

# Nuclear structure in strong magnetic fields: nuclei in the crust of a magnetar

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# Objective

Objective:

- What is the effect of the magnetic field on the composition of a neutron star?

Previous studies:

- Several studies on the EOS
- No studies for the pasta phase
- Only qualitative results on nuclei

# Objective

- ① What is the minimum field that is able to significantly alter the nuclear structure?
- ② Is this field low enough to be found in a significant proportion of neutron stars or magnetars?
- ③ Is this effect big enough to influence astrophysically relevant situations and processes, e.g. neutron star outer crust composition or final element abundances in nucleosynthesis scenarios?

# Orders of Magnitude

$\approx 0.5$  G earth's magnetic field

$\approx 10^4$  G magnetic resonance

$\approx 10^9$  G atomic cigars

$4 \cdot 10^{13}$  G electron critical field

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$10^{12} - 10^{14}$  G neutron stars

$\approx 10^{16}$  G largest observed in magnetars

$10^{17} - 10^{18}$  G largest in magnetar (theory)

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$\approx 10^{20}$  G proton critical field

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# The model: overview

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- Simplest self-consistent formulation
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Consequences of the breaking of TR symmetry by the magnetic field

- Currents
- Deformation

# The model: details

Standard RMF Lagrangian:

$$\mathcal{L} = \mathcal{L}_N + \mathcal{L}_m + \mathcal{L}_i$$

$\mathcal{L}_i$  includes two new terms:

- Coupling to orbital motion  
(only protons)

$$\mathcal{L}_{BO} = -e\bar{\psi}\frac{1}{2}(1 - \tau_3)\gamma^\mu A_\mu^{(e)}\psi$$

- Coupling to magnetic moments  
(protons and neutrons)

$$\mathcal{L}_{BM} = -i\tau_3\bar{\psi}\chi_{\tau_3}^{(e)}\psi$$

$$\chi_{\tau_3}^{(e)} = \frac{1}{4}|\kappa_{\tau_3}|\gamma_0\sigma_{\mu\nu}F^{(e)\mu\nu}$$

$$\kappa_{\tau_3} = \frac{1}{2}\mu_N g_{\tau_3}$$

$$\sigma_{\mu\nu} = [\gamma_\mu, \gamma_\nu]$$



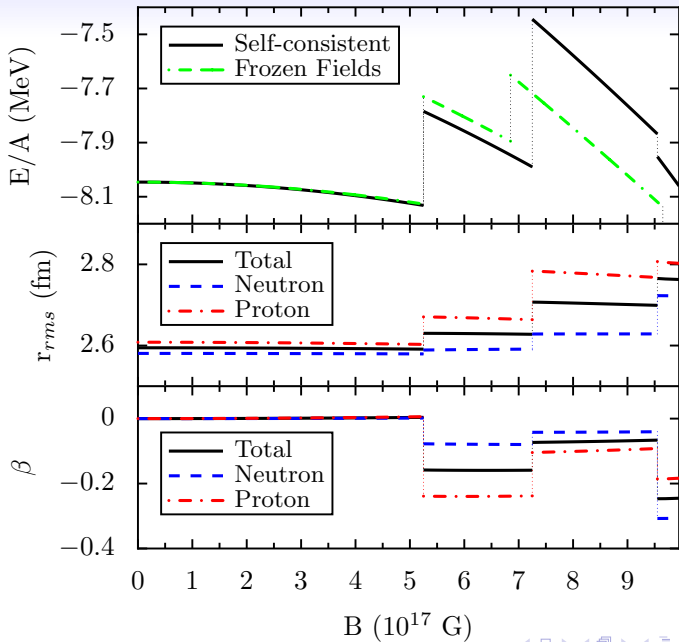
# The model: details

- Energy density functional

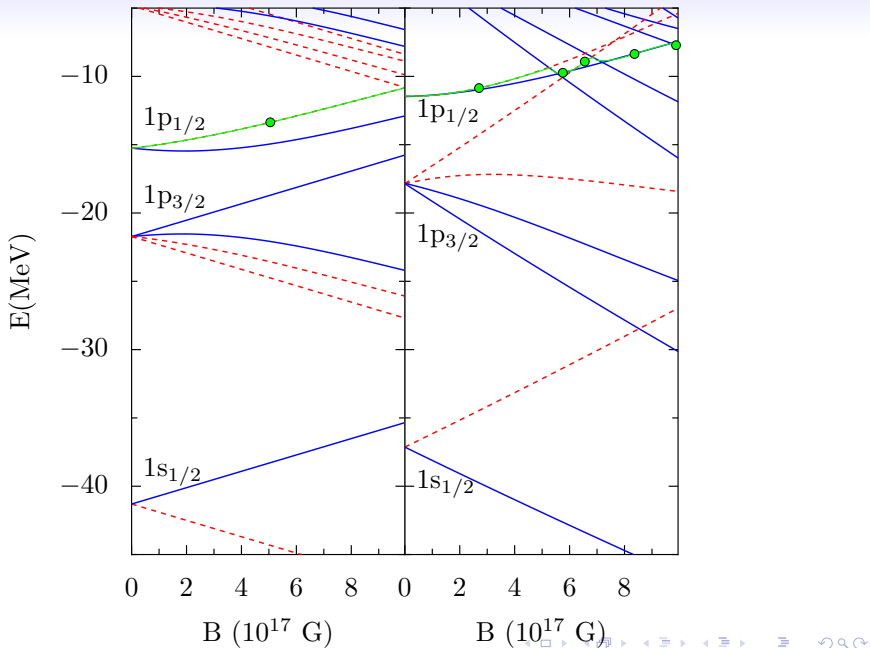
$$\begin{aligned} E_{\mathbf{B}}[\hat{\rho}, \phi] &= \text{Tr} \left[ \left( \boldsymbol{\alpha} \left( -i\nabla - e\mathbf{A}^{(e)} \right) + \beta(m + \chi_{\tau_3}^{(e)}) \right) \hat{\rho} \right] \\ &+ \sum_m \text{Tr} [(\beta \boldsymbol{\Gamma}_m \phi_m) \hat{\rho}] \\ &\pm \sum_m \int d^3r \left[ \frac{1}{2} (\partial_\mu \phi_m)^2 + \frac{1}{2} m_m^2 \phi_m^2 \right], \end{aligned}$$

- Constant magnetic field  $\mathbf{B}$  along positive z-axis
- Axial deformation, symmetry axis along  $\mathbf{B}$
- Symmetric gauge  $\rightarrow \mathbf{A} = i \left( -\frac{r|\mathbf{B}|}{2} e^{-i\theta}, +\frac{r|\mathbf{B}|}{2} e^{i\theta}, 0 \right)_{\text{stb}}$

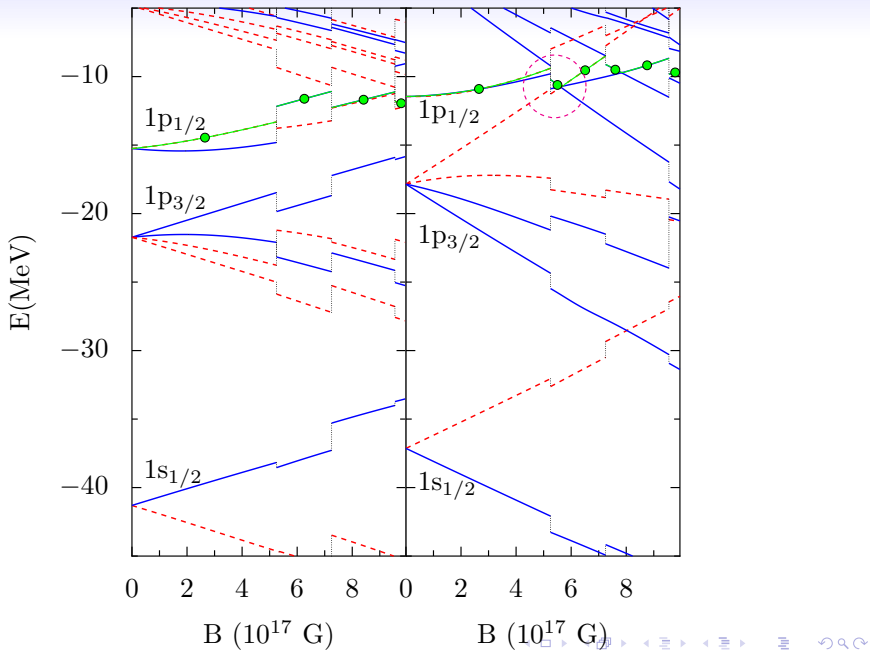
# Example: $^{16}\text{O}$

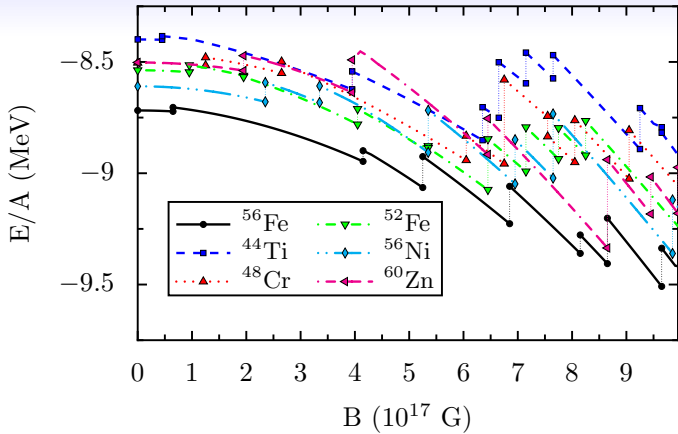


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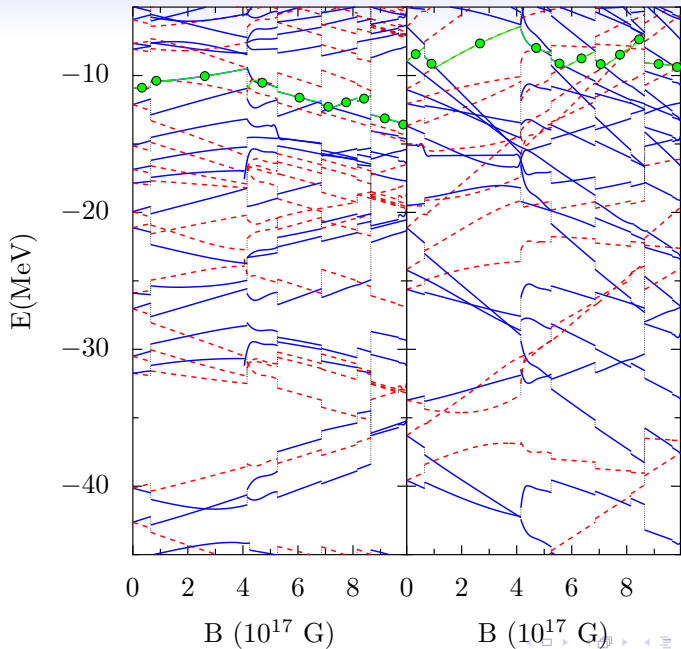


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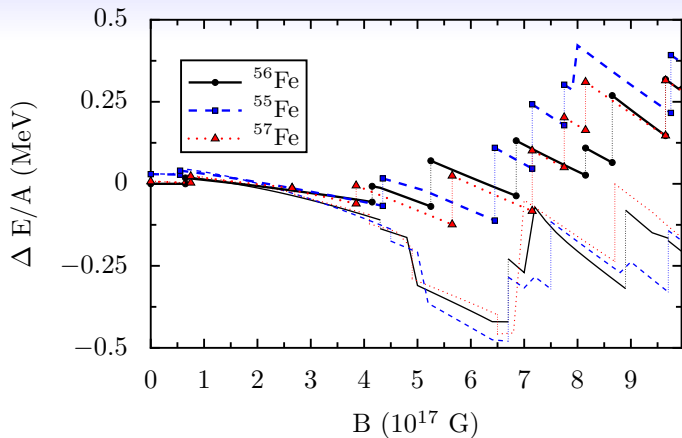


$^{56}\text{Fe}$ 

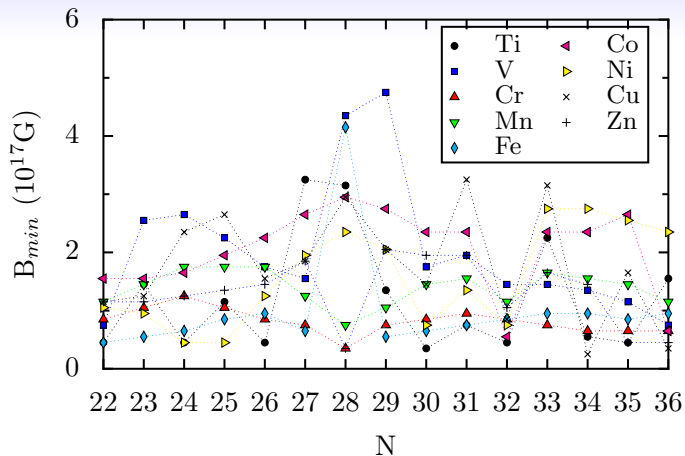
# Sample of crust nuclei around $^{56}\text{Fe}$



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# Minimum Magnetic Field





# Conclusions

- Minimum fields are of the order of  $10^{17}$ G
- Require a self-consistent formulation
- Minimum field depends a lot on the considered nucleus
- Changes in BE: from few tenths of keV to hundreths of keV