

RS MODEL EFFECTS ON B_s^0 CP-VIOLATION

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We study the impact of the Randall-Sundrum setup on the width difference $\Delta\Gamma_s$ and the CP-violating phase ϕ_s in the \bar{B}_s^0 - B_s^0 system. We find that the correction to the magnitude of the decay amplitude Γ_{12}^s is below 4% for a realistic choice of input parameters. The main modification in the $\Delta\Gamma_s/\beta_s$ -plane is caused by a new CP-violating phase in the mixing amplitude, which allows for a better agreement with the experimental results of CDF and DØ from $B_s^0 \rightarrow J/\psi\phi$ decays. The best-fit value of the CP asymmetry $S_{\psi\phi}$ can be reproduced, while simultaneously the theoretical prediction for the semileptonic CP asymmetry A_{SL}^s can enter the 1σ range.

1 Introduction

Within the search for new physics (NP) in the decay of B_s^0 -mesons, an important observable is the width difference $\Delta\Gamma_s \equiv \Gamma_L^s - \Gamma_H^s$ between the light and the heavy meson state. According to the above definition, $\Delta\Gamma_s$ happens to be positive in the Standard Model (SM). It can be computed from the dispersive and absorptive part of the \bar{B}_s^0 - B_s^0 mixing amplitude, M_{12}^s and Γ_{12}^s . To leading order in $|\Gamma_{12}^s|/|M_{12}^s|$ one finds the simple relation

$$\Delta\Gamma_s = -\frac{2\text{Re}(M_{12}^s\Gamma_{12}^{s*})}{|M_{12}^s|} = 2|\Gamma_{12}^s| \cos\phi_s. \quad (1)$$

We define the relative phase ϕ_s between the mixing and the decay amplitude according to the convention

$$\frac{M_{12}^s}{\Gamma_{12}^s} = -\frac{|M_{12}^s|}{|\Gamma_{12}^s|} e^{i\phi_s}, \quad \phi_s = \arg(-M_{12}^s\Gamma_{12}^{s*}), \quad (2)$$

for which the SM value is positive and explicitly given by¹ $\phi_s^{\text{SM}} = (4.2 \pm 1.4) \cdot 10^{-3}$. The combined experimental results of CDF and DØ² differ from the SM prediction in the $(\beta_s^{J/\psi\phi}, \Delta\Gamma_s)$ -plane by about 2σ , whereas the latest CDF results disagree by 1σ only³. Here, $\beta_s^{J/\psi\phi} \in [-\pi/2, \pi/2]$ is the CP-violating phase in the interference of mixing and decay, obtained from the time-dependent angular analysis of flavor-tagged $B_s^0 \rightarrow J/\psi\phi$ decays. In the SM it is given by¹ $\beta_s^{J/\psi\phi} = -\arg(-\lambda_t^{bs}/\lambda_c^{bs}) = 0.020 \pm 0.005$, with $\lambda_q^{bs} = V_{qb}V_{qs}^*$. In the presence of NP, $\Delta\Gamma_s$ will be modified^{4,5}. We adopt the notation of ref.⁶ and extend the SM relations according to

$$M_{12}^s = M_{12}^{s\text{SM}} + M_{12}^{s\text{NP}} = M_{12}^{s\text{SM}} R_M e^{i\phi_M}, \quad \Gamma_{12}^s = \Gamma_{12}^{s\text{SM}} + \Gamma_{12}^{s\text{NP}} = \Gamma_{12}^{s\text{SM}} R_\Gamma e^{i\phi_\Gamma}. \quad (3)$$

From (1) it follows that

$$\Delta\Gamma_s = 2|\Gamma_{12}^{s\text{SM}}| R_\Gamma \cos(\phi_s^{\text{SM}} + \phi_M - \phi_\Gamma), \quad (4)$$

where ⁷ $\Delta\Gamma_s^{\text{SM}} = (0.087 \pm 0.021) \text{ ps}^{-1}$. A further important observable is the semileptonic CP asymmetry $A_{\text{SL}}^s = \text{Im}(\Gamma_{12}^s/M_{12}^s)$. Including NP corrections, we find

$$A_{\text{SL}}^s = \frac{|\Gamma_{12}^{s\text{SM}}|}{|M_{12}^{s\text{SM}}|} \frac{R_\Gamma}{R_M} \sin(\phi_s^{\text{SM}} + \phi_M - \phi_\Gamma). \quad (5)$$

Within the SM, the leading contribution to the dispersive part of the \bar{B}_s^0 - B_s^0 mixing amplitude appears at the one loop level. If NP involves flavor-changing neutral currents (FCNCs) at tree level, these give rise to sizable corrections to the mass difference $\Delta m_{B_s} \equiv M_H^s - M_L^s = 2|M_{12}^s|$. Moreover, the presence of tree FCNCs and right-handed charged-current interactions give rise to new decay diagrams. However, the NP corrections to the absorptive part of the amplitude are suppressed by m_W^2/Λ^2 with respect to the SM contribution, where Λ is the NP mass scale. Thus, they are neglected in many NP studies.

2 RS corrections to the \bar{B}_s^0 - B_s^0 system

The Randall-Sundrum (RS) model⁸ is a five-dimensional (5D) quantum field theory (QFT) with an compactified extra-dimension of the order of the Planck length. A ‘‘warped metric’’ is used to generate hierarchies, which are non-understood in the SM. The theory is decomposed into an effective four-dimensional QFT by means of a Kaluza-Klein (KK) decomposition. This gives rise to an infinite tower of heavy copies of the SM particles. The mass scale of the first KK excitations M_{KK} is taken to be a few TeV.

We consider two different scenarios. The first one consists of the SM gauge and matter fields living in the bulk of the 5D space-time, and a Higgs doublet, which is confined to the so-called infra-red boundary of the extra dimension⁹. The second scenario features an extended symmetry group $SU(2)_L \times SU(2)_R \times U(1)_X$ of the electroweak (EW) sector, which is broken down $SU(2)_L \times U(1)_Y$ by the choice of boundary conditions of the respective gauge fields^{10,12,13}. An appropriate embedding of the fermions allows for a protection of $Z^0 b_L \bar{b}_L$ couplings¹¹.

A numerical scan across the ‘‘RS landscape’’ is performed by evaluating M_{12}^s and Γ_{12}^s for appropriate random sets of input parameters, that reproduce the quark masses, mixing angles, and CKM phase. Furthermore, bounds from the $Z^0 b_L \bar{b}_L$ coupling, the oscillation frequency Δm_{B_s} , and the observable ϵ_K , are taken into account. Details of the calculations are given in ref.¹⁵.

3 Numerical analysis

In the first panel of Figure 1 we show the RS corrections to the magnitude and CP-violating phase of the \bar{B}_s^0 - B_s^0 decay width, R_Γ and ϕ_Γ , for a set of 10000 parameter points at $M_{\text{KK}} = 2 \text{ TeV}$. The blue (dark gray) points correspond to the minimal RS model, where we plot only those that are in agreement with the $Z^0 \rightarrow b\bar{b}$ ‘‘pseudo observables’’. The orange (light gray) points correspond to the custodial extension, where the latter bound vanishes. As expected, the RS corrections to $|\Gamma_{12}^s|$ are rather small, typically not exceeding $\pm 4\%$. The corrections to the magnitude and phase of the dispersive part of the mixing amplitude, R_M and ϕ_M , are plotted in the second panel of Figure 1. Here, one should keep in mind the experimental result from the time-dependent measurement of the \bar{B}_s^0 - B_s^0 oscillation frequency¹⁶

$$\Delta m_{B_s}^{\text{exp}} = (17.77 \pm 0.10 \text{ (stat)} \pm 0.07 \text{ (syst)}) \text{ ps}^{-1}, \quad (6)$$

which is in good agreement with the SM prediction⁷ $(17.3 \pm 2.6) \text{ ps}^{-1}$. As a consequence, all points with $R_M \notin [0.718, 1.336]$ are excluded at 95% confidence level, as indicated by the dashed lines. Compared to ϕ_M , the new phase ϕ_Γ can be neglected (what we will do from now on).

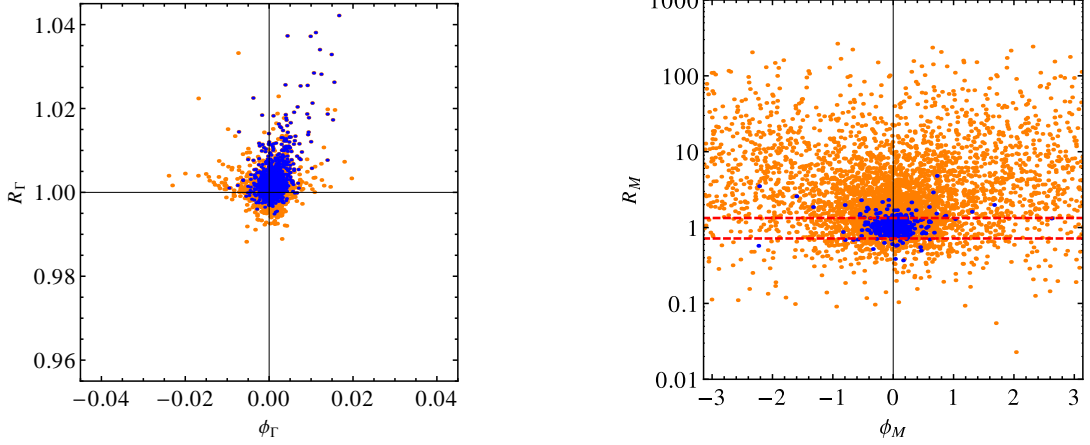


Figure 1: RS corrections to the magnitude and CP-violating phase of the \bar{B}_s^0 - B_s^0 decay amplitude, R_Γ and ϕ_Γ , as well as for the mixing amplitude, R_M and ϕ_M . Blue points correspond to the minimal, orange to the custodial RS model. The red dashed lines mark the 99% confidence region with respect to the measurement of Δm_{B_s} .

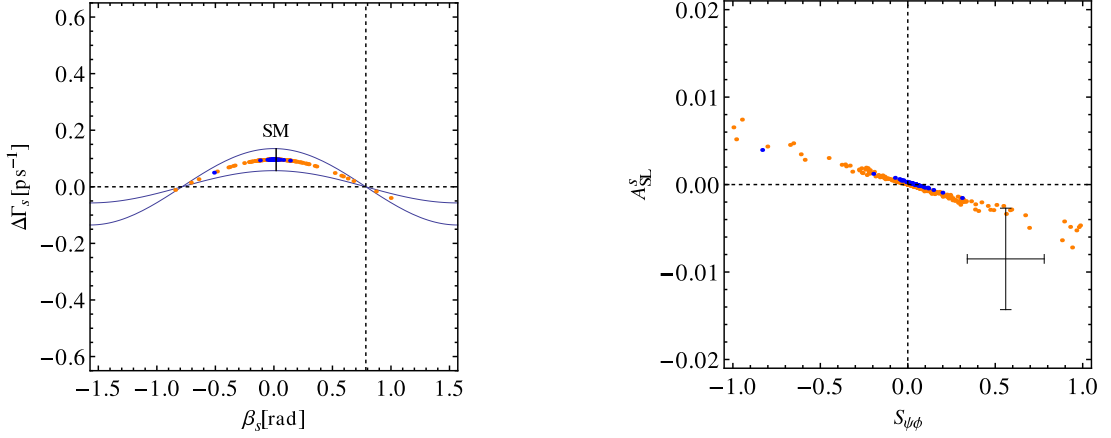


Figure 2: Left panel: Corrections within the $\Delta\Gamma_s^{\text{SM}}/\beta_s$ -plane for the minimal (blue/dark gray) and custodial (orange/light gray) RS model. Bounds from $Z^0 b\bar{b}$, Δm_{B_s} , and ϵ_K are satisfied. Right panel: Corrections within the $A_{\text{SL}}^s/S_{\psi\phi}$ -plane for the minimal and custodial RS model.

Neglecting the small SM phases, the width difference (4) can be written as

$$\Delta\Gamma_s = \Delta\Gamma_s^{\text{SM}} R_\Gamma \cos 2\beta_s, \quad (7)$$

where $2\beta_s \approx -\phi_M^{\text{RS}}$. The preliminary CDF analysis³ uses the older SM prediction¹ $\Delta\Gamma_s^{\text{SM}} = (0.096 \pm 0.039)\text{ps}^{-1}$, which we will take as central value for our calculation. Taking the more recent value will not change our conclusions. The resulting RS predictions for $\Delta\Gamma_s$ are plotted against β_s in the left panel of Figure 2. Comparing to the latest preliminary CDF results³, we conclude that the RS model can enter the 68% confidence region and come close to the best fit value. It stays below the desired value for $\Delta\Gamma_s$, as there are no sizable positive corrections to $|\Gamma_{12}^s|$.

The SM prediction⁷ $(A_{\text{SL}}^s)_{\text{SM}} = (1.9 \pm 0.3) \cdot 10^{-5}$, which is often named a_{sl}^s or a_{fs}^s in the literature, agrees with the direct measurement¹⁷ $(A_{\text{SL}}^s)_{\text{exp}} = -0.0017 \pm 0.0092$ within the (large) error. However, recent measurements of the like-sign dimuon charge asymmetry¹⁸ A_{SL}^b , which

connect A_{SL}^s to its counterpart A_{SL}^d of the B_d^0 -meson sector²⁰, imply a deviation of almost 2σ . If one neglects the tiny SM phases and the NP phase corrections related to decay, A_{SL}^s is proportional to the quantity¹⁹ $S_{\psi\phi}$, which is given by the amplitude of the time-dependent asymmetry in $B_s^0 \rightarrow J/\psi\phi$ decays, $A_{\text{CP}}^s(t) = S_{\psi\phi} \sin(\Delta m_{B_s} t)$. Setting just the NP phase in the decay to zero, one obtains the well known expression²¹ $S_{\psi\phi} = \sin(2\beta_s^{J/\psi\phi} - \phi_M)$, and thus

$$A_{\text{SL}}^s \approx - \frac{|\Gamma_{12}^{s\text{SM}}|}{|M_{12}^{s\text{SM}}|} \frac{R_\Gamma}{R_M} S_{\psi\phi}. \quad (8)$$

The RS result is shown in the right panel of Figure 2, where we have sketched the experimental favored values $S_{\psi\phi} = 0.56 \pm 0.22$ ²² and $A_{\text{SL}}^s = -0.0085 \pm 0.0058$ ¹⁷. The latter number combines the direct measurement with the results derived from the measurement of A_{SL}^b in semileptonic B -decays together with the average $A_{\text{SL}}^d = -0.0047 \pm 0.0046$ from B -factories. It is evident from the plot that the best fit value of $S_{\psi\phi}$ can be reproduced (with some tuning in the minimal RS variant), which has already been noted in ref.¹⁴. Furthermore, the custodial RS model can enter the 1σ range of the measured value of A_{SL}^s . The necessary choice of input parameters is similar to that one, which is suggested by the $\Delta\Gamma_s/\beta_s$ -confidence region.

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