

UNPARTICLES

A VIEW FROM THE HIGGS WINDOW

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★ Introduction to Unparticles

★ Exploring the Unparticle World through
the Higgs boson

Work with A. Delgado (Notre Dame), J.M. No (IFT) & M. Quiros (IFAE)

[[hep-ph/0707.4309](https://arxiv.org/abs/hep-ph/0707.4309)] [[hep-ph/0802.2680](https://arxiv.org/abs/hep-ph/0802.2680)]

HIDDEN SECTOR OF UNPARTICLES

Georgi-03

Assume a scale invariant sector of the theory

⇒ Very different from our observable sector

★ SM contains explicit mass term $V \supset -m^2 |H|^2$

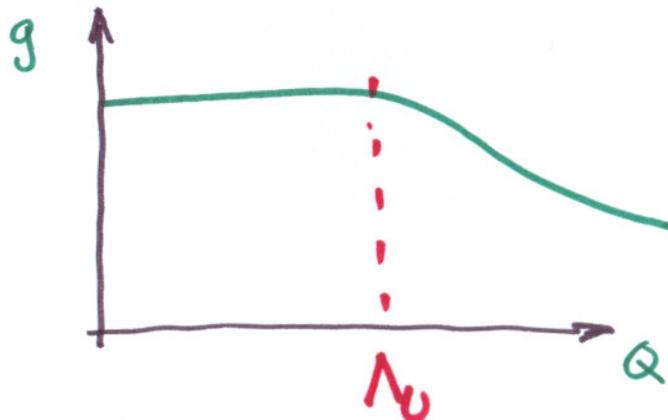
Scale inv. sector ⇒ Massless fields*

★ Quantum corrections can generate a scale (QCD)

Scale inv. sector ⇒ $\beta(g) = 0$ required & $g \neq 0$ ($g=0$ is trivial)

Ex:

Banks-Zaks
 $SU(N_c)$
 N_f flavors
tailored
to give →

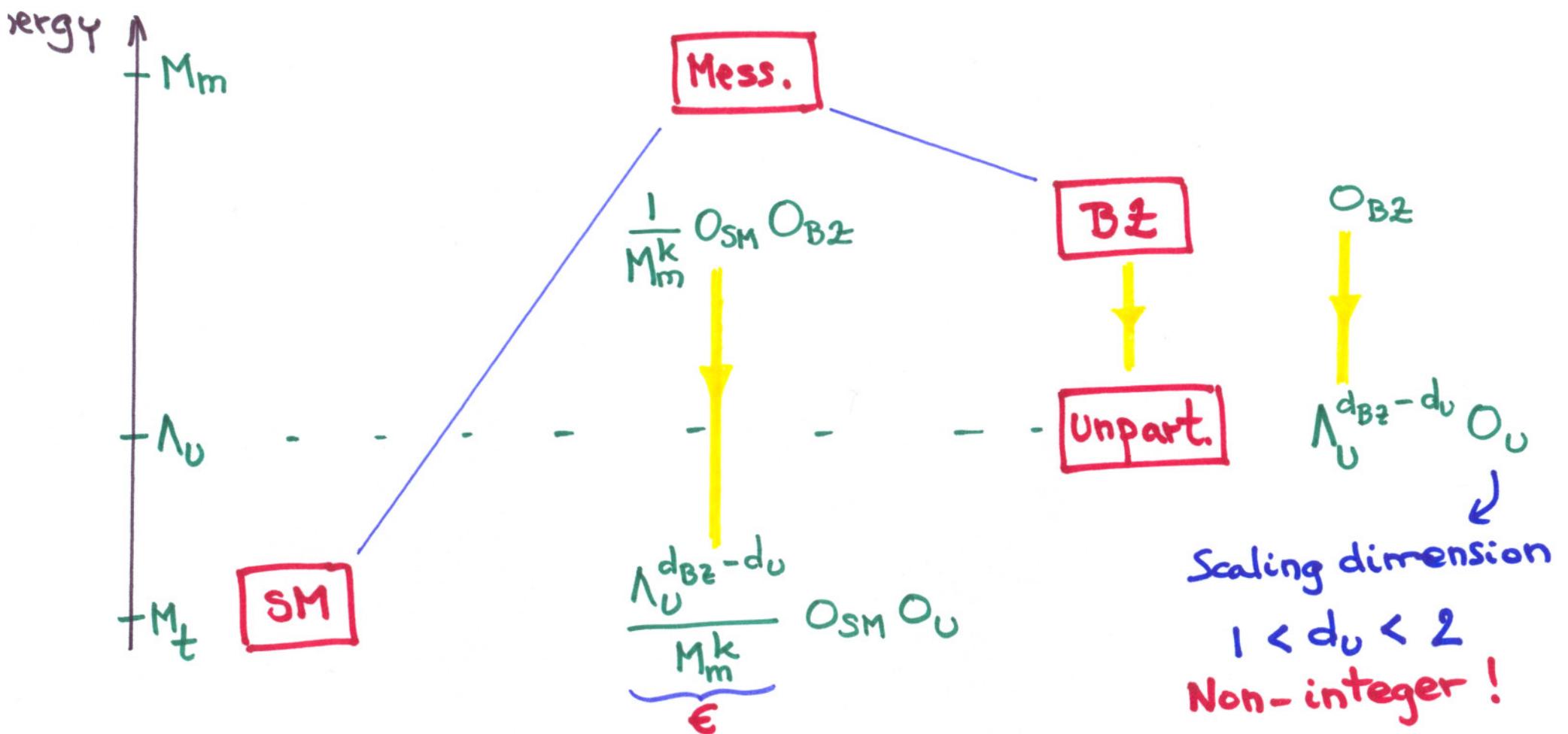


Non-trivial IR fixed point



Dimensional
"anti-transmutation"

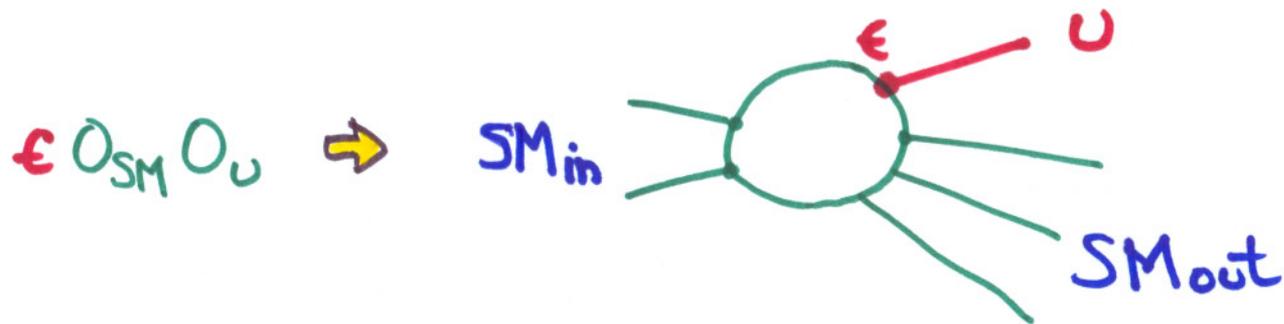
COUPLE SM TO SCALE INV. SECTOR



* d_U : O_U cannot be interpreted in terms of particles

* As ϵ small \Rightarrow Unparticle effects not seen (so far...)

UNPARTICLE PRODUCTION



At lowest order in ϵ



Missing energy & momentum

Probability $\rightarrow \epsilon^2 |\langle SM_{out} | O_{SM} | SM_{in} \rangle|^2 |\langle U | O_U | 0 \rangle|^2$

Phase space for unparticles \Rightarrow determined by scale invariance:

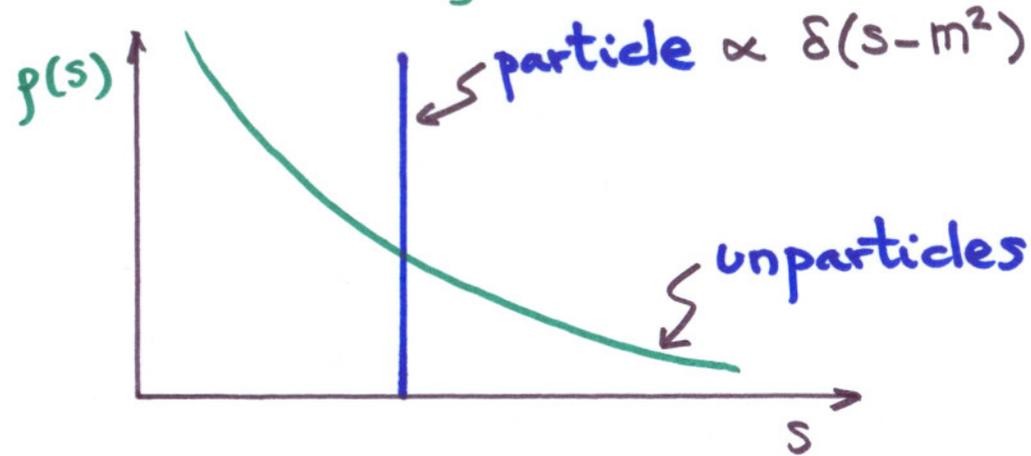
$$\langle 0 | O_U(x) O_U^\dagger(0) | 0 \rangle = \int e^{-ipx} \underbrace{|\langle U, p | O_U(0) | 0 \rangle|^2 \rho(p^2)}_{A_{d_U} \theta(p^0) \theta(p^2) (p^2)^{d_U-2}} \frac{d^4 p}{(2\pi)^4} \sim x^{-2d_U}$$

Compare phase space for n massless particles $dLIPS_n = A_n s^{n-2}$

Unparticles \sim non-integer number of massless particles!

UNPARTICLE SPECTRAL FUNCTION

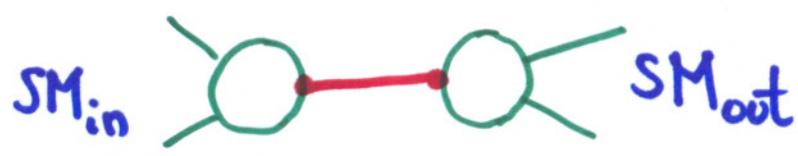
from $P(p^2) \sim (p^2)^{d_U-2} \Rightarrow f(s) \sim s^{d_U-2}$



O_U does not create particles out of the vacuum but a non-localized wave over the full range of $p^2 \Rightarrow$ **Unparticles**

\Rightarrow Lots of phenomenological implications explored.

E.g.



Interference $\sim \epsilon^2 \Delta_U(p^2)$
 strange phase inside

DECONSTRUCTING UNPARTICLES

Stephanov-05

Unparticles as infinite tower of massive scalars ϕ_n ($n=1, \dots, \infty$) with

$$M_n^2 = \Delta^2 n$$

in the limit of zero mass splitting $M_{n+1}^2 - M_n^2 = \Delta^2 \rightarrow 0$

⇒ **Scale invariant continuous mass spectrum!**

$$S = \int d^4x \sum_{n=1}^{\infty} \left[\frac{1}{2} (\partial_\mu \phi_n)^2 + \frac{1}{2} M_n^2 \phi_n^2 \right] \xrightarrow[\substack{\Delta^2 \rightarrow dM^2 \\ \phi_n \rightarrow \Delta \cdot u(M^2)}]{\Delta \rightarrow 0} S = \int d^4x \int_0^{\infty} dM^2 \left[\frac{1}{2} (\partial_\mu u)^2 + \frac{1}{2} M^2 u^2 \right]$$

$\phi_n(x) \rightarrow \lambda \phi_n(x\lambda)$ $u(M^2, x) \rightarrow u(M^2/\lambda^2, x\lambda)$

Deconstructed \mathcal{O}_U :

$$\mathcal{O} \equiv \sum_n F_n \phi_n \rightarrow \mathcal{O}_U \quad \text{for} \quad F_n^2 = \frac{A du}{2\pi} \Delta^2 (M_n^2)^{d_U-2}$$

⇒ Not clear phys. interp. (relation to RS?) but very useful tool

HIGGS AS PORTAL TO UNPARTICLE WORLD

$$k_U O_{SM} O_U \rightarrow k_U |H|^2 O_U \quad \text{Relevant operator!}$$

⇒ After EWSB this is the main source of conformal breaking

Difficulty: $\langle H \rangle = \frac{v}{2}$ induces a tadpole for O_U triggering the VEV

$$\langle O_U \rangle = - \frac{k_U v^2}{2} \int_0^\infty \frac{F^2(M^2)}{M^2} dM^2 \propto \int_0^\infty (M^2)^{d_U-3} dM^2$$

$d_U < 2 \Rightarrow$ **IR Divergence**

DEQ-07

$$\left\{ \begin{array}{l} \text{Tadpole} \sim \sum_n F_n \varphi_n \sim M_n^{d_U-2} \varphi_n \xrightarrow{M_n^2 \rightarrow 0} \infty \\ \text{Restoring mass} \sim M_n^2 \varphi_n^2 \rightarrow 0 \end{array} \right.$$

CURES FOR IR PROBLEM

(A) Add coupling $\zeta |H|^2 \sum_n \phi_n^2 \rightarrow \zeta v^2$ as **IR cutoff**
Delgado, E, Quiros

(B) Add coupling $\frac{1}{4} \lambda v (\sum_n \phi_n)^2$ work in progress
Delgado, E, Quiros, No

(A) $\rightarrow \langle O_v \rangle = -\frac{\kappa v^2}{2} \int_0^\infty \frac{F^2(M^2)}{M^2 + \zeta v^2} dM^2 \propto \int_0^\infty \frac{(M^2)^{d_b-2}}{M^2 + \zeta v^2} dM^2$
IR finite ✓

$\zeta v^2 = m_{\text{gap}}^2$ for unparticle continuum



\rightarrow Changes dramatically pheno and constraints

EFFECTS ON EWSB

Deconstructed potential

$$V = m^2 |H|^2 + \lambda |H|^4 + \frac{1}{2} \sum_n M_n^2 \varphi_n^2 + k_U |H|^2 \sum_n F_n \varphi_n + \zeta |H|^2 \sum_n \varphi_n^2$$

Min Conditions \Rightarrow

$$m^2 + \lambda \varphi^2 + k_U \sum_n F_n \varphi_n + \zeta \sum_n \varphi_n^2 = 0$$

and
$$\varphi_n = - \frac{k_U \varphi^2}{2(M_n^2 + \zeta \varphi^2)} F_n$$

Continuum limit \Rightarrow
$$m^2 + \lambda \varphi^2 - \lambda_U (\mu_U^2)^{2-d_U} \varphi^{2(d_U-1)} = 0 \quad \left[(\mu_U^2)^{2-d_U} \equiv k_U^2 \frac{A d_U}{2\pi} \right]$$

Correction looks like coming from $\delta V \propto h^{2d_U}$ between h^2 & h^4 !

\Rightarrow It's possible to have a minimum for $m^2 \geq 0$

HIGGS - UNPARTICLES INTERPLAY

(Infinite)
Mass Matrix
after EWSB



$$\begin{bmatrix} m_{h_0}^2 & A_1 & \dots & A_n & \dots \\ A_1 & M_1^2 + \xi v^2 & & & \\ \vdots & & \ddots & & \\ A_n & & & M_n^2 + \xi v^2 & \\ \vdots & & & & \ddots \end{bmatrix}$$

1. Mass gap in diagonal entries $m_g^2 = \xi v^2$

2. Mixing Higgs - Unparticles

$$A_n = v(k_v F_n + 2\xi v_n)$$

3. Pole structure? Examine $p^2 \Pi - M^2$

MIXED PROPAGATOR

$$h \text{ --- } \text{[Crossed Circle]} \text{ --- } h = \frac{h}{\quad} + \text{---} \times \overset{u}{\text{---}} \times \text{---} + \text{---} \times \text{---} \times \text{---} \times \text{---} + \dots$$

or

$$u \text{ --- } \text{[Crossed Circle]} \text{ --- } u = \frac{u}{\quad} + \text{---} \times \overset{h}{\text{---}} \times \text{---} + \text{---} \times \text{---} \times \text{---} \times \text{---} + \dots$$

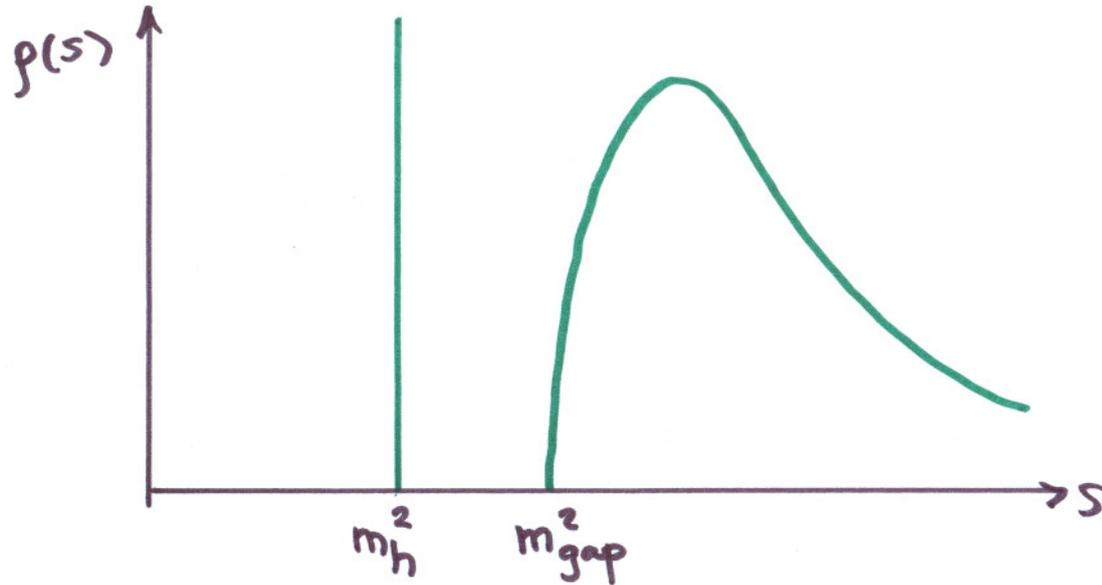
$$\Rightarrow iP(p^2)^{-1} = p^2 - m_h^2 + \omega^2(\mu^2)^{2-d_U} \int_0^\infty \frac{(M^2)^{d_U-2}}{(M^2 + m_{\text{gap}}^2 - p^2)} \left(\frac{M^2}{M^2 + m_{\text{gap}}^2} \right)^2 dM^2$$

For the pole structure \Rightarrow Spectral Function



SPECTRAL FUNCTION ($m_h < m_{\text{gap}}$)

m_h^2 shifted down from SM value $m_{h0}^2 = 2\lambda\varphi^2$ (eigenvalue repulsion)



$$\rho(s) = \frac{1}{K^2} \delta(s - m_h^2) + \theta(s - m_{\text{gap}}^2) \frac{Q_0^2(s)}{[iP(s)^{-1}]^2 + \pi^2 Q_0^4(s)}$$

ZZ COUPLING

$$\rho_{hh}(s) \equiv \langle h|s\rangle\langle s|h\rangle = |\langle h|H\rangle|^2 \delta(s-m_h^2) + \theta(s-m_g^2) |\langle h|U,s\rangle|^2$$

$$\text{Int. Eigenstates } \begin{cases} |h\rangle \\ |u,s\rangle \end{cases} \rightarrow \text{Mass Eig. } \begin{cases} |H\rangle \\ |U,s\rangle \end{cases}$$

★ HZZ coupling \Rightarrow diluted by unparticle contamination of isolated pole

Pure Higgs composition $\Rightarrow R_h \equiv \langle h|H\rangle = \frac{1}{K}$

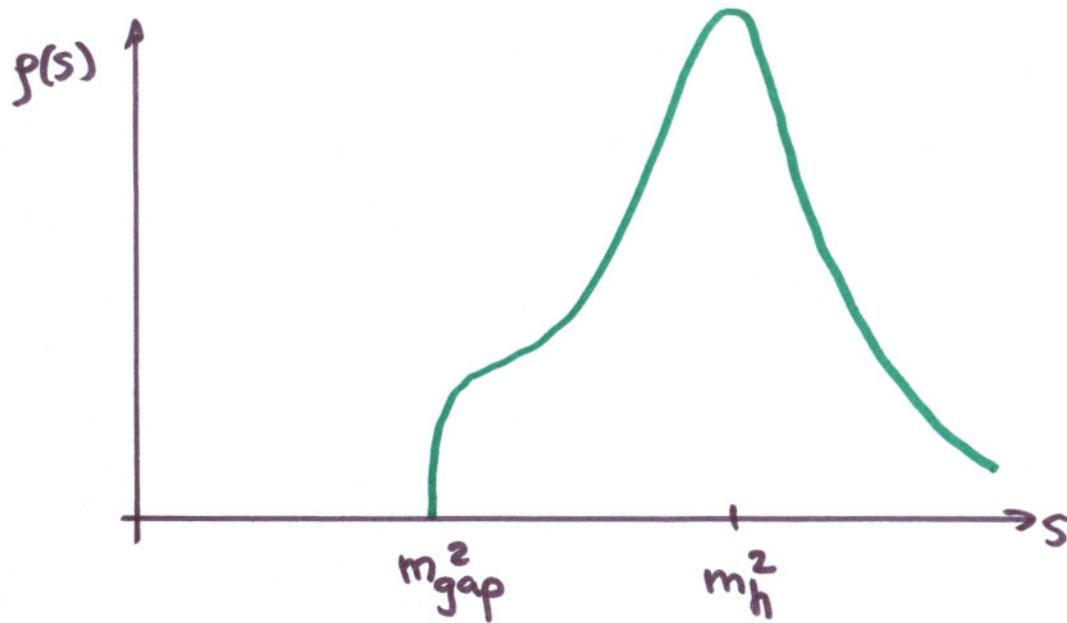
★ UZZ coupling \Rightarrow induced by Higgs contamination
Higgs diffused in unparticle continuum

$$\Rightarrow R_u(M^2) = \theta(M^2 - m_g^2) |\langle h|U, M^2\rangle|^2$$

Sum Rule $\Rightarrow R_h^2 + \int_{m_g^2}^{\infty} R_u^2(M^2) dM^2 = 1$

SPECTRAL FUNCTION ($m_h > m_{\text{gap}}$)

Higgs pole is subsumed in unparticle continuum



$$\rho(s) = \theta(s - m_{\text{gap}}^2) \frac{Q_0^2(s)}{[iP(s)^{-1}]^2 + \pi^2 Q_0^4(s)}$$

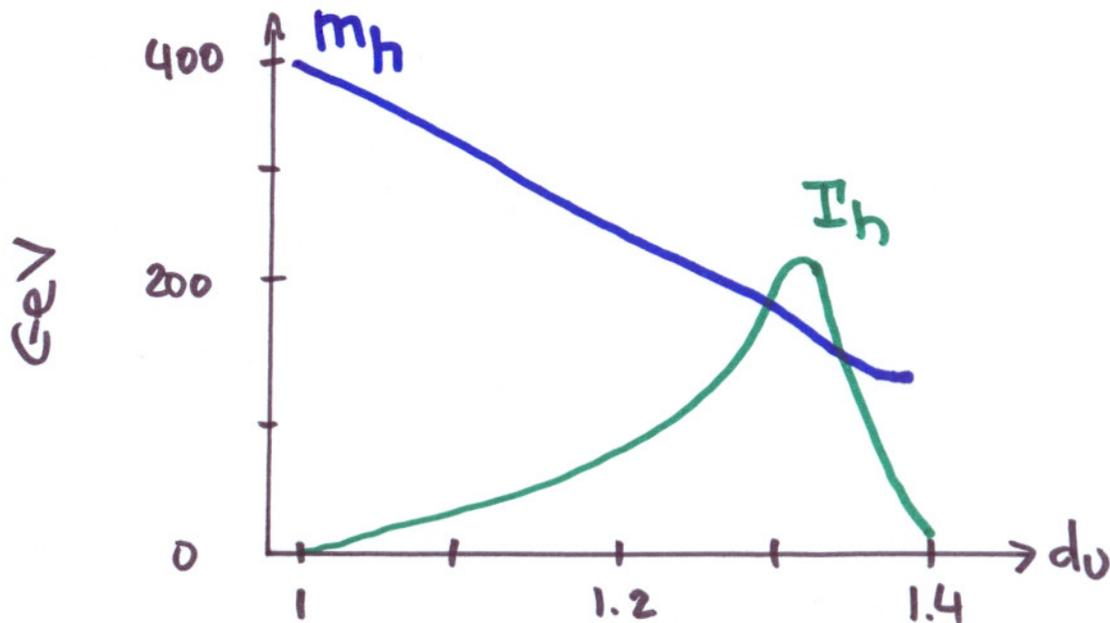
HIGGS WIDTH ($m_h > m_{\text{gap}}$)

Near Higgs pole $\Rightarrow iP(s)^{-1} \simeq (s - m_h^2) \tilde{K}$



$$\Gamma_h = \frac{\pi Q_0^2(m_h^2)}{m_h \tilde{K}^2}$$

Γ_h can be quite large



UN-CONCLUSIONS

Unparticles

extremely speculative, no good motivation?

But

keep in mind for LHC; interesting Higgs effects ...

→ $|H|^2 O_U$

can destabilize the theory but can be cured

→ EWSB

Effects as $h^2 du$!

Unconventional possibilities allowed ($m^2 \geq 0$)

EWSB →

Induces a mass gap

Dramatic impact in phenomenology and constraints

Higgs - Unparticles: Drastic changes in Higgs mass, couplings, width

Unparticles get "Higgs contaminated": UZZ