

Institut de
Fecherche sur les lois
fondamentales de

### Unité mixte de recherche 7095 : CNRS - Université Pierre et Marie Curie

# GRB spectral evolution in the internal shock model: confrontation with Fermi observations

Bosnjak Z.

CEA - Saclay

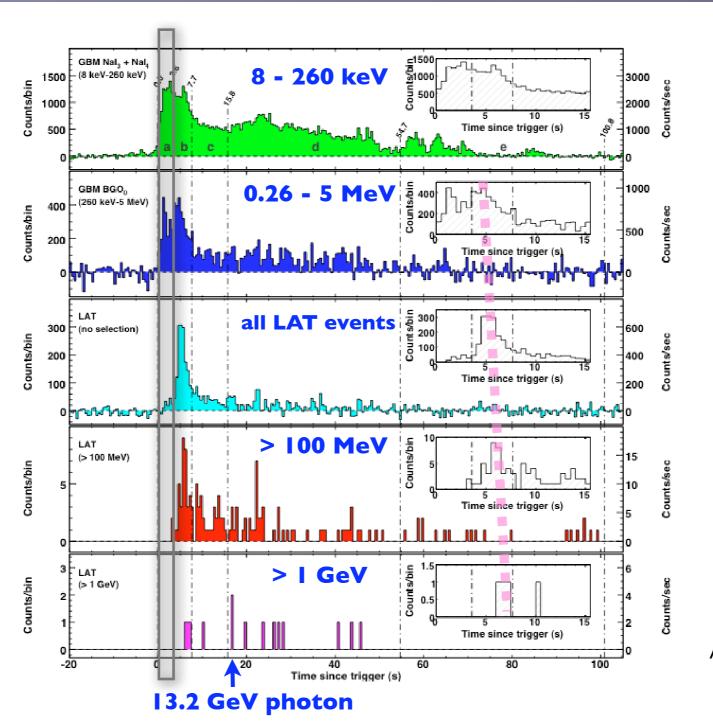
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Dubus G.

Laboratoire d'Astrophysique de l'Observatoire de Grenoble

## High energy gamma ray emission

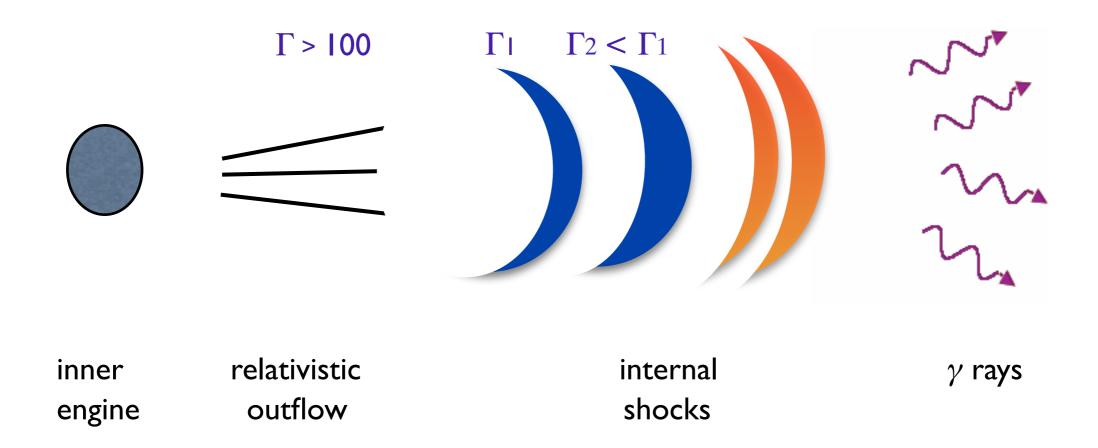


GRB 080916C Abdo et al. 2008

#### Fermi observations:

- √ Delayed onset of high energy (>100 MeV) emission (5/11 of LAT GRBs)
- √ Long lived high energy emission (9/11)
- ✓ Deviation from the usual GRB spectral models: extra component (3/11)

### Modeling: prompt high energy emission in the framework of IS

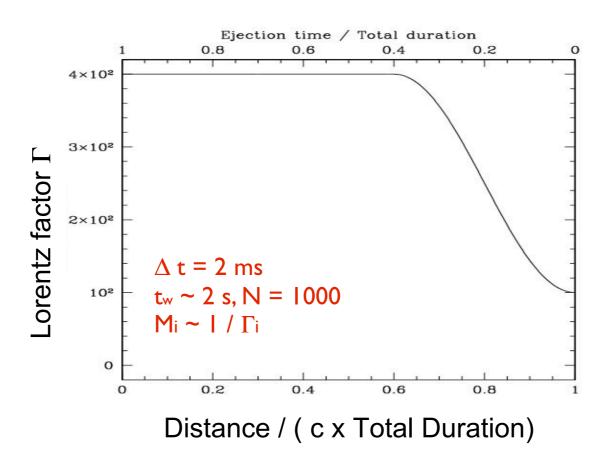


### Modeling:

- I. dynamics of internal shocks
- 2. radiative processes in the shocked medium
- 3. observed spectra and time profiles

### Dynamics of the internal shocks

Input parameters: distribution of Lorentz factors  $\Gamma(t)$ , kinetic energy rate dE/dt during the relativistic ejection, total duration of the ejection phase tw



R IS, start 
$$\sim \Gamma^2 c \text{ tvar} \sim 3 \times 10^{11} \text{cm} (\Gamma/100)^2 (\text{tvar} / 1 \text{ ms})$$
  
R IS, end  $\sim \Gamma^2 c \text{ tw} \sim 3 \times 10^{15} \text{cm} (\Gamma/100)^2 (\text{tw} / 10 \text{ s})$ 

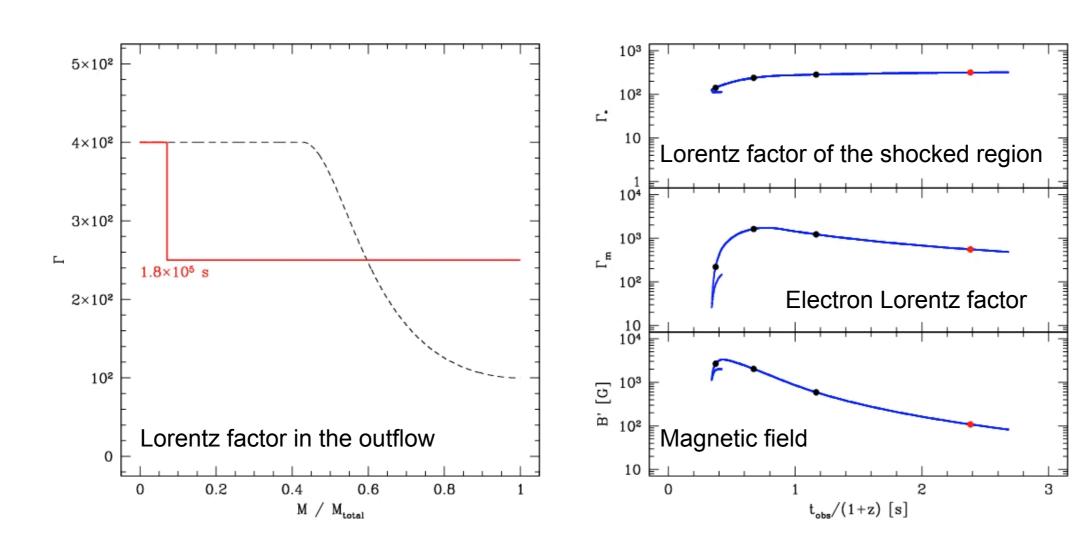
Dissipated energy: from 6% ( $\Gamma_2 / \Gamma_1 = 2$ ) to 43 % ( $\Gamma_2 / \Gamma_1 = 10$ )

#### Daigne & Mochkovitch 2000:

the simplified approach for the dynamics has been confirmed by a comparison with a full hydrodynamical calculation

### Dynamics of the internal shocks

Input parameters: distribution of Lorentz factors  $\Gamma(t)$ , kinetic energy rate dE/dt during the relativistic ejection, total duration of the ejection phase tw



### Dynamics of the internal shocks

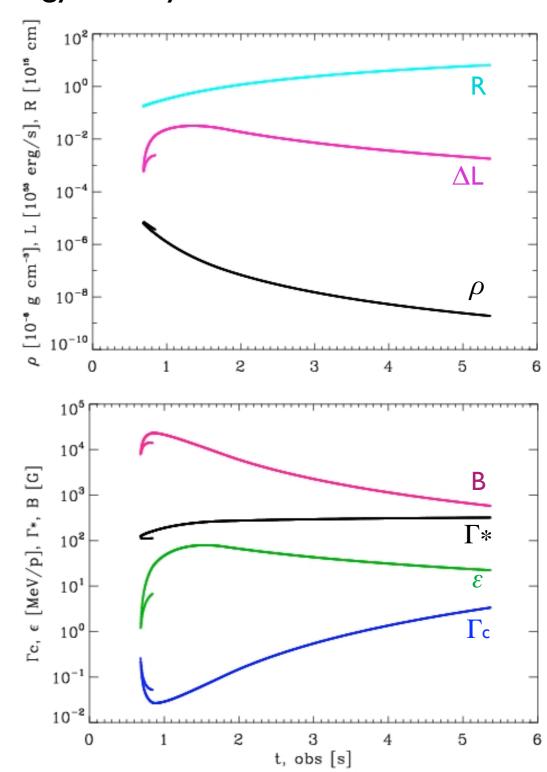
Physical conditions in the shocked medium: Lorentz factor  $\Gamma^*$ , comoving density  $\rho^*$ , comoving specific energy density  $\varepsilon^*$ 

Dissipated energy is distributed between protons, electrons (fraction  $\varepsilon_e$ ) and magnetic field (fraction  $\varepsilon_B$ )

Relativistic electron density:

$$n'(\Gamma_e, t'=0) \propto \Gamma_e^{-p} \qquad \Gamma_e \geq \Gamma_m$$

 $\zeta$  < 1 of all electrons is accelerated into a powerlaw distribution with slope -p (e.g. Bykov & Meszaros 1996)



Assumption: instantaneous shock acceleration

\* Adiabatic cooling timescale:  $t'ex = R / \Gamma^* c$  (comoving frame)

\* Radiative timescale: t`rad

t'rad << t'ex high radiative efficiency

Electron and photon distributions evolve strongly with time!

Assumption: instantaneous shock acceleration

- \* Adiabatic cooling timescale:  $t'ex = R / \Gamma^* c$  (comoving frame)
- \* Radiative timescale: t`rad

t`rad << t`ex high radiative efficiency

Electron and photon distributions evolve strongly with time!

The present version of the code follows the time evolution of the electron density n'e ( $\Gamma'e$ , t') and the photon density  $n'\nu'$  (t') including the following processes:

- adiabatic cooling (spherical expansion)
- synchrotron
- inverse Compton
- synchrotron self-absoprtion
- \*  $\gamma\gamma$  annihilation

Not included:

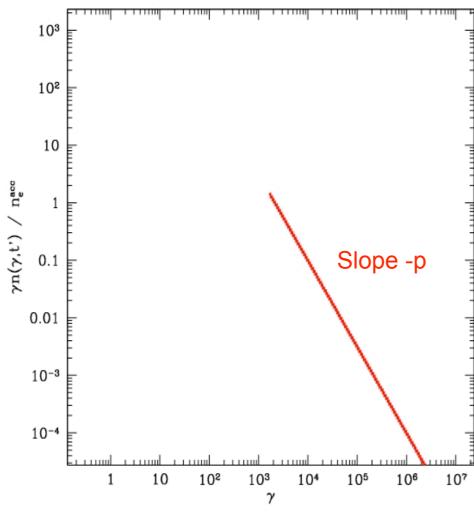
- \* emission from secondary leptons
- \* IC in optically thick regime (Comptonization)

#### **ELECTRONS:**

$$\frac{\partial n'}{\partial t'}(\Gamma'_{e},t') = -\frac{\partial}{\partial \Gamma'_{e}} \left[ \left( \frac{d\Gamma'_{e}}{dt'} \Big|_{syn+ic} + \frac{d\Gamma'_{e}}{dt'} \Big|_{ad} \right) n'(\Gamma'_{e},t') \right]$$

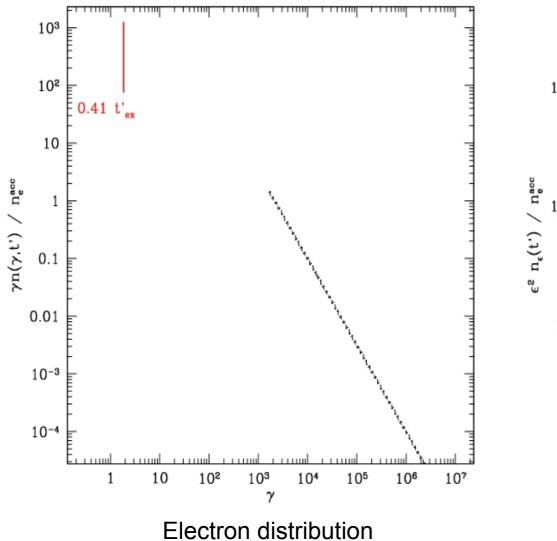
PHOTONS: 
$$\frac{\partial n'_{\mathbf{v}}}{\partial t'} = \int n'(\Gamma'_{e}, t') P_{syn+ic}(\Gamma'_{e}) d\Gamma'_{e} - cn'_{\mathbf{v}} \int n'(\Gamma'_{e}, t') \sigma_{abs}(\Gamma'_{e}, \mathbf{v}) d\Gamma'_{e} - cn'_{\mathbf{v}} \int_{\mathbf{v}' > \frac{(m_{e}c^{2})^{2}}{h^{2}\mathbf{v}}} n'_{\mathbf{v}'}(t') \sigma_{\gamma\gamma}(\mathbf{v}, \mathbf{v}') d\mathbf{v}'$$

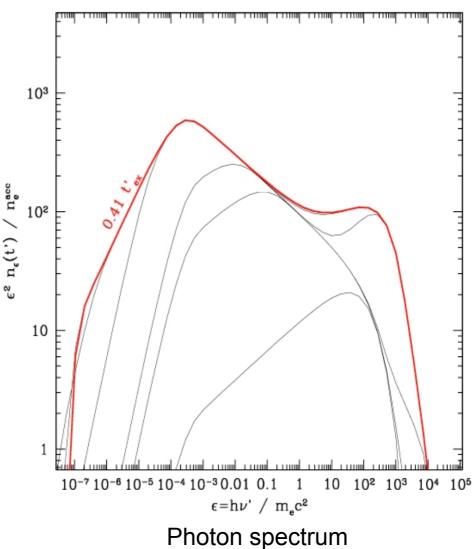
Radiation: the time evolution of electrons and photons in the comoving frame is solved (time-dependent radiative code)



**Electron distribution** 

Radiation: the time evolution of electrons and photons in the comoving frame is solved (time-dependent radiative code)





Radiation: the time evolution of electrons and photons in the comoving frame is solved (time-dependent radiative code)

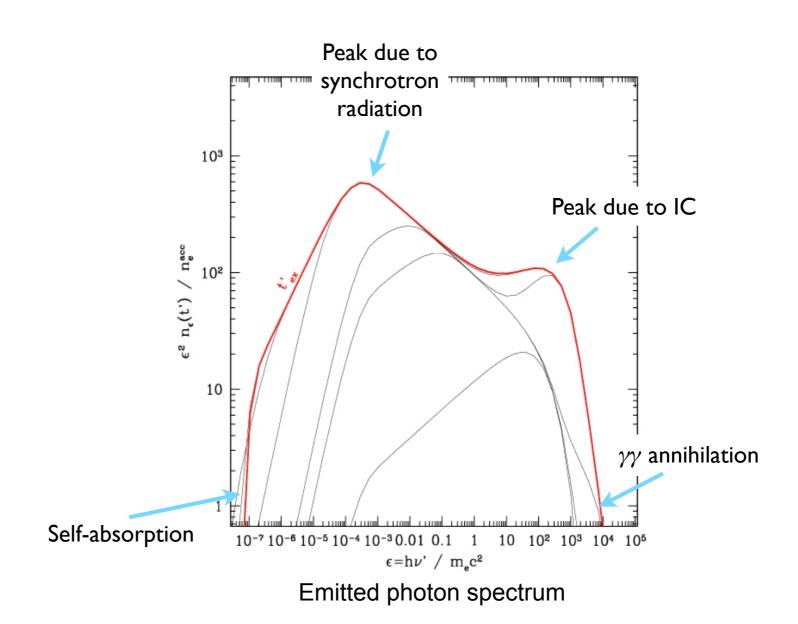
Comptonization parameter  $Y = L_{ic} / L_{syn}$ 

#### IC dominant:

low frequency synchrotron peak Thomson regime

#### Synchrotron dominant:

high frequency synchrotron peak Klein-Nishina regime



Radiation: the time evolution of electrons and photons in the comoving frame is solved (time-dependent radiative code)

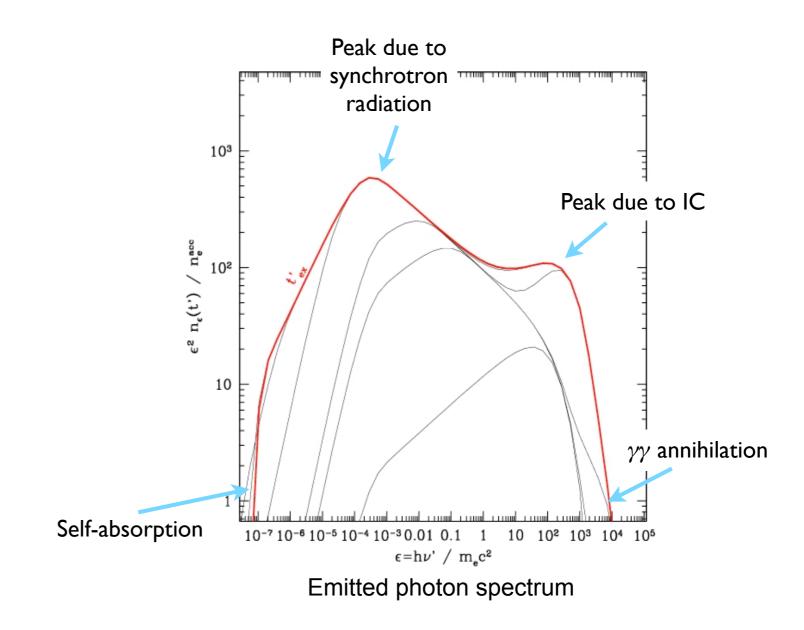
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#### IC dominant:

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high frequency synchrotron peak Klein-Nishina regime



This calculation is done at all times along the propagation of each shock wave

All the contributions are added together to produce a synthetic gamma-ray burst (spectrum+lightcurve)

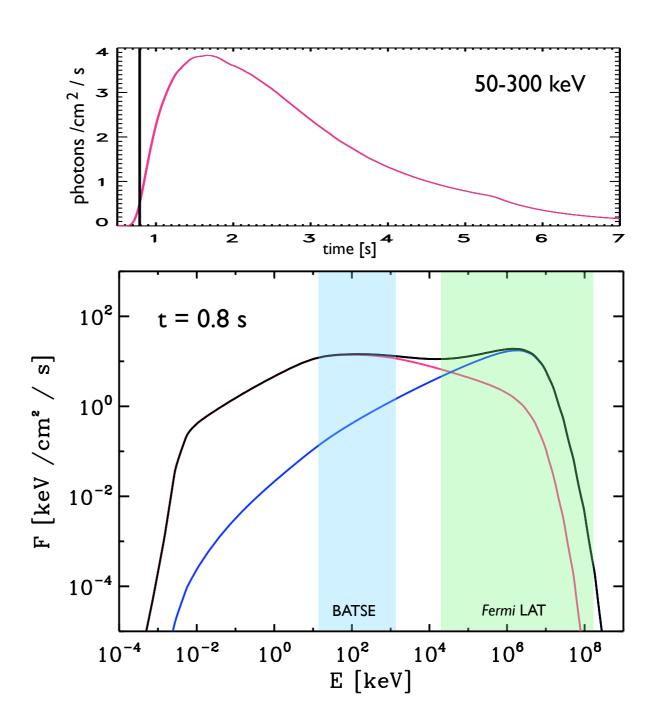
### Observed spectra and time profiles

The observed spectra and the light curves are computed from the comoving emission by integration over equal-arrival time surfaces.

- \* relativistic effects (Doppler factor)
- \* geometry (curvature of the emitting surface)
- \* cosmological effect (redshifts)

Instantaneous observed spectrum:

synchrotron inverse Compton total



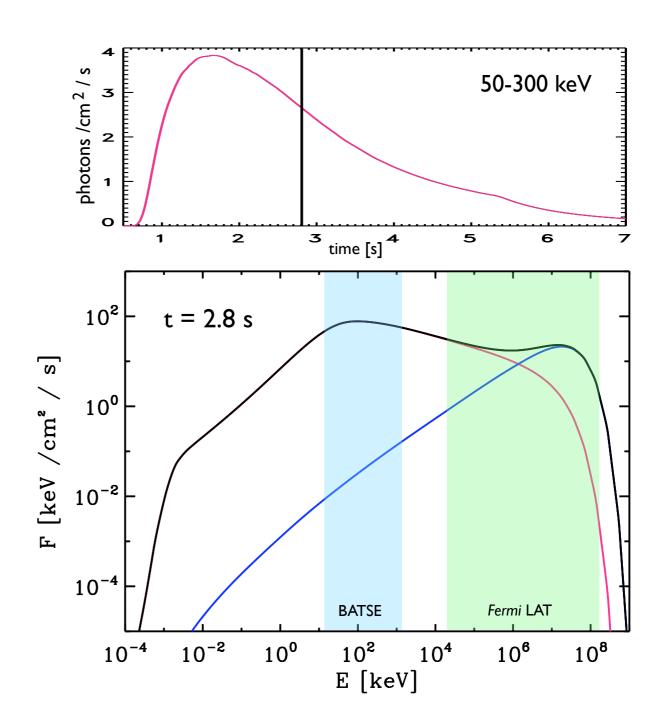
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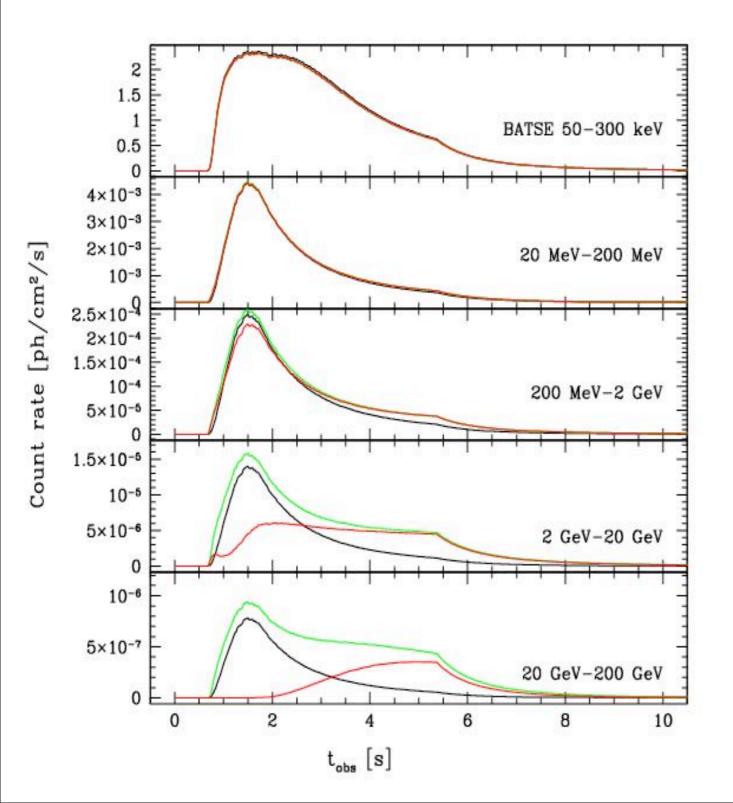
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Instantaneous observed spectrum:

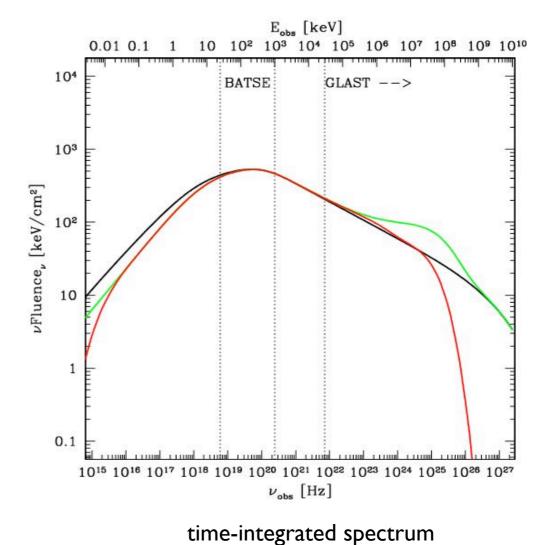
synchrotron inverse Compton total



### Observed spectra and time profiles



# synchrotron only synchrotron + IC synchrotron + IC + $\gamma\gamma$ annihilation + self absorption



2 possibilities for the dominant process in the keV-MeV range

1. SYNCHROTRON

2. INVERSE COMPTON

2 possibilities for the dominant process in the keV-MeV range

#### 1. SYNCHROTRON

High  $\Gamma_m$  requires that only a fraction of the electrons is accelerated (<10%)

High B: no IC component at high energy

Low B: IC component at high energy

#### 2. INVERSE COMPTON

2 possibilities for the dominant process in the keV-MeV range

#### 1. SYNCHROTRON

High  $\Gamma$ m requires that only a fraction of the electrons is accelerated (<10%)

High B: no IC component at high energy

Low B: IC component at high energy

#### 2. INVERSE COMPTON

All electrons are accelerated

Synchrotron component at low energy

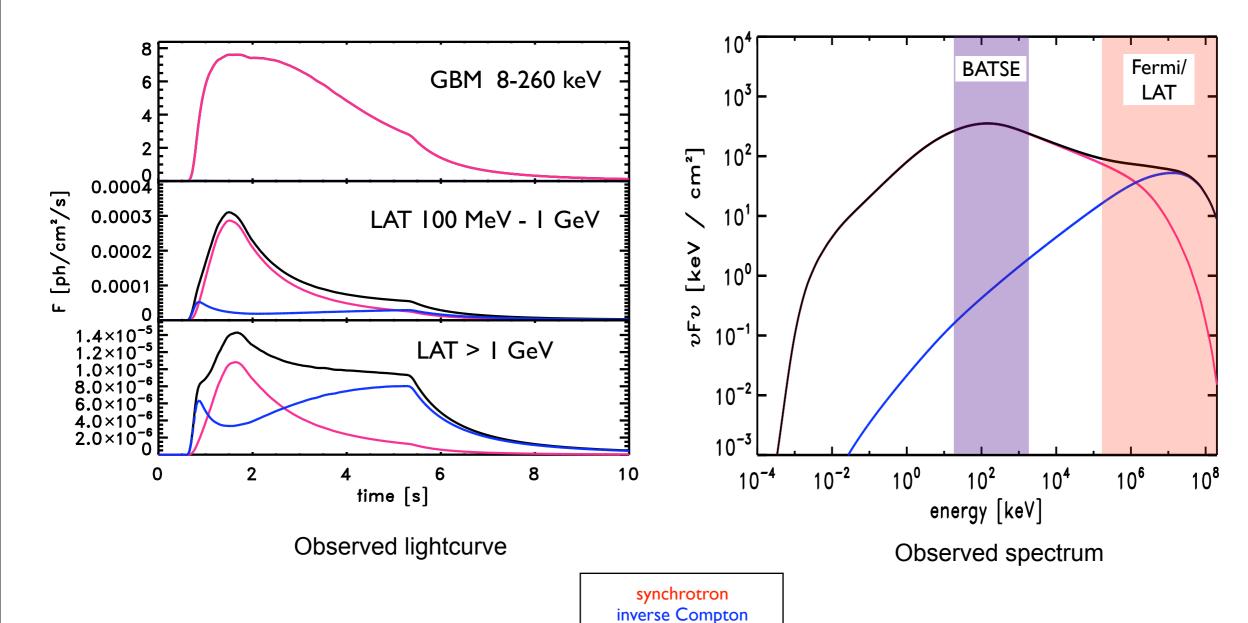
Second inverse Compton peak at high energy

A steep electron slope (p>3) is required to have two well defined peaks

#### SYNCHROTRON CASE

#### high magnetic field

dE/dt = 5 x 
$$10^{53}$$
erg s,  $\varepsilon_{B} = \varepsilon_{e} = 1/3$ ,  $\zeta = 0.003$ , p = 2.5, z=1

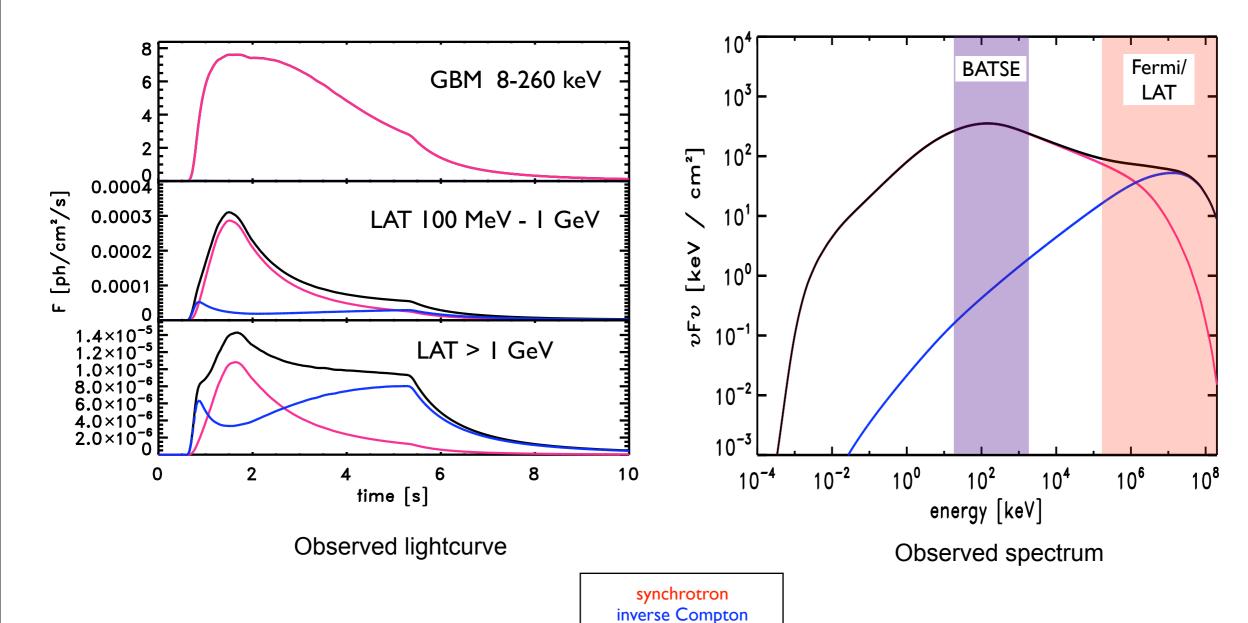


total

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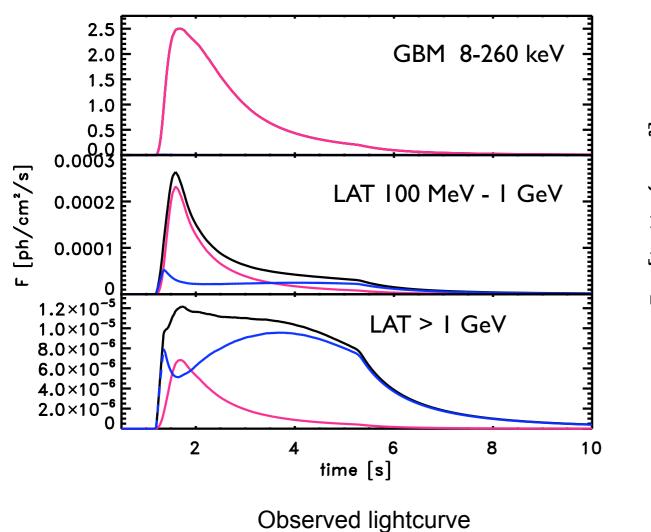


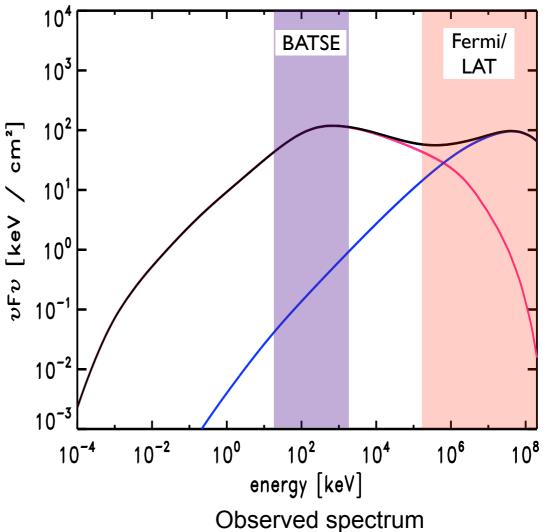
total

#### SYNCHROTRON CASE

#### low magnetic field

dE/dt = 5 x 
$$10^{53}$$
erg s,  $\varepsilon_{\rm B}$  = 0.003,  $\varepsilon_{\rm e}$  = 1/3,  $\zeta$  = 0.003, p = 2.5, z=1



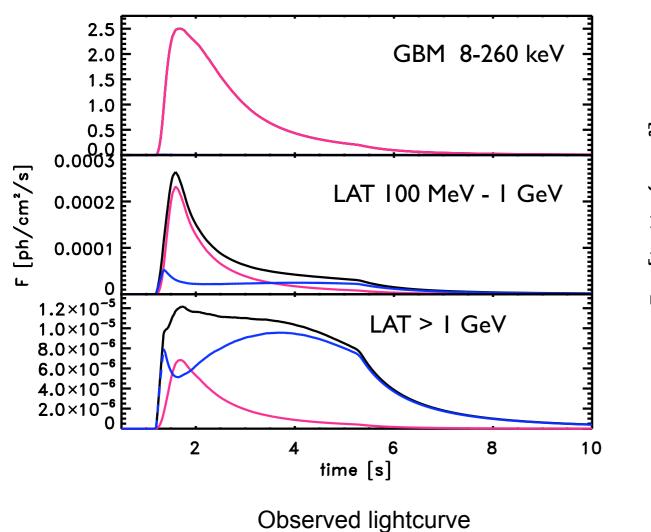


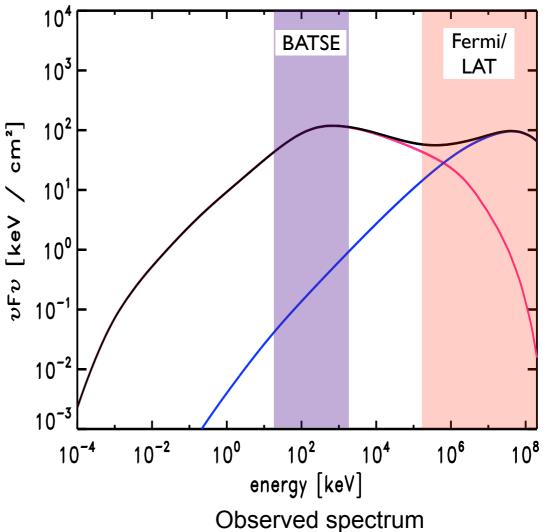
synchrotron inverse Compton total

#### SYNCHROTRON CASE

#### low magnetic field

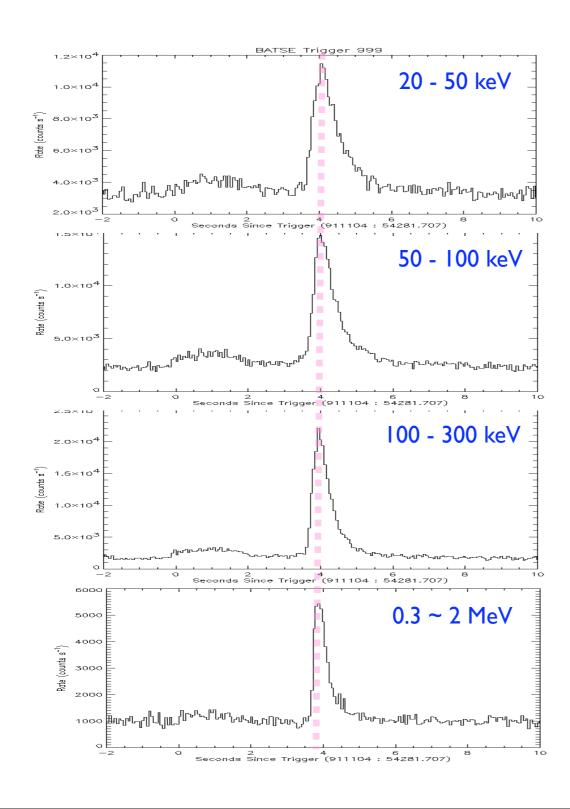
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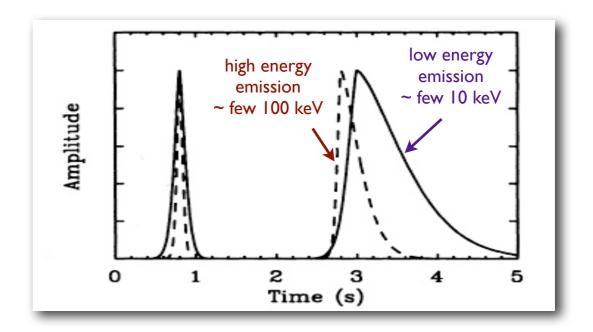




synchrotron inverse Compton total

# Emission processes and temporal profile: sub MeV range

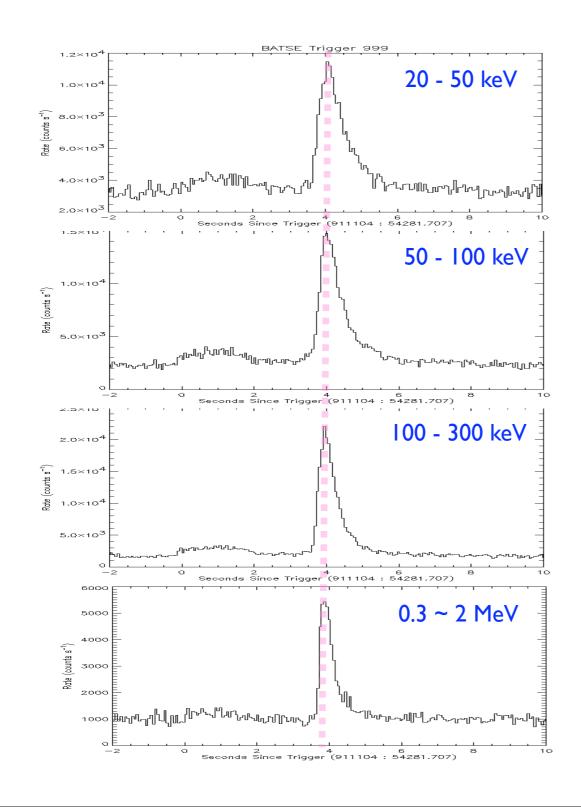


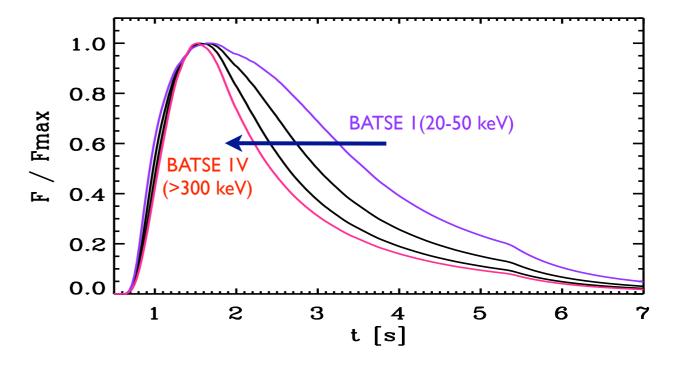


Norris et al. 1996 (BATSE GRBs): asymmetry/energy-shift paradigm

# Emission processes and temporal profile: sub MeV range

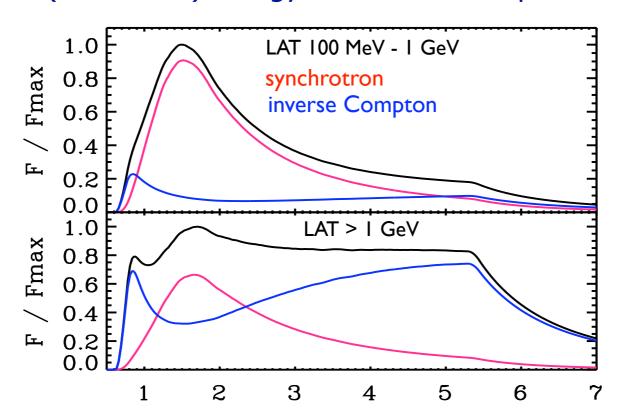
#### Model: synchrotron emission dominant in sub-MeV range





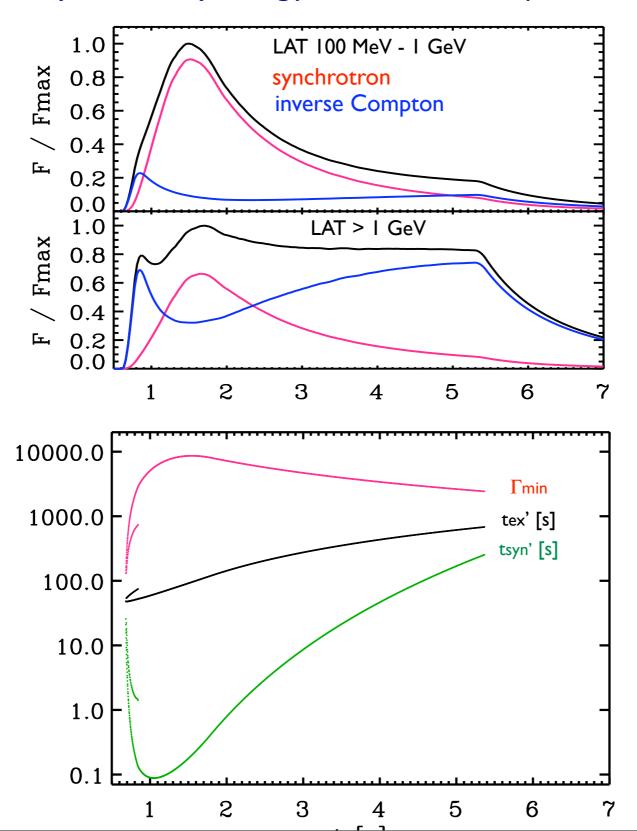
# Emission processes and temporal profile: >100 MeV bands

Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC



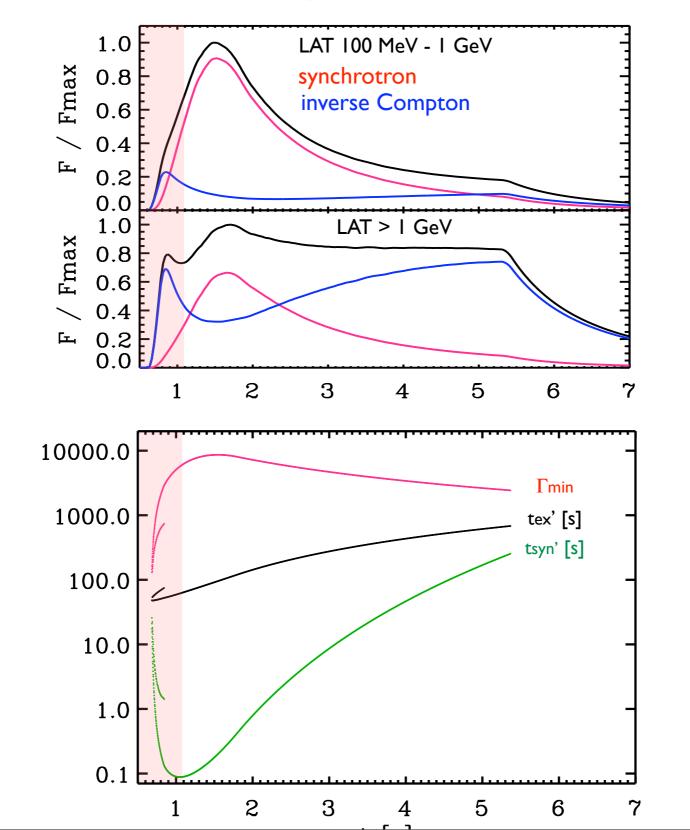
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## Emission processes and temporal profile: >100 MeV bands

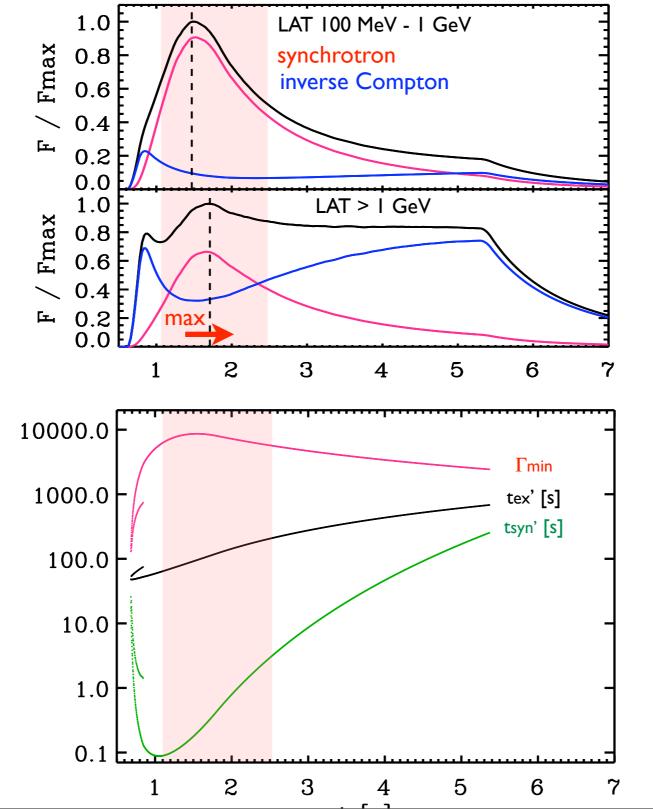
Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC



 $\begin{array}{c} \text{weak shock} \\ \epsilon^* \text{ low} \\ \text{moderate } \Gamma \text{m} \Rightarrow \text{large tsyn'} \\ \text{R small} \Rightarrow \text{tex'} \cong \text{R}/\Gamma^* \text{c small} \\ \text{tsyn'} \leq \text{tex'} \Rightarrow \text{large efficiency of IC} \end{array}$ 

# Emission processes and temporal profile

Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC



shock becomes stronger

I'm increases ⇒ tsyn' decreases

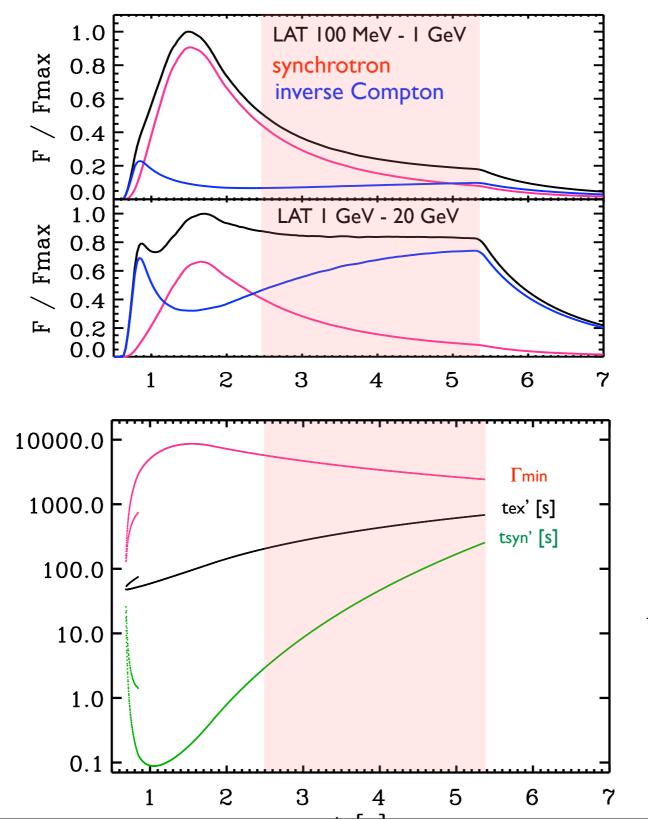
R, tex' increase

tsyn' << tex' ⇒ low efficiency of IC

dominant synchrotron component

## Emission processes and temporal profile

Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC



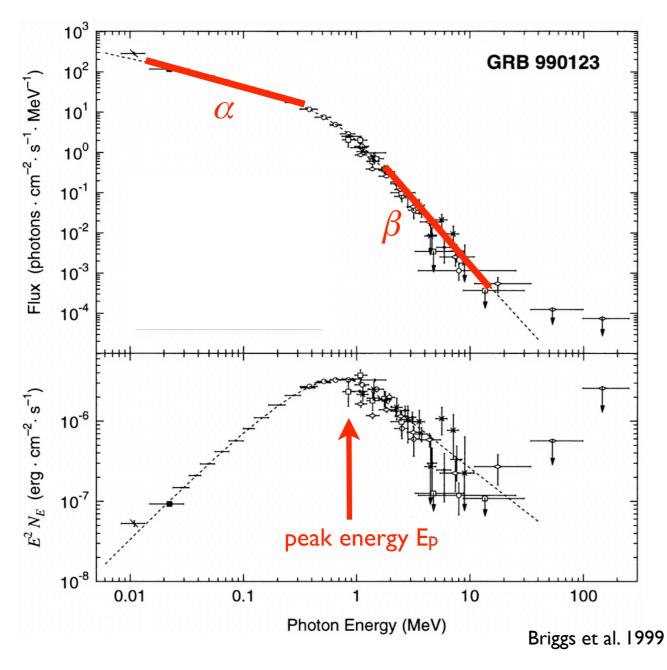
tail of the pulse:

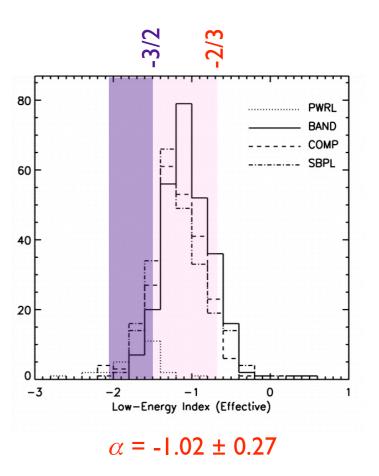
B decreases  $\Rightarrow$  tsyn' increases

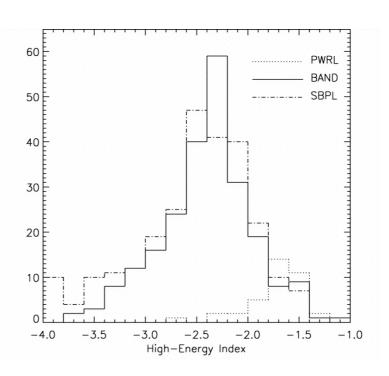
tsyn'  $\leq$  tex'  $\Rightarrow$  increased efficiency of IC

IC component dominant in GeV

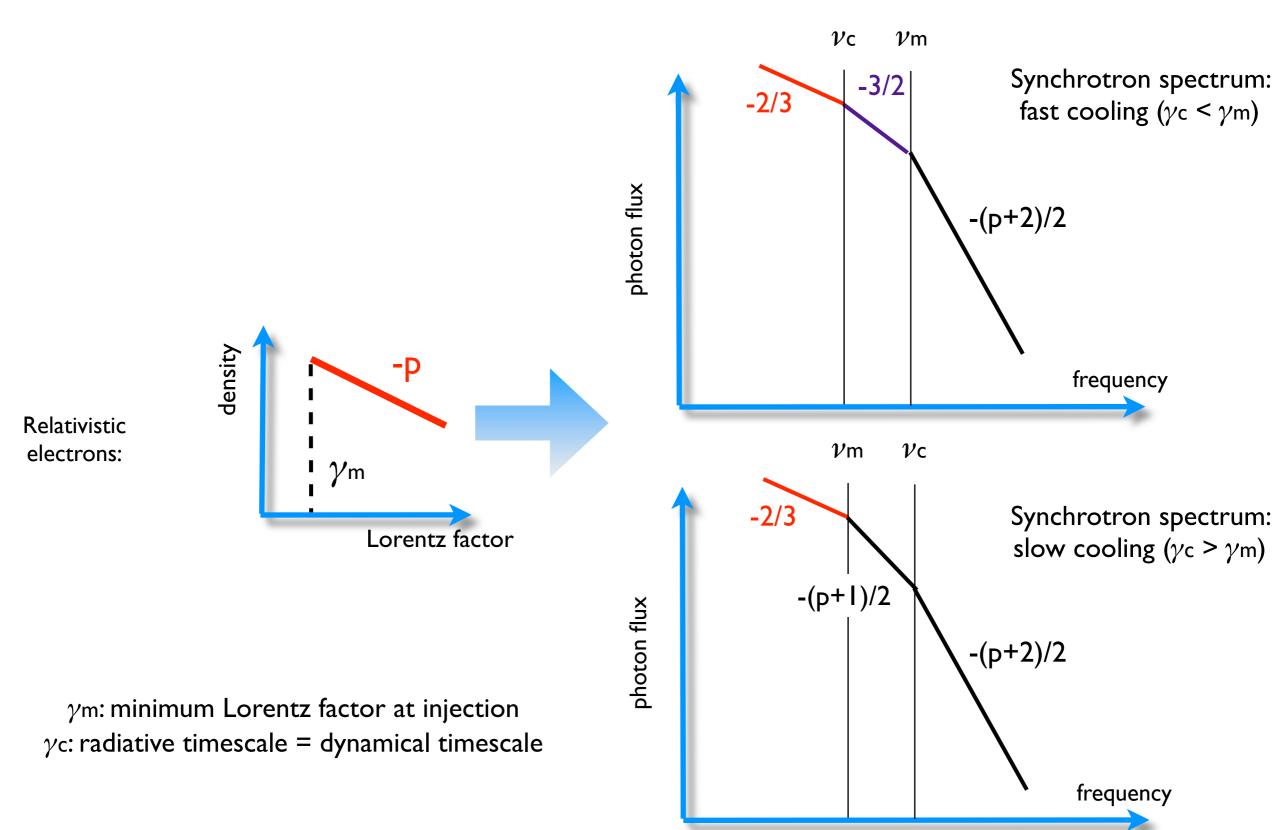


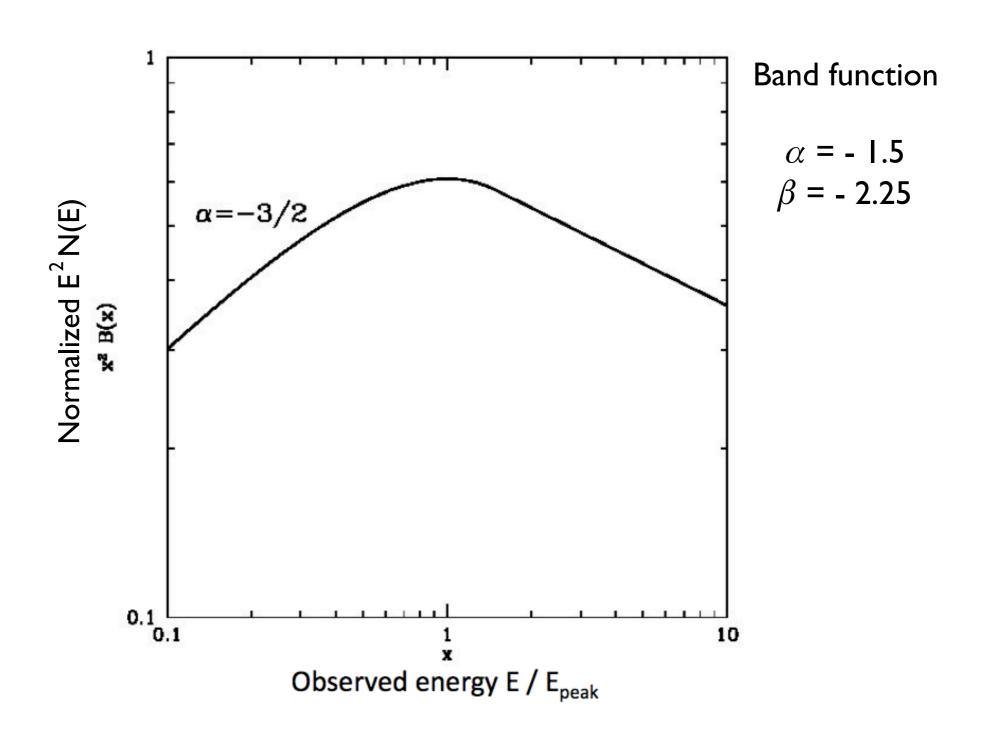


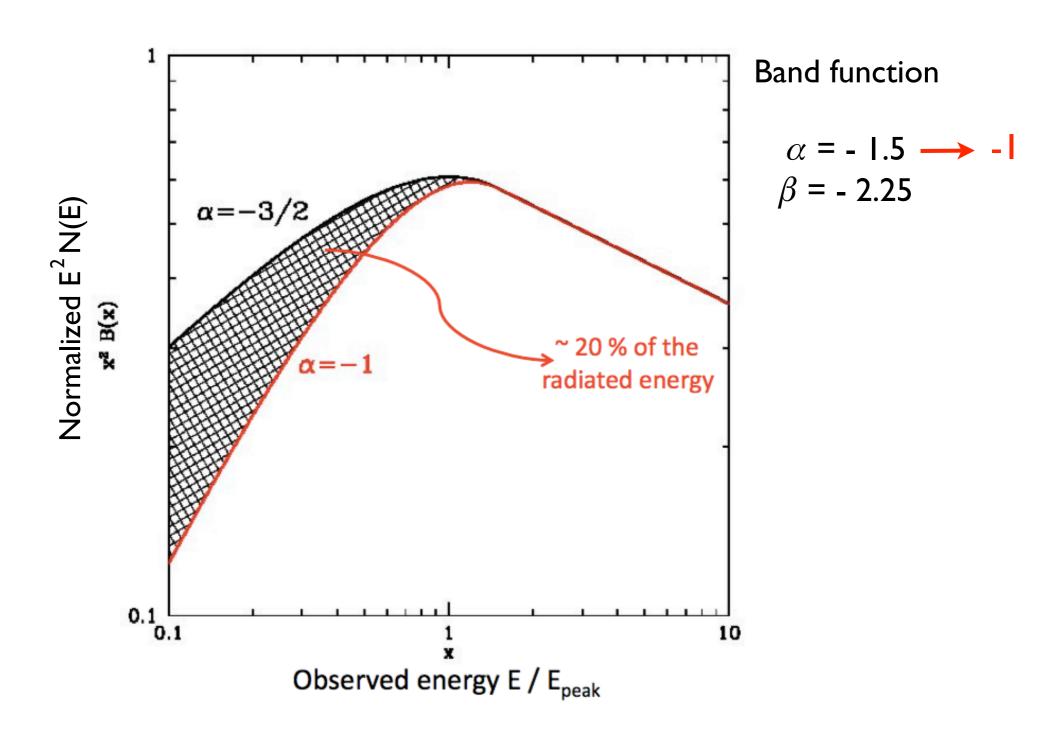




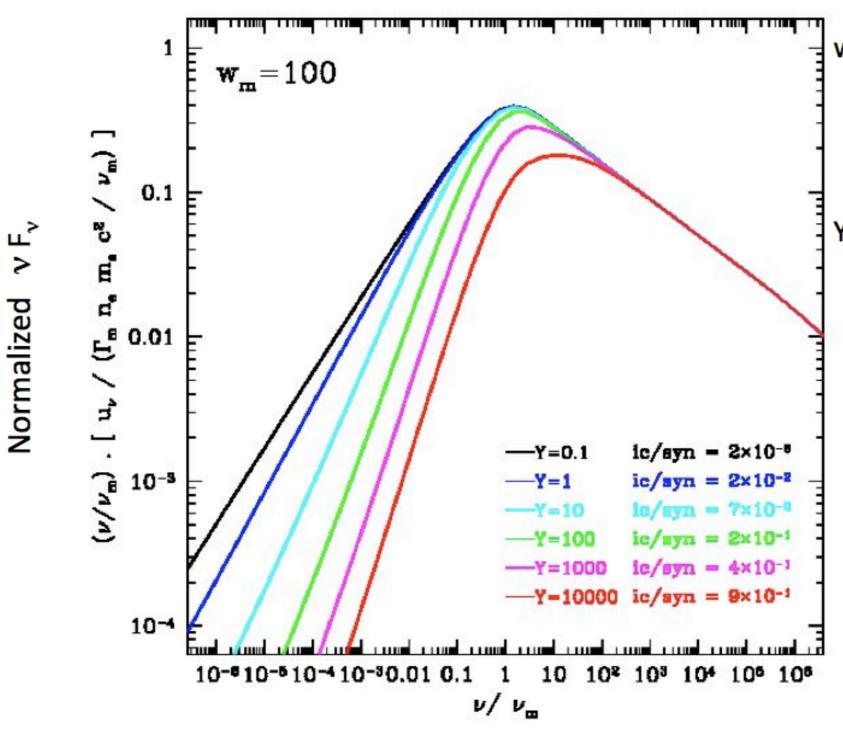








Inverse Compton scatterings in Klein-Nishina regime have an impact on the synchrotron slope  $\alpha$ 



w<sub>m</sub>: importance of KN

$$w_{
m m} = \Gamma_{
m m} rac{h 
u_{
m m}'}{m_{
m e} c^2}$$

Y: importance of IC vs syn

$$Y = rac{4}{3} au_{
m T} \Gamma_{
m m} \Gamma_{
m c} \simeq rac{\epsilon_{
m e}}{\epsilon_{
m B}}$$

Thomson regime: the electron cooling rate due to IC scatterings remains proportional to  $\gamma^2$  as for the synchrotron power

KN regime: the electron cooling rate due to IC depends on  $\gamma$ 

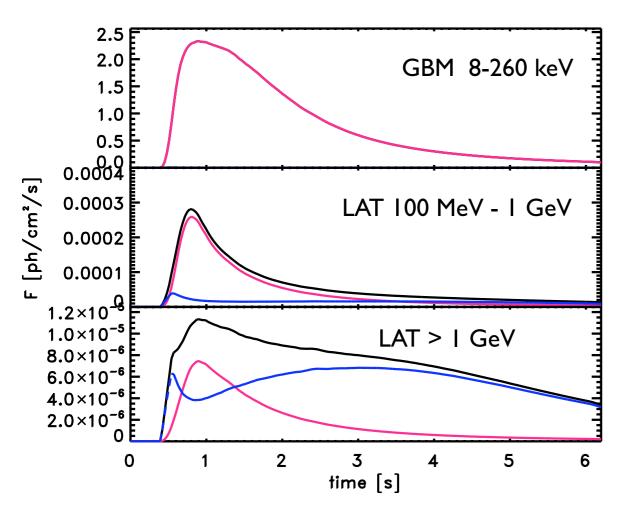
Frequency

Exact calculation with synchrotron + IC only (no adiabatic cooling, synchrotron self-absorption,  $\gamma\gamma$  annihilation)

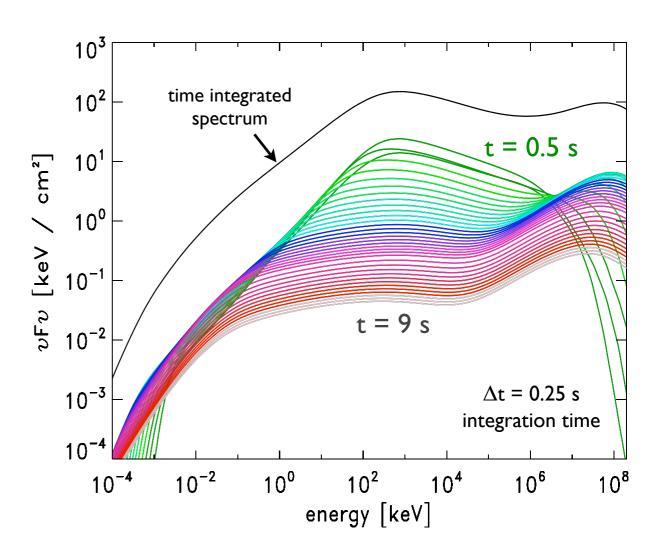
#### SYNCHROTRON CASE

#### low magnetic field

dE/dt = 5 x 
$$10^{53}$$
erg s,  $\varepsilon_{\rm B}^{-1}$   $\varepsilon_{\rm B}^{-1}$  = 0.0005,  $\varepsilon_{\rm e}^{-1}$  = 1/3,  $\zeta$  = 0.002, p = 2.5, z=1



Observed lightcurve

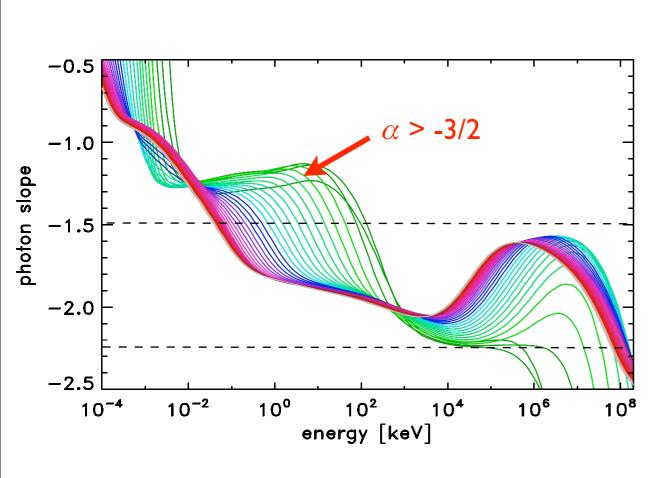


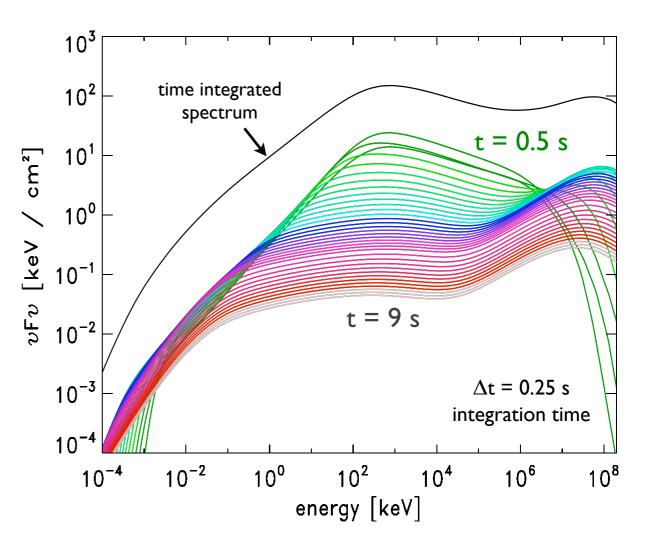
Time resolved spectra

#### SYNCHROTRON CASE

#### low magnetic field

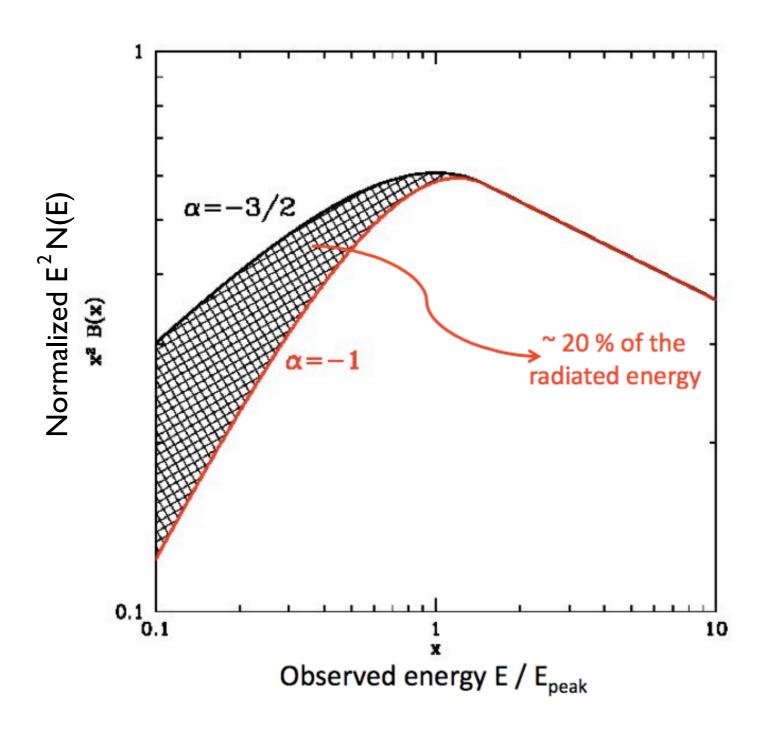
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low energy spectral slope  $\alpha = -3/2$  of the fast cooling synchrotron spectrum

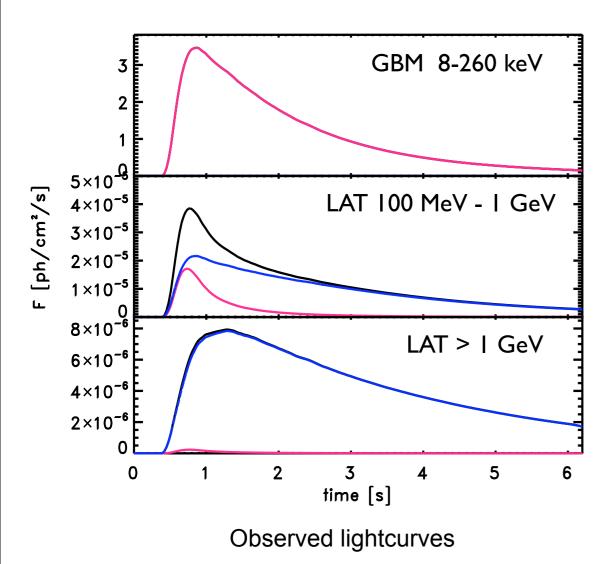
Time resolved spectra

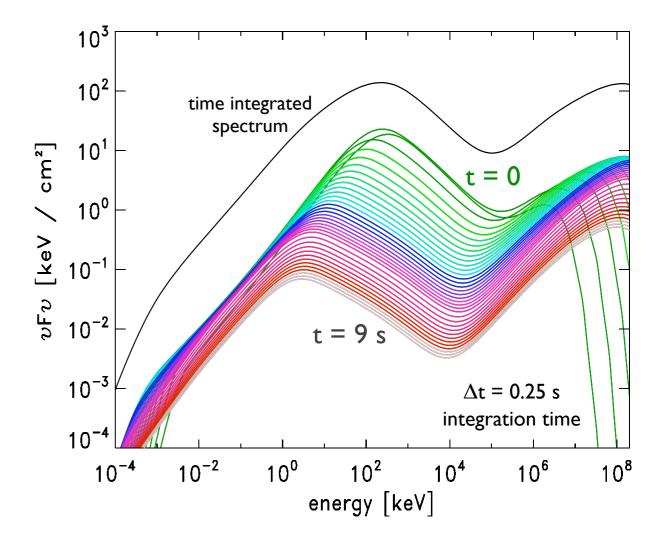


Spectral peak energy is also affected by IC scatterings!

#### SYNCHROTRON CASE

low magnetic field with  $\zeta$  varying and a steep slope p of the e- distribution dE/dt = 5 x 10<sup>53</sup>erg s,  $\varepsilon_{\rm R}$  = 0.0005,  $\varepsilon_{\rm p}$  = 1/3,  $\zeta_{\rm max}$  = 0.0025, p = 3.5, z=1

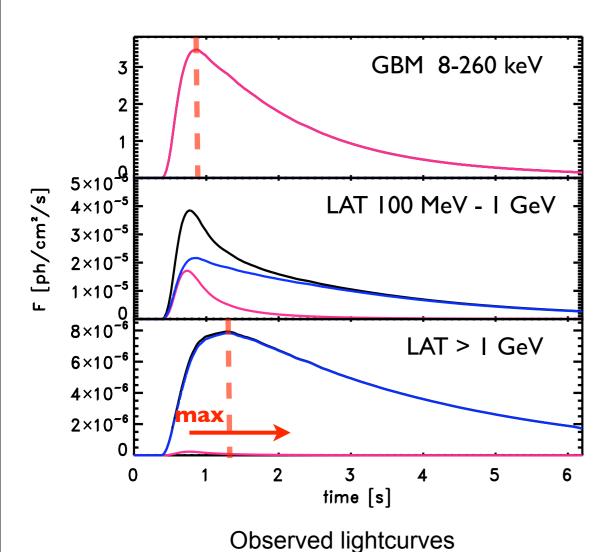


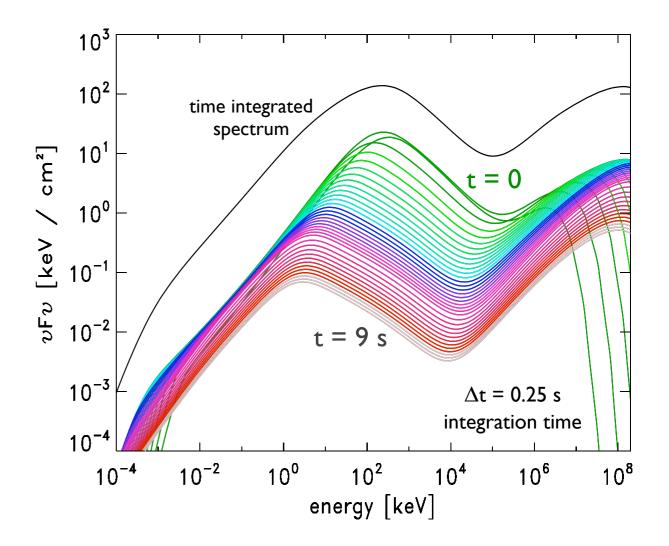


Time resolved spectra

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low magnetic field with  $\zeta$  varying and a steep slope p of the e- distribution dE/dt = 5 x 10<sup>53</sup>erg s,  $\varepsilon_R$  = 0.0005,  $\varepsilon_R$  = 1/3,  $\zeta_{max}$  = 0.0025, p = 3.5, z=1

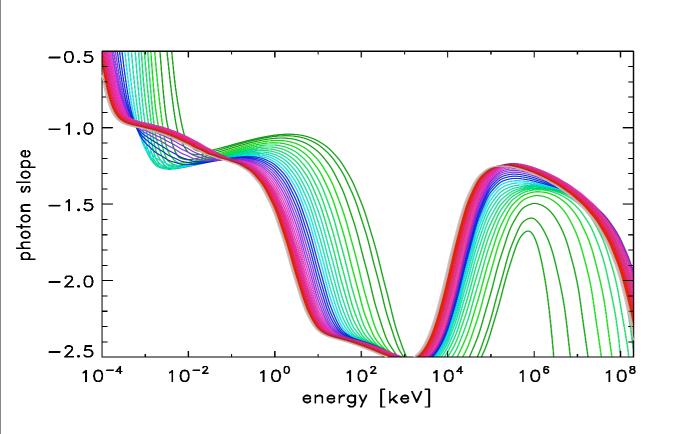


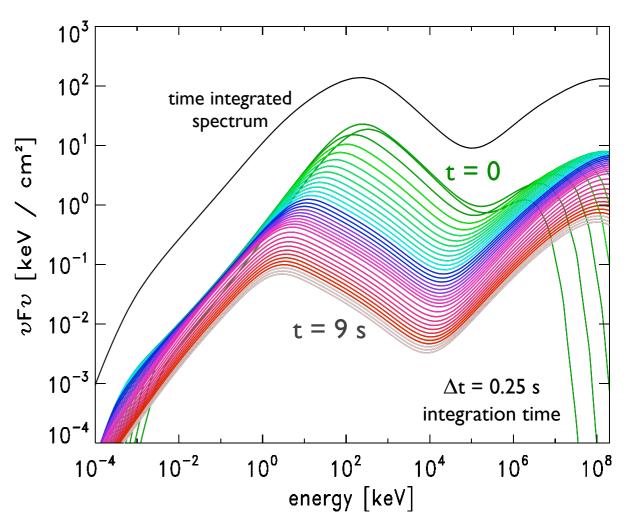


Time resolved spectra

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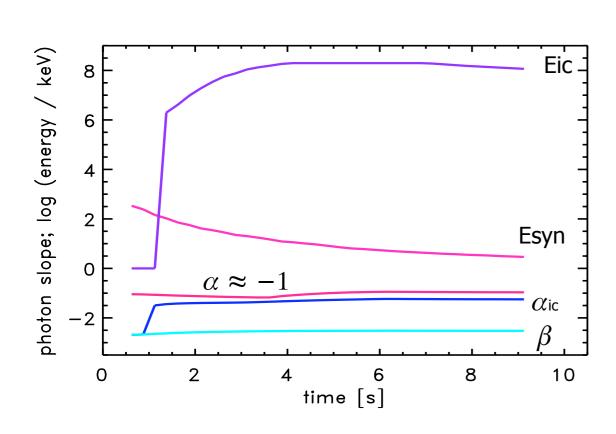


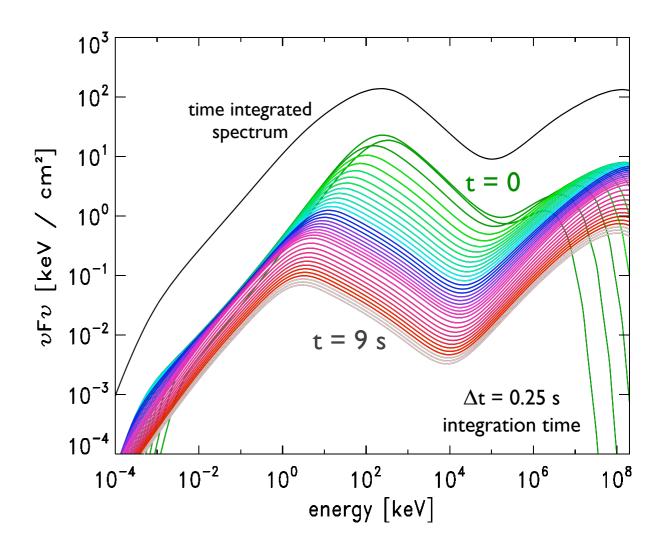


Time resolved spectra

#### SYNCHROTRON CASE

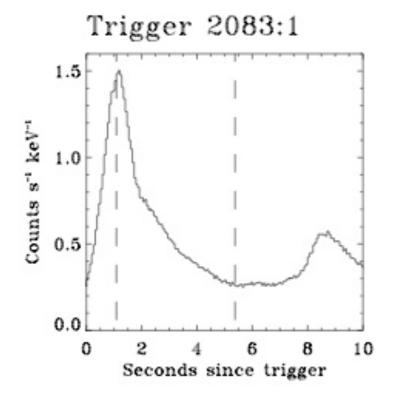
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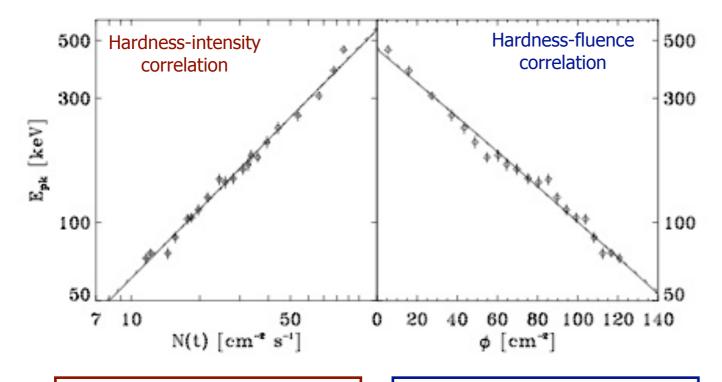


Time resolved spectra

### Spectral and temporal behavior: HIC & HFC



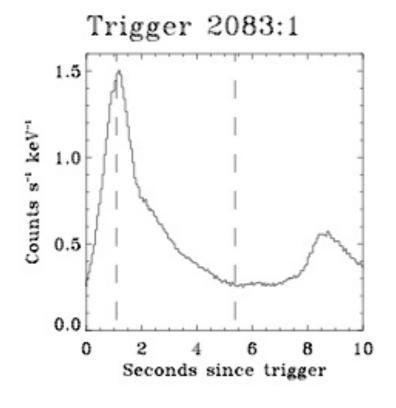
Kargatis 1995 Liang & Kargatis 1996 Ryde & Svensson 2002



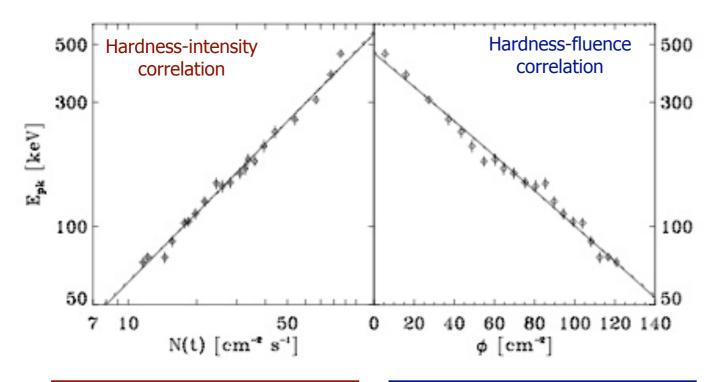
$$E_{\rm pk}(N) = E_{\rm pk, \ 0} (N/N_{\rm 0})^{\delta}$$

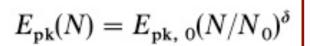
$$E_{\rm pk}(\Phi) = E_{\rm pk,\,max} \, e^{-\Phi/\Phi_0}$$

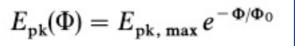
### Spectral and temporal behavior: HIC & HFC

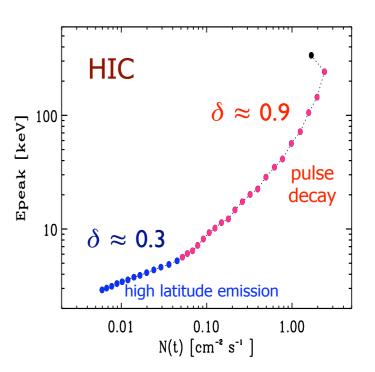


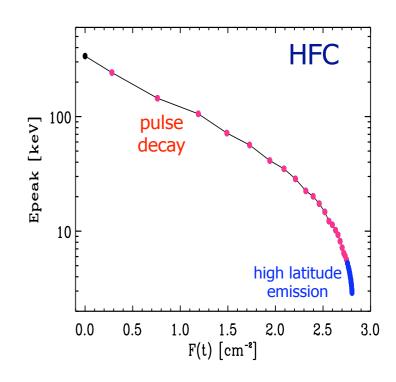
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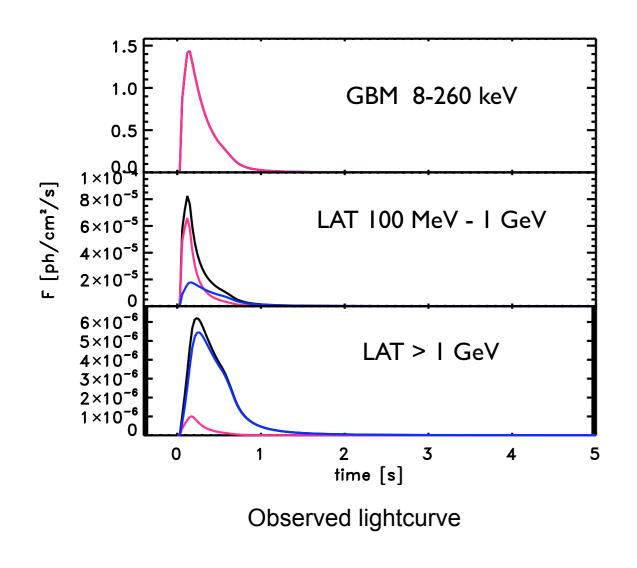


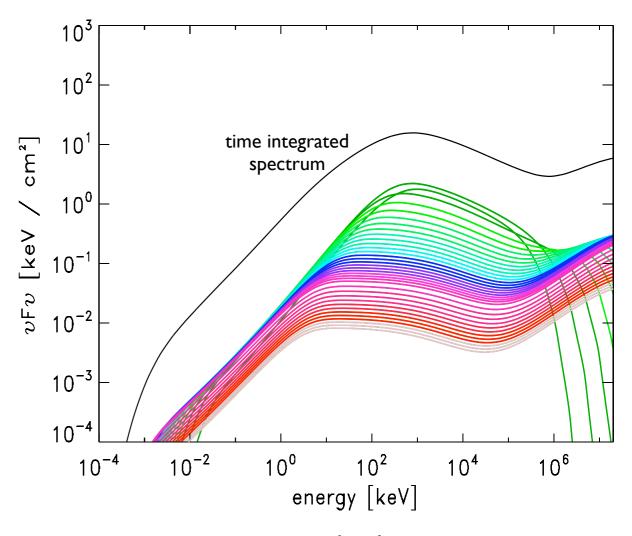
### Modeling of short pulses, multi-peaked bursts...

#### Short GRBs are expected to be harder in our model:

Ris  $\sim$  c t<sub>var</sub>  $\Gamma^2 \Rightarrow$  closer to the source if the burst evolves on a shorter time scale

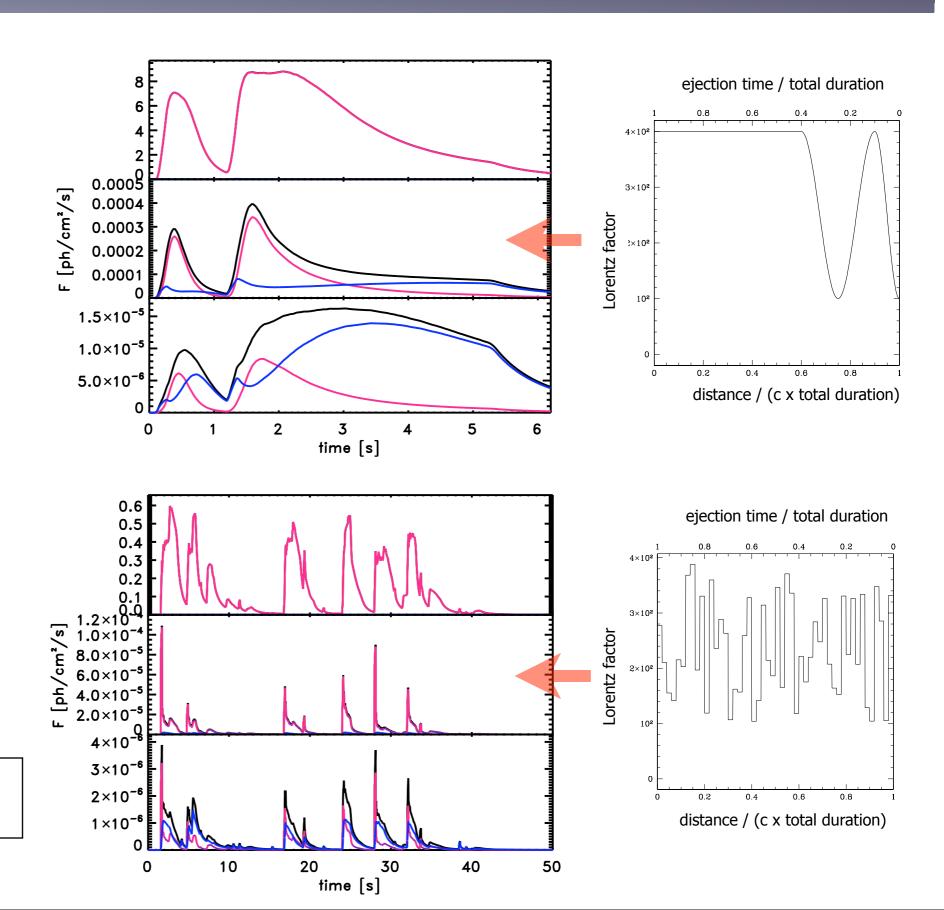
The equipartition B is stronger and Esyn is larger in the dissipation region





Time resolved spectra

### Modeling of short pulses, multi-peaked bursts...



synchrotron inverse Compton total

### Summary

We developed modeling tools to compute the GRB prompt emission from internal shocks in a time-dependent way in different spectral bands, including the high-energy gamma rays

The exploration of the parameter space shows that we can expect two classes of broad-band spectra, which correspond to different physical conditions in the shocked region: "synchrotron case" (where the dominant process in Fermi-GBM range is synchrotron radiation) and "inverse Compton case" (where the synchrotron component peaks at low energy and dominant process in GBM range is inverse Compton)

Fermi GRB observations favor the "synchrotron case". We investigated the effect of the different radiation processes on spectral parameters and light curve properties. We are currently modeling in detail the broad-band spectra of Fermi GRBs