DIMUON CP ASYMMETRY IN B DECAYS

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OUTLINE

- Brief introduction to CP viol' in neutral B_(s) meson systems
 - DØ di-muon charge asymmetry
 - CP viol' in $B_s \rightarrow J/\psi \phi$
- Implications for Physics Beyond the SM
- Conclusions

(See also U. Nierste talk this morning)

EXPERIM. STATUS: THE SM



- SM agrees with data
- CKM unitarity well tested → triangle closes!
- Non-degenerate triangle → CP violation!
- Now look for deviations as sign of New Physics

FCNCs @ 1-loop \rightarrow small in the SM \rightarrow good places to look for NP

decay

CP VIOLATION IN $B_{(s)}$ - $\overline{B}_{(s)}$ SYSTEM



oscillate





• Width difference: absorptive part (long distance)



CP VIOLATION IN $B_{(s)}$ - $\overline{B}_{(s)}$ SYSTEM



B→µ v X, e v X (tree level W[±] exchange, no CP viol')

$$\mathcal{A}_{SL} \equiv \mathcal{A}_{flav.spec.}$$

CP VIOLATION IN $B_{(s)}$ - $\overline{B}_{(s)}$ SYSTEM



Probes CP viol' in the interference between mixing and decay:

$$\mathcal{A}_{f_{CP}}(t) = \frac{d\Gamma/dt \left[\bar{B} \to f_{CP}\right] - d\Gamma/dt \left[B \to f_{CP}\right]}{d\Gamma/dt \left[\bar{B} \to f_{CP}\right] + d\Gamma/dt \left[B \to f_{CP}\right]}$$

 $\mathcal{A}_f(t) = S_f \sin(\Delta m t) - C_f \cos(\Delta m t)$

For B_d use: $J/\psi K_S$ $S_{J/\psi K_S} = sin(2\beta)$ For B_s use: $J/\psi \varphi$ $S_{J/\psi \varphi} = sin(2\beta_s) = -sin(\varphi_s)$

MESON MIXING

• 5 parameters:

 $M_{H,L}^{q} \Gamma_{H,L}^{q} \phi_{q} = \arg(-M_{12}^{q}/\Gamma_{12}^{q})$

• 6 observables:

$$\frac{M_H + M_L}{2} \quad \frac{\Gamma_H + \Gamma_L}{2} \quad \Delta M \quad \Delta \Gamma \quad \mathcal{A}_{SL} \quad \sin(2\beta_q)$$

• 1 relation among them:

$$\mathcal{A}_{SL}^{s} = -\frac{|\Delta\Gamma_{s}|}{\Delta M_{s}} \frac{S_{J/\psi\phi}}{\sqrt{1 - S_{J/\psi\phi}^{2}}}$$

and similarly for B_d

absence of direct CP viol' only assumption, valid even with NP

RECENT MEASUREMENTS: Asl

DØ measures charge asymmetry in muons: $N^{++} - N^{--}$

$$\frac{N_{\mu}^{++} - N_{\mu}^{--}}{N_{\mu}^{++} + N_{\mu}^{--}}$$

- Backgrounds due to in-flight meson decays, matter effects, etc. are removed
- Remainder interpreted as coming from B_{d,s} decays:

 $A^{b}_{SL} \sim 0.5 A^{d}_{SL} + 0.5 A^{s}_{SL}$

no b identification (could be due to something else)



 3.1σ discrepancy (w / A_{SL}^s and A_{SL}^d meas')

Recent Measurements: β_s vs. $\Delta\Gamma_s$

- CDF and DØ measured time dep' CP asym in $B_s \rightarrow J/\psi \phi$
- SM pred': $\beta_s \sim 10^{-2}$
- pre-summer 2010 combination:
 - 2.1σ discrepancy with the SM



Recent Measurements: β_s vs. $\Delta\Gamma_s$

- Both CDF and DØ updated their results in the summer
- Discrepancy with SM is decreased (more on CDF side)
- CDF+DØ combination not available yet (we use D0 2.8fb⁻¹ + CDF 5.2fb⁻¹ φ_s 1D likelihood)



WHAT CAN IT BE?

- Both $A_{SL}{}^{b}$ and $S_{\psi\varphi}$ may be statistical fluctuations...
- A_{SL}^b may be something totally different (new) not related to B physics...
- A_{SL}^{b} and $S_{\psi\varphi}$ may be due to NP in B physics

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Compatibility?

CONSISTENCY CHECK

• From A_{SL}^{b} , A_{SL}^{d} , ΔM_{s} one gets:

$$|\Delta \Gamma_s| \sim \left[(0.28 \pm 0.15) \text{ps}^{-1} \right] \sqrt{1 - S_{\psi\phi}^2} / S_{\psi\phi}$$

- To be compared with the best fit point of $(|\Delta\Gamma_s|, S_{\psi\phi}) \sim (0.15 \text{ps}^{-1}, 0.5)$
- Also consistent if using new CDF+D0 data

Non trivial test for NP! → assume the discrepancies are due to NP and study consequences

Modify only the mixing amplitude:

$$\frac{M_{12}^q}{M_{12}^{\text{SM},q}} = 1 + \frac{M_{12}^{\text{NP},q}}{M_{12}^{\text{SM},q}} = 1 + h_q \,\mathrm{e}^{\mathrm{i}2\sigma_q}$$

In the fit: other 4 real params on top of ϱ , η from the Unit. Triang

$$\Delta m_q = \Delta m_q^{\text{SM}} \left| 1 + h_q e^{2i\sigma_q} \right|,$$

$$\Delta \Gamma_s = \Delta \Gamma_s^{\text{SM}} \cos \left[\arg \left(1 + h_s e^{2i\sigma_s} \right) \right],$$

$$A_{\text{SL}}^q = \text{Im} \left\{ \Gamma_{12}^q / \left[M_{12}^{q,\text{SM}} (1 + h_q e^{2i\sigma_q}) \right] \right\},$$

$$S_{\psi K} = \sin \left[2\beta + \arg \left(1 + h_d e^{2i\sigma_d} \right) \right],$$

$$S_{\psi \phi} = \sin \left[2\beta_s - \arg \left(1 + h_s e^{2i\sigma_s} \right) \right].$$

(CKMFitter package used)

U.4

0.3

0.2

0.1

0.0

Loops are contaminated from NP \rightarrow use tree level observables to determine ϱ , η :





- Favored regions are for $h_d{\sim}0.25~\sigma_d{\sim}110^\circ$ and $h_s{\sim}0.5~\sigma_s{\sim}120^\circ$ or $h_s{\sim}1.8~\sigma_s{\sim}100^\circ$
- σ_s and σ_d are very close \rightarrow same origin?
- $h_s=0$ is disfavored at 2.6 σ , $h_d=0$ at $<2\sigma$

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- h_s and h_d not strongly correlated
- $h_s=h_d=0$ (SM) disfavored at 3.3σ
- $h_s = h_d$ is not ruled out but data prefer $h_s \gtrsim h_d$



NEW PHYSICS MODELS?

- h_{s,d}, σ_{s,d} is a general parameterization.
 Which types of NP models can account for this?
- Are new sources of flavor viol' needed?
- Are new sources of CP viol' needed?
- How to get $h_s \gtrsim h_d$?

The smallest perturbation of the SM picture:

MINIMAL FLAVOR VIOLATION (MFV)

- New particles & new interactions? Yes
- New Flavor viol' sources at *low energy*? No

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_{i} c_i(y_u, y_d) \frac{\mathcal{O}_i}{\Lambda^2}$$

FLAVOR & CP VIOL' IN (G)MFV

 What about new CP phases? → Can be flavorblind phases (e.g. in SUSY the phases coming from gaugino masses, etc.)

E.g.:
$$c(y_u^{\dagger}y_u, y_d^{\dagger}y_d) \frac{(\bar{b}_L \gamma^{\mu} s_L)(\bar{b}_L \gamma^{\mu} s_L)}{\Lambda^2}$$

$$c(y_{u}^{\dagger}y_{u}, y_{d}^{\dagger}y_{d}) = a_{1}(y_{u}^{\dagger}y_{u})_{32}^{2} + a_{2}(y_{u}^{\dagger}y_{u})_{32}(y_{u}^{\dagger}y_{u}y_{d}^{\dagger}y_{d})_{32} + a_{3}(y_{u}^{\dagger}y_{u})_{32}(y_{d}^{\dagger}y_{d}y_{u}^{\dagger}y_{u})_{32} + \dots$$

Some of the a_i can be complex \rightarrow new sources of CPV $\Lambda \sim$ few hundreds GeV \div few TeV to account for the data (depends on y_b)

FLAVOR & CP VIOL' IN MFV

- What about $h_s \gg h_d$?
- $m_s \gg m_d \rightarrow \text{down yukawa coupling can}$ do it! (that's the only way in MFV...)
- Requires one Y_d insertion:
- "scalar" operator involving b_Ls_R, b_Ld_R (similar to an Higgs exchange):

E.g.: $\left[\bar{Q}_3(A_d^m A_u^n Y_d)_{3i} d_i\right] \left[\bar{d}_3(Y_d^{\dagger} A_d^{l,\dagger} A_u^{p,\dagger})_{3i} Q_i\right]$ $A_{u,d} \equiv Y_{u,d} Y_{u,d}^{\dagger}$

 $h_d/h_s \propto m_d/m_s \sim O(5\%) \rightarrow$ testable prediction

NEW PHYSICS MODELS?

CAVEAT: flavor blind phases → potential problems with EDMs

Electric Dipole Moments are generated by different operators, but in specific models the size of the contrib' to EDMs and to CPV in B_s mixing is correlated

(model building gymnastics required, various working examples in the literature)

FLAVOR & CP VIOL' BEYOND SM

- What about non-MFV models?
 - New sources of flavor viol' and new flavordependent CP phases
 - Can account easily for the data BUT if new sources generate h_s ≠ h_d then constraints from Kaon and D-meson mixing may be non-trivial (Model building needed...)

Problem with $\Delta\Gamma_s$?

- The SM prediction for $\Delta\Gamma_s$ is in some tension with the DØ (slightly more than $1\sigma...$)
- The data prefers a ΔΓ_s larger by a factor ~2-2.5 (shallow profile, new CDF+DØ data should lower the "tension" to a factor of 1.x)
- Could be a fluctuation in the data
- Could be a problem with the SM prediction
- It's very hard to account a shift in ΔΓ_s with new physics but possible (e.g. need a new light state)
 (Bai, Nelson 2010)

CONCLUSIONS

- The dimuon CP asymmetry may be due to BSM physics finally showing up in the flavor sector.
- Very easy to account for the data (→ quite generic)
- Data does not exclude MFV (at least its general definition)
- More exp' input needed → available later this year!

ON THE LHC FEEDBACK

"...most important of all, Switzerland is now the home of the Largest Superconducting Supercollider which was moved from Texas to Switzerland after the '94 budget, but we still have smaller ones in the United States. The first significant discovery came [...] last year. They may offer a key on how life began after the Big Bang. [...] The smallest subatomic particle, the muon, it turns out, that based on their findings, which will be confirmed or contradicted when the Swiss machine is up and going, [...] [It] turns out that there's slightly more positive than negative muons in all of our atoms which will justify the faith of all the believers of the world, make you more optimistic and gives us an explanation for how we might have all come to this moment from the primordial slack."

William J. Clinton, World Economic Forum, Davos, 2011