

(Q-PYTHIA: a) Monte Carlo Framework for Jet Quenching

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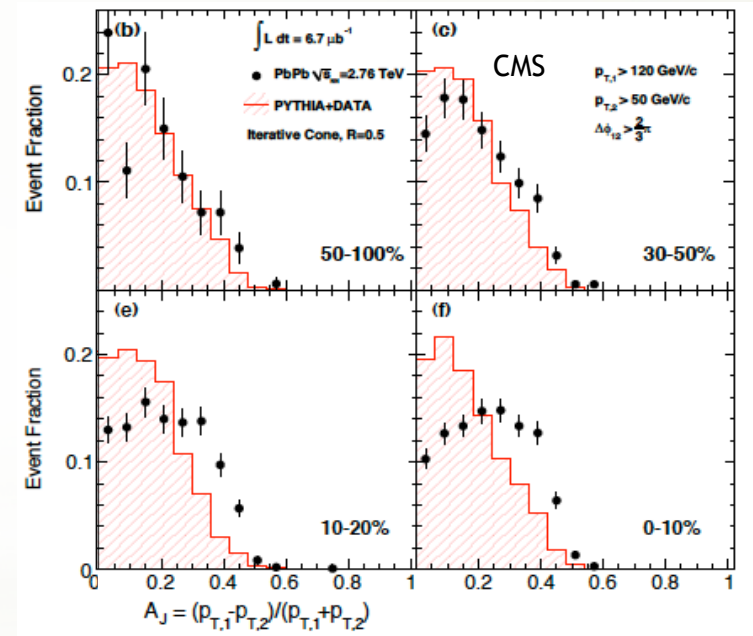
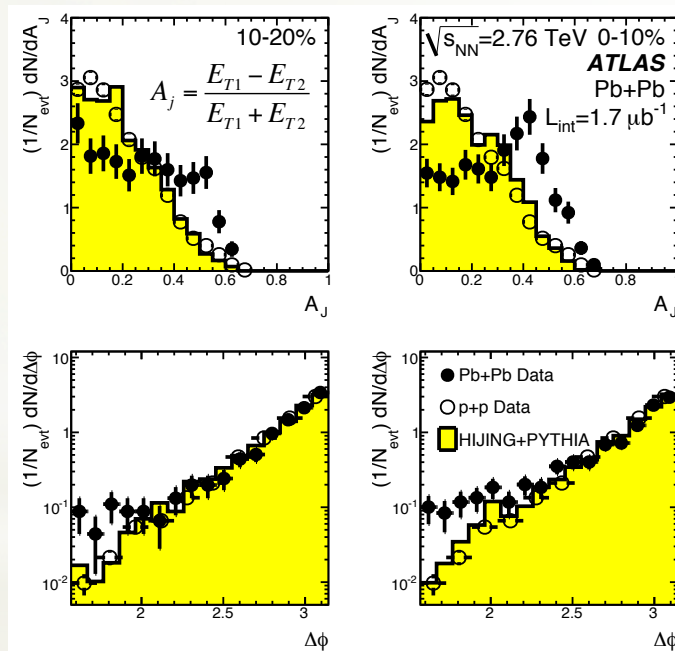
High-pt Probes of High Density QCD at the LHC

May 30th - June 1st, 2011

École Polytechnique, Palaiseau

Introduction

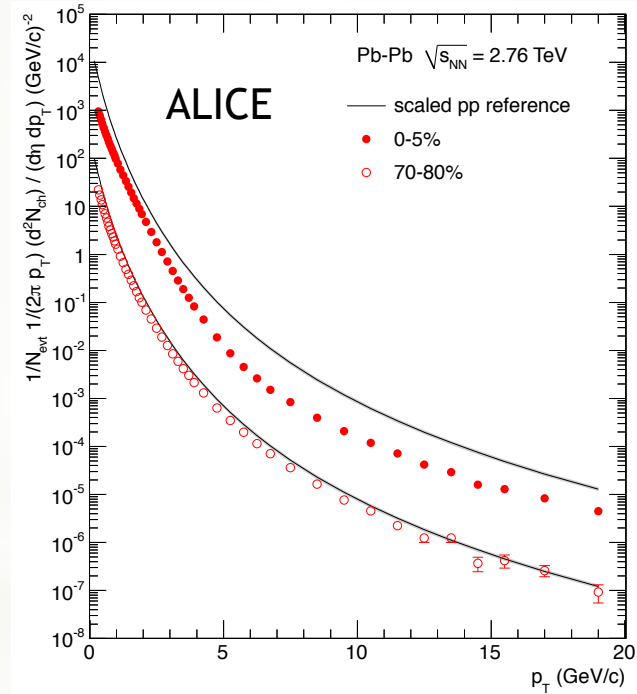
- ✦ Jet Quenching has been established as a fundamental tool to study hot matter in HIC: already at the LHC!!!



Di-jet Asymmetry

Introduction

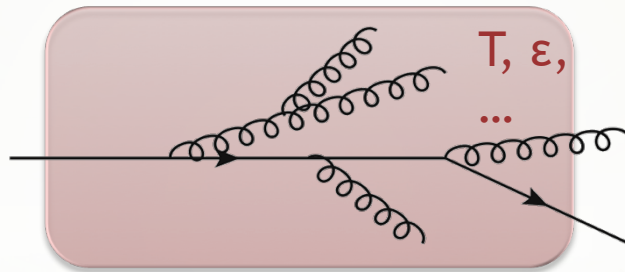
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Suppression of high- p_T spectrum

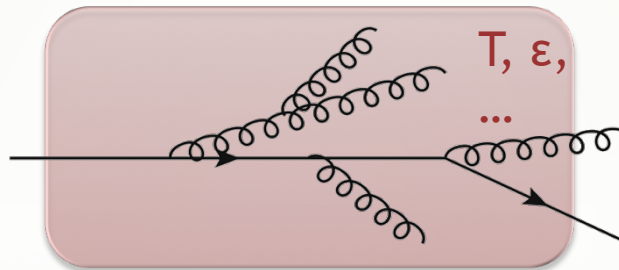
Introduction

- ✦ How does the medium modify a jet?
 - ✦ Radiative energy loss of a high-energy parton
 - ✦ Modification of the standard QCD radiation pattern
 - ✦ Energy loss $\longleftrightarrow T, \epsilon, \dots$



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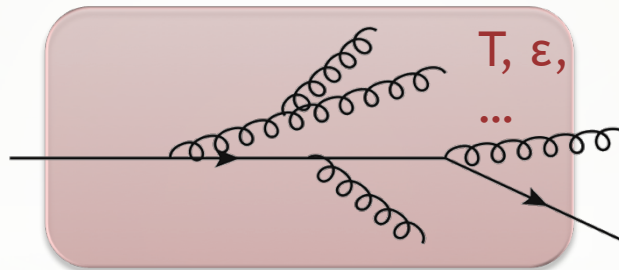
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 - ✦ high-energy approximation
 - ✦ assumptions on multiple emissions

Introduction

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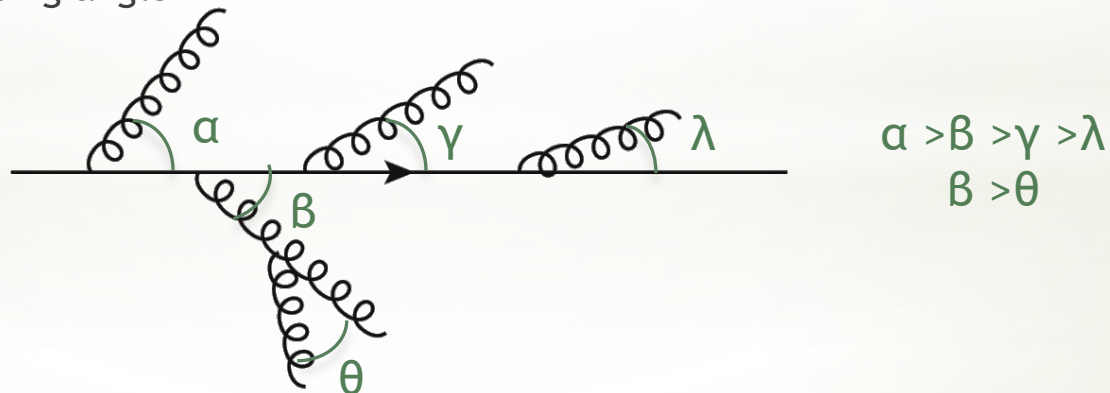
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Monte Carlo
(PYQUEN, YaJem, Q-PYTHIA/
HERWIG, HIJING, MARTINI,...)

Jet Formation

- ★ Vacuum (PYTHIA, HERWIG, SHERPA):
 - ★ Process of branching characterized by $P_{a \rightarrow bc}$: splitting functions
 - ★ Each parton characterized by some virtuality scale, Q^2 ($t=m^2, p_t^2, \theta$: all of them equivalent at high energies)
 - ★ Evolution downwards in Q^2 ('time' ordering)
 - ★ Color coherence effects essential
 - ★ Ordering of subsequent independent emissions in terms of decreasing angle



Jet Formation

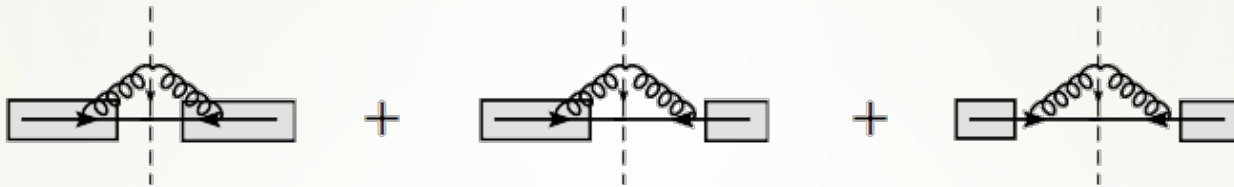
- ✦ Medium (Q-PYTHIA):
 - ✦ Medium-induced gluon radiation taken as the main ingredient
 - ✦ Presence of a medium
 - ✦ Time should play a role as an ordering variable
 - ✦ Ordering variable for multiple gluon emission?
 - ✦ Assume $Q^2 = m^2$ but eventually correct for the finite formation time of the gluons
 - ✦ Independence of multiple gluon emission when re-scattering with the medium as in vacuum

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 - ✦ Assume $Q^2 = m^2$ but eventually correct for the finite formation time of the gluons
 - ✦ Independence of multiple gluon emission when re-scattering with the medium as in vacuum
 - ✦ Does not consider:
 - ✦ Recoil (elastic energy loss)
 - ✦ Modification of the color structure of the shower by exchanges with the medium
 - ✦ Back-reaction
 - ✦ In-medium hadronization

Splitting Functions

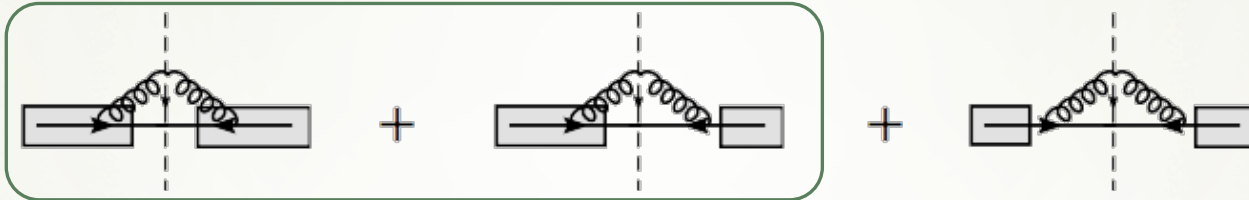
✦ Total medium-induced radiation spectrum



$$\frac{dI^{\text{tot}}}{dz d\mathbf{p}_T^2} = \frac{dI^{\text{vac}}}{dz d\mathbf{p}_T^2} + \frac{dI^{\text{med}}}{dz d\mathbf{p}_T^2}$$

Splitting Functions

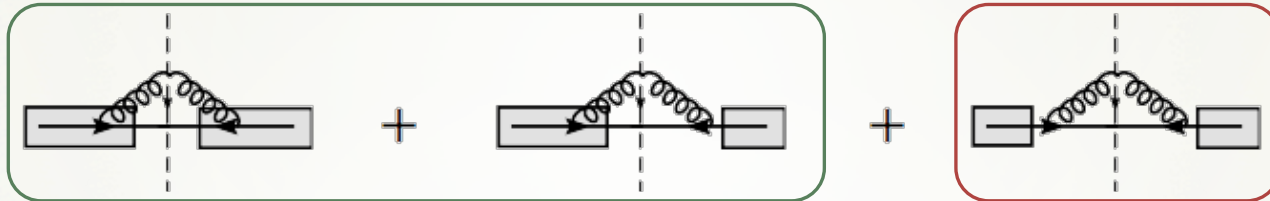
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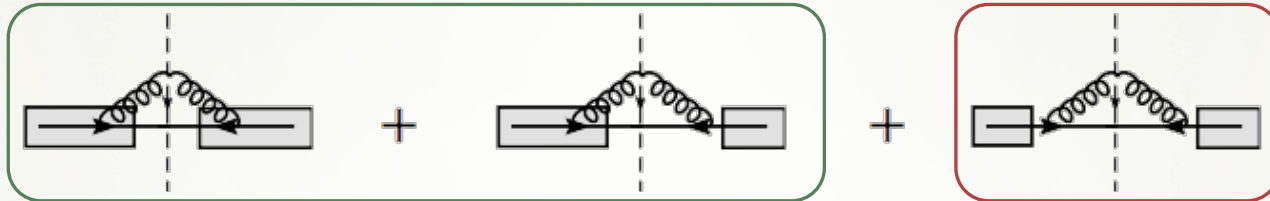


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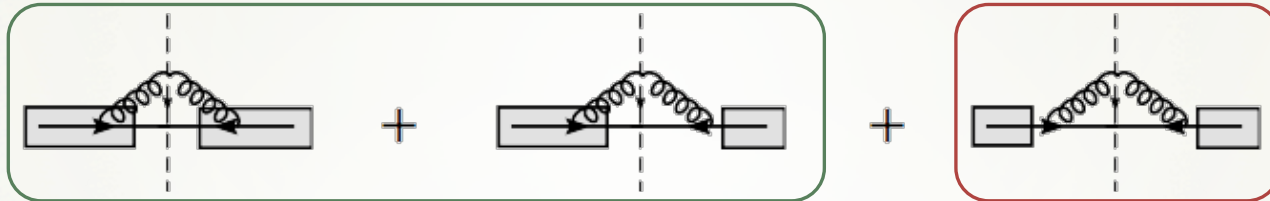
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$$\Delta P(z, t) \simeq \frac{2\pi t}{\alpha_s} \frac{dI^{\text{med}}}{dz dt}$$

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- ★ Extension to hard emission as in vacuum

Sudakov Form Factor

✦ Implementation of the t-ordered final state radiation routine in PYTHIA:

✦ Probabilistic interpretation of DGLAP evolution:

$$D(x, t) = \underbrace{\Delta(t)D(x, t_0)}_{\text{No splitting between } t_0 \text{ and } t} + \underbrace{\Delta(t) \int_{t_0}^t \frac{dt_1}{t_1} \frac{1}{\Delta(t_1)} \int \frac{dz}{z} P(z) D\left(\frac{x}{z}, t_1\right)}_{\text{Contribution when some finite amount of radiation is present}}$$

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$$\Delta(t_1) = \exp \left[- \int_{t_0}^{t_1} \frac{dt'}{t'} \int_{z^-}^{z^+} dz \frac{\alpha_s(t_1)}{2\pi} P(z) \right]$$

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$$P_{\text{tot}}(z) = P_{\text{vac}}(z) + \Delta P(z, t, \hat{q}, L, E)$$

Sudakov Form Factor

✦ Medium-modified Sudakov factor

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\hat{q}

Sudakov Form Factor

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★ Medium parameters:

- ★ Length (L)
- ★ Transport coefficient (\hat{q})

Sudakov Form Factor

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Transformed  into

$$\boxed{\omega_c = \frac{1}{2} \hat{q} L^2} \quad \boxed{\hat{q} L}$$

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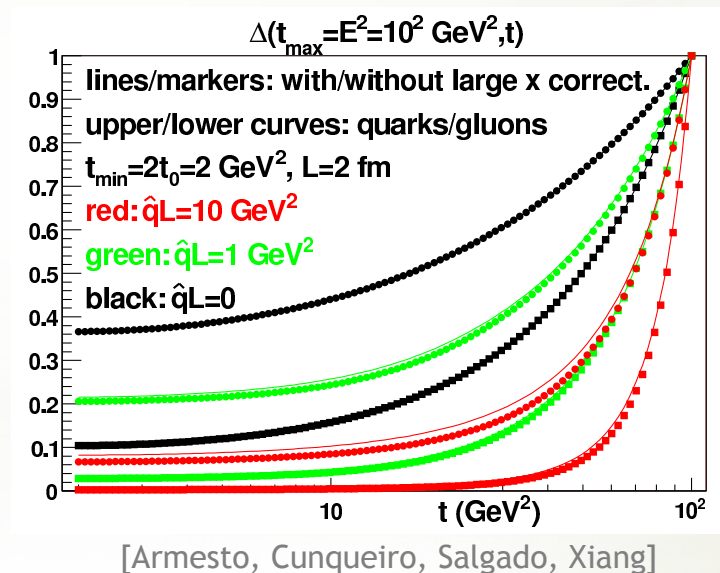
★ Length (L)

★ Transport coefficient (\hat{q})

Transformed  into

$$\omega_c = \frac{1}{2} \hat{q} L^2 \quad \hat{q} L$$

Suppression of the Sudakov:
more radiation in medium than
in vacuum

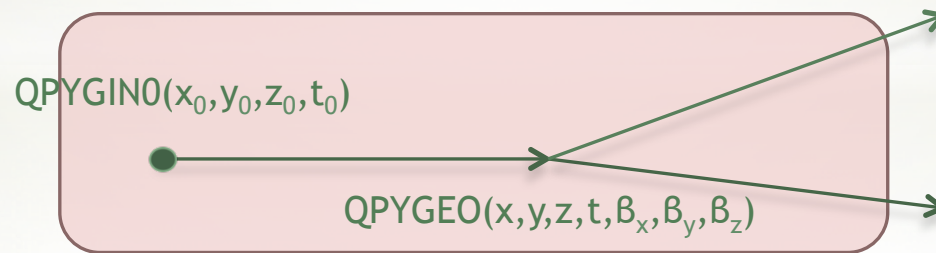


Program

- ✦ Q-PYTHIA is the usual PYTHIA-6.1.18 with a modified final-state radiation:
 - ✦ Only modification: t-ordered FSR routine PYSHOW
 - ✦ Additional auxiliary routines (black box) + two routines that can be modified by the user:

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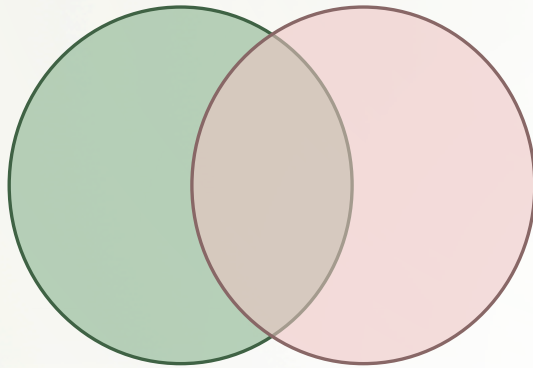
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 - ✦ QPYGIN0(x_0, y_0, z_0, t_0): user specifies the position and time of the hard scattering, to be called once per NN collision
 - ✦ QPYGEO($x, y, z, t, \beta_x, \beta_y, \beta_z, QHL, OC$) computes the parameters (QHL, OC) for a parton located at (x, y, z, t) moving along the direction defined by $(\beta_x, \beta_y, \beta_z)$; medium to be specified by the user (some defaults available)



Medium with:
 $QHL = \hat{q}L$
 $OC = \omega_c$

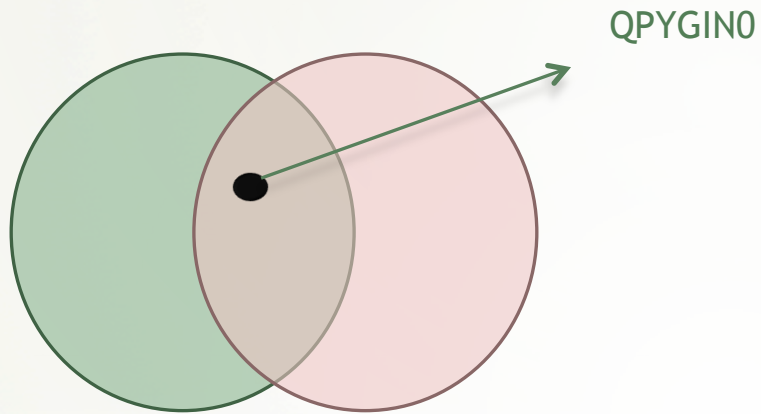
Program

✦ Usage:



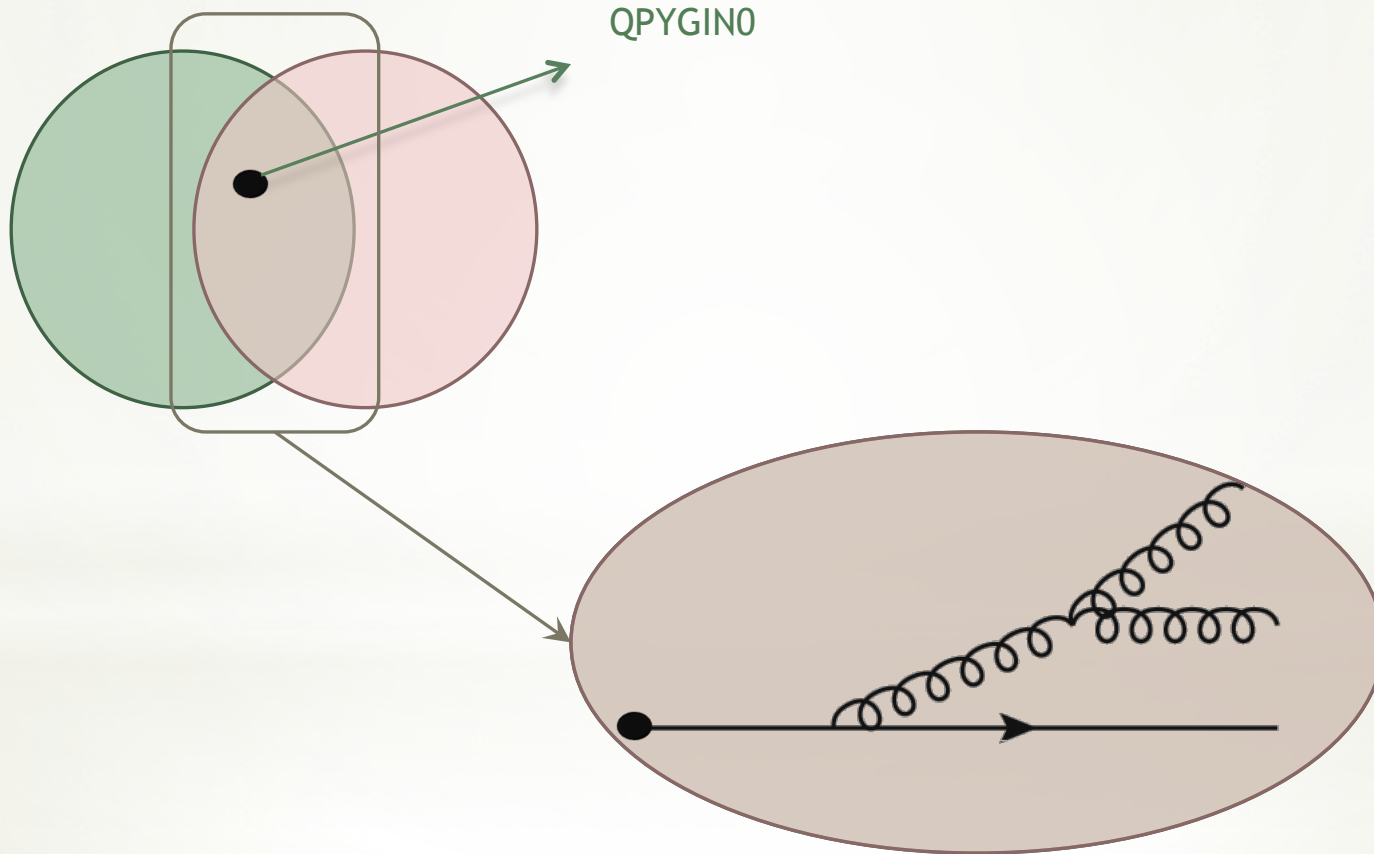
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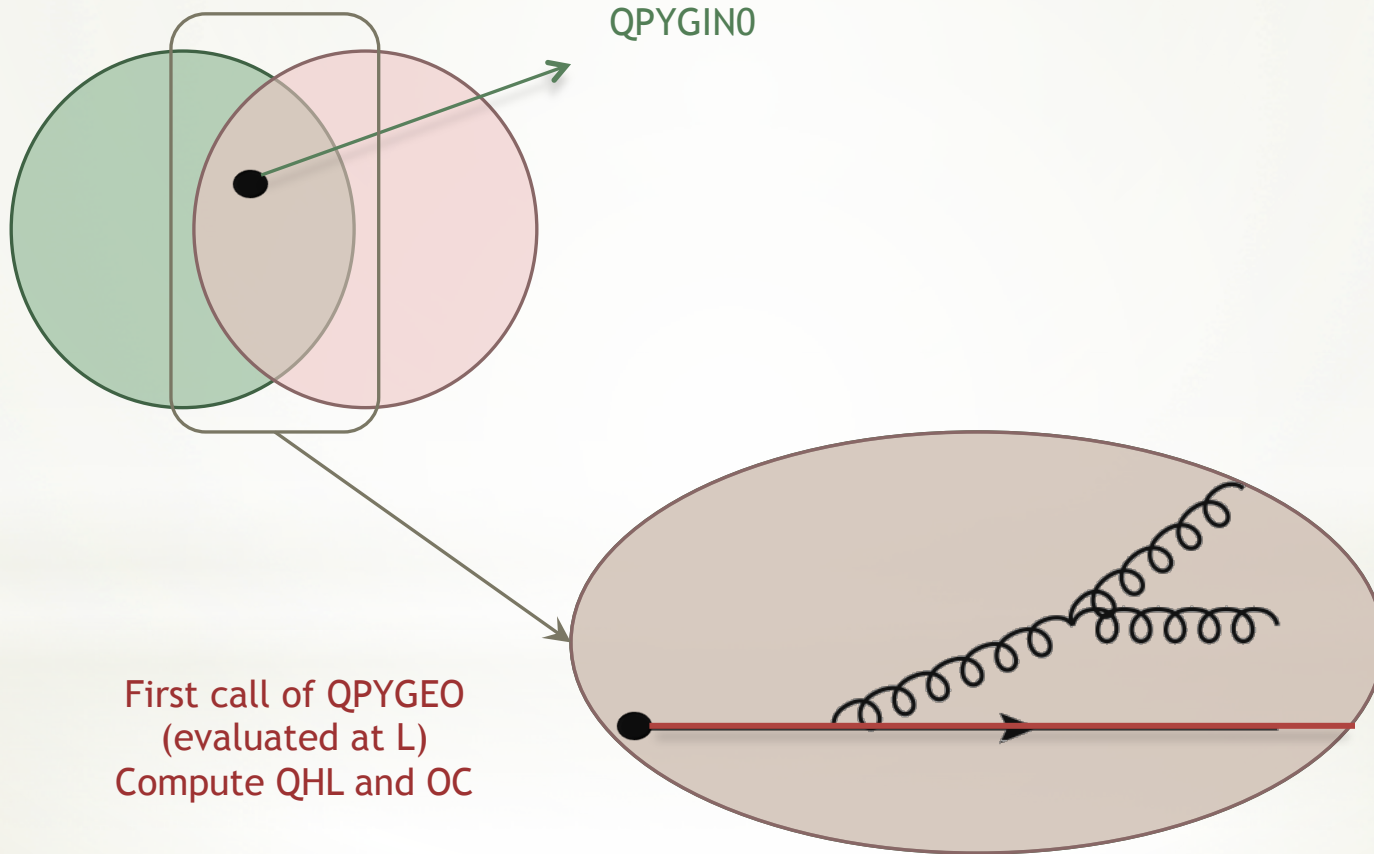
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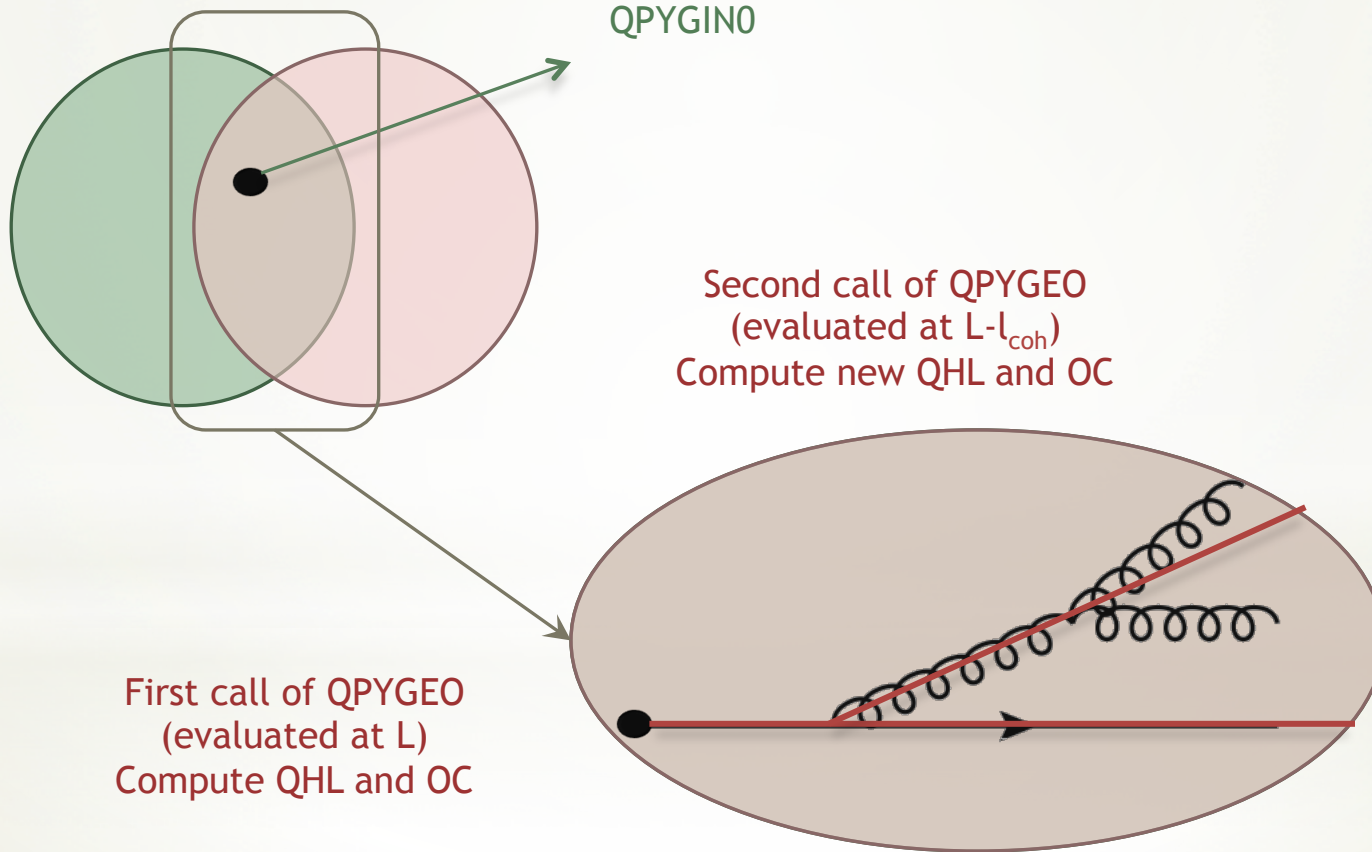
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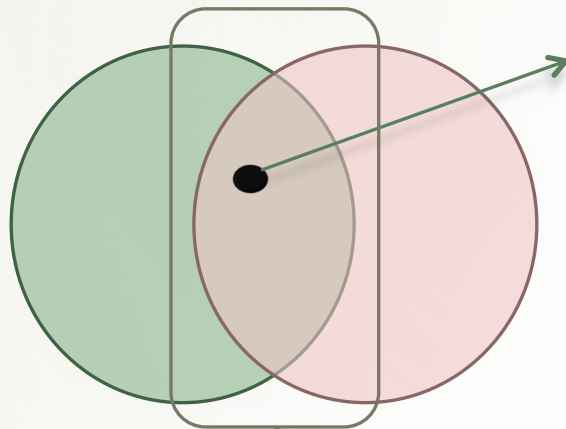
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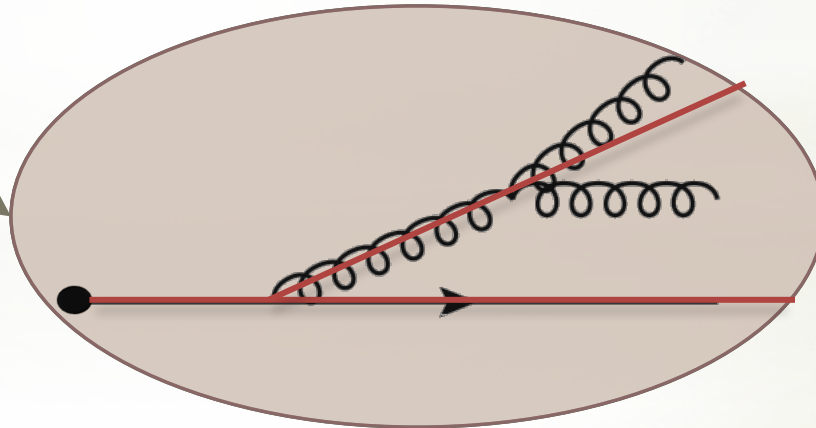


QPYGINO

Moreover:
Energy degradation
considered at each
splitting (ΔP depends
on E)

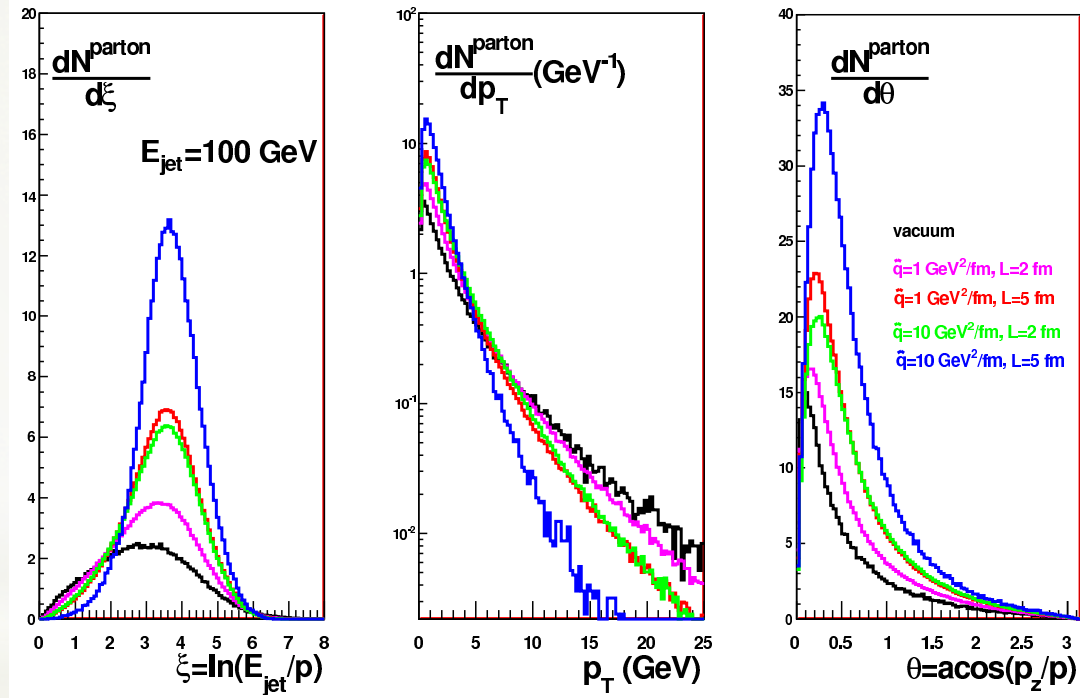
Second call of QPYGEO
(evaluated at $L-l_{\text{coh}}$)
Compute new QHL and OC

First call of QPYGEO
(evaluated at L)
Compute QHL and OC



Results

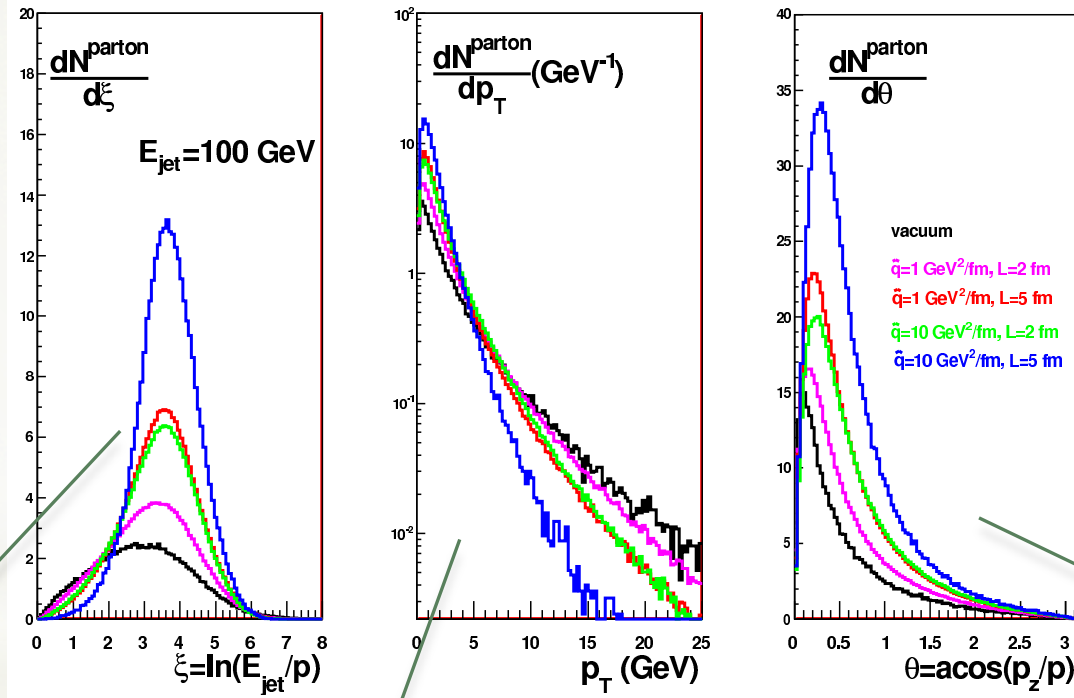
★ Results at parton level:



[Armesto, Cunqueiro, Salgado]

Results

★ Results at parton level:



Suppression of
low- ξ particles
Enhancement of
large- ξ

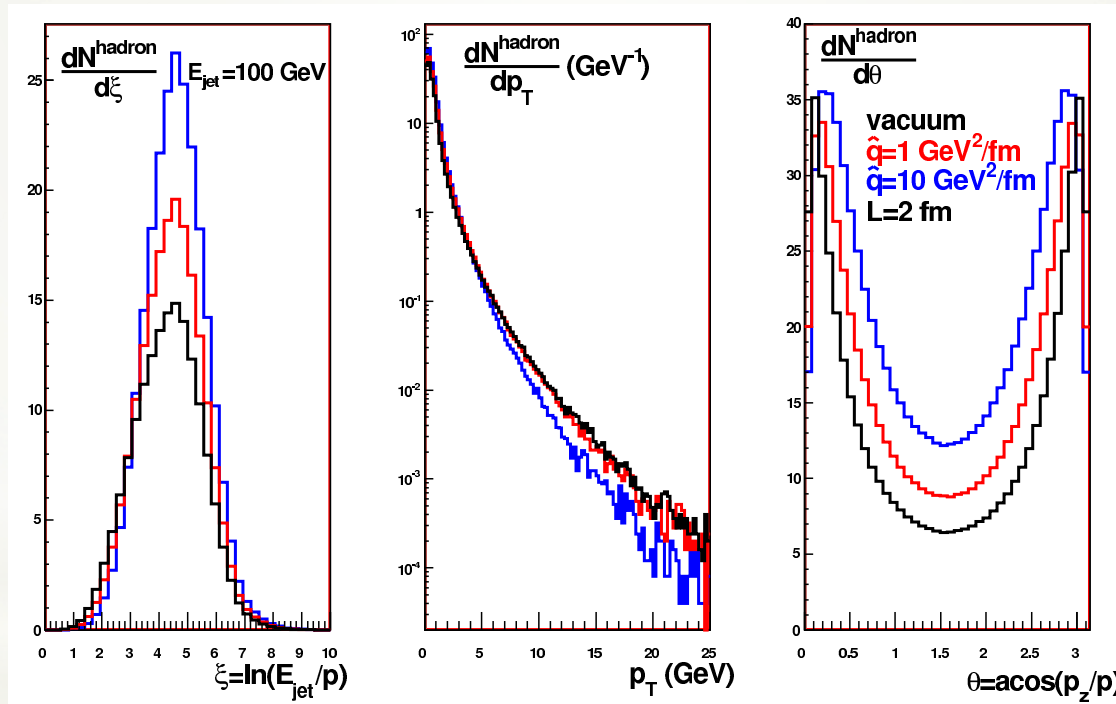
Suppression of high- p_T particles (E conservation) and
enhancement of intermediate- p_T

Broadening of
angular
distribution

[Armesto, Cunqueiro, Salgado]

Results

★ Results at hadron level:

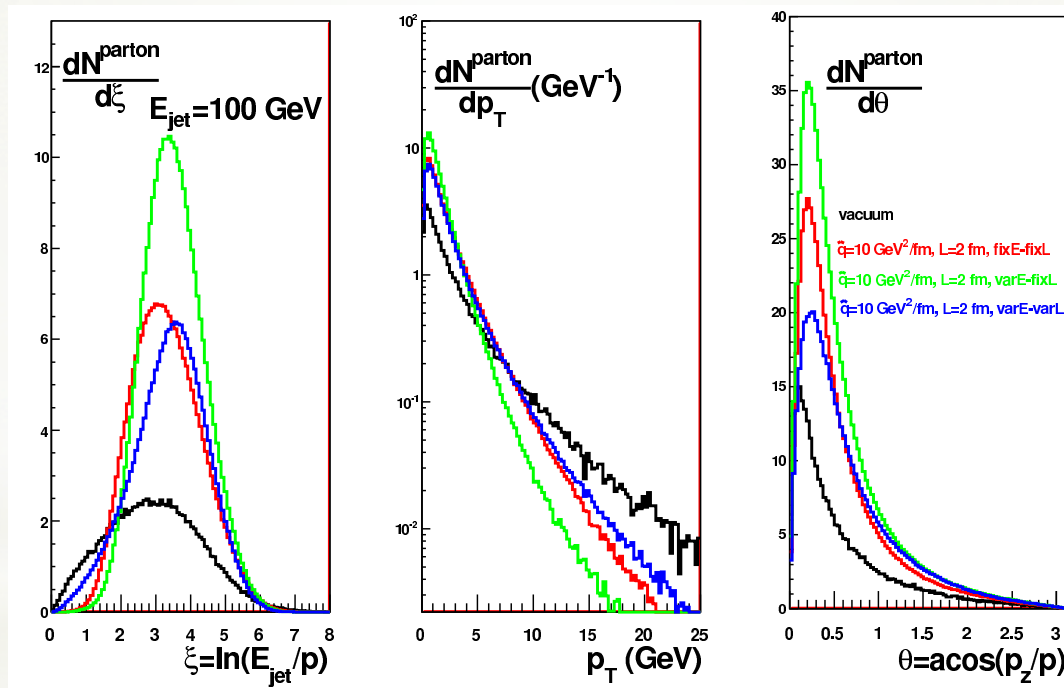


[Armesto, Cunqueiro, Salgado]

Medium effects are diminished by hadronization: low momentum partons do not produce hadrons

Results

- ★ Medium length and energy degradation effects on evolution

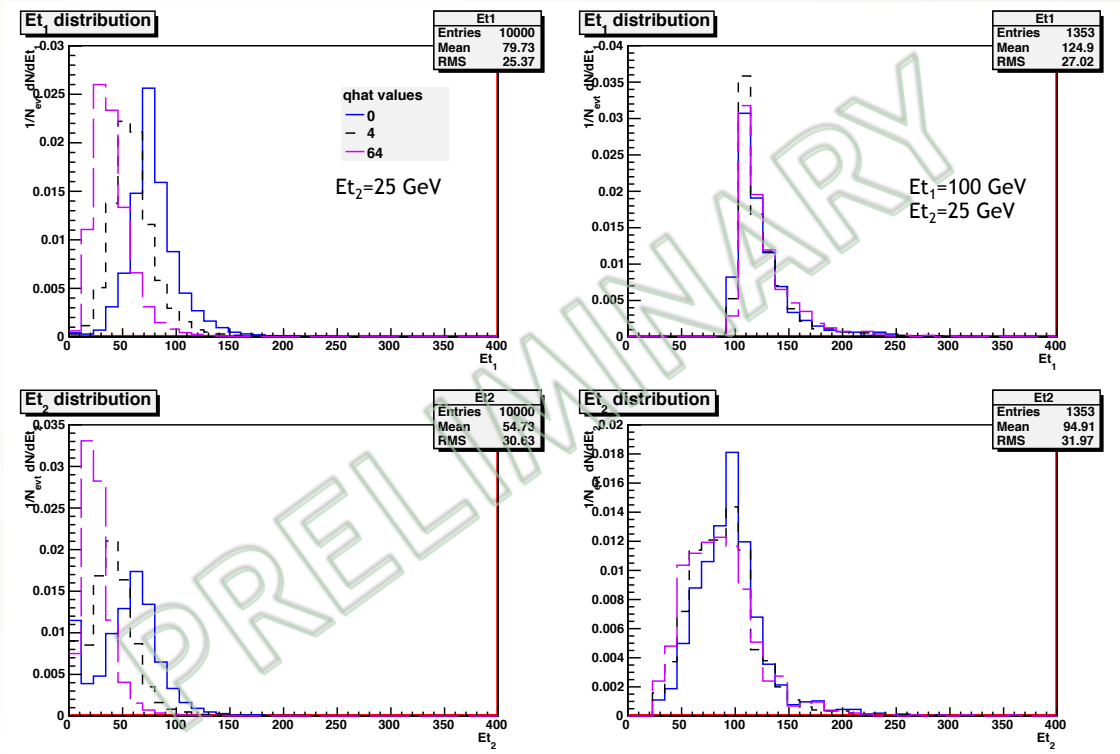


[Armesto, Cunqueiro, Salgado]

Medium effects are largest considering energy degradation but no evolution in length

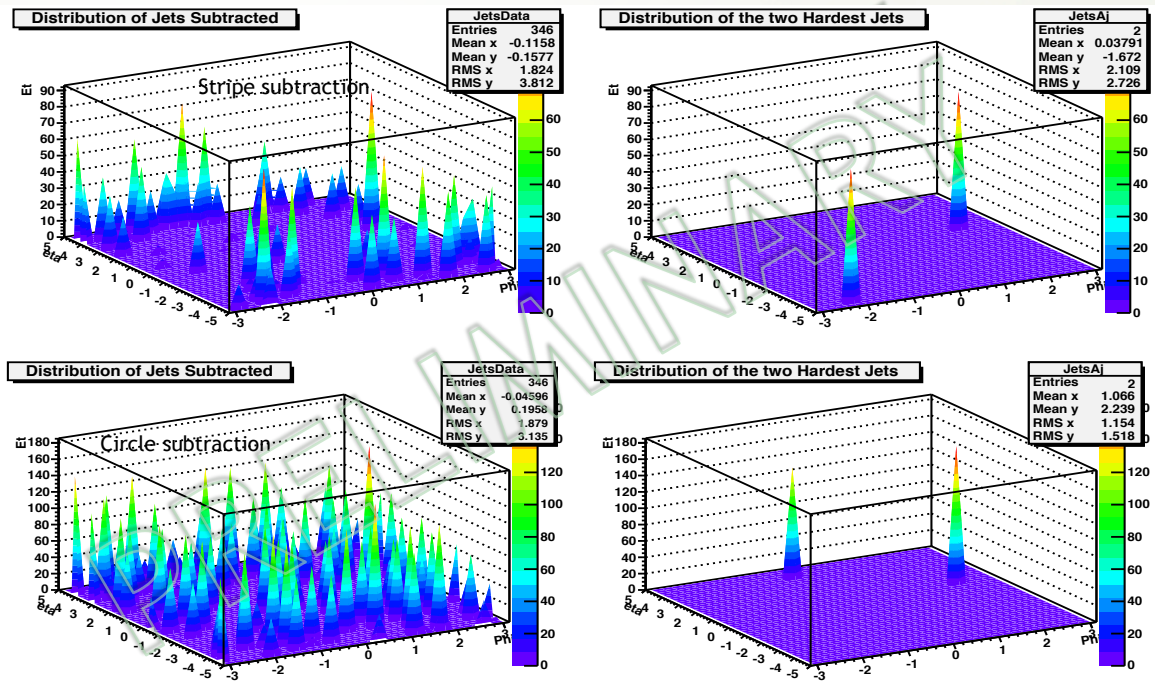
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- ✦ Some on-going studies on jets:
 - ✦ Problems/Difficulties on the theory side:
 - ✦ Cut in p_T and definition of MC truth



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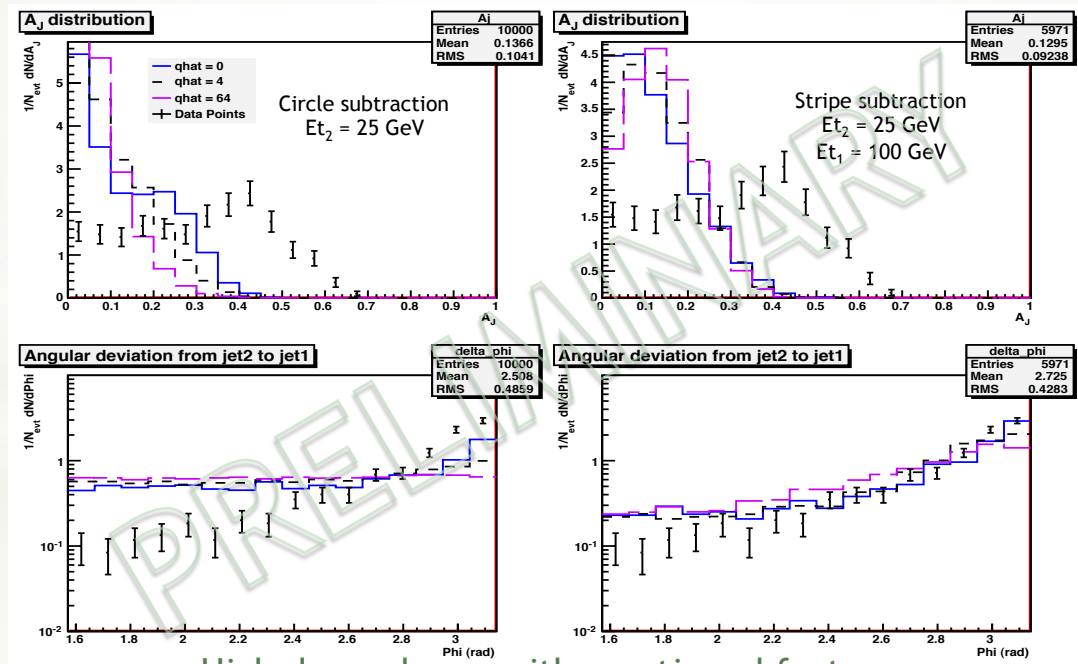
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See also:
Cacciari, Salam, Soyez, 2011
He, Vitev, Zhang, 2011
Lokhtin, Belyaev, Snigirev, 2011



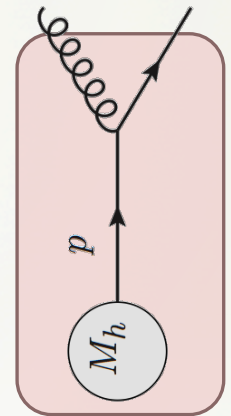
High dependence with mentioned factors

On-going Improvements

- ★ Medium-induced gluon radiation spectrum only for small-x

$$\omega \frac{dI^{med}}{d\omega d^2\mathbf{k}_\perp} \Big|_{x \rightarrow 0} = \frac{\alpha_s C_F}{(2\pi)^2 \omega} 2\text{Re} \left\{ \frac{1}{\omega} \int dy_+ d\bar{y}_+ d\mathbf{x}_\perp e^{-i\mathbf{k}_\perp \cdot \mathbf{x}_\perp} e^{-\frac{1}{2} \int d\xi n(\xi) \sigma(\mathbf{x}_\perp)} \right. \\ \left. \frac{\partial}{\partial \mathbf{y}_\perp} \frac{\partial}{\partial \mathbf{x}_\perp} \mathcal{K}(y_+, \mathbf{y}_\perp = \mathbf{0}; \bar{y}_+, \mathbf{x}_\perp) + \right. \\ \left. + 2 \frac{\mathbf{k}_\perp}{\mathbf{k}_\perp^2} \int dy_+ d\mathbf{x}_\perp e^{-i\mathbf{x}_\perp \cdot \mathbf{k}_\perp} \frac{\partial}{\partial \mathbf{y}_\perp} \mathcal{K}(y_+, \mathbf{y}_\perp = \mathbf{0}; L_+, \mathbf{x}_\perp) \right\}$$

$$k = xp \quad q = (1-x)p$$



On-going Improvements

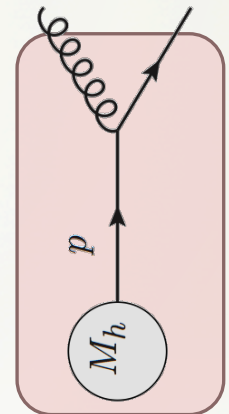
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$$k = xp \quad q = (1-x)p$$



- ★ Medium modified splitting functions extended by comparison with the vacuum ones

- ★ e.g: $q \rightarrow q$

$$x = 1 - z$$

$$P_{vac}(z) \simeq \frac{2C_F}{1-z}$$

Results

$$P_{q \leftarrow q}(z) = 2C_F \frac{1+z^2}{1-z}$$

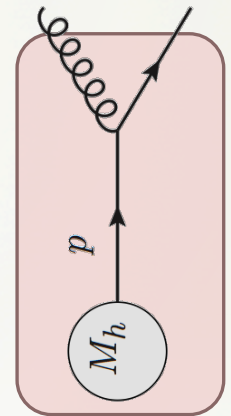
Vacuum SF

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Vacuum SF



$$(1+z^2)P_{vac}$$

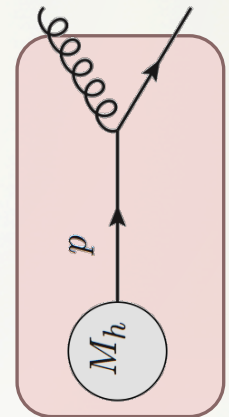
$$(1+z^2)\Delta P$$

On-going Improvements

- ✦ Instead of extending, compute the radiation spectrum in the opposite limit ($x \sim 1$):

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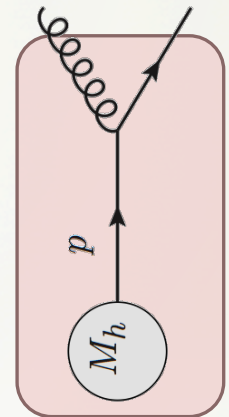
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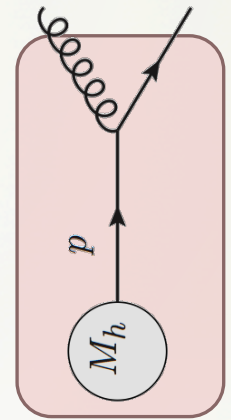


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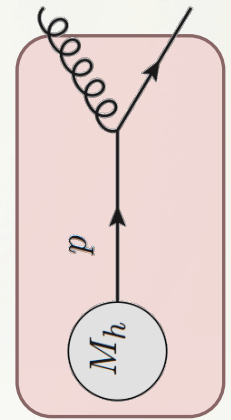
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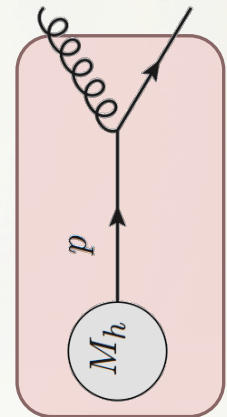
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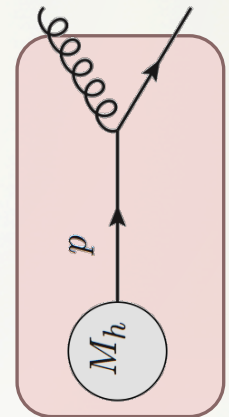
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- ✦ Include corrections in the Q-PYTHIA by matching the two functions

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✦ Q-PYTHIA download:

✦ <http://igfae.usc.es/qatmc/>

END
Thanks!