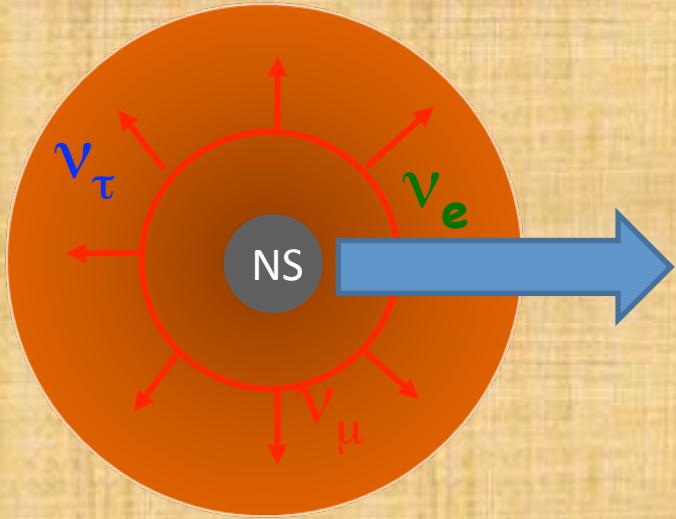


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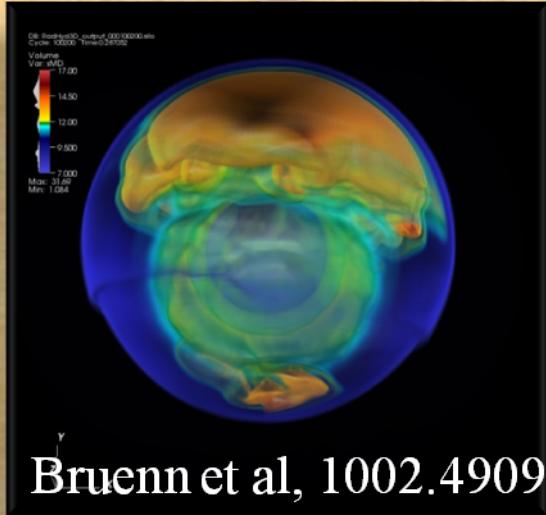


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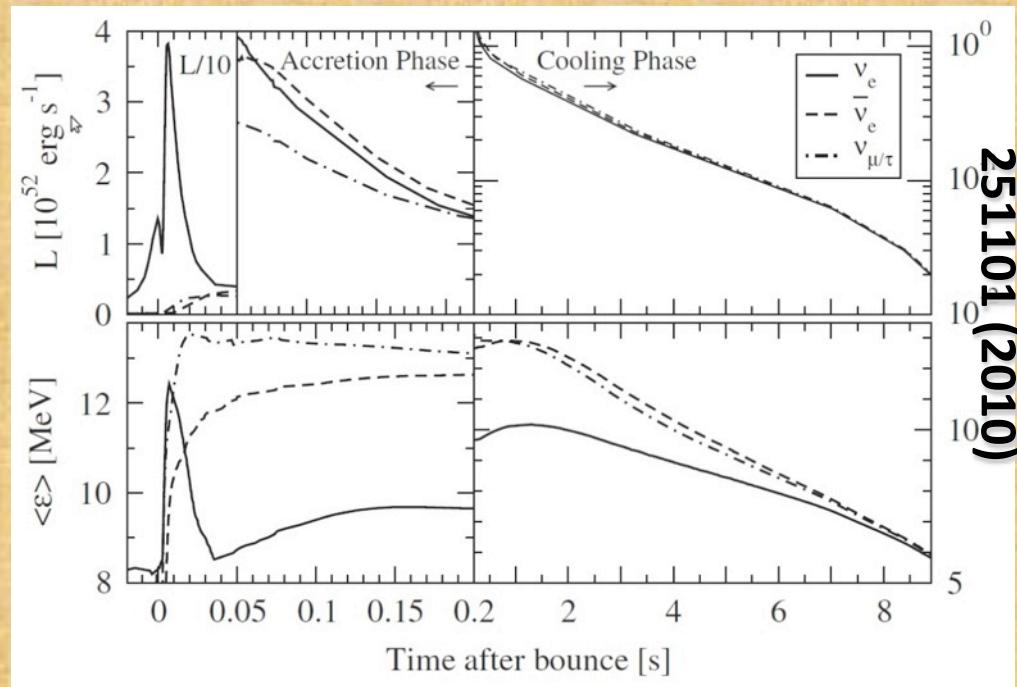
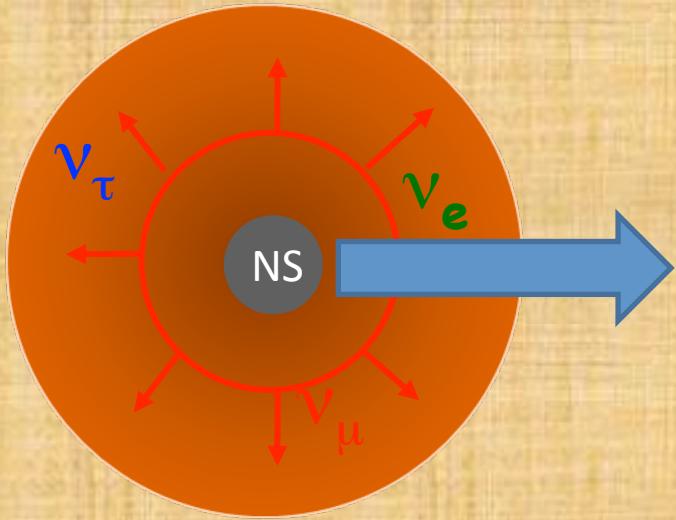


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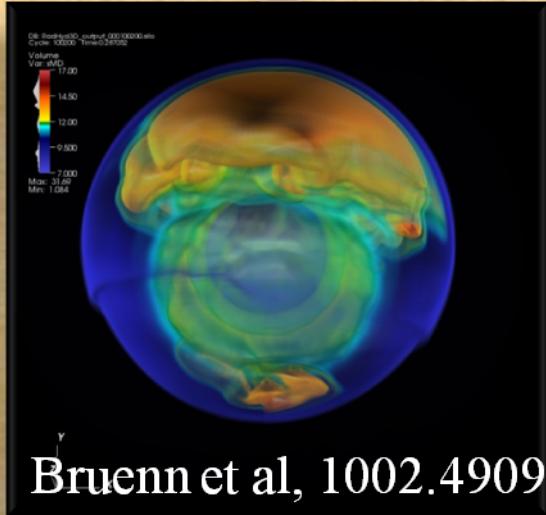
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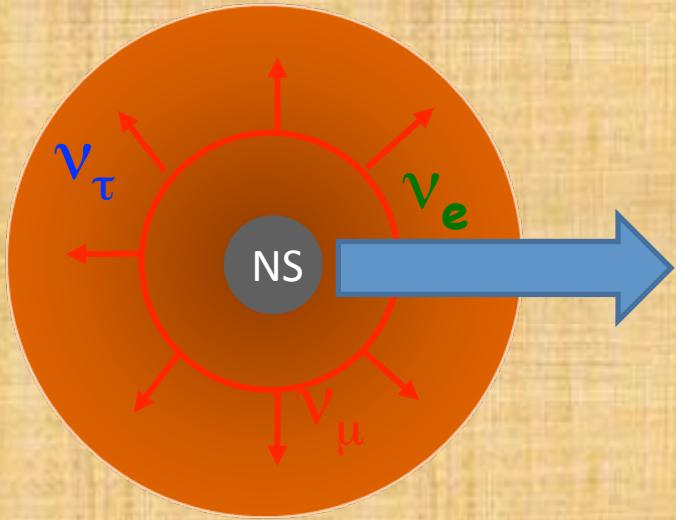
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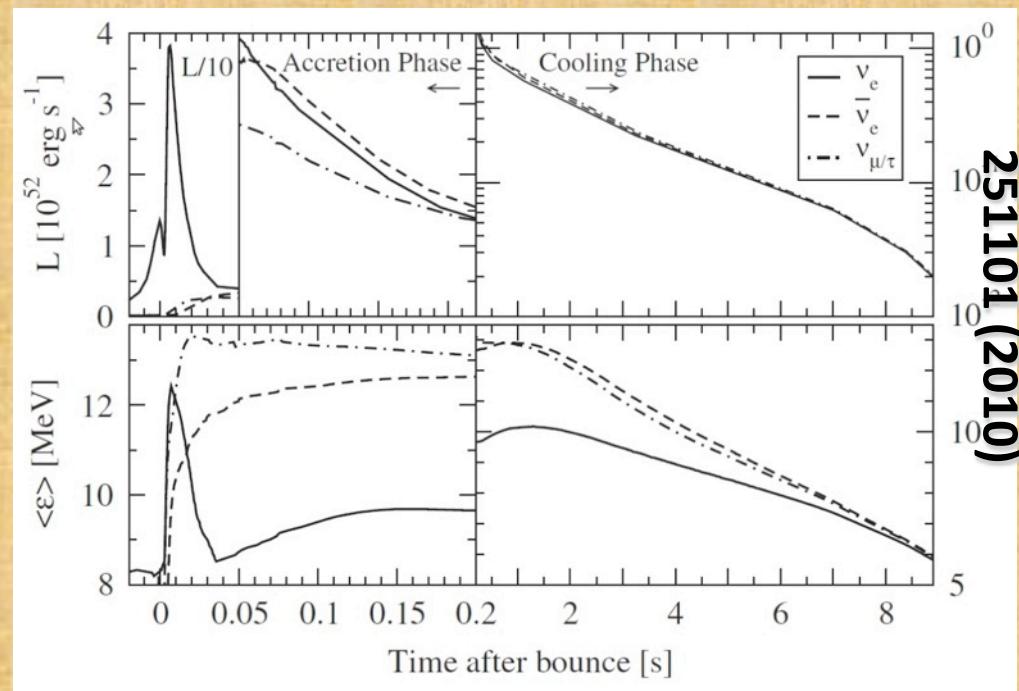


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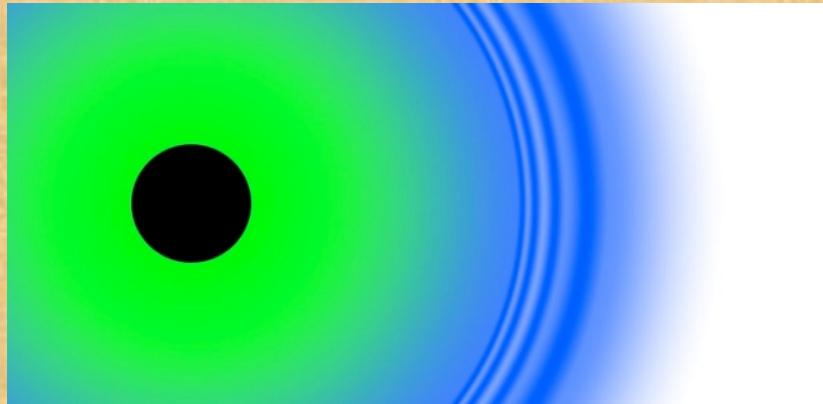


FUTURE OBSERVATIONS
IMPORTANT !

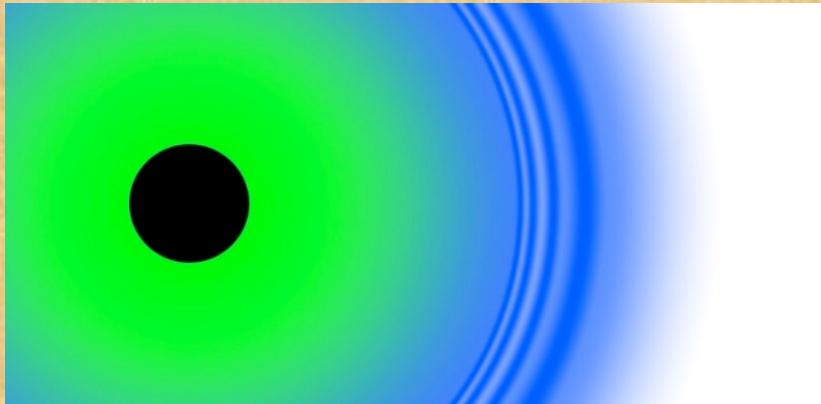


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Neutrino flavour conversion in SN

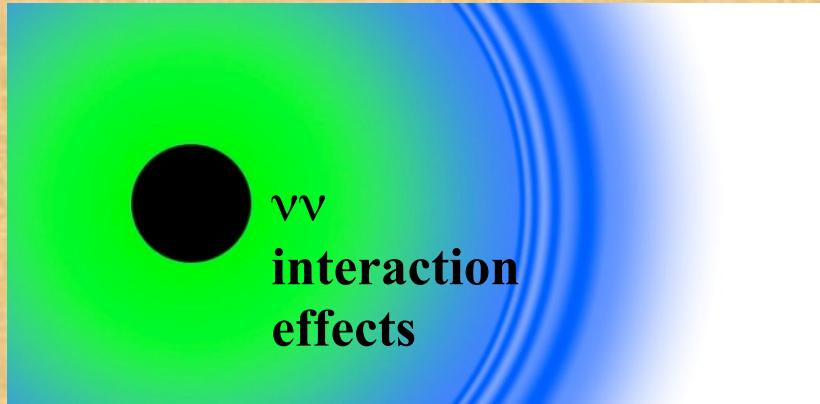


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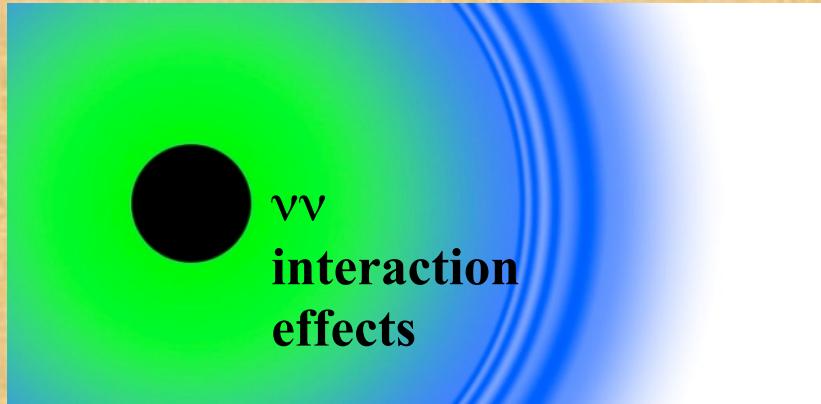
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Collective phenomena emerge.

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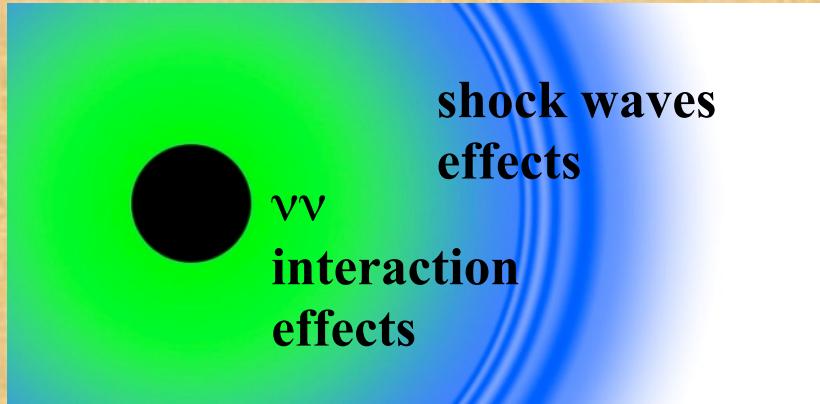
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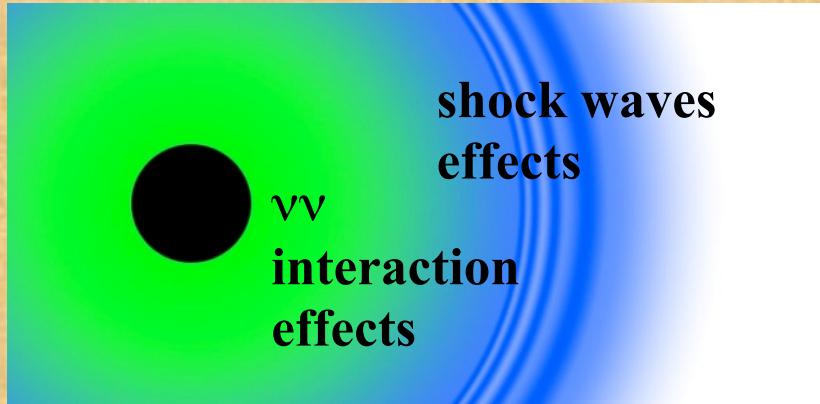
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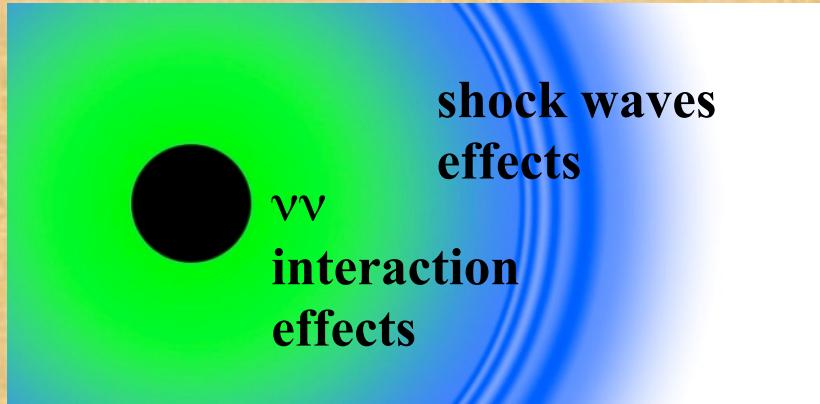
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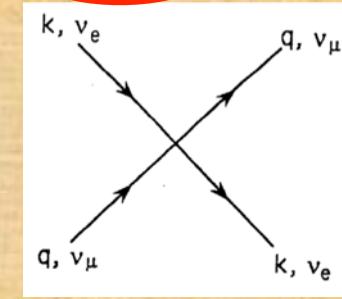
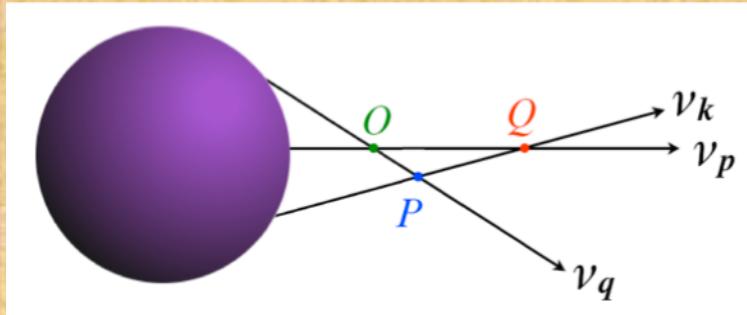
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The ν - ν interaction

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$$H = H_{\text{vacuum}} + H_{\text{matter}} + H_{\nu\nu}(\rho_\nu)$$

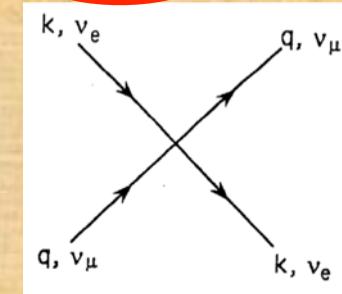
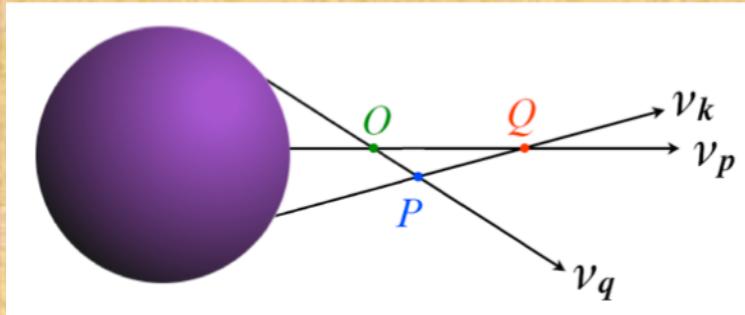


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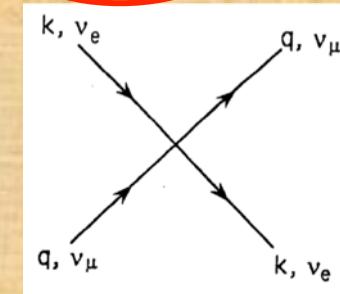
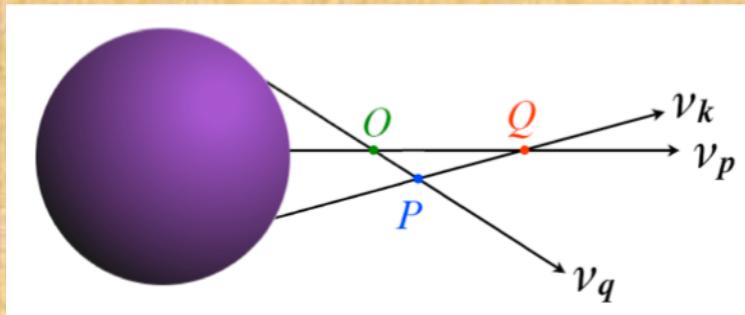
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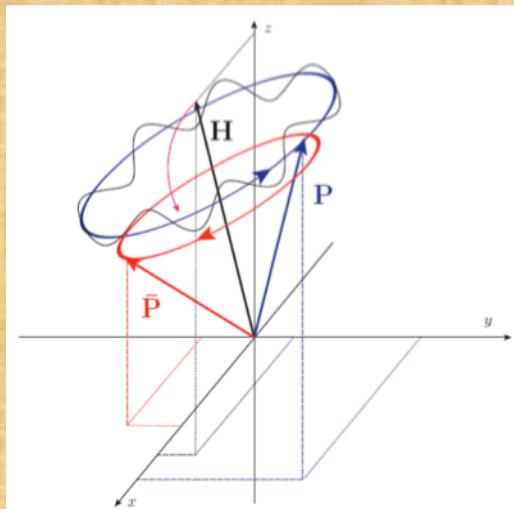
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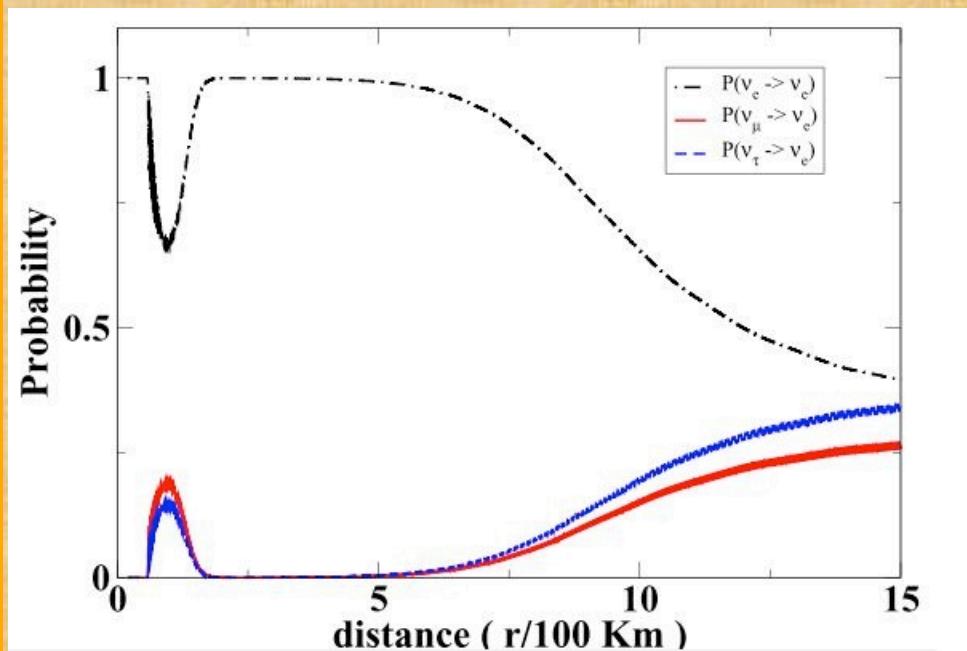


- Three regimes are identified :
 - The synchronization regime
 - The bipolar oscillations
 - The spectral split

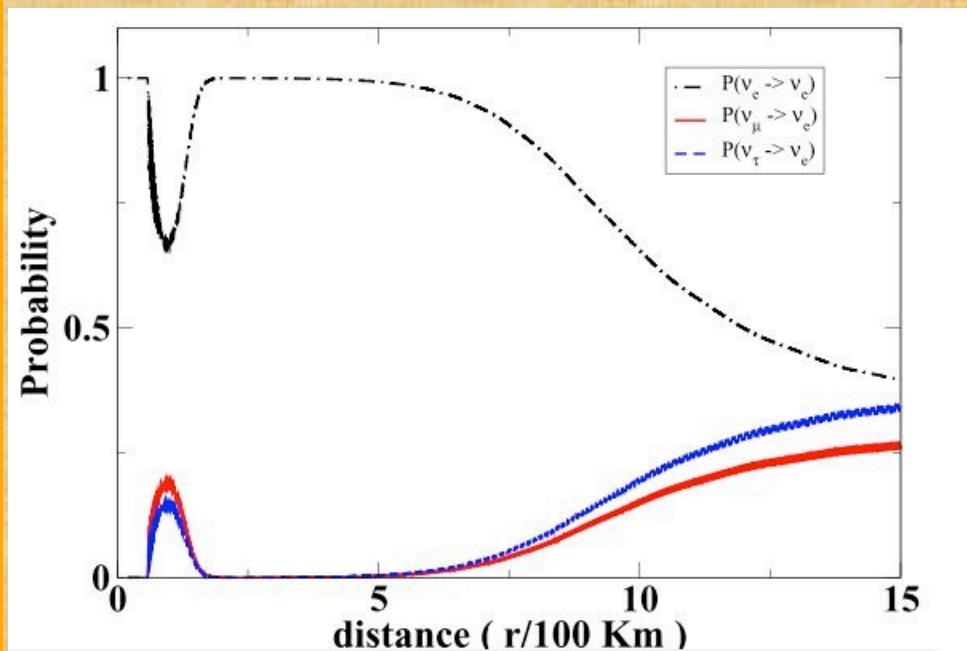
Duan, Fuller, Qian PRD74 (2006), PRL 97 (2006),
Hannestad, Raffelt, Sigl, Wong, PRD 74 (2006),
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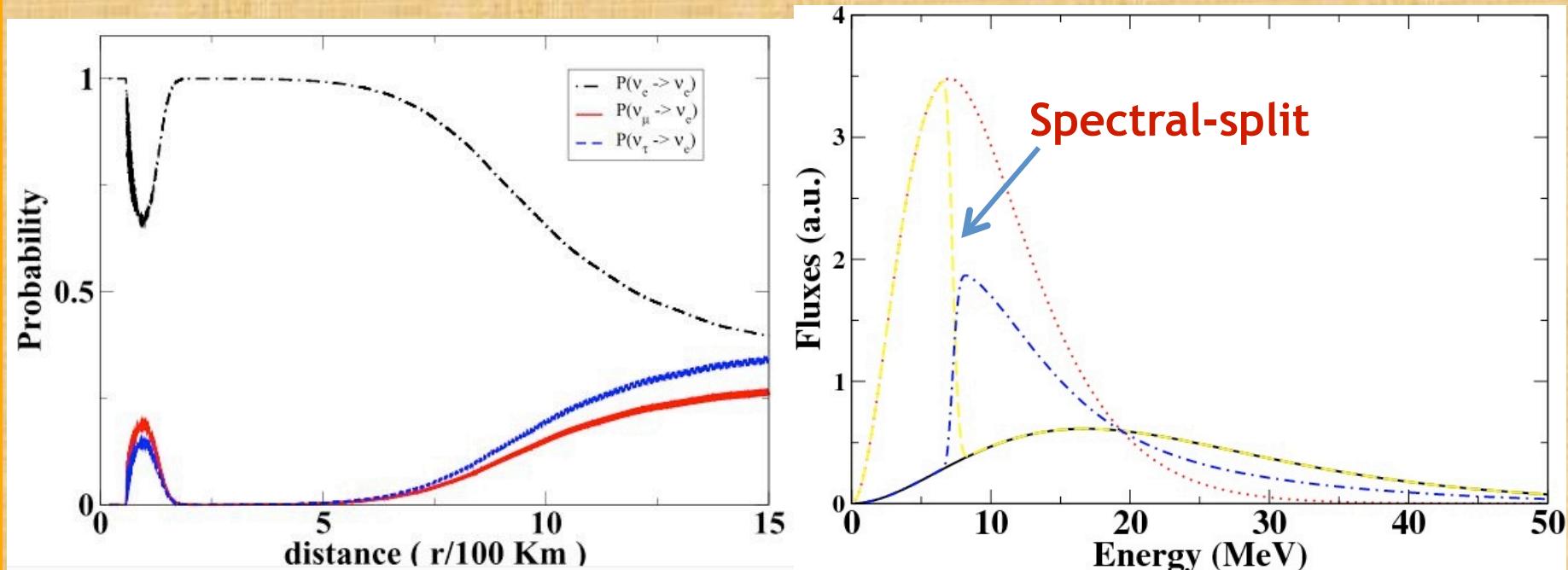


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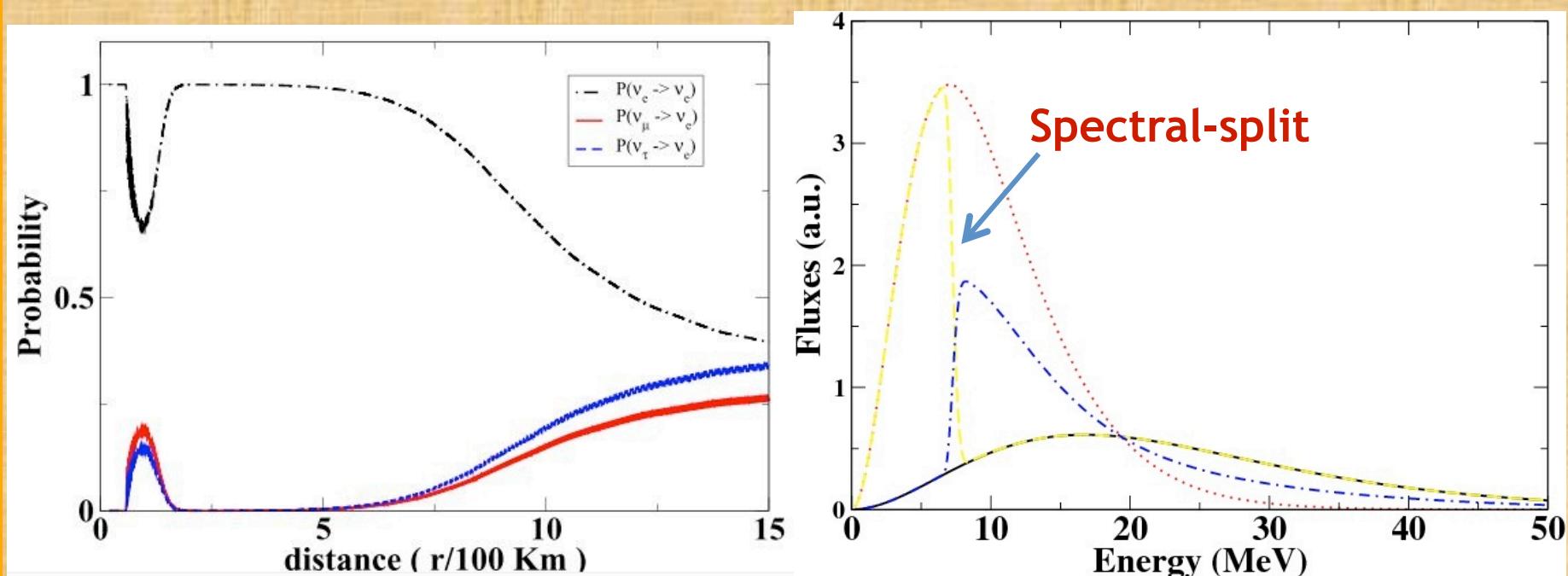
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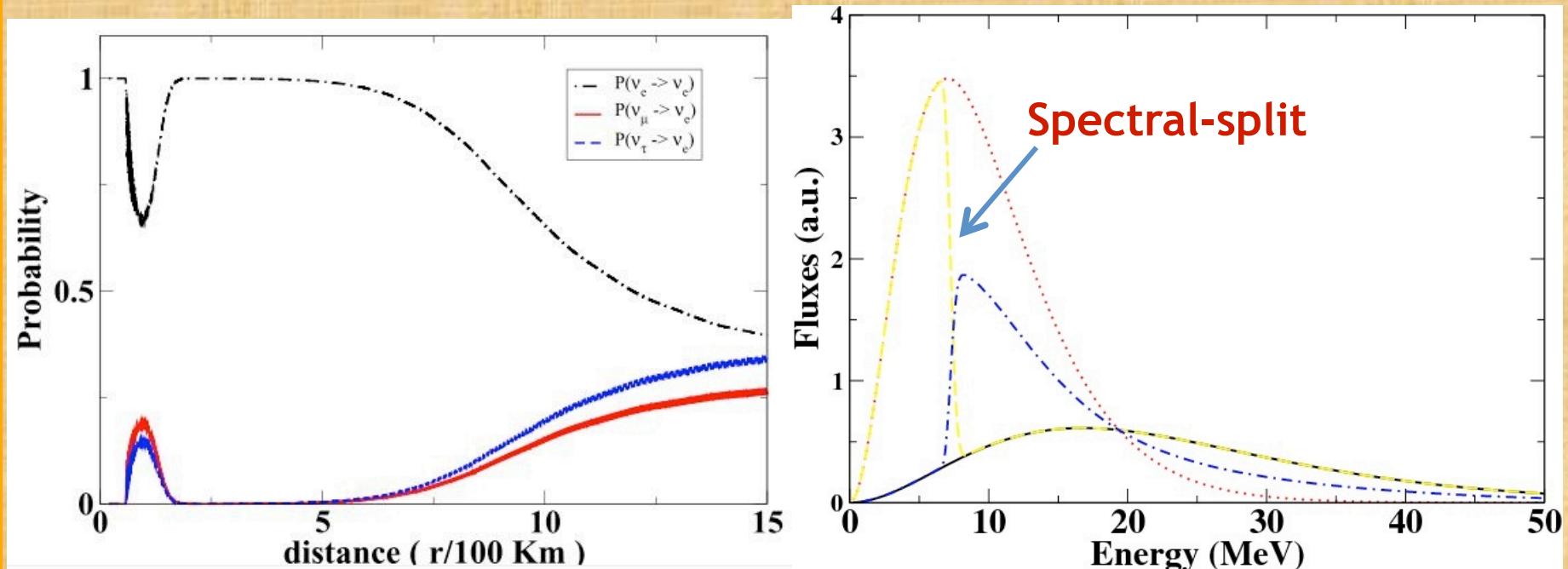
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We have discovered that the spectral split is similar
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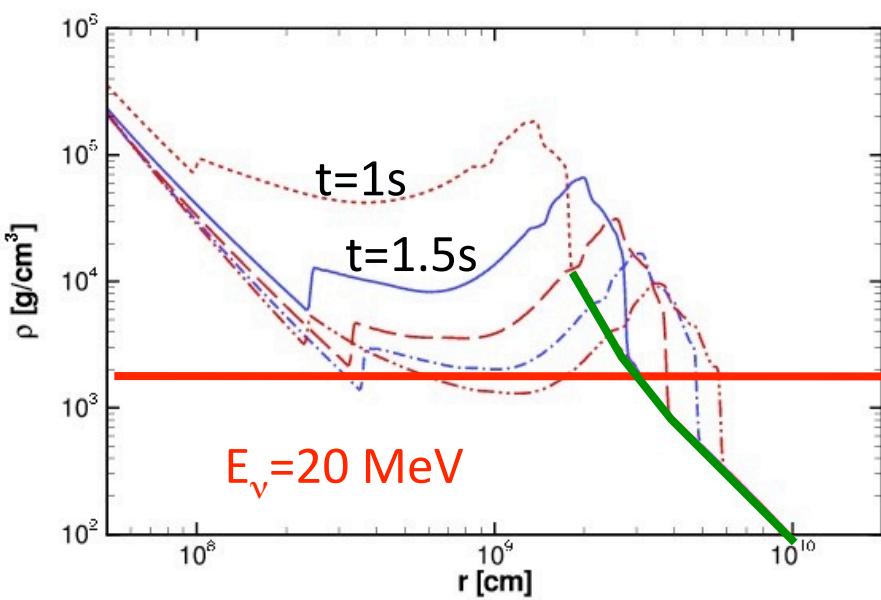
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The shock wave effects

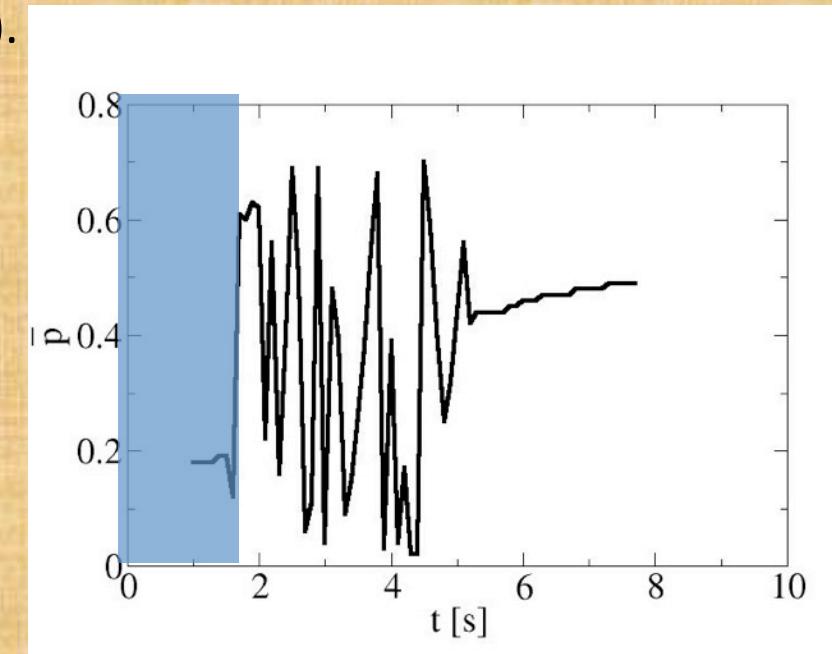
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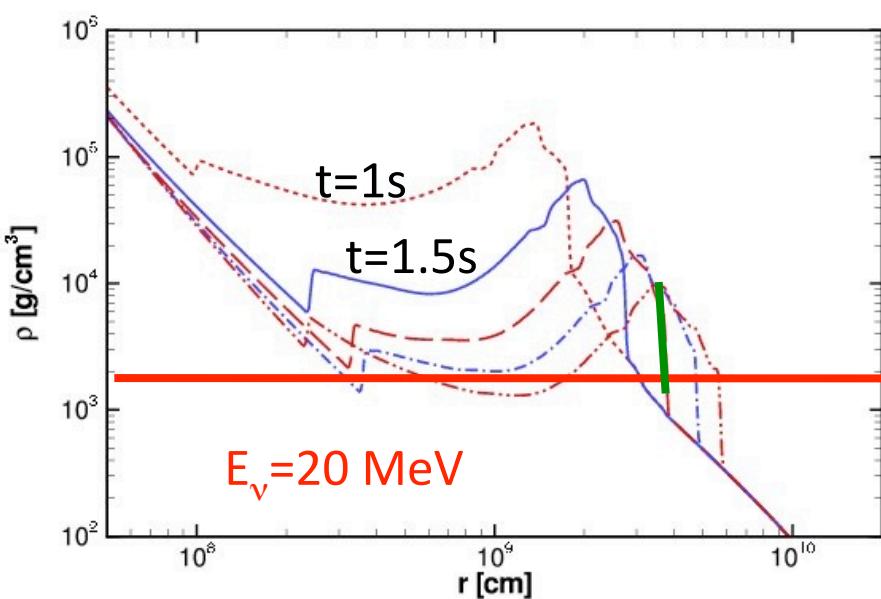
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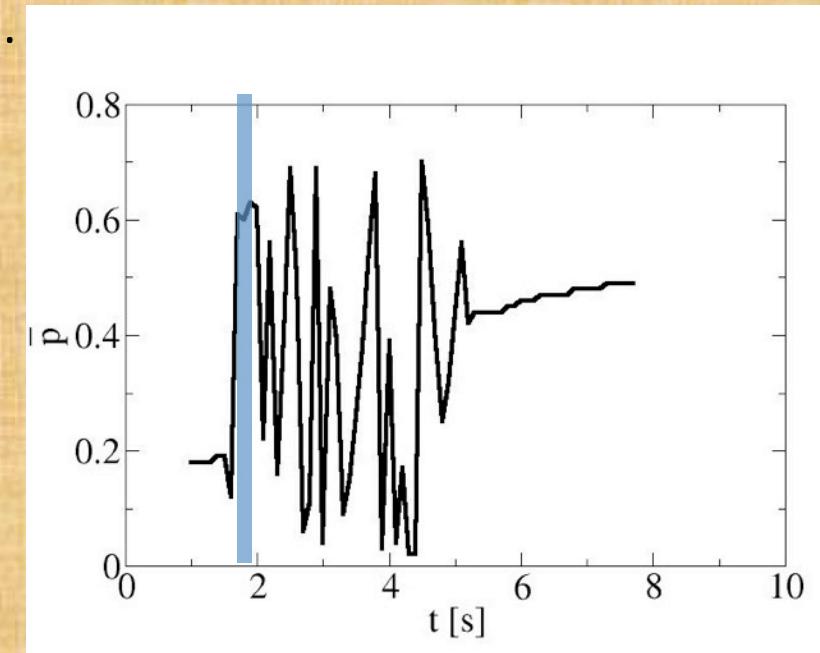
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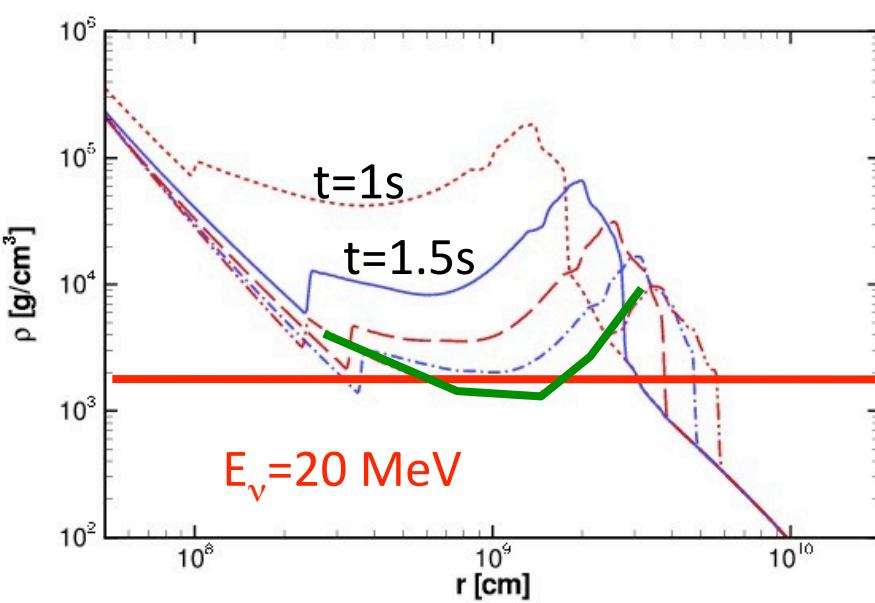
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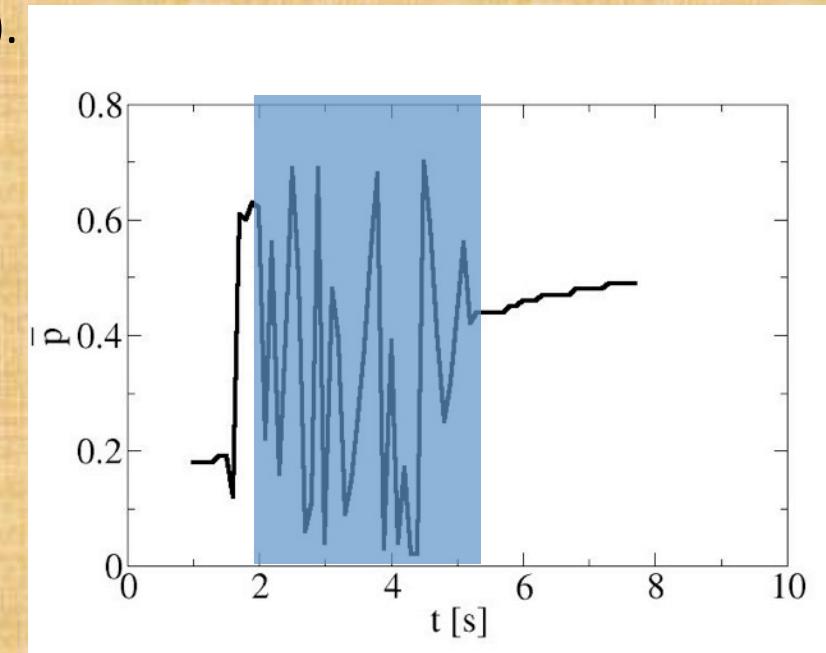
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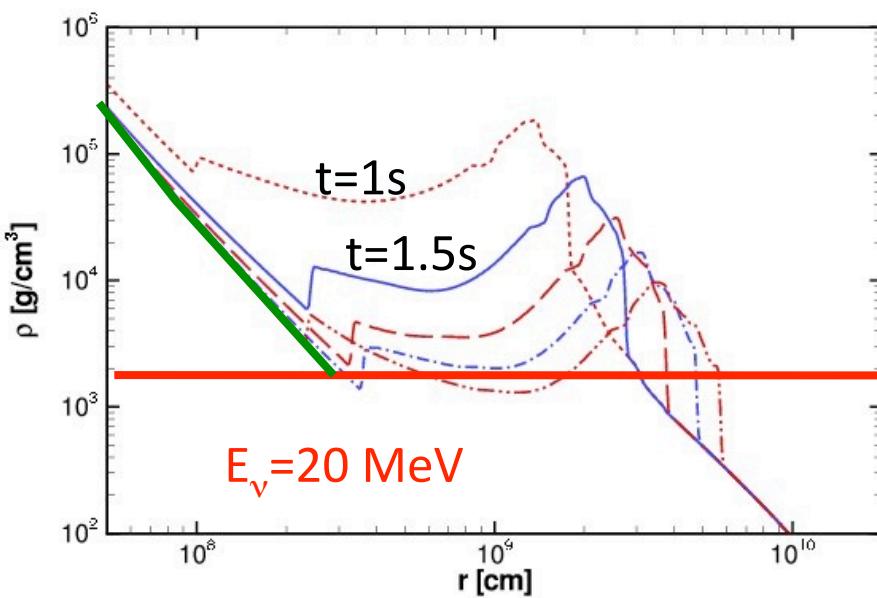
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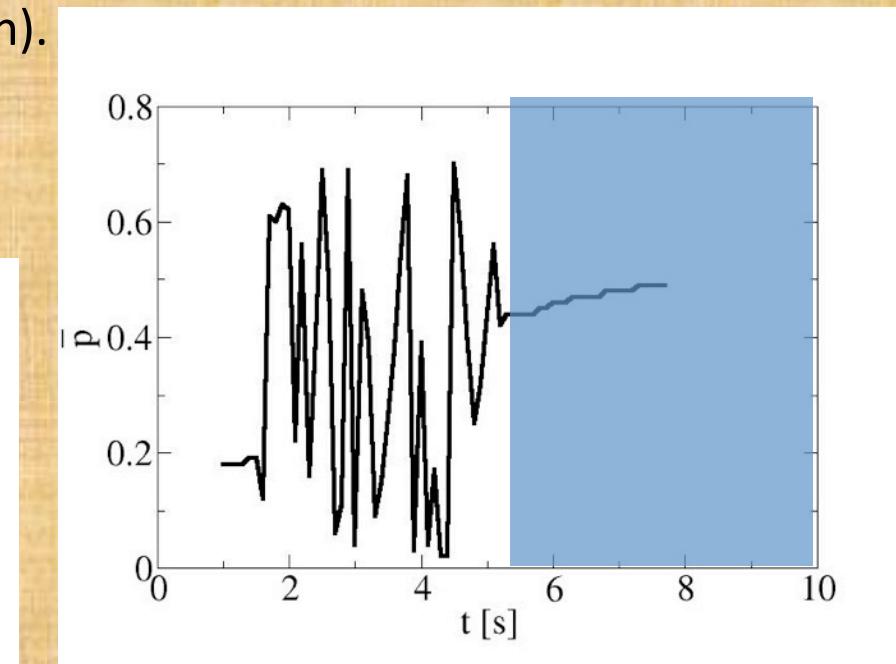
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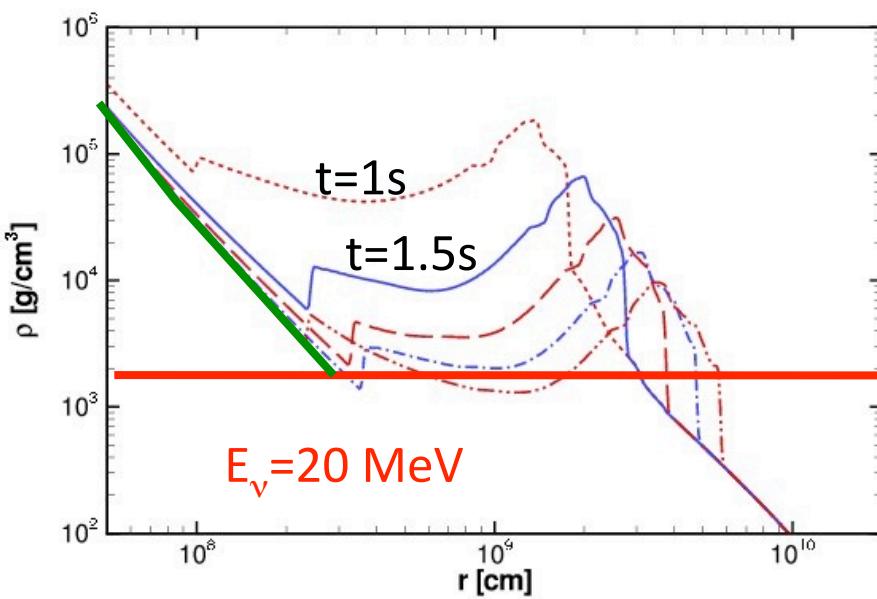
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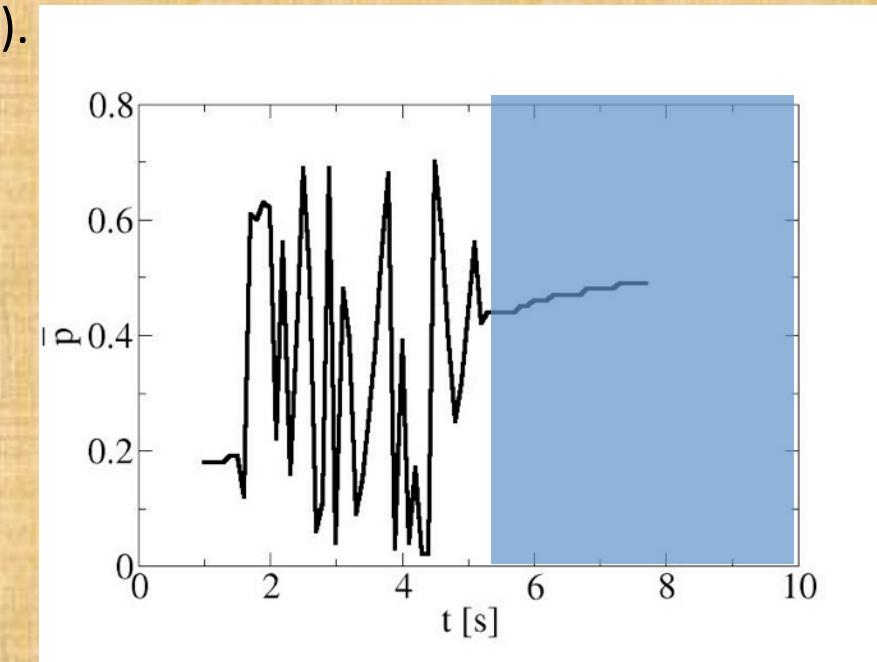
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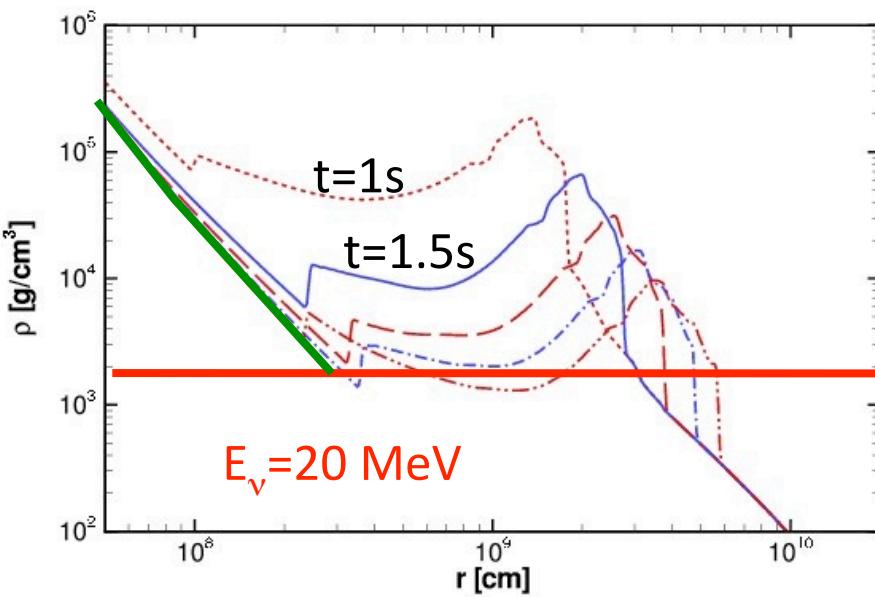
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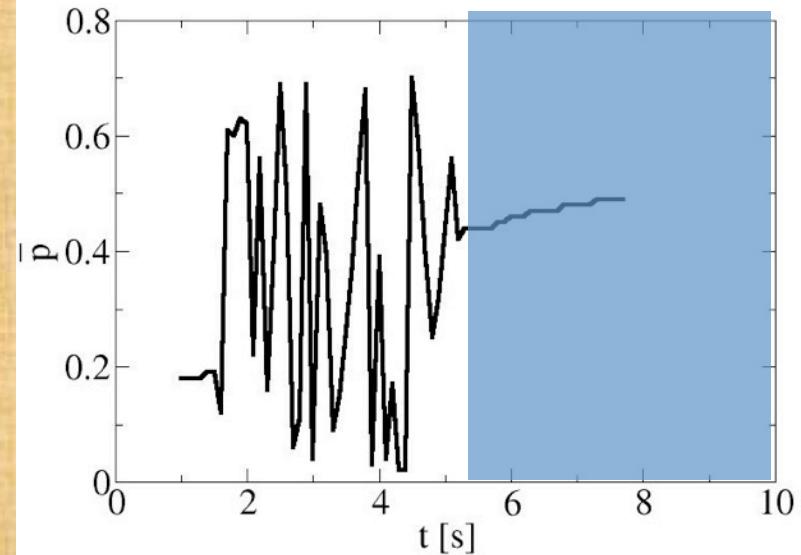
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inverted hierarchy, large θ_{13}



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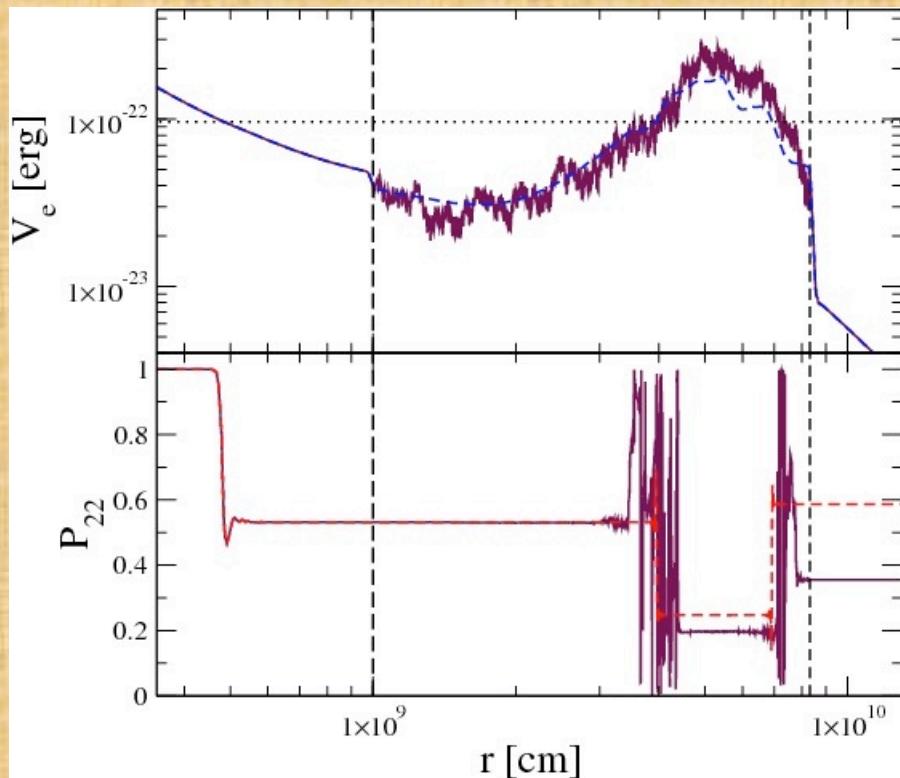
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Turbulence effects

Loreti et al, PRD 52 (1995); Balantekin et al, PRD 54 (1006); Friedland and Gruzinov, hep/0607244, Fogli et al JCAP 0606 (2006); Kneller arXiv: 1004.1288.

Calculation of instantiations of the neutrino amplitudes- not of average probabilities - in presence of shock waves and turbulence.

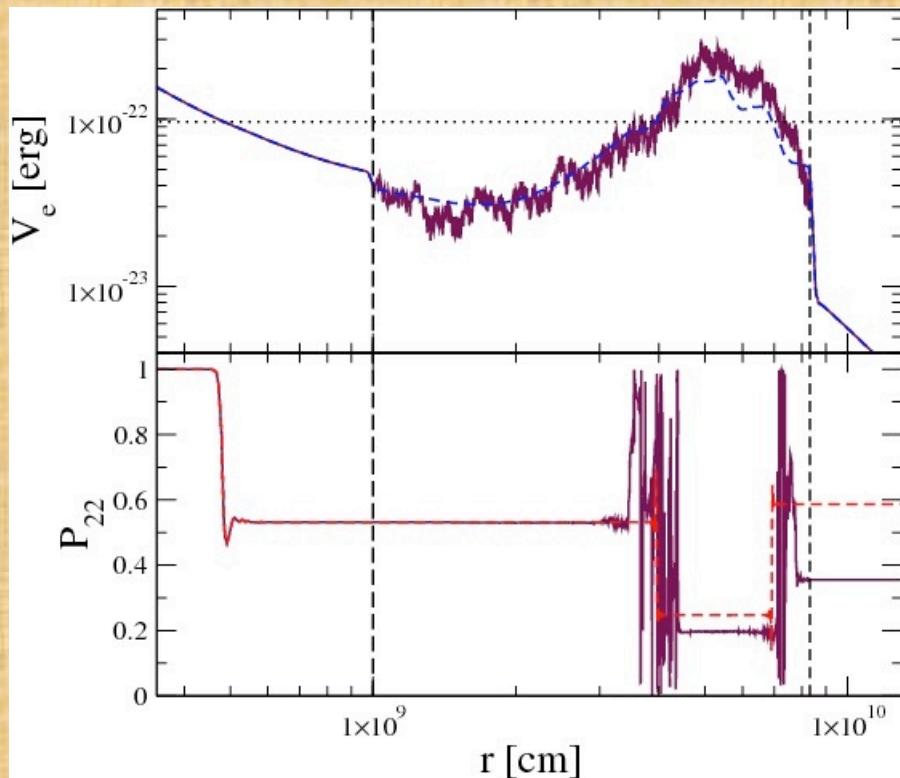


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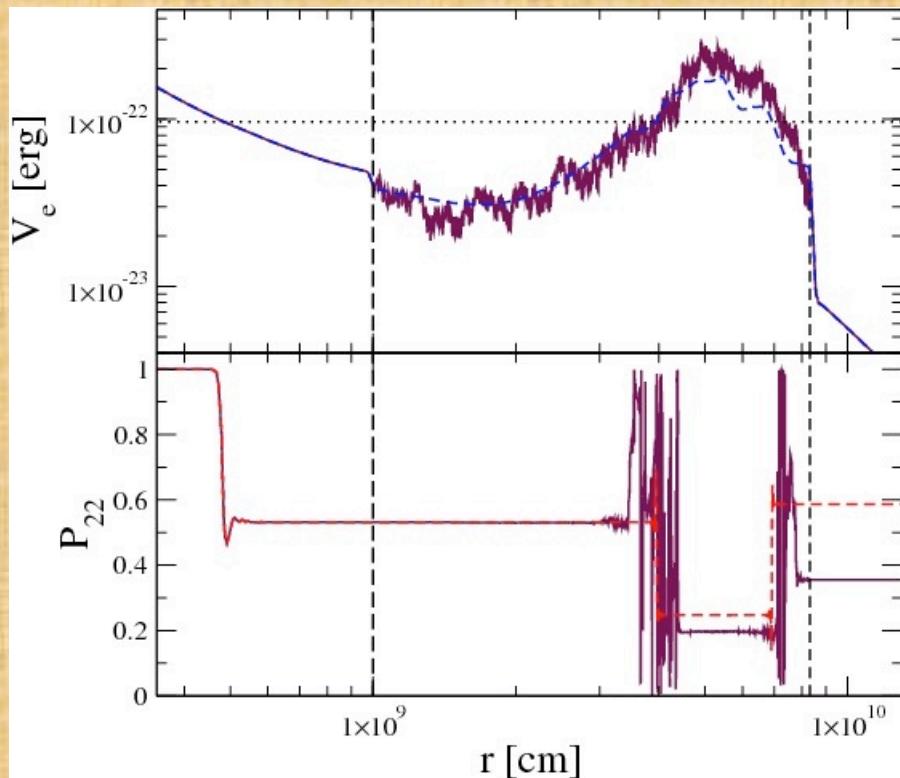
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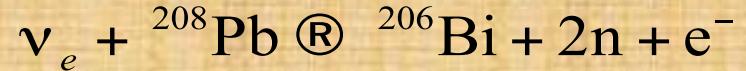
TRANSITION FROM shock-wave to turbulence dominated regime, as the fluctuation amplitude increases.

Kneller and Volpe, PRD (2010) , arXiv: 1006.0913 .

The HALO project



80 tonnes of Pb, HALO-2: 1kt



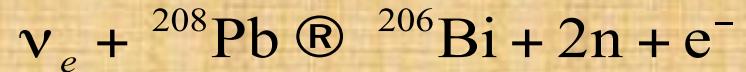
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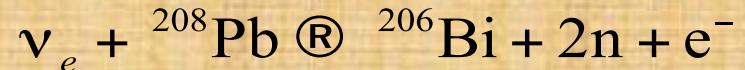
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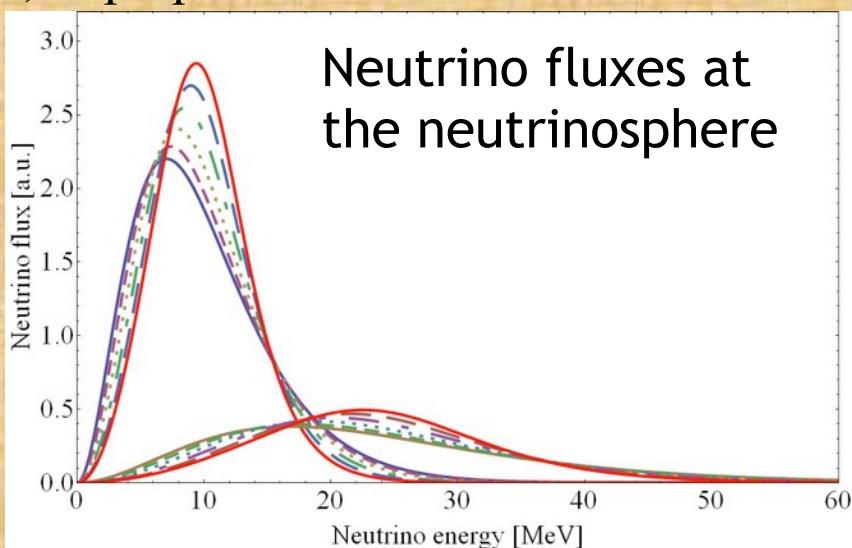


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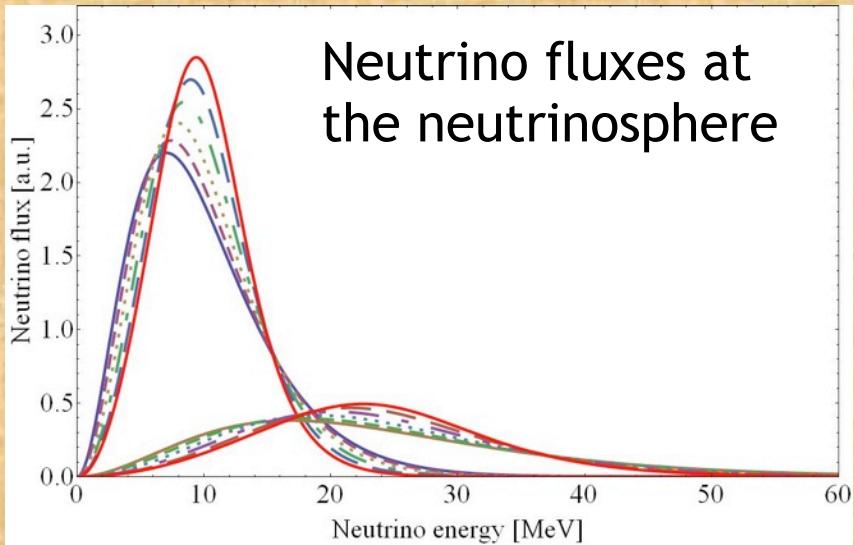
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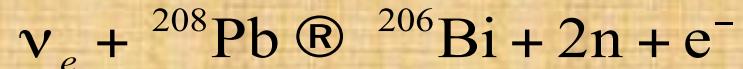
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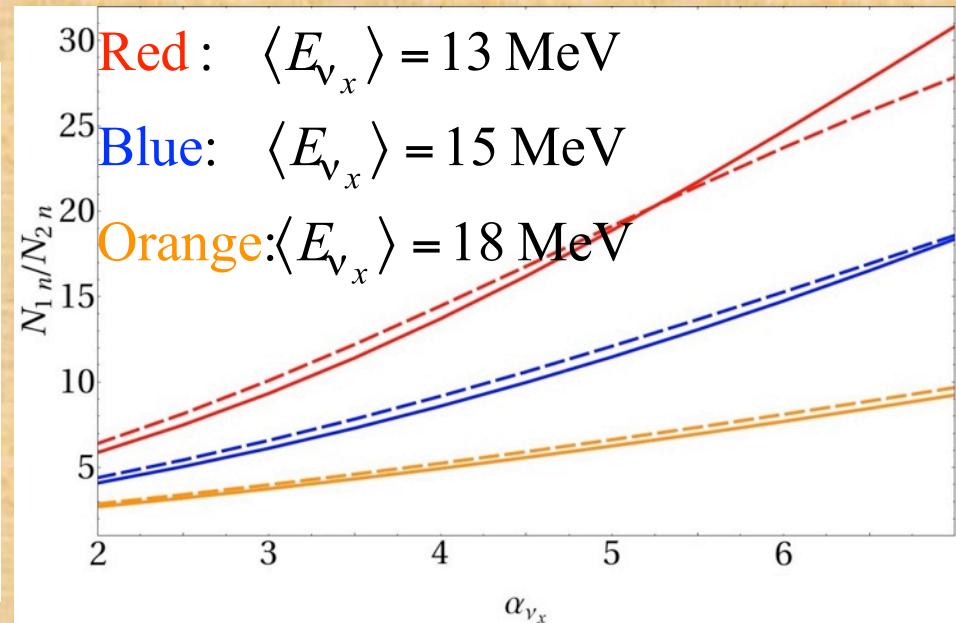
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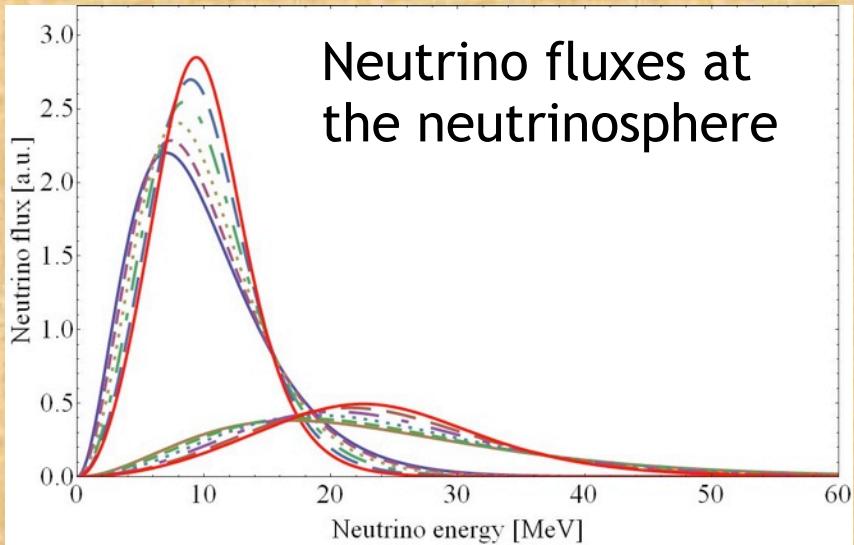
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The HALO project



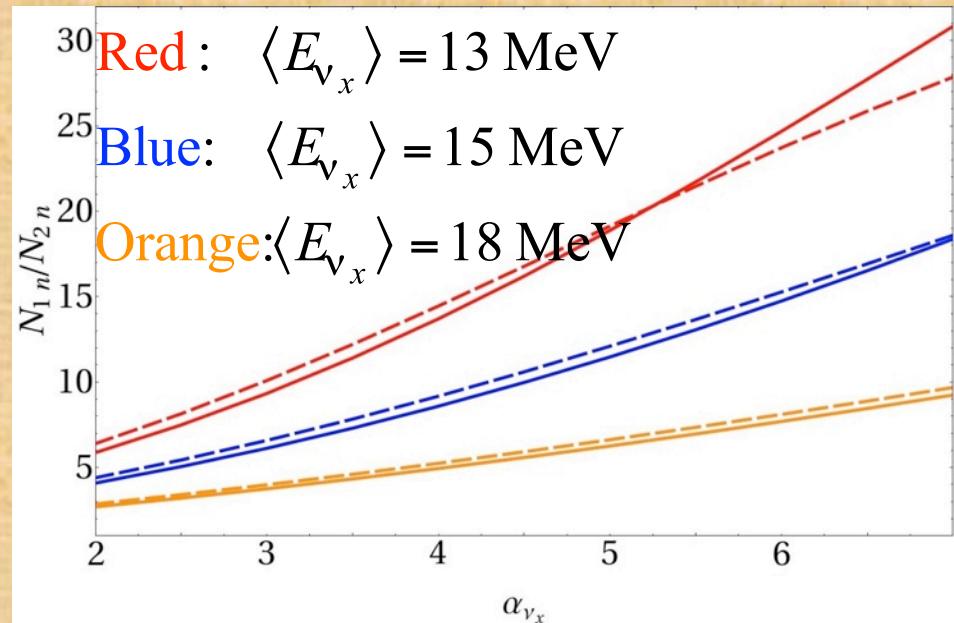
C. Volpe et, D. Vaananen
, in preparation.



80 tonnes of Pb, HALO-2: 1kt



Important to have different energy thresholds.



Extracting the primary ν -flux pinching parameter

CP violation effects in astrophysics and cosmology (BBN epoch)

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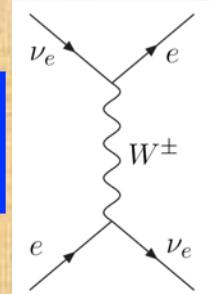
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matterterm

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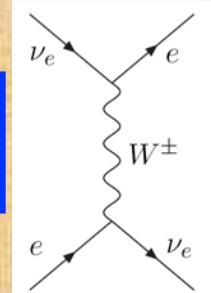
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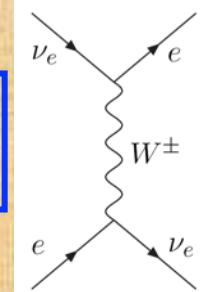
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factorizes out easily and gives:



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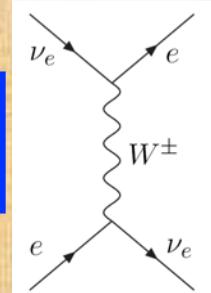
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the T23 basis

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CP violation effect in SN

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$$\tilde{H}(\delta) = S^\dagger \tilde{H}(\delta = 0) S$$

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Evolution operator
in the T_{23} basis



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↑
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At tree level no CP effects, however ...

Conditions for CP violation effects

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Hamiltonian does not
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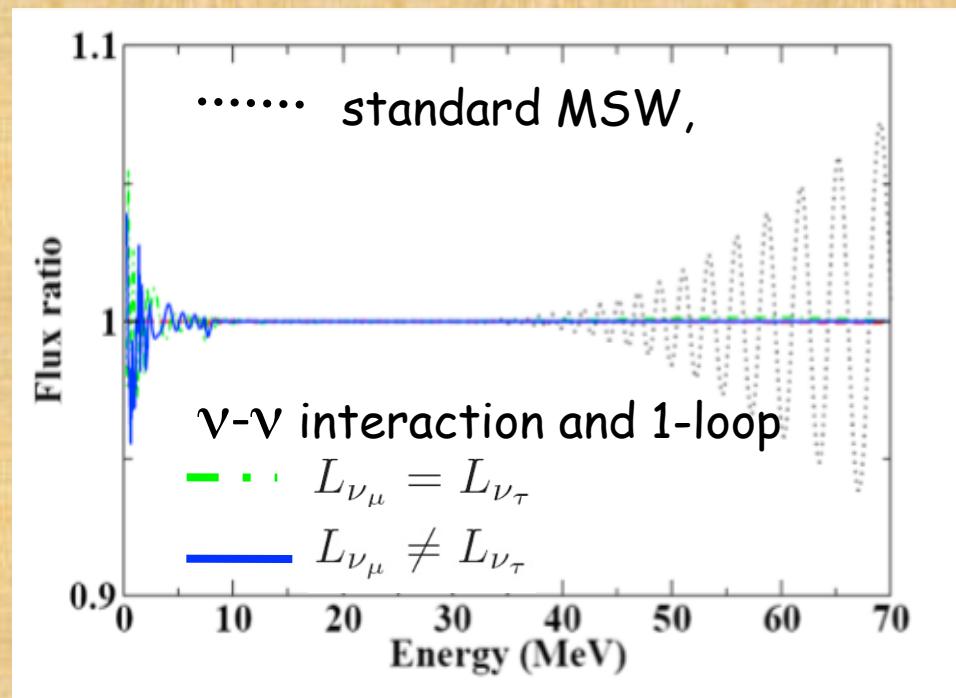
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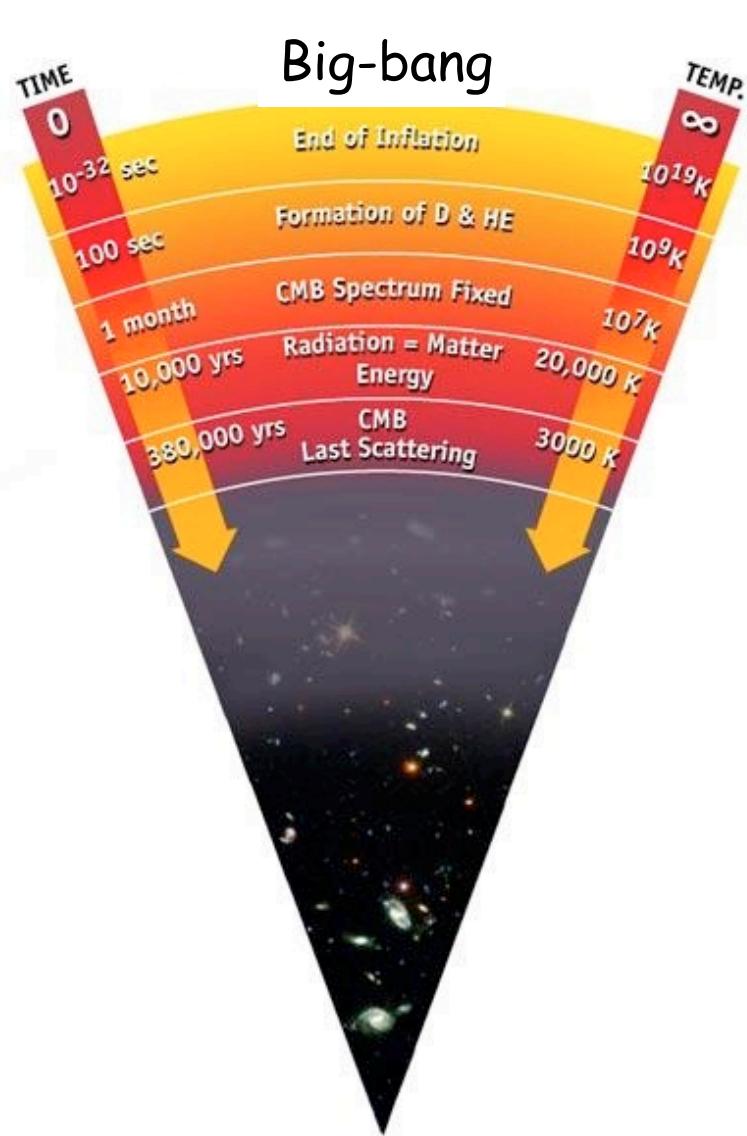
THERE CAN BE CP-VIOLATION EFFECTS IN SUPERNOVAE.

Numerical results

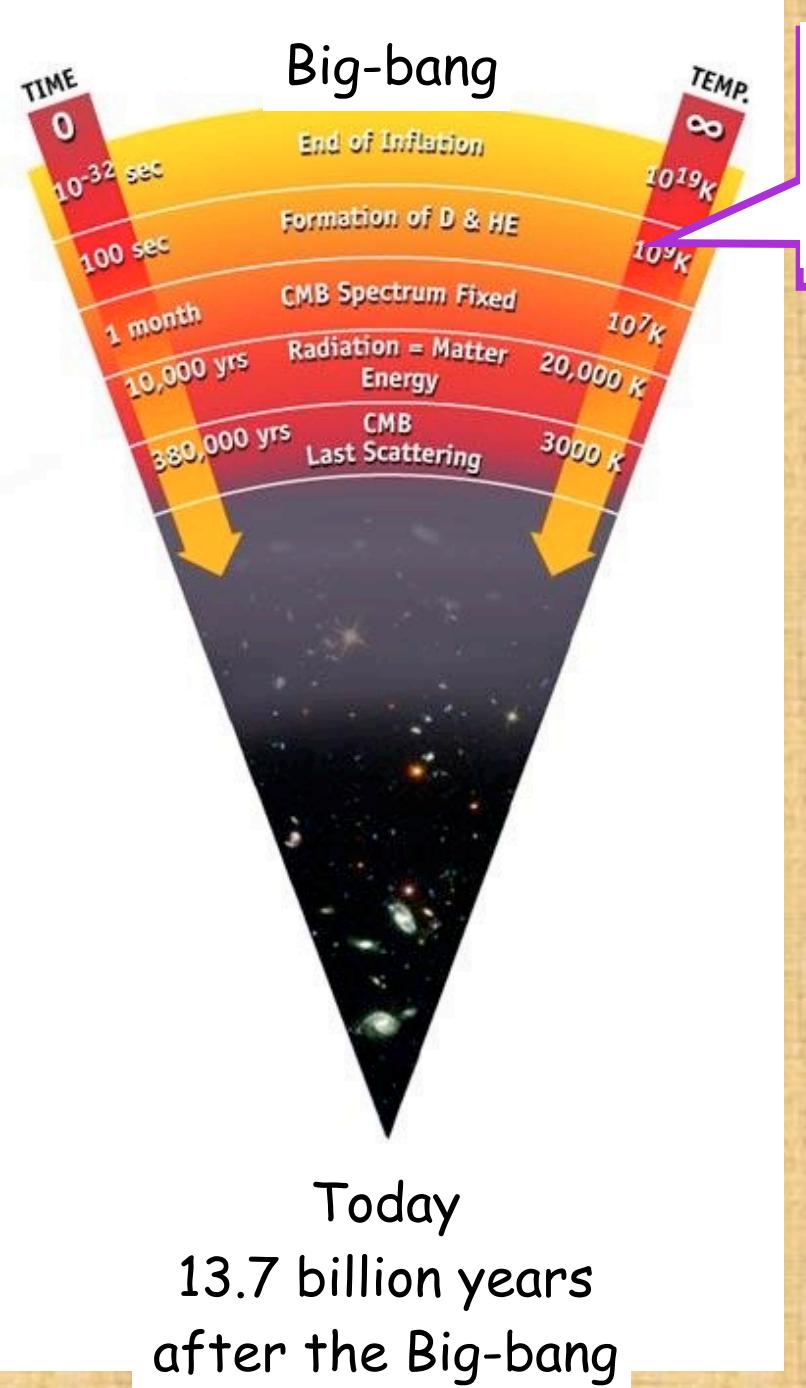


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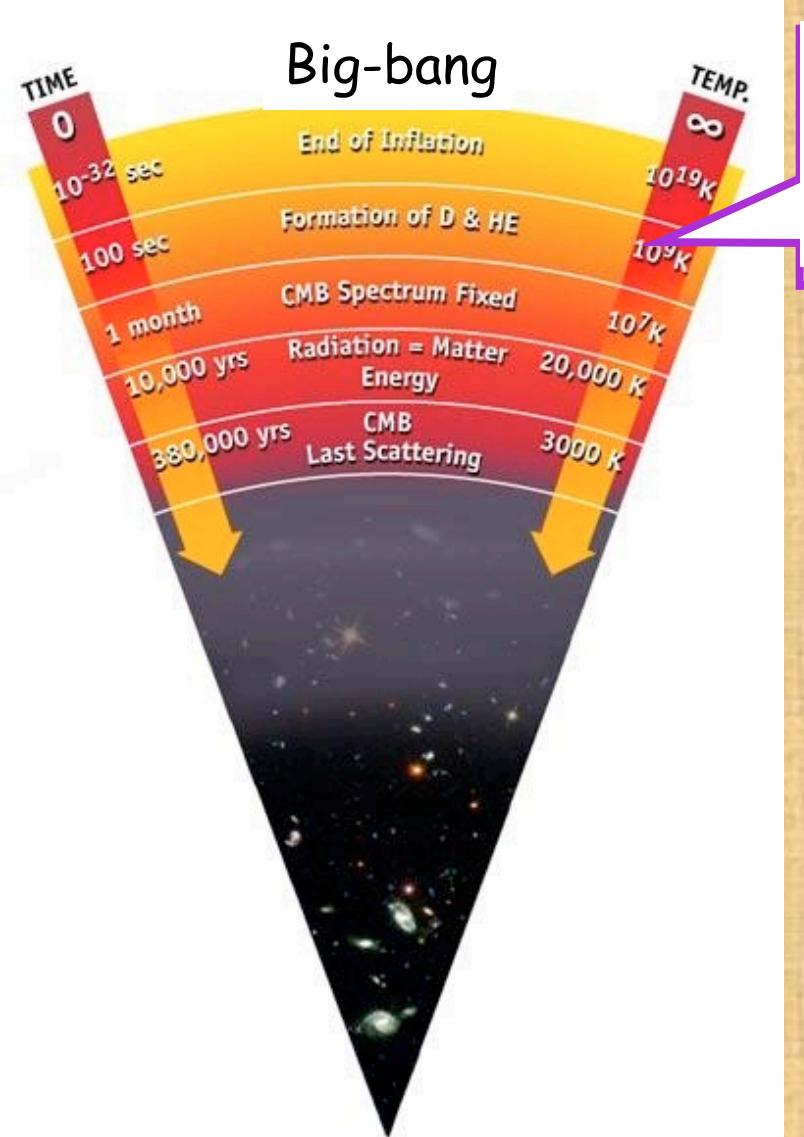
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Neutrino properties leave an imprint on the neutron/proton ratio that sets the abundance of light elements.

Several properties have been studied in the past (ex. sterile neutrinos, non-standard interactions, ...)

A.D. Dolgov, Phys. Rept. 370, 333
(2002), hep-ph/0202122.

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Calculations including mixing, e.g.

Abazajian, Beacom, Bell, PRD 66 (2002) 013008,
Mangano et al, Nucl. Phys. B 729 (2005) 221,
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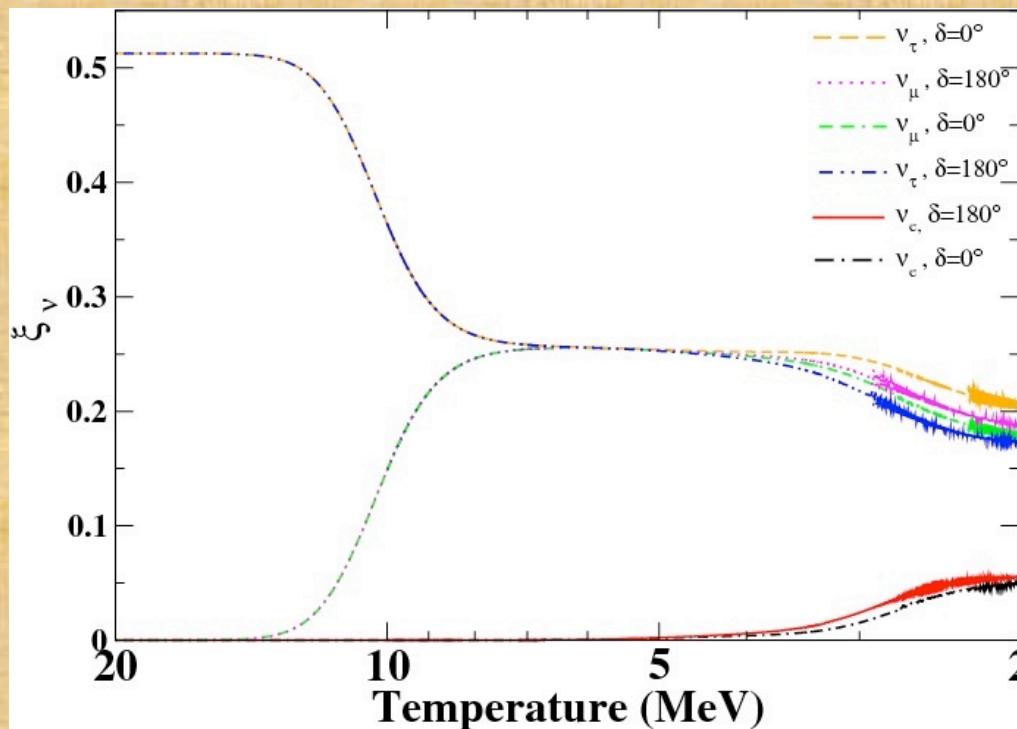
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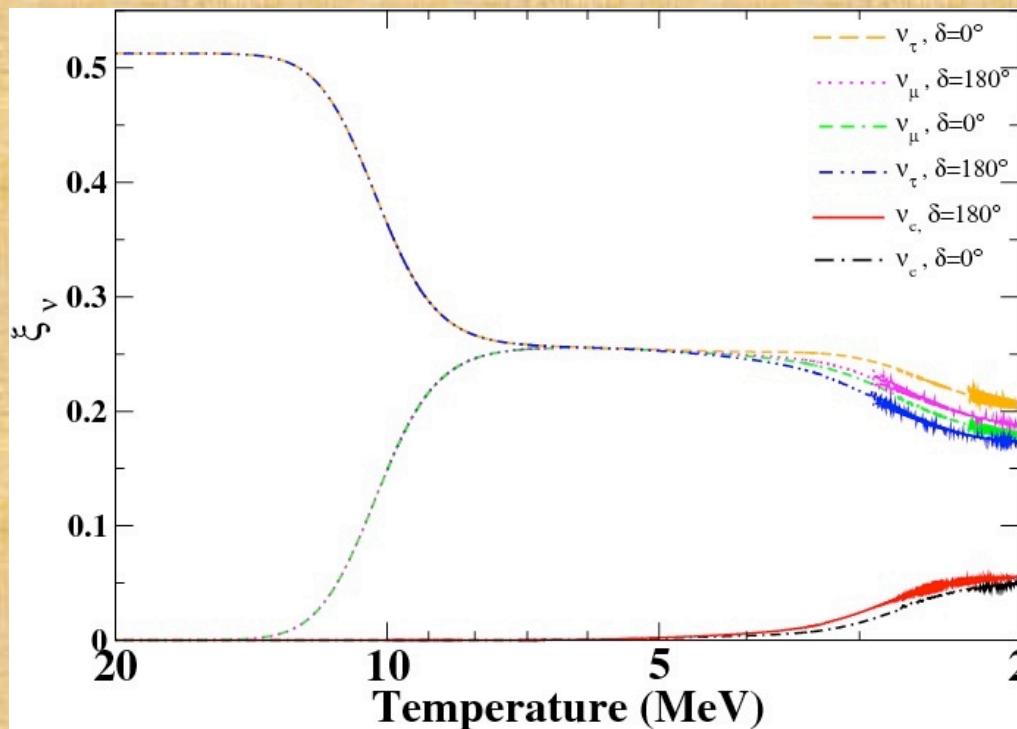
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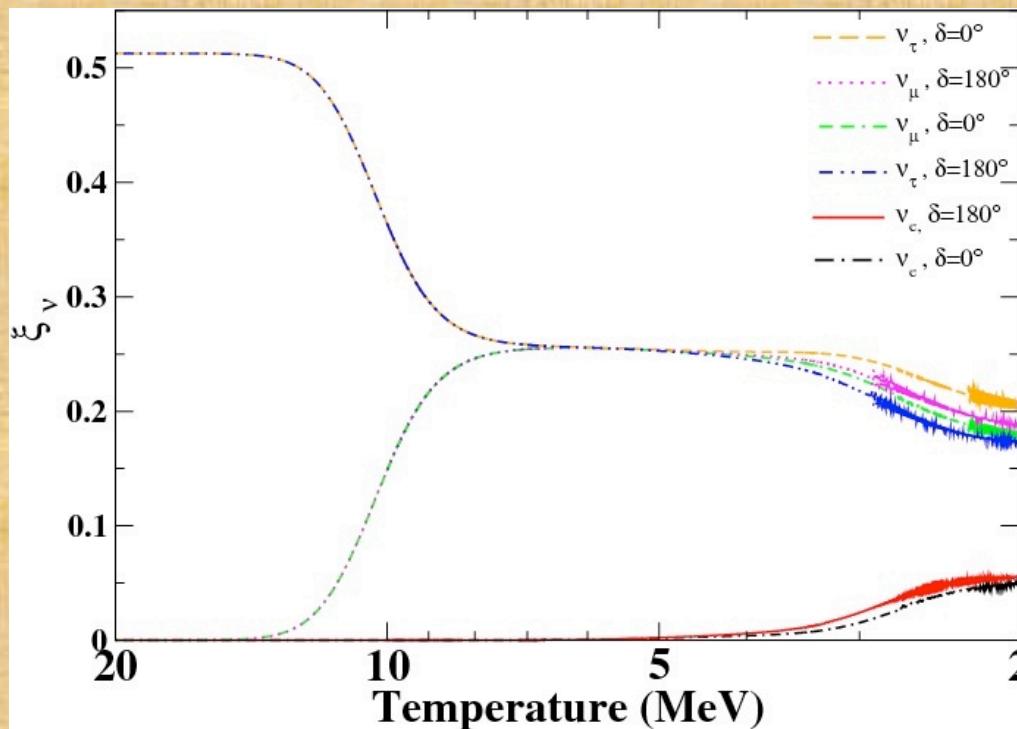


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*In D. Horn's Master Thesis we are now
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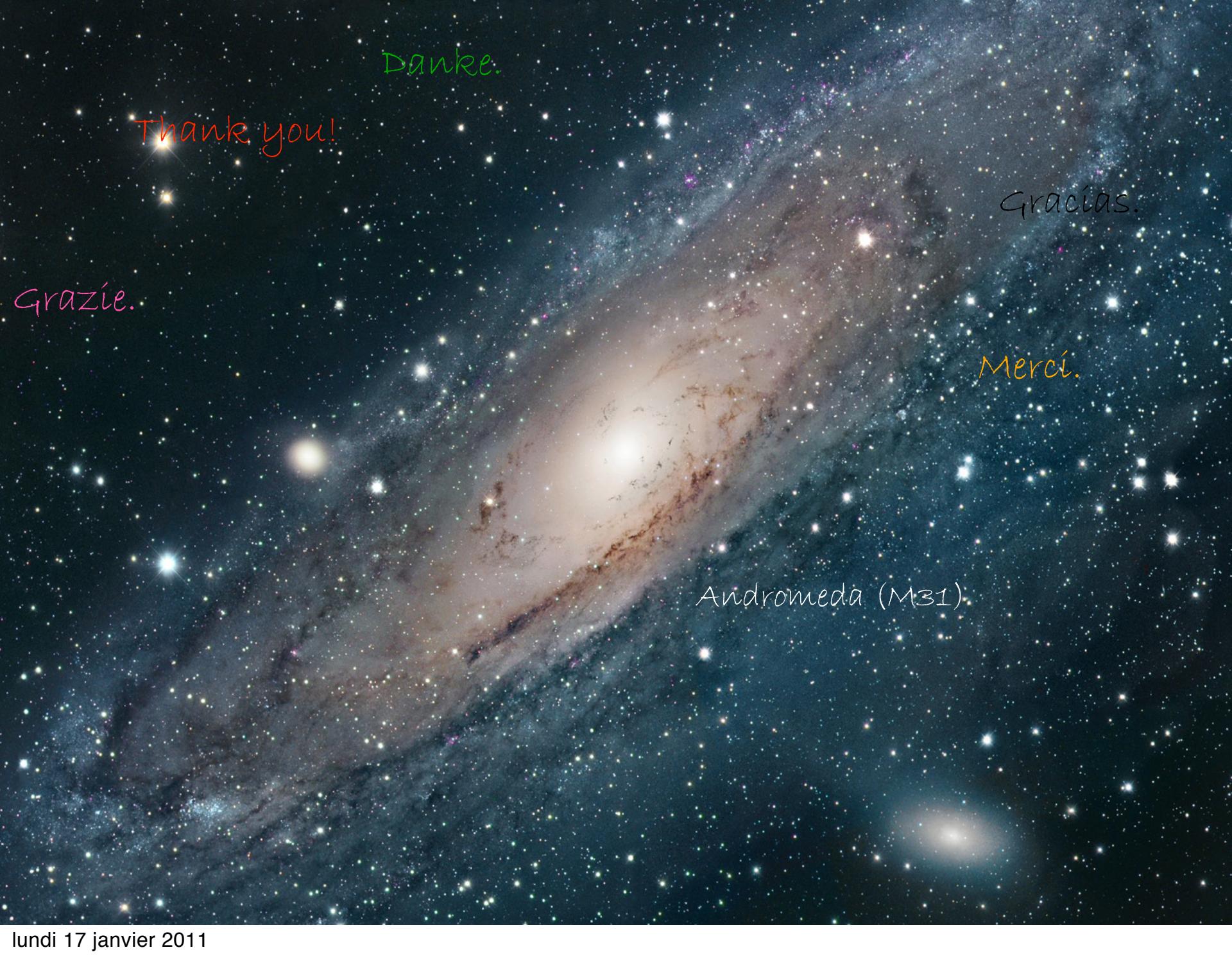
Conclusions and Perspectives



We have set the basis for the exploration of CP violating effects in media.

In particular, radiative corrections and beyond the Standard Model physics can engender CP effects on the neutrino fluxes.

Further work is needed to study the possible impact e.g. on the r-process and in other Environments, such as the Early Universe.



Danke.

Thank you!

Gracias.

Grazie.

Merci.

Andromeda (M31)

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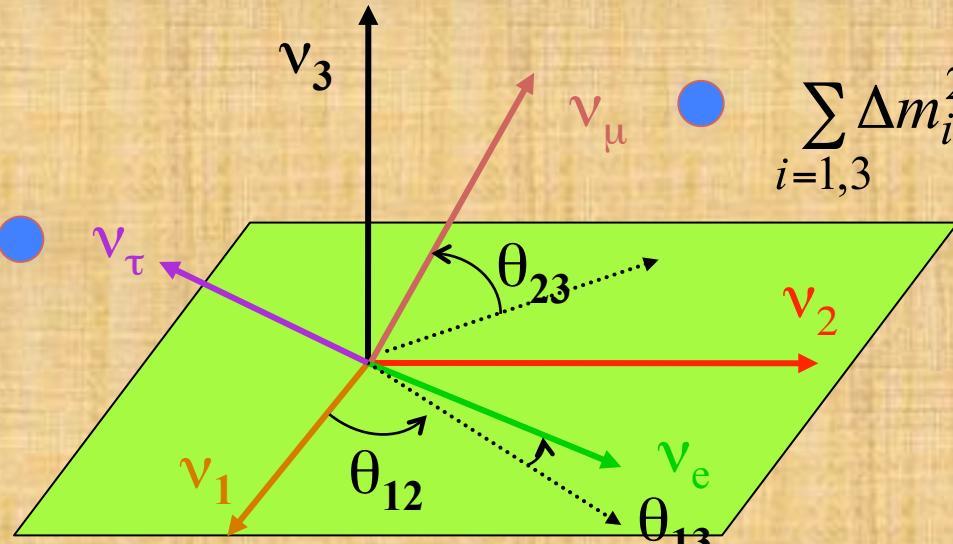
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at the European Center for Theoretical Physics in Trento
« Neutrinos in Nuclear-, Particle and Astrophysics »
April-June 2010 See <http://www.ect.it/>

THE 3-flavours OSCILLATION PARAMETERS

- In the case of three families, there are three mass eigenstates (ν_1, ν_2, ν_3) and three flavour eigenstates $(\nu_e, \nu_\mu, \nu_\tau)$.



$$\Delta m_{21}^2$$

$$\Delta m_{32}^2$$

$$\sum_{i=1,3} \Delta m_i^2 = m_2^2 - m_1^2 + m_3^2 - m_2^2 + m_1^2 - m_3^2$$

Only two Δm^2 are independent.

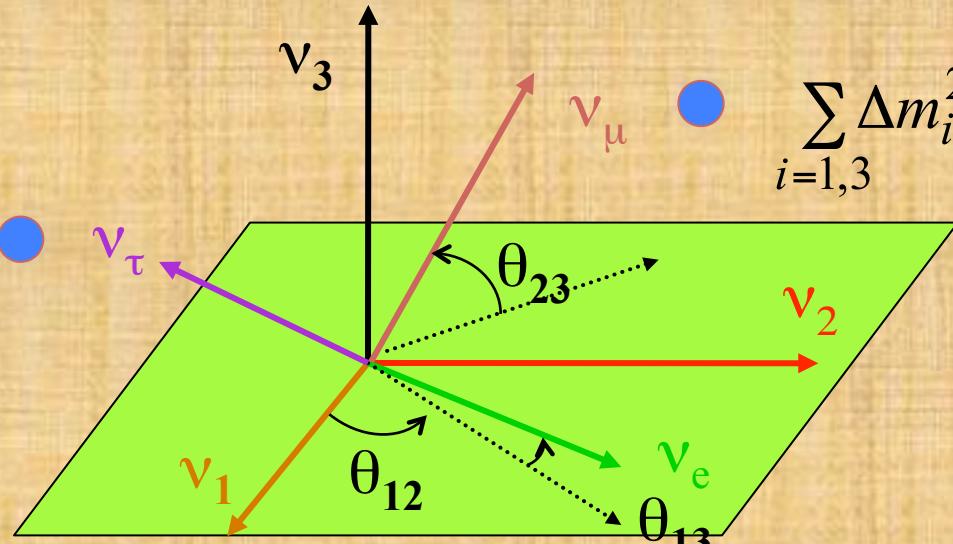
The two basis are related by a unitary matrix, called the Maki-Nakagawa-Sakata-Pontecorvo (MNSP) matrix.

$$c_{ij} = \cos \theta_{ij} \quad s_{ij} = \sin \theta_{ij}$$

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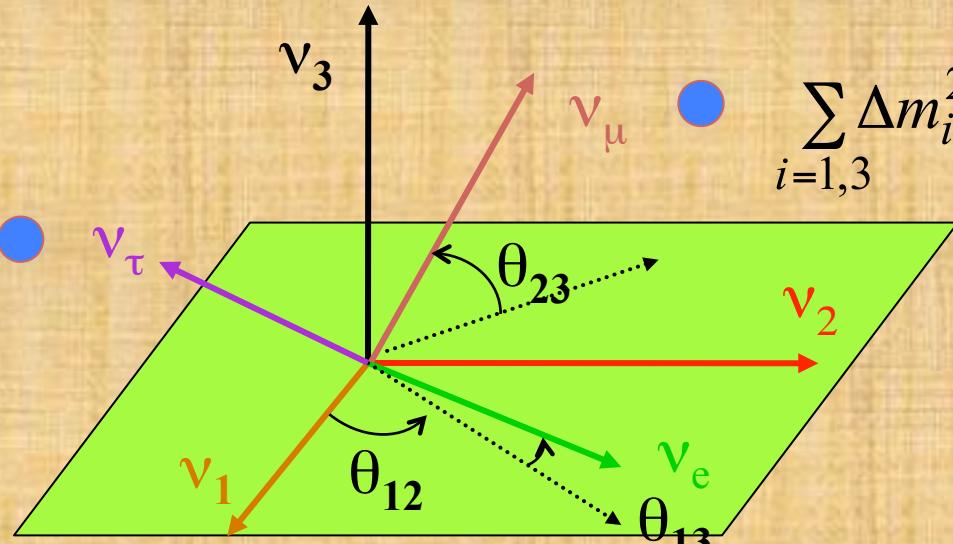
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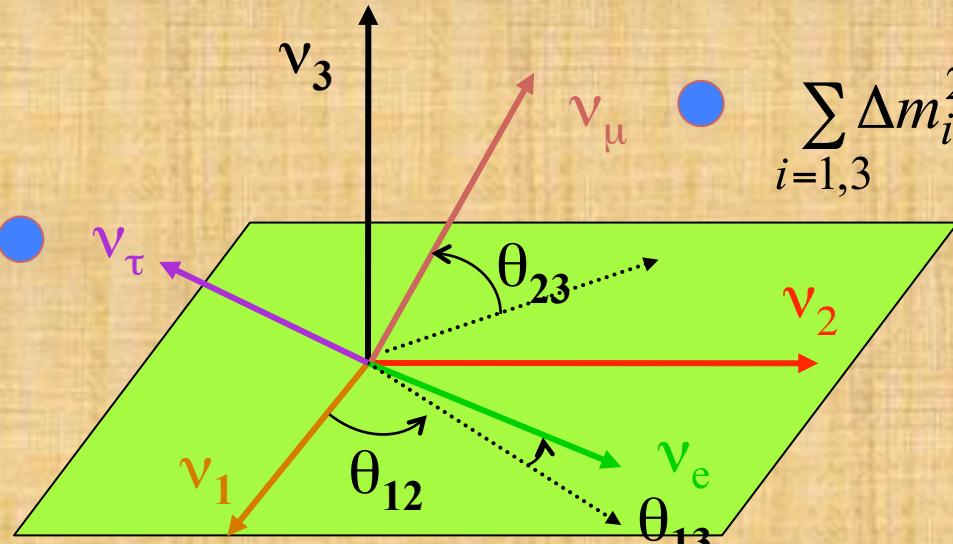
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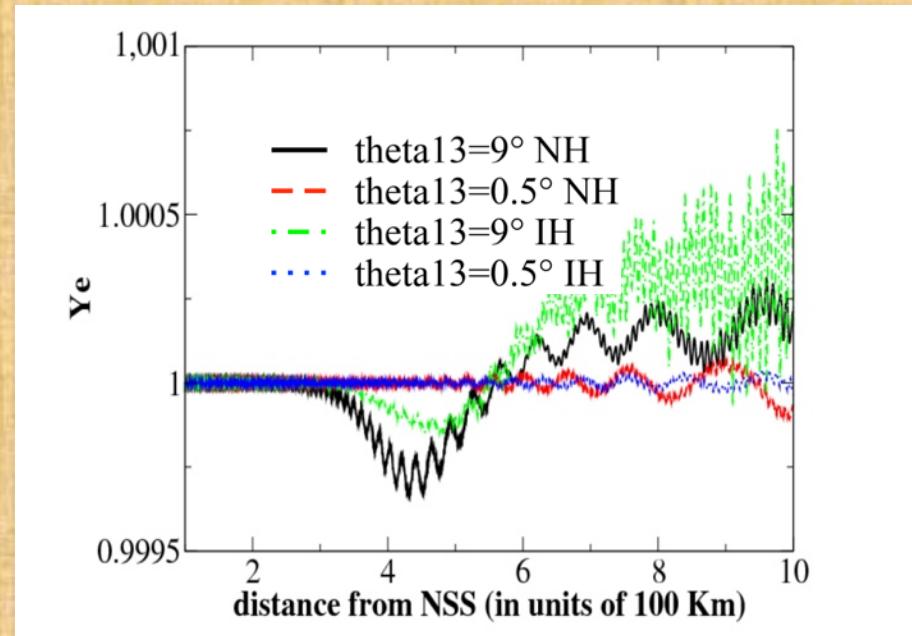
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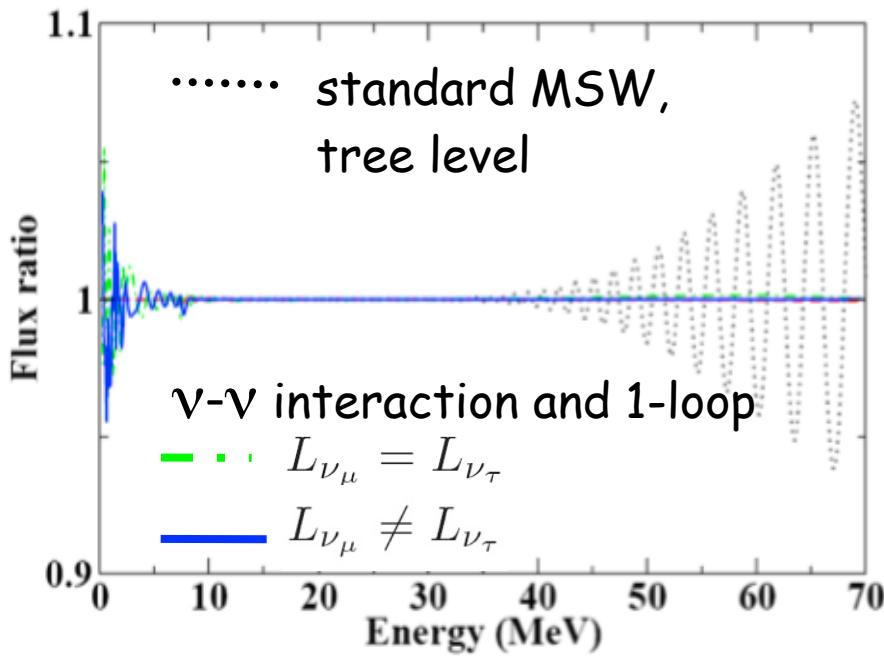
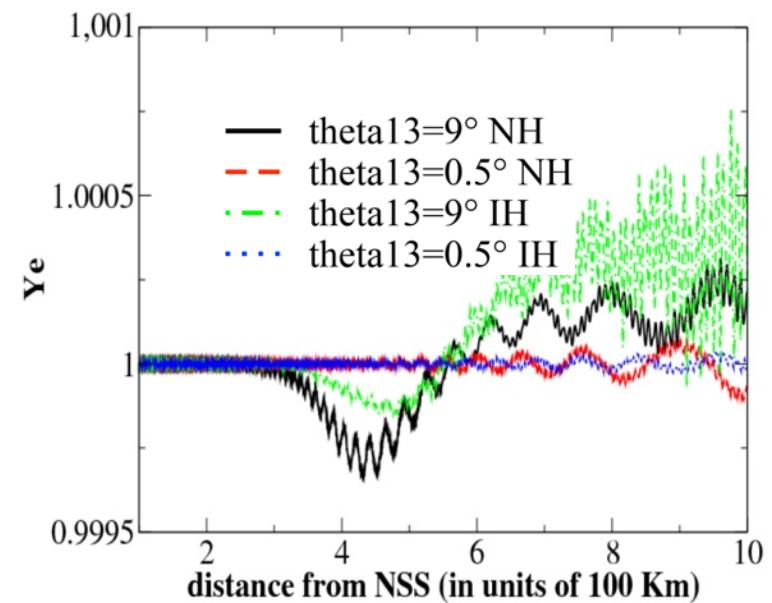
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