

Quarkonia Propagation & Collectivity in the QGP : Elastic Diffusions, Energy losses and elliptic flow $v_2(\phi)$

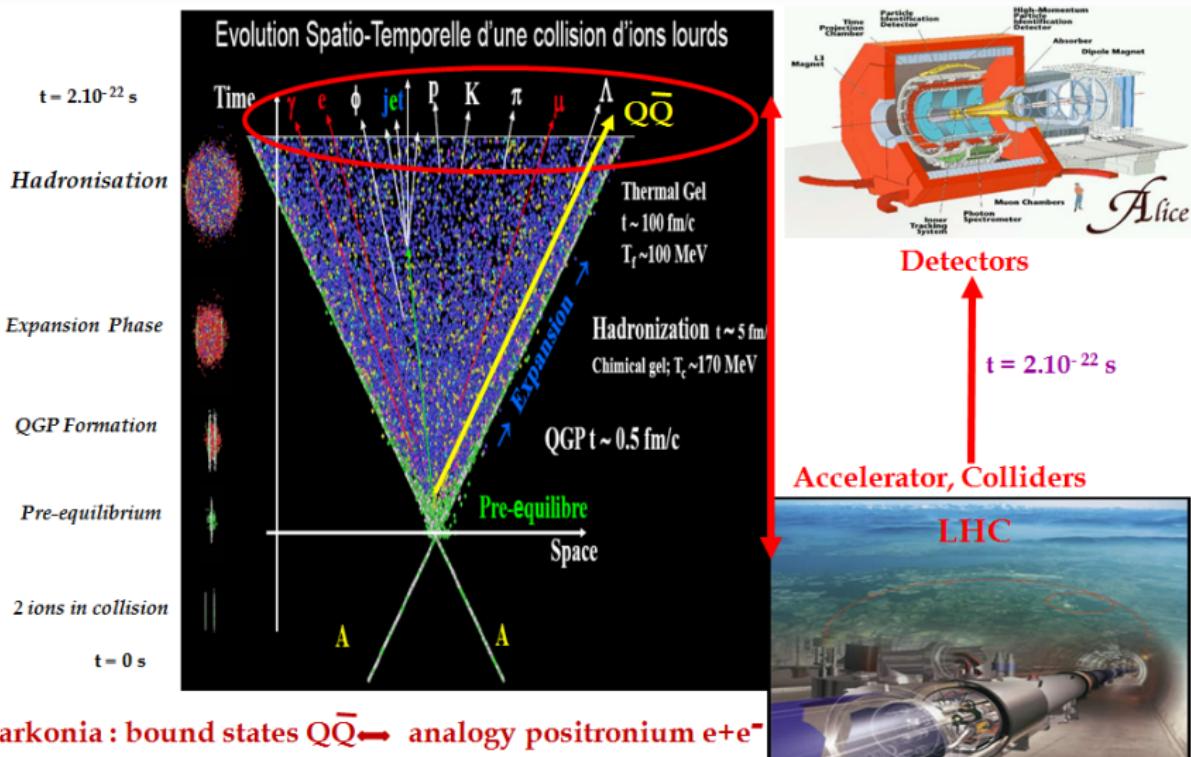
Hamza Berrehrah

Supervisors: P.B. Gossiaux & J. Aichelin

Subatech-Nantes

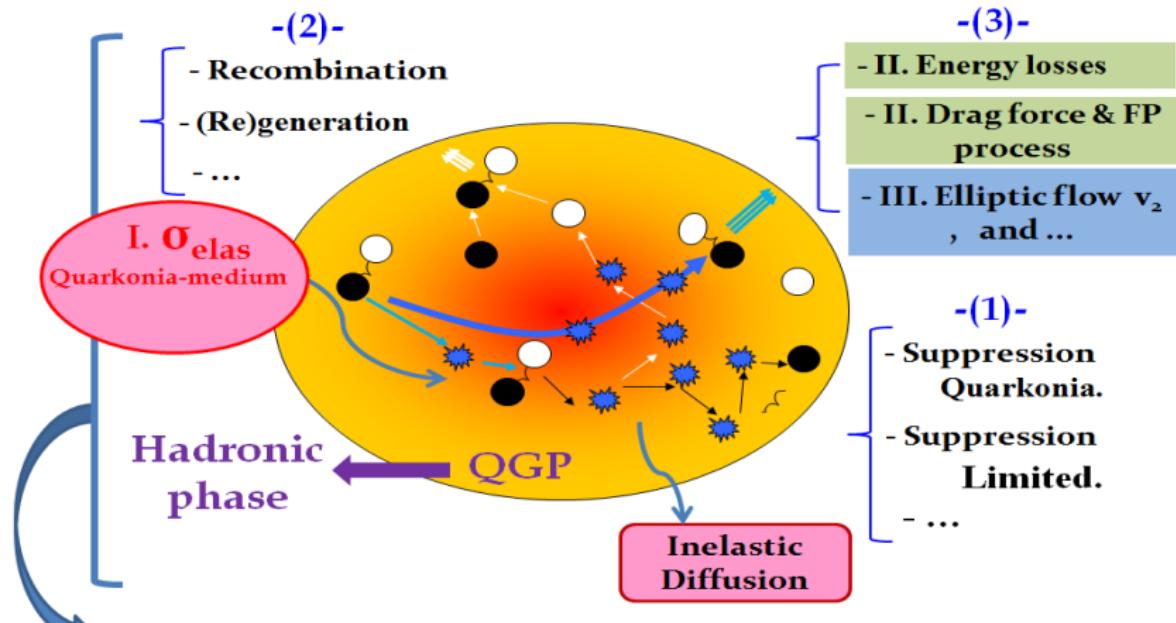
AGThNuc 2010 ~ Wednesday 20th Octobre 2010

QGP Probe : bound states of Quarkonia ...

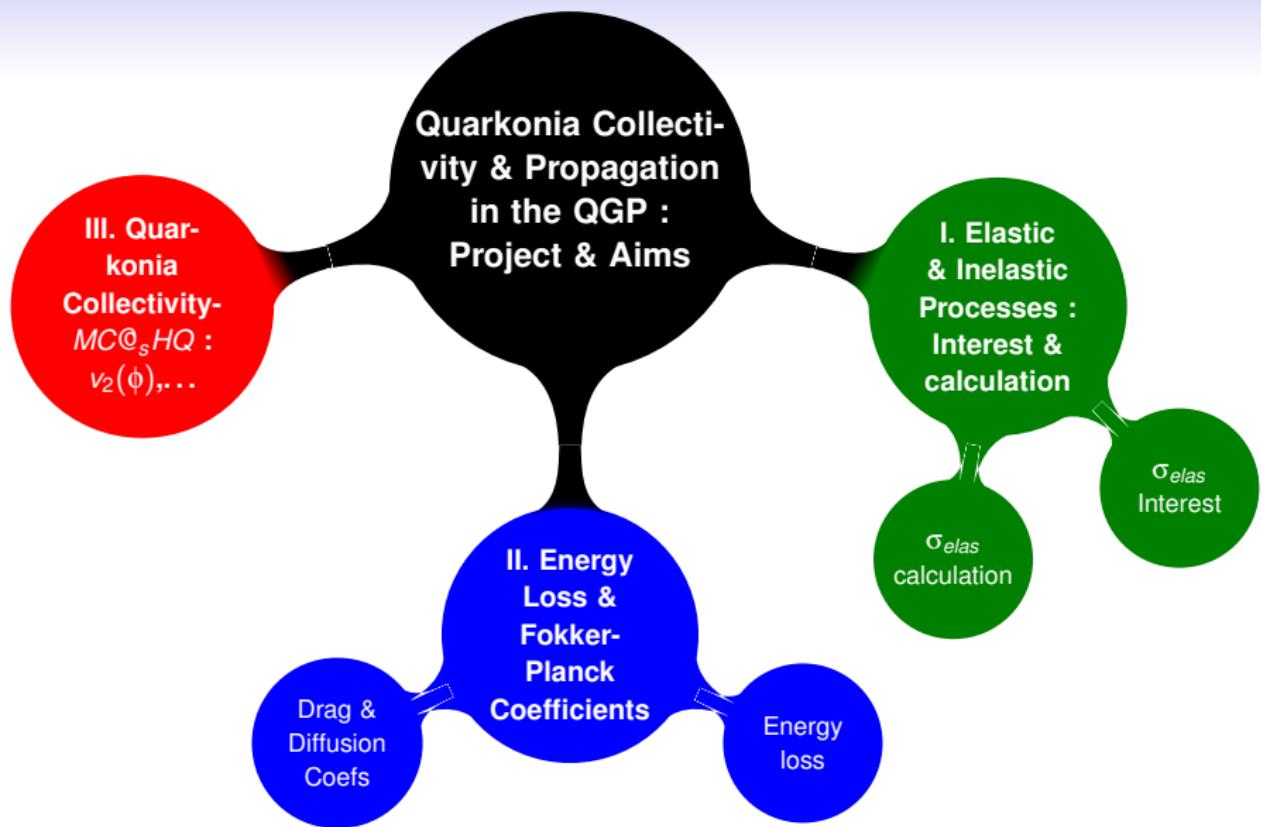


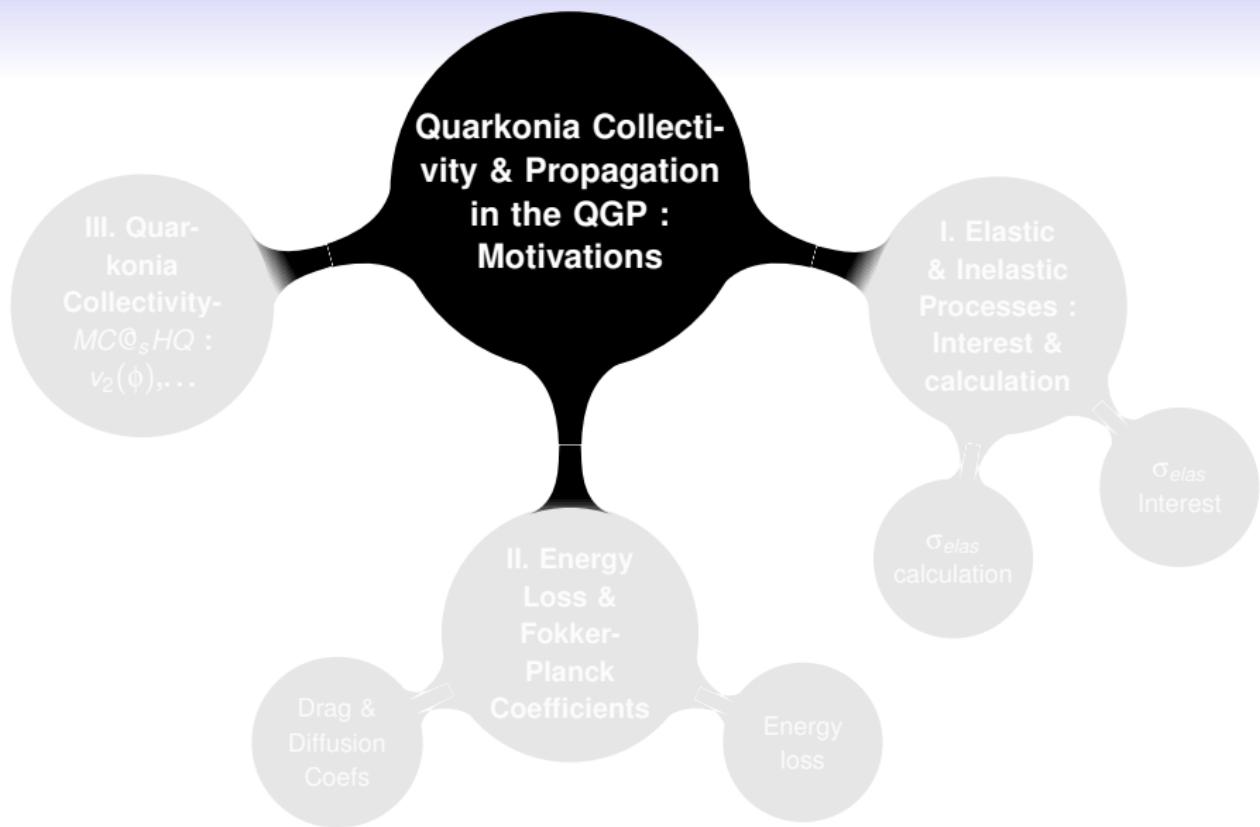
Quarkonia propagation and collectivity in QGP -AGThNuc 2010-H.Berrerah

Project : quarkonia propagation & collectivity ...

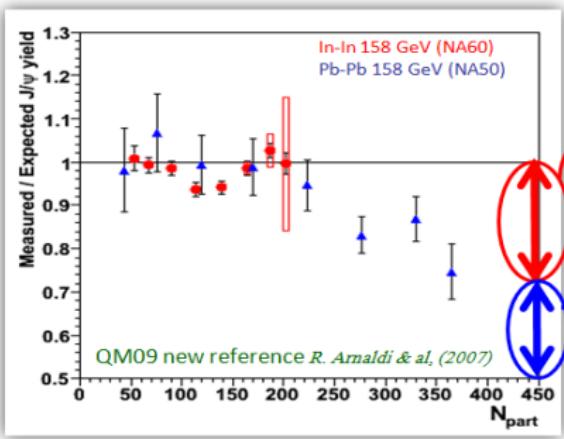


Quarkonia Propagation & Collectivity in a thermalized QGP → QGP Proprieties



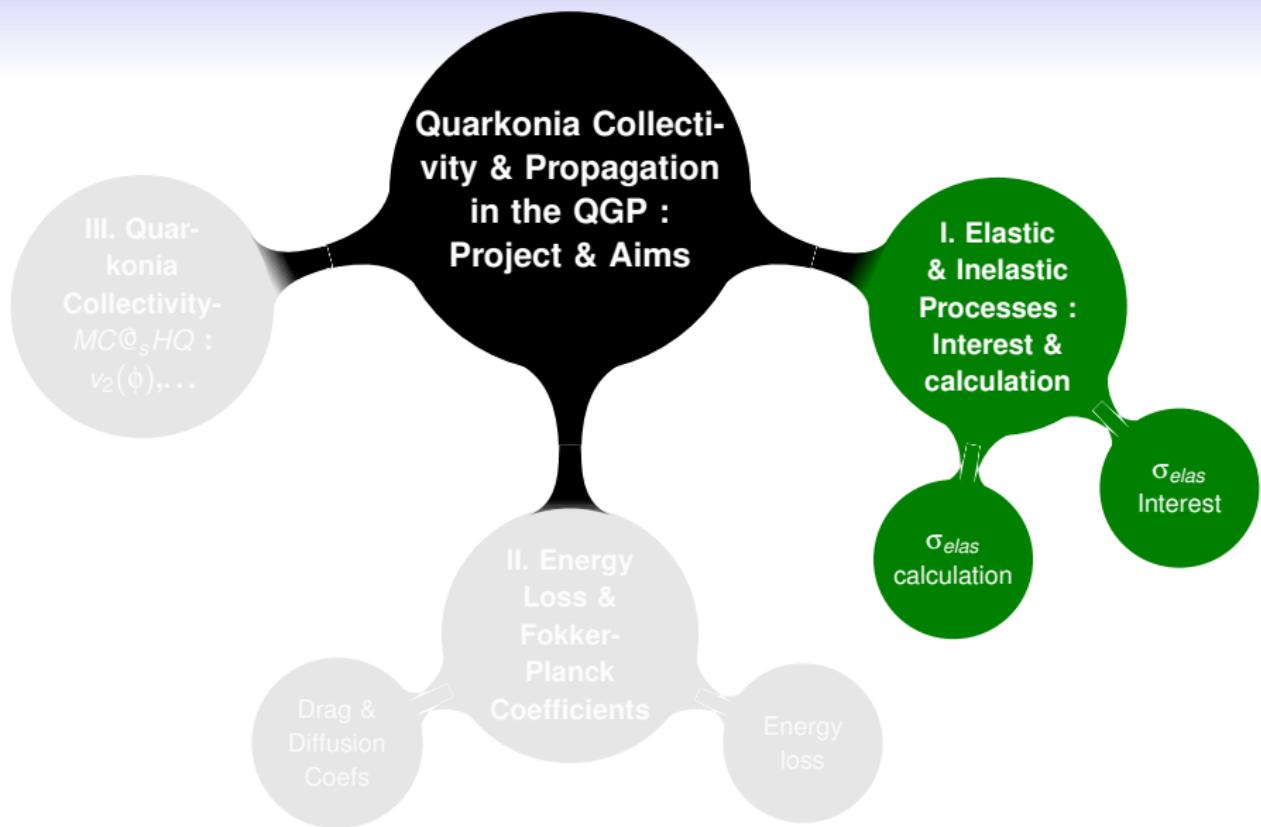


Motivations ...



- Suppressed (studied with σ_{inel})
- Focus on J/ψ remaining properties modified in the plasma during the scattering J/ψ-hadron, J/ψ-gluons...
- σ_{elas} : Elliptic flow, Energy losses ...

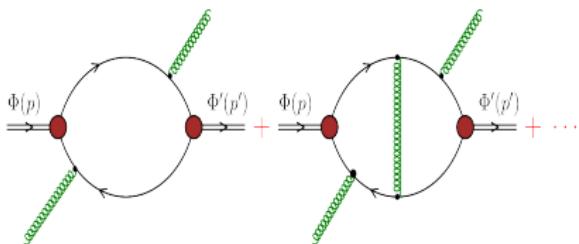
- + J/ψ Suppression studied for 20 years,
... But No consistent results
- + Recent Results of RHIC ... No significant additional suppression expected has been observed for J/ψ by increasing energy...
- + Large fluctuations over the value of σ_{elas}



Elastic & Inelastic Process

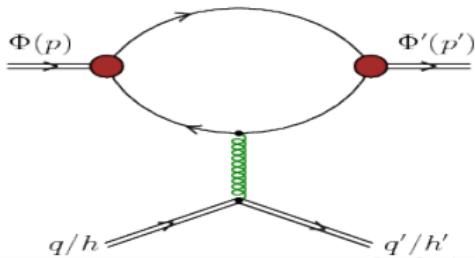
1 Elastic Process

- Compton Diffusion



Most important process

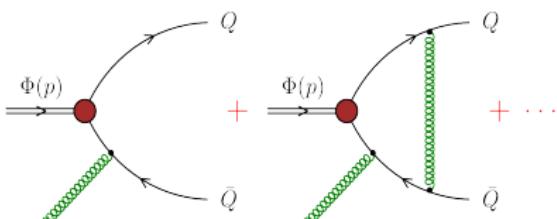
- Hadron Diffusion



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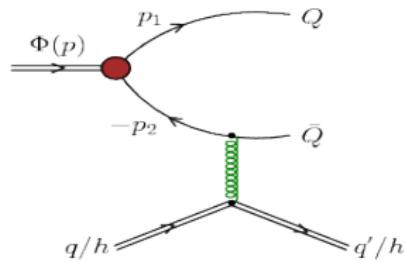
2 Inelastic Process

- Gluon Dissociation



Dominant process

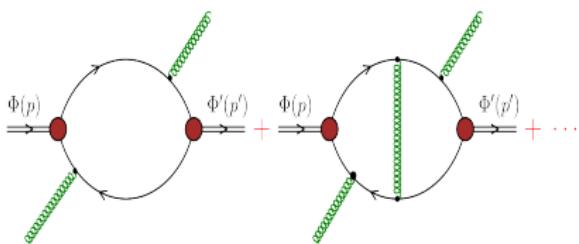
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Elastic & Inelastic Process

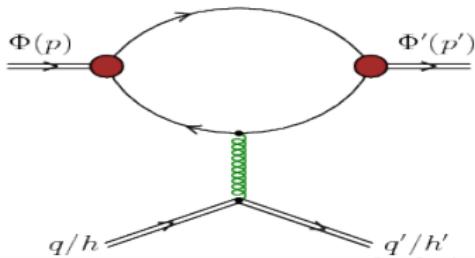
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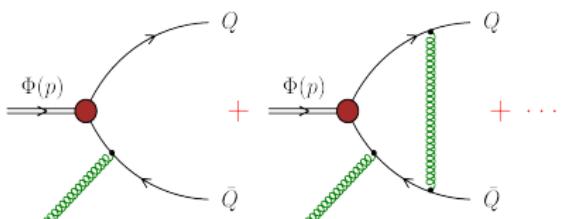
Most important process

- Hadron Diffusion



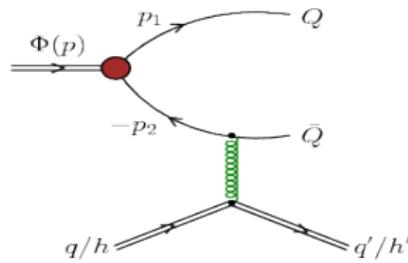
2 Inelastic Process

- Gluon Dissociation



Dominant process

- Hadron Dissociation



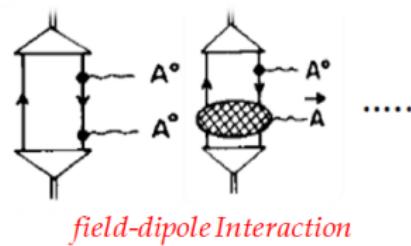
σ_{elas} calculation : How ?

● Formalisms for σ_{elas} calculation

① Historical Calculation on $Q\cdot h/g \sigma_{elas}$

① Bhanot and Peskin formalism (79)

- From OPE (operator product expansion)
- Binding energy = $e_0 \gg LQCD$



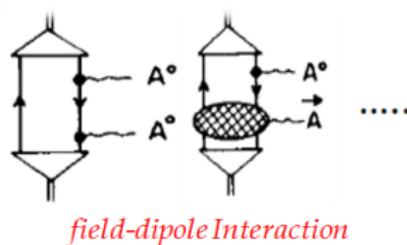
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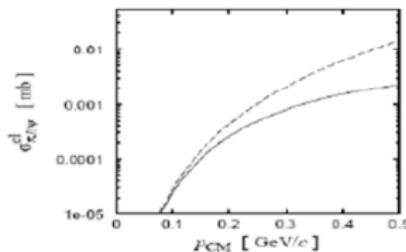
① Bhanot and Peskin formalism (79)

- From OPE (operator product expansion)
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② Kharzeev and Fujii (99), Povh & Hüfner

- Short-distance QCD calculation
- Optic theorem ...



Kharzeev and Fujii (99)

σ_{elas} calculation : How ?

- Formalisms for σ_{elas} calculation

- ② Rederivation of Peskin formula using Bethe-Salpeter
- ③ Other equivalent method in 1st order in pQCD based on :
 - Factorization Theorem + Bethe-Salpeter Amplitudes

Factorization Formula

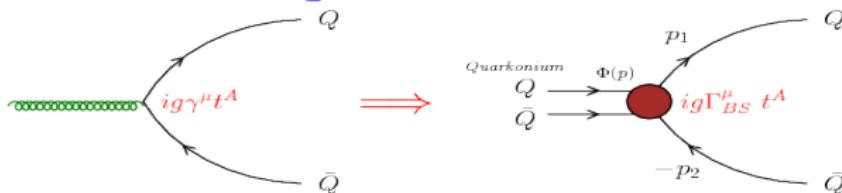
$$\sigma_{\phi h}(v) = \int_0^1 dx \sigma_{\phi g}(xv) \times g(x)$$

- σ total of scattering quarkonia $\phi - h$
- Perturbative Cross Section of $\phi - g$ scattering \Rightarrow BS Formalism
- Distribution function of gluons in the hadron

Y. Oh, S. Kim, S. Hwang Lee, (02)

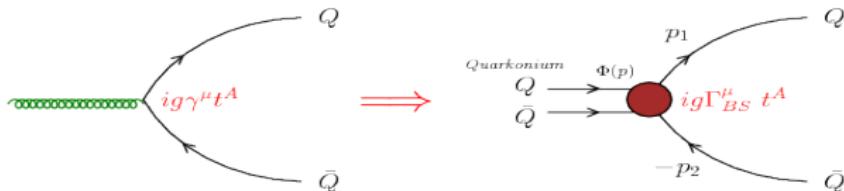
Bethe Salpeter Formalism

① Goal : Bethe Salpeter Vertex



Bethe Salpeter Formalism

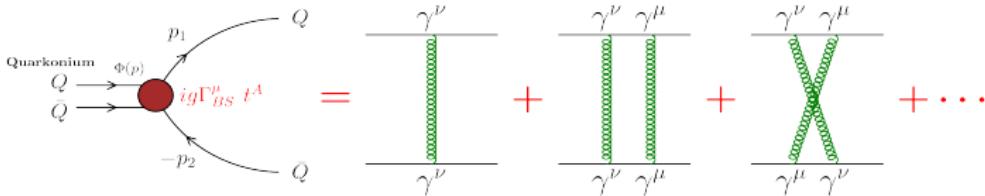
① Goal : Bethe Salpeter Vertex



② Bethe Salpeter Equation

$$\mathcal{M} = \mathcal{M} + \int \mathcal{V} \mathcal{G} \mathcal{V} + \int \int \mathcal{V} \mathcal{G} \mathcal{V} \mathcal{G} \mathcal{V} + \dots + \left(\int \mathcal{V} \mathcal{G} \right)^n + \dots = \frac{\mathcal{V}}{1 - \int \mathcal{V} \mathcal{G}}$$

\mathcal{V} : kernel, \mathcal{M} : amplitude, \mathcal{G} : propagateur

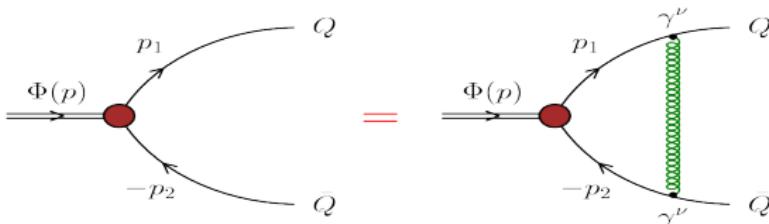


Bound states : produce a pole in \mathcal{M} .

$\Rightarrow \mathcal{M}$ eigenvector Γ satisfies : $\Gamma = \int_k \mathcal{V}(p, k, P) \mathcal{G}(k, P) \Gamma(k, P)$

Bethe Salpeter Formalism

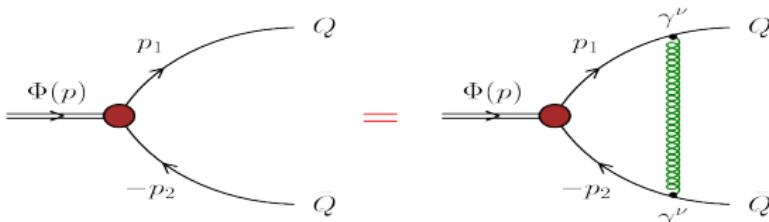
② Bethe Salpeter Equation



$$\Gamma(p, P) = iC_{color} \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 + i\eta} \gamma_\nu \frac{1}{p_1 + k - m + i\eta} \Gamma(p+k, P) \frac{1}{-p_2 + k - m + i\eta} \gamma_\nu$$

Bethe Salpeter Formalism

② Bethe Salpeter Equation



$$\Gamma(p, P) = iC_{color} \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 + i\eta} \gamma^\nu \frac{1}{p_1 + k - m + i\eta} \Gamma(p+k, P) \frac{1}{-p_2 + k - m + i\eta} \gamma^\nu$$

③ Bethe Salpeter Vertices (Case of quarkonium in the rest frame)

• Instantaneous Interaction

$$\Gamma_I(E, \vec{p}) \approx \frac{-ie_0(2e_0 - E)}{\pi E} \gamma^0 \frac{I + \gamma_0 - \vec{p}'/m}{2} \cdot \phi_I^{+-}(\vec{p}) \cdot \frac{I - \gamma_0 - \vec{p}'/m}{2} \cdot \gamma^0$$

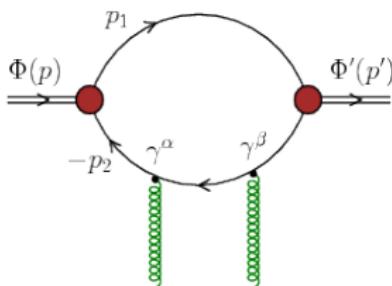
- $\phi_I^{+-}(\vec{p}) = \phi_{space}(\vec{p}) \times \phi_{spin}^{jj}$: instantaneous wave function for the bound state

σ_{elas} Results & Discussion

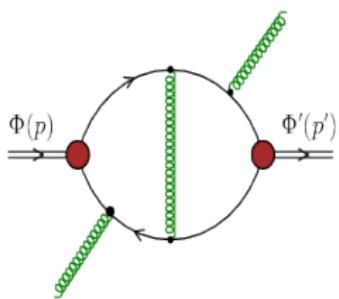
① Compton Diffusion Process $J/\psi - g$

① 2 gluons attached, "LO"

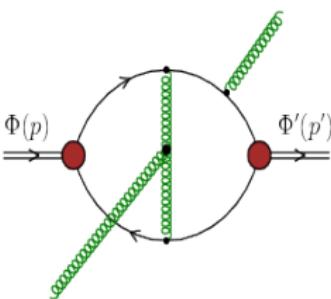
⇒ 6 diagrams ($bb\parallel$, bbX , $tt\parallel$, ttX , tb , bt)



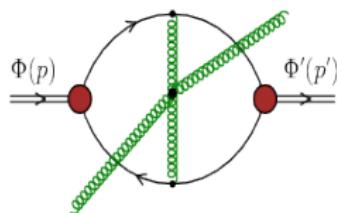
② 3 gluons attached, "S \mathcal{N} LO"



⇒ 4 diagrams
($bb\parallel$, btX , $tt\parallel$, tbX)



⇒ 7 diagrams
(gluon emitted in each line)

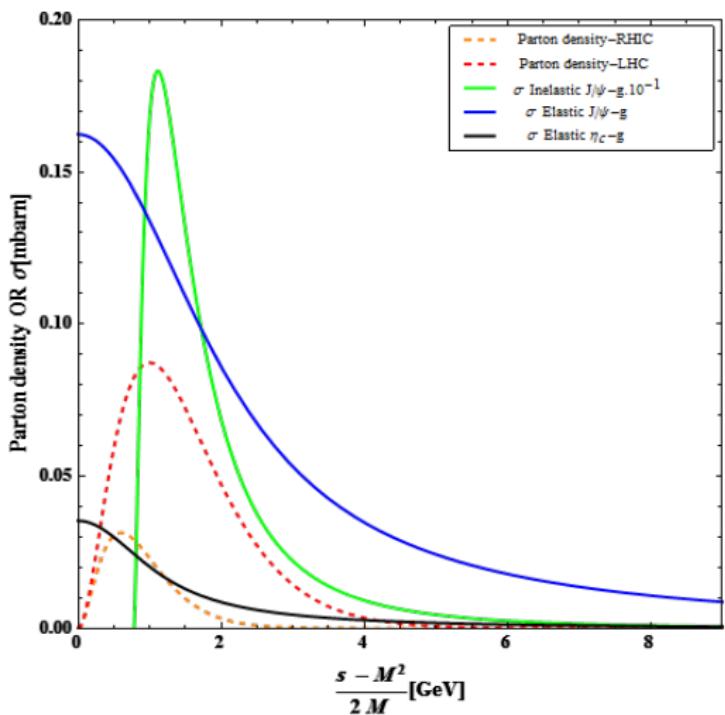


⇒ 1 diagram

σ_{elas} Interest & Discussion

① $J/\psi - g$: Gluon Dissociation vs Compton Diffusion

"LO" Diagrams



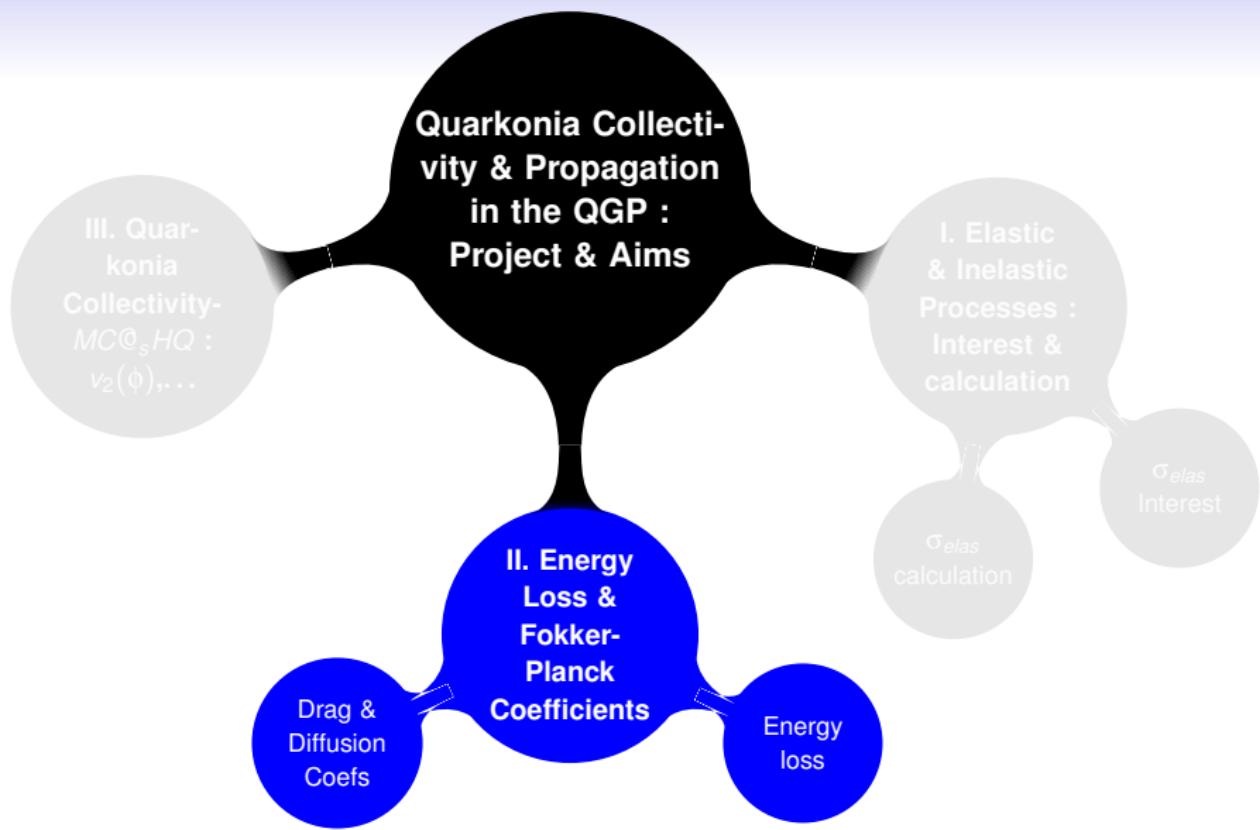
$$R = \int d^3k n_{mb}(k) \sigma_{elas/inel}(s(k))$$

$$\hookrightarrow R = \int_{M^2}^{+\infty} ds \tilde{n}_{mb}(k) \sigma_{elas/inel}(s)$$

- Inelastic cross section has a threshold
- Quantities measured are convoluted by MB Distribution
- Overlap σ_{elas} and MB ditribution larger than σ_{inel} and MB.

Y. Oh, S. Kim, S. Houngh Lee, (02)

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Energy losses & Transport Coefficient

① Energy losses, given by Bjorken ($\phi(M, E, p) \rightsquigarrow "i" (m, e, q)$)

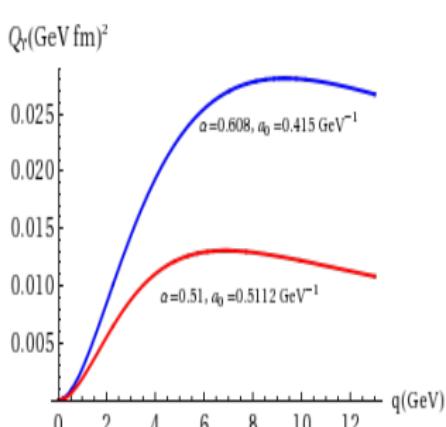
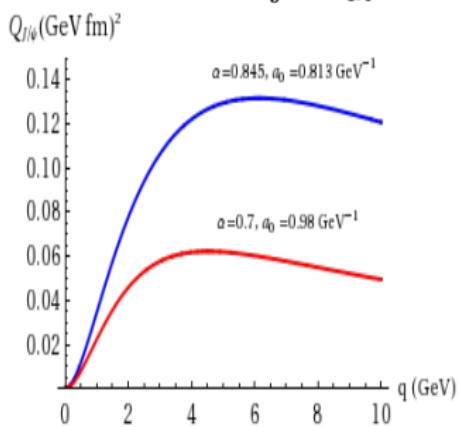
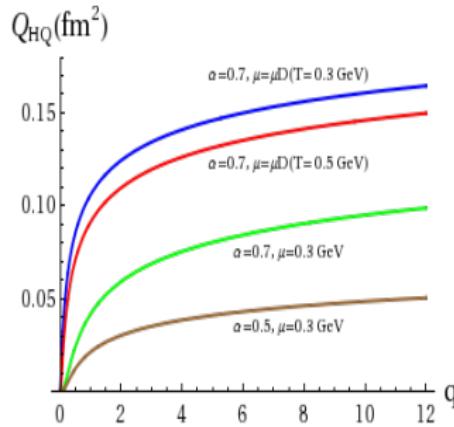
$$\frac{dE}{d\tau} = \frac{dE}{dt} \times \frac{E}{M} = \sum_i \int d^3q n_i(\vec{q}) \frac{\sqrt{(p \cdot q)^2 - M^2 m^2}}{Me} \int dt \frac{d\sigma_{elas}}{dt} (E' - E),$$

Energy losses & Transport Coefficient

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② *Transport coefficient* $\hat{Q}(s) = \int \frac{d\sigma_{elas}}{dt} t dt$: cf. Part I

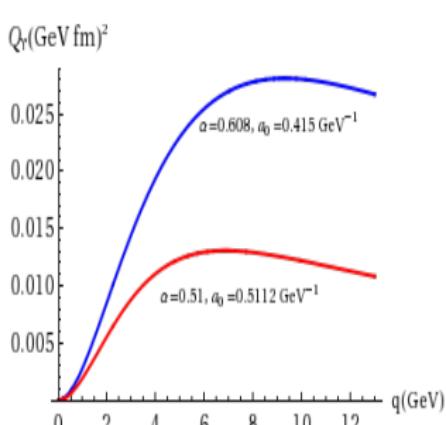
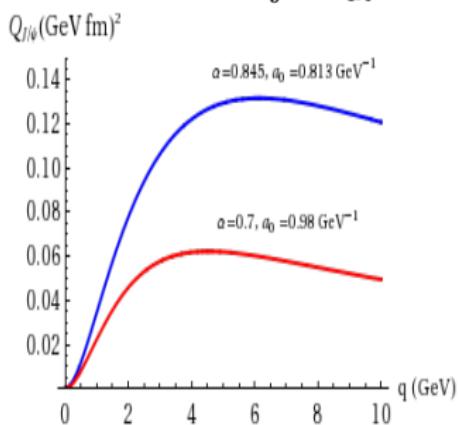
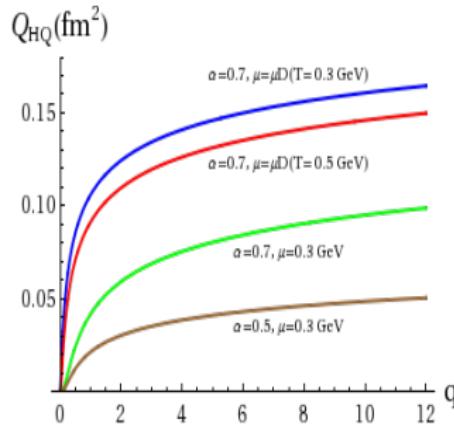


Energy losses & Transport Coefficient

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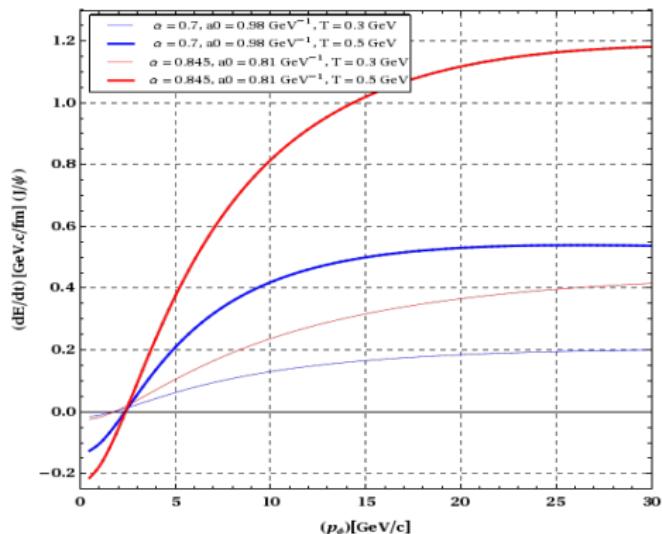
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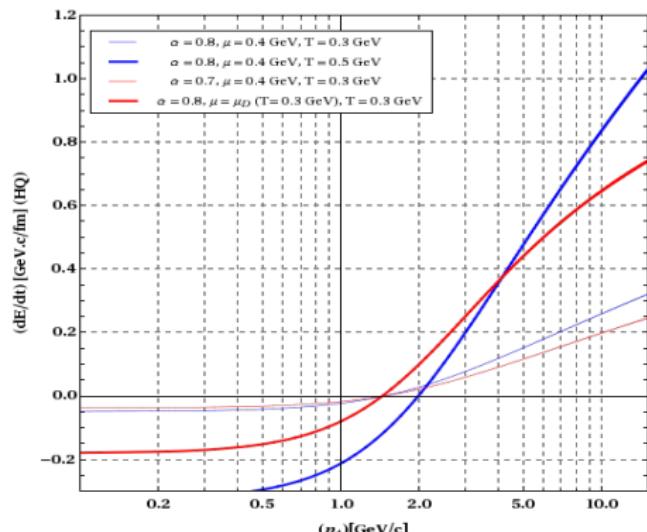


Energy losses (HQ , J/ψ , Υ)

1 Quarkonia (J/ψ)



3 Heavy Quark (C)



- Behaviour : decrease at $p \nearrow$

- Two regimes, $P \approx \sqrt{MT} \Rightarrow \frac{dE_{J/\psi}}{dt} = 0$

- $\frac{dE}{dt}$ variation vs α and T , $\frac{dE}{dt} \nearrow$ with $T \nearrow$

- Behaviour : log increase vs p

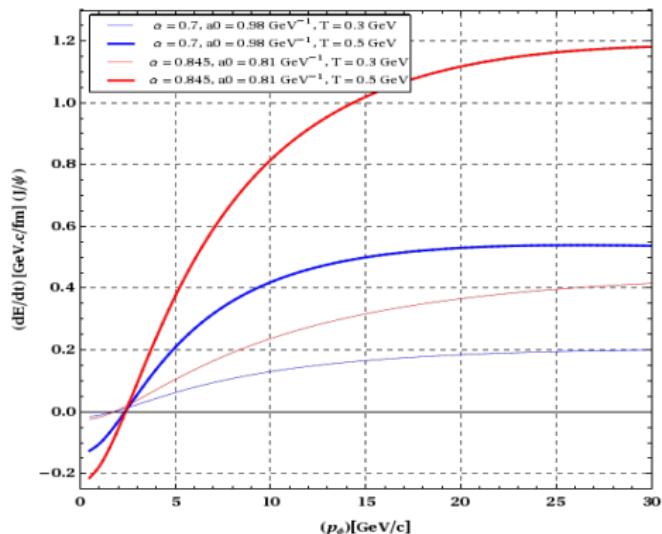
- $P \approx \sqrt{MT} \Rightarrow \frac{dE_{HQ}}{dt} = 0$

- $\frac{dE}{dt} \nearrow$ with $T \nearrow$

Energy losses ($\mathcal{H}Q, J/\psi, \Upsilon$)

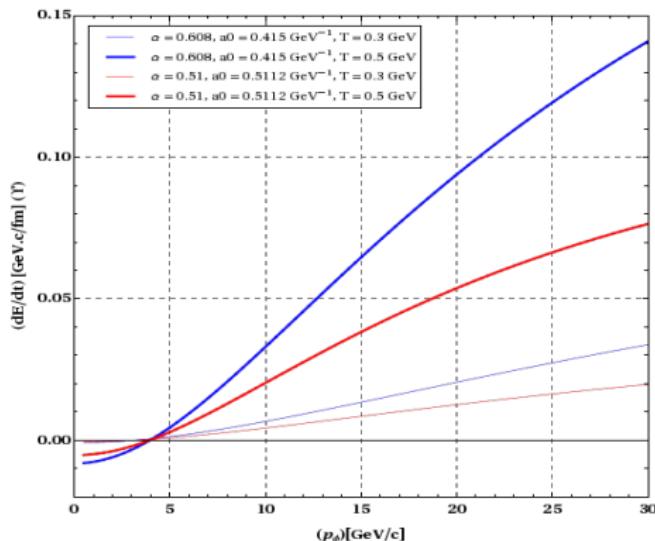
1

Quarkonia (J/ψ)



3

Quarkonia (Υ)



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- Behaviour : decrease at $p \nearrow$

- Two regimes, $P \approx \sqrt{MT} \Rightarrow \frac{dE_{\Upsilon}}{dt} = 0$

- $\frac{dE}{dt}$ variation vs α and T , $\frac{dE}{dt} \nearrow$ with $T \nearrow$

Drag Force & coefficient ($HQ, J/\psi, \Upsilon$)

① **Drag Coefficient**, for: $(\phi(M, E, p) \rightsquigarrow "i" (m, e, q))$

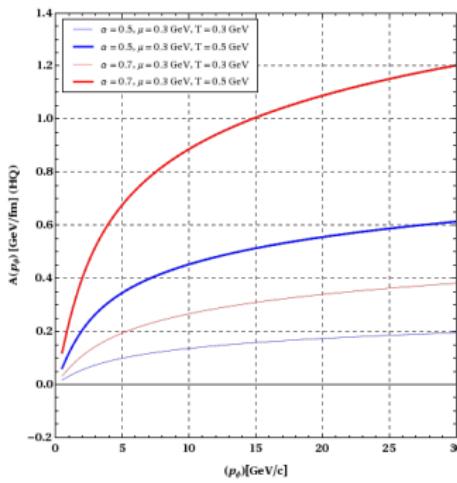
$$A_i = \frac{d < p >}{dt} = \sum_i \int d^3 q n_i(\vec{q}) \frac{\sqrt{(p \cdot q)^2 - M^2 m^2}}{E e} \int dt \frac{d\sigma_{elas}}{dt} \frac{\langle (\vec{P} - \vec{P}') \cdot \vec{P} \rangle}{\|\vec{P}\|},$$

Drag Force & coefficient ($HQ, J/\psi, \Upsilon$)

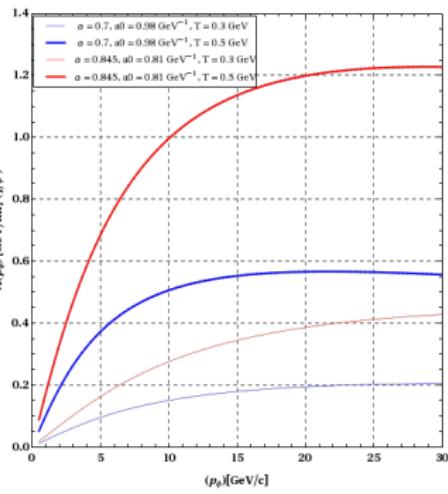
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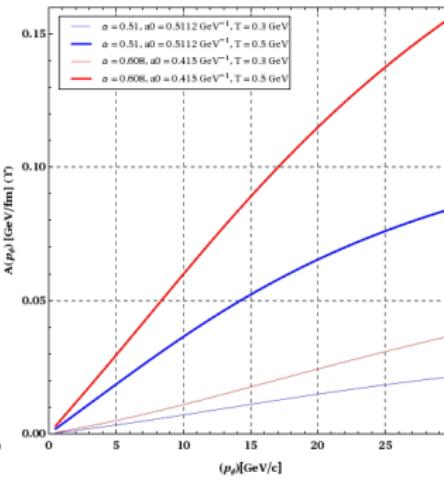
① **Heavy Quark**



② **Quarkonia (J/ψ)**



③ **Quarkonia (Υ)**



Diffusion coefficient ($HQ, J/\psi, \Upsilon$)

① Diffusion Coefficient, for: $\phi(M, E, p) \rightsquigarrow "i" (m, e, q)$

$$B(E) = \int_E^{+\infty} dE' A_i(E') \times \frac{E'}{P'} e^{-(E'-E)/T}, \text{ with: } B_{\perp} = B_{\parallel} = B, \quad B \leftrightarrow A \text{ relation ,}$$

- Fokker-Planck equation : $\frac{\partial f}{\partial t} = \frac{\partial}{\partial p_i} \left(A_i f + \frac{\partial}{\partial p_i} B_{ij} f \right) = -\vec{\nabla}_p \cdot \vec{p},$ (homogenous background)
- Einstein relation : $[\vec{A}f + \vec{\nabla}_p(Bf)]_i = 0, \quad f = e^{-E/T},$ (stationary, **D.B Walton, J.Rafelski(99)**)

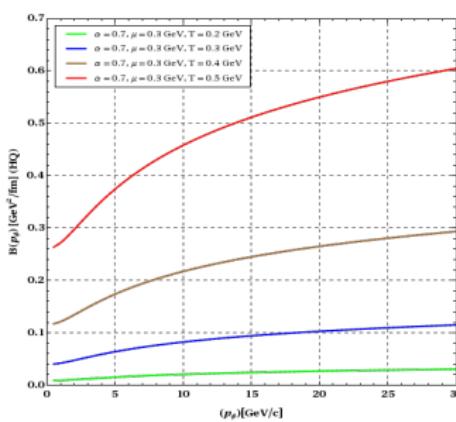
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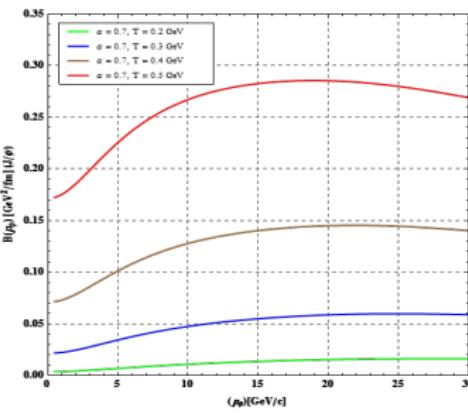
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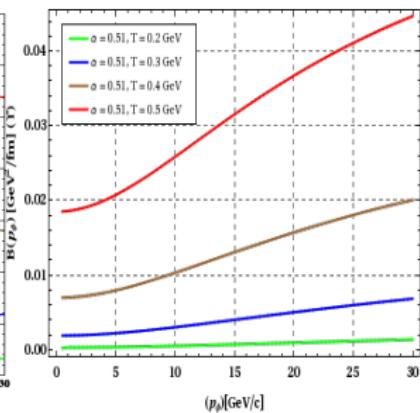
① Heavy Quark



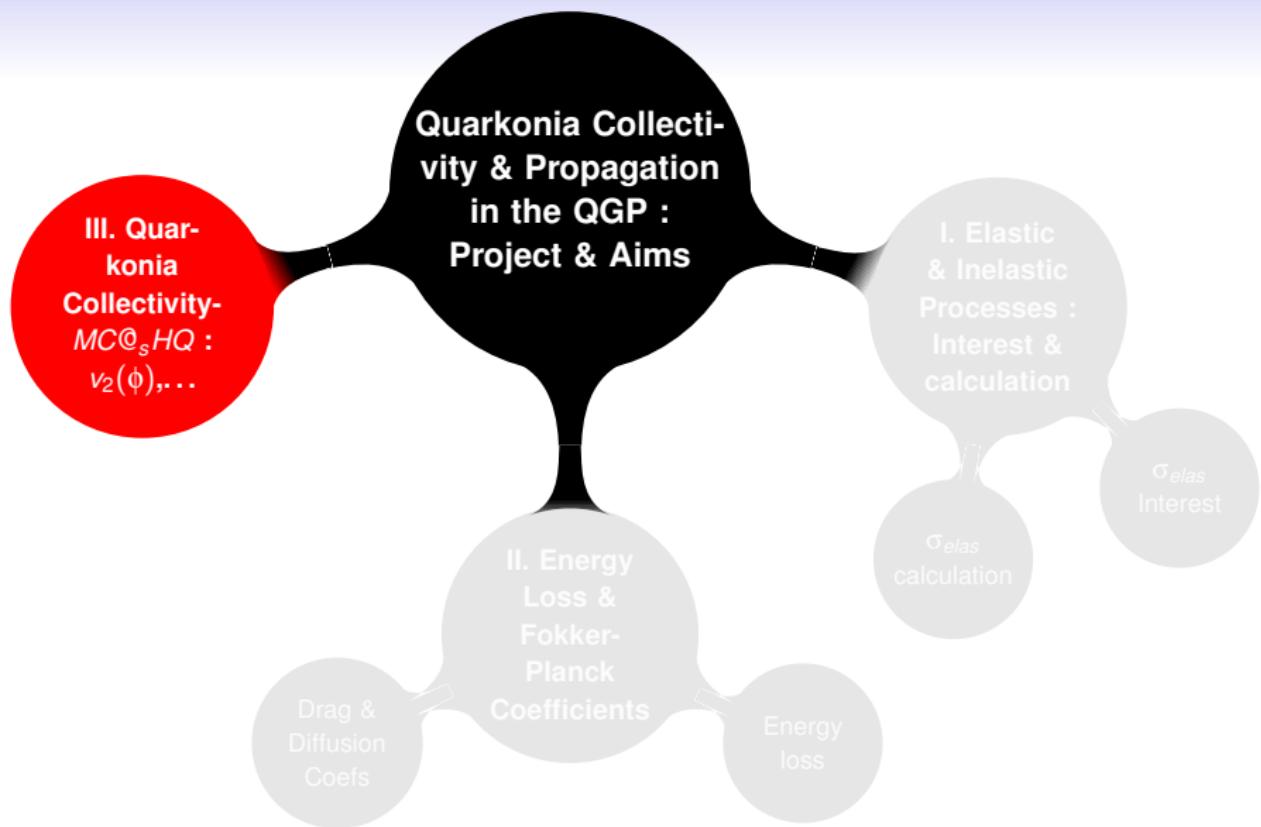
② Quarkonia (J/ψ)



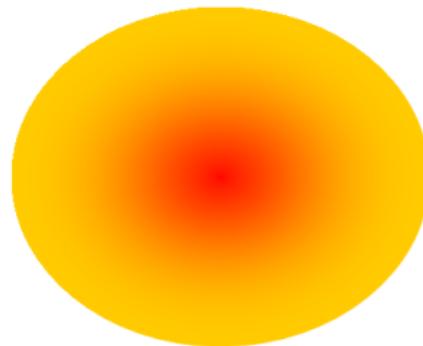
③ Quarkonia (Υ)



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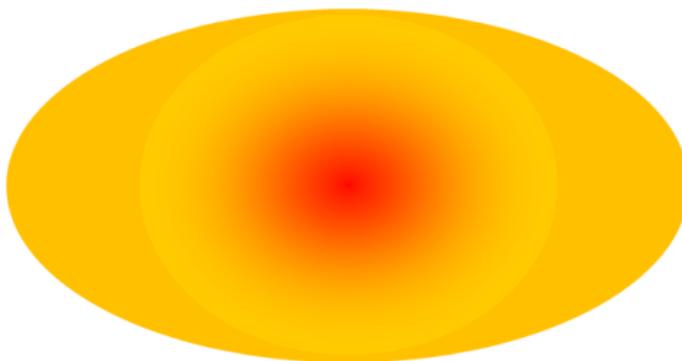


$MC@sHQ$ - as a cartoon



: QGP

$\mathcal{MC@sHQ}$ - as a cartoon

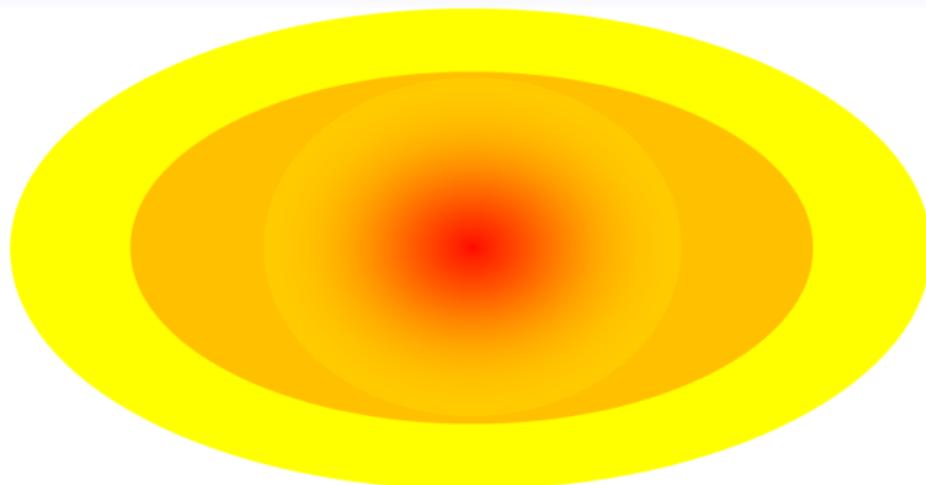


: QGP



: Mixed phase

$MC@sHQ$ - as a cartoon



: QGP

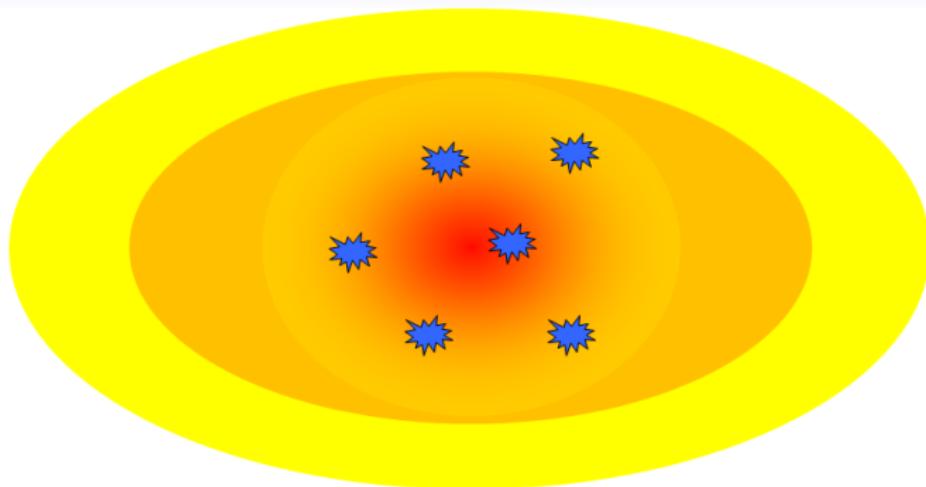


: Mixed phase



: Hadronic phase (not taken into account)

$\mathcal{MC@sHQ}$ - as a cartoon



: QGP



: hard collisions in initial NN collisions

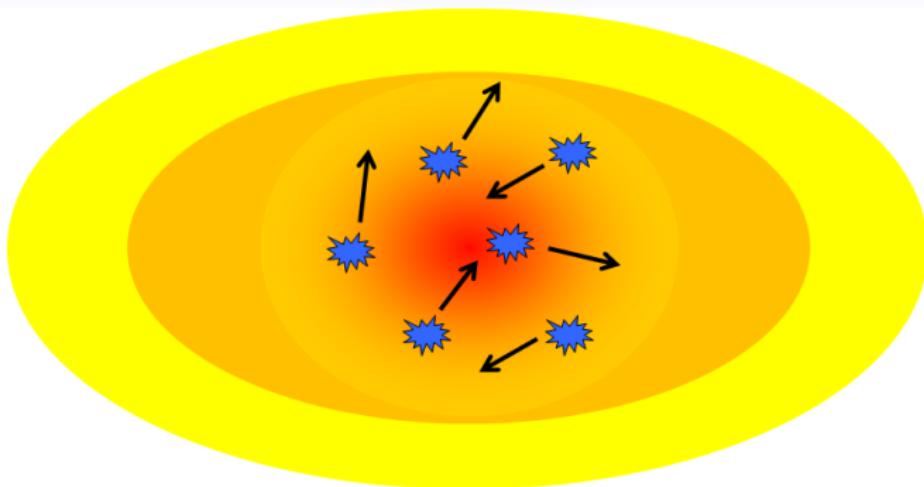


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: Hadronic phase (not taken into account)

$MC@sHQ$ - as a cartoon



: QGP



: Mixed phase



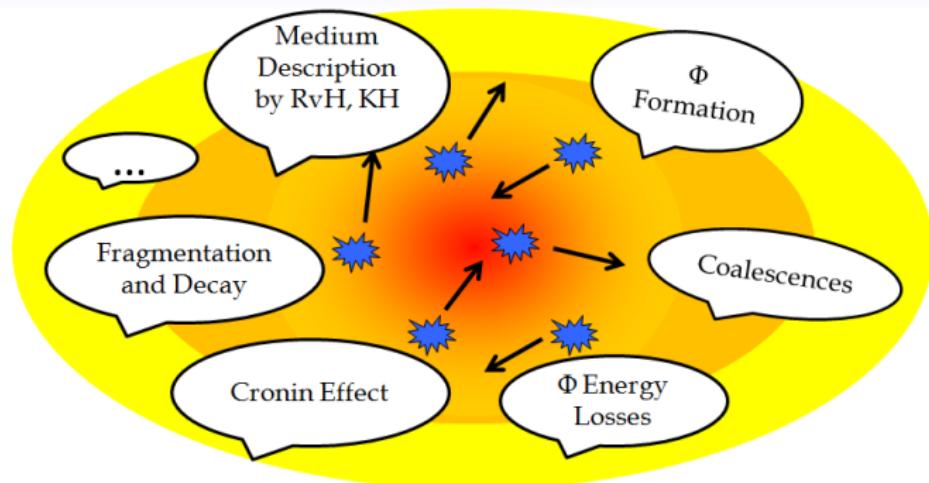
: Hadronic phase (not taken into account)



: hard collisions in initial NN collisions

: Φ evolution according to
Fokker-Planck

$MC@sHQ$ - as a cartoon



: QGP



: Mixed phase



: Hadronic phase (not taken into account)



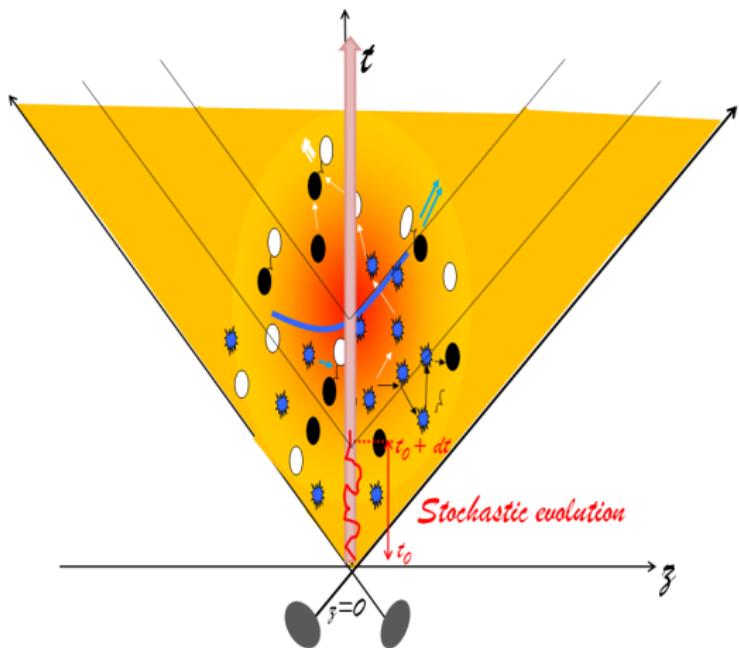
: hard collisions in initial NN collisions



: Φ evolution according to
Fokker-Planck

$MC@sHQ$ -Keywords

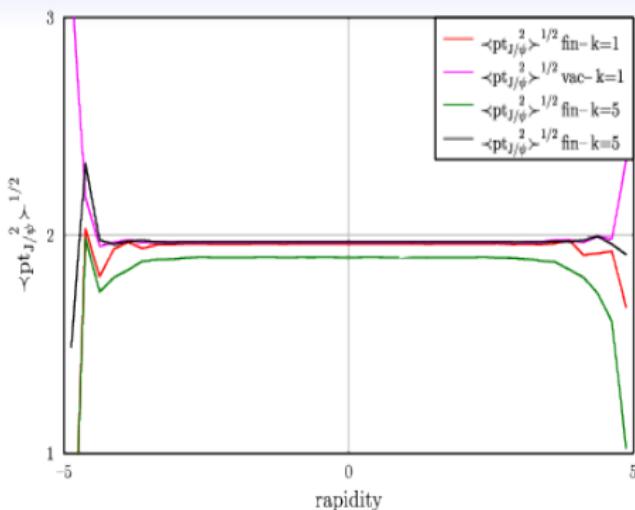
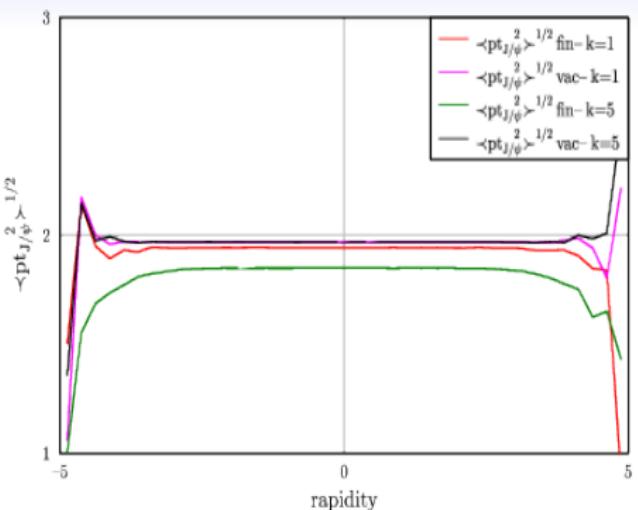
① Stochastic evolution of $Q\bar{Q}$



② $MC@sHQ$ in few words

- **Stochastic Evolution** : Boltzman evolution or FP evolution
- **J/ ψ suppression** :- by T, - cold nuclear effects, -Inelastic processes.
- Au-Au/Pb-Pb collisions at $\sqrt{s} = 200$ GeV or $\sqrt{s} = 5.5$ TeV
- Coalescence and fragmentation mechanism ...
- Plasma type and ...
- **study many observables** : $v_2(\phi)$, $\langle p_t^2 \rangle$, $\frac{dN}{d\phi_{rel}}$, R_{AA} , p_t spectra, y spectra, ϕ -spectra ...

Mean $\langle p_t^2 \rangle^{1/2}$ of J/ψ vs y



- Au-Au, $\sqrt{s} = 200\text{GeV}$, Central collisions

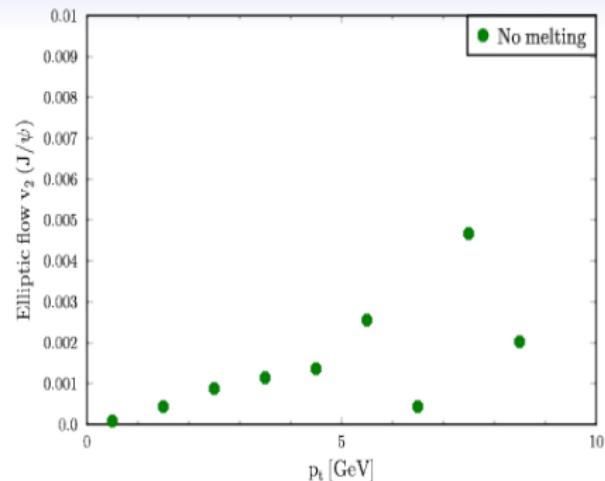
- No Melting, tuning factor $k=1,5$.

- Au-Au, $\sqrt{s} = 200\text{GeV}$, Central collisions

- Melting ($T_c (J/\psi, \Upsilon) = 0.3\text{ GeV}$), $k=1,5$

- Interaction with the medium \Rightarrow reduce their mean $p_t \prec p_{tJ/\psi}^2 \succ^{1/2}$
- Less visible effect with Melting & more visible with $\nearrow \sigma_{elas}$

J/ψ Elliptic flow vs p_t at mid rapidity



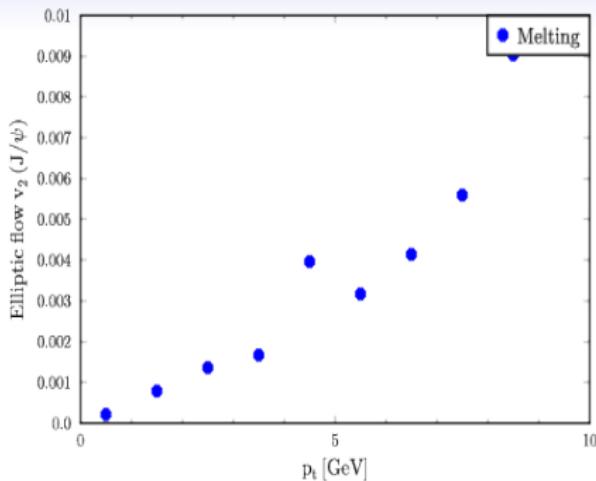
- Au-Au, $\sqrt{s} = 200 \text{ GeV}$, Min bias collisions

- No Melting, 50M collisions

- Non zero elliptic flow

- Dissociation by T $\Rightarrow v_2 \nearrow \dots$ Recombination ?

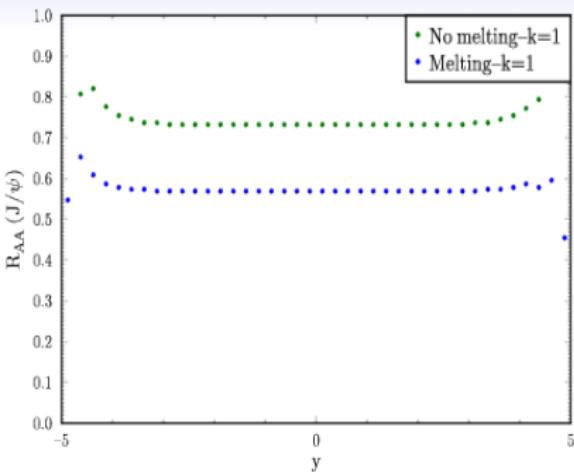
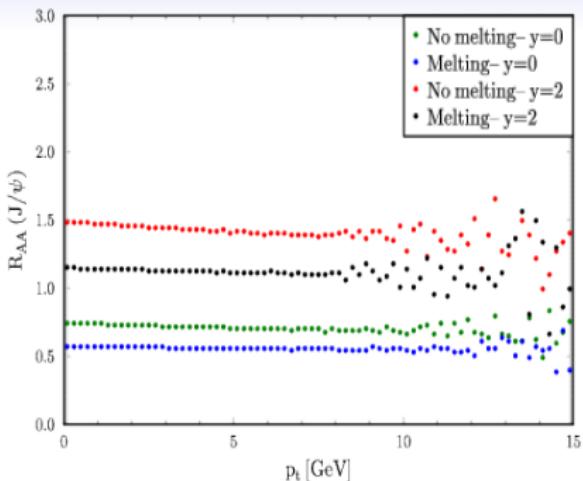
- Influence of elastic processes, increase of σ_{elas} $\Rightarrow v_2 \nearrow$



- Au-Au, $\sqrt{s} = 200 \text{ GeV}$, Min bias collisions

- Melting ($T_c (J/\psi, \Upsilon) = 0.3 \text{ GeV}$), 50M collisions

R_{AA} of J/ψ vs p_t at mid/average rapidity & vs y



- Au-Au, $\sqrt{s} = 200 \text{ GeV}$, Min bias collisions

- No Melting, 50M collisions

- Dissociation by T $\Rightarrow R_{AA}(J/\psi) \searrow$

- Dissociation by T $\Rightarrow R_{AA}(J/\psi)$ variation vs $p_t \approx$ small

- Influence of elastic processes, increase of $\sigma_{elas} \Rightarrow v_2 \nearrow$?

- Au-Au, $\sqrt{s} = 200 \text{ GeV}$, Min bias collisions

- Melting ($T_c(J/\psi, \Upsilon) = 0.3 \text{ GeV}$), 50M collisions

Quarkonia Collectivity & Propagation in the QGP : Conclusions

III. Quarkonia Collectivity-
 $MC@_sHQ$:
 $v_2(\phi), \dots$

- ϕ in hydro code $MC@_sHQ$
- effect of σ_{elas} in $\langle pt_{J/\psi}^2 \rangle$
- effect of σ_{elas} in v_2

Drag &
Diffusion
Coefs

II. Energy
Loss &
Fokker-
Planck
Coefficients

- pQCD ϕ Energy loss
- FP Coefs ↗ with T ↗

I. Elastic & Inelastic Processes : Interest & calculation

σ_{elas}
calculation

σ_{elas}
Interest

Energy
loss

- Develop BS Formalism
- pQCD σ_{elas} & σ_{inel} calculation
- Elastic process Interest

Quarkonia Collectivity & Propagation in the QGP : Perspectives

III. Quarkonia Collectivity-
 $MC@_sHQ$:
 $v_2(\phi), \dots$

- ϕ in hydro code $MC@_sHQ$
- Inclusion of recombination
- ...
- $J/\psi, \Upsilon \dots$ other observable

Drag &
Diffusion
Coefs

II. Energy
Loss &
Fokker-
Planck
Coefficients

- Elastic & inelastic interaction rate
- $E_{loss}, A_{Drag} (\phi - q) \dots$

Energy
loss

I. Elastic & Inelastic
Processes :
Interest &
calculation

σ_{elas}
calculation

σ_{elas}
Interest

- Investigate the role of ϕ wave function
- fine structure of BS vertex

Contents

1 Global Project : quarkonia propagation

2 σ_{elas} calculation & results

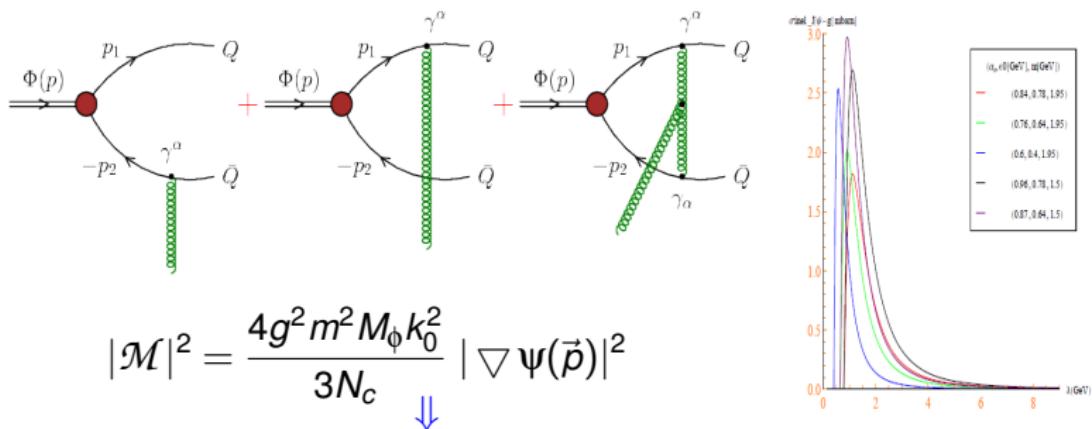
3 Energy loss & Fokker-Planck Coefficient

4 Quarkonia collectivity : $v_2(\phi)$

5 Conclusion & Outlook

σ_{inel} Results & Discussion

① Gluon Dissociation Process $J/\psi - g$



$$|\mathcal{M}|^2 = \frac{4g^2 m^2 M_\phi k_0^2}{3N_c} |\nabla \psi(\vec{p})|^2$$

↓

$$\sigma_{\phi g}(\lambda) = \frac{128g^2}{3N_c} a_0^2 \frac{(\lambda/\varepsilon_0 - 1)^{3/2}}{(\lambda/\varepsilon_0)^5}$$

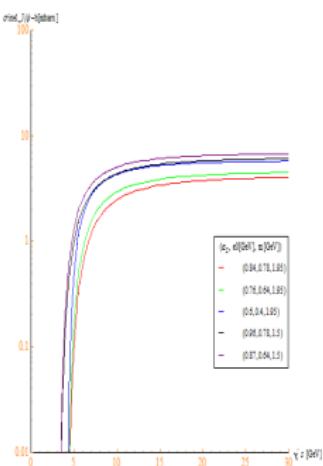
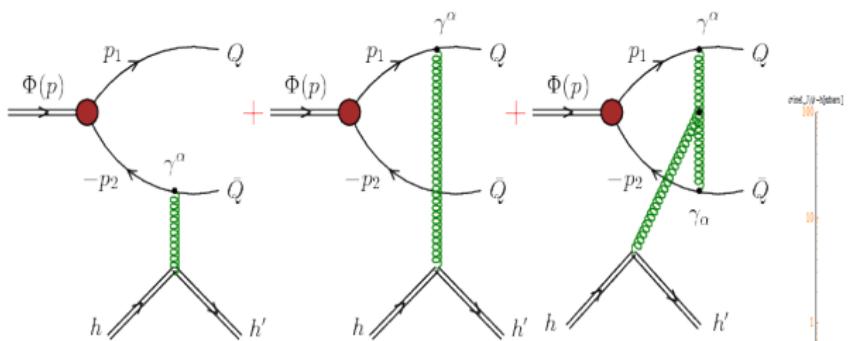
with : $\lambda = \frac{q \cdot k}{M_\phi} = \frac{s - M_\phi^2}{2M_\phi}$

- Dependence σ_{inel} vs ε et M

Oh, S. Kim, S. Hwang Lee, (02)

σ_{inel} Results & Discussion

② Hadron Dissociation Process $J/\psi - h$



Factorization Theorem

$$\sigma_{\phi h}(v) = \int_0^1 dx \sigma_{\phi g}(xv) \times g(x)$$

- Cross section of $J/\psi - h$ ●

- Cross section of $J/\psi - g$ ●

- $\mathcal{GDF}: g(x) = 0.5(\eta+1) \frac{(1+x)^\eta}{x}, \eta = 5(BP)$ ●

- Dependence σ_{inel} vs ϵ et m

Oh, S. Kim, S. Hwang Lee, (02)

σ_{elas} Results & Discussion

② Compton Diffusion Process $J/\psi - g$

② 2 gluons attached, "LO"

- Soft Gluons ($k \approx mg^4$): columbic case

$$\mathcal{M}(k, k' \approx mg^4) \approx -2\alpha g^2 \frac{\delta^{ab}}{2N_c} \varepsilon_{\lambda 1}(k_1) \cdot \varepsilon_{\lambda 2}(k_2)$$

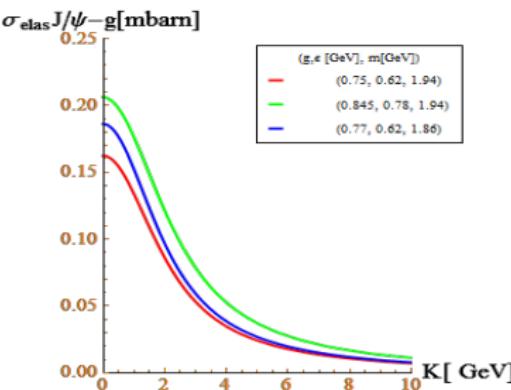
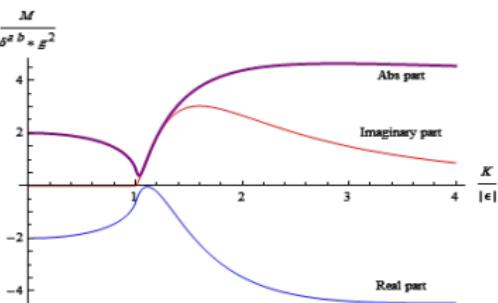
- Opening of the imaginary part for $K < |\varepsilon|$
- Opening of the imaginary part for $K = |\varepsilon|$

- Hard Gluons ($k \approx mg^2$)

$$\mathcal{M}(k, k' \equiv mg^2) \approx -\frac{4g^2 \delta_{ij} \delta^{ab}}{N_c} \times \frac{1}{\left(1 + \left(\frac{a_0 |k-k'|}{4}\right)^2\right)^2}$$

- Form Factor

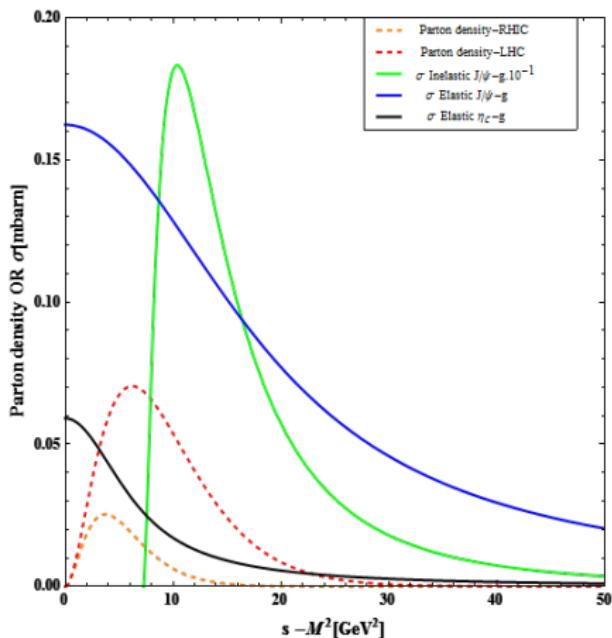
$$\Rightarrow \sigma_{elas} \approx \frac{g^4}{2\pi m^2 N_c^2} \times \frac{1 + \frac{K^2}{8m|\varepsilon|} + \frac{1}{2} \left(\frac{K^2}{8m|\varepsilon|}\right)^2}{\left(1 + \frac{K^2}{8m|\varepsilon|}\right)^3}$$



$\sigma_{elas}, \sigma_{inel}$ Results & Discussion

① $J/\psi - g$: Gluon Dissociation vs Compton Diffusion

"LO" Diagrams



$$R = \int d^3k n_{mb}(k) \sigma_{elas/inel}(s(k))$$

$$\rightarrow R = \int_{M^2}^{+\infty} ds \tilde{n}_{mb}(k) \sigma_{elas/inel}(s)$$

- → Parton density at LHC energy
- → Parton density at RHIC energy
- → σ Inelastic $J/\psi - g$
- → σ Elastic $J/\psi - g$
- → σ Inelastic $\eta_c - g$

Diffusion coefficient ($HQ, J/\psi, \Upsilon$)

① Diffusion Coefficient, for: $\phi(M, E, p) \rightsquigarrow "i" (m, e, q)$

$$B(E) = \int_E^{+\infty} dE' A_i(E') \times \frac{E'}{P'} e^{-(E'-E)/T}, \text{ with: } B_\perp = B_\parallel = B \quad B \leftrightarrow A \text{ relation ,}$$

- Fokker-Planck equation : $\frac{\partial f}{\partial t} = \frac{\partial}{\partial p_i} \left(A_i f + \frac{\partial}{\partial p_i} B_{ij} f \right) = -\vec{\nabla}_p \cdot \vec{p},$ (homogenous background)
- Einstein relation : $[\vec{A}f + \vec{\nabla}_p(Bf)]_i = 0, f = e^{-E/T},$ (stationary case) ○

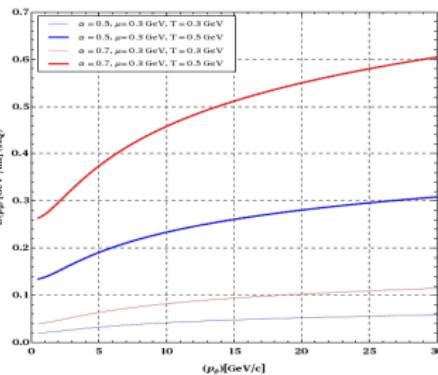
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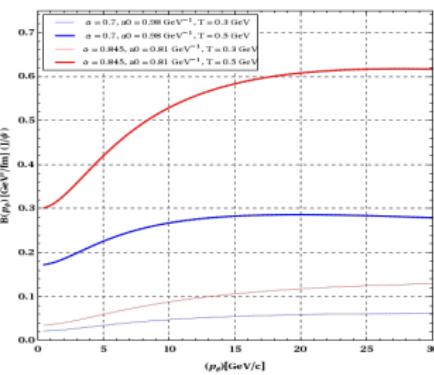
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① Heavy Quark



② Quarkonia (J/ψ)



③ Quarkonia (Υ)

