

Probing time-dependent cosmic birefringence with LiteBIRD

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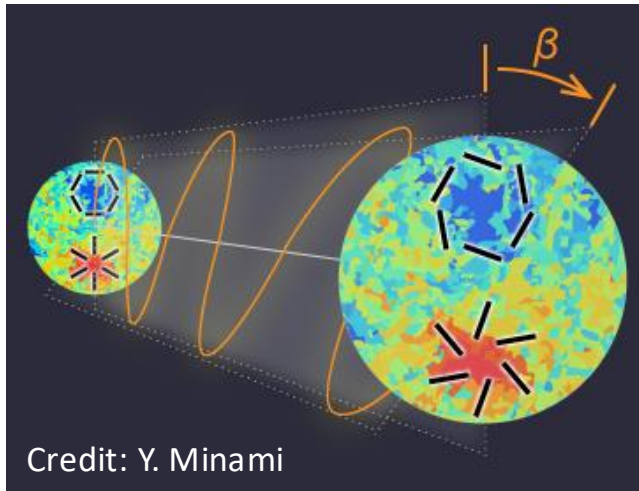
(working with people in IPMU CMB group (Baptiste, Marta, Ryosuke, Tomo, Yuji, ...))



2026.7.3 @ Challenges at the CMB Frontier (Chino, Nagano)

Isotropic polarization rotation in CMB

Lue, Wang & Kamionkowski (1999); Feng+ (2005,2006); Liu, Lee & Ng (2006); ...



■ Parity-violating interactions

$$\mathcal{L}_{\text{int}} \supset \frac{1}{4} g_{\phi\gamma} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

convert E- and B-mode polarization as

$$\begin{pmatrix} E_{\ell m} \\ B_{\ell m} \end{pmatrix}^{\text{obs}} = \begin{pmatrix} \cos(2\beta) & -\sin(2\beta) \\ \sin(2\beta) & \cos(2\beta) \end{pmatrix} \begin{pmatrix} E_{\ell m} \\ B_{\ell m} \end{pmatrix}^{\text{CMB}}$$

$\uparrow$ observed $\uparrow$ intrinsic

■ It produces a **parity-odd EB correlation**

$$C_{\ell}^{EB,o} = \frac{1}{2} \sin(4\beta) \left(C_{\ell}^{EE,\text{CMB}} - C_{\ell}^{BB,\text{CMB}} \right) + \cos(4\beta) C_{\ell}^{EB,\text{CMB}}$$

(note: β is assumed to be constant)

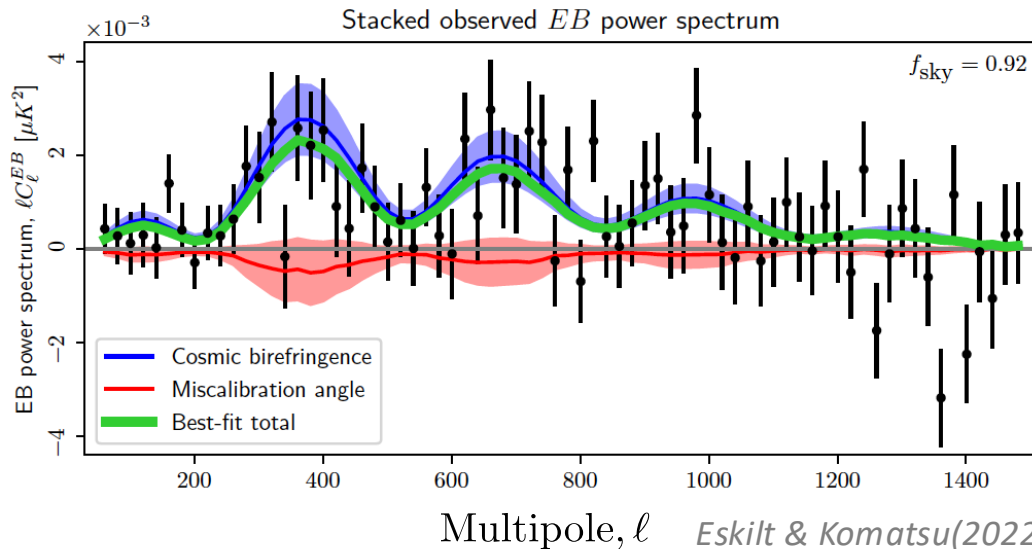
$\uparrow$ assuming 0

Measurements of cosmic birefringence

- Nonzero isotropic cosmic birefringence (ICB) angle was reported by *Planck* data:

PR3: $\beta = 0.35 \pm 0.14$ deg *Minami, Komatsu (2020)*;

PR4: $\beta = 0.30 \pm 0.11$ deg *Diego-Palazuelos+ (2022)*;

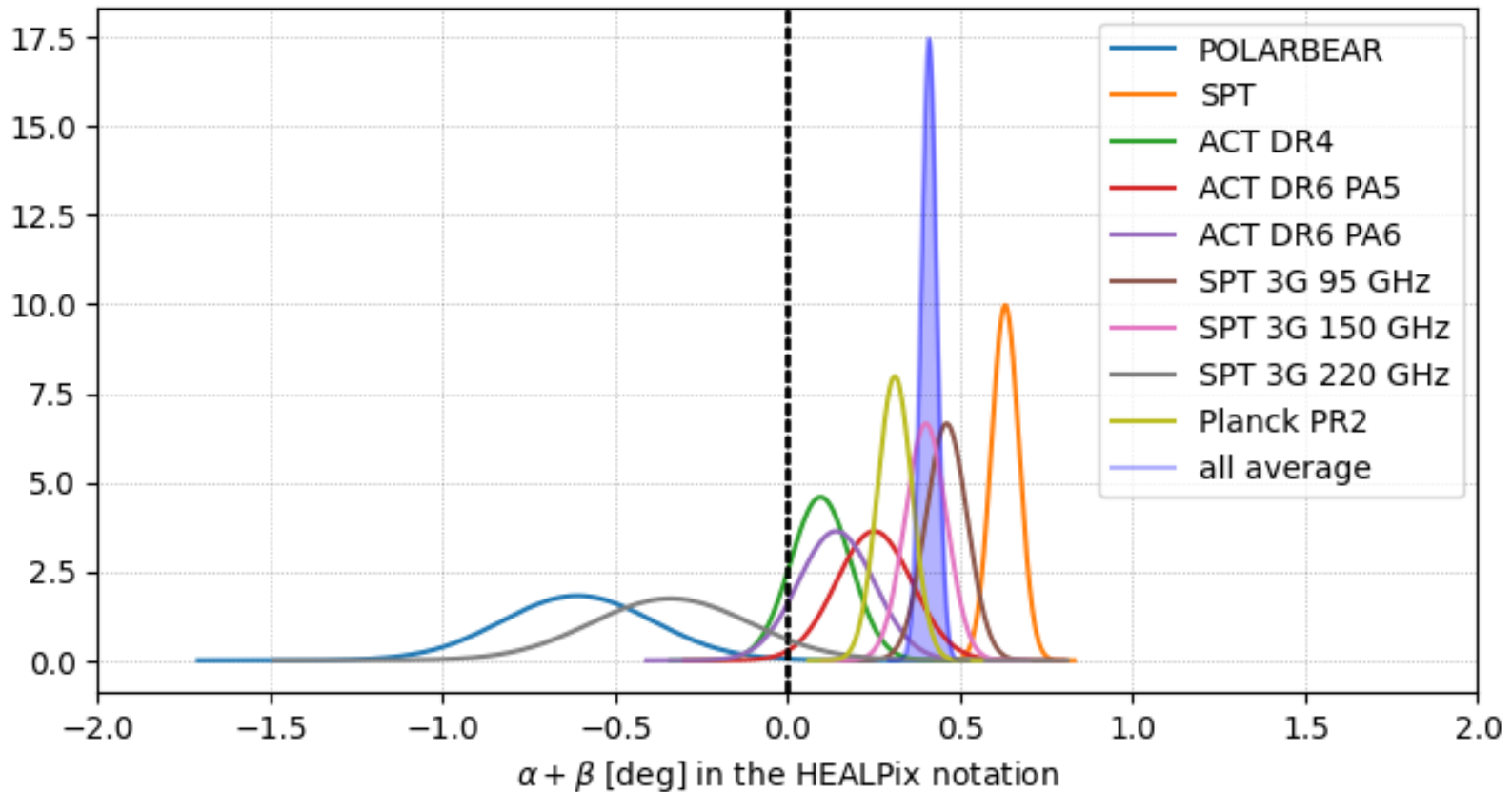


- *Planck*/*WMAP* joint analysis:
Eskilt & Komatsu (2022);

$\beta = 0.34 \pm 0.09$ deg (3.6σ)

Other recent measurements overview

Credit: Fumihiro Naokawa

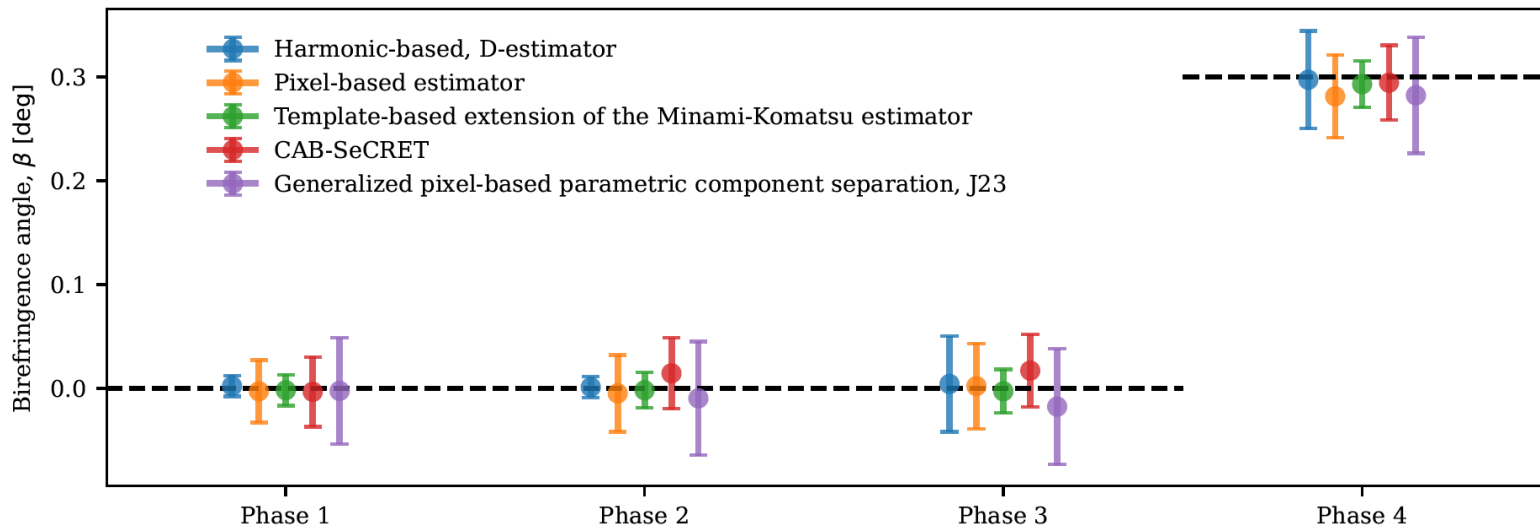


✂ the calibration method as in Planck PR3, PR4 & WMAP is not used in these measurements



Future forecast with LiteBIRD

de la Hoz, Diego-Parazuelos, Errard, Gruppuso, Jost, Sullivan, Bortolami, Chinone, Hergt, Komatsu, Minami, IO, Paoletti, Scott, Vielva+ (2025)



Pipeline	Phase 1 $\beta [\times 10^{-2} \text{ deg}]$	Phase 2 $\beta [\times 10^{-2} \text{ deg}]$	Phase 3 $\beta [\times 10^{-2} \text{ deg}]$	Phase 4 $\beta [\times 10^{-2} \text{ deg}]$
<i>D</i> -estimator	0.2 ± 1.0	0.1 ± 1.0	0.4 ± 4.6	29.7 ± 4.7
Pixel-based estimator	-0.3 ± 3.0	-0.5 ± 3.7	0.2 ± 4.1	28.1 ± 4.0
Template-based MK	-0.2 ± 1.5	-0.2 ± 1.7	-0.3 ± 2.1	29.3 ± 2.2
CAB-SeCRET	-0.4 ± 3.4	1.4 ± 3.4	1.7 ± 3.5	29.4 ± 3.6
J23	-0.2 ± 5.1	-1.0 ± 5.5	-1.8 ± 5.6	28.2 ± 5.6

Implications for axion cosmology

(axion-like particle)

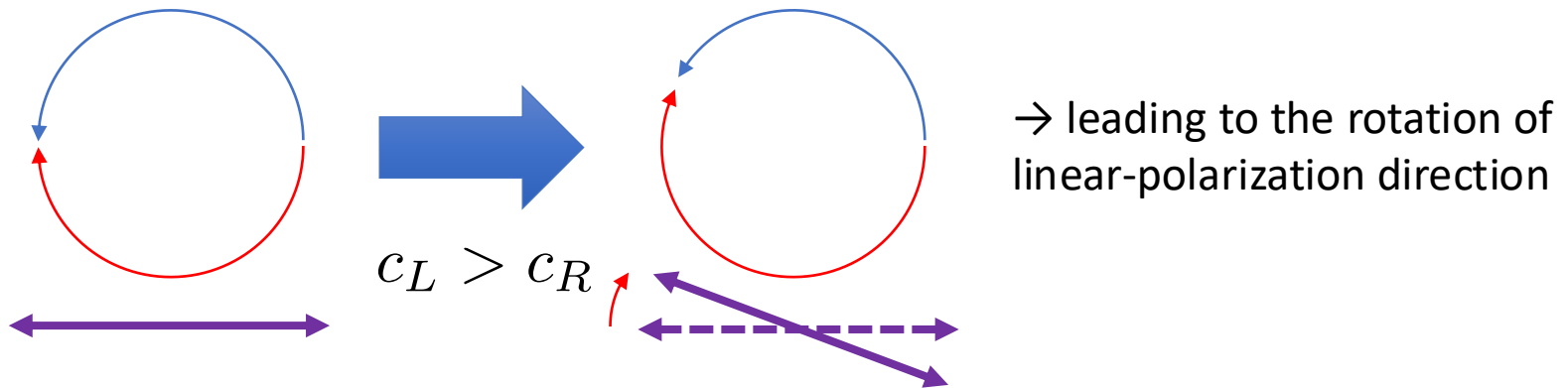
Cosmic birefringence by axion (ALP)

Harari & Sikivie (1992); Carroll (1998); ...

$$\mathcal{L} \supset \frac{1}{4} g_{\phi\gamma} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- Dispersion relation of left- and right-handed polarization is modified:

$$\ddot{A}_k^{\text{L/R}} + \omega_{\text{L/R}}^2 A_k^{\text{L/R}} = 0, \quad c_{\text{L/R}} \equiv \frac{\omega_{\text{L/R}}}{k} = \sqrt{1 \pm \frac{g_{\phi\gamma} \dot{\phi}}{k}}$$

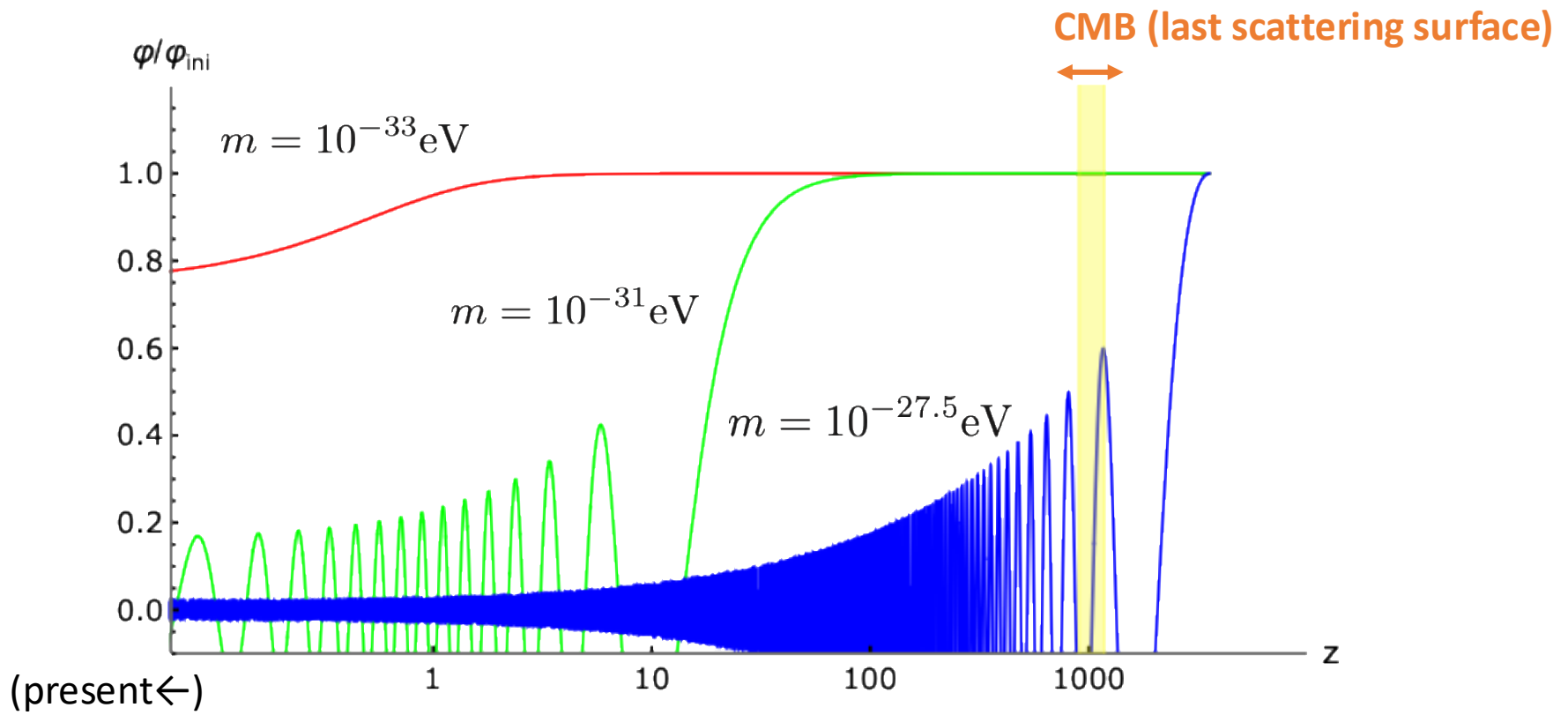


- Rotation angle is given by the integration of phase velocity difference

$$\beta = \frac{1}{2} \int_{t_{\text{emit}}}^{t_{\text{obs}}} dt (\omega_L - \omega_R) = \frac{g_{\phi\gamma}}{2} \int_{t_{\text{emit}}}^{t_{\text{obs}}} dt \dot{\phi} = \frac{g_{\phi\gamma}}{2} [\phi(t_{\text{obs}}) - \phi(t_{\text{emit}})]$$

Birefringence from ALP field dynamics

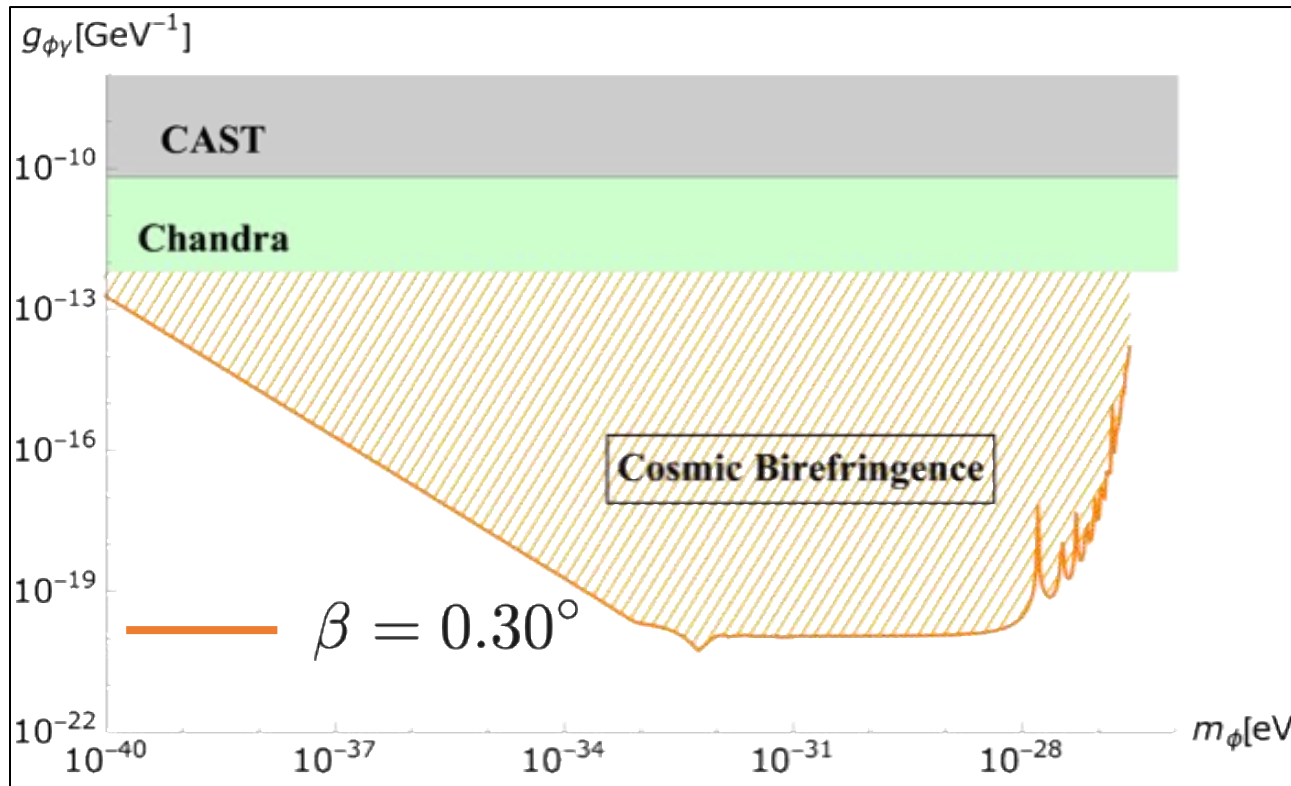
■ EOM: $\ddot{\phi} + 3H\dot{\phi} + m^2\phi = 0$



$\beta = \frac{g_{\phi\gamma}}{2}(\phi(t_0) - \langle\phi(t_{\text{LSS}})\rangle)$ ■ Field displacement depends on ALP mass

ALP discovery space

$$\beta = \frac{g_{\phi\gamma}}{2} (\phi(t_0) - \langle \phi(t_{\text{LSS}}) \rangle)$$



➤ Less ALP abundance

$$\frac{m_{\phi}^2 \phi^2 / 2}{\rho_c} = \Omega_{\phi} \ll 1$$

➤ $n\pi$ phase ambiguity

$$\beta = 0.30^{\circ} + n\pi$$

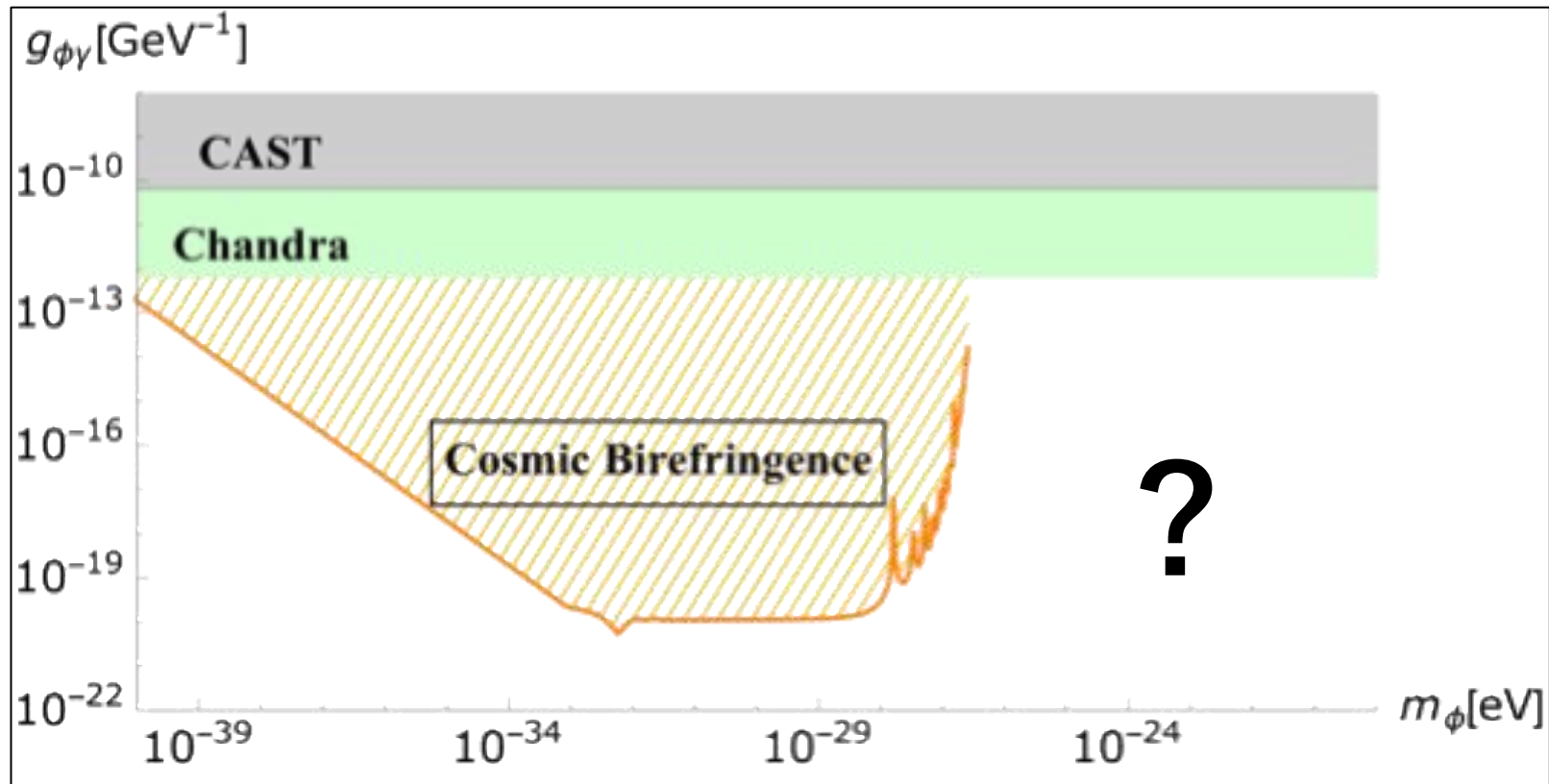
*Naokawa, Namikawa,
Murai, IO, Kamada, PRL (2024);*

ALP abundance

$$\Omega_{\phi, \text{max}} \simeq \begin{cases} 0.69 & (m \lesssim 10^{-33} \text{eV}) \\ 6 \times 10^{-3} h^{-2} & (10^{-32} \text{eV} \lesssim m \lesssim 10^{-25} \text{eV}) \end{cases}$$

ALP discovery space

$$\beta = \frac{g_{\phi\gamma}}{2} (\phi(t_0) - \langle \phi(t_{\text{LSS}}) \rangle)$$



How about fuzzy DM mass range?: $m_{\phi} \gtrsim 10^{-25} \text{eV}$

■ Mixed CDM+FDM simulations

→ ALP undergoes **nonlinear clustering**

Schwabe, Gosenca, Behrens, Niemeyer, Easter (2020); Laguë, Schwabe, Hložek, Marsh, Rogers (2023);...

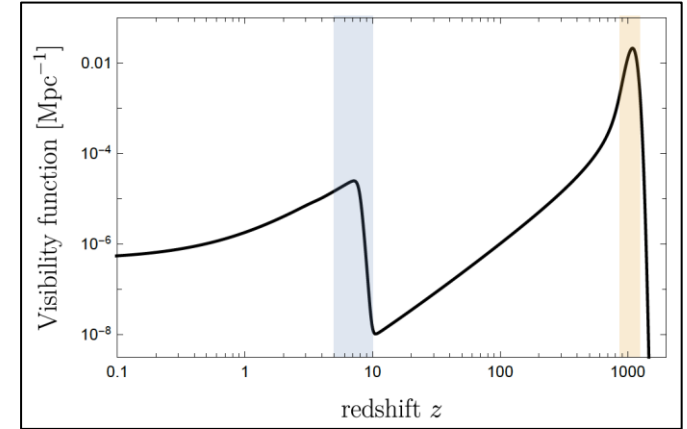
ALP oscillatory effects at CMB epoch

$$m_\phi \gg H_{\text{CMB}} \sim 10^{-29} \text{ eV}$$

- Stokes parameters with ALP coupling

$$Q \pm iU \propto \int dt' g(t') \exp [\pm 2i\beta(t')]$$

$g(t') : \text{vigibility function}$



Credit: Nakatsuka, Namikawa, Komatsu (2022)

Rotation angle

$$\begin{aligned} \beta &= \frac{g_{\phi\gamma}}{2} [\phi(t_0) - \phi(t_{\text{LSS}})] \\ &\propto a(t')^{-3/2} \\ &= \frac{g_{\phi\gamma}}{2} [\phi_0 \cos(mt_0 + \delta_0) - \phi_{\text{LSS}}(t') \cos(mt' + \delta(t'))] \end{aligned}$$

$$\propto e^{\pm i g_{\phi\gamma} \phi(t_0)} \int dt' g(t') \exp [\mp i g_{\phi\gamma} \phi_{\text{CMB}}(t') \cos(mt' + \delta(t'))]$$

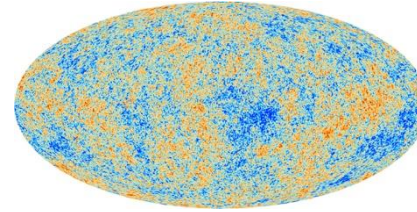
$$= \langle e^{\mp i g_{\phi\gamma} \phi(t_{\text{LSS}})} \rangle$$

CMB washout effect

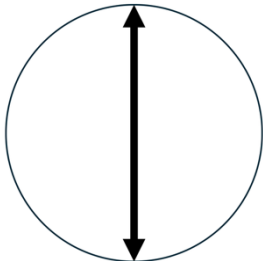
Finelli & Galaverni (2008);

Fedderke, Graham, Rajendran (2019);

Murai (2024); Zhang, Ferreira, IO, Namikawa (2024);



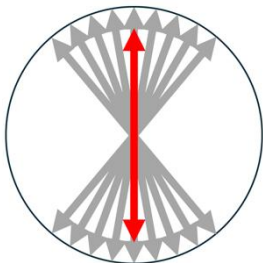
Last scattering surface



Standard polarization magnitude (w/o ALP coupling)



$$\langle e^{\mp i g_{\phi\gamma} \phi(t_{\text{LSS}})} \rangle \sim 1 - g_{\phi\gamma}^2 \langle \phi_{\text{LSS}}^2 \rangle / 4 \quad m\Delta T \gg 1$$

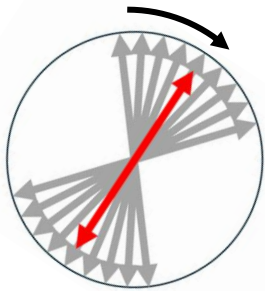


**Reduced polarization magnitude (w ALP rapid oscillation)
by the superposition of randomly-oriented polarization**

To be summarized...

- Stokes parameters with ALP-photon coupling for mass: $10^{-25} \text{eV} \lesssim m_\phi \lesssim 10^{-23} \text{eV}$

$$Q \pm iU \sim \left(1 - \frac{g_{\phi\gamma}^2 \langle \phi_{\text{LSS}}^2 \rangle}{4} \right) \exp \left[\pm i g_{\phi\gamma} \phi_0 \cos(m_\phi t_0 + \delta_0) \right] (Q_0 \pm iU_0)$$



Washout effect

Polarization rotation effect

$$= \pm 2i\beta \quad (\text{cosmic birefringence})$$

- For ALP mass as (fuzzy) dark matter

$$m_\phi \gtrsim 10^{-23} \text{ eV}$$

$$(t_{\text{osc}} \lesssim \text{years})$$

**Polarization rotation becomes time-varying
in real time measurements**

(frequency = DM mass/(2π))

Constraining on axion parameter space

■ Axion dark matter energy density: $\rho_\phi = \frac{1}{2}m_\phi^2\phi^2 \longleftrightarrow \phi = \frac{\sqrt{2\rho_\phi}}{m_\phi}$

■ Convert polarization magnitude into coupling constant:

$$\beta = \frac{g_{\phi\gamma}}{2}\phi_0 \longleftrightarrow g_{\phi\gamma} = \sqrt{2}\beta m_\phi (\kappa \rho_{\text{DM,local}})^{-1/2}$$

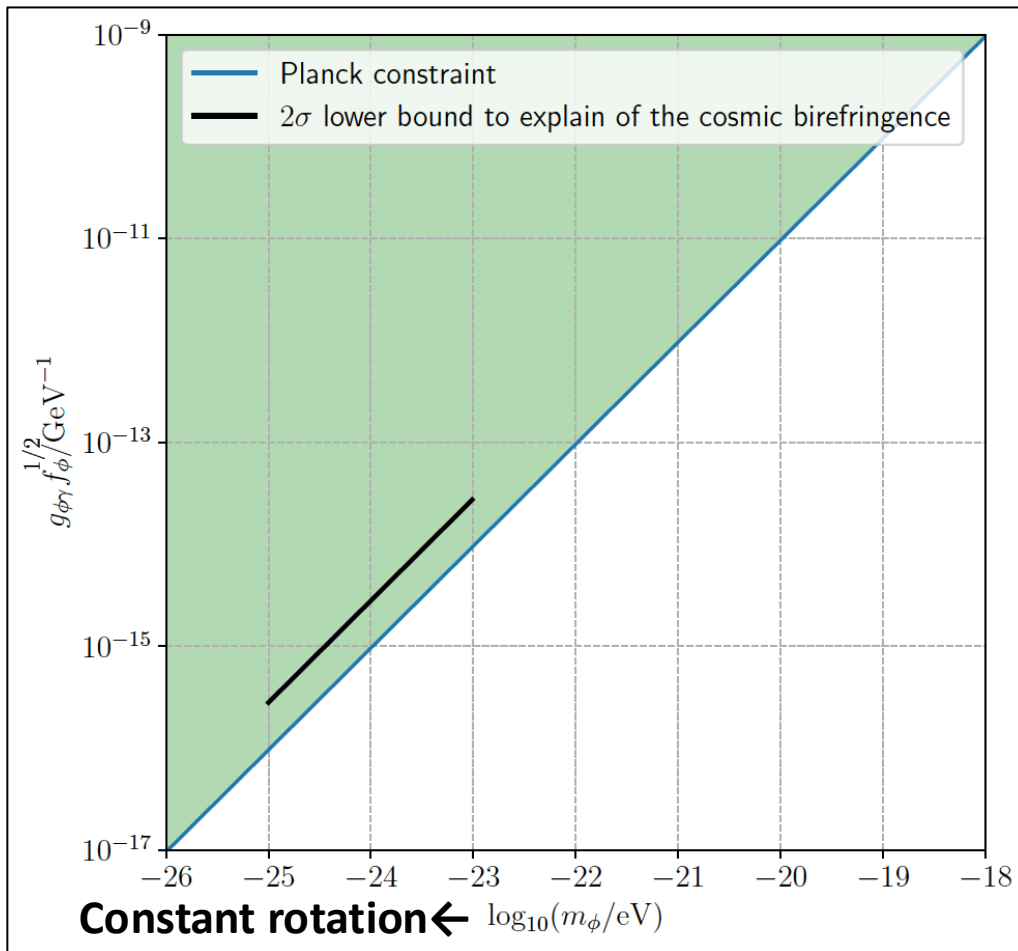
$$\kappa \equiv \frac{\rho_{\phi,\text{local}}}{\rho_{\text{DM,local}}} \quad : \text{ axion dark matter fraction}$$

$$g_{\phi\gamma}\kappa^{1/2} \simeq 4.9 \times 10^{-14} \text{GeV}^{-1} \left(\frac{\beta}{0.3\text{deg}} \right) \left(\frac{m_\phi}{10^{-23}\text{eV}} \right) \left(\frac{\rho_{\text{DM,local}}}{0.3\text{GeV/cm}^3} \right)^{-1/2}$$

Washout effect v.s. Cosmic birefringence

Zhang, Ferreira, IO, Namikawa (2024)

$$(\beta \sim 0.3^\circ)$$



Comparable, but slightly excluded

Note

We assume that
cosmological DM fraction

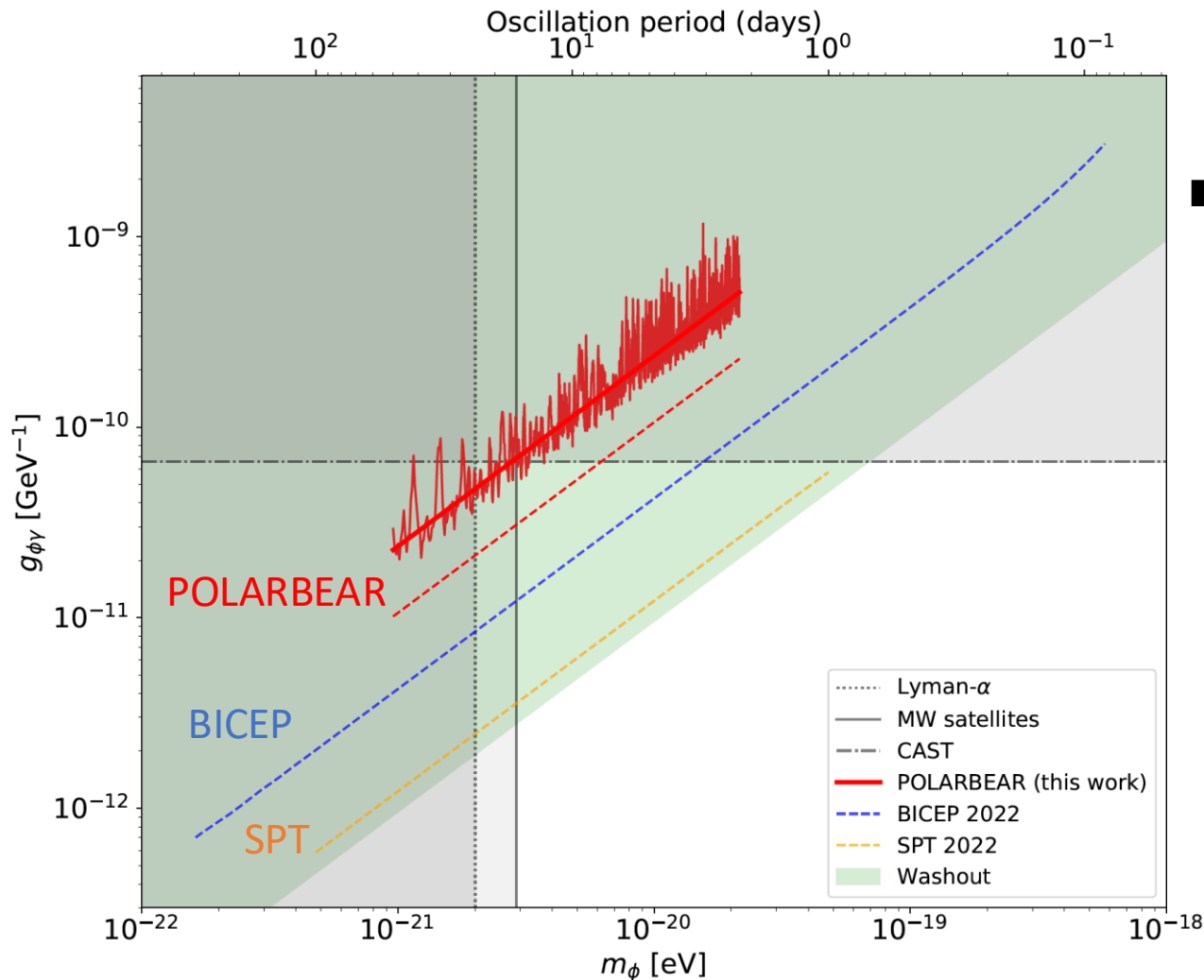
$$f_\phi \equiv \frac{\Omega_\phi}{\Omega_m}$$

equals to local DM fraction

$$\kappa \equiv \frac{\rho_{\phi, \text{local}}}{\rho_{\text{DM, local}}}$$

(comparing with observing time)

Time-dependent birefringence constraints



POLARBEAR 2023;

■ Fit a sinusoidal model and extract limits on the coupling constant

■ 100% DM fraction is assumed

$$(f_{\phi} = \kappa = 1)$$

■ Stochastic axion distribution is taken into account (POLARBEAR)

SPT-3G v.s. LiteBIRD for birefringence test

(not full survey)

SPT-3G

LiteBIRD

Angular scales	high ℓ	low ℓ ($2 \leq \ell \leq 200$)
Angular resolution	1.2'	30'
Map depth	6.6 $\mu\text{K} \cdot \text{arcmin}$ (?)	2.5 – 5.0 $\mu\text{K} \cdot \text{arcmin}$
Sky fraction	$f_{\text{sky}} \simeq 0.036$ (1500deg ²)	$f_{\text{sky}} = 0.6 - 0.7$ (full sky survey)
Observation time for oscillation period	<i>SPT-3G (2022)</i> 12 hours – 100 days	? ? 1 day – 3 years

(very naive) Sensitivity estimation

■ Birefringence signal: $C_\ell^{EB,o} = \frac{1}{2} \sin(4\beta) \left(C_\ell^{EE,CMB} - C_\ell^{BB,CMB} \right)$

$$\beta \ll 1, \quad C_\ell^{BB} \ll C_\ell^{EE} \quad \beta \simeq \frac{C_\ell^{EB,o}}{2C_\ell^{EE,CMB}}$$

Yadav+(2009); Gluscevic, Kamionkowski (2010);...

■ Uncertainty of rotation angle: $\sigma_\beta^{-2} = \sum_\ell \sigma_{\beta,\ell}^{-2} = \sum_\ell \frac{(2C_\ell^{EE,CMB})^2}{\text{Var}(C_\ell^{EB,o})}$

$$C_\ell^{EE,CMB} = \frac{2\pi}{\ell(\ell+1)} D_\ell^{EE} \quad N_\ell^{EE} = N_\ell^{BB} = (\Delta_P)^2 \exp[\ell(\ell+1)\sigma_b^2]$$

$$\rightarrow \sigma_\beta^{-2} \simeq 16\pi^2 f_{\text{sky}} \sum_\ell \frac{2\ell+1}{\ell^2(\ell+1)^2} \frac{(D_\ell^{EE})^2}{(\Delta_P)^4 \exp[2\ell(\ell+1)\sigma_b^2]}$$

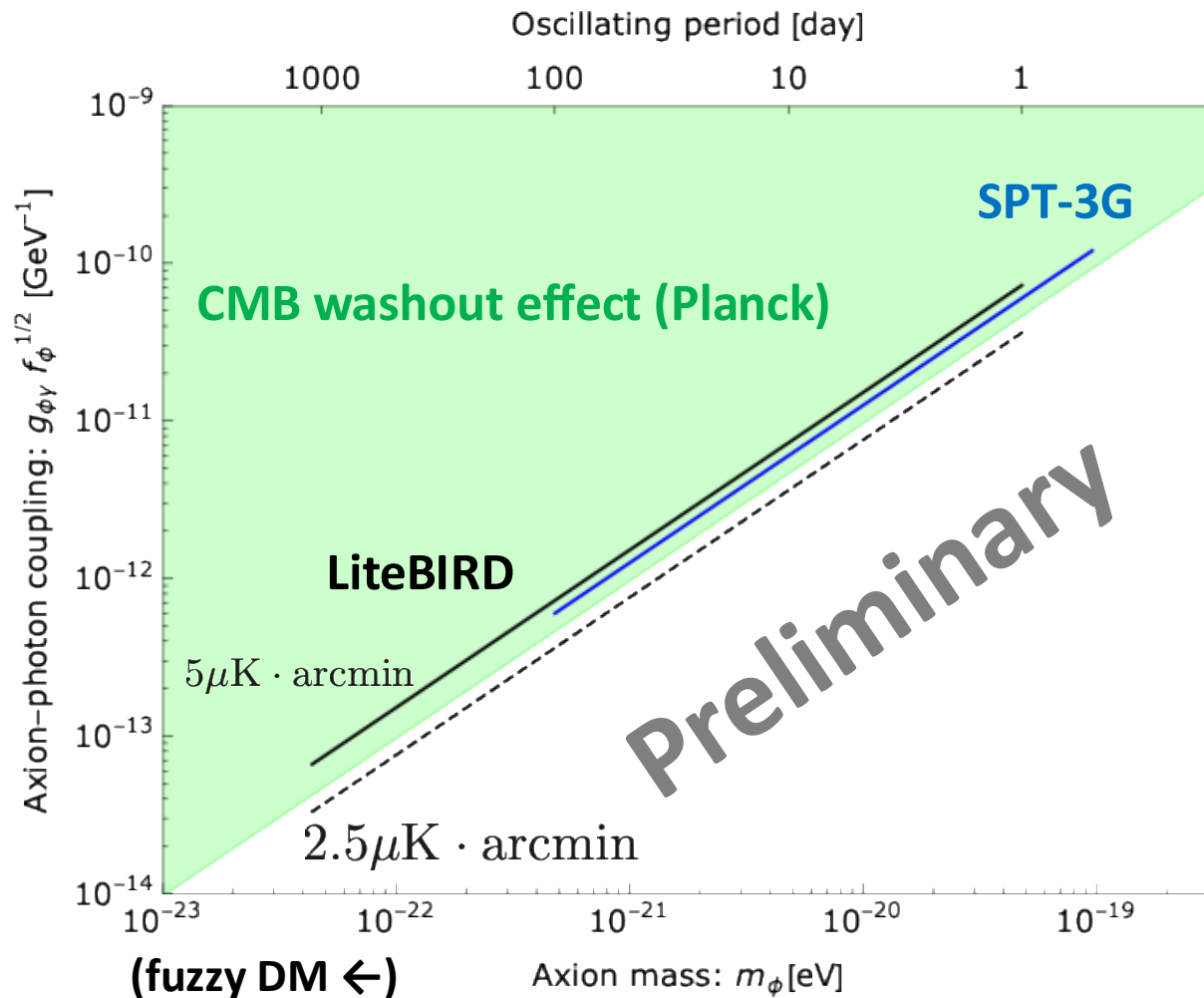
■ 2σ upper bound on angle magnitude for LiteBIRD:

$$\sigma_b \equiv \theta_{\text{FWHM}} / \sqrt{8 \ln 2}$$

$$A_{95} = \sqrt{2\sigma_\beta} \sim 0.09 - 0.18 \text{ deg}$$

Potential constraint with LiteBIRD

(IO+, in preparation)



■ Assuming the same DM fraction: $f_{\phi} = \kappa$

■ Stochastic effect on DM local density is not concerned in this plot

Summary & Outlook

- Cosmic birefringence is a phenomenon of polarization rotation effect in the presence of parity-violating physics, such as axion-photon interaction.
- For heavier mass range, oscillatory motion of axion field washes out the CMB polarization and puts an upper bound on the axion photon coupling constant.
- Probing time dependent cosmic birefringence is also important for the dark matter search. We made a very naïve sensitivity estimation expected by LiteBIRD.
- LiteBIRD potentially covers lower mass ranges than that with SPT-3G, POLARBEAR, BICEP, and improves the constraint on the parameter space.
- More realistic estimation is necessary.

Thank you very much!