

# RS MODEL EFFECTS ON $B_s^0$ CP-VIOLATION

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We study the impact of the Randall-Sundrum setup on the width difference  $\Delta\Gamma_s$  and the CP-violating phase  $\phi_s$  in the  $\bar{B}_s^0$ - $B_s^0$  system. We find that the correction to the magnitude of the decay amplitude  $\Gamma_{12}^s$  is below 4% for a realistic choice of input parameters. The main modification in the  $\Delta\Gamma_s/\beta_s$ -plane is caused by a new CP-violating phase in the mixing amplitude, which allows for a better agreement with the experimental results of CDF and DØ from  $B_s^0 \rightarrow J/\psi\phi$  decays. The best-fit value of the CP asymmetry  $S_{\psi\phi}$  can be reproduced, while simultaneously the theoretical prediction for the semileptonic CP asymmetry  $A_{\text{SL}}^s$  can enter the  $1\sigma$  range.

## 1 Introduction

Within the search for new physics (NP) in the decay of  $B_s^0$ -mesons, an important observable is the width difference  $\Delta\Gamma_s \equiv \Gamma_L^s - \Gamma_H^s$  between the light and the heavy meson state. According to the above definition,  $\Delta\Gamma_s$  happens to be positive in the Standard Model (SM). It can be computed from the dispersive and absorptive part of the  $\bar{B}_s^0$ - $B_s^0$  mixing amplitude,  $M_{12}^s$  and  $\Gamma_{12}^s$ . To leading order in  $|\Gamma_{12}^s|/|M_{12}^s|$  one finds the simple relation

$$\Delta\Gamma_s = -\frac{2\text{Re}(M_{12}^s\Gamma_{12}^{s*})}{|M_{12}^s|} = 2|\Gamma_{12}^s| \cos\phi_s. \quad (1)$$

We define the relative phase  $\phi_s$  between the mixing and the decay amplitude according to the convention

$$\frac{M_{12}^s}{\Gamma_{12}^s} = -\frac{|M_{12}^s|}{|\Gamma_{12}^s|} e^{i\phi_s}, \quad \phi_s = \arg(-M_{12}^s\Gamma_{12}^{s*}), \quad (2)$$

for which the SM value is positive and explicitly given by<sup>1</sup>  $\phi_s^{\text{SM}} = (4.2 \pm 1.4) \cdot 10^{-3}$ . The combined experimental results of CDF and DØ<sup>2</sup> differ from the SM prediction in the  $(\beta_s^{J/\psi\phi}, \Delta\Gamma_s)$ -plane by about  $2\sigma$ , whereas the latest CDF results disagree by  $1\sigma$  only<sup>3</sup>. Here,  $\beta_s^{J/\psi\phi} \in [-\pi/2, \pi/2]$  is the CP-violating phase in the interference of mixing and decay, obtained from the time-dependent angular analysis of flavor-tagged  $B_s^0 \rightarrow J/\psi\phi$  decays. In the SM it is given by<sup>1</sup>  $\beta_s^{J/\psi\phi} = -\arg(-\lambda_t^{bs}/\lambda_c^{bs}) = 0.020 \pm 0.005$ , with  $\lambda_q^{bs} = V_{qb}V_{qs}^*$ . In the presence of NP,  $\Delta\Gamma_s$  will be modified<sup>4,5</sup>. We adopt the notation of ref.<sup>6</sup> and extend the SM relations according to

$$M_{12}^s = M_{12}^{s\text{SM}} + M_{12}^{s\text{NP}} = M_{12}^{s\text{SM}} R_M e^{i\phi_M}, \quad \Gamma_{12}^s = \Gamma_{12}^{s\text{SM}} + \Gamma_{12}^{s\text{NP}} = \Gamma_{12}^{s\text{SM}} R_\Gamma e^{i\phi_\Gamma}. \quad (3)$$

From (1) it follows that

$$\Delta\Gamma_s = 2|\Gamma_{12}^{s\text{SM}}| R_\Gamma \cos(\phi_s^{\text{SM}} + \phi_M - \phi_\Gamma), \quad (4)$$

where <sup>7</sup>  $\Delta\Gamma_s^{\text{SM}} = (0.087 \pm 0.021) \text{ ps}^{-1}$ . A further important observable is the semileptonic CP asymmetry  $A_{\text{SL}}^s = \text{Im}(\Gamma_{12}^s/M_{12}^s)$ . Including NP corrections, we find

$$A_{\text{SL}}^s = \frac{|\Gamma_{12}^{s\text{SM}}|}{|M_{12}^{s\text{SM}}|} \frac{R_\Gamma}{R_M} \sin(\phi_s^{\text{SM}} + \phi_M - \phi_\Gamma). \quad (5)$$

Within the SM, the leading contribution to the dispersive part of the  $\bar{B}_s^0$ - $B_s^0$  mixing amplitude appears at the one loop level. If NP involves flavor-changing neutral currents (FCNCs) at tree level, these give rise to sizable corrections to the mass difference  $\Delta m_{B_s} \equiv M_H^s - M_L^s = 2|M_{12}^s|$ . Moreover, the presence of tree FCNCs and right-handed charged-current interactions give rise to new decay diagrams. However, the NP corrections to the absorptive part of the amplitude are suppressed by  $m_W^2/\Lambda^2$  with respect to the SM contribution, where  $\Lambda$  is the NP mass scale. Thus, they are neglected in many NP studies.

## 2 RS corrections to the $\bar{B}_s^0$ - $B_s^0$ system

The Randall-Sundrum (RS) model<sup>8</sup> is a five-dimensional (5D) quantum field theory (QFT) with an compactified extra-dimension of the order of the Planck length. A ‘‘warped metric’’ is used to generate hierarchies, which are non-understood in the SM. The theory is decomposed into an effective four-dimensional QFT by means of a Kaluza-Klein (KK) decomposition. This gives rise to an infinite tower of heavy copies of the SM particles. The mass scale of the first KK excitations  $M_{\text{KK}}$  is taken to be a few TeV.

We consider two different scenarios. The first one consists of the SM gauge and matter fields living in the bulk of the 5D space-time, and a Higgs doublet, which is confined to the so-called infra-red boundary of the extra dimension<sup>9</sup>. The second scenario features an extended symmetry group  $SU(2)_L \times SU(2)_R \times U(1)_X$  of the electroweak (EW) sector, which is broken down  $SU(2)_L \times U(1)_Y$  by the choice of boundary conditions of the respective gauge fields<sup>10,12,13</sup>. An appropriate embedding of the fermions allows for a protection of  $Z^0 b_L \bar{b}_L$  couplings<sup>11</sup>.

A numerical scan across the ‘‘RS landscape’’ is performed by evaluating  $M_{12}^s$  and  $\Gamma_{12}^s$  for appropriate random sets of input parameters, that reproduce the quark masses, mixing angles, and CKM phase. Furthermore, bounds from the  $Z^0 b_L \bar{b}_L$  coupling, the oscillation frequency  $\Delta m_{B_s}$ , and the observable  $\epsilon_K$ , are taken into account. Details of the calculations are given in ref.<sup>15</sup>.

## 3 Numerical analysis

In the first panel of Figure 1 we show the RS corrections to the magnitude and CP-violating phase of the  $\bar{B}_s^0$ - $B_s^0$  decay width,  $R_\Gamma$  and  $\phi_\Gamma$ , for a set of 10000 parameter points at  $M_{\text{KK}} = 2 \text{ TeV}$ . The blue (dark gray) points correspond to the minimal RS model, where we plot only those that are in agreement with the  $Z^0 \rightarrow b\bar{b}$  ‘‘pseudo observables’’. The orange (light gray) points correspond to the custodial extension, where the latter bound vanishes. As expected, the RS corrections to  $|\Gamma_{12}^s|$  are rather small, typically not exceeding  $\pm 4\%$ . The corrections to the magnitude and phase of the dispersive part of the mixing amplitude,  $R_M$  and  $\phi_M$ , are plotted in the second panel of Figure 1. Here, one should keep in mind the experimental result from the time-dependent measurement of the  $\bar{B}_s^0$ - $B_s^0$  oscillation frequency<sup>16</sup>

$$\Delta m_{B_s}^{\text{exp}} = (17.77 \pm 0.10 \text{ (stat)} \pm 0.07 \text{ (syst)}) \text{ ps}^{-1}, \quad (6)$$

which is in good agreement with the SM prediction<sup>7</sup>  $(17.3 \pm 2.6) \text{ ps}^{-1}$ . As a consequence, all points with  $R_M \notin [0.718, 1.336]$  are excluded at 95% confidence level, as indicated by the dashed lines. Compared to  $\phi_M$ , the new phase  $\phi_\Gamma$  can be neglected (what we will do from now on).

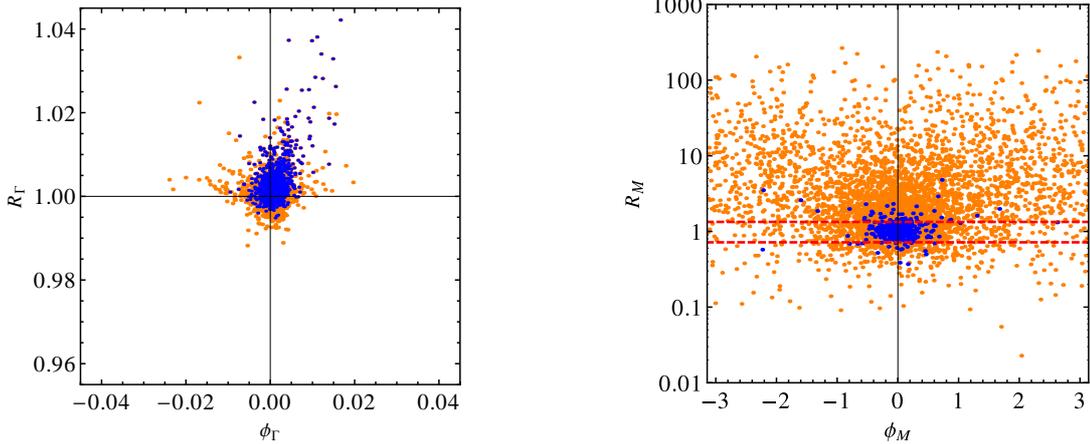


Figure 1: RS corrections to the magnitude and CP-violating phase of the  $\bar{B}_s^0$ - $B_s^0$  decay amplitude,  $R_\Gamma$  and  $\phi_\Gamma$ , as well as for the mixing amplitude,  $R_M$  and  $\phi_M$ . Blue points correspond to the minimal, orange to the custodial RS model. The red dashed lines mark the 99% confidence region with respect to the measurement of  $\Delta m_{B_s}$ .

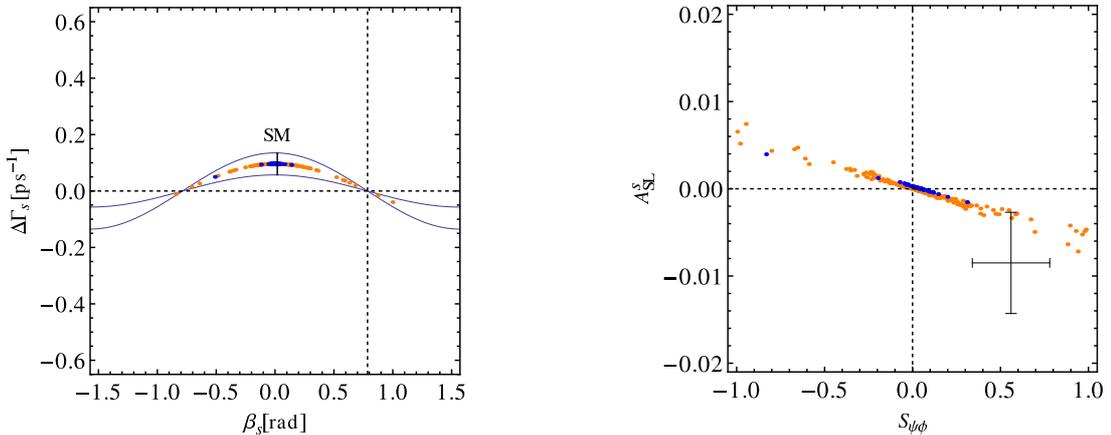


Figure 2: Left panel: Corrections within the  $\Delta\Gamma_s^{\text{SM}}/\beta_s$ -plane for the minimal (blue/dark gray) and custodial (orange/light gray) RS model. Bounds from  $Z^0 b\bar{b}$ ,  $\Delta m_{B_s}$ , and  $\epsilon_K$  are satisfied. Right panel: Corrections within the  $A_{\text{SL}}^s/S_{\psi\phi}$ -plane for the minimal and custodial RS model.

Neglecting the small SM phases, the width difference (4) can be written as

$$\Delta\Gamma_s = \Delta\Gamma_s^{\text{SM}} R_\Gamma \cos 2\beta_s, \quad (7)$$

where  $2\beta_s \approx -\phi_M^{\text{RS}}$ . The preliminary CDF analysis<sup>3</sup> uses the older SM prediction<sup>1</sup>  $\Delta\Gamma_s^{\text{SM}} = (0.096 \pm 0.039)\text{ps}^{-1}$ , which we will take as central value for our calculation. Taking the more recent value will not change our conclusions. The resulting RS predictions for  $\Delta\Gamma_s$  are plotted against  $\beta_s$  in the left panel of Figure 2. Comparing to the latest preliminary CDF results<sup>3</sup>, we conclude that the RS model can enter the 68% confidence region and come close to the best fit value. It stays below the desired value for  $\Delta\Gamma_s$ , as there are no sizable positive corrections to  $|\Gamma_{12}^s|$ .

The SM prediction<sup>7</sup>  $(A_{\text{SL}}^s)_{\text{SM}} = (1.9 \pm 0.3) \cdot 10^{-5}$ , which is often named  $a_{\text{sl}}^s$  or  $a_{\text{fs}}^s$  in the literature, agrees with the direct measurement<sup>17</sup>  $(A_{\text{SL}}^s)_{\text{exp}} = -0.0017 \pm 0.0092$  within the (large) error. However, recent measurements of the like-sign dimuon charge asymmetry<sup>18</sup>  $A_{\text{SL}}^b$ , which

connect  $A_{\text{SL}}^s$  to its counterpart  $A_{\text{SL}}^d$  of the  $B_d^0$ -meson sector<sup>20</sup>, imply a deviation of almost  $2\sigma$ . If one neglects the tiny SM phases and the NP phase corrections related to decay,  $A_{\text{SL}}^s$  is proportional to the quantity<sup>19</sup>  $S_{\psi\phi}$ , which is given by the amplitude of the time-dependent asymmetry in  $B_s^0 \rightarrow J/\psi\phi$  decays,  $A_{\text{CP}}^s(t) = S_{\psi\phi} \sin(\Delta m_{B_s} t)$ . Setting just the NP phase in the decay to zero, one obtains the well known expression<sup>21</sup>  $S_{\psi\phi} = \sin(2\beta_s^{J/\psi\phi} - \phi_M)$ , and thus

$$A_{\text{SL}}^s \approx - \frac{|\Gamma_{12}^{s\text{SM}}|}{|M_{12}^{s\text{SM}}|} \frac{R_\Gamma}{R_M} S_{\psi\phi}. \quad (8)$$

The RS result is shown in the right panel of Figure 2, where we have sketched the experimental favored values  $S_{\psi\phi} = 0.56 \pm 0.22$ <sup>22</sup> and  $A_{\text{SL}}^s = -0.0085 \pm 0.0058$ <sup>17</sup>. The latter number combines the direct measurement with the results derived from the measurement of  $A_{\text{SL}}^b$  in semileptonic  $B$ -decays together with the average  $A_{\text{SL}}^d = -0.0047 \pm 0.0046$  from  $B$ -factories. It is evident from the plot that the best fit value of  $S_{\psi\phi}$  can be reproduced (with some tuning in the minimal RS variant), which has already been noted in ref.<sup>14</sup>. Furthermore, the custodial RS model can enter the  $1\sigma$  range of the measured value of  $A_{\text{SL}}^s$ . The necessary choice of input parameters is similar to that one, which is suggested by the  $\Delta\Gamma_s/\beta_s$ -confidence region.

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