Recent topics on Electroweak Symmetry Breaking

(Alternatives to the MSSM)

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## In most people's mind:

- Are there really serious alternatives to the MSSM?
- The MSSM is the perfect candidate for physics beyond the SM

It has become the orthodoxy

Relies on the existence of a fundamental Higgs + symmetry to keep it stable

My role here: Devil's Advocate

Defend the Alternatives:

- I) No Higgs (such as Technicolor = a copy of QCD at the TeV)
- 2) Composite Higgs (Pseudo-Goldstone Higgs = similar to Kaons in QCD)

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SUSY 98: TC was stabbed twice (from Electroweak precision tests)



## Veneziano summary talk:

To conclude, the score on precision tests puts the SSM first, with the SM itself (with a light Higgs and some additional intermediate scale) a close second. <u>Technicolour theories</u> appear to lag far behind and... there is not much else in the race. I would conclude that, if The MSSM gained its present status after LEP I, where it left behind its main competitors (e.g. technicolor)



But after LEP I, it came LEP II and Tevatron...

In the MSSM the Higgs or the sparticles were expected to be seen! (nothing expected from the alternatives) But after LEP I, it came LEP II and Tevatron...

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But Nothing was seen!

But after LEP I, it came LEP II and Tevatron...

In the MSSM the Higgs or the sparticles were expected to be **seen!** (nothing expected from the alternatives)

**Example of how** much is ruled out after LEP/Tevatron (MSSM with universal masses)



It is however clear that susy has the bonus of being a predictable theory... not the case for  $TC \approx QCD$ 

Recent progress: explicit weakly-coupled examples

• Extra dimensional Higgsless models

Csaki,Grojena,Pilo,Terning

• Little Higgs

Arkani-Hamed, Cohen, Katz, Nelson

 Holographic Higgs: Extra dimensional Composite Higgs models

Agashe, Contino, AP



Predictive models!

### Still their main obstacle is the S-T ellipse

### To any alternative model, one must ask "Where it is in the ellipse?"



### Minimal 5D Higgsless theory



Csaki,Grojena,Pilo,Terning

## Minimal 5D Higgsless theory



## Reason for generic large effects on S:

In 5D WW-scattering unitarized by KK-resonances:



This KK-states modify the gauge boson propagators:

$$W^{(n)}$$
  
 $W \longrightarrow B$ 

Similar situation to Technicolor theories!

 $W^{(n)}$  = techni-rho

## Better situation for the second alternative: composite PGB Higgs:

**WW-scattering unitarized by a Higgs + KK-resonances:** 



## Minimal model of a 5D composite Higgs



Agashe, Contino, A.P.

Why this symmetry breaking pattern?

We are in 5D: 
$$A_M = (A_\mu, A_5)$$

Massless boson spectrum:

- $A_{\mu}$  of  $SU(2)_L \otimes U(1)_Y = SM$  Gauge bosons
- A<sub>5</sub> of SO(5)/SO(4) = 2 of SU(2)<sub>L</sub> = SM Higgs

$$\hookrightarrow$$
 **Higgs-gauge unification**

Hosotani mechanism

Higgs mass protected by 5D gauge invariance!

$$A_5 \rightarrow A_5 + \partial_5 \theta$$
 shifts as a PGB

# Minimal model of a 5D composite Higgs



## Little Higgs with T-parity

(as in susy virtual effects at one-loop)



Keeping the analogy of Veneziano, the race now looks more like this...

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As in susy, I consider worthy to look for these alternatives at the LHC

# What to expect at the LHC and, maybe, ILC?

I) Extra resonances around TeV with SM quantum numbers: W', Z', t', b', ...

2) Non-elementary Higgs: Its couplings will differ from the SM Higgs

# 1) Direct searches:

New resonances







Possible to see up to 2-3 TeV





feasible to see up to 1-2 TeV

# 2) Indirect searches:

Modifications of the Higgs properties

(... in the case it is there)

### Model independent approach:

Find the effective theory after integrating out the heavy states:

 $\mathcal{L}_{SM+H}$  + higher dimensional operators

(the equivalent of the pion chiral lagrangian in QCD)

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Find the effective theory after integrating out the heavy states:



#### **DIMENSION-6 OPERATORS**

Suppressed by f (the analog of  $f_{\pi}$  in QCD):

$$\frac{c_{H}}{2f^{2}}\partial^{\mu}\left(H^{\dagger}H\right)\partial_{\mu}\left(H^{\dagger}H\right) + \frac{c_{T}}{2f^{2}}\left(H^{\dagger}\overline{D^{\mu}}H\right)\left(H^{\dagger}\overline{D^{\mu}}H\right)$$
$$-\frac{c_{6}\lambda}{f^{2}}\left(H^{\dagger}H\right)^{3} + \left(\frac{c_{y}y_{f}}{f^{2}}H^{\dagger}H\bar{f}_{L}Hf_{R} + \text{h.c.}\right)$$

 $c_H, c_T, c_6, c_y$ : model-dependent coefficients

DIMENSION-6 OPERATORS  
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$$tested at LEP:$$
  
 $T-parameter c_T = 0 ext{ if the BSM sector is custodial invariant}}$   
Suppressed by  $f$  (the analog of  $f_{\pi}$  in QCD):  
 $\frac{c_H}{2f^2} \partial^{\mu} (H^{\dagger}H) \partial_{\mu} (H^{\dagger}H) + \frac{c_T}{2f^2} \left(H^{\dagger}\widehat{D^{\mu}}H\right) \left(H^{\dagger}\widehat{D}_{\mu}H\right)$   
 $-\frac{c_6\lambda}{f^2} (H^{\dagger}H)^3 + \left(\frac{c_y y_f}{f^2}H^{\dagger}H\bar{f}_LHf_R + ext{h.c.}\right)$ 

The rest, not tested yet!

 $c_H, c_T, c_6, c_y$ : model-dependent coefficients

Measuring the compositeness of the Higgs:



### Definite modifications of Higgs decay widths:

$$\begin{split} \Gamma\left(h \to f\bar{f}\right)_{\rm SILH} &= \Gamma\left(h \to f\bar{f}\right)_{\rm SM} \left[1 - \frac{\xi}{\xi} (2c_y + c_H)\right] \\ \Gamma\left(h \to W^+W^-\right)_{\rm SILH} &= \Gamma\left(h \to W^+W^{(*)-}\right)_{\rm SM} \left[1 - \frac{\xi}{\xi} \left(c_H - \frac{g^2}{g_\rho^2} \hat{c}_W\right)\right] \\ \Gamma\left(h \to ZZ\right)_{\rm SILH} &= \Gamma\left(h \to ZZ^{(*)}\right)_{\rm SM} \left[1 - \frac{\xi}{\xi} \left(c_H - \frac{g^2}{g_\rho^2} \hat{c}_Z\right)\right] \\ \Gamma\left(h \to gg\right)_{\rm SILH} &= \Gamma\left(h \to gg\right)_{\rm SM} \left[1 - \frac{\xi}{\xi} \operatorname{Re}\left(2c_y + c_H + \frac{4y_t^2 c_g}{g_\rho^2 I_g}\right)\right] \\ \Gamma\left(h \to \gamma\gamma\right)_{\rm SILH} &= \Gamma\left(h \to \gamma\gamma\right)_{\rm SM} \left[1 - \frac{\xi}{\xi} \operatorname{Re}\left(\frac{2c_y + c_H}{1 + J_\gamma/I_\gamma} + \frac{c_H - \frac{g^2}{g_\rho^2} \hat{c}_W}{1 + I_\gamma/J_\gamma} + \frac{\frac{4g^2}{g_\rho^2} c_\gamma}{I_\gamma + J_\gamma}\right)\right] \\ \Gamma\left(h \to \gamma Z\right)_{\rm SILH} &= \Gamma\left(h \to \gamma Z\right)_{\rm SM} \left[1 - \frac{\xi}{\xi} \operatorname{Re}\left(\frac{2c_y + c_H}{1 + J_\gamma/I_\gamma} + \frac{c_H - \frac{g^2}{g_\rho^2} \hat{c}_W}{1 + I_Z/J_Z} + \frac{4c_{\gamma Z}}{I_Z + J_Z}\right)\right] \end{split}$$

### Contribution to the coefficients of the dim-6 operators from explicit models:



From EWPT at LEP:  $m_{\rho} > 2 \ TeV \longrightarrow f > 200 \ GeV$ 

### **Deviations from the SM:**



Visible at LHC?





... certainly if they are of order 20-40%

ILC would be a perfect machine to test these scenarios: effects could be measured up to a few %

### Best test of composite Higgs: WW-scattering



even that the Higgs is light, it grows with s

$$\mathcal{A}\left(Z_{L}^{0}Z_{L}^{0} \to W_{L}^{+}W_{L}^{-}\right) = \mathcal{A}\left(W_{L}^{+}W_{L}^{-} \to Z_{L}^{0}Z_{L}^{0}\right) = -\mathcal{A}\left(W_{L}^{\pm}W_{L}^{\pm} \to W_{L}^{\pm}W_{L}^{\pm}\right) = \frac{c_{H}s}{f^{2}},$$
$$\mathcal{A}\left(W^{\pm}Z_{L}^{0} \to W^{\pm}Z_{L}^{0}\right) = \frac{c_{H}t}{f^{2}}, \quad \mathcal{A}\left(W_{L}^{+}W_{L}^{-} \to W_{L}^{+}W_{L}^{-}\right) = \frac{c_{H}(s+t)}{f^{2}},$$
$$\mathcal{A}\left(Z_{L}^{0}Z_{L}^{0} \to Z_{L}^{0}Z_{L}^{0}\right) = 0.$$

Difficult to see. From Higgsless studies possible to see if  $\frac{c_H v^2}{f^2} \sim 0.5 - 0.7$ 

2 Higgs-production also grows with s:



 $\mathcal{A}\left(Z_L^0 Z_L^0 \to hh\right) = \mathcal{A}\left(W_L^+ W_L^- \to hh\right) = \frac{c_H s}{f^2}.$ 

Challenging!

# Conclusions

- "There is life" beyond the MSSM
- Alternatives based on either (i) No Higgs, or (ii) Composite/PGB Higgs
- WW-amplitudes unitarized by (i) extra states (e.g. KK-states) and, in case (ii) by a non-elementary "Higgs"

Worthy to be explored at the LHC

Signals:

- W',Z'-type resonances: Quite generic
- Gluonic resonances: Cleanest signature
- Fermionic resonances: Lightest states "partners" of the top
- Top+Higgs couplings different from the SM, and strong WW-scattering at high E

although also exotic states, e.g. Q=5/3